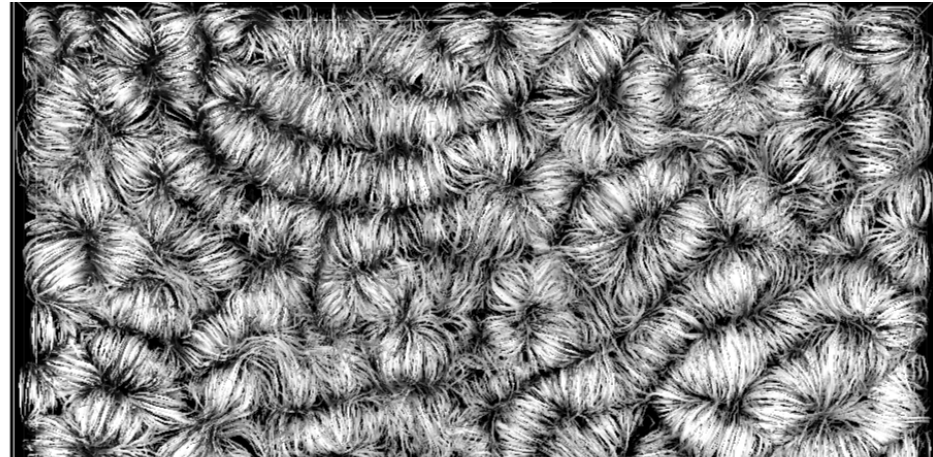
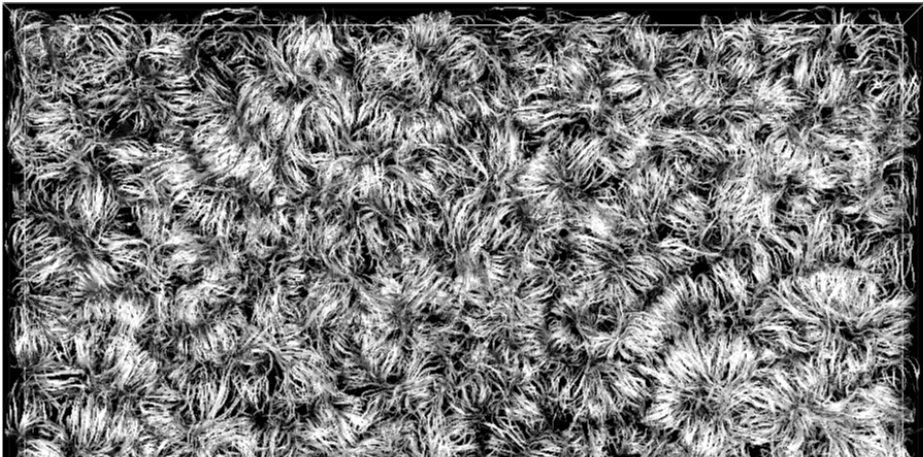


Turbulent superstructures in Rayleigh-Bénard convection

Jörg Schumacher
Technische Universität Ilmenau
Germany

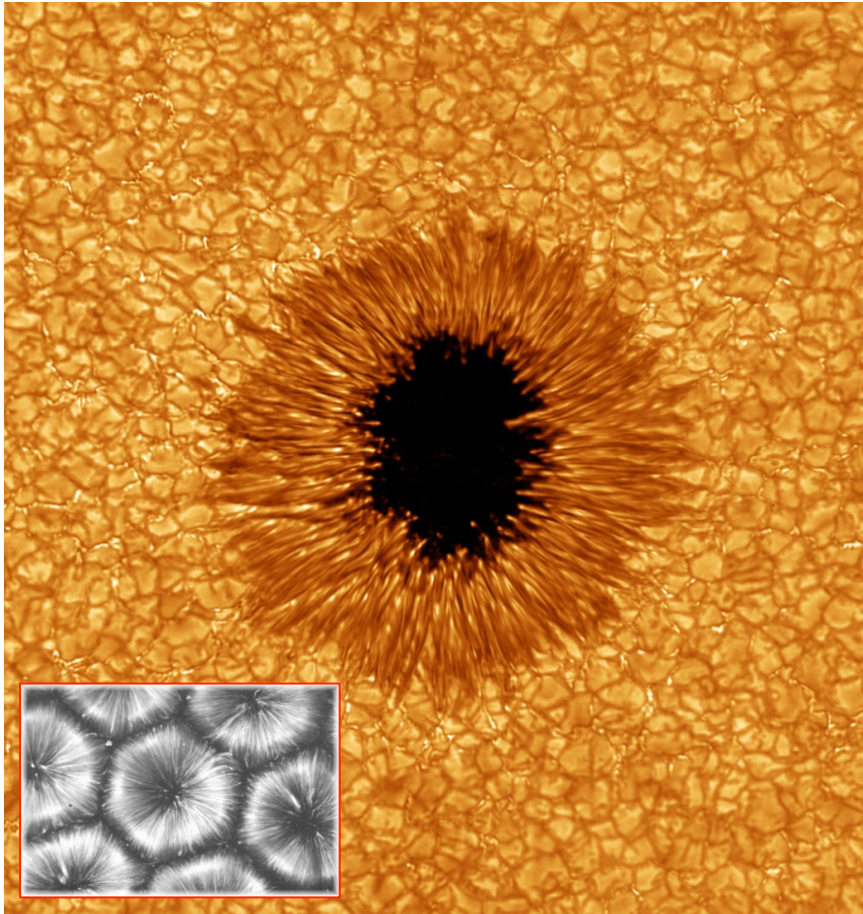
Martin Ender (TU München)
Kathrin Padberg-Gehle (Leuphana Universität Lüneburg)
Ambrish Pandey (TU Ilmenau)
Janet D. Scheel (Occidental College Los Angeles)
Christiane Schneide (Leuphana Universität Lüneburg)
Rüdiger Westermann (TU München)



Turbulent convection in nature

Solar Granulation & Supergranulation

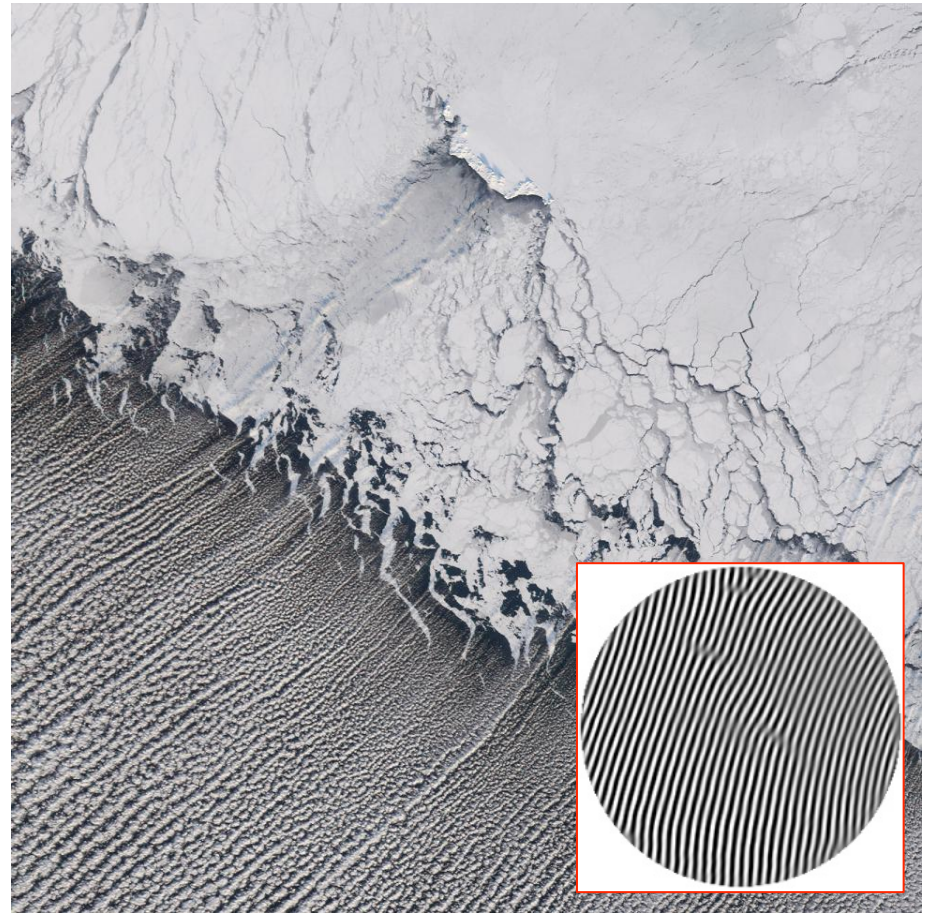
$$Ra \gtrsim 10^{22} \quad Pr \lesssim 10^{-3}$$



Big Bear Solar Observatory

Cloud Streets over the Bering Sea

$$Ra \sim 10^{18} \quad Pr = 0.7$$



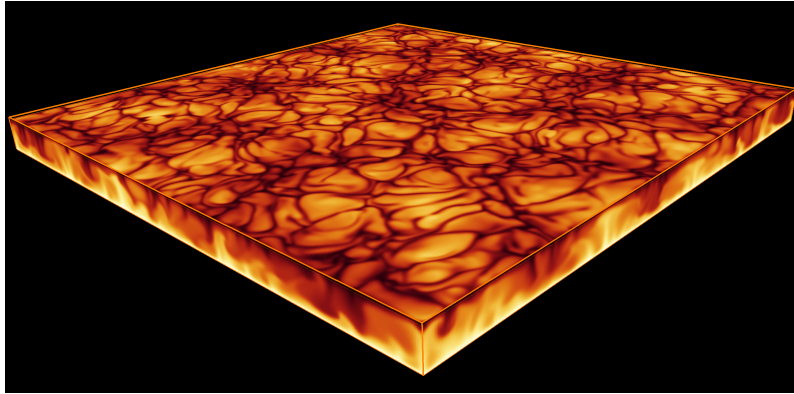
NASA Modis Mission

Questions

- Do large-scale patterns exist in turbulent Rayleigh-Bénard convection (RBC)?
- Are these patterns relics from the onset or the weakly nonlinear regime of RBC?
- What are the typical spatial and temporal scales of these patterns and how do they depend on Ra and Pr ?
- Are the pattern scales related to the boundary layer dynamics?
- Can we develop simplified models for the large-scale patterns?
- What is their role for turbulent transport?

