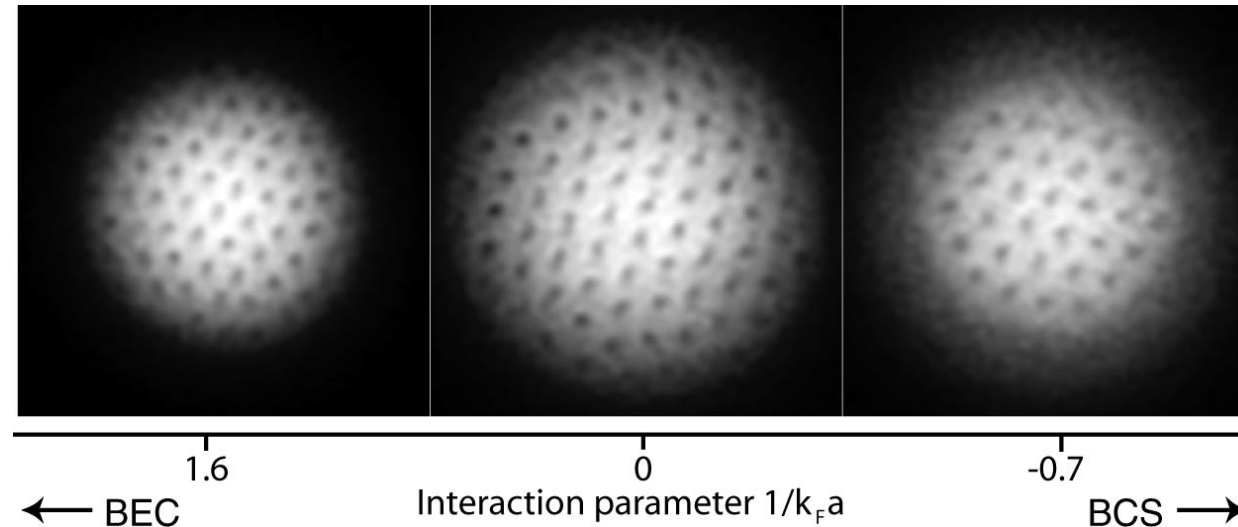


Superfluid Fermi Gases

Martin Zwierlein

Massachusetts Institute of Technology
Center for Ultracold Atoms



Some reviews

Varennna Notes:

on Fermi Gases:

Ketterle, Zwerlein, *Making, Probing and Understanding Ultracold Fermi Gases*

<http://arxiv.org/abs/0801.2500>

Zwerlein, *Thermodynamics of Strongly Interacting Fermi Gases*,

Varennna School of Physics “Enrico Fermi” 2014, vol 191

Stefano Giorgini, Lev P. Pitaevskii, Sandro Stringari

The theory of Fermi gases

<http://arxiv.org/abs/0706.3360>

Immanuel Bloch, Jean Dalibard, Wilhelm Zwerger

Many-Body Physics with Ultracold Gases:

<http://arxiv.org/abs/0704.3011>

Lecture Notes “The BEC-BCS crossover and the Unitary Fermi Gas”

Edited by W. Zwerger, Springer, 2012

How cold is ultracold?

Atoms move at:

$\sim 1 \text{ mm/s}$

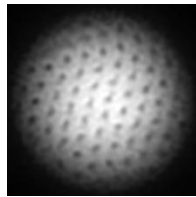
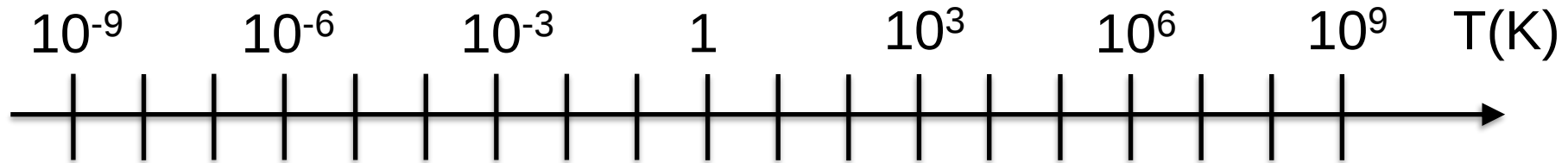


$\sim 100 \text{ m/s}$



$\sim 10^6 \text{ m/s}$

NY – Paris in 10
seconds



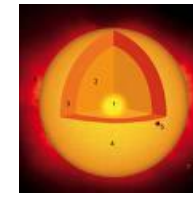
Ultracold atom
experiments



Outer
space



Your
living
room



Center
of the
sun



Supernova
explosion

Particles behave as waves



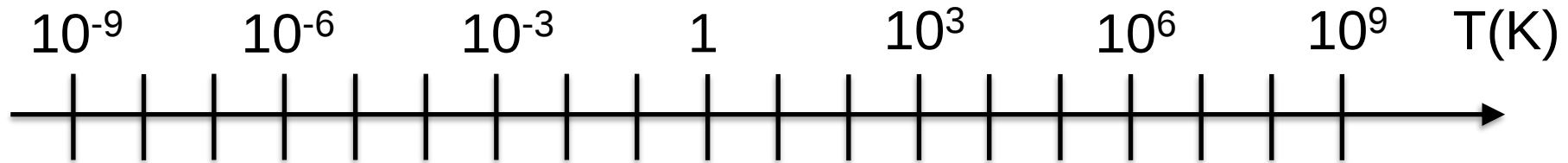
*Louis
de Broglie*

$$\lambda = \frac{h}{mv}$$

Planck's constant

Mass

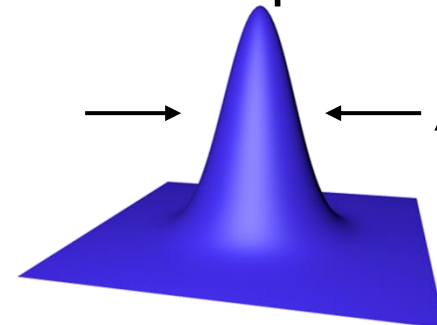
Velocity



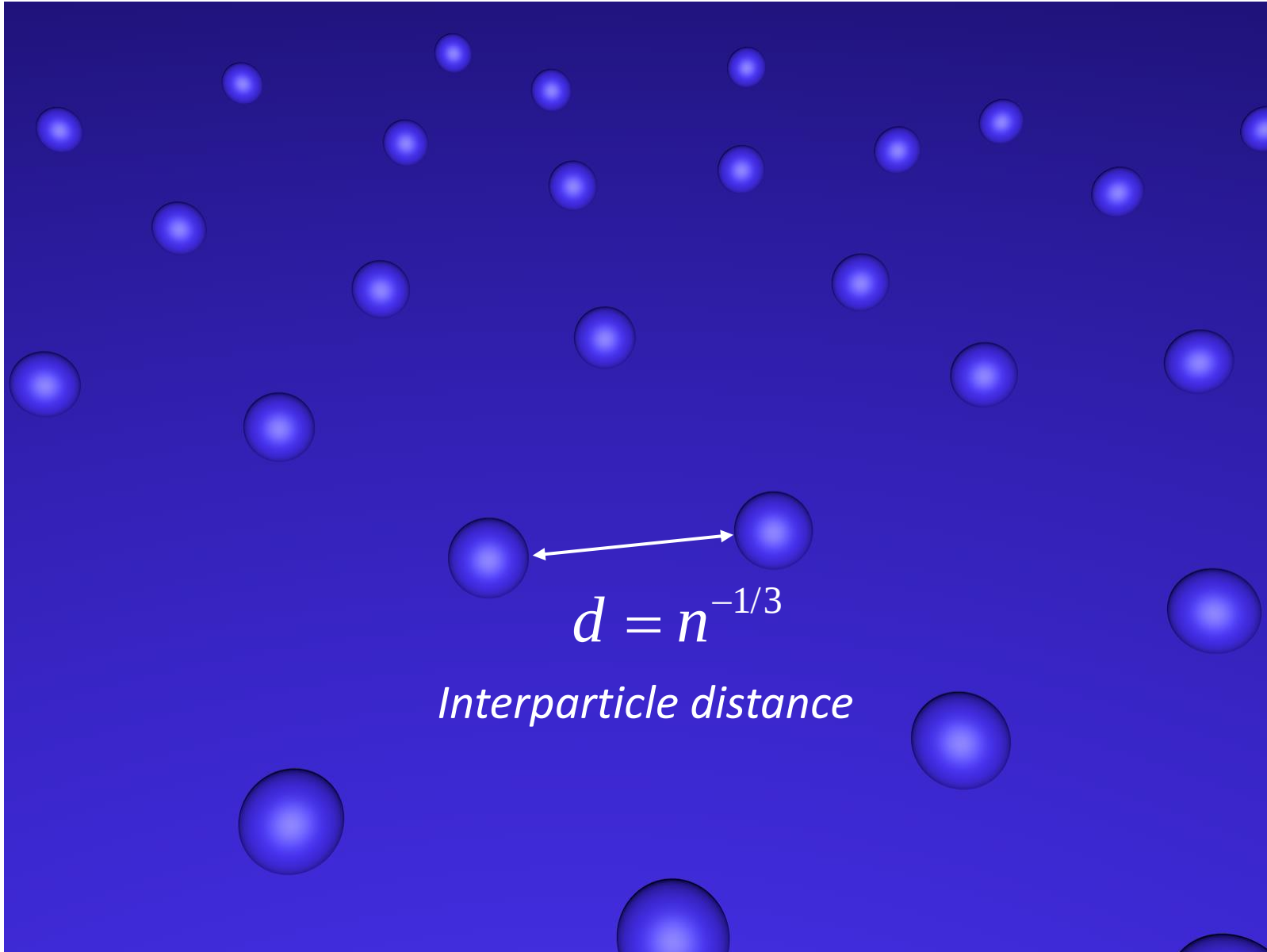
*Werner
Heisenberg*

Temperature = Uncertainty of velocity²

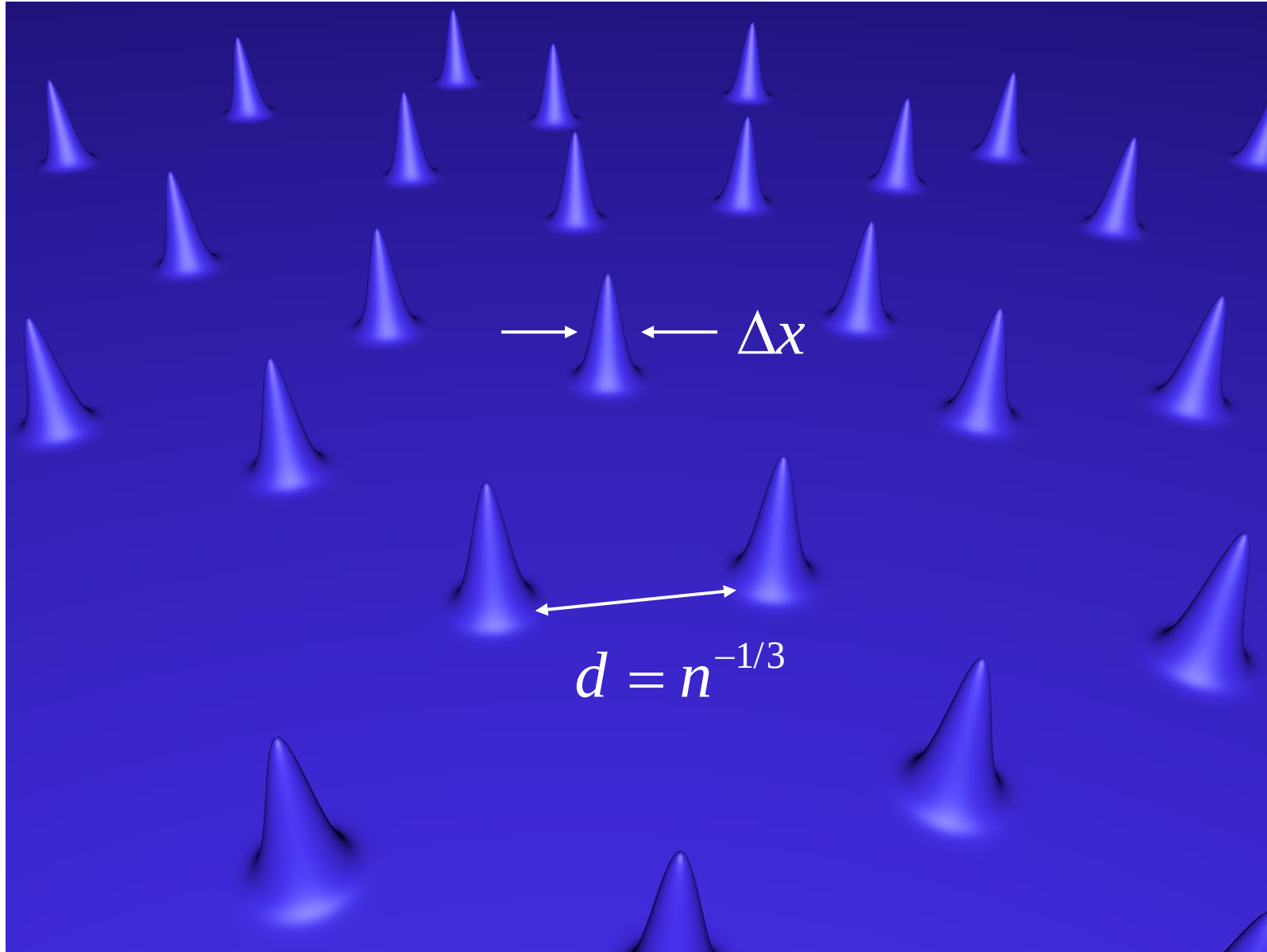
Size of a wave packet = Uncertainty of position

A 3D plot of a wave packet, represented as a blue bell-shaped curve on a flat base.
$$\Delta x \approx \frac{\hbar}{\Delta p} \approx \frac{\hbar}{\sqrt{mT}}$$

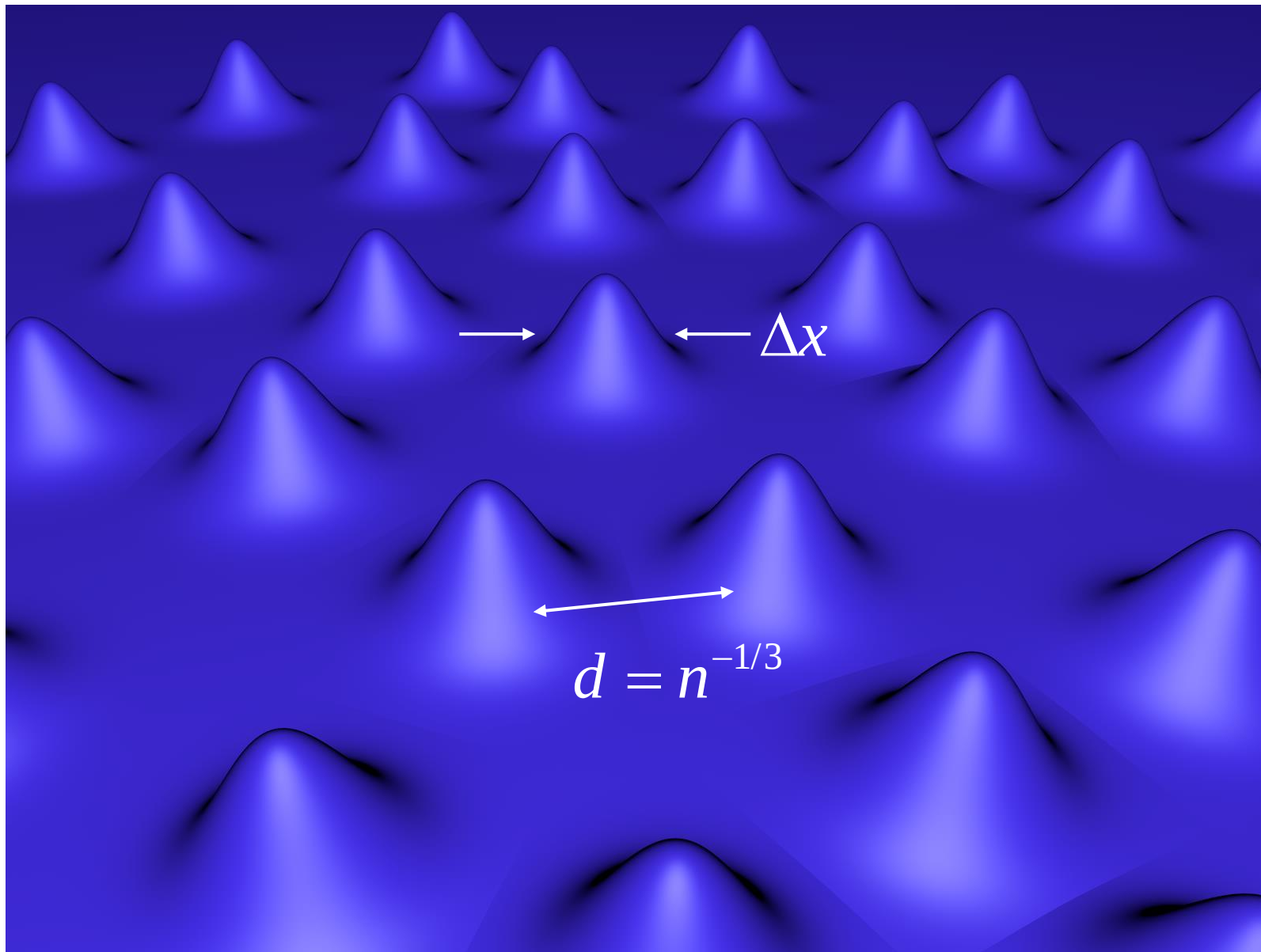
Particles...



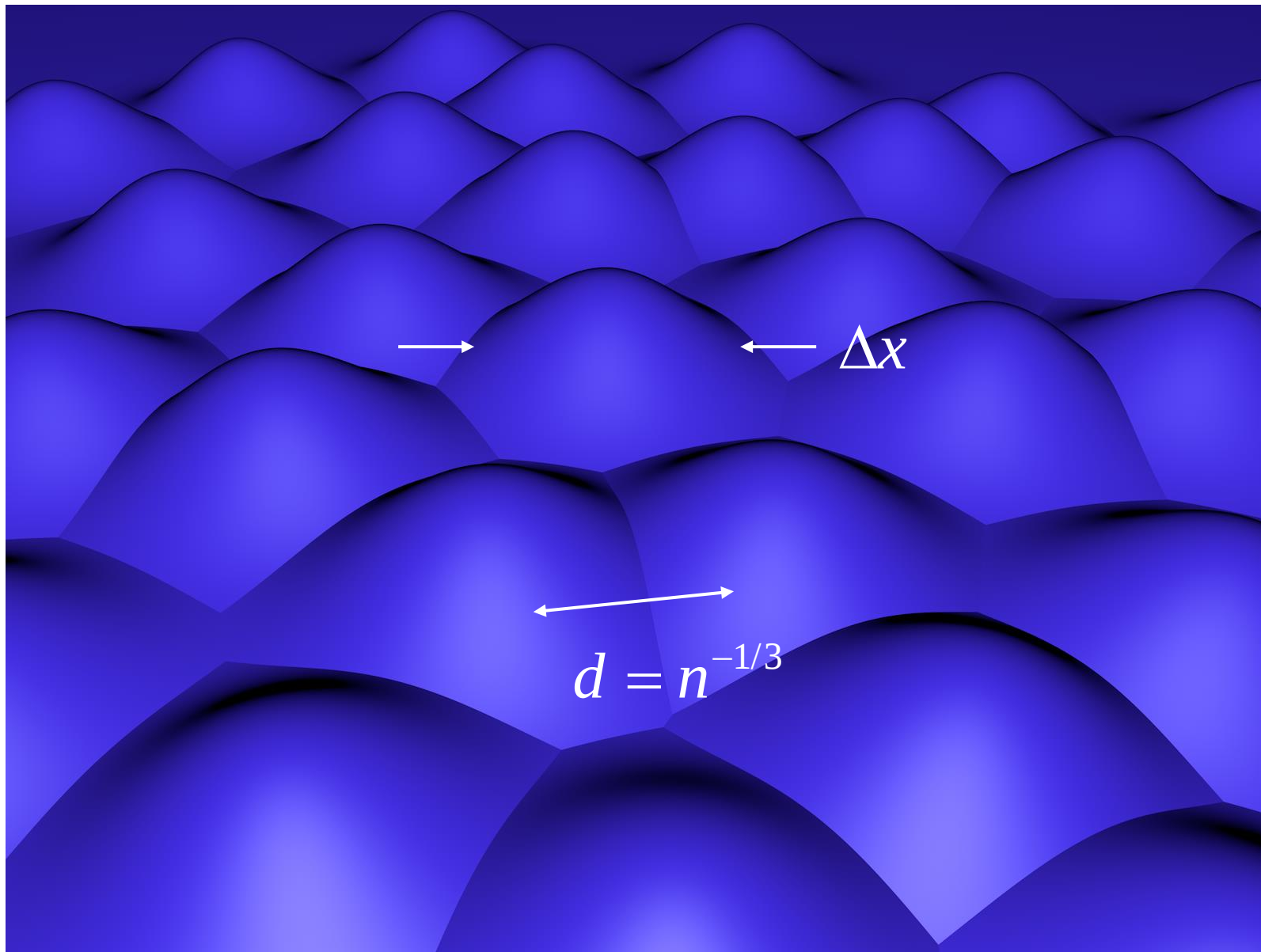
...behave as waves



When does wave mechanics matter?



When does wave mechanics matter?



Bosons versus Fermions

Fermions (unsociable):

Half-Integer Spin

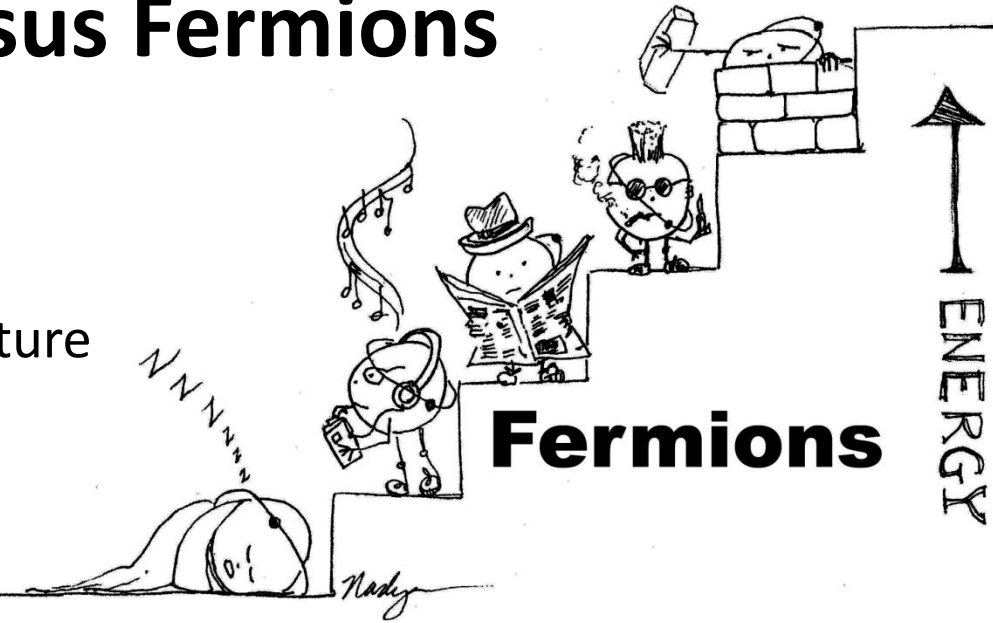
Pauli blocking $\rightarrow$ Form Fermi sea

No phase transition at low Temperature



Fermi Feb. 1926

Dirac Oct. 1926



Bosons (sociable):

Integer Spin

Can share quantum states

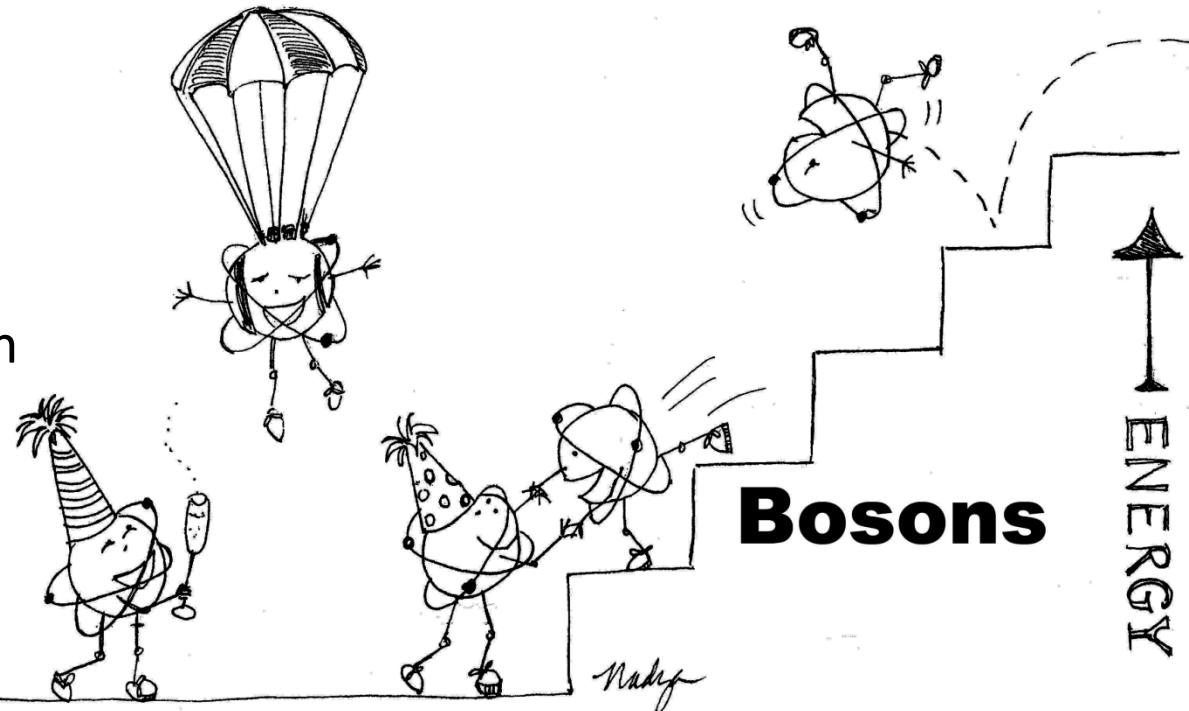
At low temperatures:

Bose-Einstein condensation



Bose 1924

Einstein 1924/25



Bosons



N bosons sharing one and the same macroscopic matter wave

(Artist's conception)

Fermions

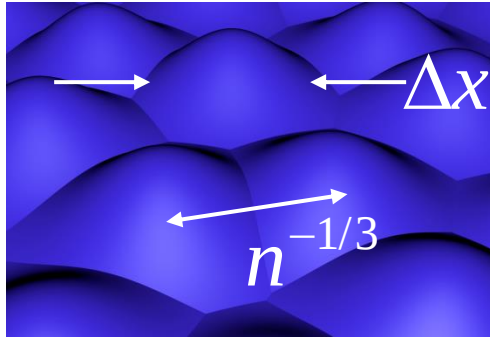


N fermions avoiding each other

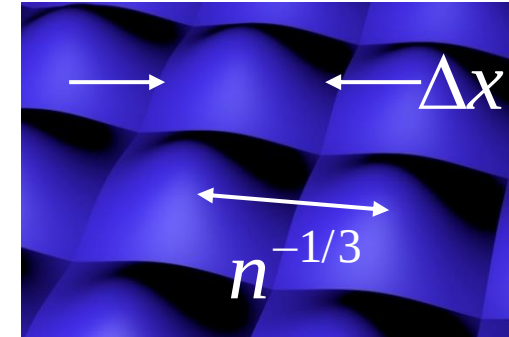
(Artist's conception)

Condition for quantum degeneracy

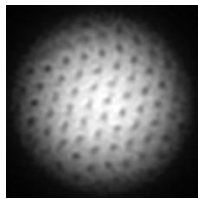
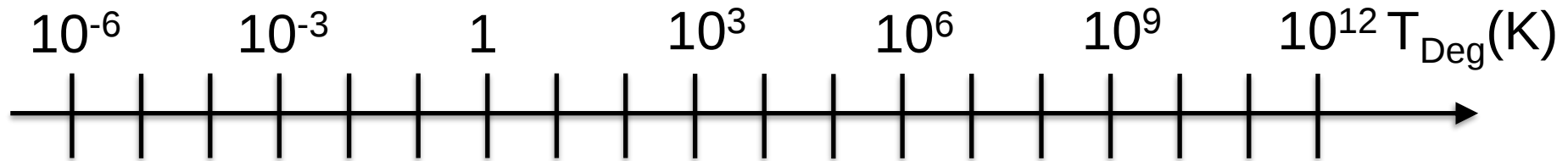
Position uncertainty $\sim$ Interparticle spacing



$$\Delta x = \frac{\hbar}{\Delta p} \approx n^{-1/3}$$



$$k_B T_{\text{Degeneracy}} \approx \frac{\hbar^2}{m} n^{2/3} \approx E_F \quad \text{Fermi energy}$$



