Settling-driven Instabilities in Mammatus Clouds

S. Ravichandran, Eckart Meiburg and Rama Govindarajan

Turbulence from angstroms to light years









Mammatus Clouds

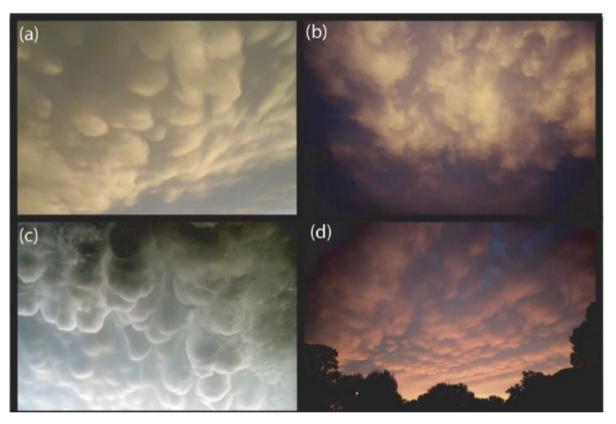


Fig. 2 from Schultz et al. 2006

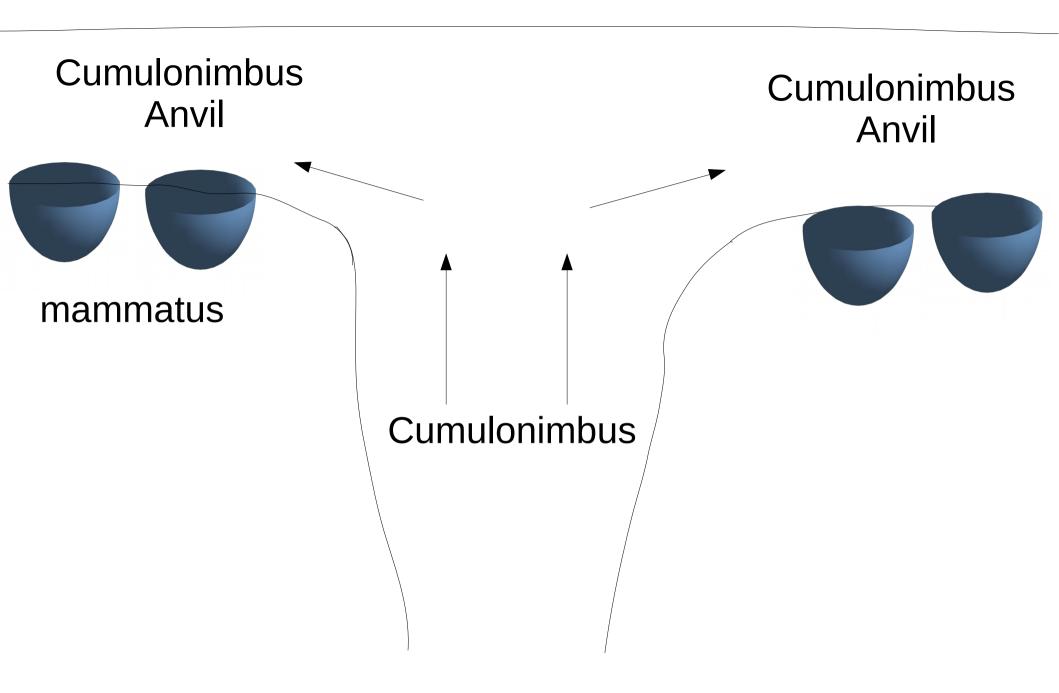
Asperitas Clouds



In Cincinnati, Ohio by Ron Steele on August 3, 2015.



A dramatic asperitas cloud. (WMO International Cloud Atlas/Gary McArthur)



Typical conditions

REVIEW

The Mysteries of Mammatus Clouds: Observations and Formation Mechanisms

David M. Schultz,*,+ Katharine M. Kanak,* Jerry M. Straka,* Robert J. Trapp,[@]
Brent A. Gordon,[&] Dusan S. Zrnić,⁺ George H. Bryan,** Adam J. Durant,⁺⁺
Timothy J. Garrett,** Petra M. Klein,* and Douglas K. Lilly[@]@

Schultz et al., Journal of the Atmospheric Sciences 63 (2006)

Mammatus clouds: a primer

- Altitude: 2-8 km
- Horizontal extent: O(100m 1km)
- Velocities: 1-3 m/s
- Lobe-like structures are smooth
- Lifetime: O(10 minutes)
- Typically found under cumulonimbus anvils
- Can also be found under
 - volcanic ash clouds
 - jet contrails

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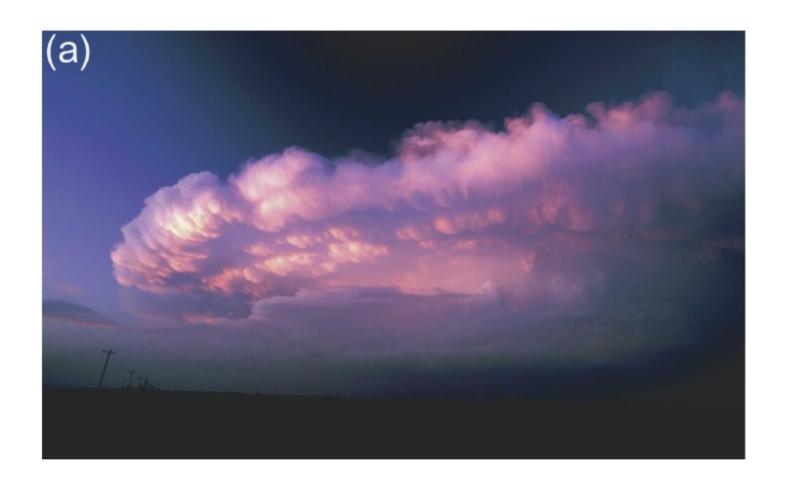
Mammatus clouds: a primer

Several possible mechanisms suggested:

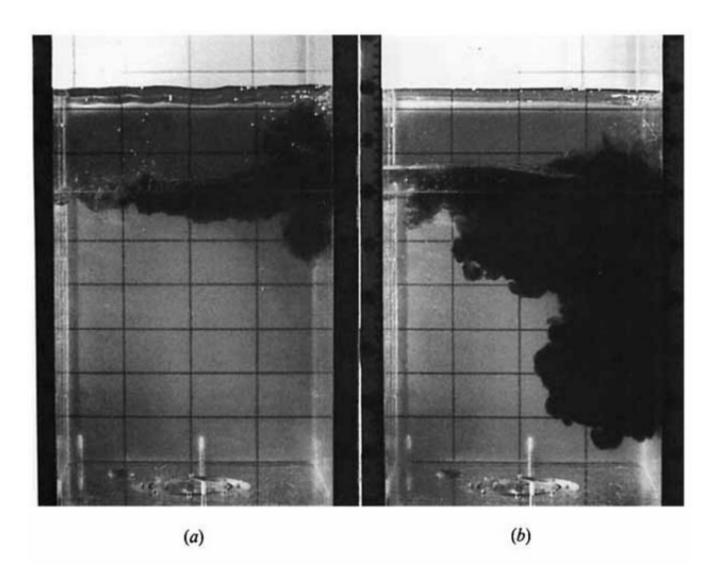
- (re-)circulation associated with cumulonimbus convection
- Radiation
- Gravity waves
- Kelvin-Helmholtz
- Rayleigh-Taylor
- Detrainment at cloud-base
- Settling- or precipitation-driven

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Shear (=>mixing) => Mammatus Clouds



Mixing-driven stratocumulus cloud-top instabilities



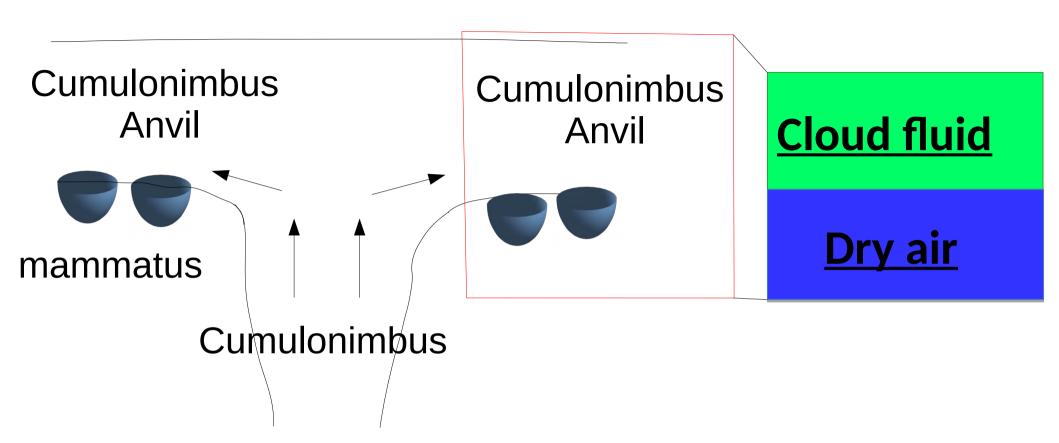
Shy and Breidenthal, J. Fluid Mech. 214 (1990)

Mammatus clouds: a primer

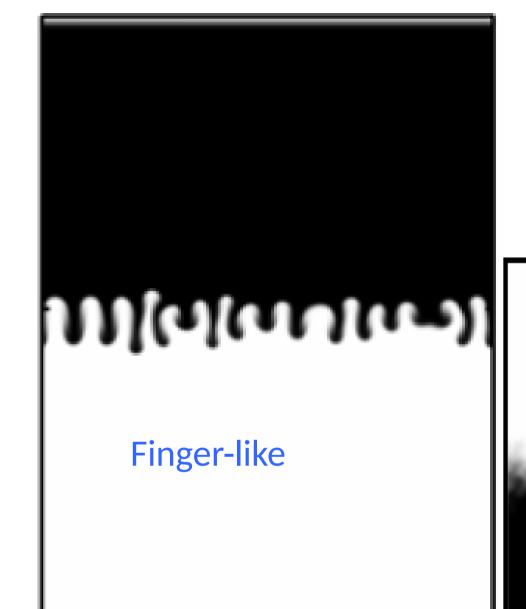
Several possible mechanisms suggested:

- (re-)circulation associated with cumulonimbus convection
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Typical conditions



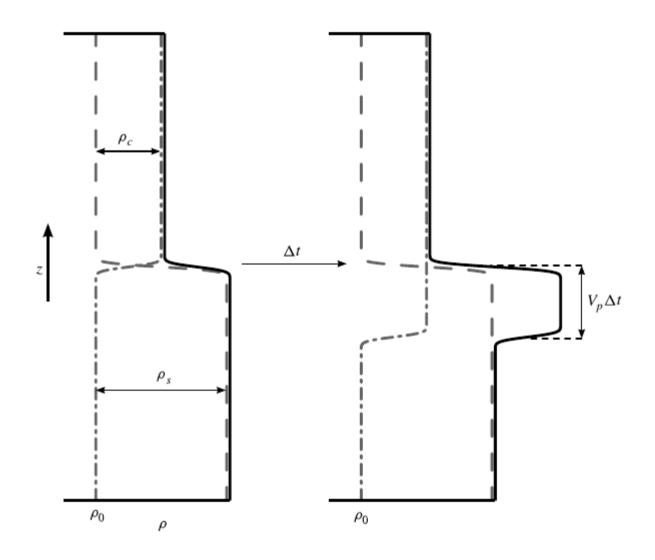
Burns & Meiburg JFM 2012 JFM 2015

String-like

Sediment laden fresh water

Salt water

Density overhang



Silt Vs Salt

Parameter space:

Cloud fluid

Saturated vapour
Droplets of liquid water
Temperature T₀

Dry air Temperature T₀

$$T_0$$
 ΔT

$$r_s^0 = r_s^0 \left(T_0 \right)$$

$$n, a_0 \mid \tau_s r_l$$

$$Fr^2 = \frac{U^2}{g\mathcal{L}\Delta T/T_0} \equiv 1$$

$$Pr_i$$
 i = 1,2,3

$$L_1 = \frac{L_v r_s^0}{C_p \Delta T} \qquad L_2 = \frac{L_v \Delta T}{R_v T_0^2}$$