Numerical simulations

Fischer, J. Comp. Phys. 1997; Emran, Scheel & JS, New J. Phys. 2013



- 3d Boussinesq equations in closed rectangular cell
- 25:25:1
 $0.005 \leq Pr \leq 70$
 $2300 \leq Ra \leq 10^7$
- Spectral element method Nek5000
Small-scale gradients at walls



Jülich

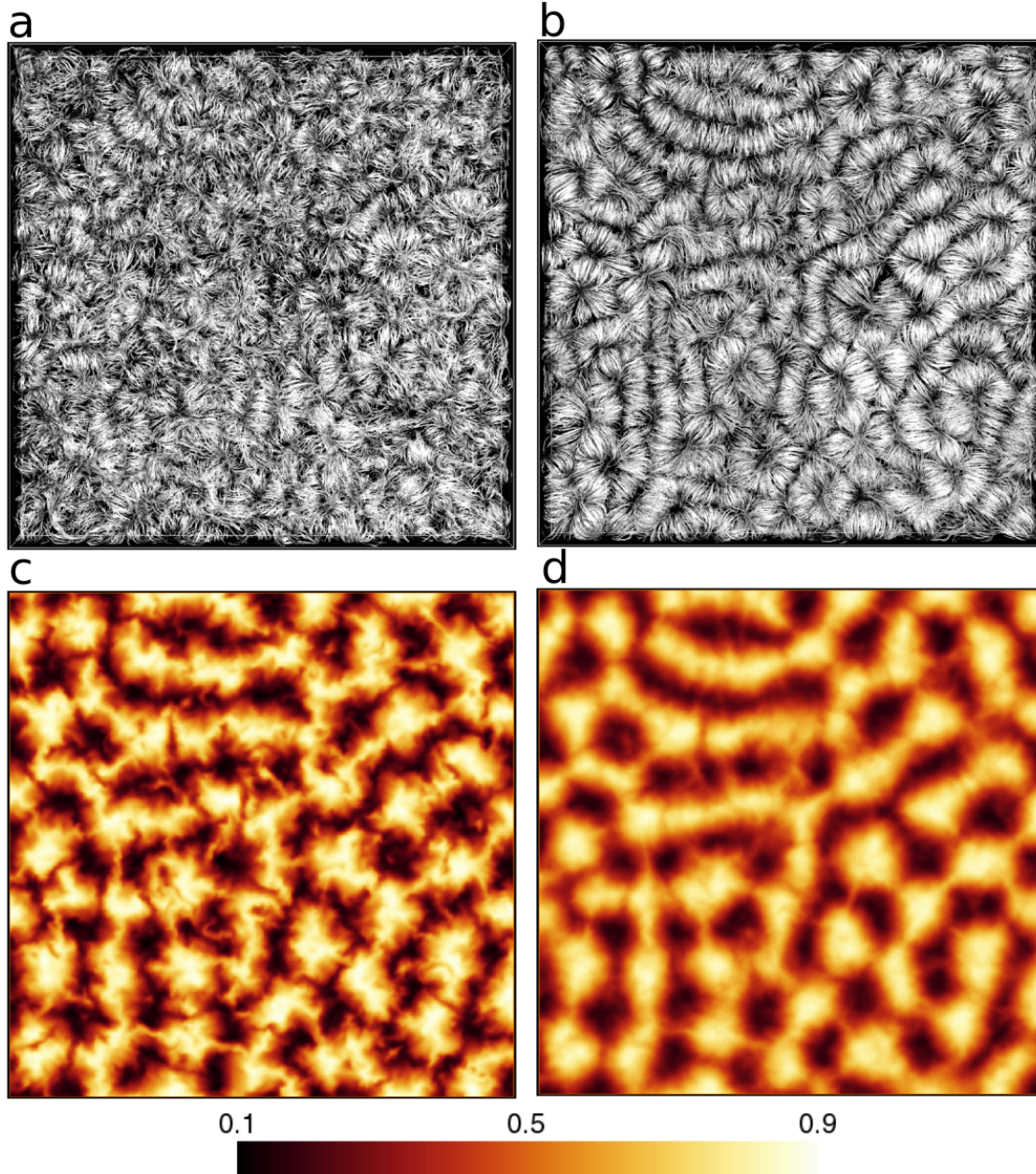


Garching

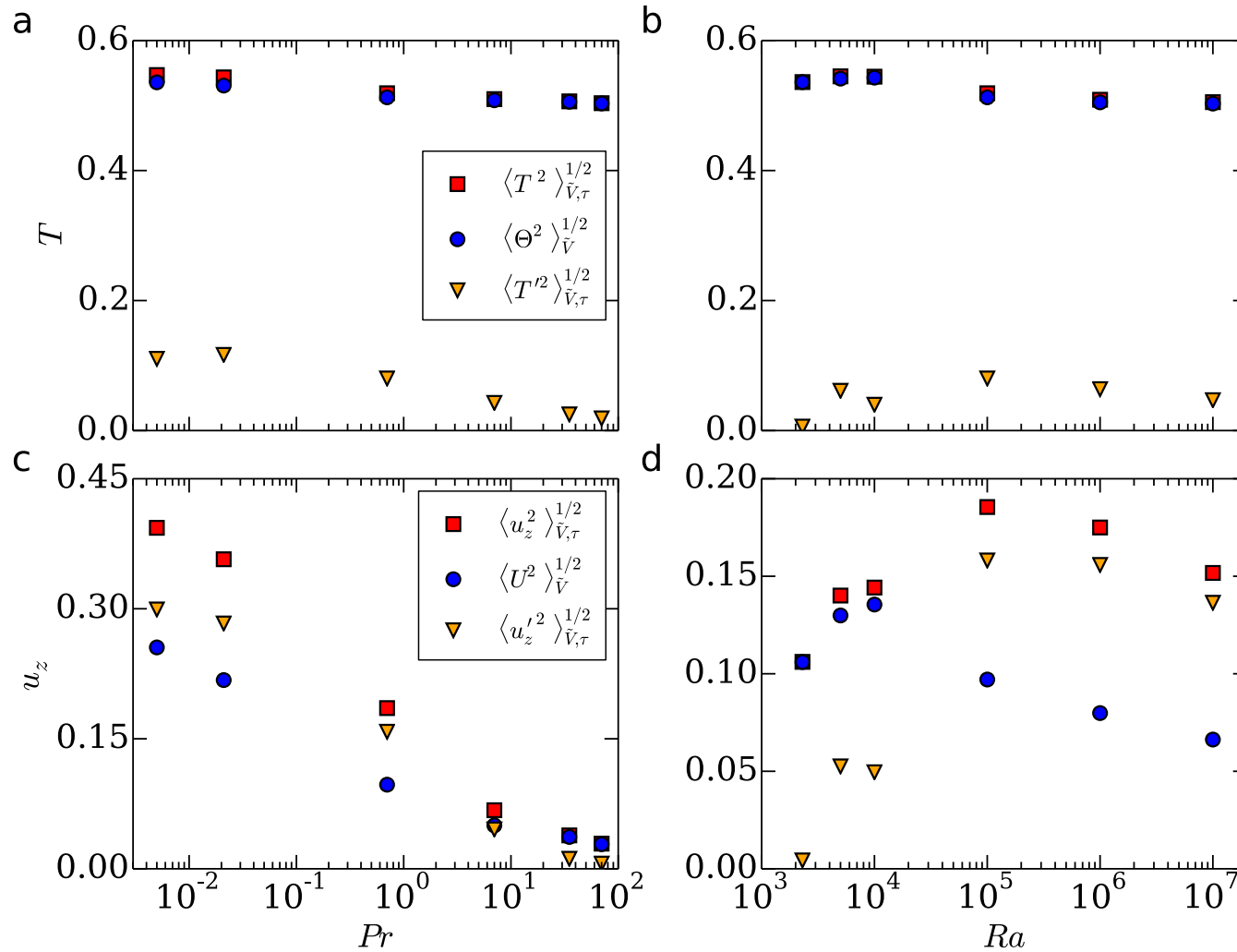
rendered on SuperMUC by LRZ

Instantaneous fields at $Ra=10^5$

$Pr=0.021$



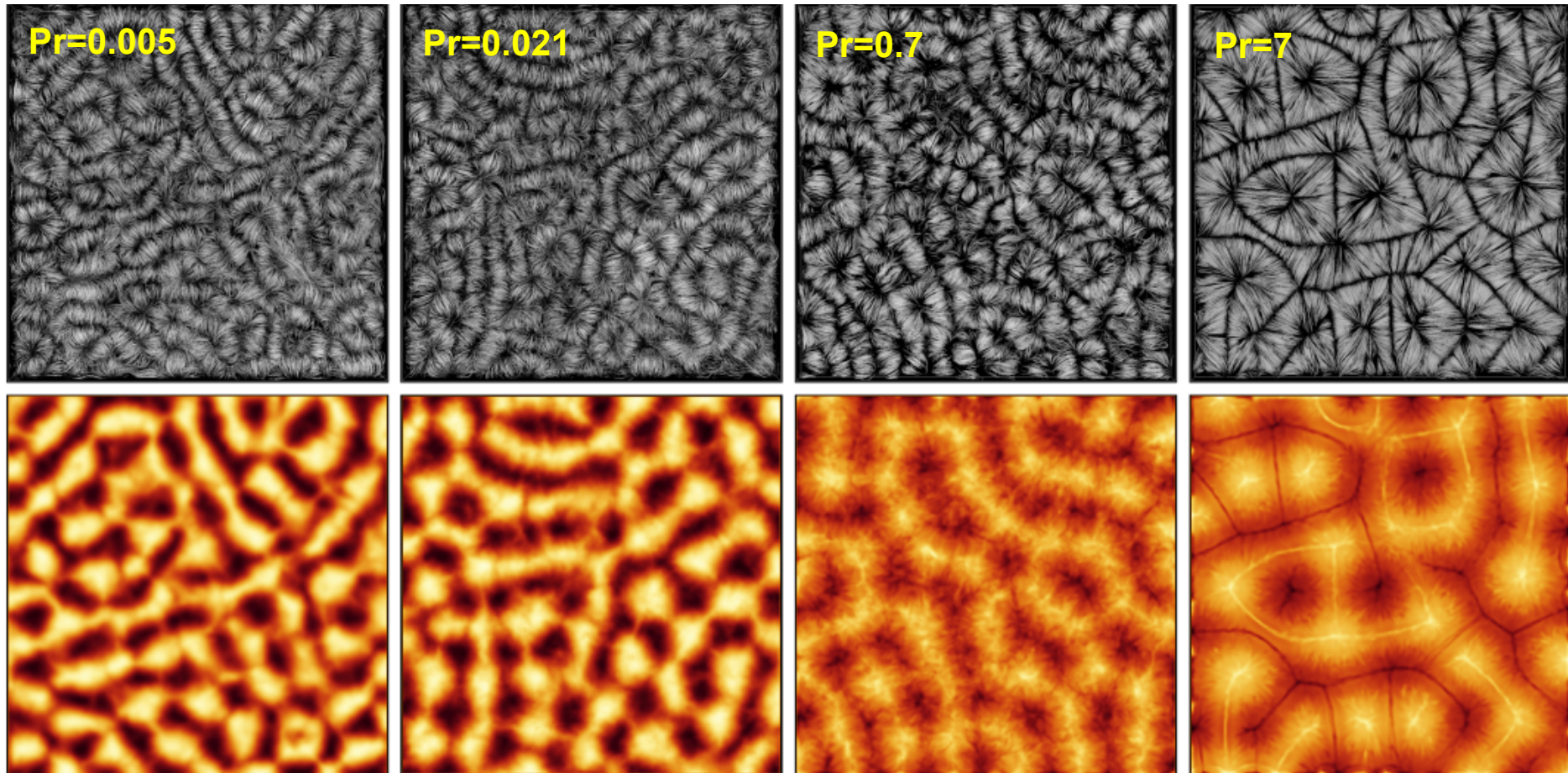
Magnitude of superstructures



$$u_z(\mathbf{x}, t) = U(\mathbf{x}) + u'_z(\mathbf{x}, t)$$

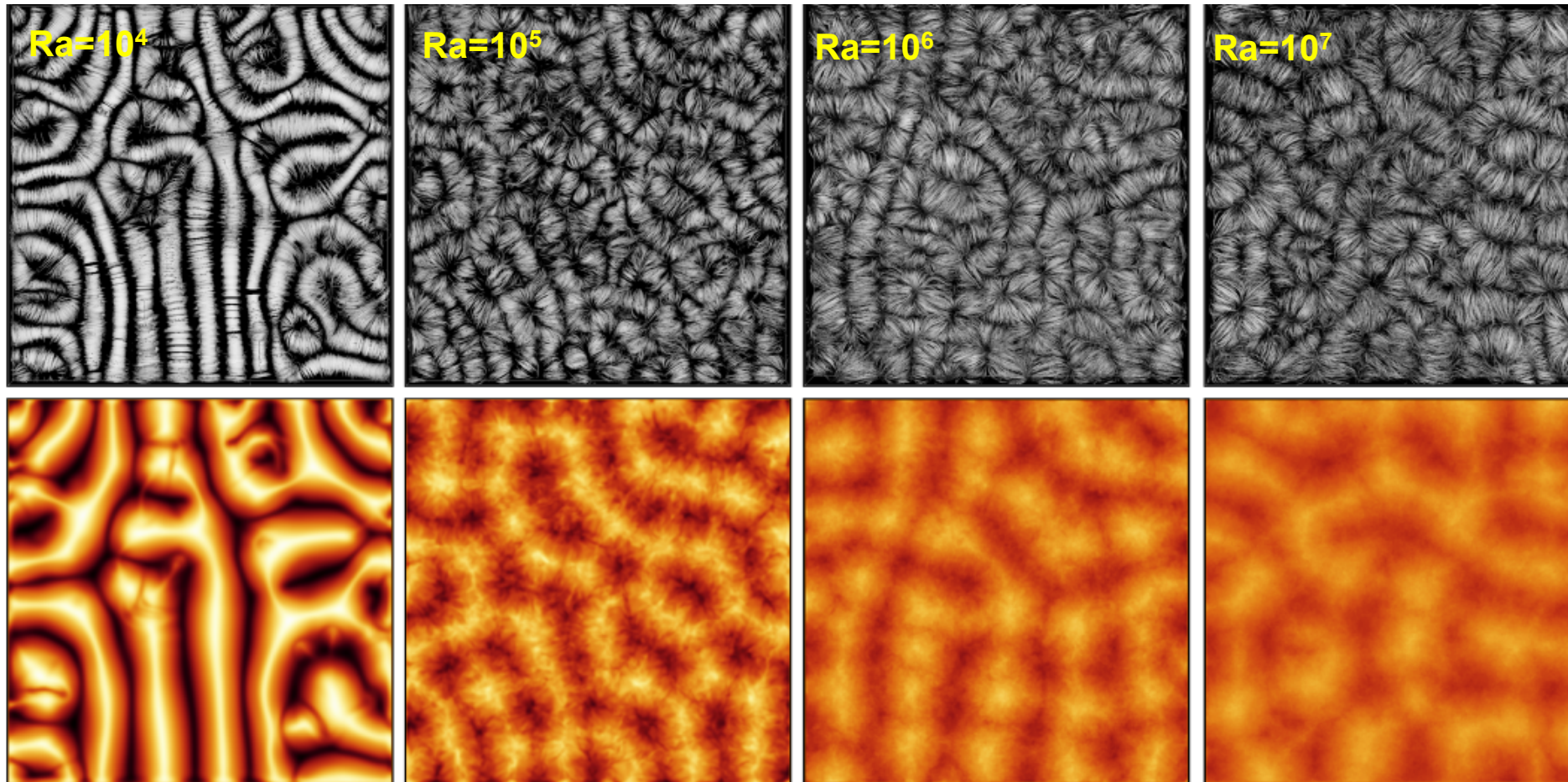
$$T(\mathbf{x}, t) = \Theta(\mathbf{x}) + T'(\mathbf{x}, t)$$

Time-averaged fields at $Ra=10^5$



Regular patterns (= [turbulent superstructures](#)) are detected for all Pr

Time-averaged fields at $Pr=0.7$