Ultracold
atomic gases



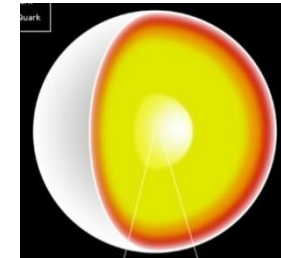
Liquid Helium



Metals

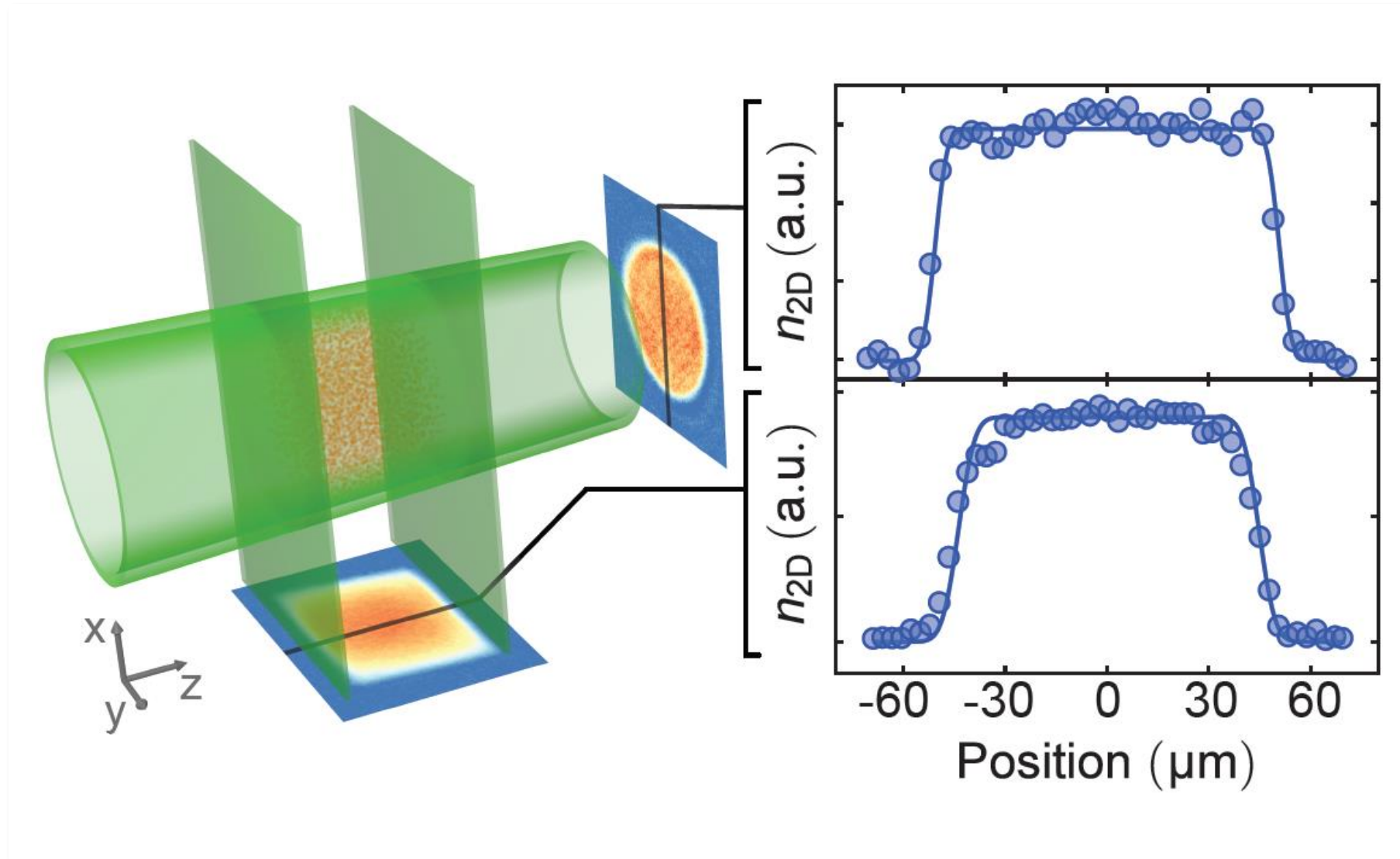


White dwarf

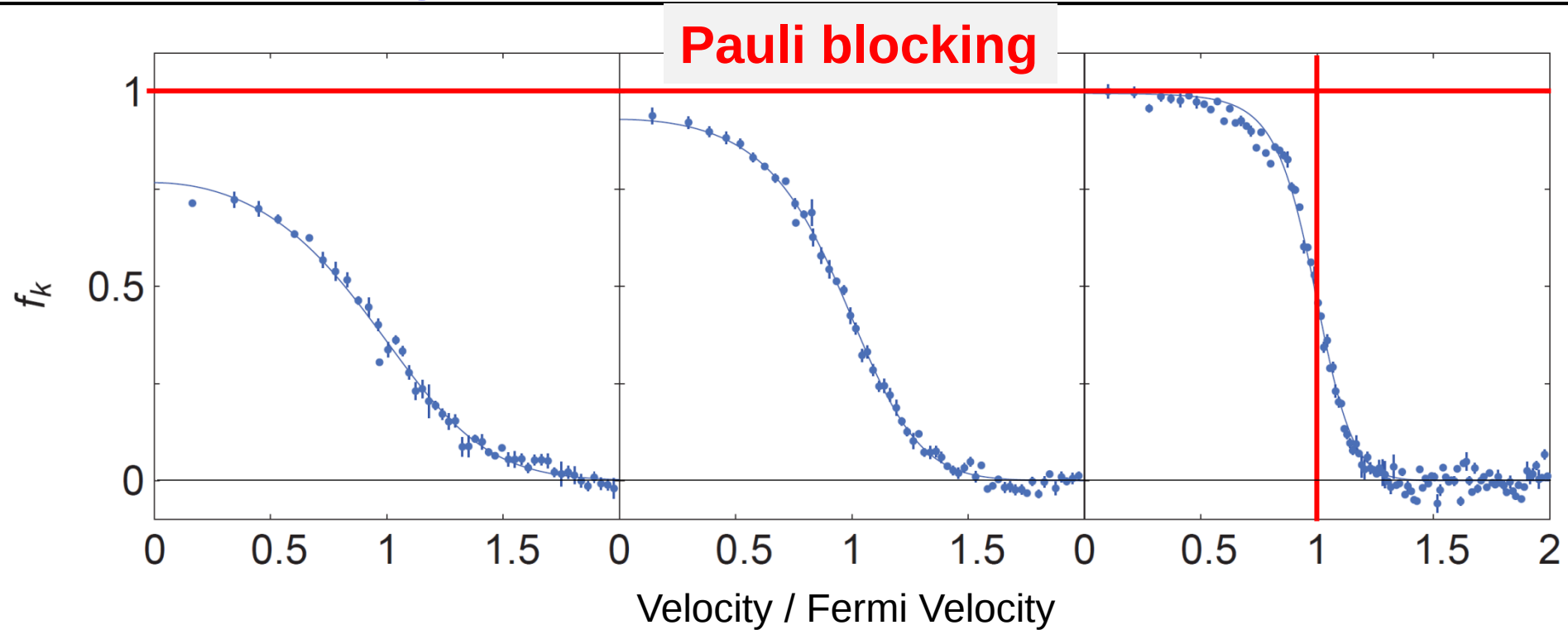


Neutron star

Fermions in a Box



Measuring the Fermi-Dirac distribution



$$f_k = \frac{1}{e^{\beta(\varepsilon_k - \mu)} + 1}$$

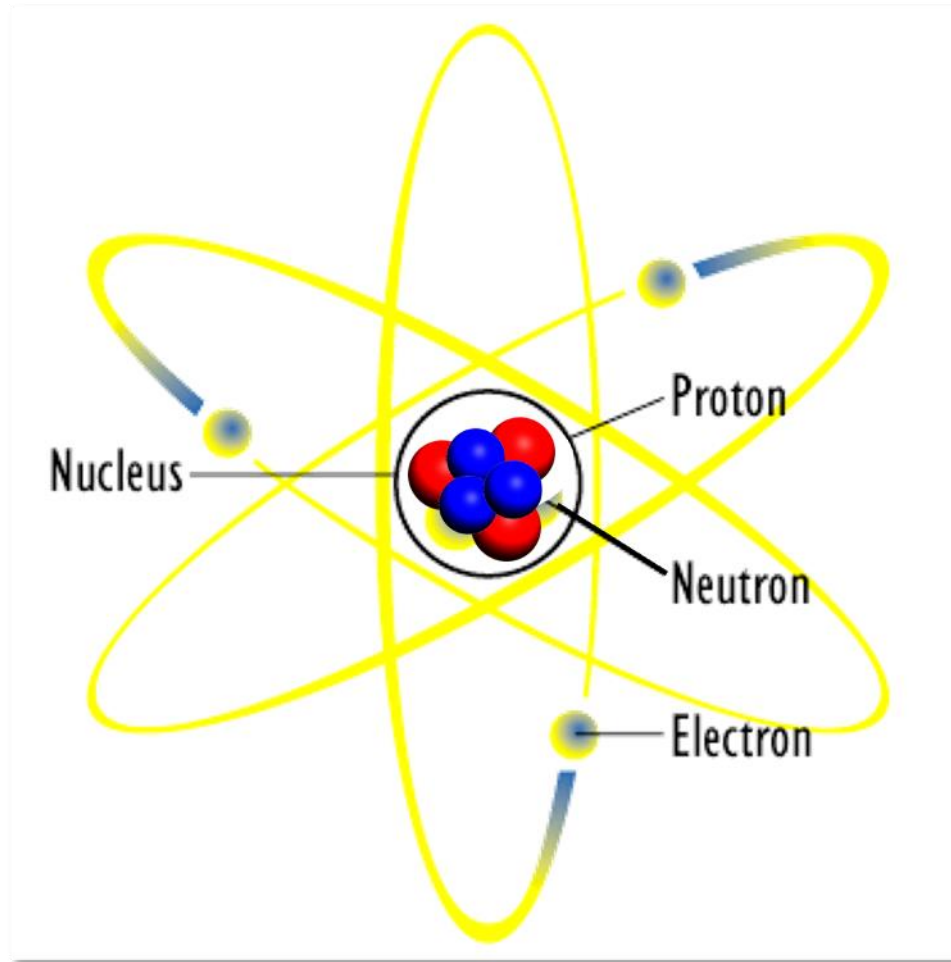
**90th
anniversary**

**Fermi surface
formation**

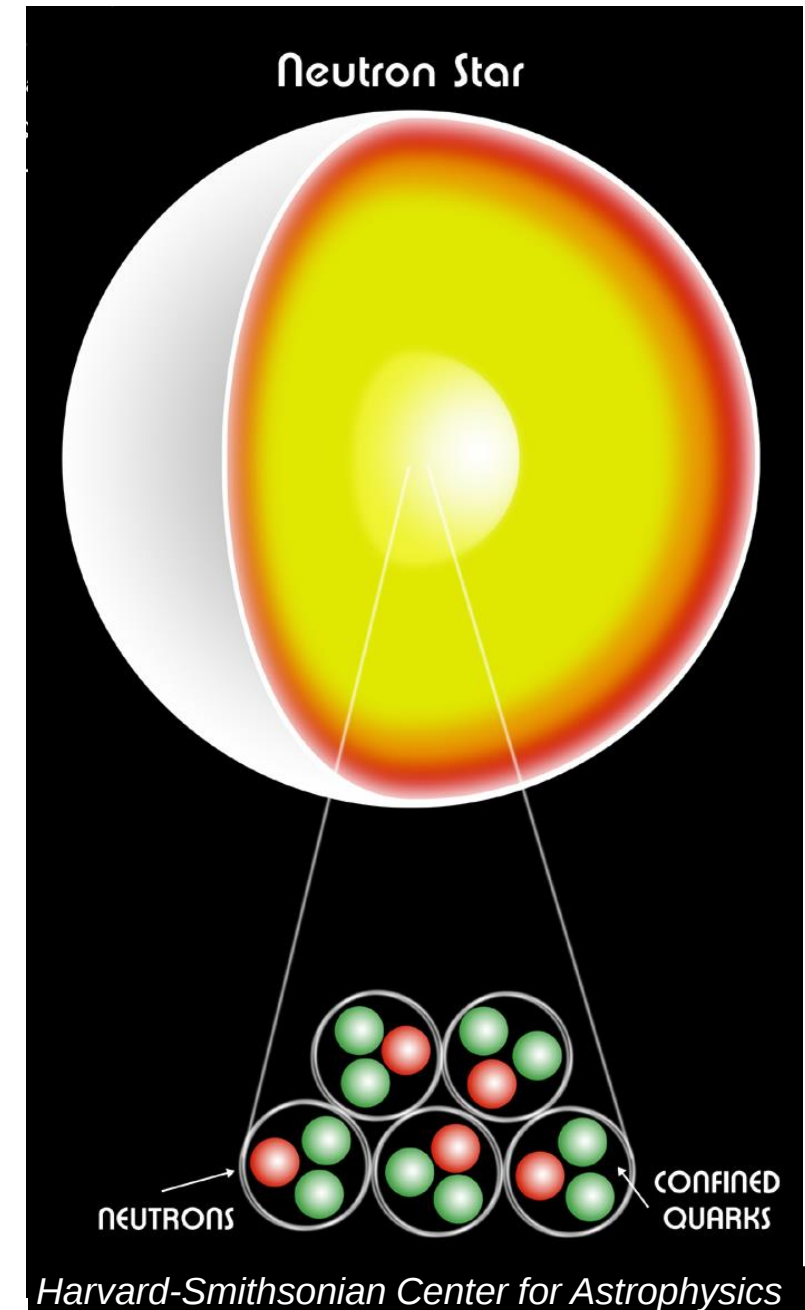
See also: Drake et al., PRA 2012, selectively probe the central portion of an inhomogeneous gas.

Z. Yan, P. Patel, B. Mukherjee, Z. Hadzibabic, T. Yefsah, J. Struck, MWZ, PRL 2017

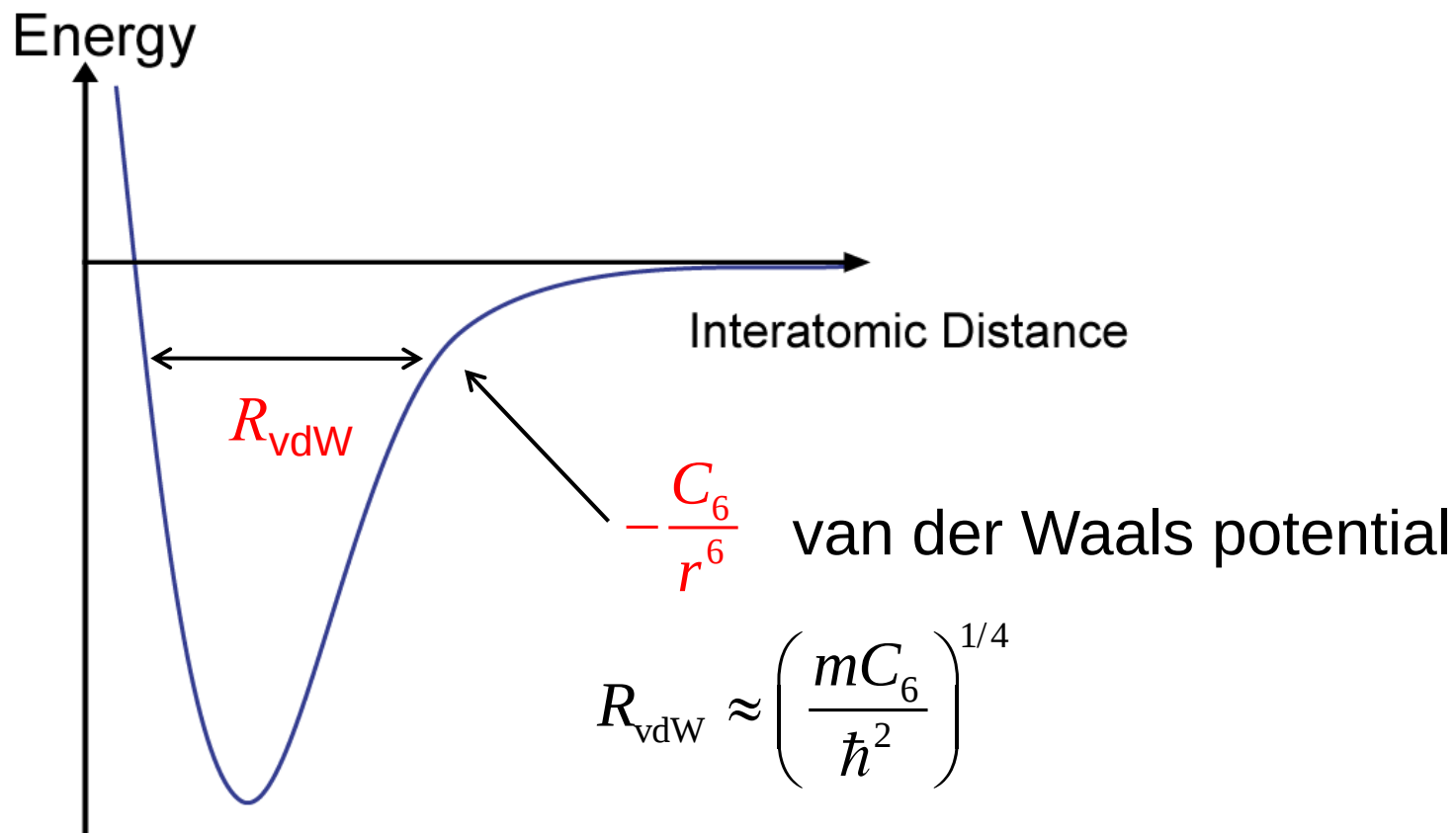
Fermions – The Building Blocks of Matter



Lithium-6

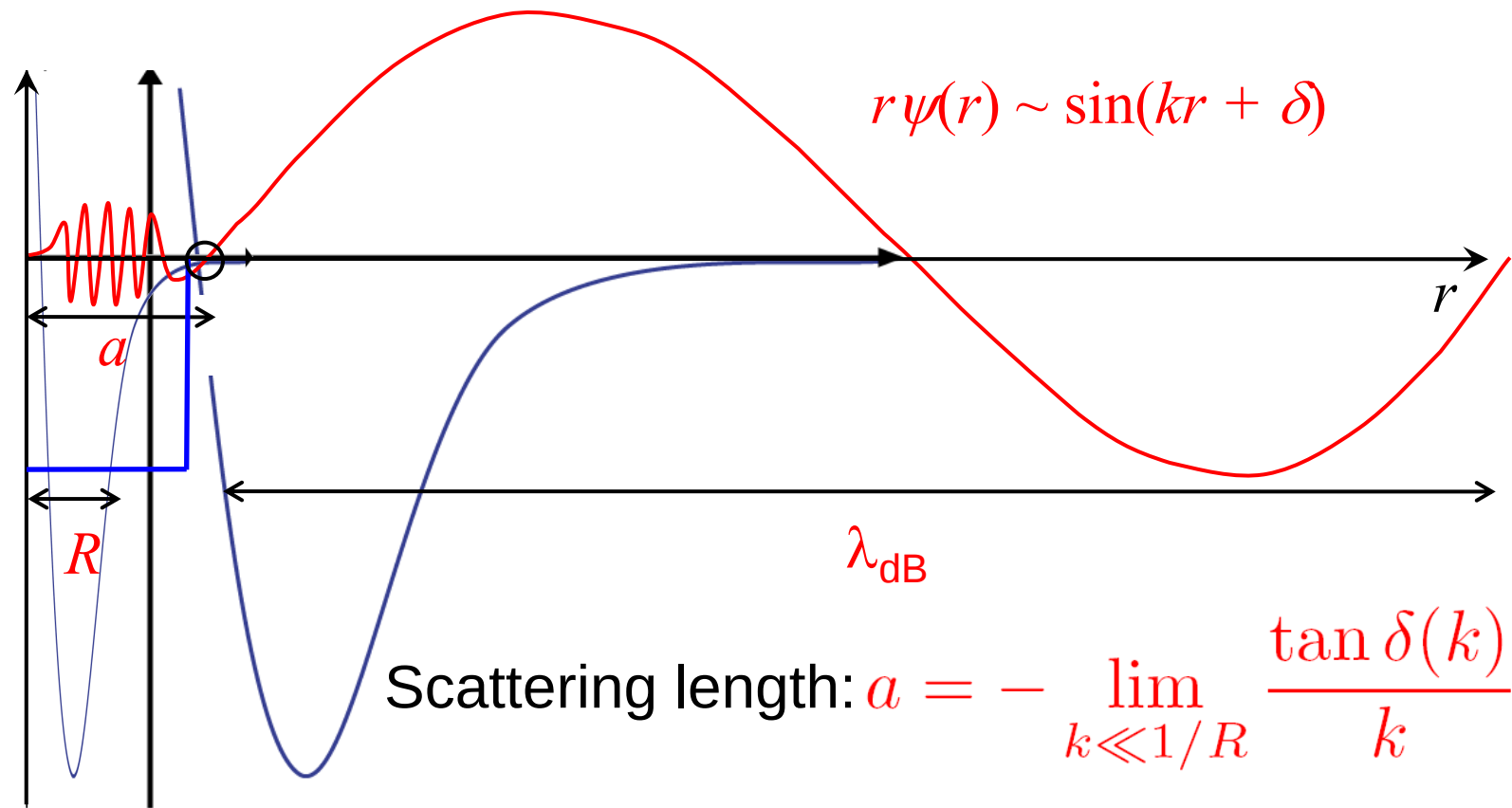


Interatomic interactions



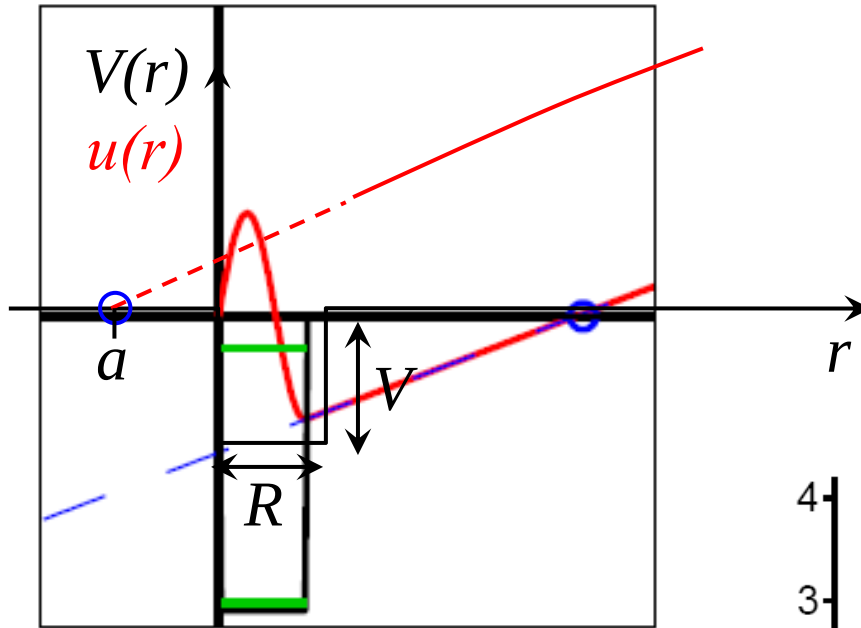
- For Alkali atoms: $R_{\text{vdW}} \sim 50\text{-}200 \text{ a}_0$
 - Ultracold collisions: $\lambda_{\text{dB}} \approx 1\mu\text{m} \gg R$
- atoms do not probe the details of the potential

Interatomic interactions



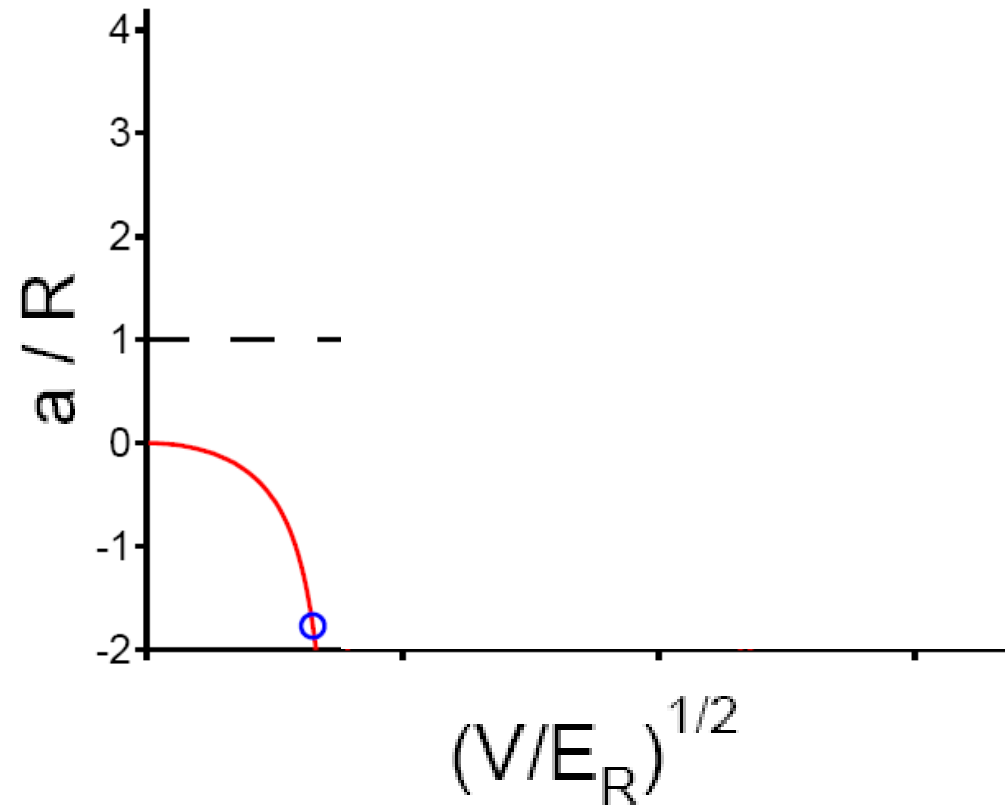
- For Alkali atoms: $R_{vdW} \sim 50-200 a_0$
 - Ultracold collisions: $\lambda_{dB} \approx 1 \mu m \gg R$
- atoms do not probe the details of the potential

Scattering Resonances



$$u(r) = \sin(k(r - a)) \quad r > R$$

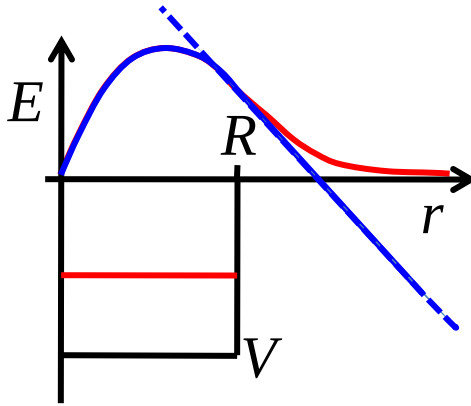
(s-wave) scattering length



Tunable Interactions

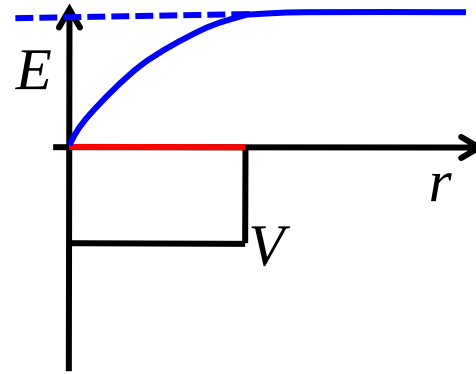
Vary interaction strength between spin up and spin down

Example: tunable square well (with $k_F R \ll 1$):



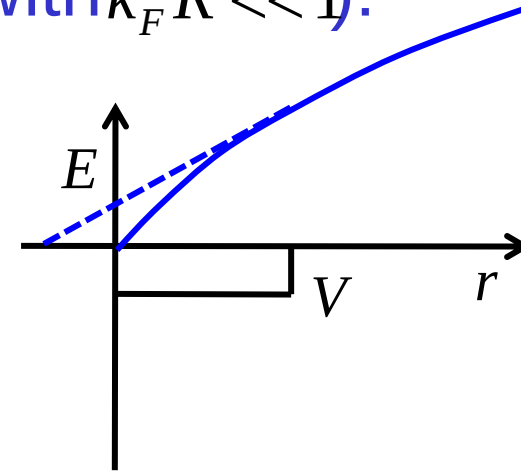
strong attraction
deep bound state

$$a > 0$$



Resonance
bound state appears

$$a \rightarrow \pm\infty$$



weak attraction
no bound state

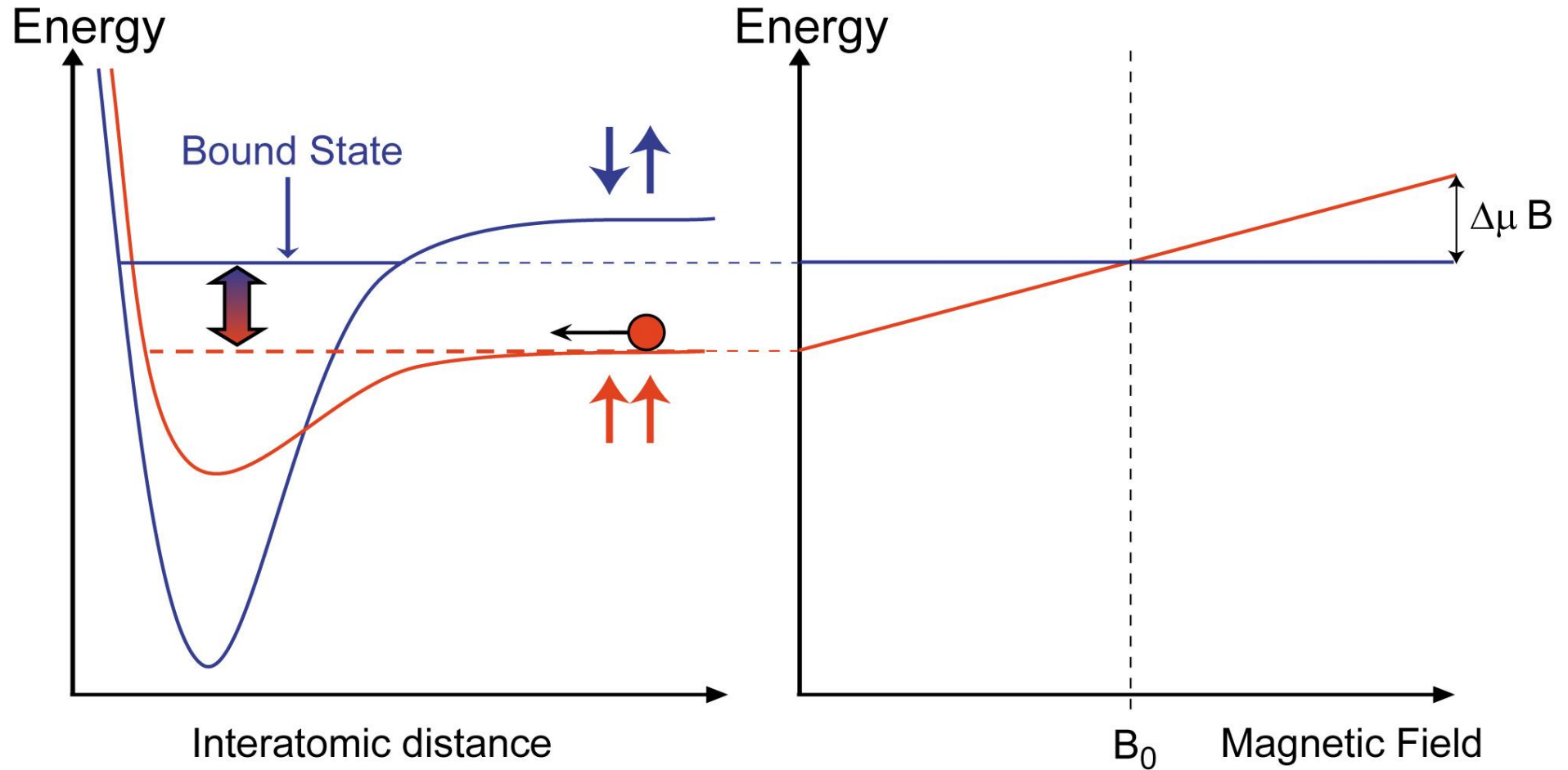
$$a < 0$$

scattering length

Important for Many-Body
Physics:

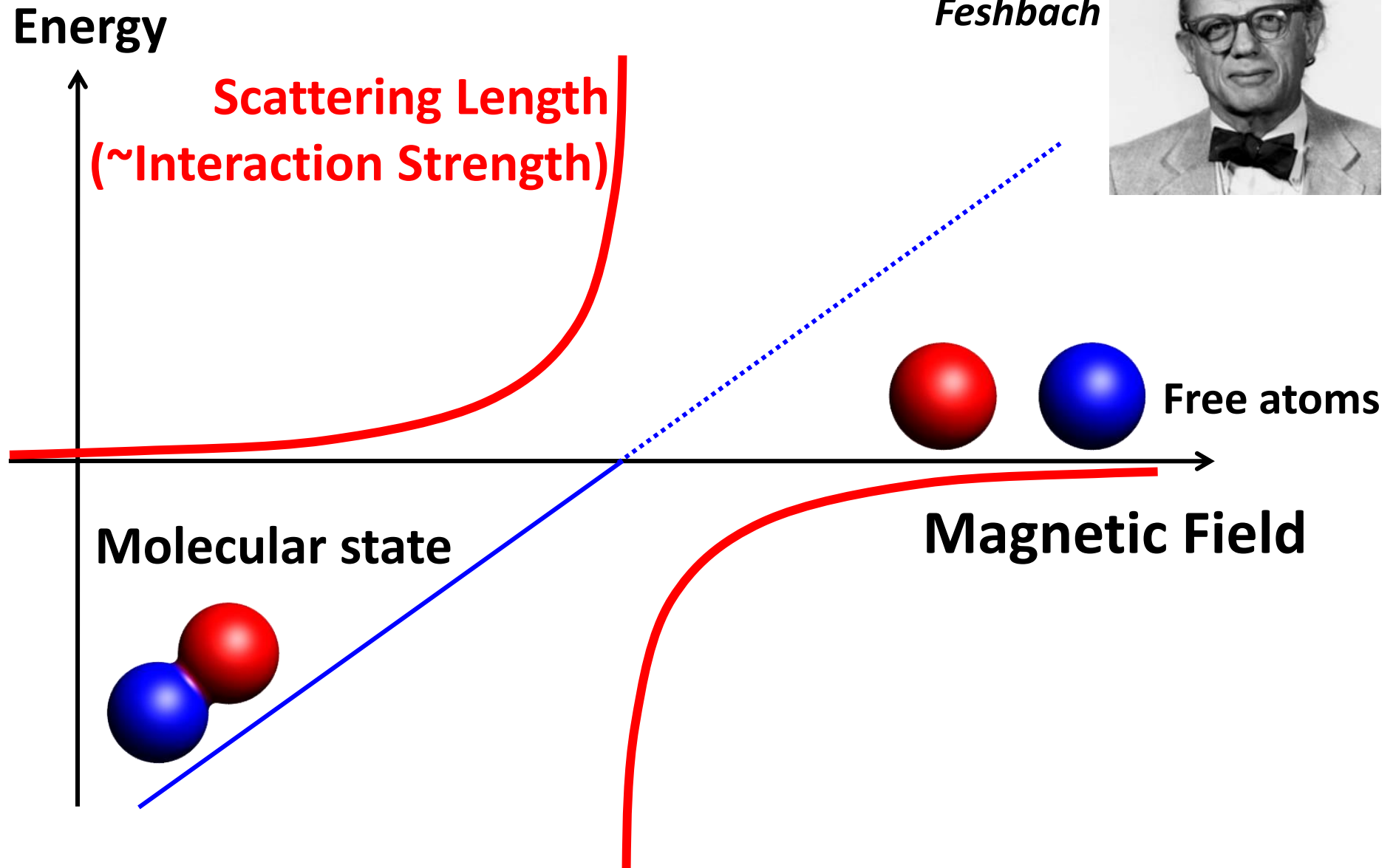
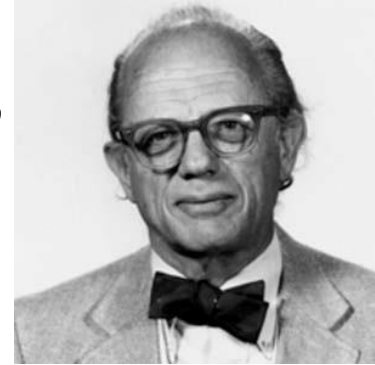
$$\frac{\text{Interparticle distance}}{\text{Scattering length}} = \frac{1}{k_F a}$$