Model Initial Conditions

Cloud fluid
Saturated vapour
Droplets of liquid water
Temperature T₀

Dry air Temperature T₀

$$\theta = (T - T_0) / \Delta T$$

$$= 0$$

$$\omega = 0$$

$$r_v = \begin{cases} 1 & y > 0 \\ 0 & y < 0 \end{cases}$$

$$r_l = \begin{cases} r_l^0 & \delta y^0 > y > 0 \\ 0 & \text{otherwise.} \end{cases}$$

$$Fr^2 \equiv 1$$
 $Re Pr_i$

Thermodynamic parameters

$$\theta = \left(T - T_0\right) / \Delta T$$

$$r_s(T) \approx \exp(L_2\theta)$$

$$L_2 = \frac{L_v \Delta T}{R_v T_0^2}$$

$$L_1 = \frac{L_v r_s^0}{C_p \Delta T}$$

$$T_0$$

$$T_0 \qquad r_s^0 = r_s^0 \left(T_0 \right)$$

$$\Delta T$$

Mechanical parameters

$$r_l^0 \sim na_0^3$$

$$v_p \sim a_0^2$$

(n, a₀)

$$\tau_s \sim (na_0)^{-1}$$

Nondimensional Equations

$$\frac{D\omega}{Dt} = \frac{1}{Fr^2} \frac{\partial}{\partial x} \theta + \frac{1}{Re} \nabla^2 \omega$$

$$\frac{D\theta}{Dt} = \frac{1}{Re \cdot Pr_1} \nabla^2 \theta - L_1 E$$

$$E = -\mathcal{H} \frac{1}{\tau_s} \left(\frac{r_v}{r_s} - 1 \right)$$

$$\frac{Dr_v}{Dt} = \frac{1}{Re \cdot Pr_2} \nabla^2 r_v + E$$

$$\frac{Dr_l}{Dt} = v_p \frac{\partial r_l}{\partial y} + \frac{1}{Re \cdot Pr_3} \nabla^2 r_l - E$$

$$E = -\mathcal{H}\frac{1}{\tau_s} \left(\frac{r_v}{r_s} - 1 \right)$$

$$\mathcal{H} = \begin{cases} 1 & \text{saturated, or} \\ 0 & \text{unsaturated.} \end{cases}$$

Clausius-Clapeyron $r_s\left(T\right) = r_s^0 exp\left(\frac{L_v}{R_v}\left(\frac{1}{T_0} - \frac{1}{T}\right)\right) \approx \exp\left(L_2\theta\right)$

Linear Stability, 1D Nonlinear evolution

- Standard approach: linearise the equations, including the thermodynamics
- Problem: frozen-flow approximation is wrong.
 Instability time-scale ~ phase-change time-scale
- Our approach: evolve the (~)equations in 1D, get the density profile, do the stability on the net density profile

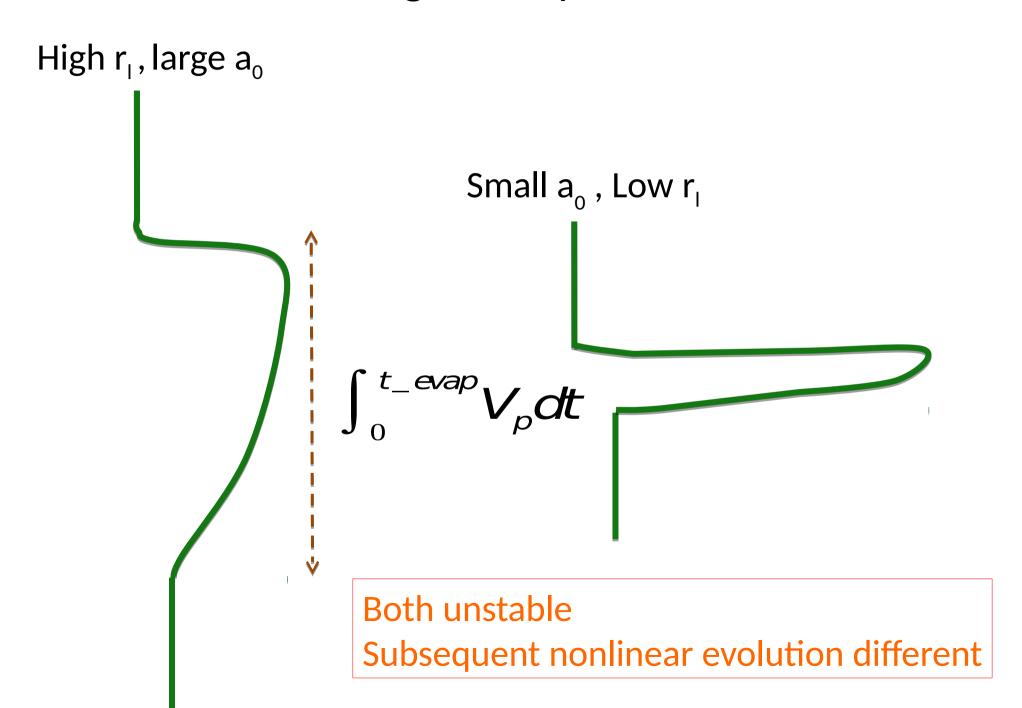
1D Nonlinear evolution

ullet Allows for parametric study of $\left(r_l^0,\delta y^0, au_s,v_p
ight)$

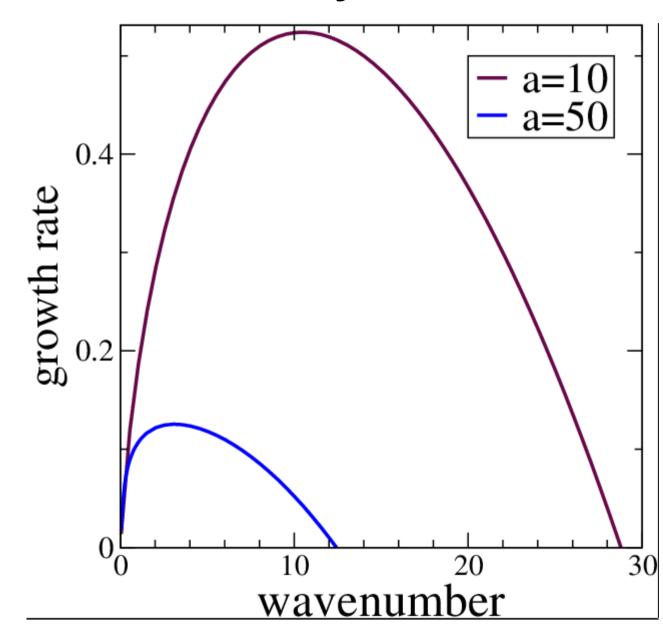
$$egin{align} rac{\delta y}{\delta y^0} &= \delta y \left(rac{\delta y^0}{v_p au_s}, r_l^0
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$$t_{evap} \sim \tau_s \delta y / v_p$$