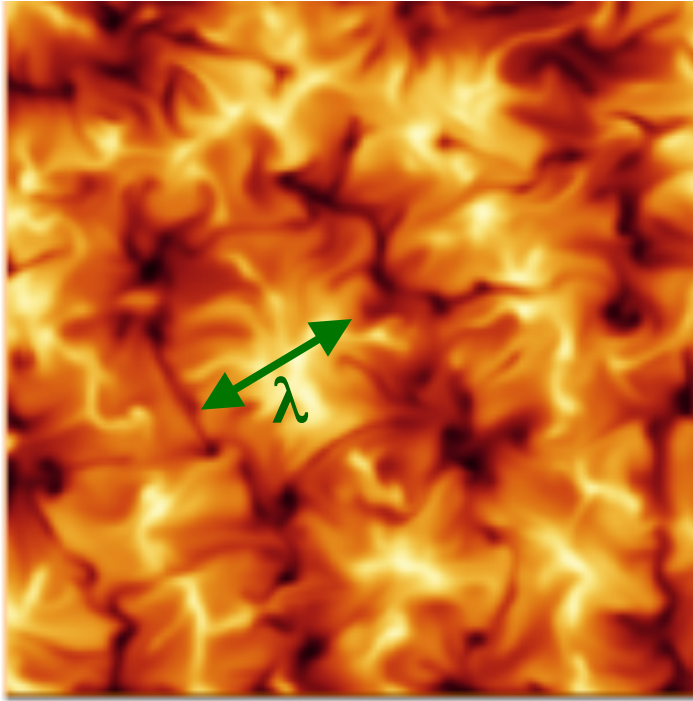


Regular patterns (= **turbulent superstructures**) are detected for all Ra

What are the typical scales and how do they depend on Ra , Pr ?

Scale separation

Fast & Small Scales

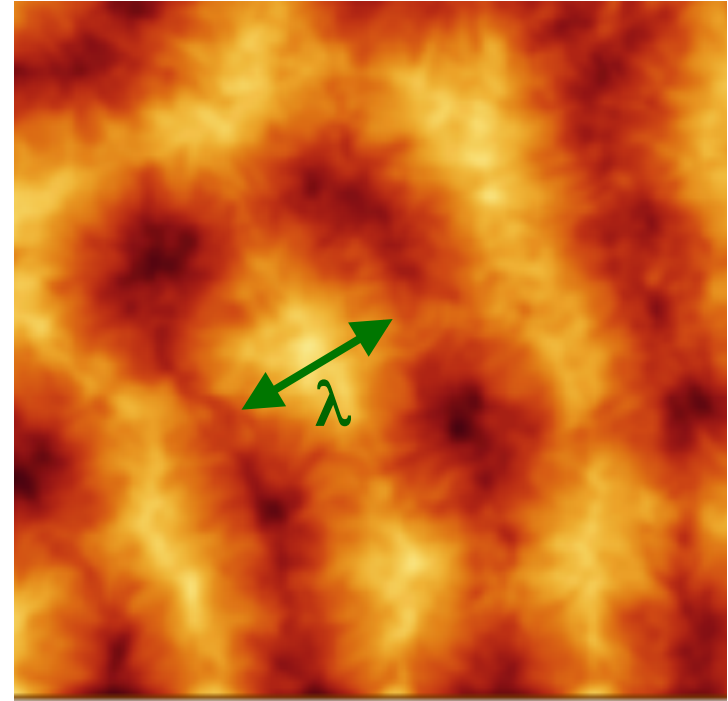


$t < \tau \quad \ell < \lambda$



$\langle \cdot \rangle_\tau$

Slow & Large Scales



$t > \tau \quad \ell > \lambda$

$$U(x, y; \tau, t_0) = \frac{1}{\tau} \int_{t_0 - \tau/2}^{t_0 + \tau/2} u_z(x, y, z = 1/2, t') dt'$$
$$\Theta(x, y; \tau, t_0) = \frac{1}{\tau} \int_{t_0 - \tau/2}^{t_0 + \tau/2} \theta(x, y, z = 1/2, t') dt'.$$

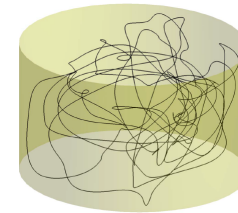
Time scales in RB convection

- Free-fall time $T_f = \frac{H}{U_f} = \frac{H}{\sqrt{g\alpha\Delta TH}}$

- Diffusion time $t_d = \frac{H^2}{\kappa} = \sqrt{RaPr} T_f \longrightarrow t_d^\perp = \frac{(\Gamma H)^2}{\kappa} = \Gamma^2 t_d$