Feshbach resonances: Tuning the interactions

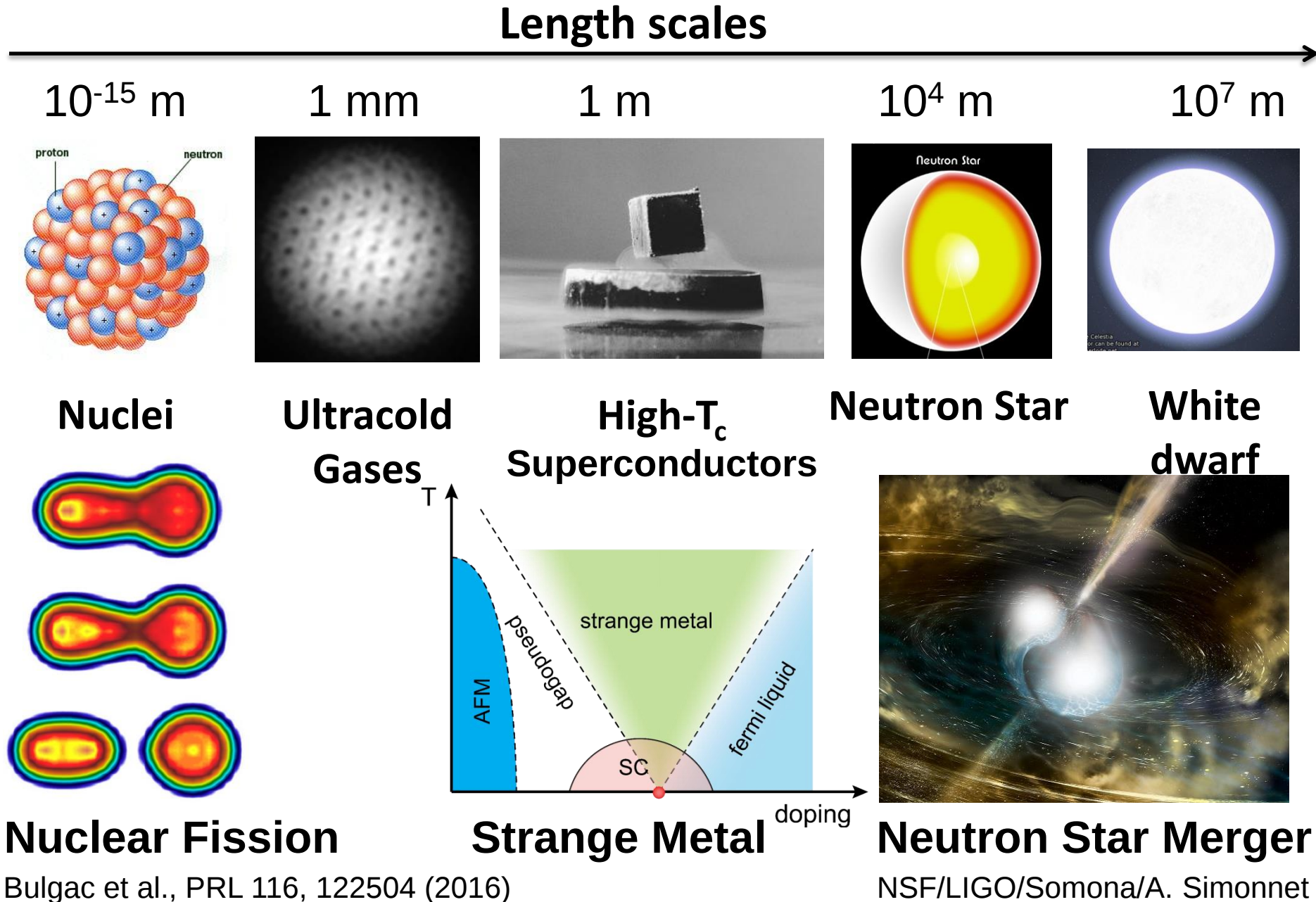


Feshbach resonances: Tuning the interactions

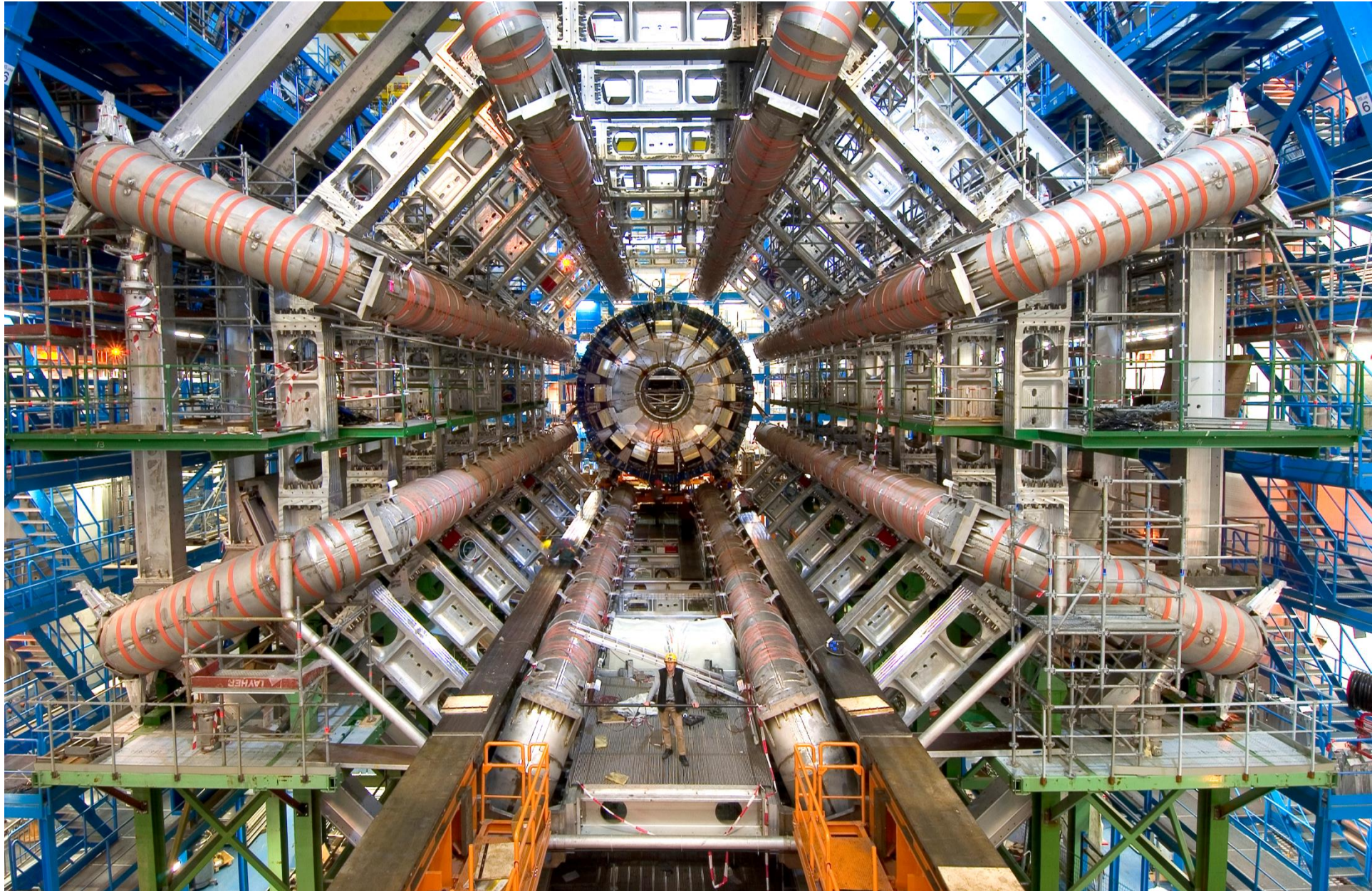
*Herman
Feshbach*

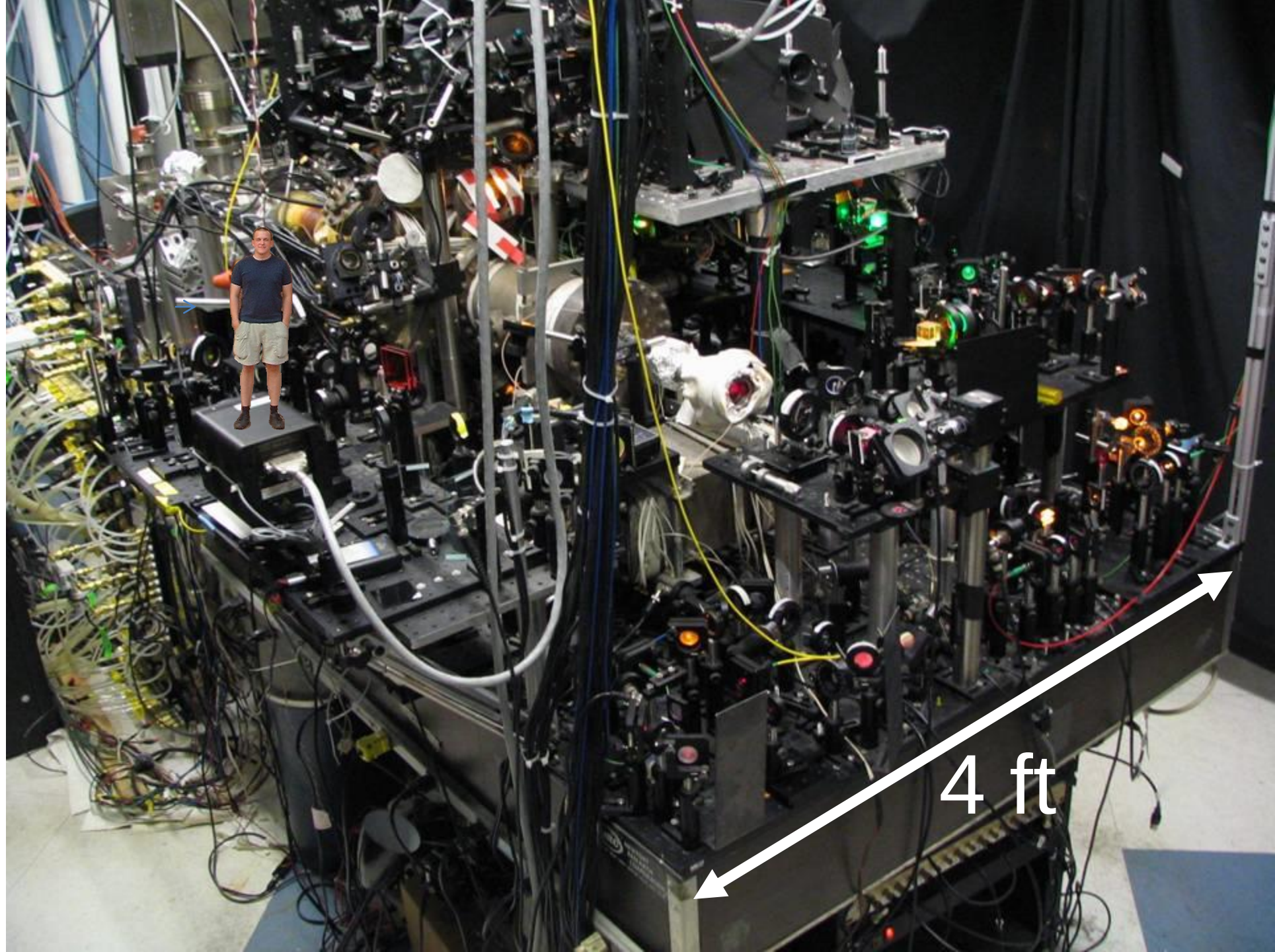


Strongly Interacting Fermi Systems



Large Hadron Collider (LHC)

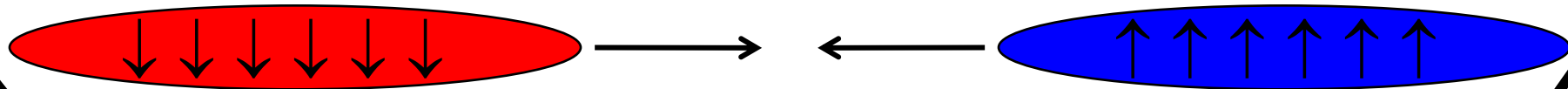




4 ft

Little Fermi Collider (LFC)

A ↓ Fermi gas collides with a ↑ cloud
with resonant interactions



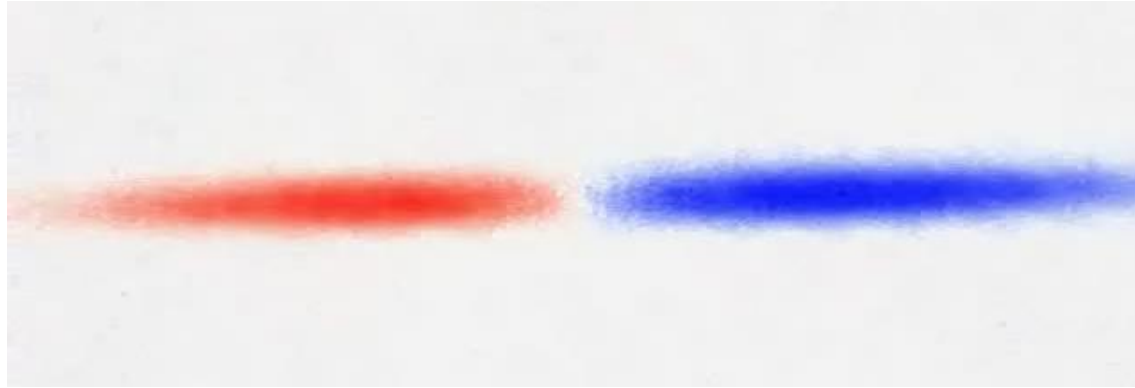
Collision at 100 peV

23 orders of magnitude smaller than LHC

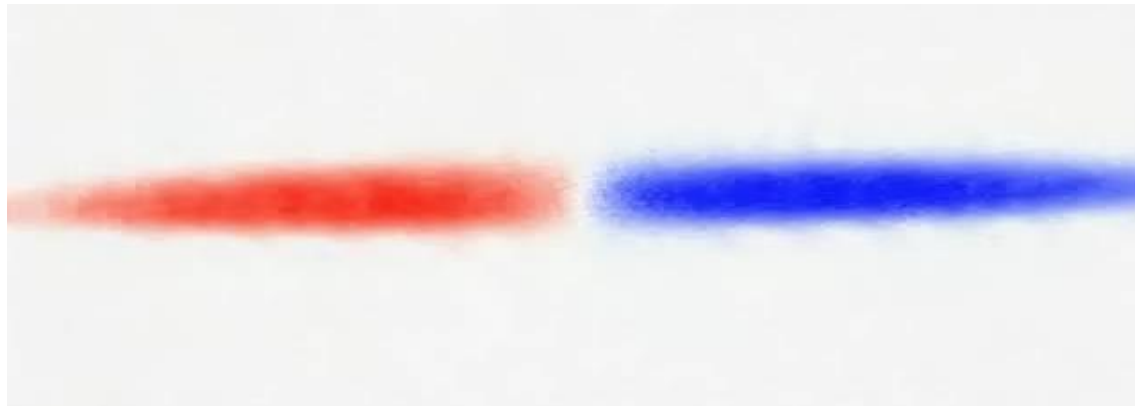
Harmonic
Trap

Little Fermi Collider (LFC)

Without Interactions

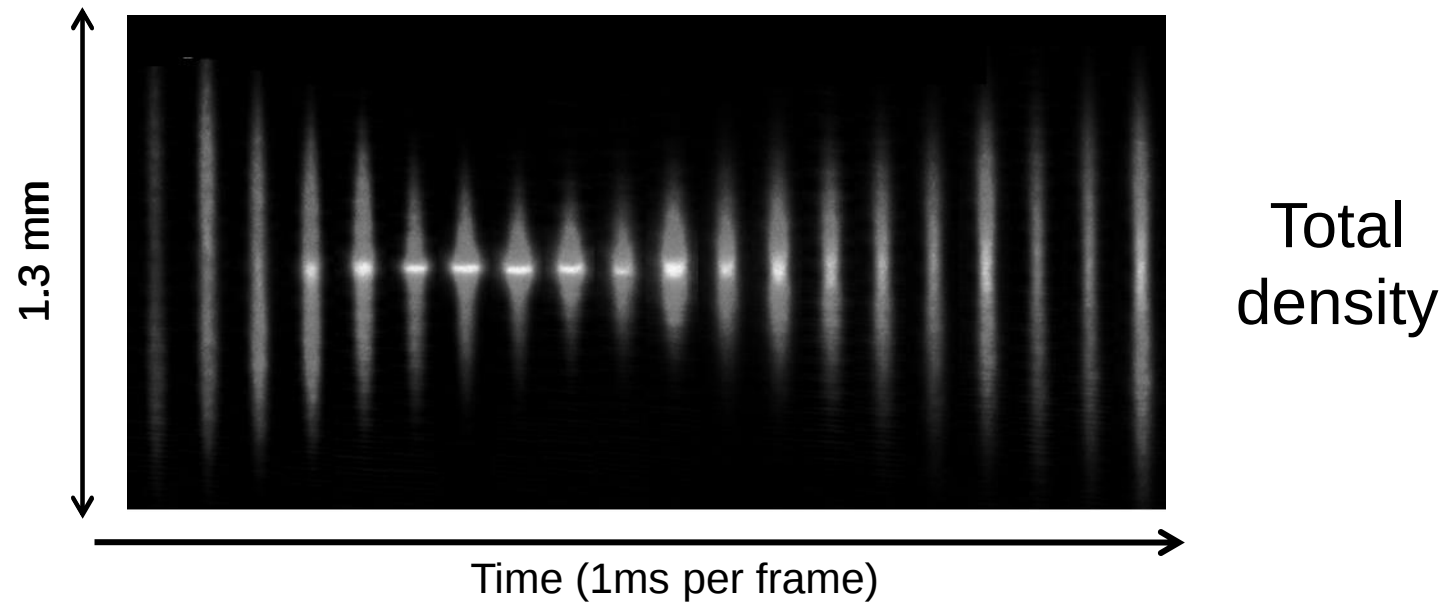
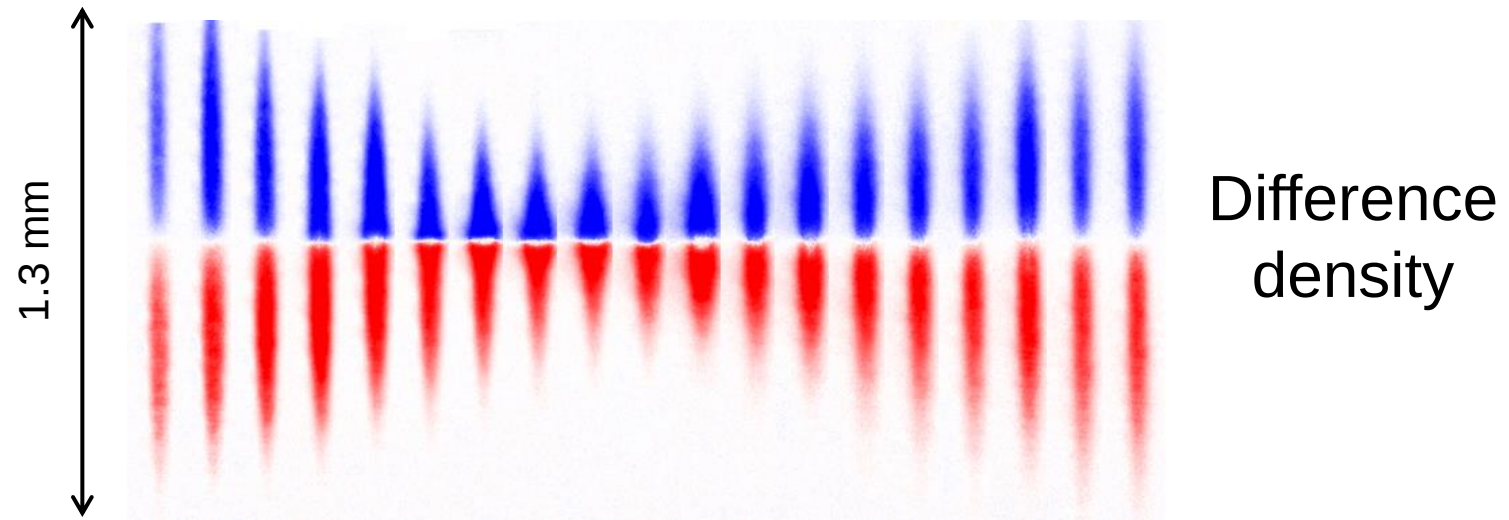


Resonant Interactions

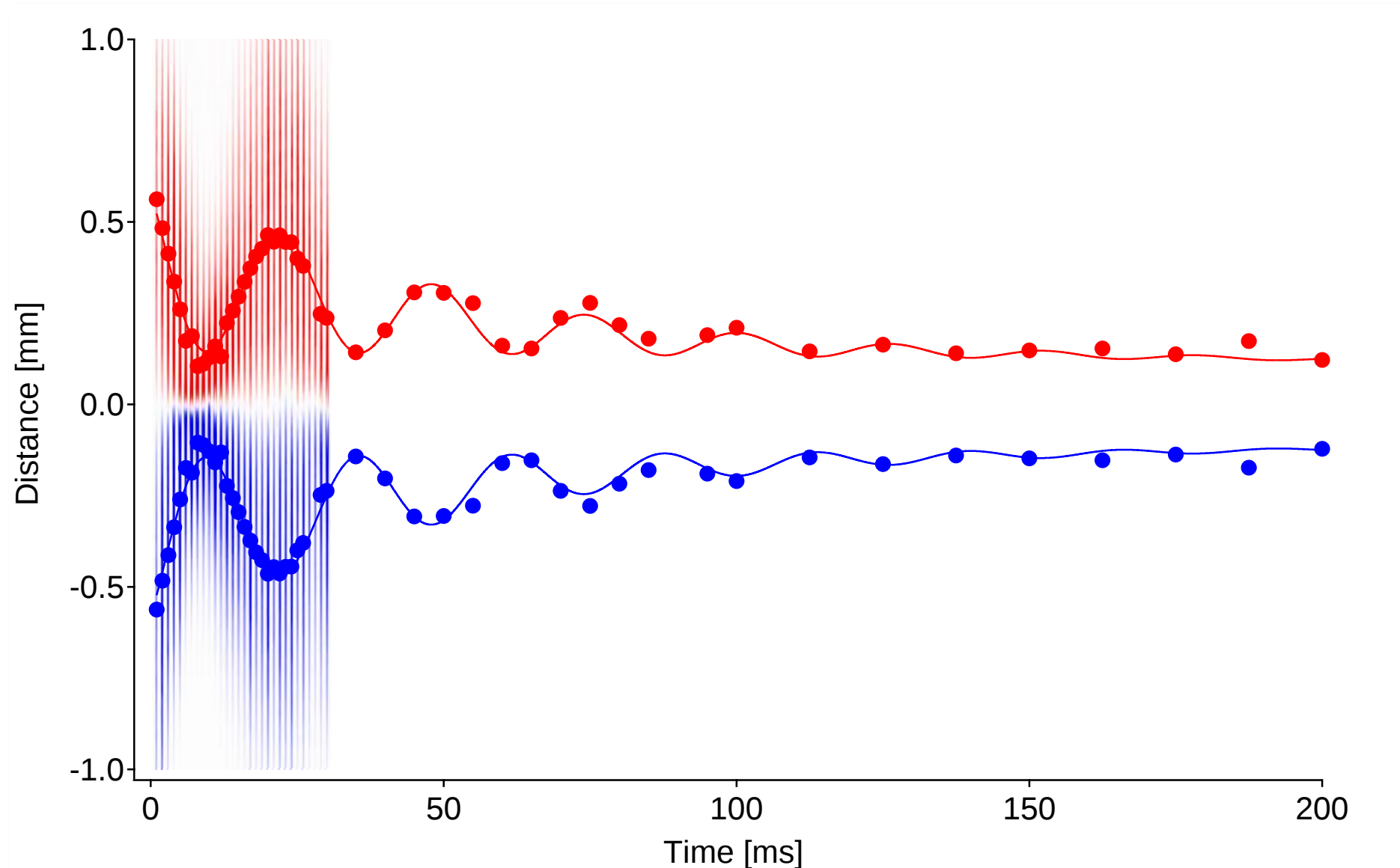


The bouncing gas

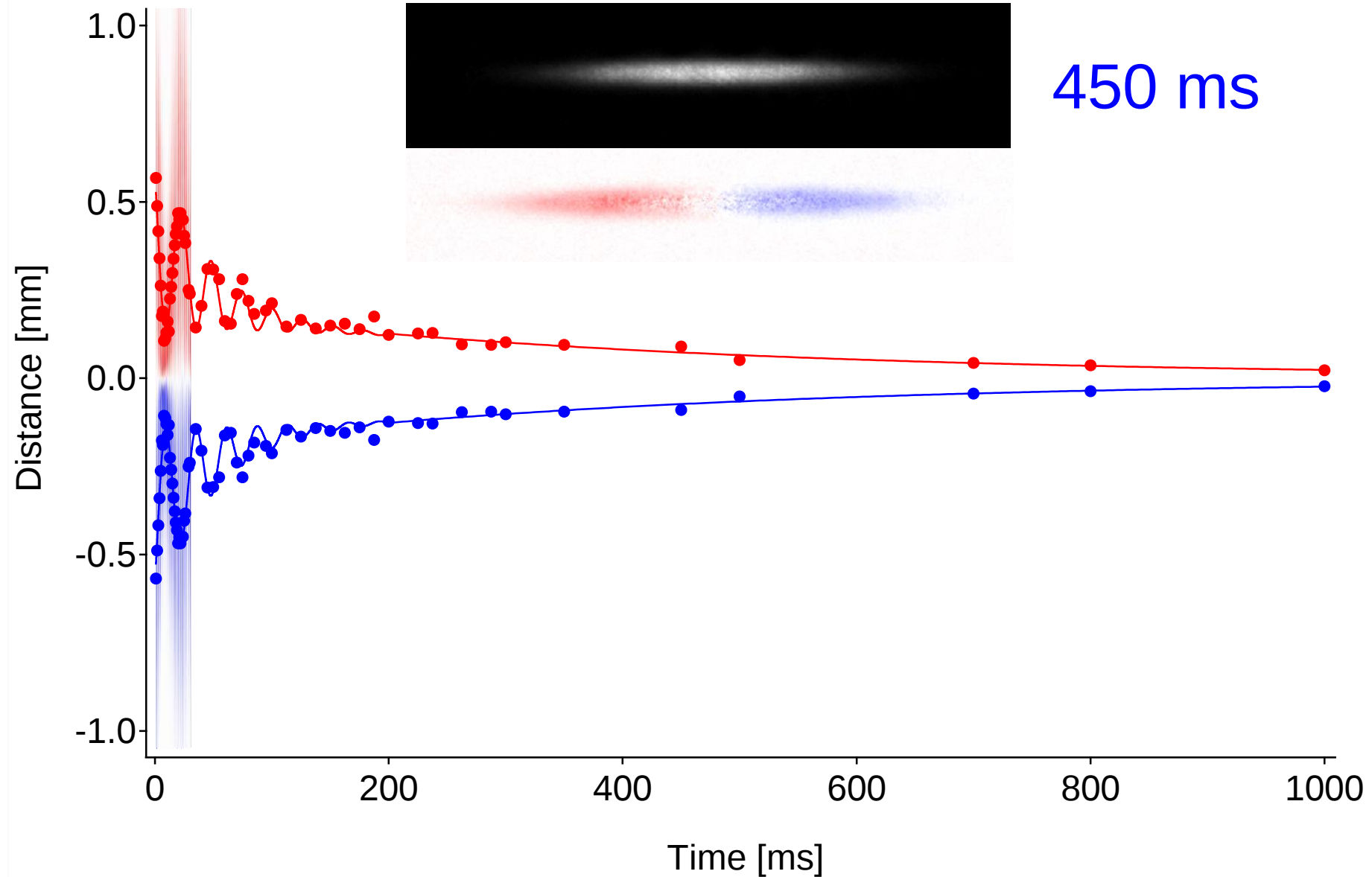
First collision



Later times



Much later times



Quantum limit of spin diffusion

Mean free path $\sim$ Interparticle spacing d

Diffusion constant:

$D \sim$ mean free path $\times$ average velocity

$$\cancel{d} \times \frac{\hbar}{m\cancel{d}}$$

$$D \sim \frac{\hbar}{m} = \frac{\text{Planck's constant}}{\text{Particle mass}} = \frac{(0.1 \text{ nm})^2}{1 \text{ fs}}$$

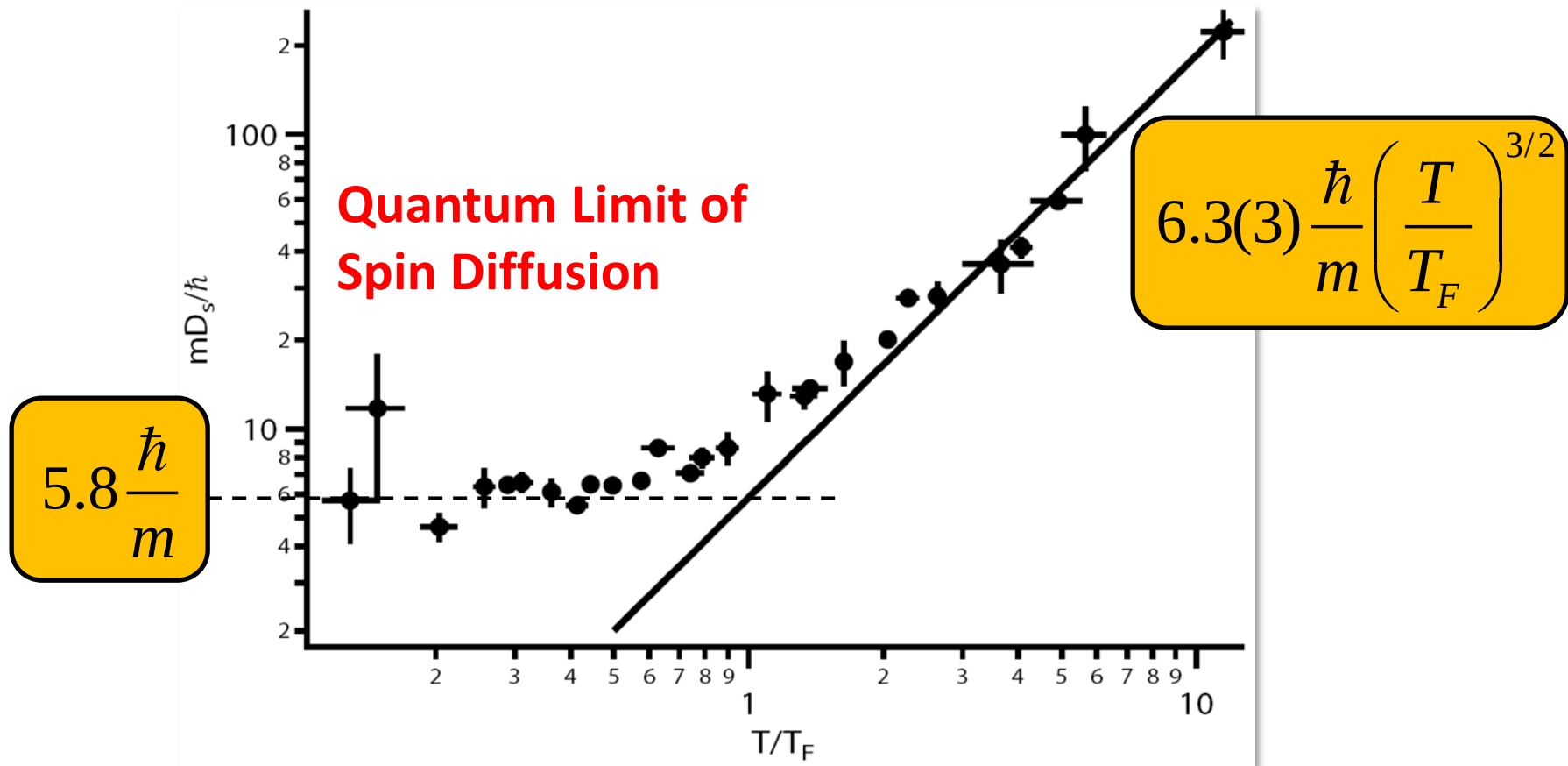
$\rightarrow$ Quantum Limit of Diffusion