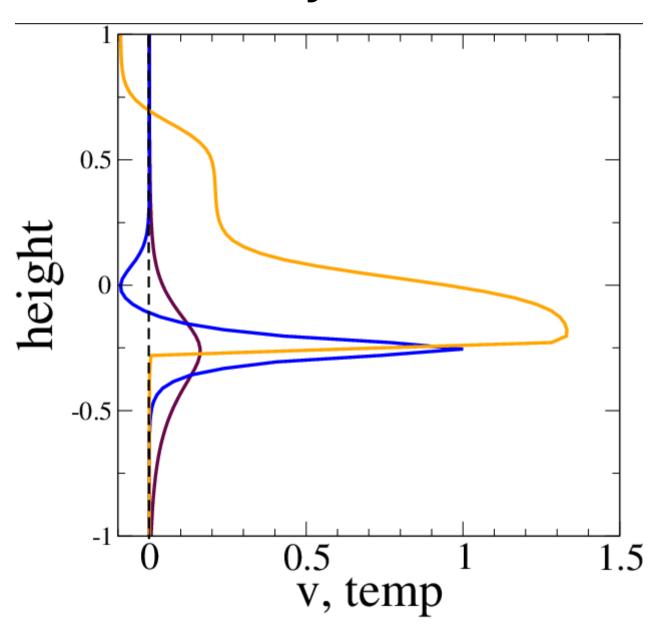
Evolving density contours



Stability results



Stability results



2D Numerical Simulations

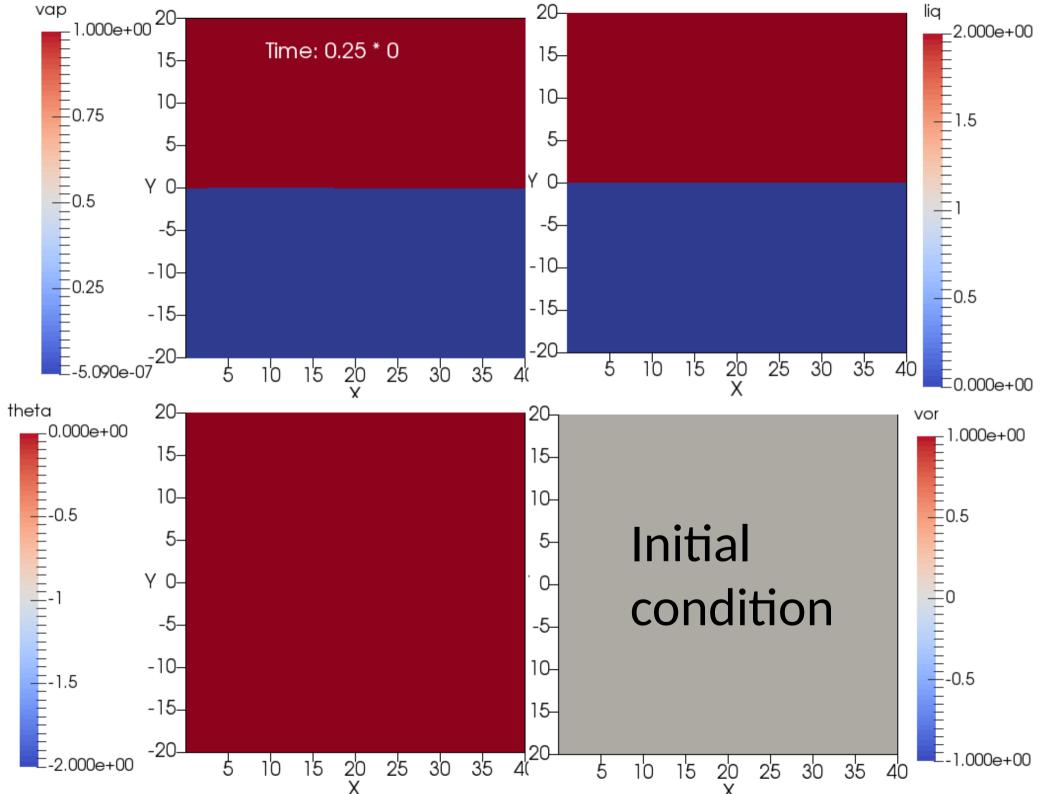
Message:

Settling velocity matters more than liquid density

$$v_p \sim a_0^2$$

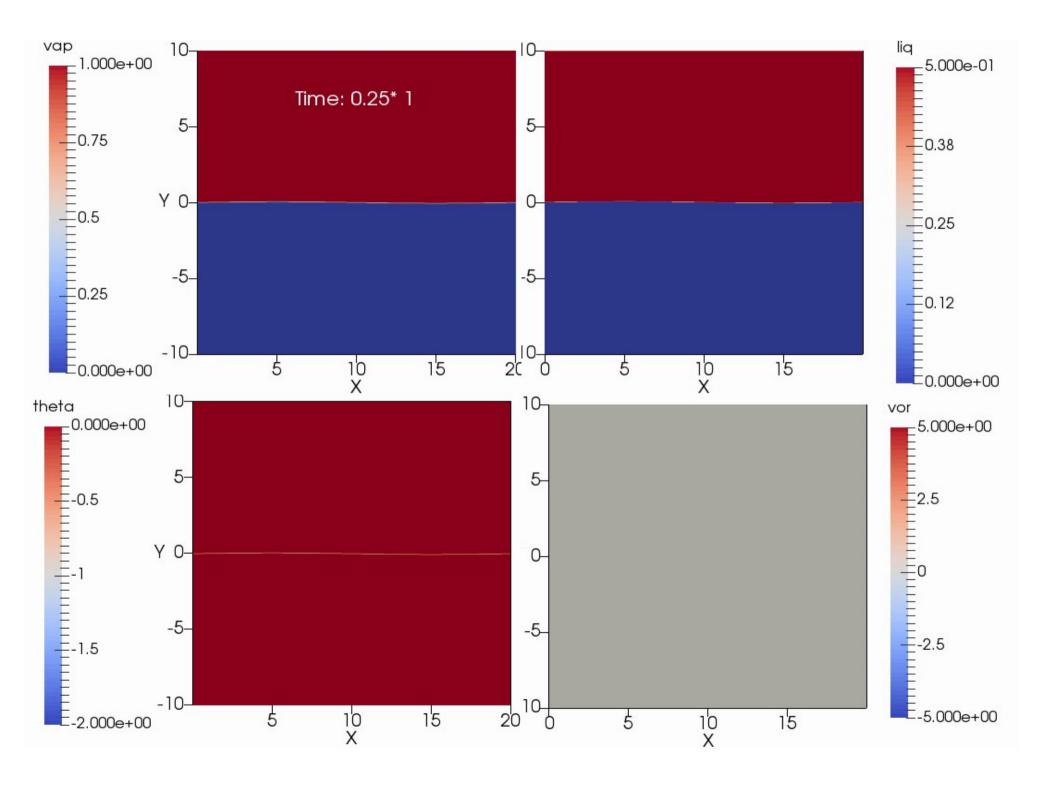
$$r_l^0 \sim na_0^3$$

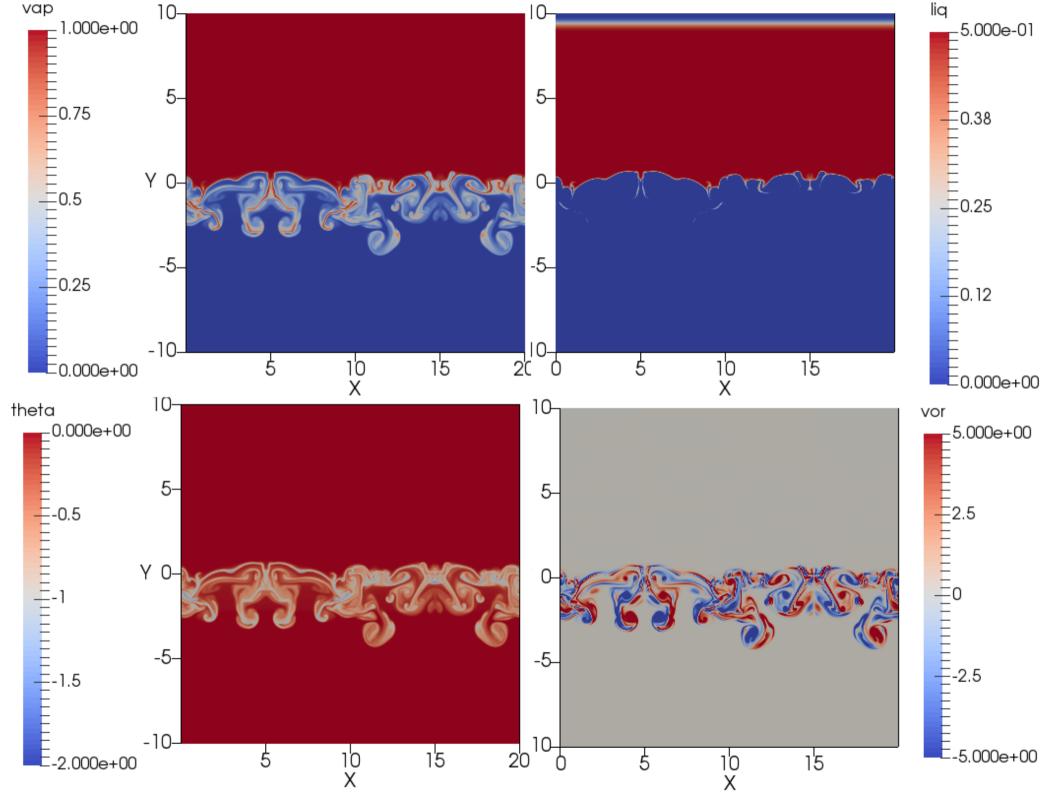
Simulations: small/large r_1^0 , small/large a_0