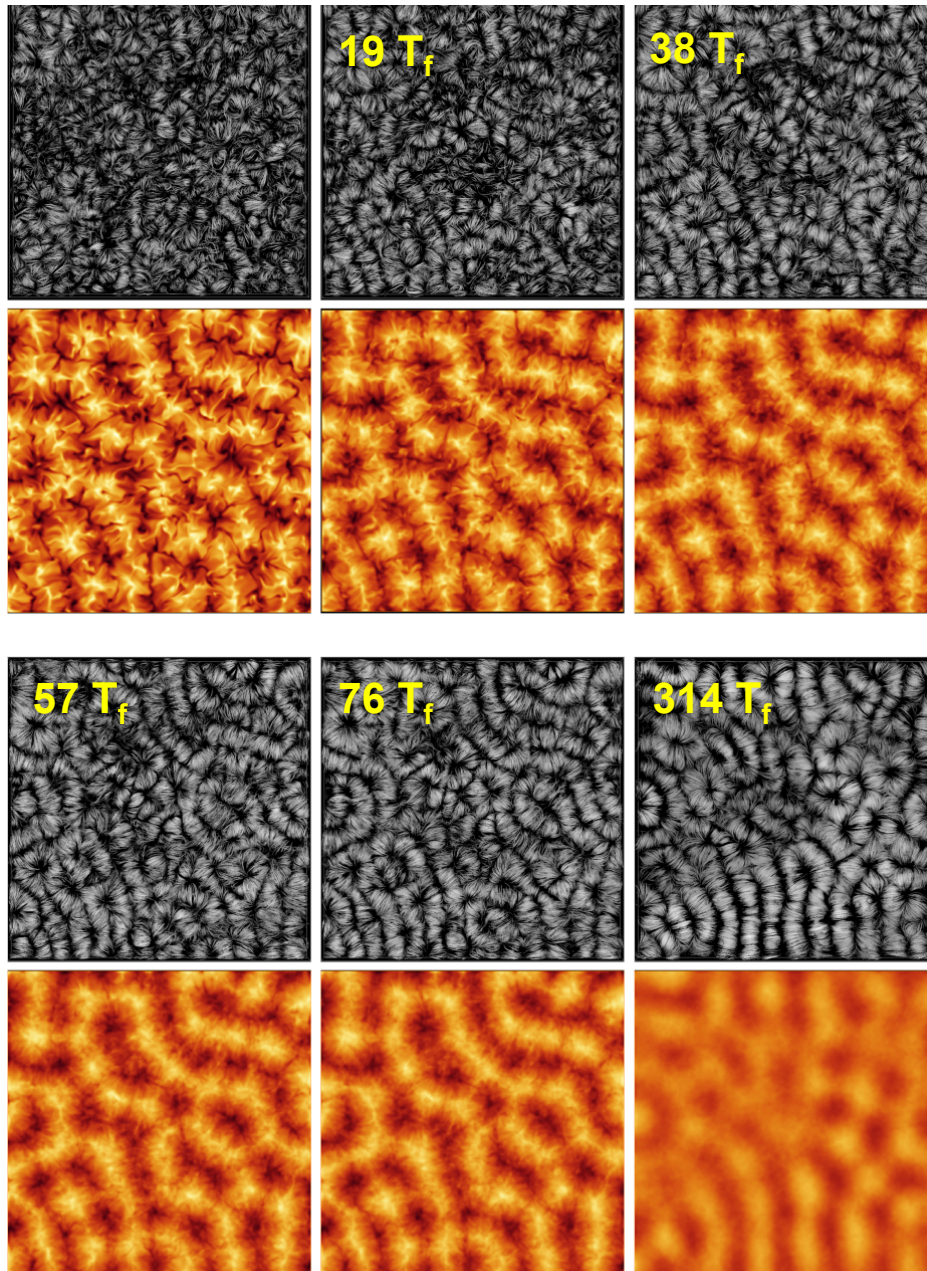
- Viscous time $t_v = \frac{H^2}{\nu} = \sqrt{\frac{Ra}{Pr}} T_f$

- Lagrangian turnover time $10 - 20 T_f$
(Emran & JS, Phys. Rev. E 2010)



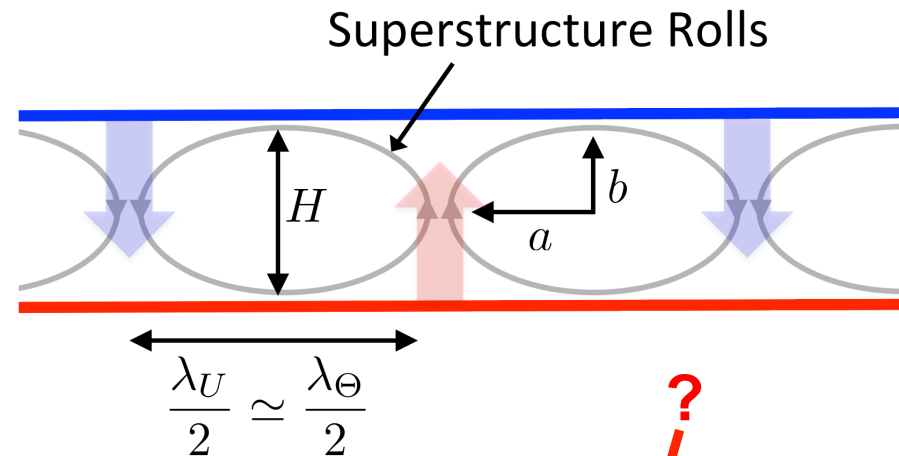
Pr at Ra=10 ⁵	T _f	t _d	t _v	t _g = Ra ^{1/2} T _f
0.005	1	22	4472	316
0.021	1	46	2182	316
0.7	1	265	378	316
7	1	837	120	316
70	1	2646	38	316

What is the appropriate time interval?



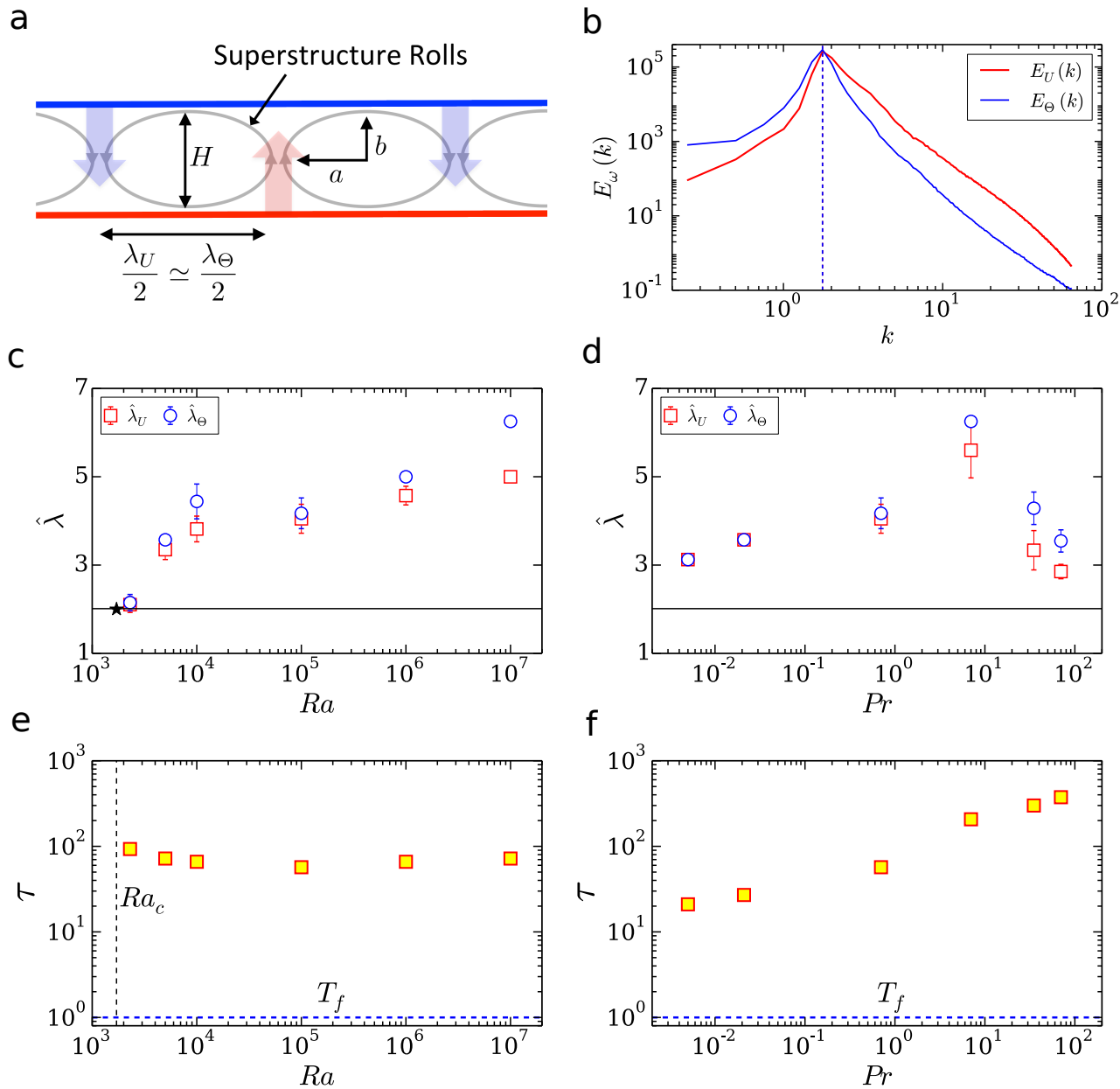
Mean turnover time τ in a superstructure roll

$$T_f \ll \tau \ll \max(t_v, t_d)$$

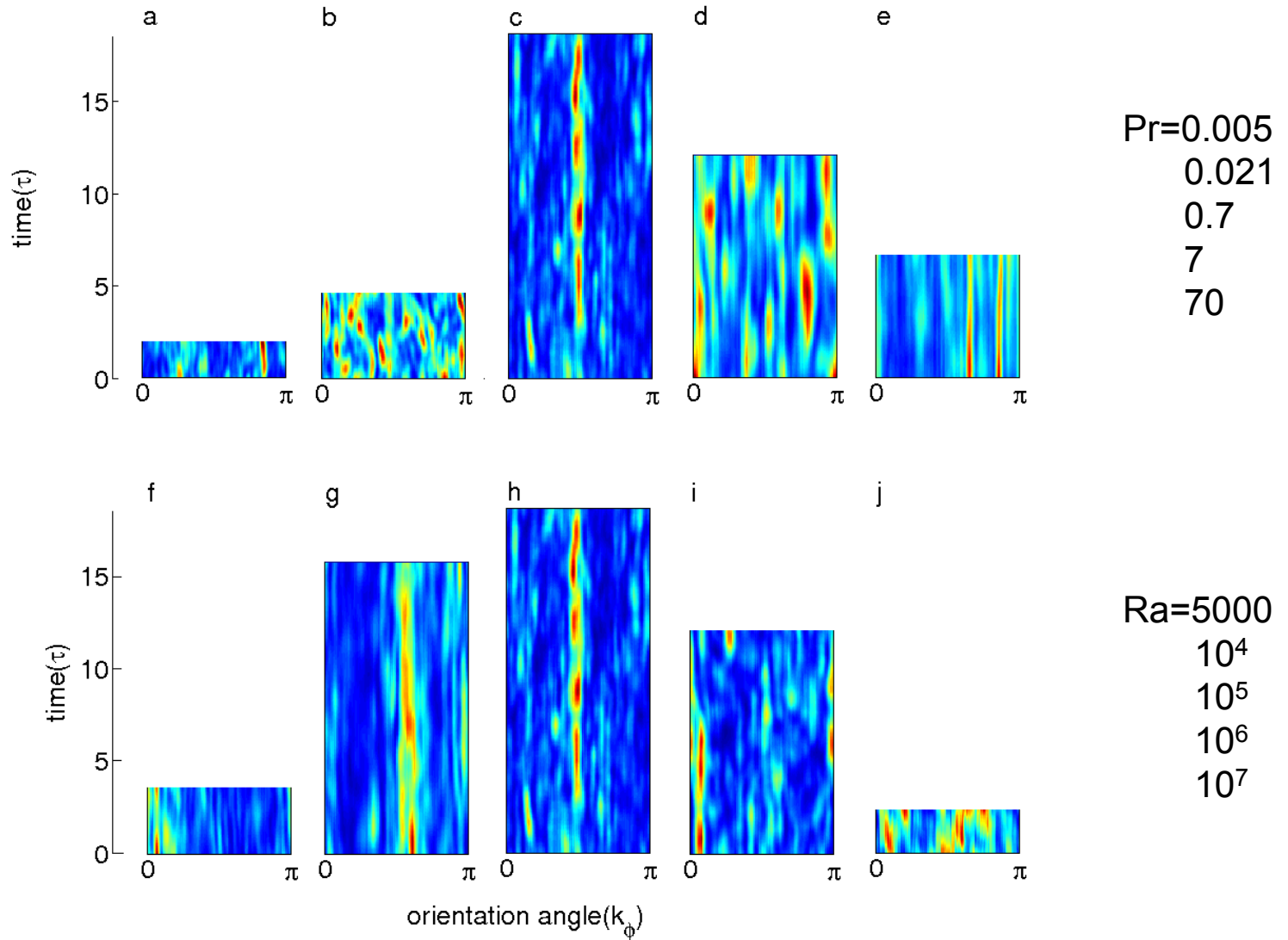


$$\tau(Ra, Pr) \sim \frac{\ell}{u_{rms}} = \frac{\pi \left(\frac{1}{4} \lambda_U + \frac{1}{2} H \right)}{\langle u_i^2 \rangle_{V,t}^{1/2}}$$