In a hot relativistic fluid (e.g. Quark-Gluon Plasma): $D \sim \frac{\hbar c^2}{T}$
 $mc^2 \rightarrow T$

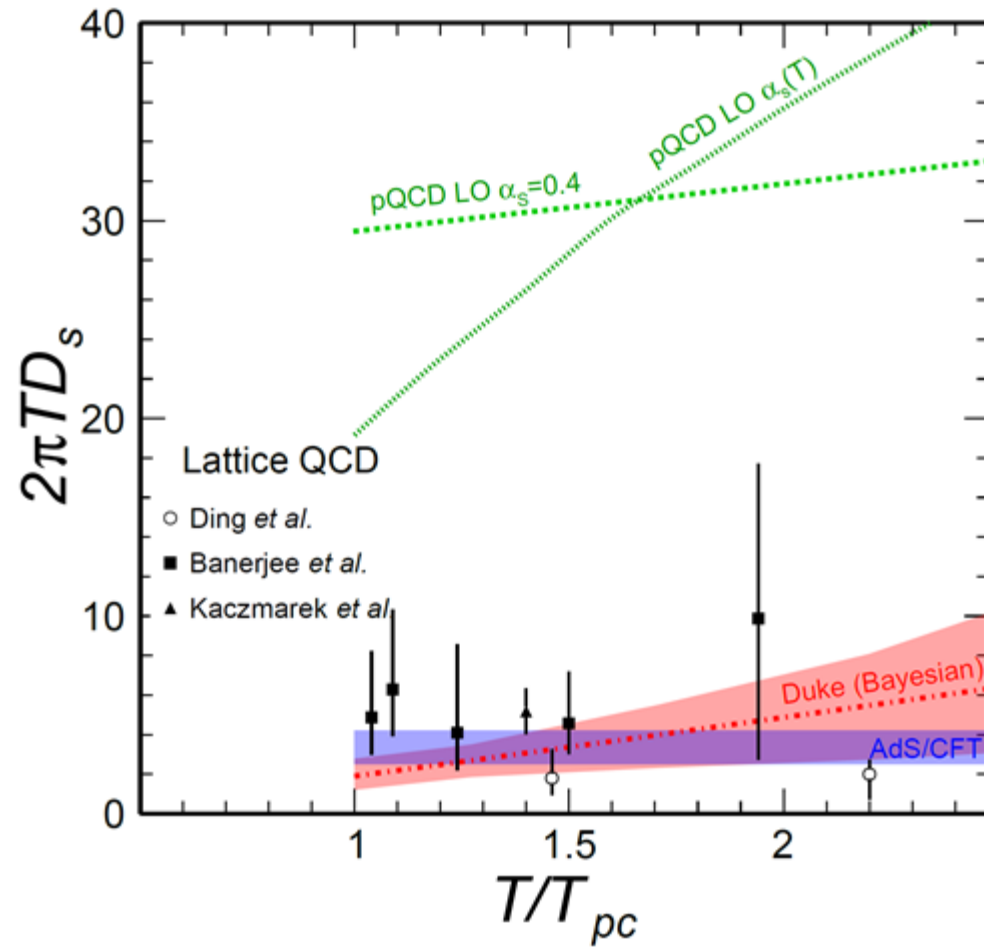
Spin Diffusion vs Temperature

Spin current = $-D \cdot \text{Spin density gradient}$

Universal high-T behavior:



Diffusivity in the Quark-Gluon Plasma



**Can Fermi Gases become
superfluid?**

Superconductivity

Electrons are Fermions

Discovery of superconductivity 1911



Heike
Kamerlingh-Onnes
1911

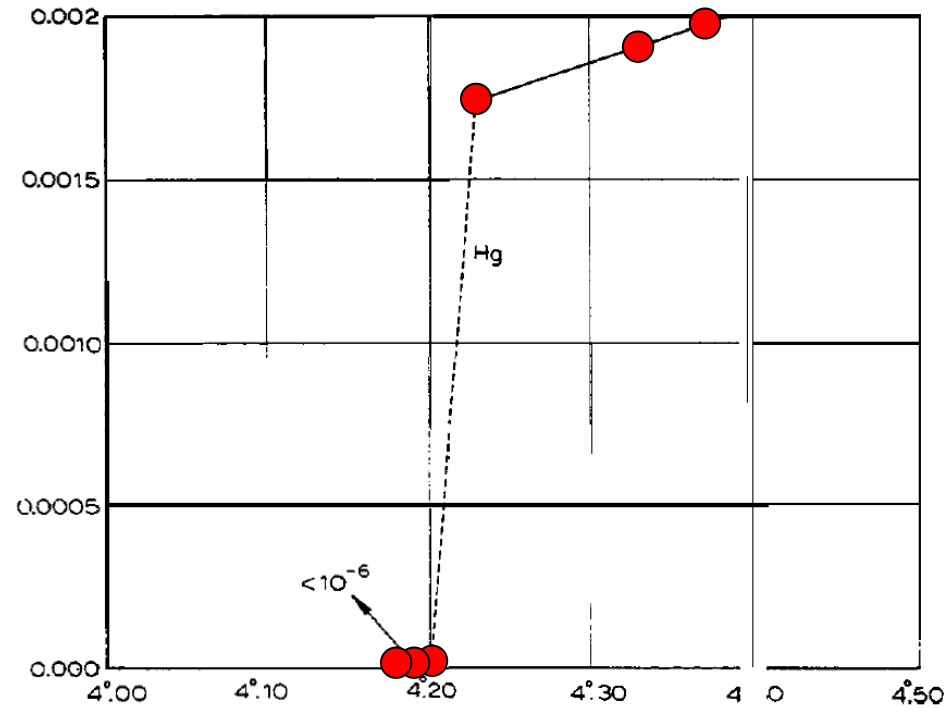


Fig. 17.

Fermionic Superfluidity

Condensation of Fermion Pairs

- **Helium-3** (Lee, Osheroff, Richardson 1971)
- **Superconductors: *Charged* superfluids of electron pairs**
Frictionless flow $\Leftrightarrow$ Resistance-less current



John Bardeen



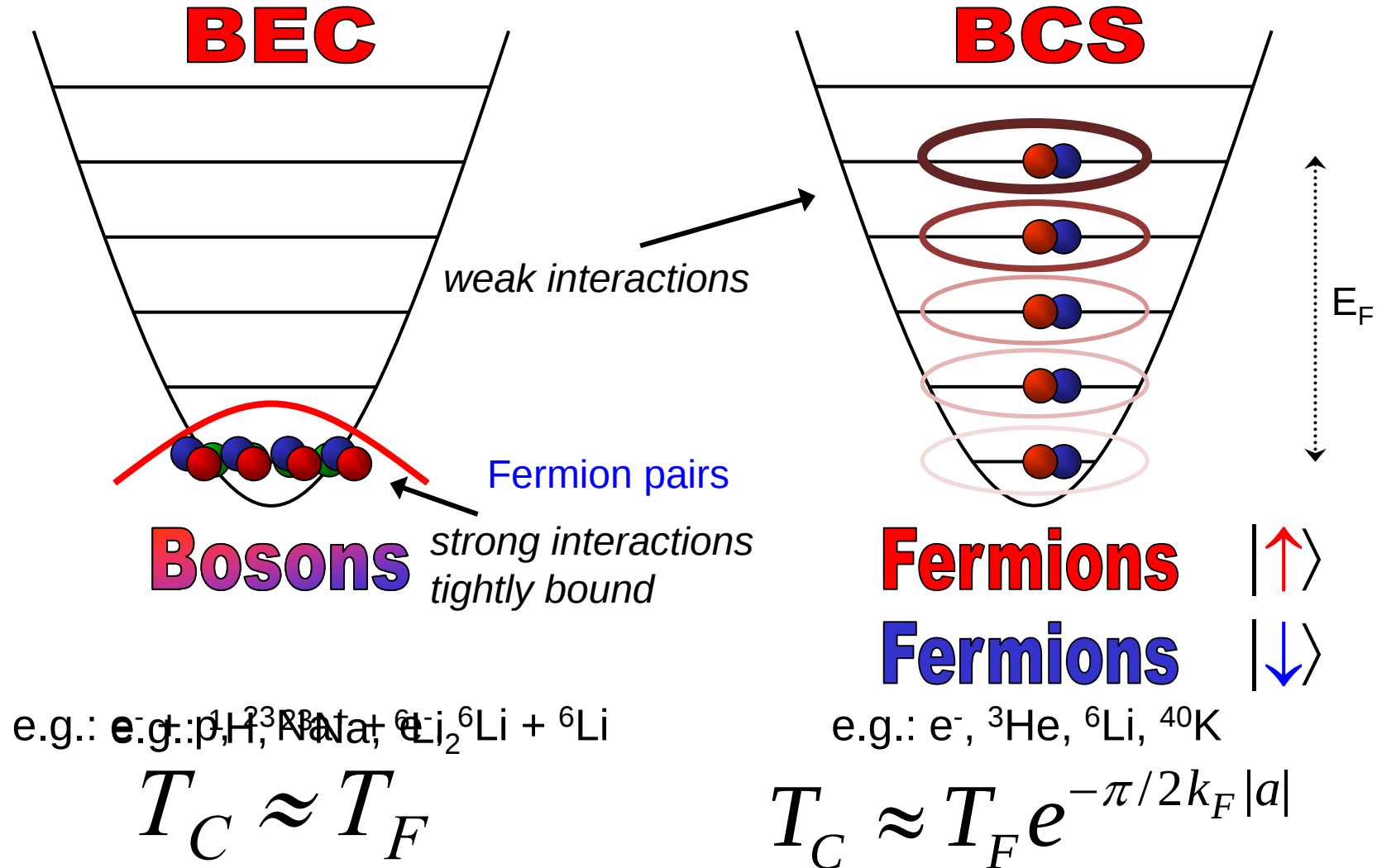
Leon N. Cooper



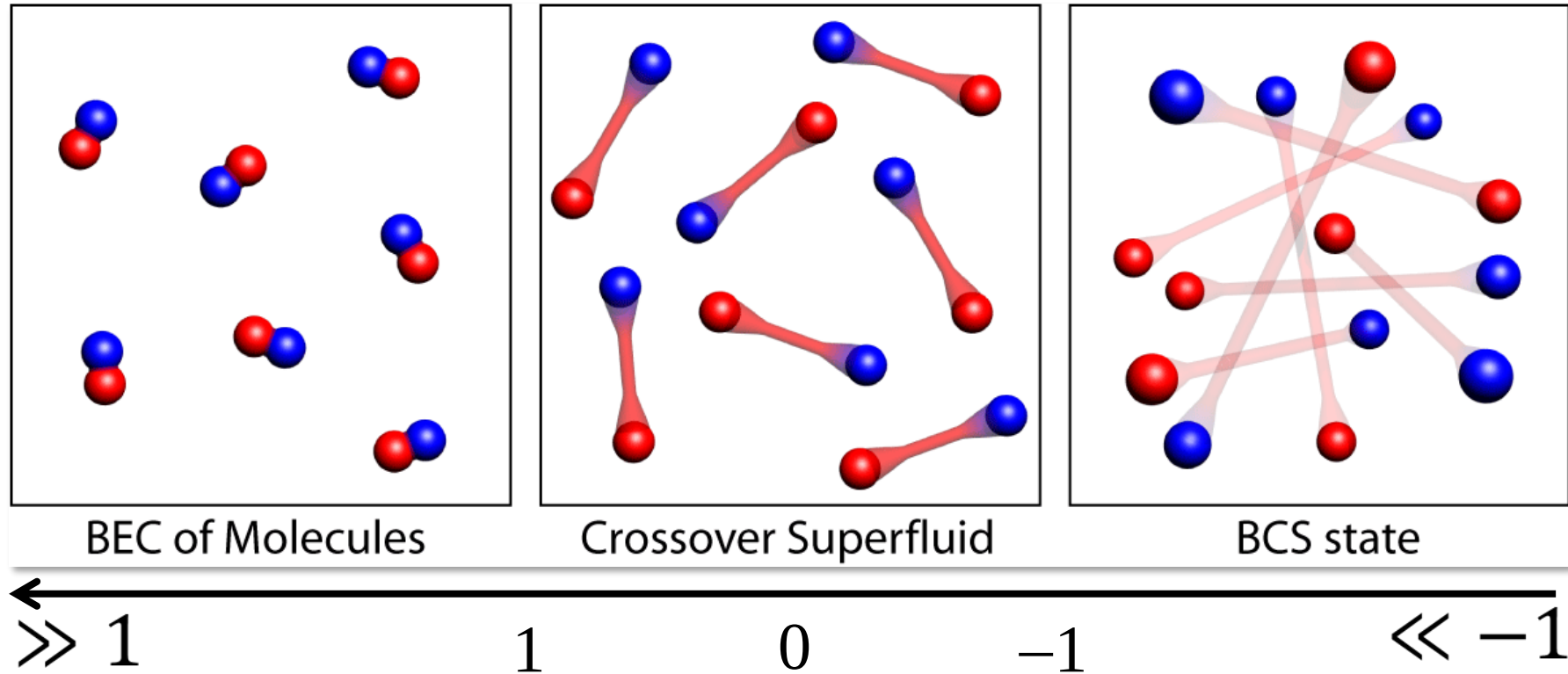
John R. Schrieffer

- **Neutron stars**
In the core: **Quark superfluid**

Bosons vs Fermions



From BEC to BCS



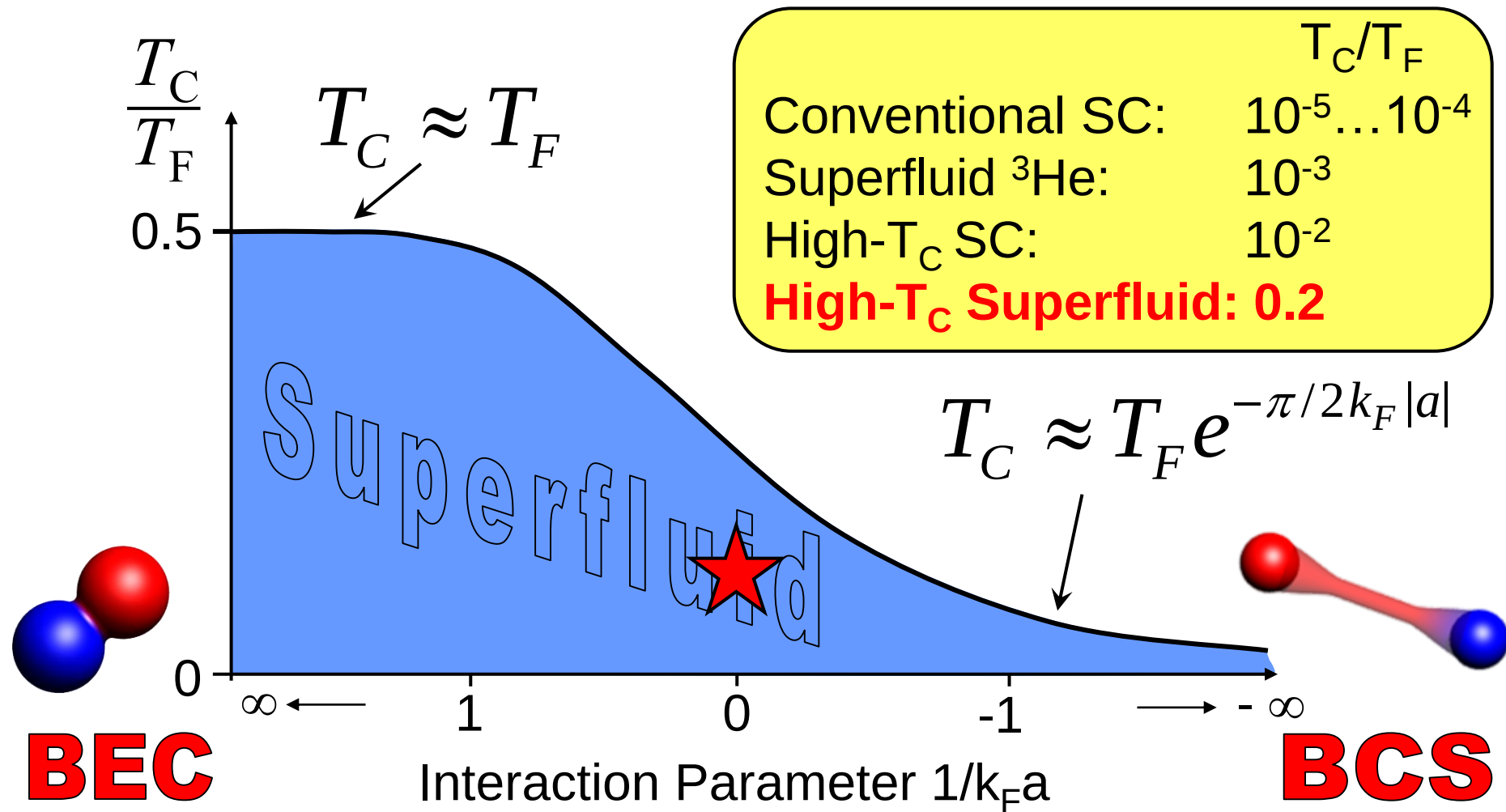
Weakly Interacting Bosons $(k_F a)^{-1} = \frac{\text{Interparticle Distance}}{\text{Scattering Length}}$

→ Strongly Interacting Bosons

→ Strongly Interacting Fermions

→ Weakly Interacting Fermions

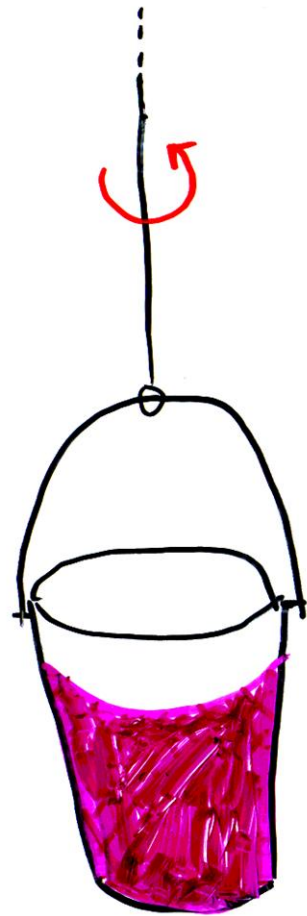
Critical Temperature for Fermionic Superfluidity



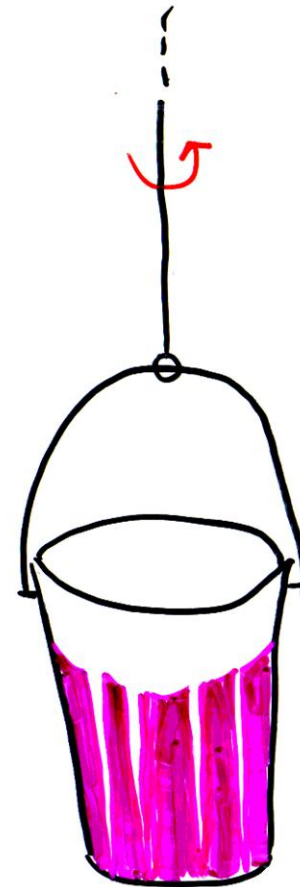
Scaled to the density of electrons in solids:

Superconductivity far above room temperature!

Rotating Fluids



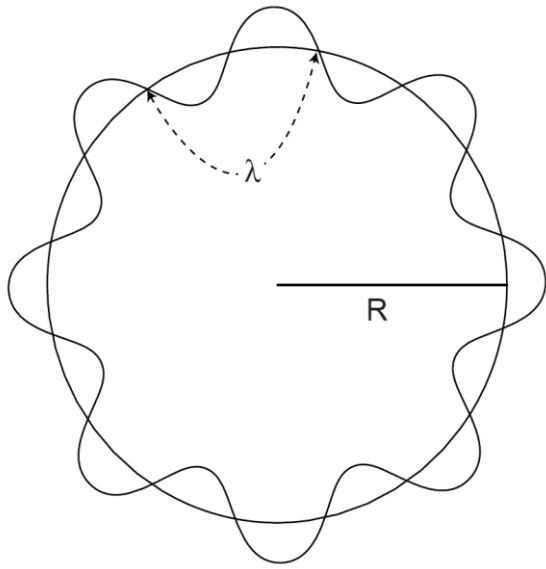
Normal



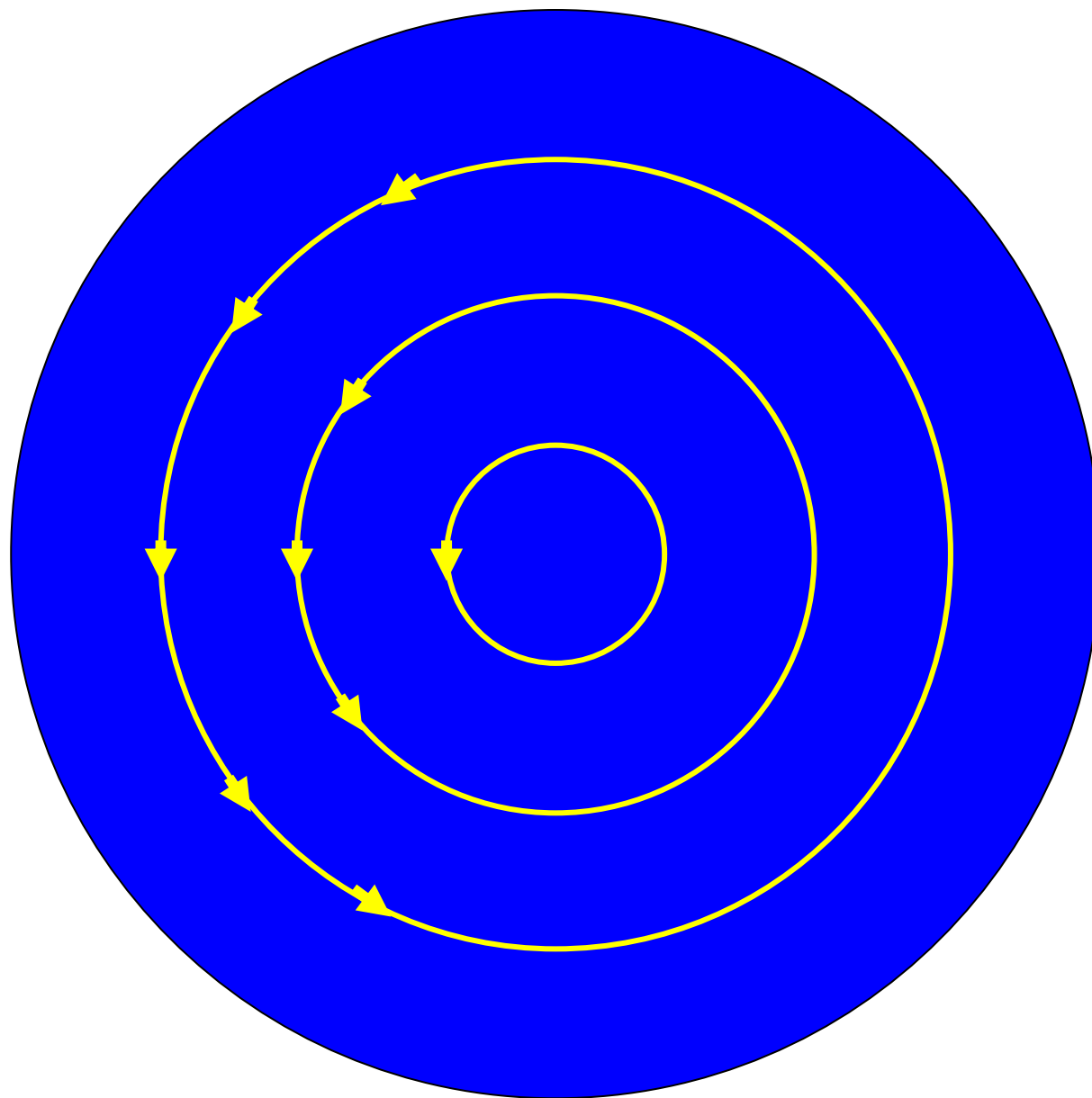
Quantum

Fluid

Rotating superfluid



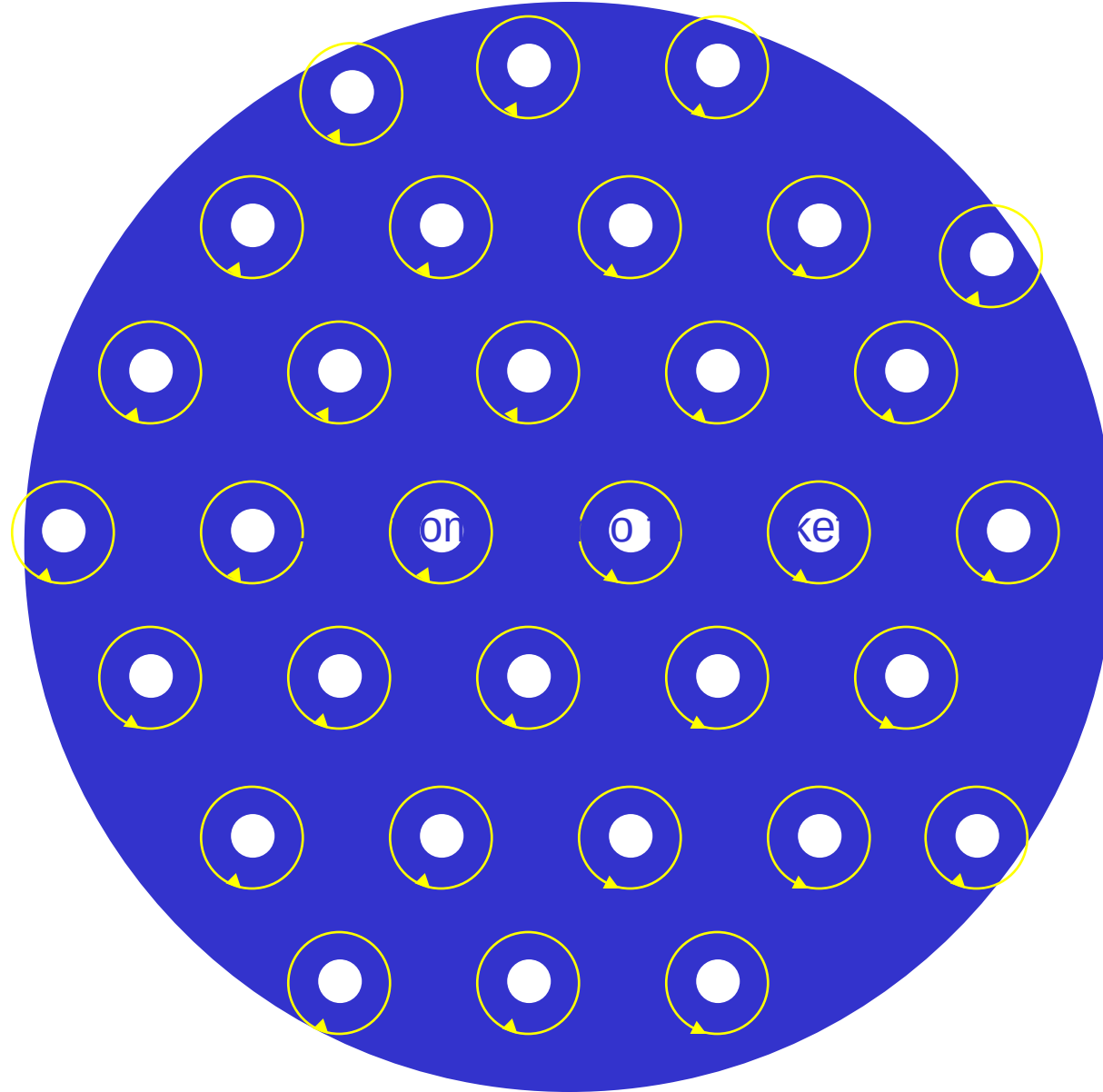
Does it rotate like this?



Like this!