Small r_1^0 , small a_0

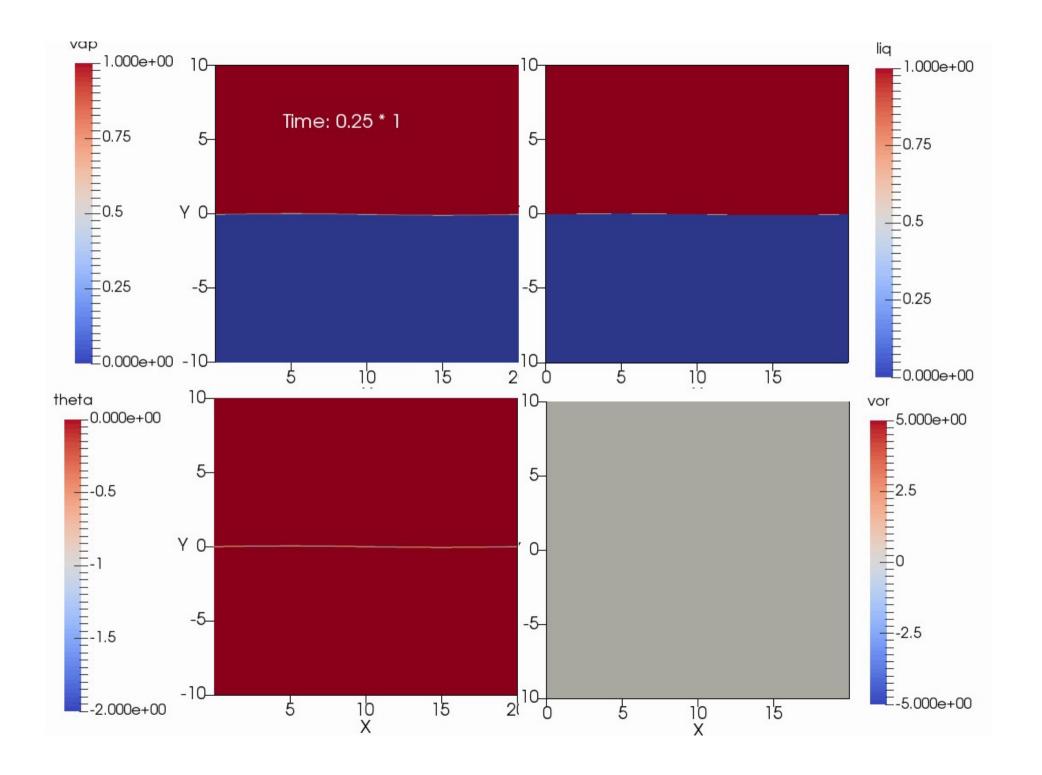
$$r_1^0 = 0.5$$
, $a_0 = 10 \mu$

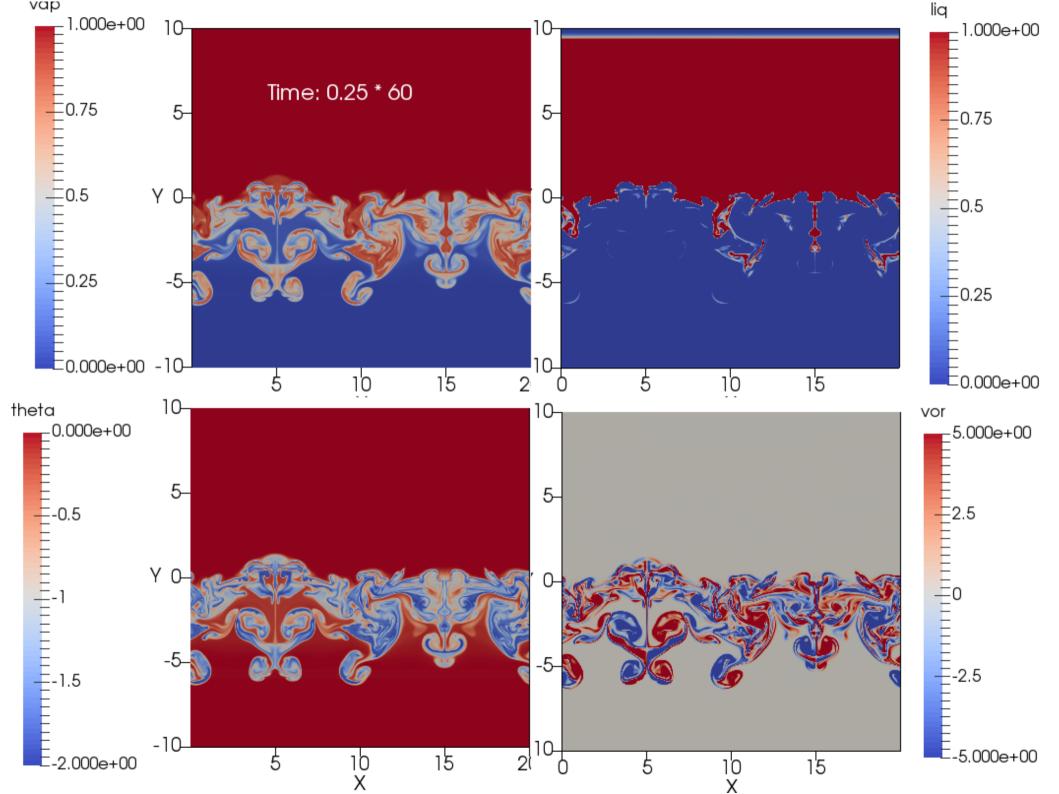




large r_1^0 , small a_0

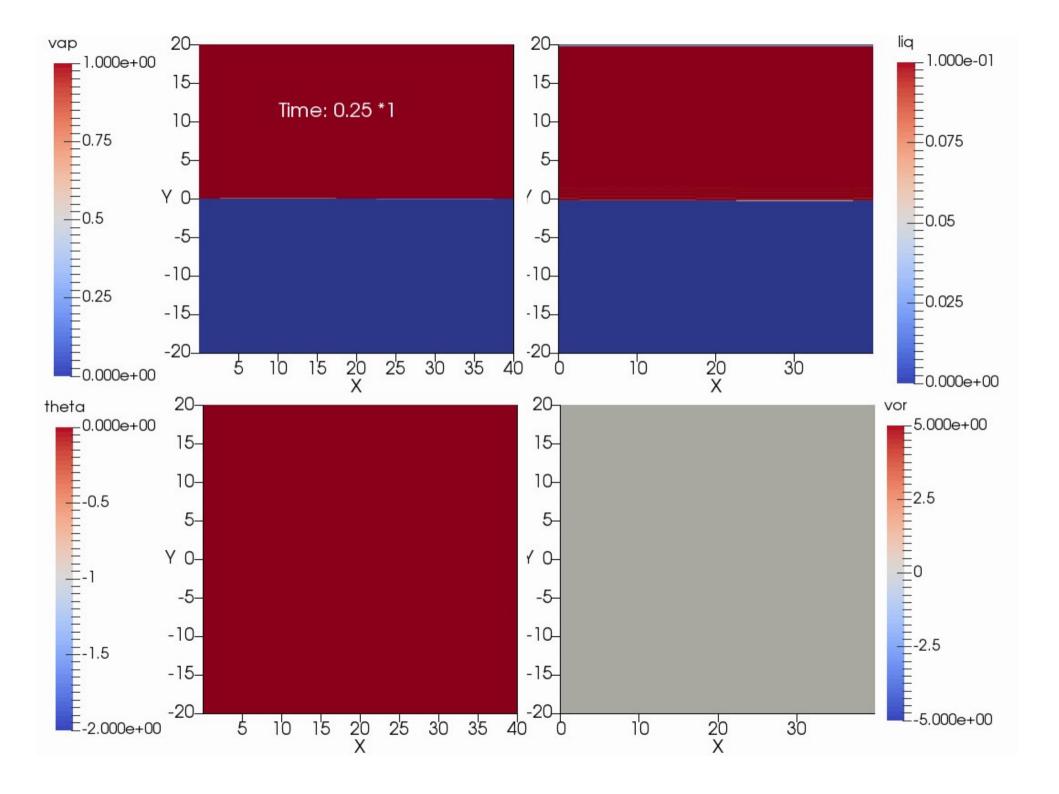
$$r_1^0 = 2.0$$
, $a_0 = 10 \mu$





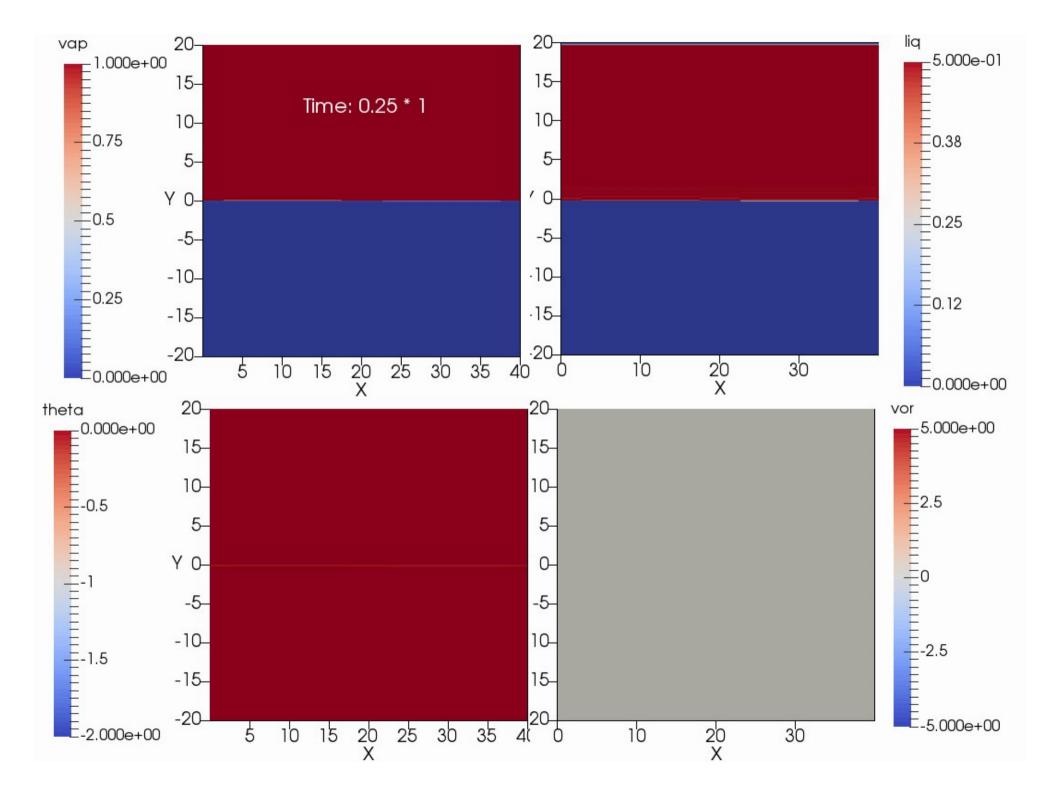