Pattern scales of u_z and θ in midplane



Slow evolution of superstructures

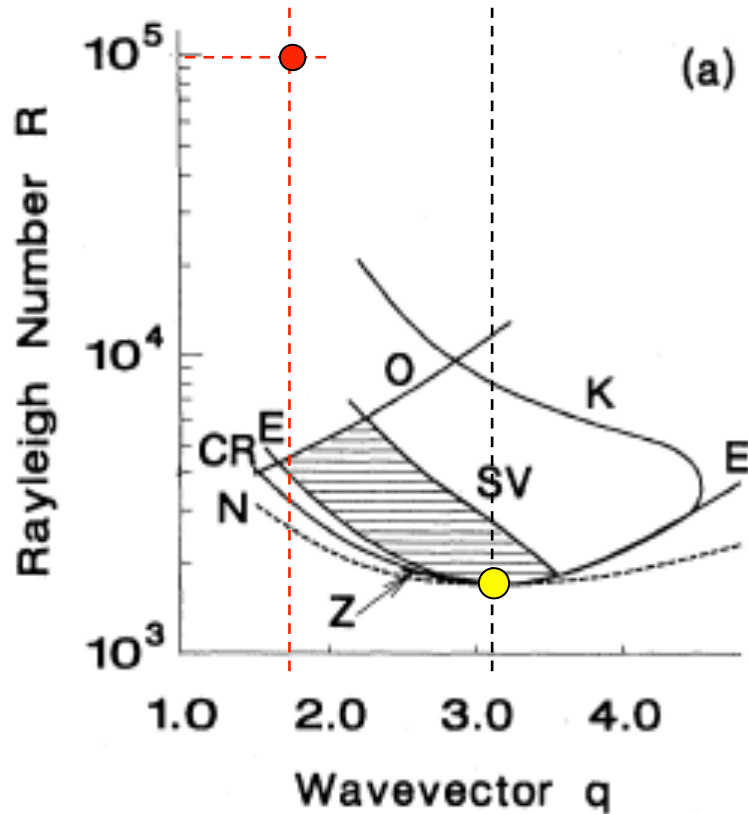


$$E_\Theta(k_\phi; \tau, t_0) = \frac{1}{k_m} \int_0^{k_m} |\hat{\Theta}(k, k_\phi; \tau, t_0)|^2 dk$$

Busse balloon

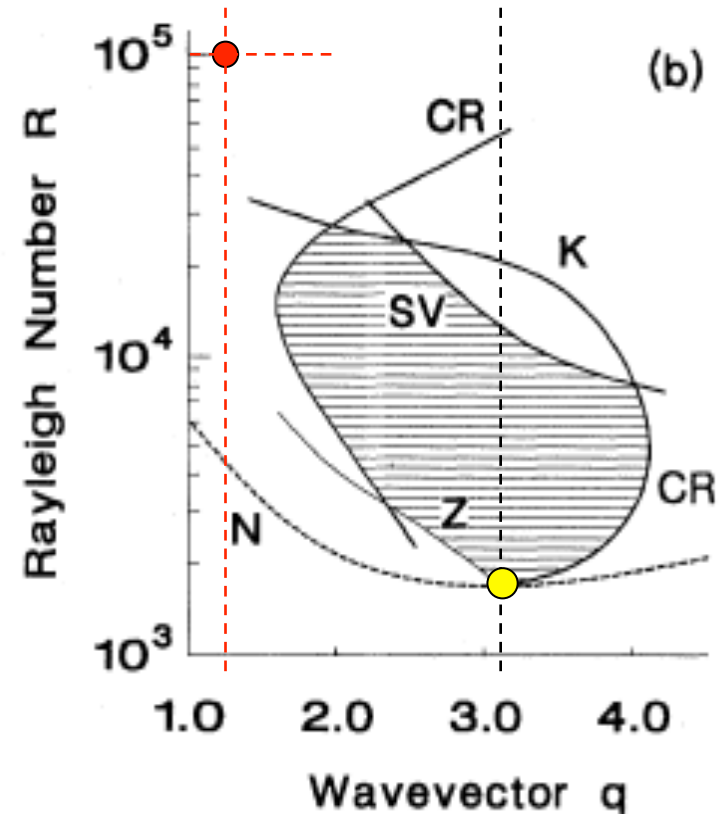
Busse, *Rep. Prog. Phys.* 1978; Cross & Hohenberg, *Rev. Mod. Phys.* 1993

$Pr=0.71$



Knot or oscillating
instability

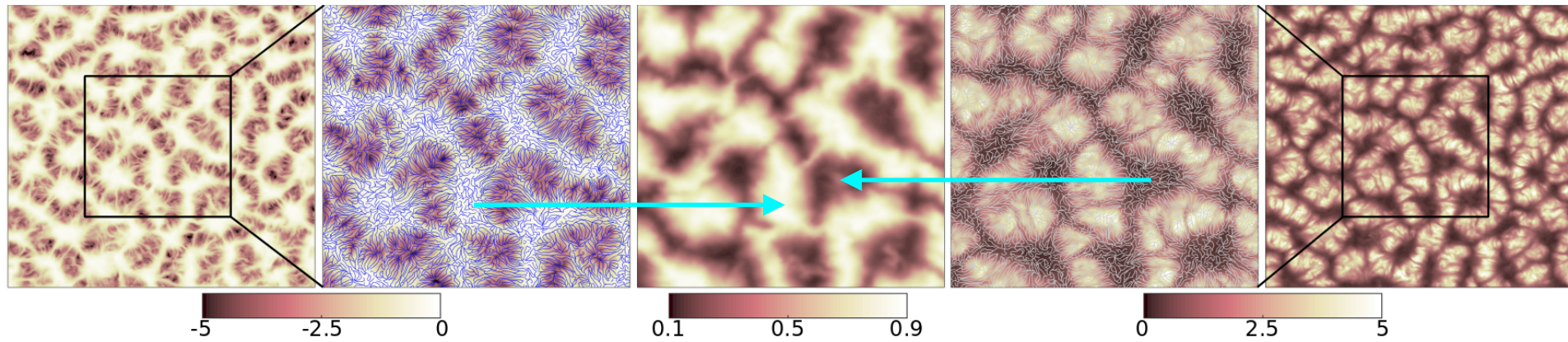
$Pr=7$



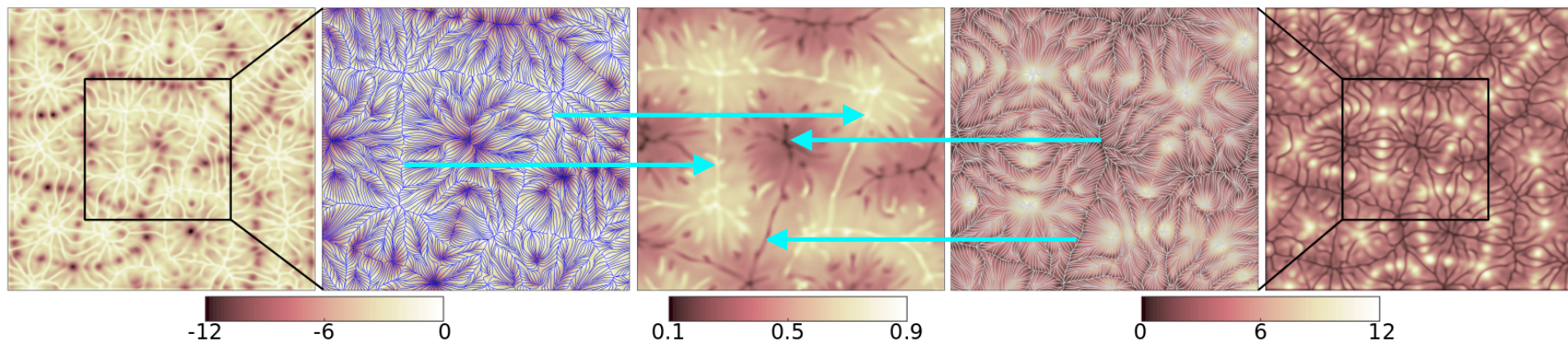
Knot or
cross-roll instability

Connection to boundary layer dynamics

Pr=0.005



Pr=7



$$\left. \frac{\partial T}{\partial z} \right|_{z=0}$$

$$\vec{s} = \left(\left. \frac{\partial u_x}{\partial z} \right|_{z=0}, \left. \frac{\partial u_y}{\partial z} \right|_{z=0} \right)$$

$$T \Big|_{z=\frac{1}{2}}$$

$$\vec{s} = \left(\left. \frac{\partial u_x}{\partial z} \right|_{z=1}, \left. \frac{\partial u_y}{\partial z} \right|_{z=1} \right)$$