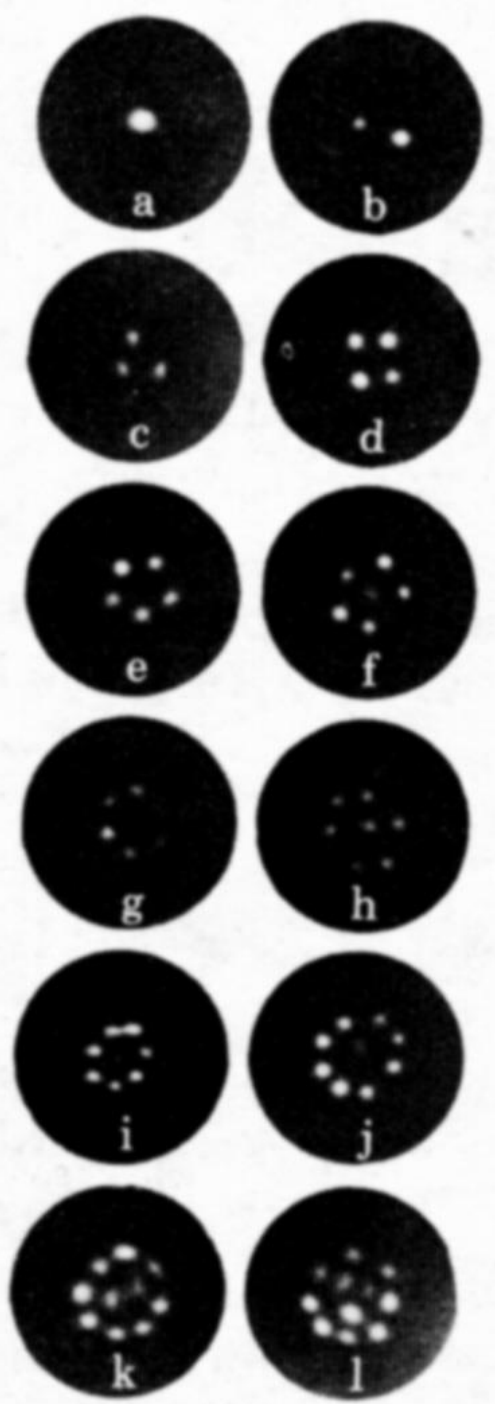
Aleksei A.
Abrikosov



Abrikosov lattice (triangular lattice)

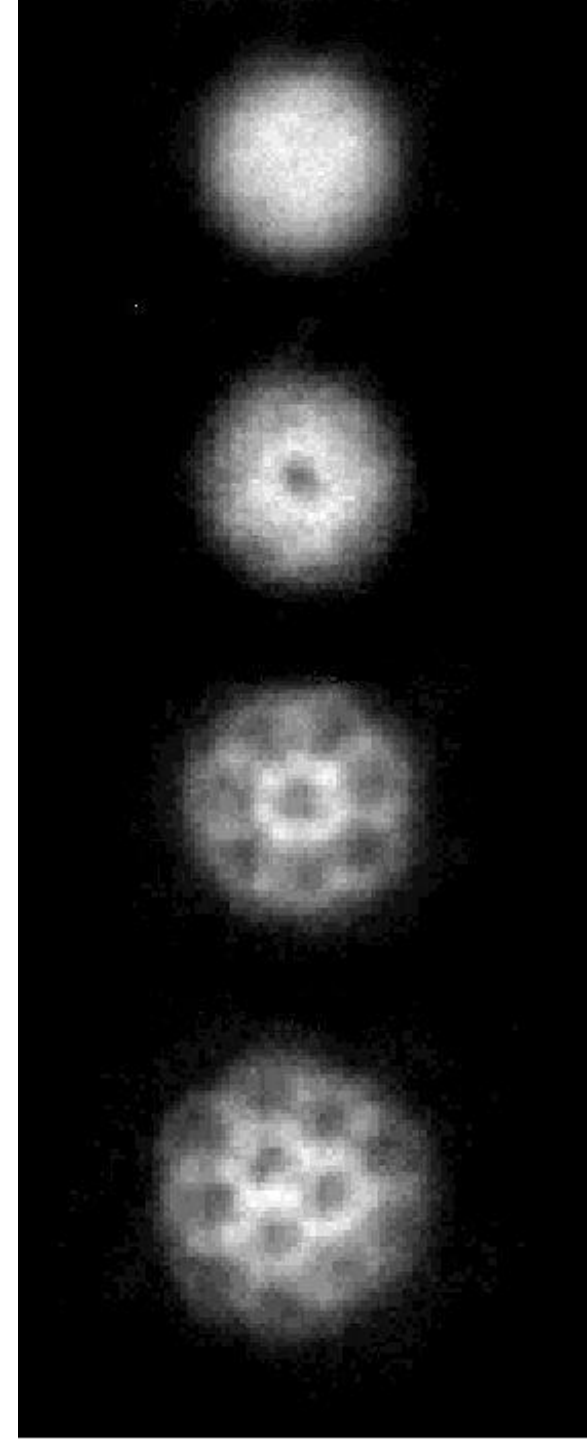
Vortex Arrays *in Bosonic Gases / Fluids*

Berkeley
(R.E. Packard, 1979)
Helium-4



ENS
(J. Dalibard, 2000)
Rubidium BEC

*Also: Phase engineering
of single vortices in BEC:
JILA (1999)*



THE DIRECT OBSERVATION OF INDIVIDUAL FLUX LINES IN TYPE II SUPERCONDUCTORS

U. ESSMANN and H. TRÄUBLE

*Institut für Physik am Max-Planck-Institut für Metallforschung, Stuttgart and
Institut für theoretische und angewandte Physik der Technischen Hochschule Stuttgart*

Received 4 April 1967

Neutral superfluids under rotation

$$\vec{F} = 2m\vec{v} \times \vec{\omega}$$

Coriolis force in rotating frame



Superconductors in magnetic field

$$\vec{F} = q\vec{v} \times \vec{B}$$

Lorentz Force

U. Essmann and H. Träuble,
Physics Letters A, **24**, 526 (1967)

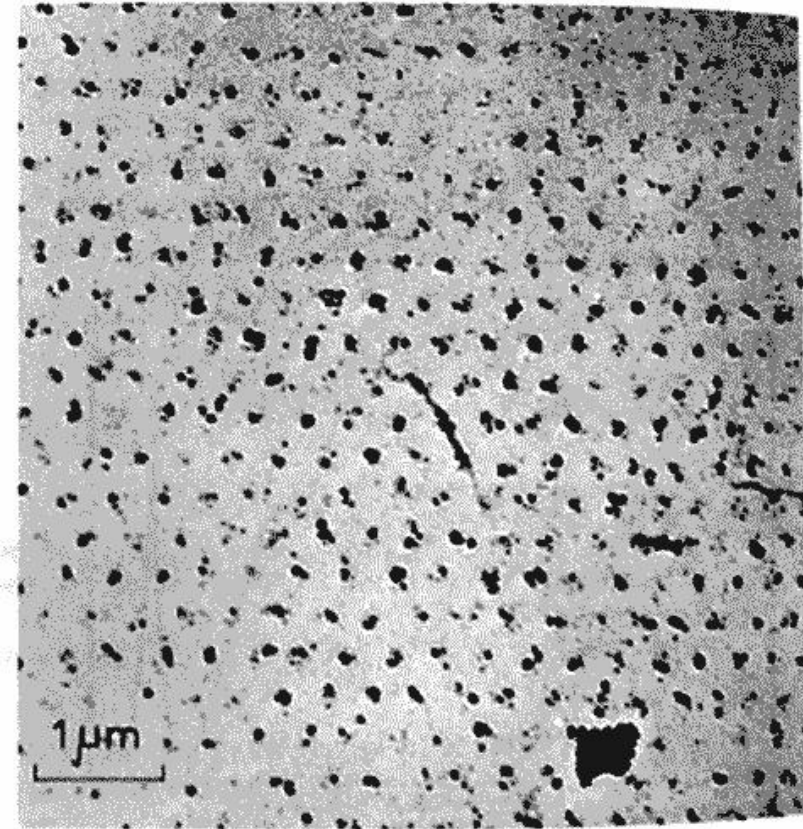
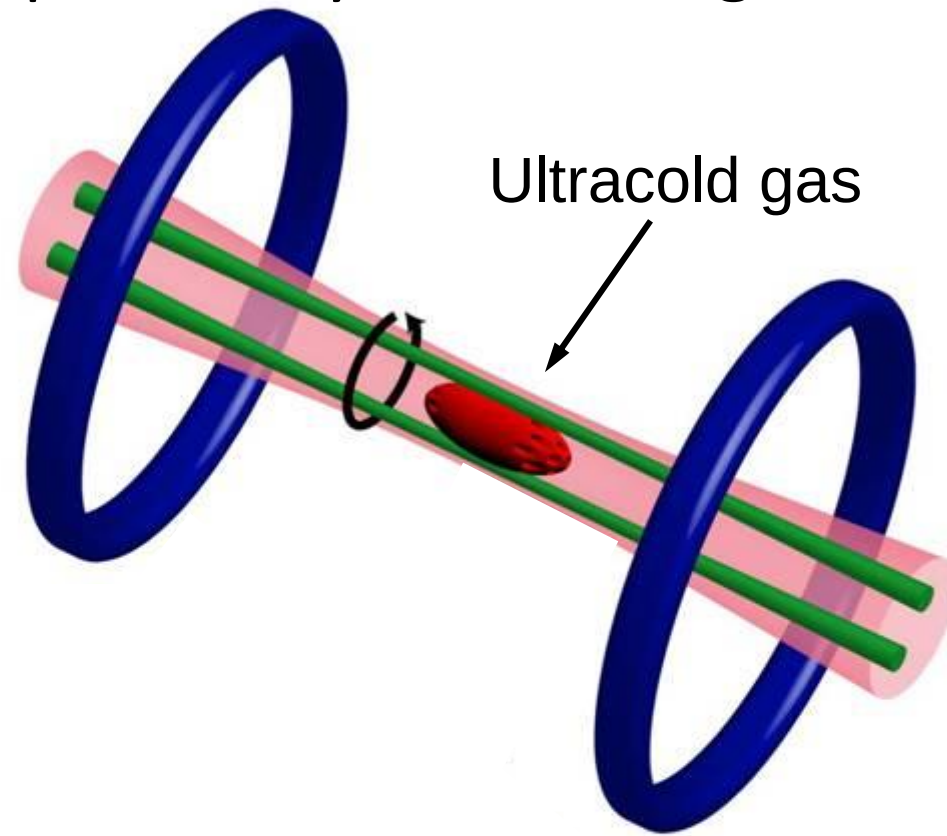
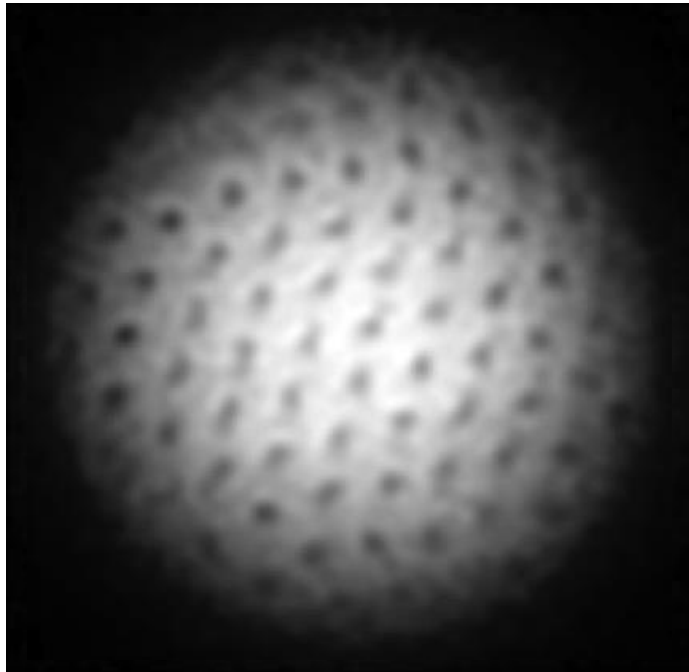
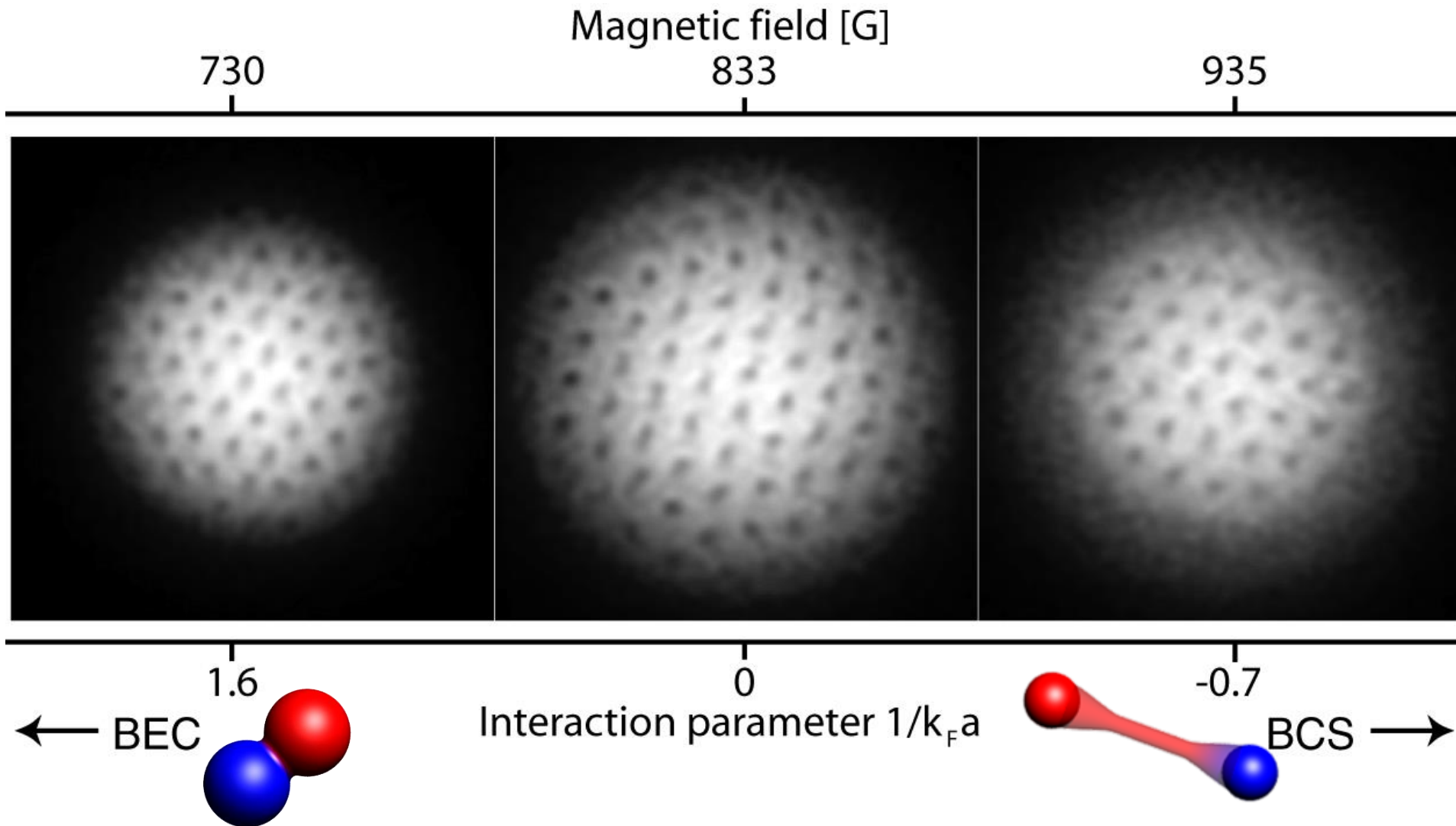


Fig. 1. "Perfect" triangular lattice of flux lines on the surface of a lead-4at% indium rod at 1.1°K. The black dots consist of small cobalt particles which have been stripped from the surface with a carbon replica.

Demonstration of superfluidity in a Fermi gas



Vortex lattices in the BEC-BCS crossover



M.W. Zwierlein, J.R. Abo-Shaeer, A. Schirotzek, C.H. Schunck, W. Ketterle,
Nature 435, 1047-1051 (2005)

The Unitary Fermi Gas

Only two length scales:

Interparticle spacing

De Broglie wavelength

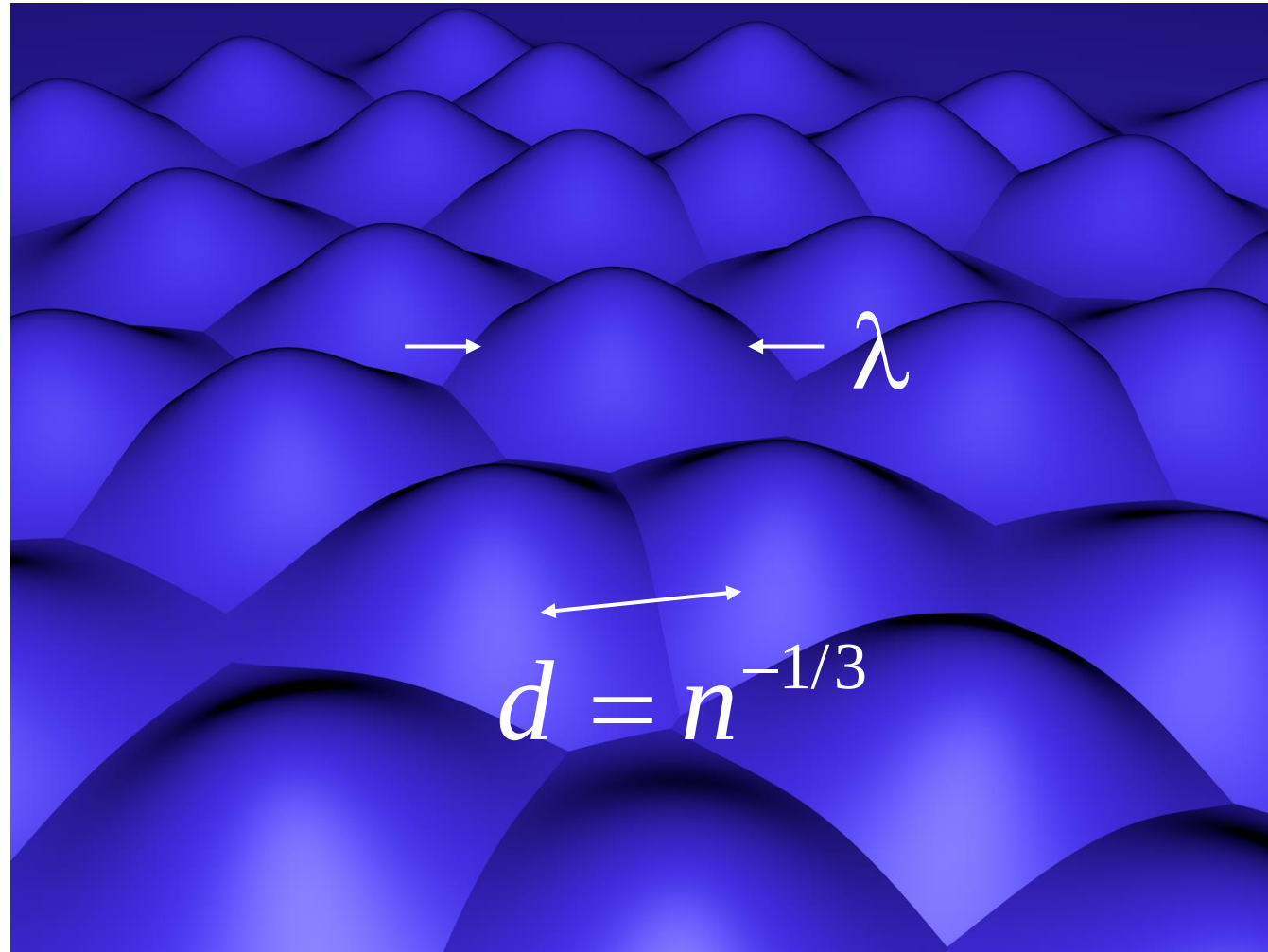
$$n^{-1/3}$$
$$\lambda_{\text{dB}}$$

Only two corresp. energy scales:

Fermi energy

Temperature

$$E_F$$
$$k_B T$$



The Unitary Fermi Gas

Only two length scales:

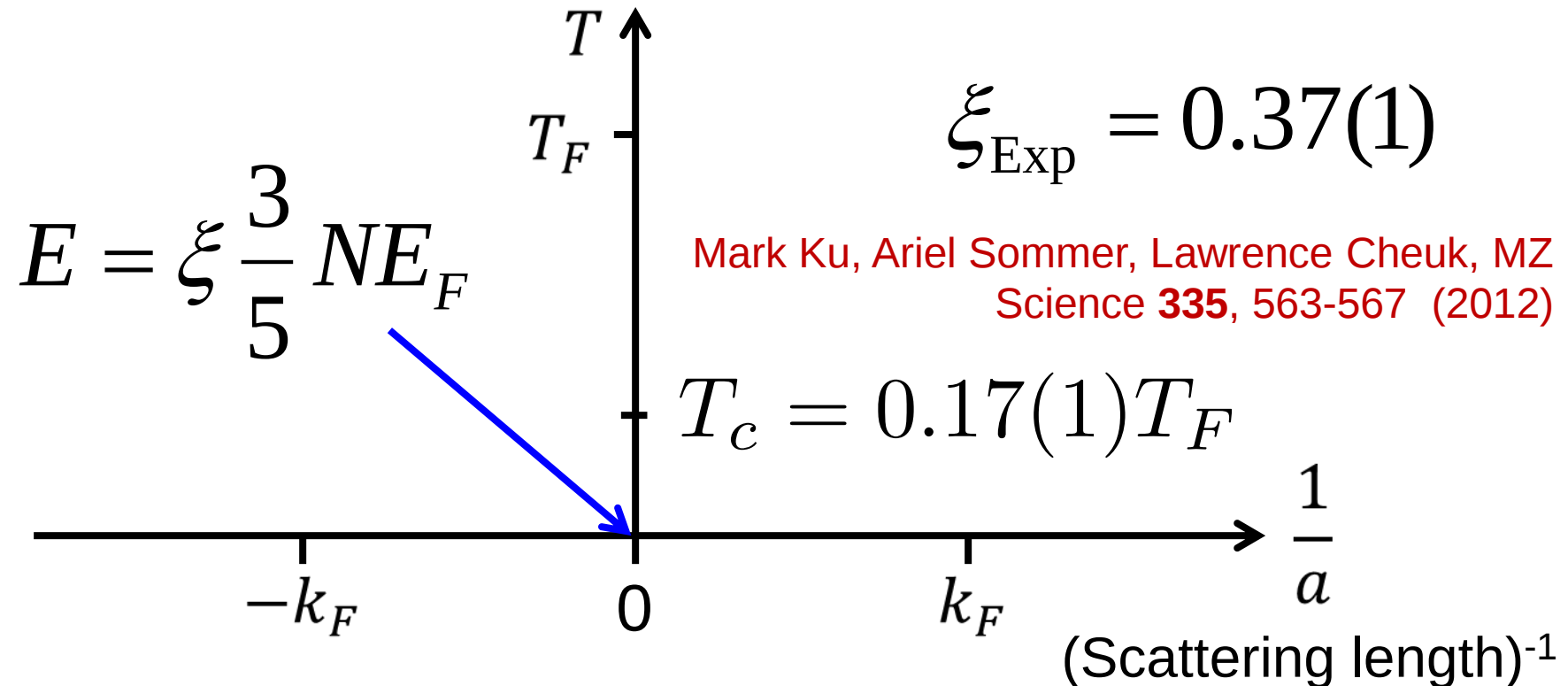
Interparticle spacing $n^{-1/3}$
De Broglie wavelength λ_{dB}

Only two corresp. energy scales:

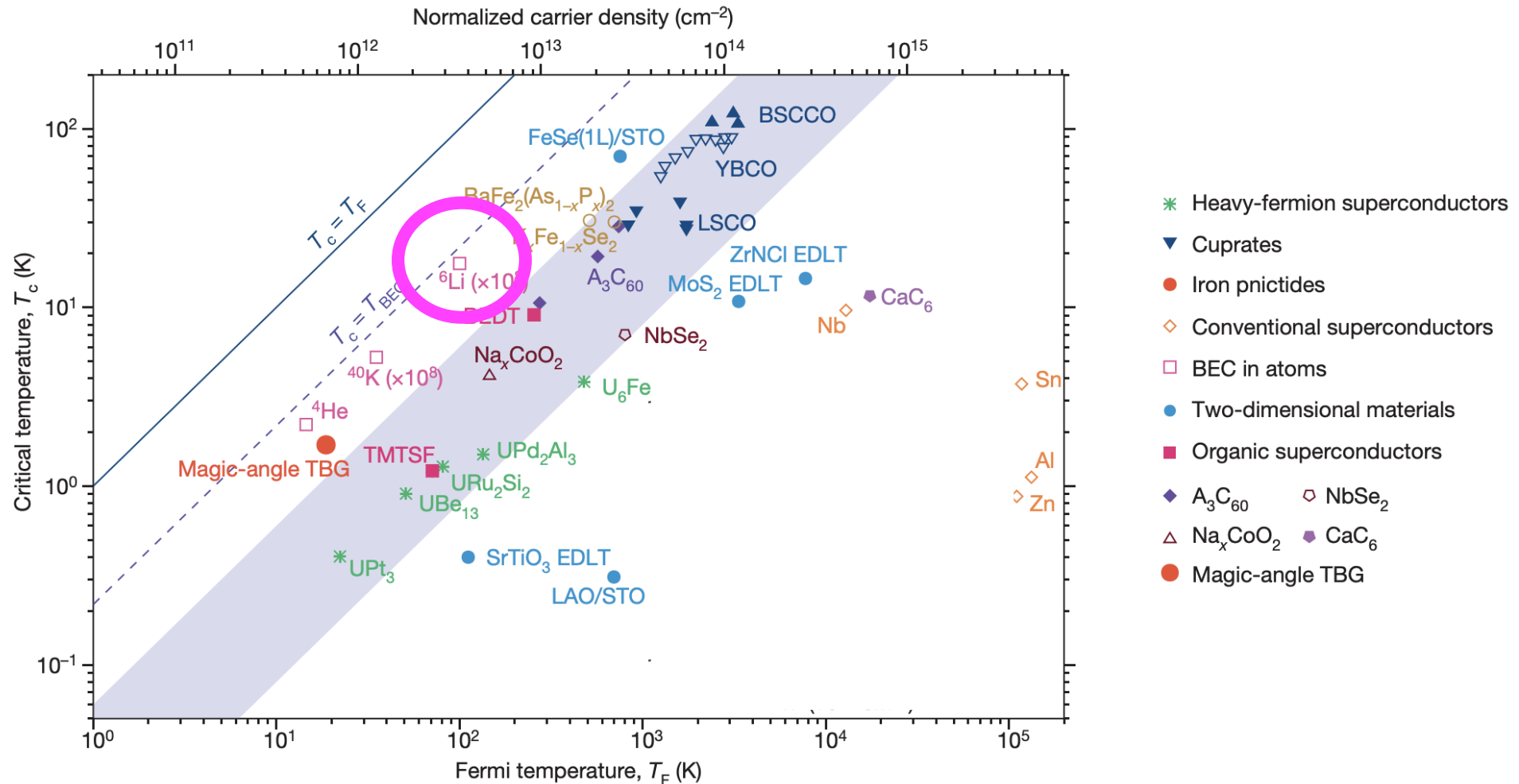
Fermi energy E_F
Temperature $k_B T$

→ Universal equation of state

Duke/NC State, ENS, JILA, Innsbruck, Swinburne, MIT



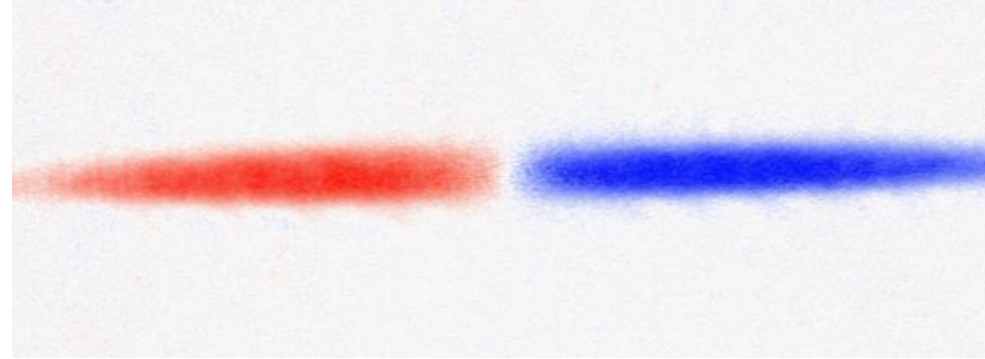
Comparison to other dilute Fermi systems



From Cao et al., Jarillo-Herrero, Nature 2018

Transport in Strongly Interacting Systems

e.g. Unitary Gas @ Feshbach Resonance



Mean free path $\sim$ Interparticle spacing d

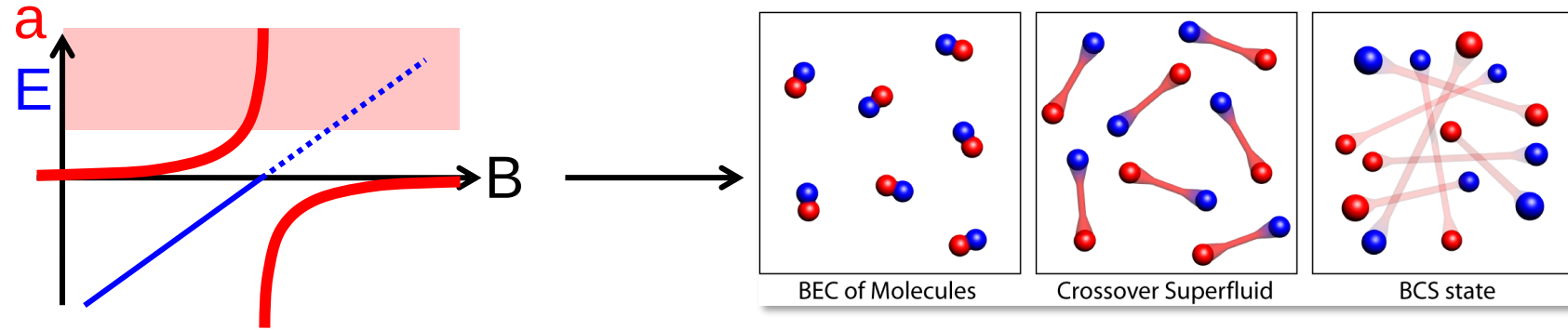
Average velocity $v \propto \frac{\hbar}{md}$

Quantum Limit of Diffusion (charge, spin, momentum, thermal)

$$D \propto vl \propto \frac{\hbar}{m}$$

Sound in strongly interacting Fermi Gases

Feshbach Resonances $\rightarrow$ e.g. BEC-BCS Crossover



Quantum Gases $\rightarrow$ Quantum Fluids



Parth
Patel



Zhenjie
Yan



Biswaroop
Mukherjee

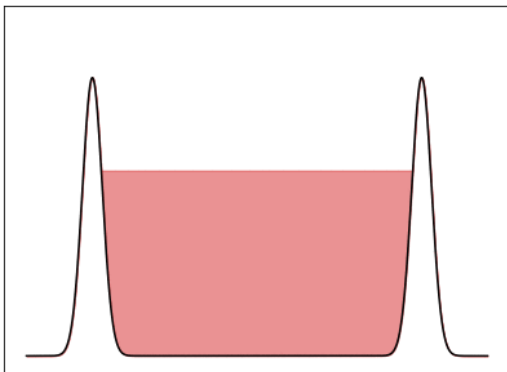
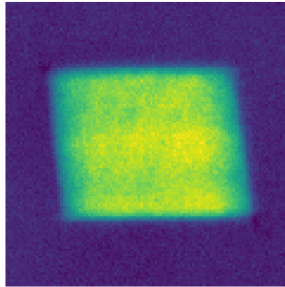
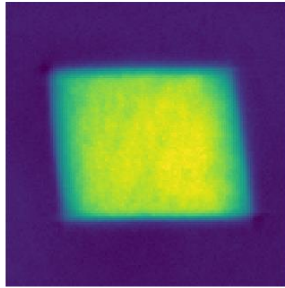
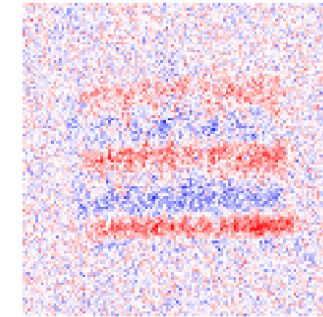
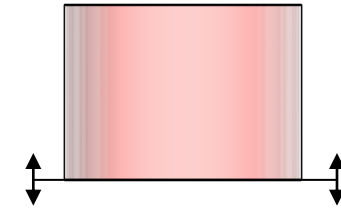
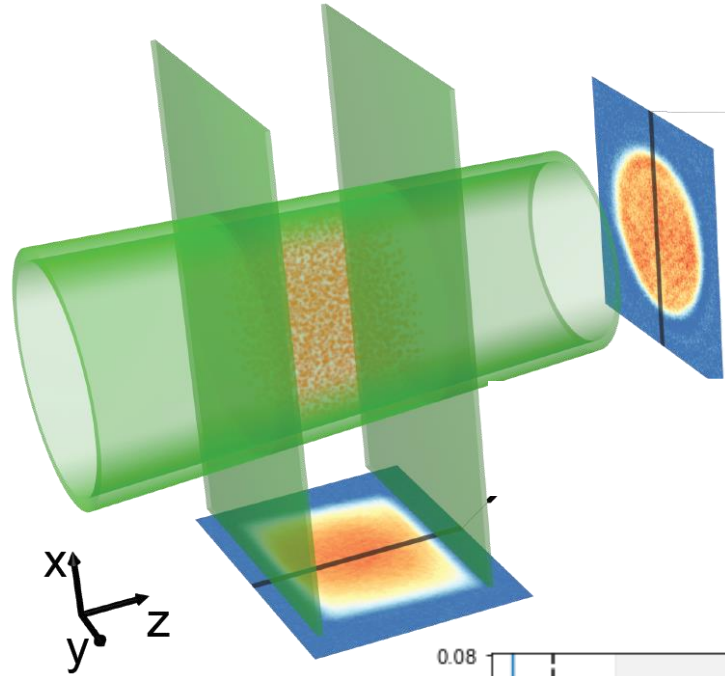