Very Small r₁°, large a₀

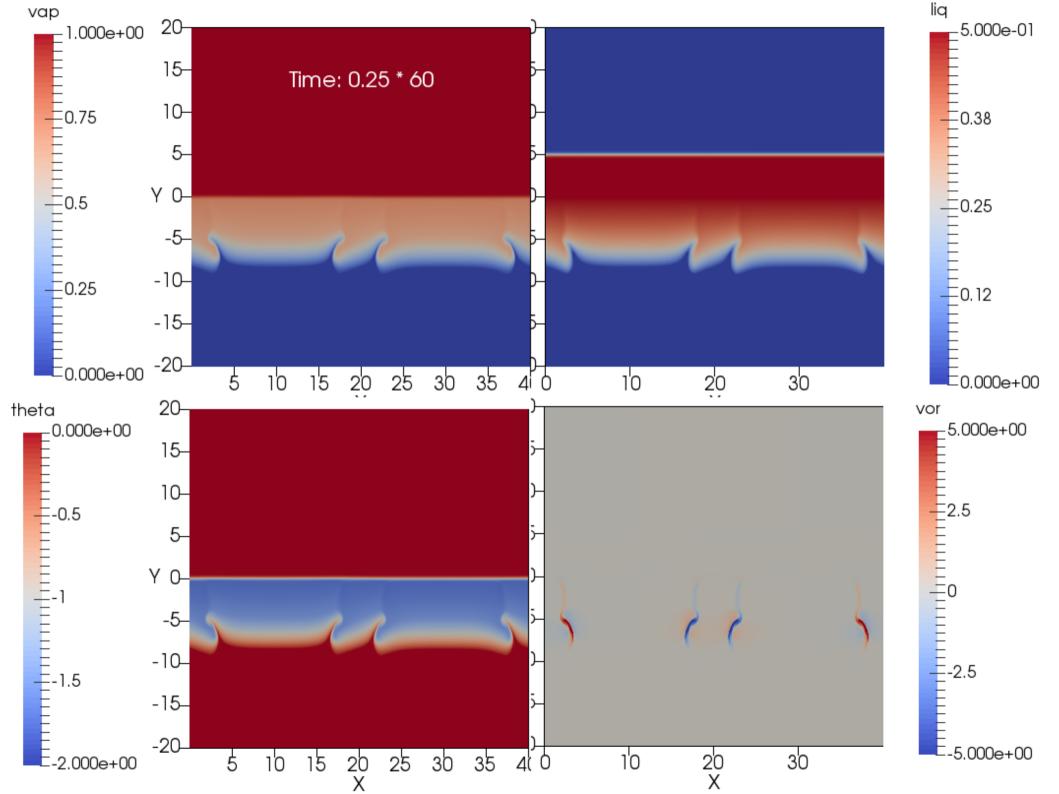
$$r_1^0 = 0.1$$
, $a_0 = 50 \mu$



Small r₁⁰, large a₀

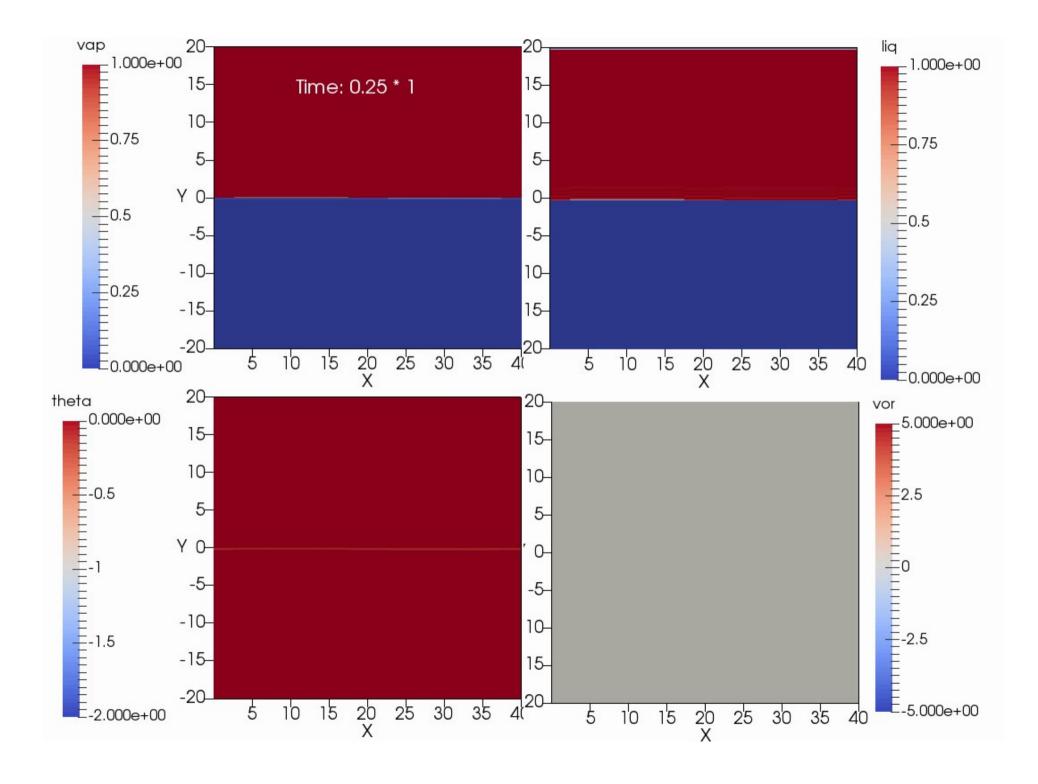
$$r_1^0 = 0.5$$
, $a_0 = 50 \mu$

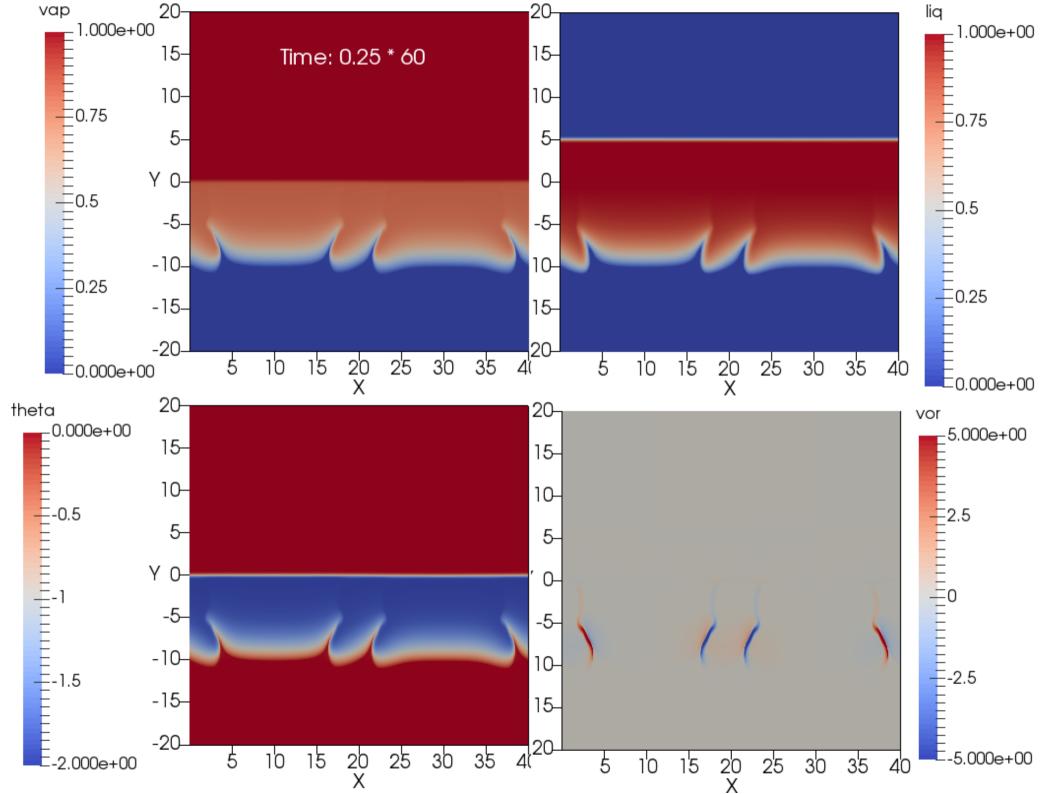




Large r₁°, large a₀

$$r_1^0 = 1.0$$
, $a_0 = 50 \mu$





Message:

Settling velocity matters more than liquid density

Summary

A minimal model for mammatus clouds

- More to be done:
 - Shear/Mixing ("detrainment instability"): Asperitas?
 - 3D simulations
 - Compare predictions with observations

20/04/18