$$-\left. \frac{\partial T}{\partial z} \right|_{z=1}$$

Bottom Plate

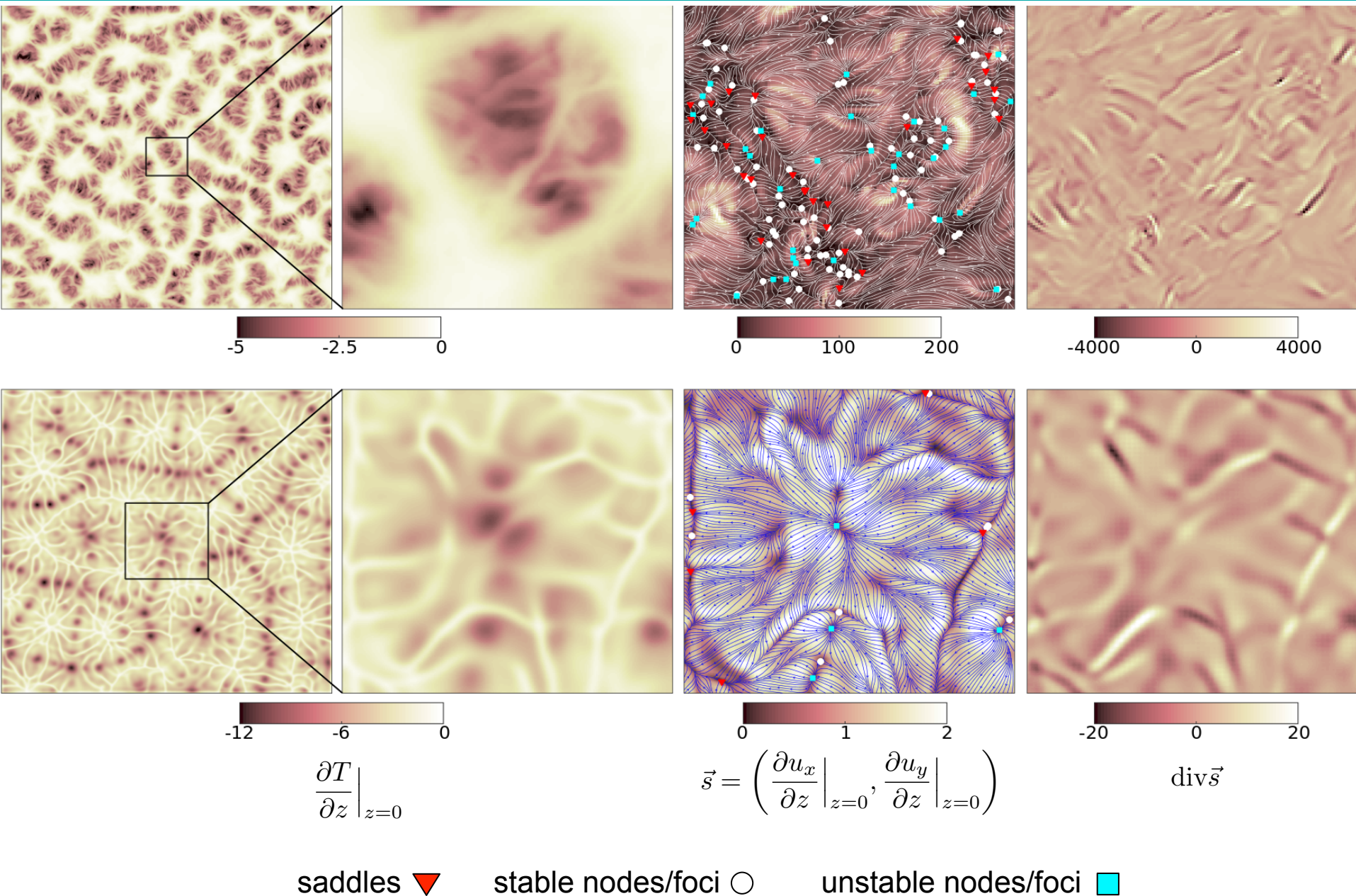
Bottom Plate
(zoom)

Midplane
(zoom)

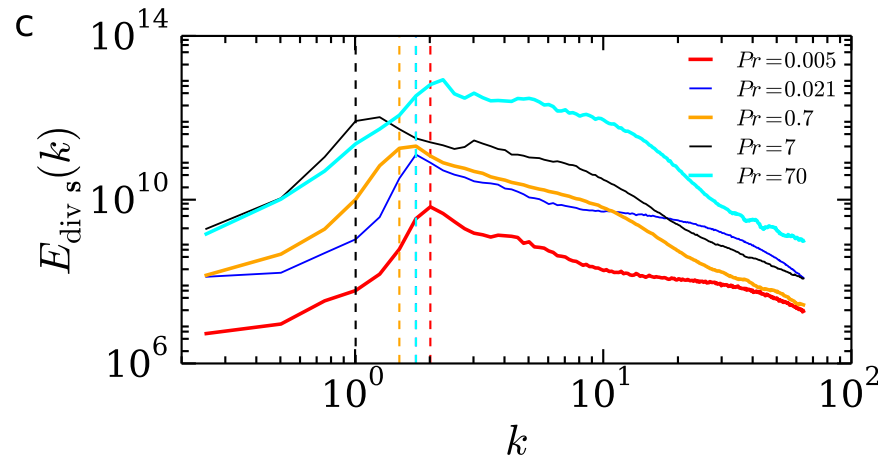
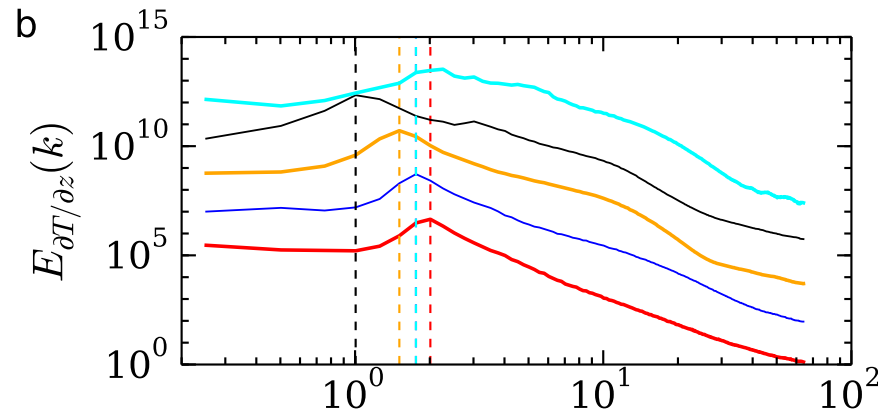
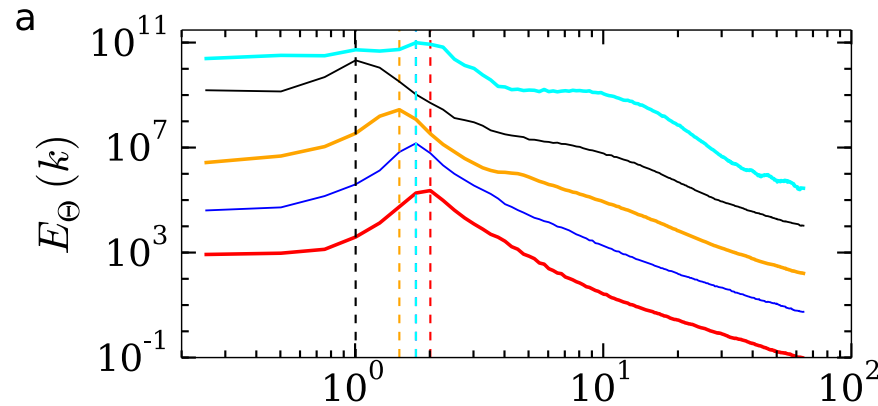
Top Plate
(zoom)

Top Plate

Skin friction field as a blueprint for plume formation



Comparison of azimuthal spectra



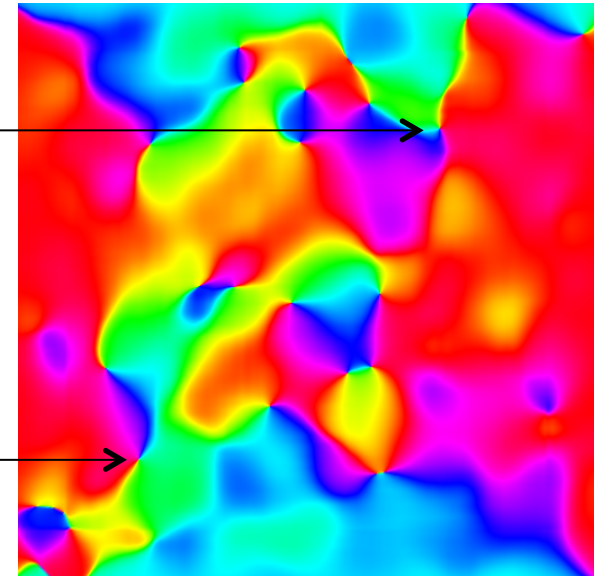
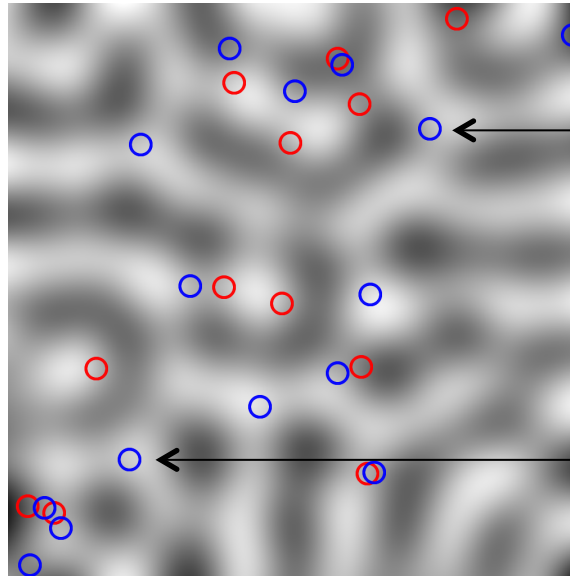
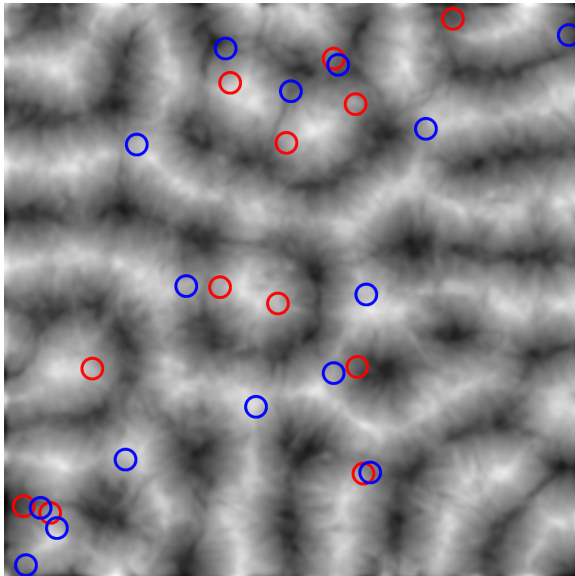
Peak wavenumbers
of all three time-
averaged spectra do
agree