Dr. Richard
Fletcher

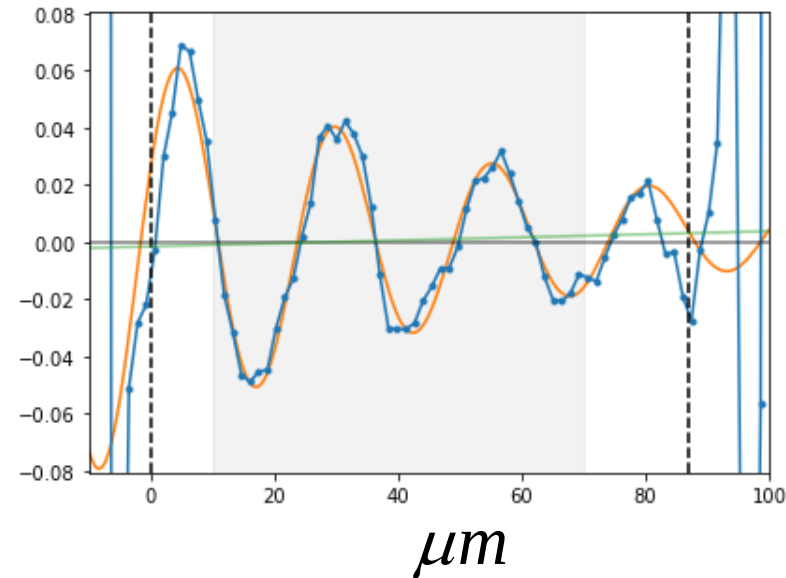


Dr. Julian
Struck

Creating Sound Waves in a Box

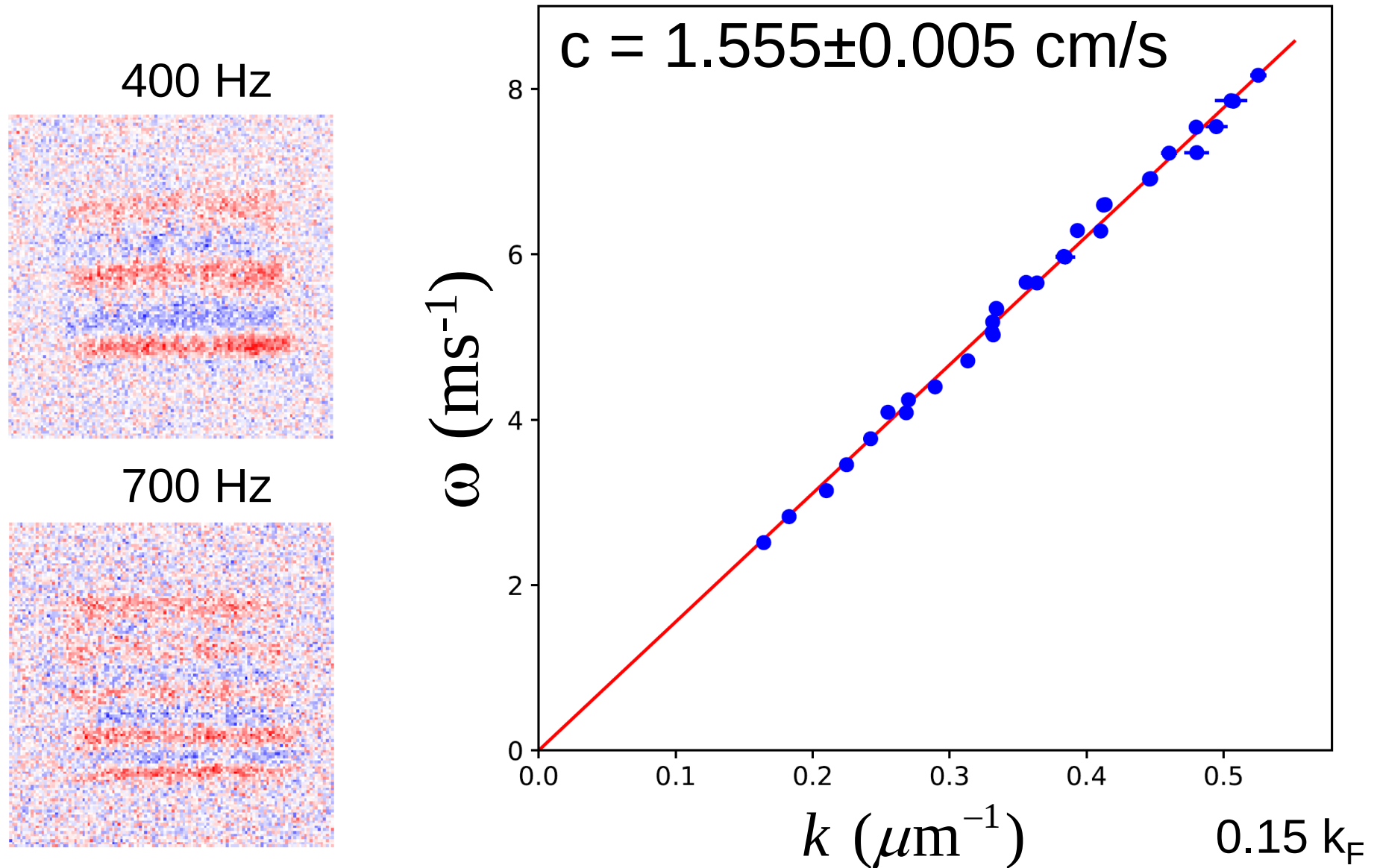


$$\frac{\delta\rho}{\rho}$$



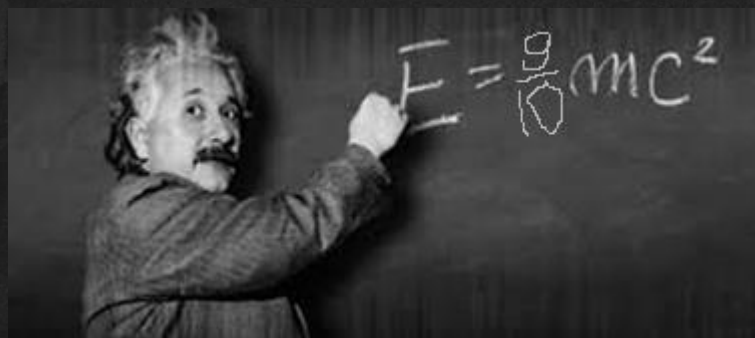
See also: 2D Bose gas, ENS (Dalibard),
3D Bose gas, Cambridge (Hadzibabic), Fermi gas, NCState (Thomas), 2D Fermi Gas, Hamburg (Moritz)

Dispersion relation



$$E = \frac{9}{10} m c^2$$

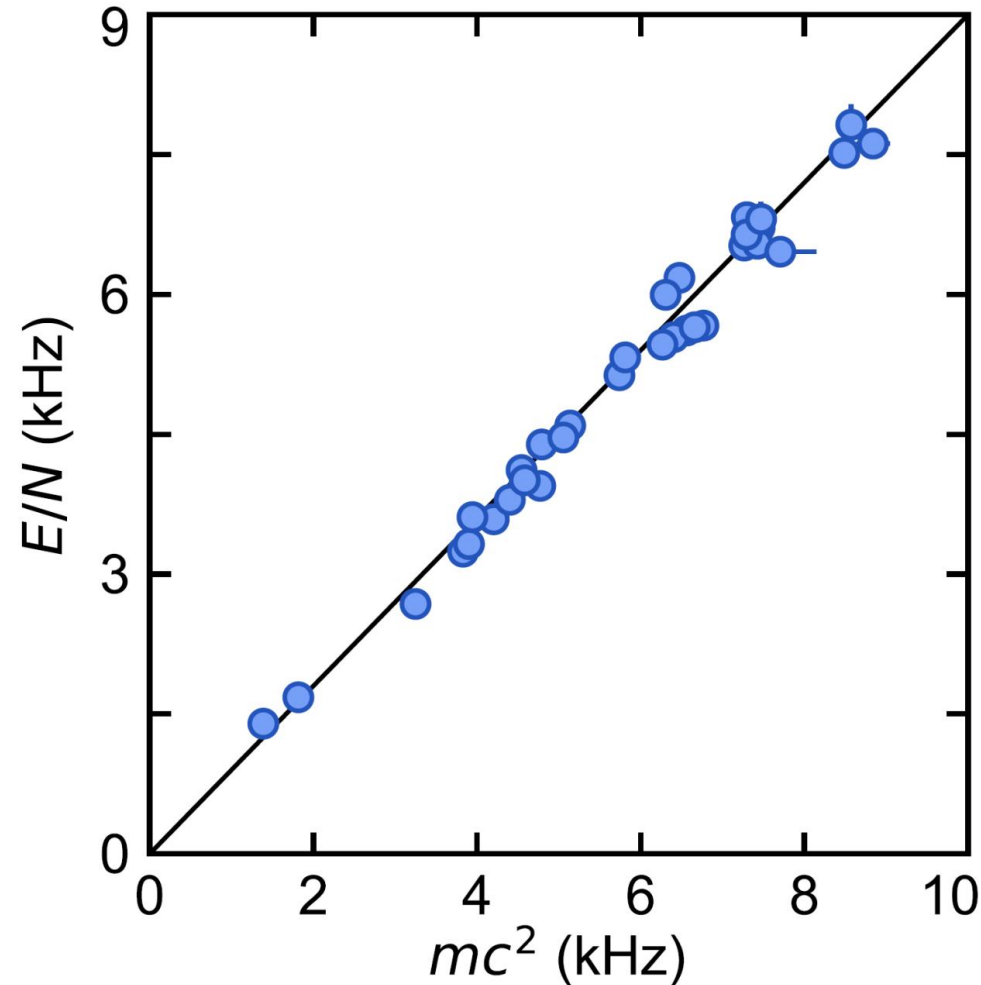
?



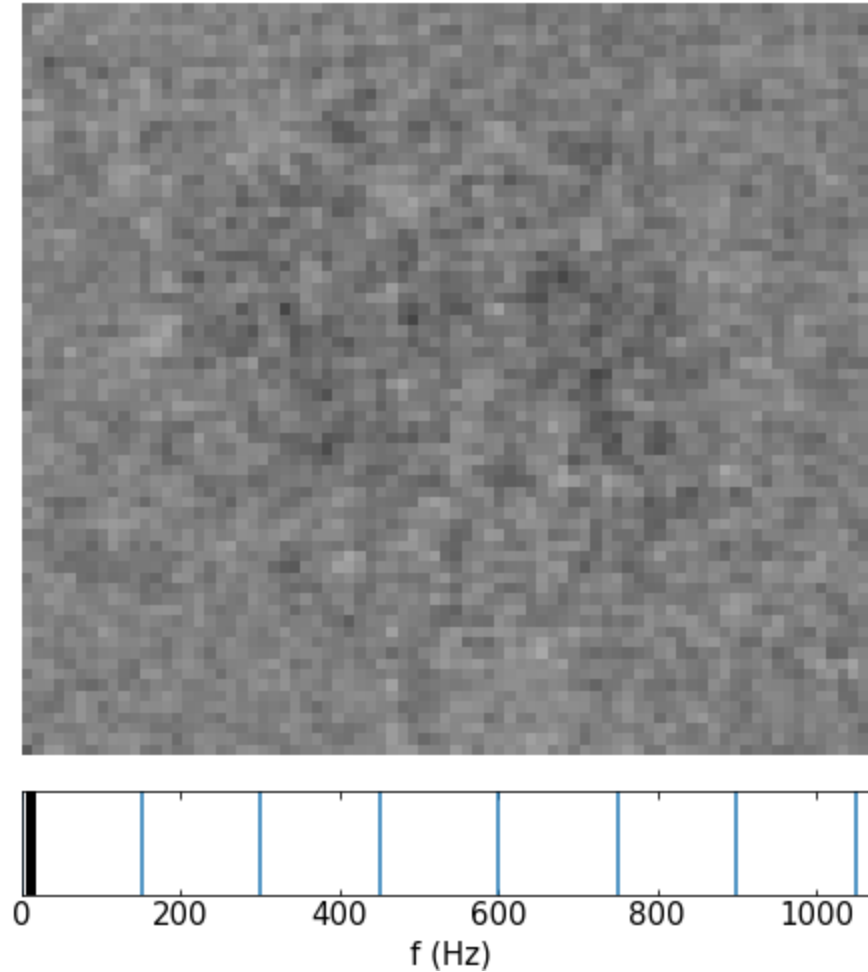
$$E = \frac{9}{10} m c^2$$

At unitarity:

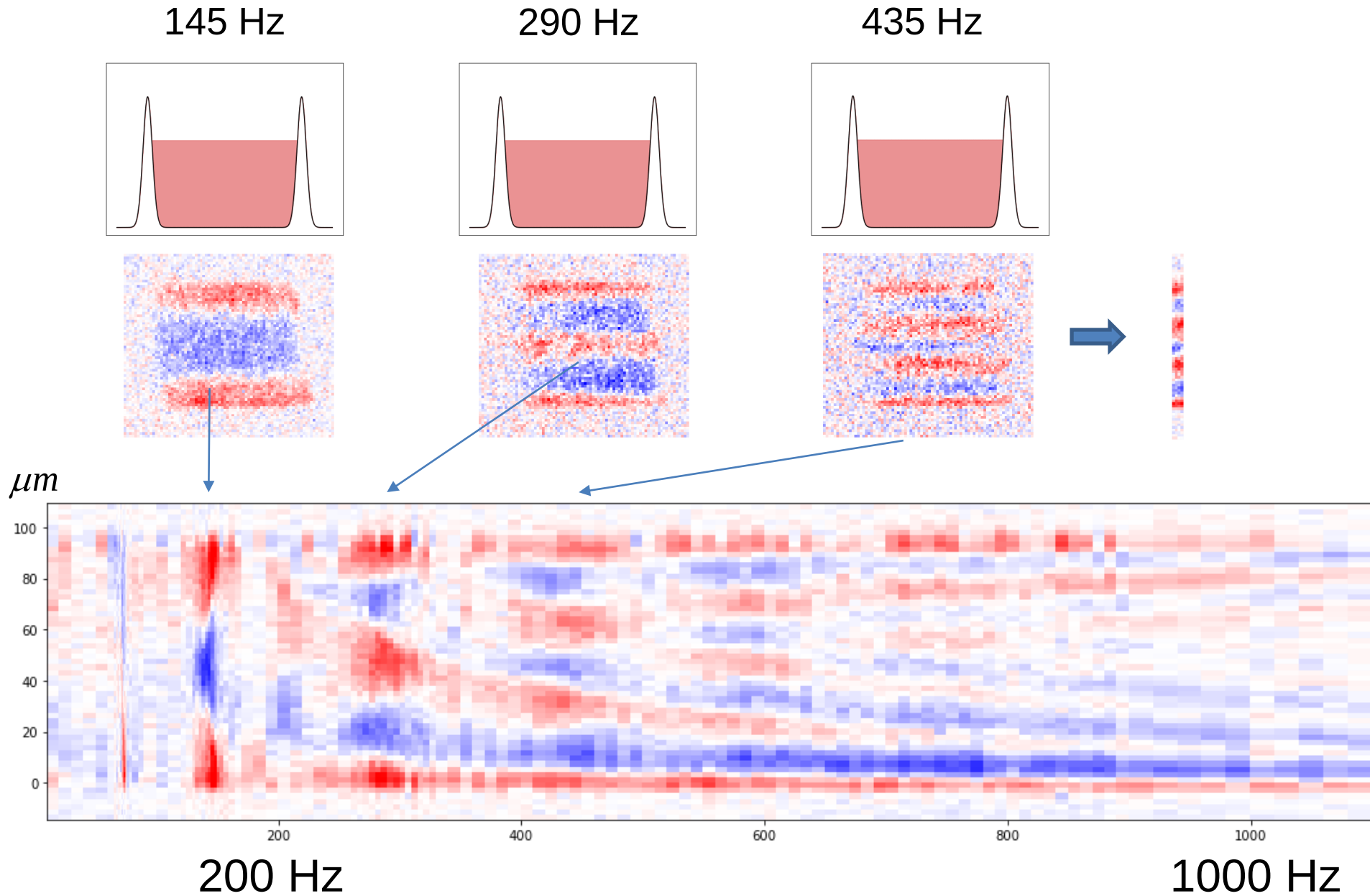
$$m c^2 = \left. \frac{\partial p}{\partial n} \right|_s = f(s) \left. \frac{\partial p_0}{\partial n} \right|_s = \frac{5}{3} \frac{p}{n} = \frac{10}{9} \frac{E}{N}$$



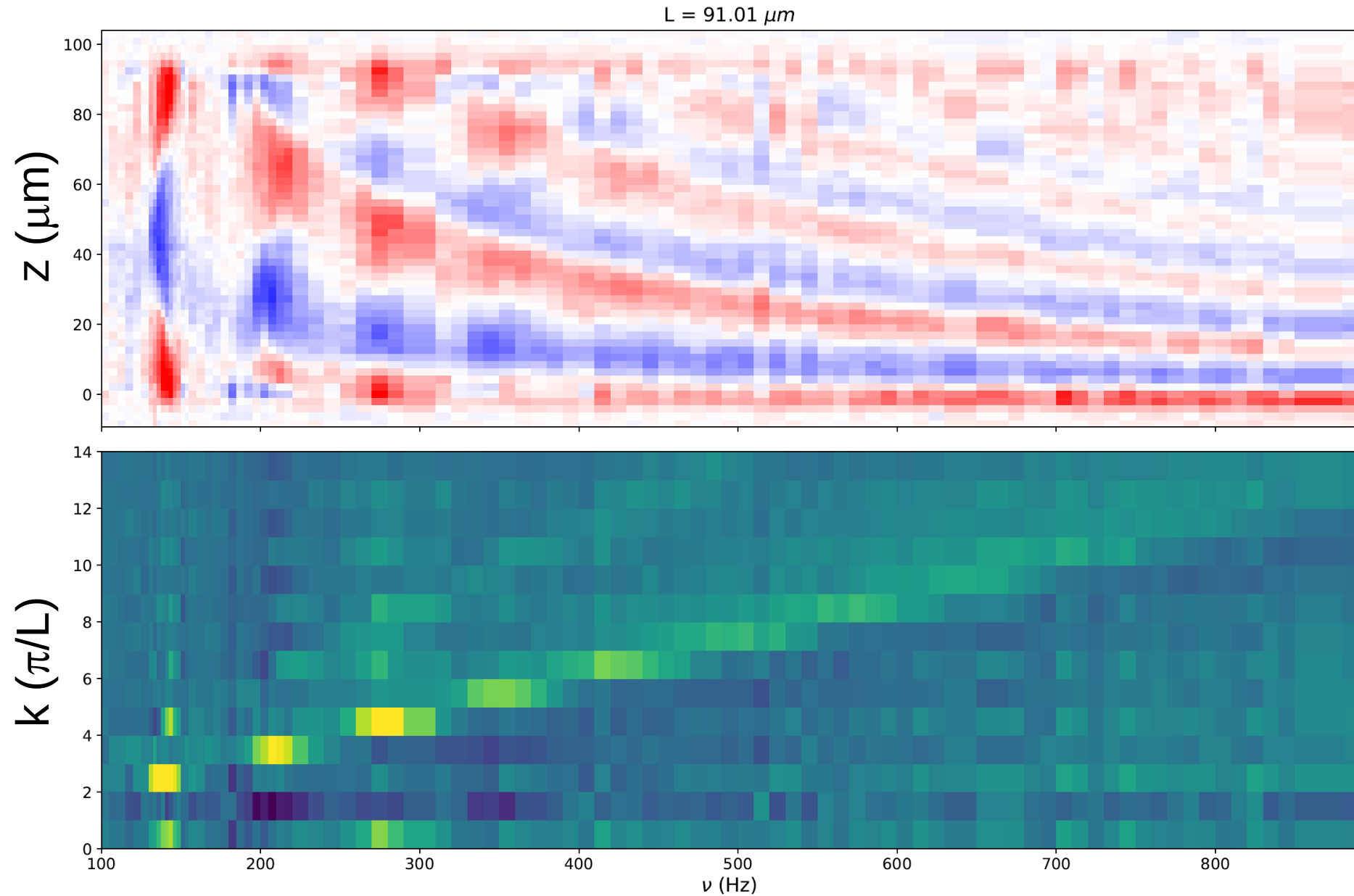
Resonant Modes



Resonant Modes



One-sided shaking: Even and odd



Unitary Hydrodynamics

Mass $\frac{\partial \rho}{\partial t} + \frac{\partial j}{\partial z} = 0$ Continuity equation

Momentum $\frac{\partial j}{\partial t} + \frac{\partial p}{\partial z} = \frac{4}{3} \eta \frac{\partial^2 v}{\partial^2 z}$ Navier-Stokes equation
viscosity damps fluid

No bulk viscosity (ζ) due to scale invariance

Heat $\frac{\partial s}{\partial t} + \frac{\partial (vs)}{\partial z} = \frac{\kappa}{T} \frac{\partial^2 T}{\partial^2 z}$ **thermal conduction**
increases entropy

Sound Diffusivity
*due to viscosity &
thermal conductivity*

$$D = \frac{4}{3} \frac{\eta}{\rho} + \frac{4}{15} \frac{\kappa T}{P}$$

Dispersion relation

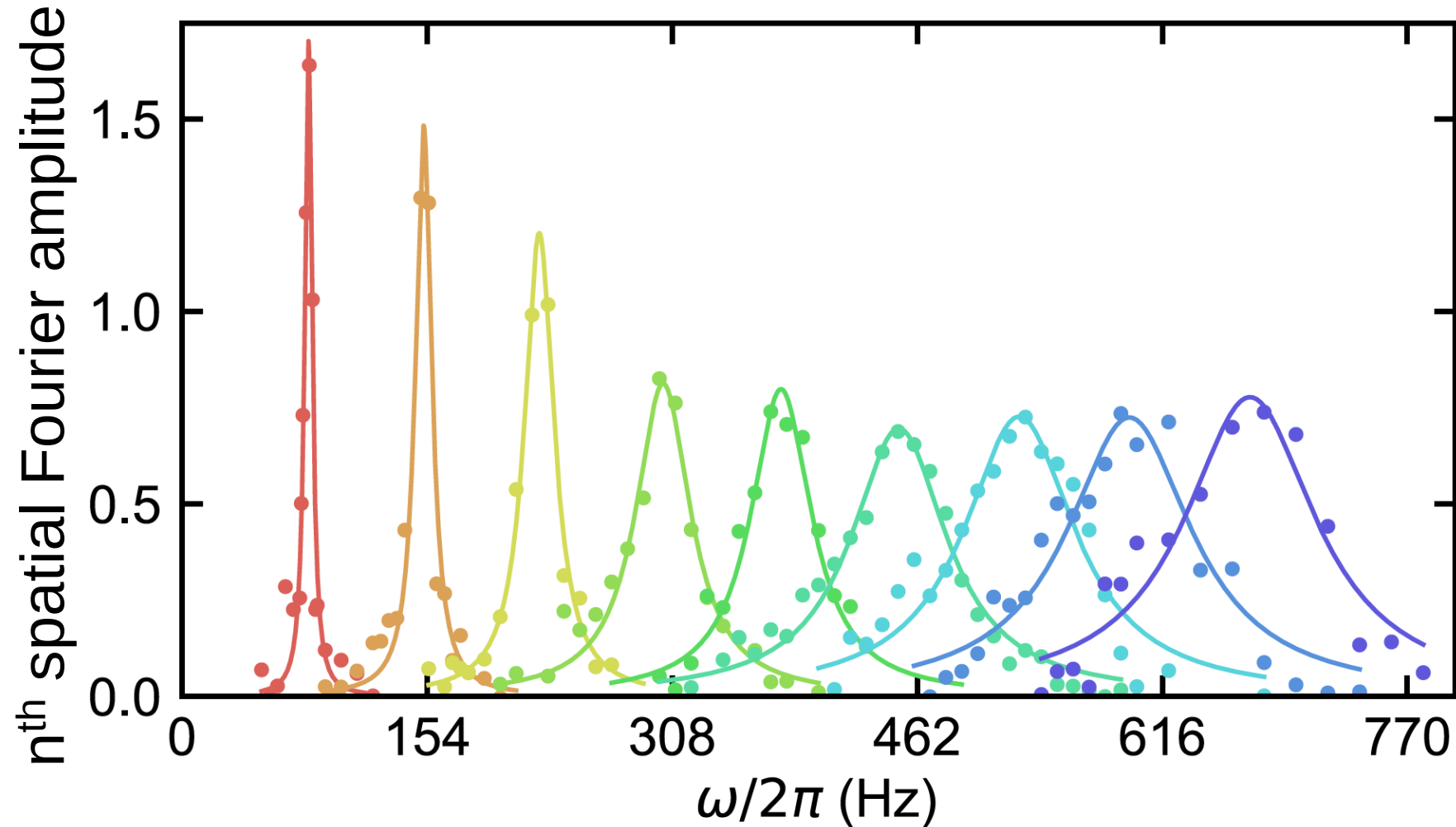
$$\omega^2 = c^2 k^2 + i \omega D k^2$$

Damping rate:

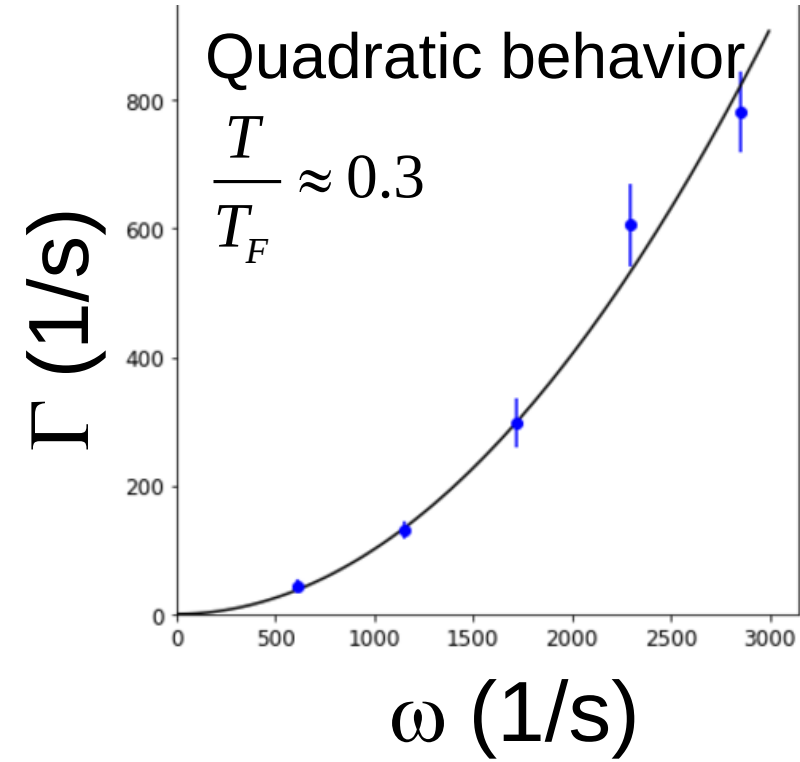
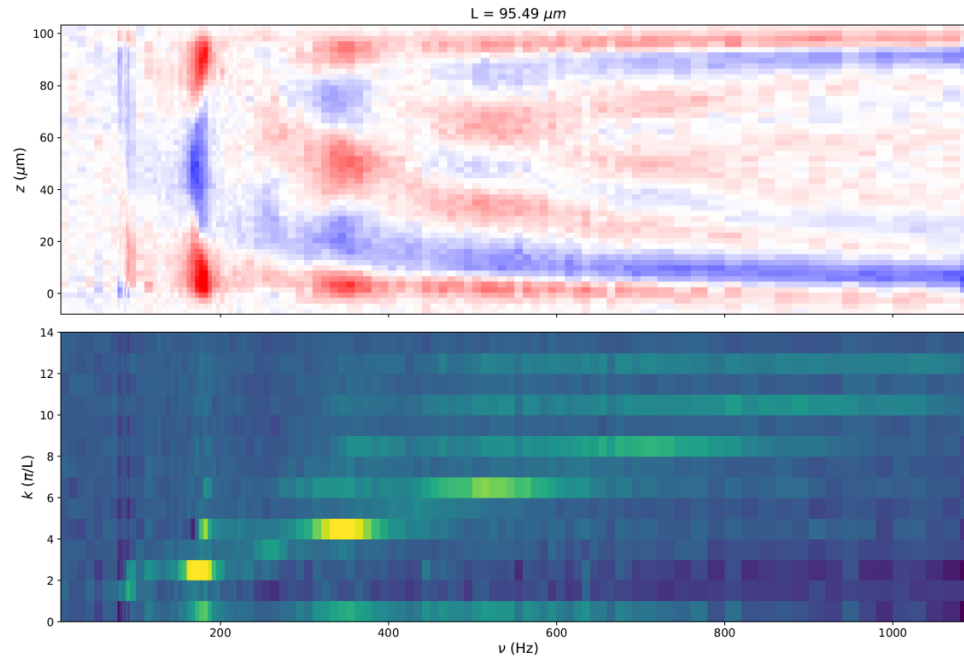
$$\Gamma = D k^2$$

Sound resonances

A direct measurement of the density response function



Sound Diffusivity from Sonogram Peaks



$$D = 2.1 \frac{\hbar}{m}$$

Quantum limited sound diffusion

Spin diffusion: Sommer et al., MIT 2011

Viscosity: Schaefer, Thomas (e.g. Science 2011)

Theory: Enss, Haussmann, Zwerger, Taylor, Randeria et al.

Viscosity of Superfluid Helium-4

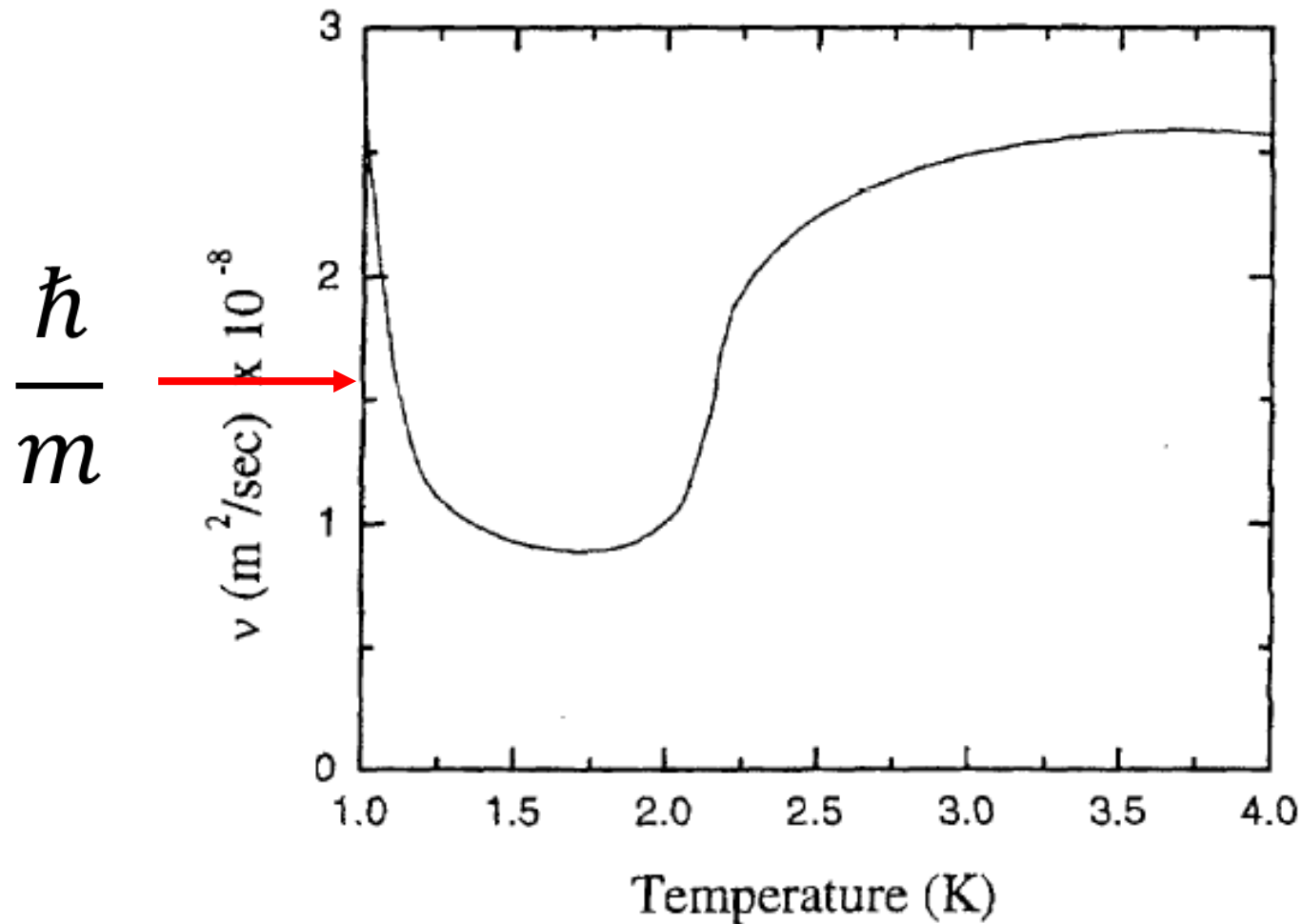
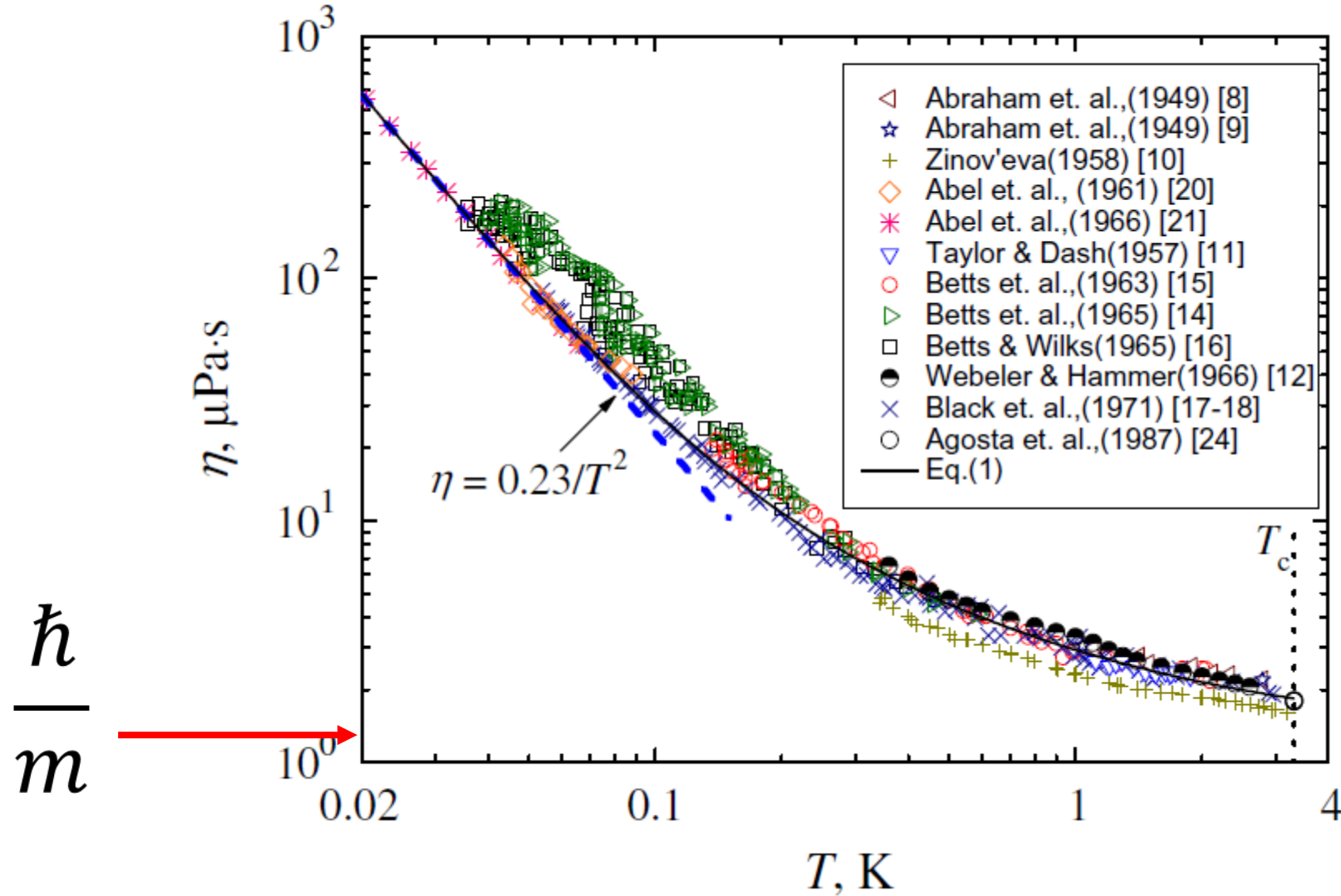


FIG. 12.3. The recommended values of the kinematic viscosity of liquid ^4He , $\nu = \eta/\rho$, as a function of temperature at the saturated vapor pressure.