Defects of superstructure patterns

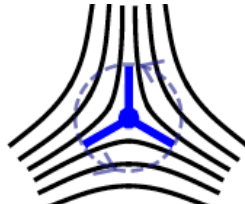
Time-averaged T

Additional Gauss filter

Phase



+1/2
Wedge



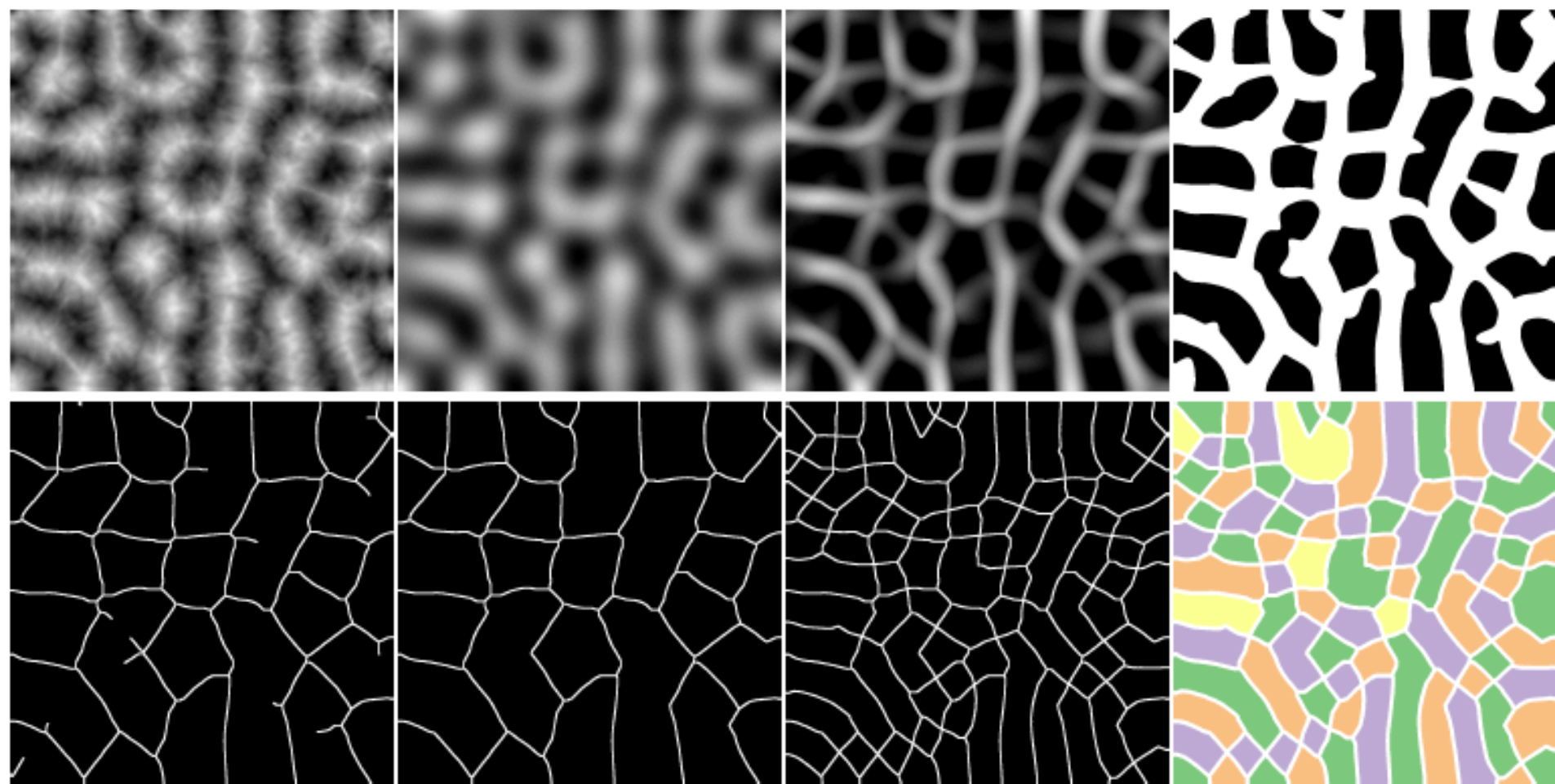
-1/2
Trisector

$$\phi = \frac{1}{2} \arctan \frac{2I_{xy}}{I_{xx} - I_{yy}}$$

Topological points defects detectable in superstructure patterns

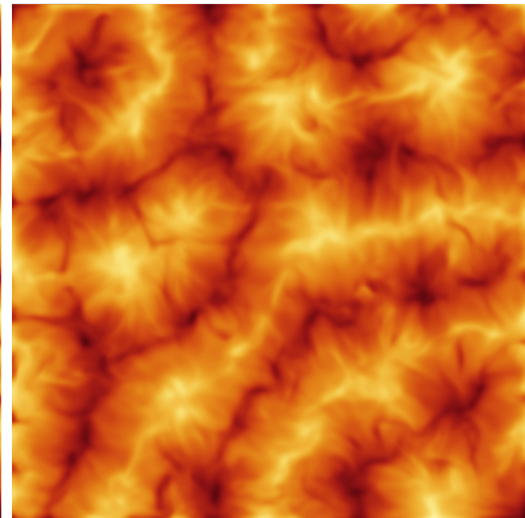
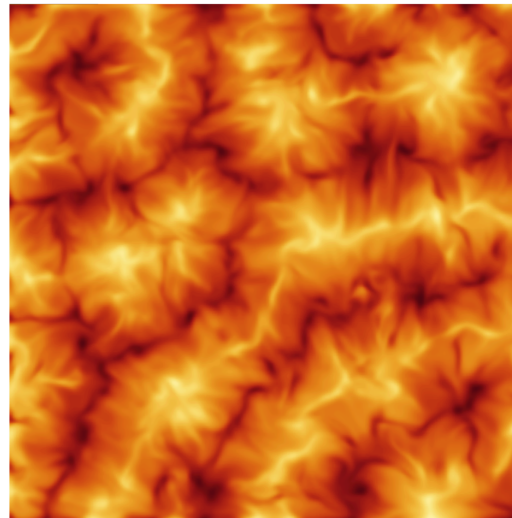
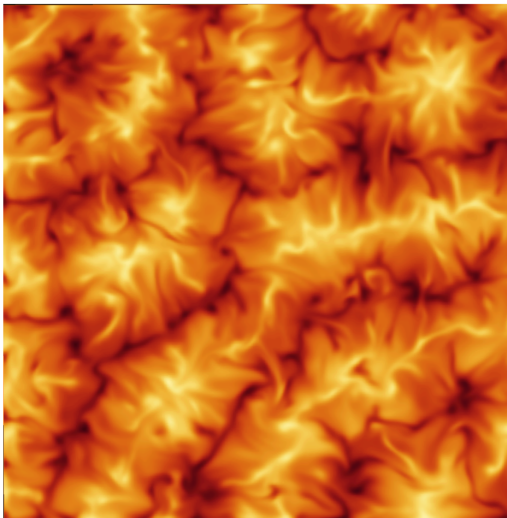
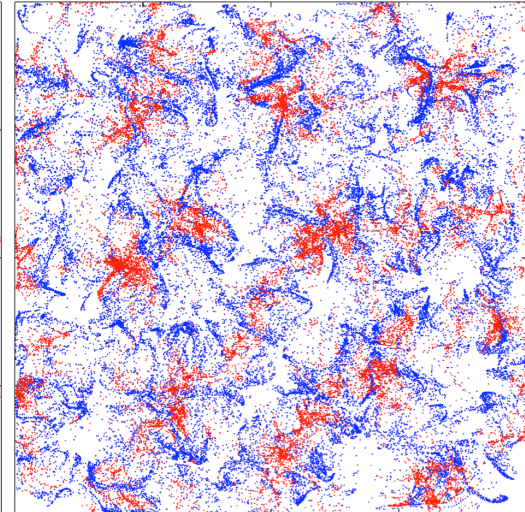
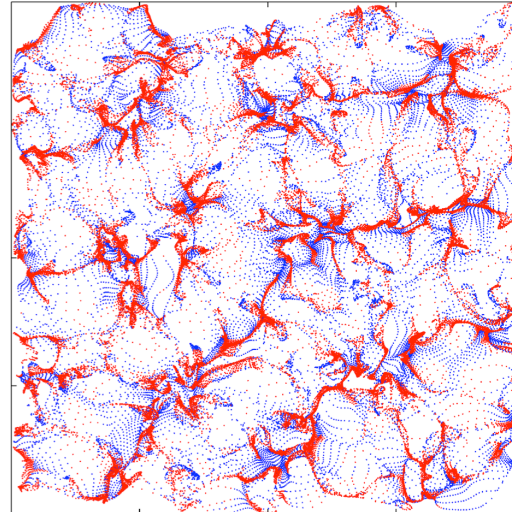
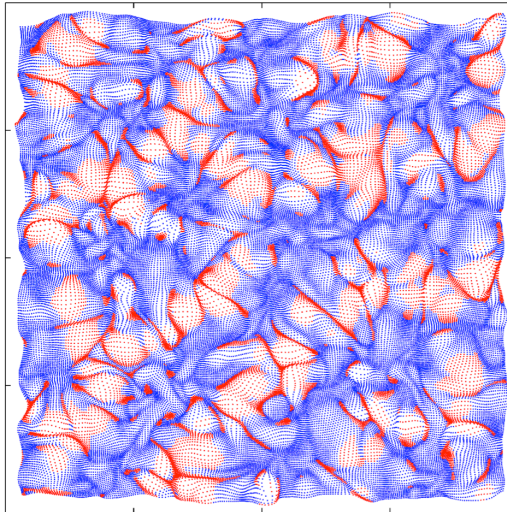
Circulation rolls

Ender, Weiss, Pandey, JS & Westermann, submitted to Eurographics 2018



Separation of different circulation rolls segments on the basis of T
to study mixing and transport on long time scales

Lagrangian analysis of superstructures

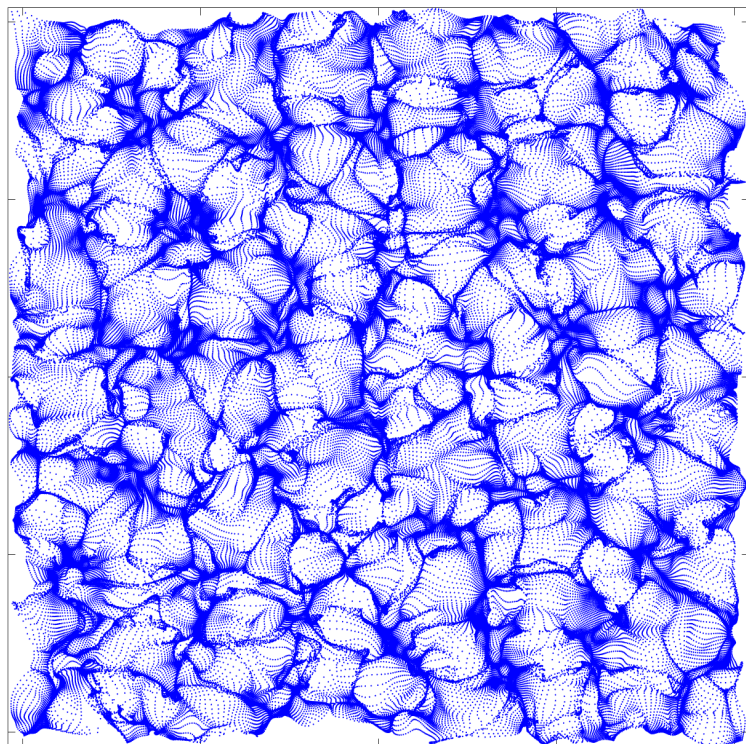
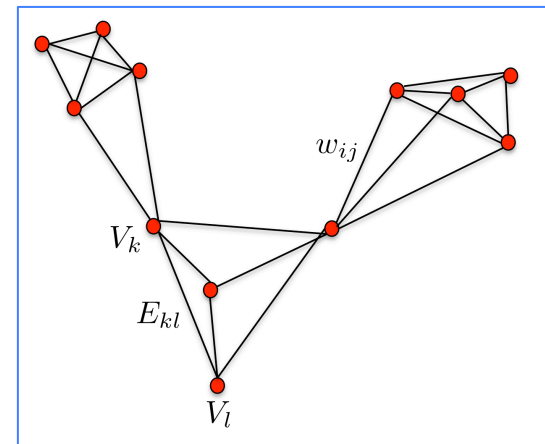


From single tracers & pairs to tracer clusters in extended domains

Approach by graph theory

Shi & Malik, *IEEE Trans. Patt. Anal. Machine Intell.* 2000

- Lagrangian tracer ensemble is composed to a space-time graph
- Superstructure detection = Image segmentation
= **Balanced cut problem for a graph**



- $G = \{V, E, W\}$

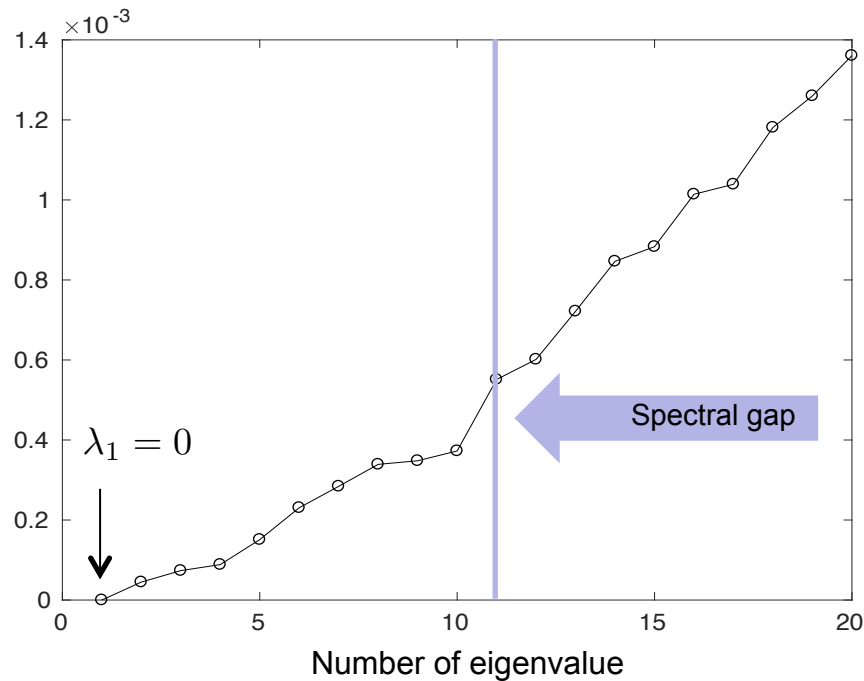
- $r_{ij}(t_1) = \frac{1}{t_1} \int_0^{t_1} |\vec{x}_i(t) - \vec{x}_j(t)| dt$

- $w_{ij} = \begin{cases} K \gg 1 & : i = j \\ \frac{1}{r_{ij}} & : i \neq j \end{cases}$

- $L\vec{y} = \lambda D\vec{y} \quad \longrightarrow \quad D^{-1/2} L D^{-1/2} \vec{z} = \lambda \vec{z}$

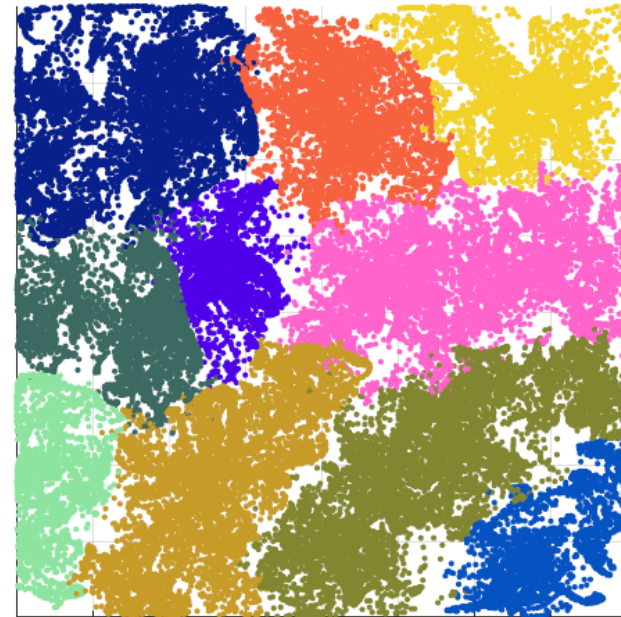
Cutting the graph into cluster

Eigenvalue spectrum of graph Laplacian

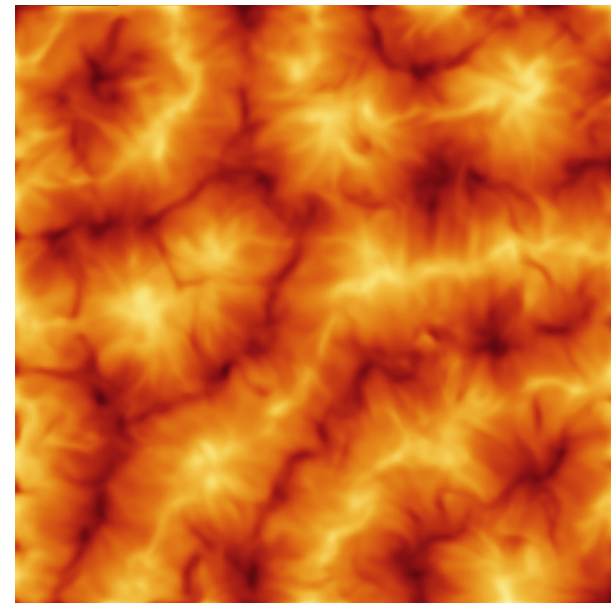


$$r_{ij} \leq \epsilon \Rightarrow w_{ij} = \frac{1}{r_{ij}} \quad r_{ij} > \epsilon \Rightarrow w_{ij} = 0$$

Clusters from spectral analysis of graph Laplacian agree with superstructures



11
spectral
clusters



Time-
averaged
temperature
in midplane

Finding the right cutoff scale ϵ

Scheel & JS, J. Fluid Mech. 2014

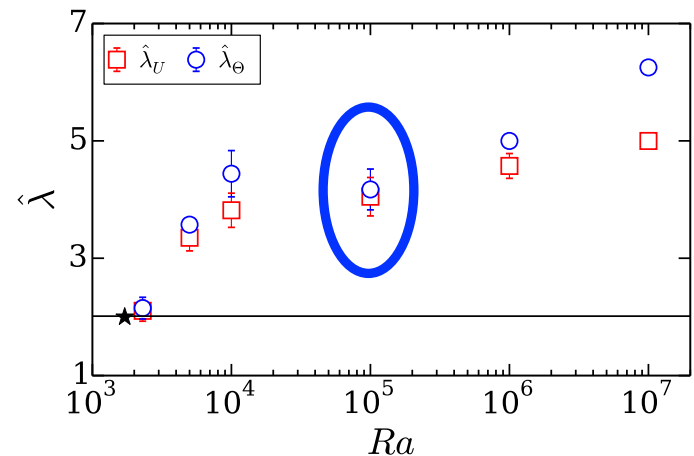
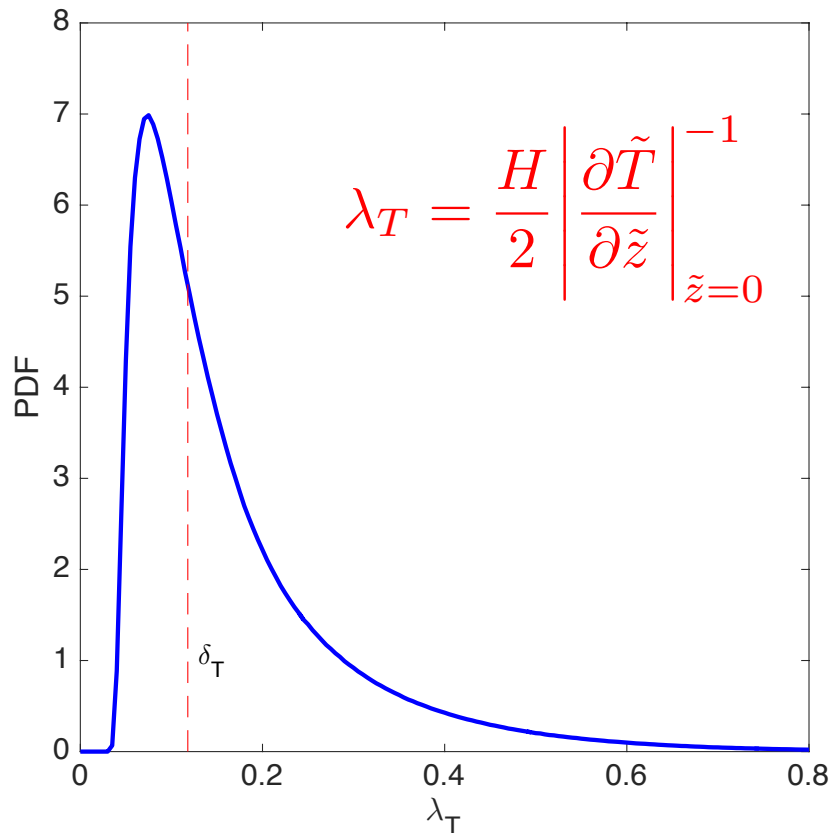
$$r_{ij} \leq \epsilon \Rightarrow w_{ij} = \frac{1}{r_{ij}}$$

$$r_{ij} > \epsilon \Rightarrow w_{ij} = 0$$

Local thermal boundary layer thickness



Characteristic superstructure scale



$$\lambda_{\Theta} \approx 2$$

Summary

- Large-scale patterns are detected in turbulent RBC for all Ra and Pr once the fast small-scale fluctuations are removed
- Characteristic mean pattern scale depends on Pr and Ra in present parameter range
- Patterns evolve gradually on time scales that are significantly larger than the free-fall time, but significantly smaller than diffusion time
- Pattern scales are correlated to characteristic scale of skin friction field at plates and thus to formation of plume ridges
- Changes of superstructure patterns are associated with the generation and annihilation of topological point defects
- Spectral analysis of graph Laplacian reveals superstructure pattern scales in Lagrangian frame



Leibniz-Rechenzentrum
der Bayerischen Akademie der Wissenschaften



<http://lanl.arxiv.org/abs/1801.04478>