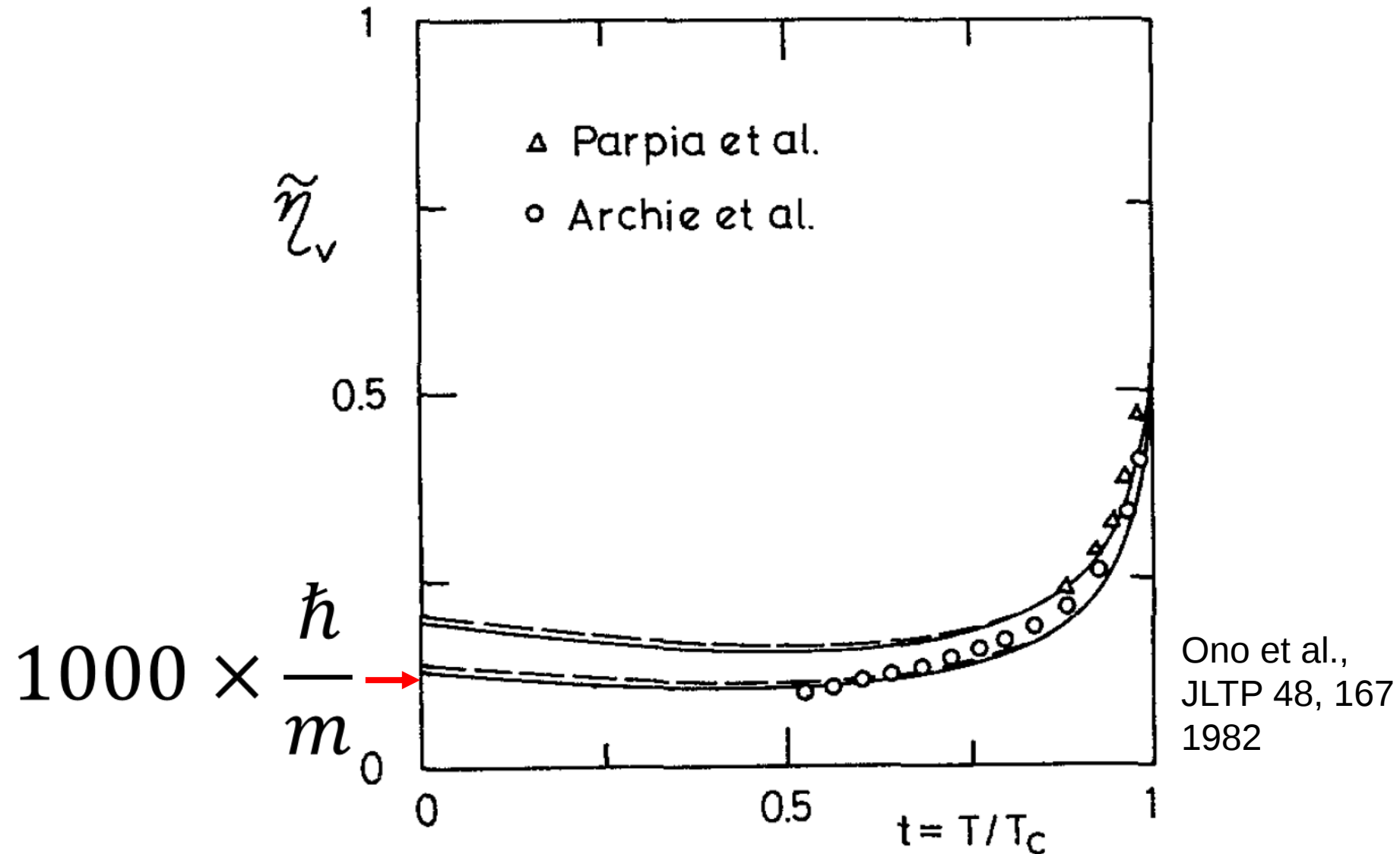
Viscosity of Superfluid Helium-3

Helium-3 above T_c : Fermi Liquid behavior

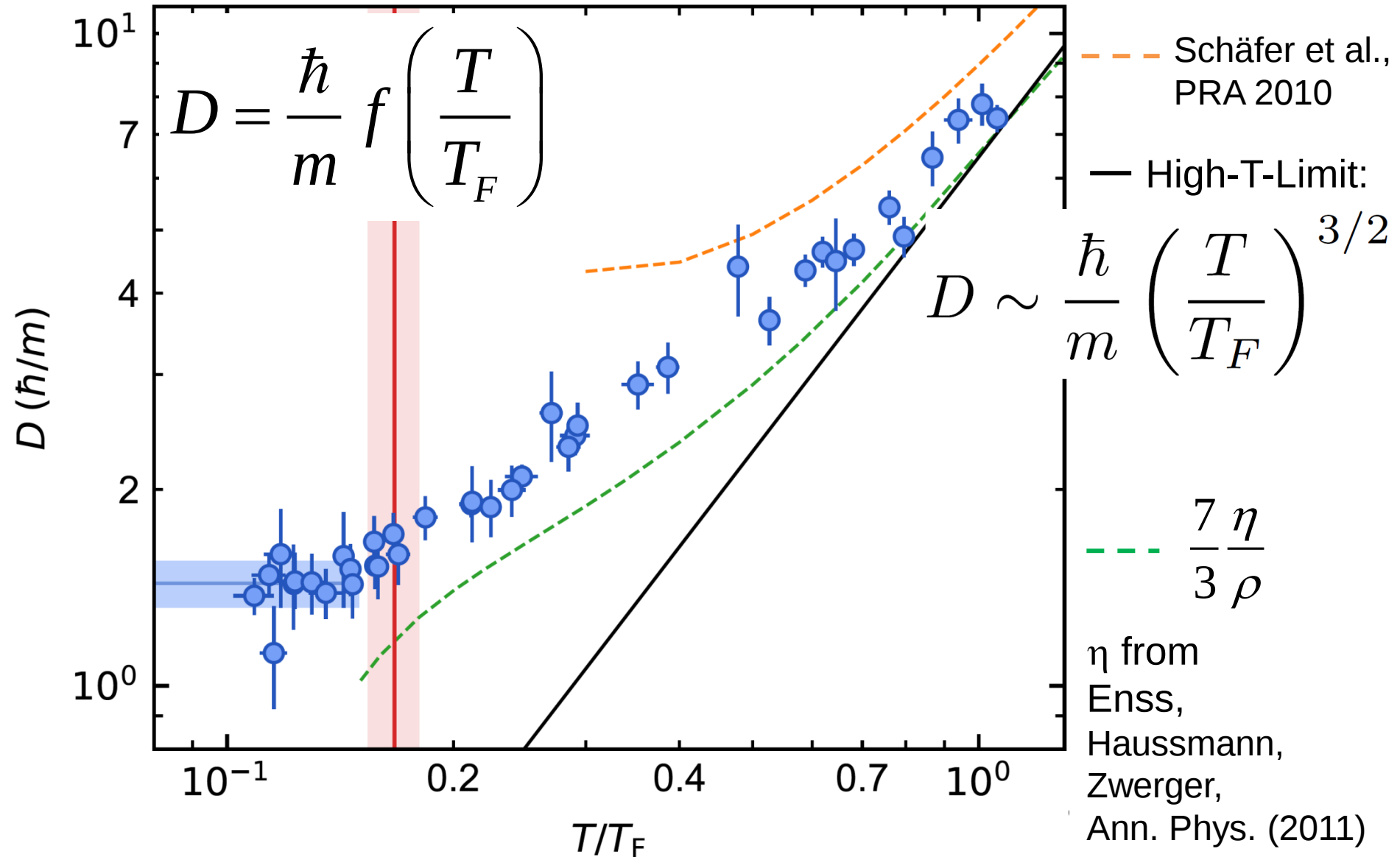


Viscosity of Superfluid Helium-3

Helium-3 below T_c : Sudden decrease, then constant



Quantum Limited Sound Diffusion



Comparison to state of the art theory

Quantum critical thermal transport in the unitary Fermi gas
Bernhard Frank, Wilhelm Zwerger, Tilman Enss
Phys. Rev. Research 2, 023301 (2020)

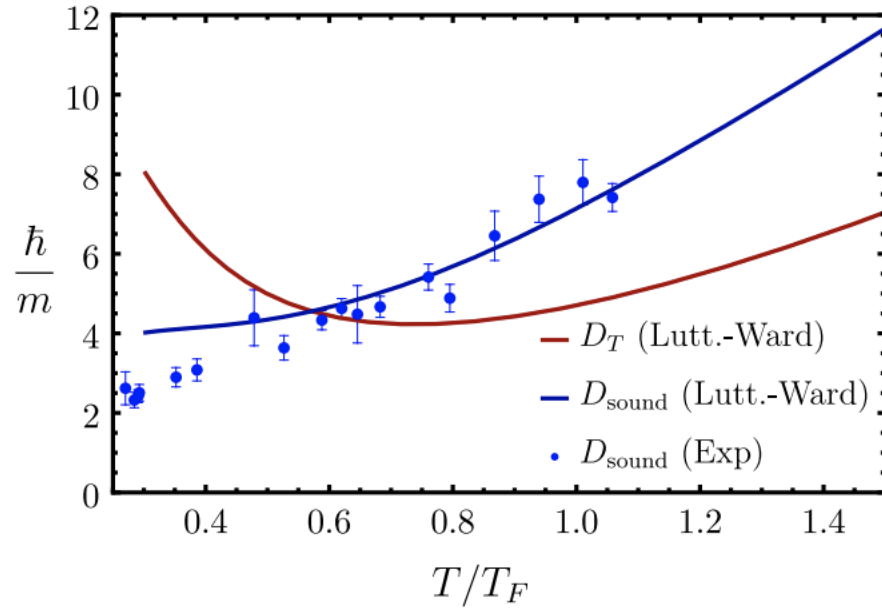


FIG. 1. Thermal diffusivity D_T (red) and sound diffusivity D_{sound} (blue) vs temperature T/T_F in the quantum critical regime of the unitary Fermi gas. Theoretical results from Luttinger-Ward calculations are shown in comparison with sound diffusion measurements [23].

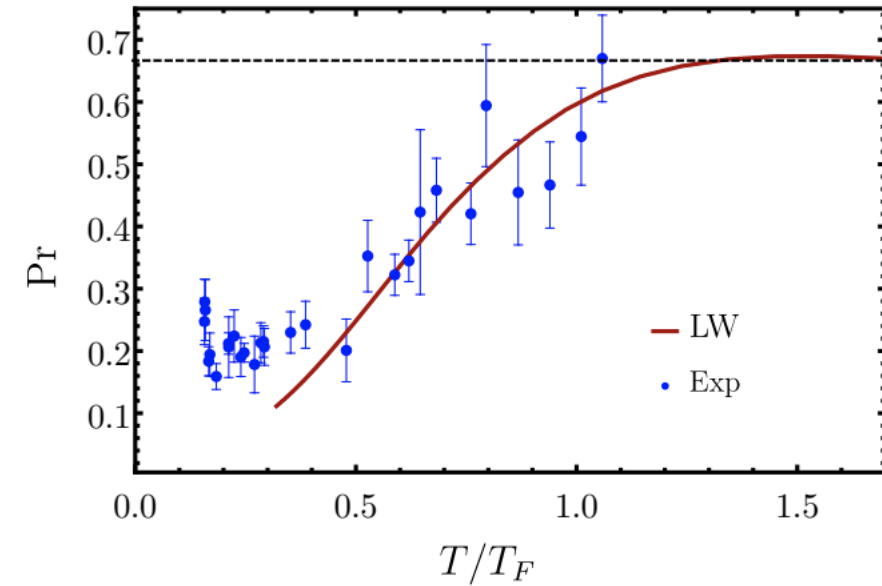


FIG. 8. Prandtl number $\text{Pr} = D_\eta/D_T$ vs temperature T/T_F from Luttinger-Ward calculations (red line) and sound attenuation measurements [23]; the dashed line marks the high-temperature limit $\text{Pr} = 2/3$.

Unitary Two-Fluid Hydrodynamics

For 1D flow, linearized:



Laszlo Tisza

$$\frac{\partial \rho}{\partial t} + \frac{\partial j}{\partial z} = 0$$

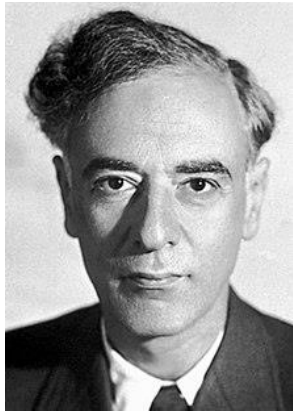
Continuity equation

$$\frac{\partial j}{\partial t} + \frac{\partial p}{\partial z} = \frac{4}{3} \eta \frac{\partial^2 v_n}{\partial^2 z}$$

Navier-Stokes eq.

viscosity

damps normal fluid



Lev Landau

$$\frac{\partial s}{\partial t} + \frac{\partial (v_n s)}{\partial z} = \frac{\kappa}{T} \frac{\partial^2 T}{\partial^2 z}$$

Entropy moving
with normal fluid,

thermal conduction

increases entropy

Superfluid flow:

$$\frac{\partial v_s}{\partial t} + \frac{\partial \mu/m}{\partial z} = \rho_s \zeta_3 \frac{\partial^2 (v_s - v_n)}{\partial^2 z}$$

bulk viscosity

normal-superfluid
interconversion

No other transport coefficients ($\zeta_1, \zeta_2, \zeta_4$)

+Thermodynamics relating p, ρ, T, s, μ

due to scale invariance (Son, Stringari&Pitaevskii)

First and Second Sound

Two fluids → Two sound modes

Review: Stringari & Pitaevskii
arXiv:1510.01306

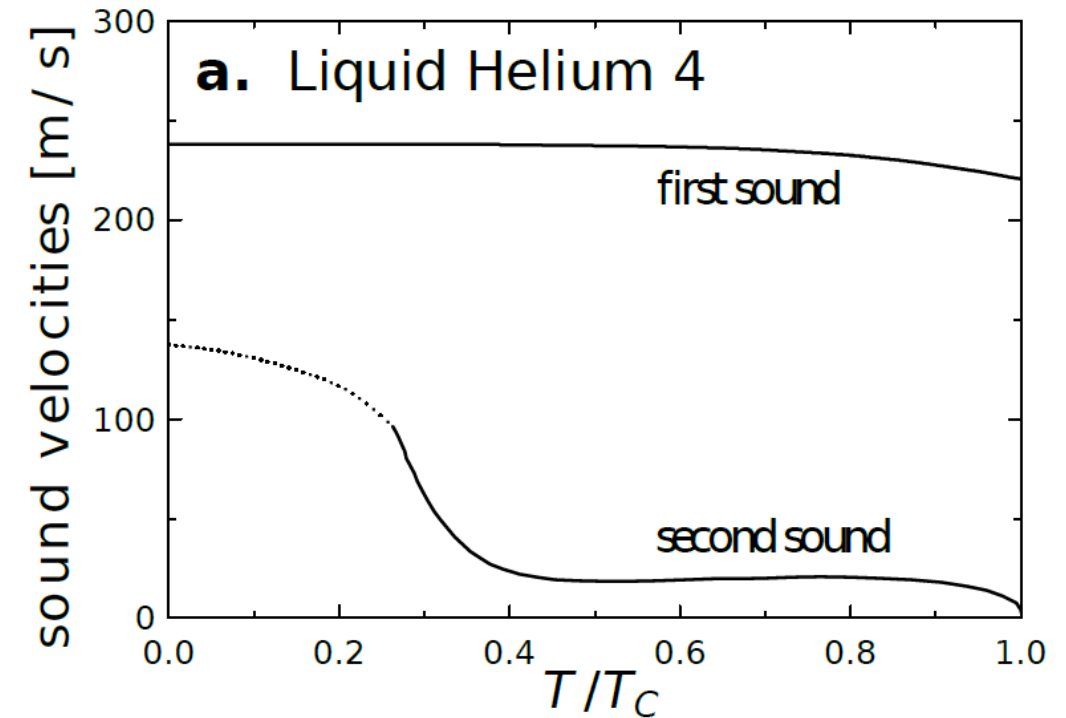
For small expansivity: Density and Heat

$$c_{10}^2 = \left. \frac{\partial p}{\partial \rho} \right|_{S,N} = \frac{1}{\rho \kappa_s}$$

First sound, a density wave
normal and superfluid
oscillate in phase

$$c_{20}^2 = \frac{\rho_s}{\rho_n} \left. \frac{\partial T}{\partial \sigma} \right|_\rho \frac{\sigma^2}{\gamma} = \frac{\rho_s}{\rho_n} \frac{\sigma^2 T}{c_p}$$

Second sound, a heat wave,
normal and superfluid oscillate out of phase

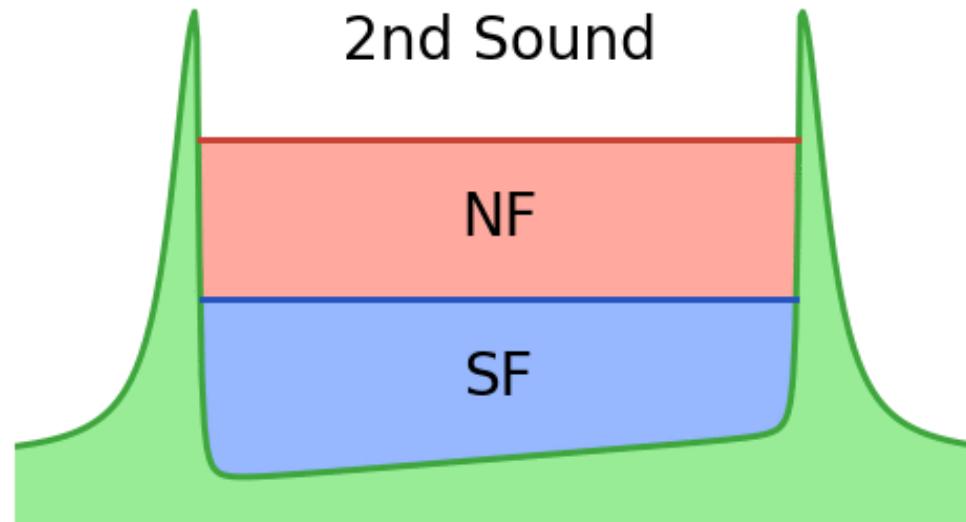


Second sound in a box

(Isentropic) expansivity at unitarity $-\frac{1}{\rho} \frac{\partial \rho}{\partial T} \Big|_s = -\frac{3}{2} \frac{1}{T}$ always non-zero

→ Can drive temperature wave using potential acting on density

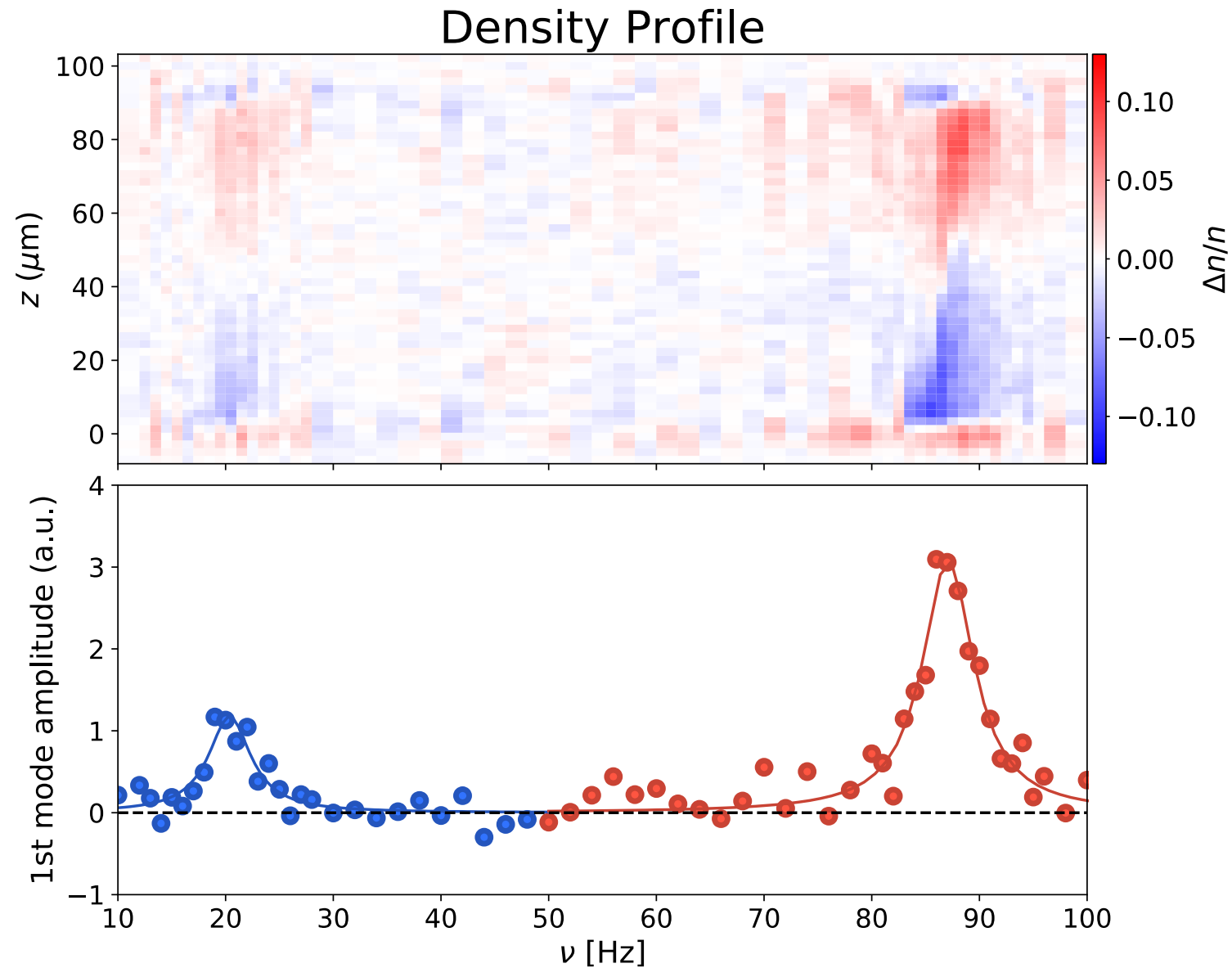
Apply oscillating gradient to the box



Observation of Second Sound in Quasi-1D geometry:
Sidorenkov et al., Grimm, Stringari & Piatevskii, Nature 2013

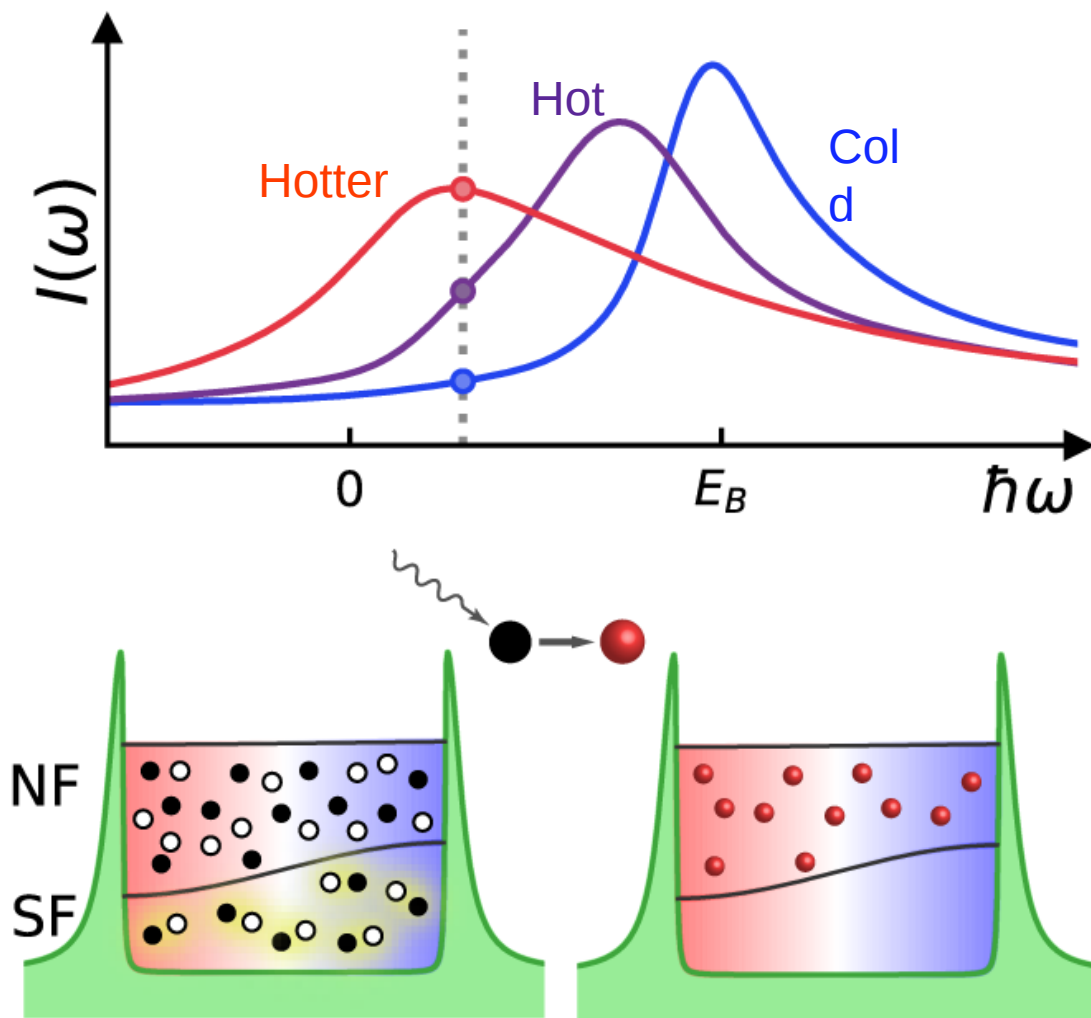
On bosons: weakly interacting situation: JILA (Debbie Jin 1996), MIT
Hydrodynamic gas: van der Straten

Second Sound seen in Density

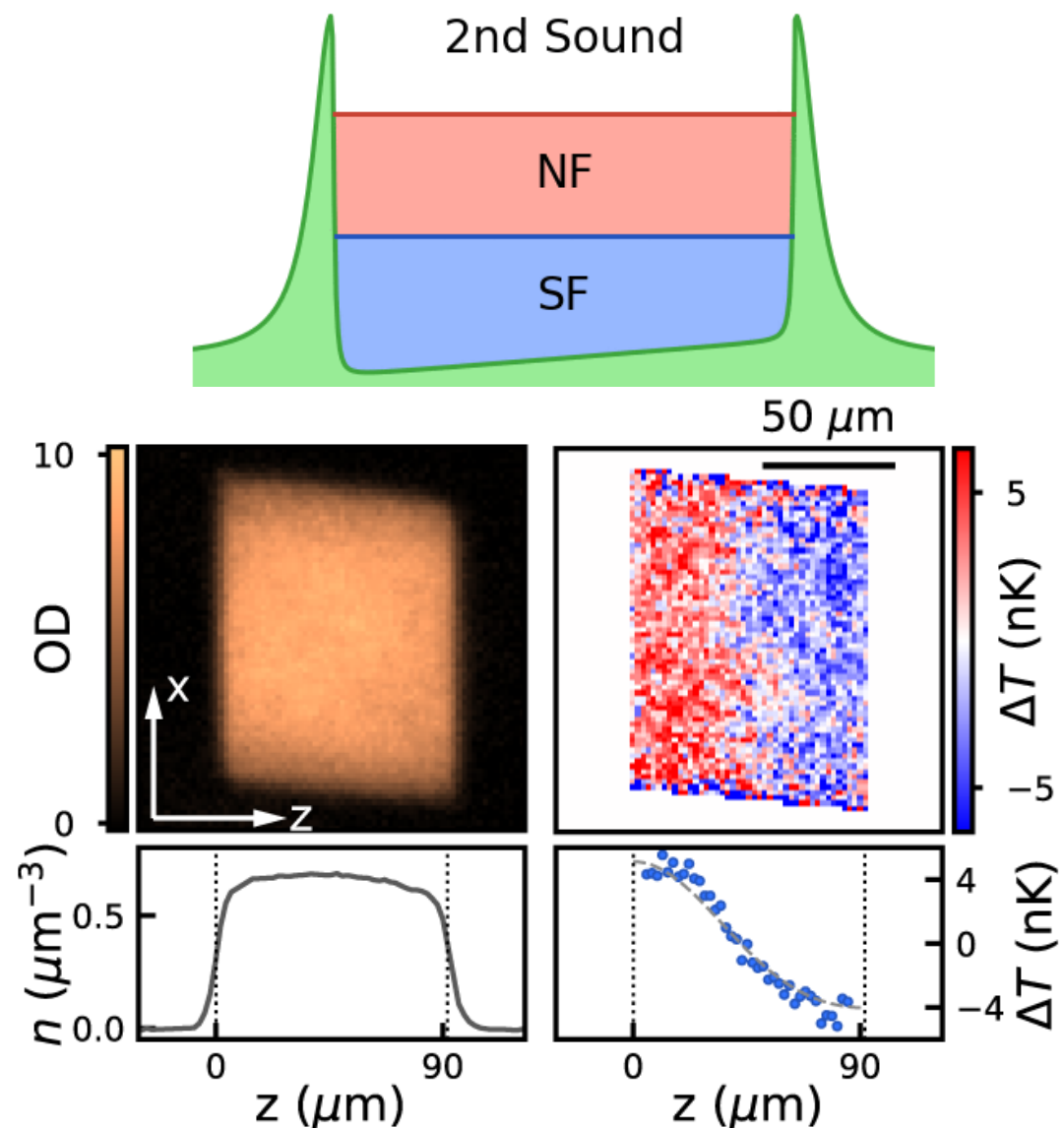


A local thermometer for heat transport

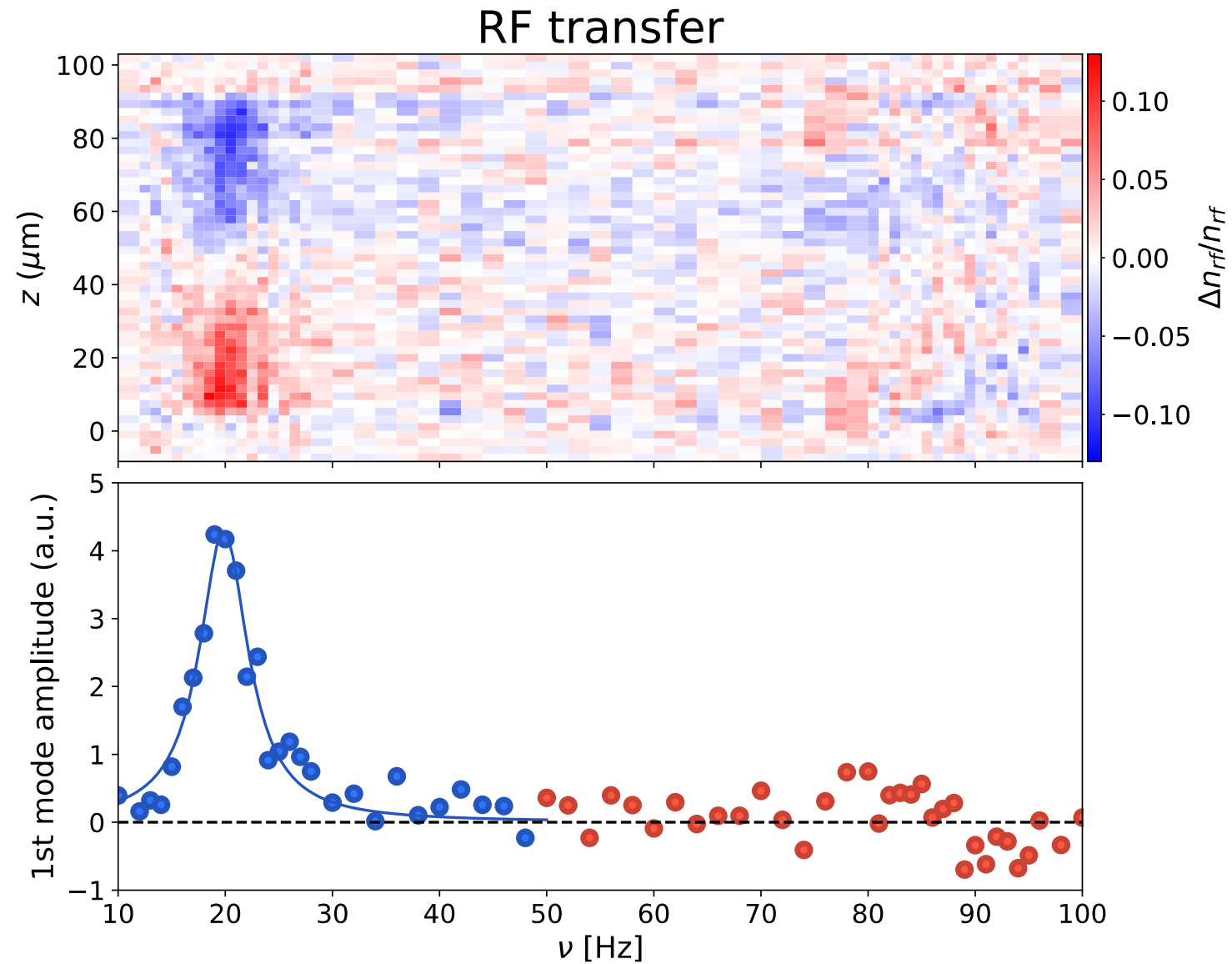
RF absorption spectra



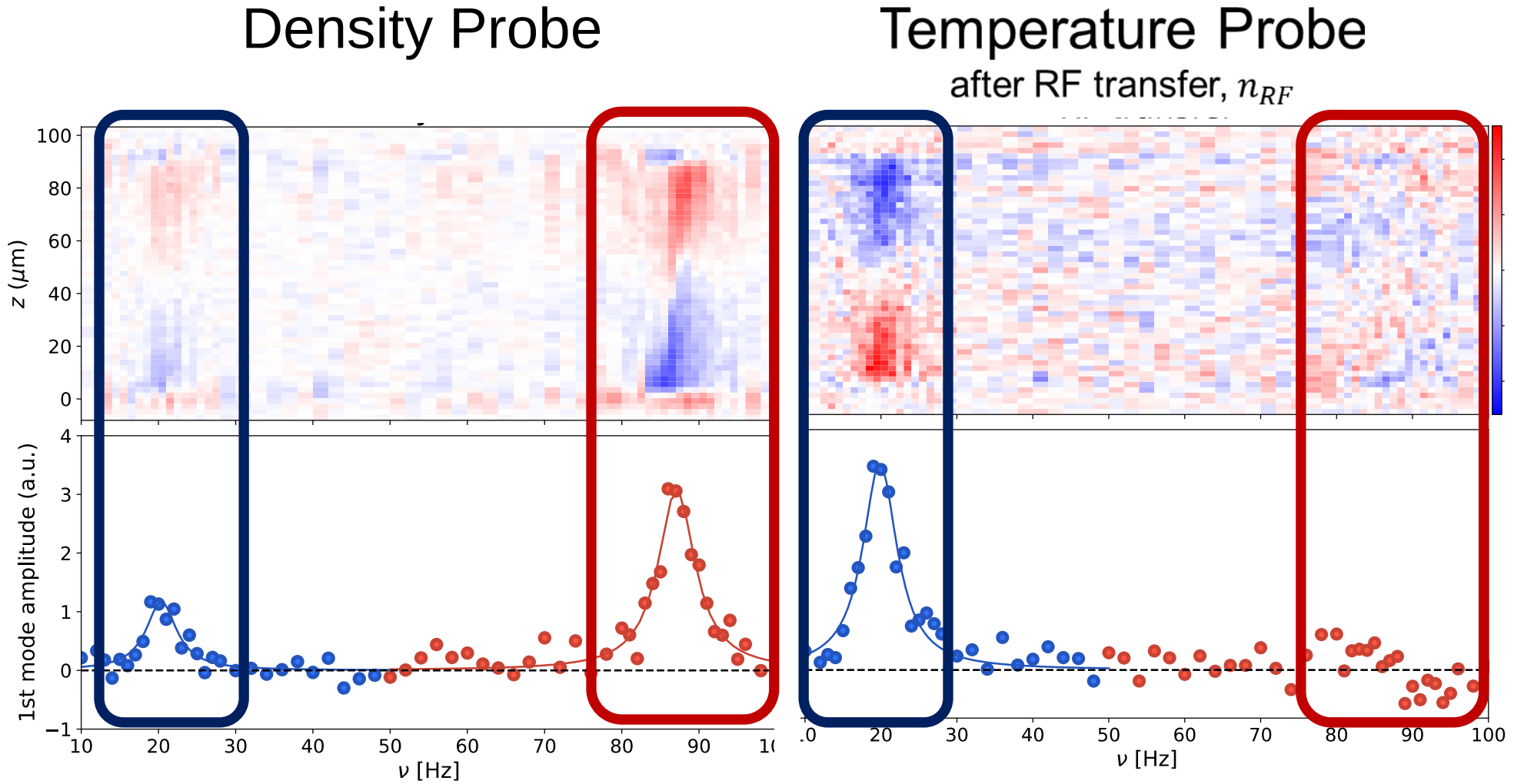
2nd Sound Creation



Second Sound with local thermometer

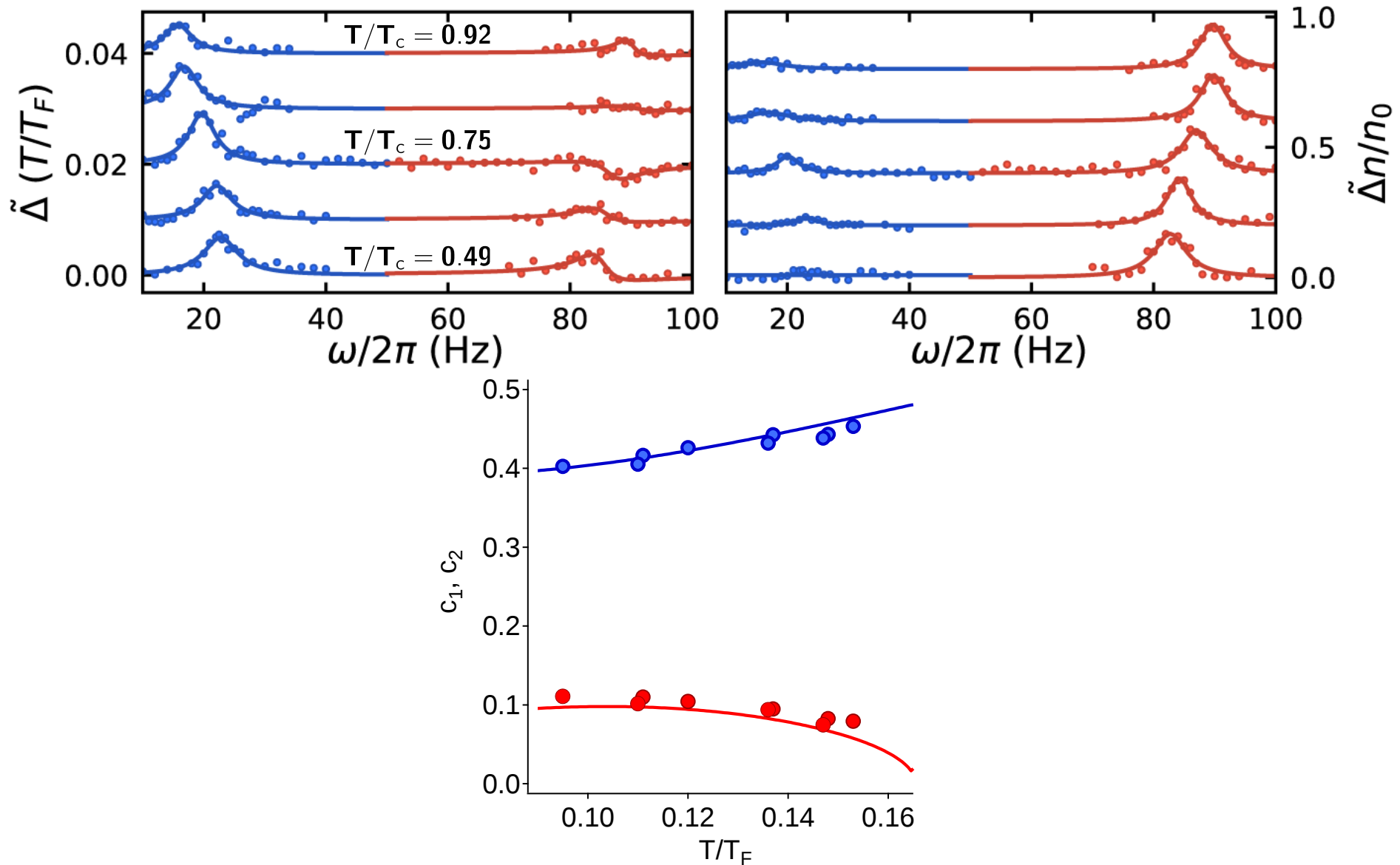


First and Second Sound



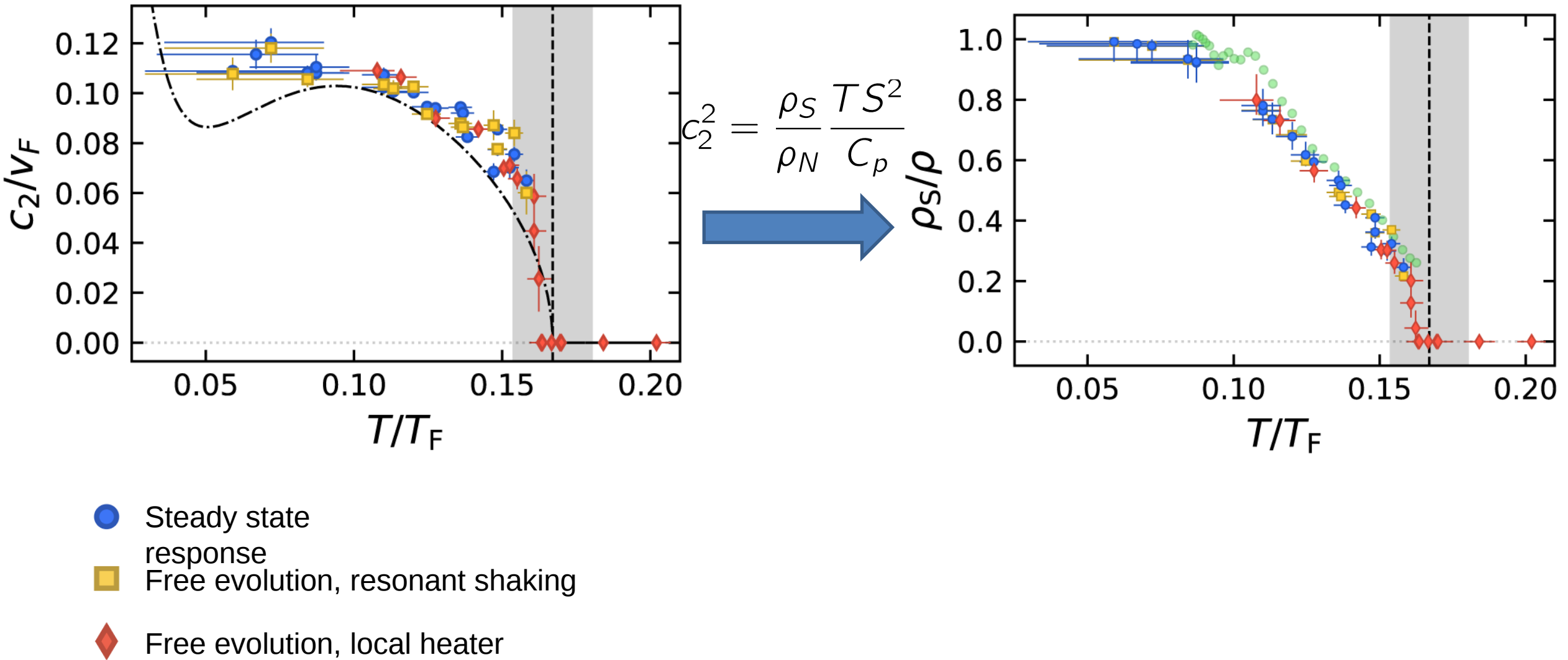
→ Speeds and Decay rates of First and Second sound

Speed of first and second sound



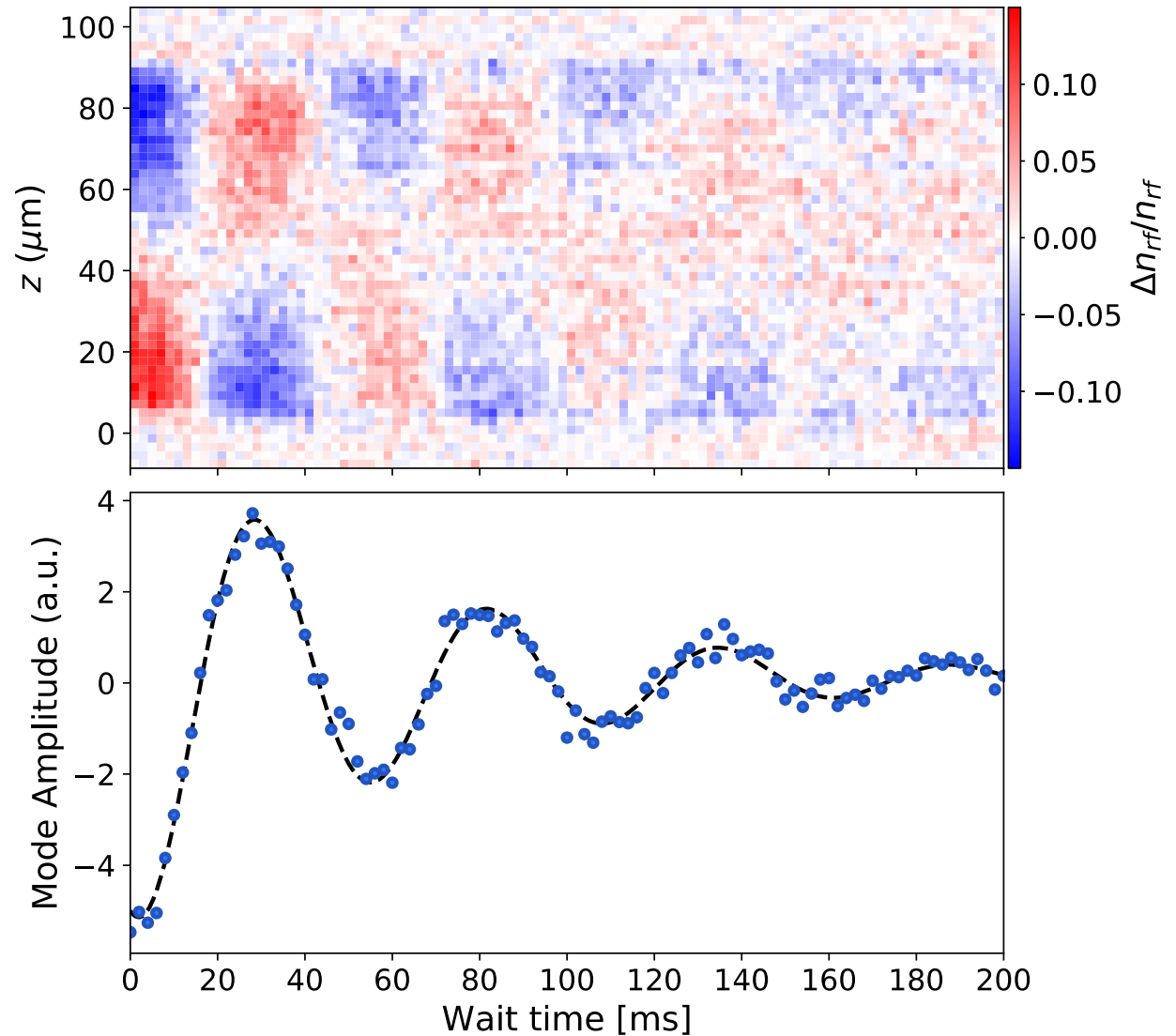
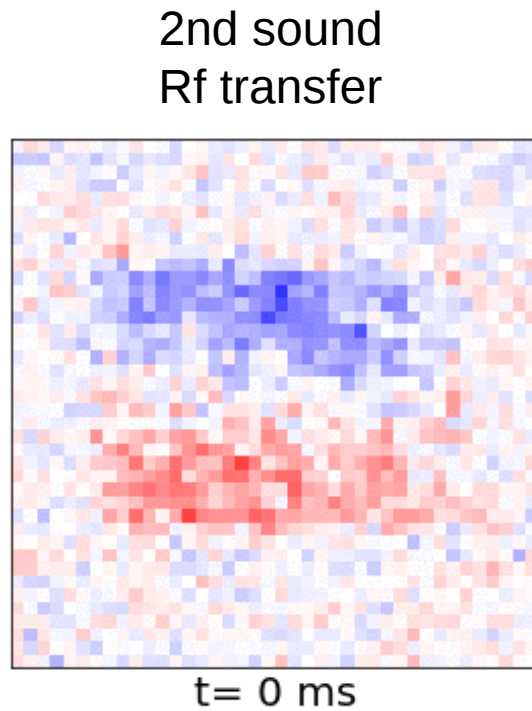
Z. Yan, P. Patel, B. Mukherjee, R. Fletcher, MZ, in prep. (2021)

Superfluid fraction

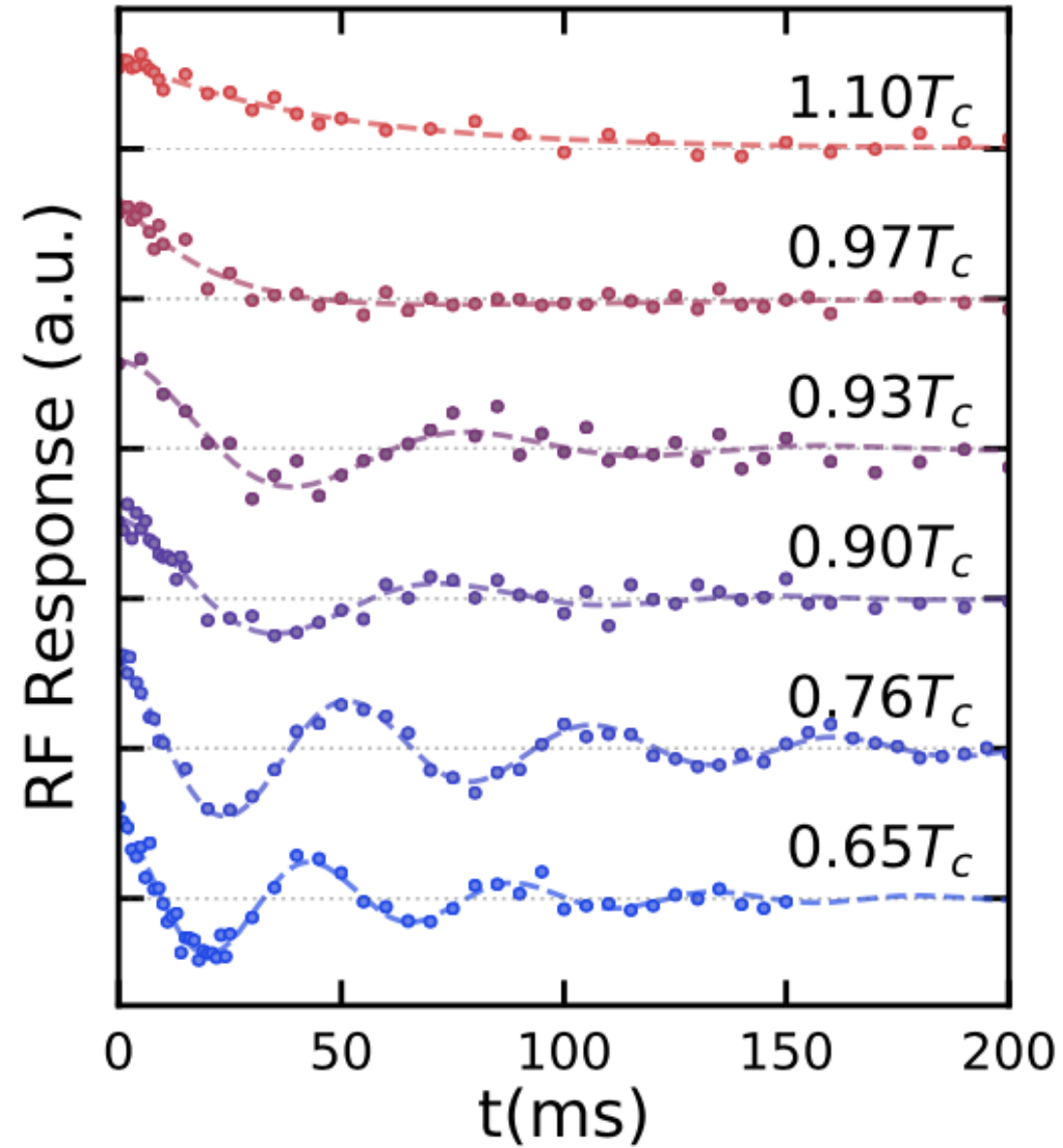


Damping of Second Sound

Temperature Probe

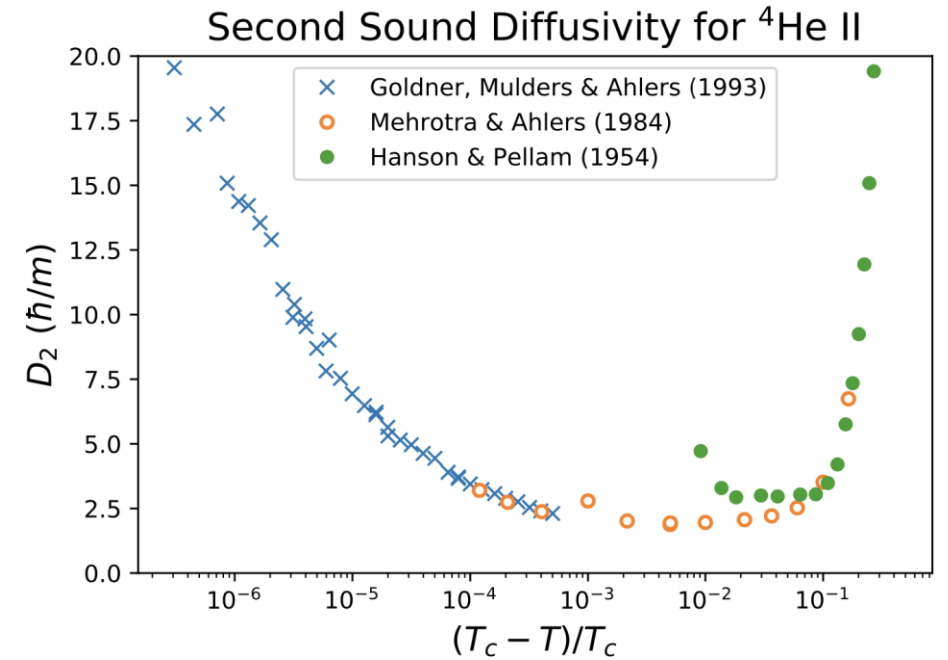
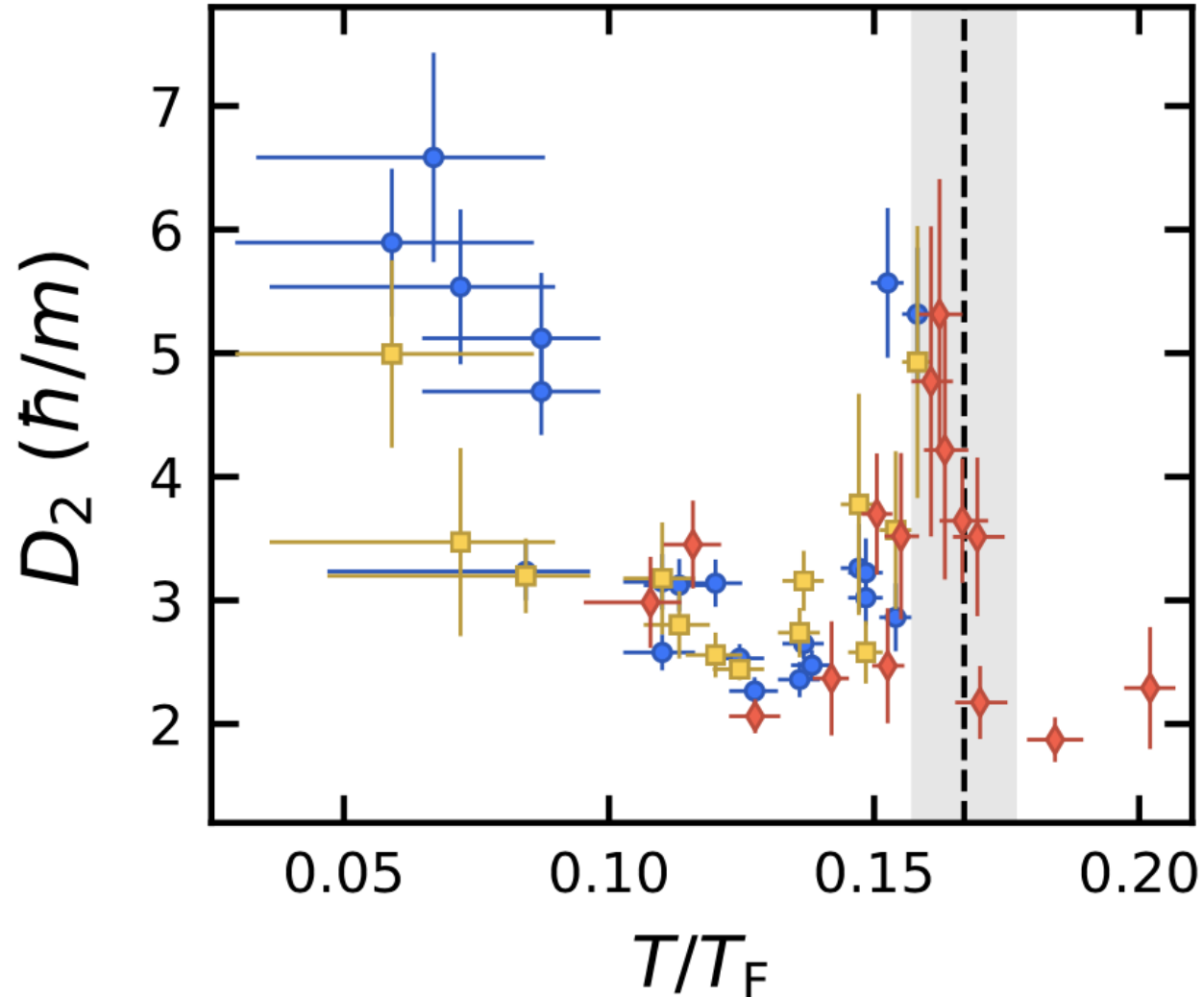


Transition from normal to superfluid



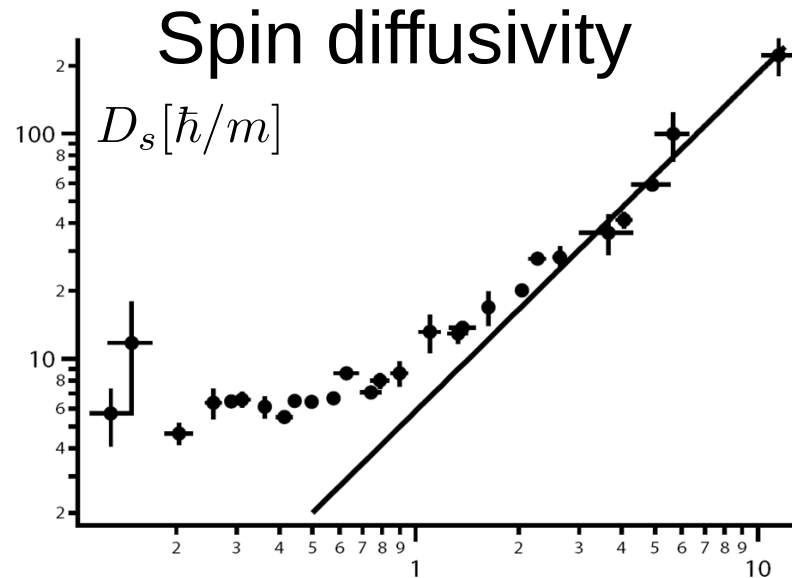
Diffusivity of Second Sound

$$D_2 = \Gamma_2 / k^2$$



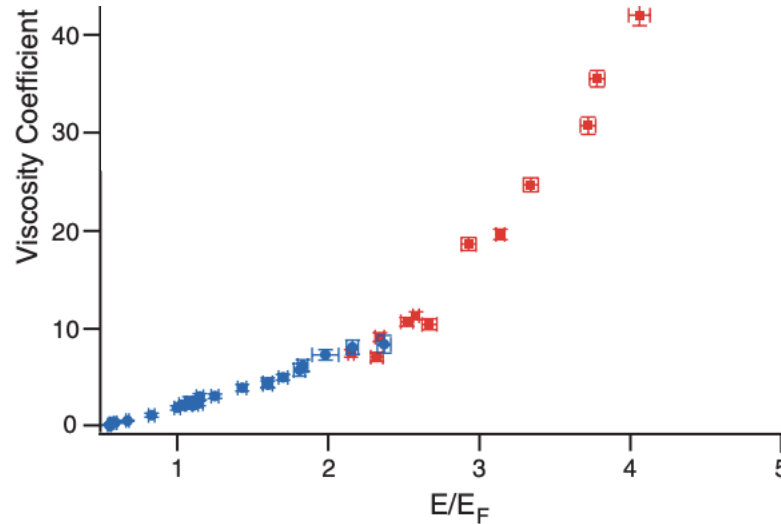
$$D_2 = \frac{\kappa}{c_p} + \frac{\rho_s}{\rho_n} \left(\xi_3 \rho + \frac{4}{3} \frac{\eta}{\rho} \right)$$

Diffusivities in the Unitary Gas

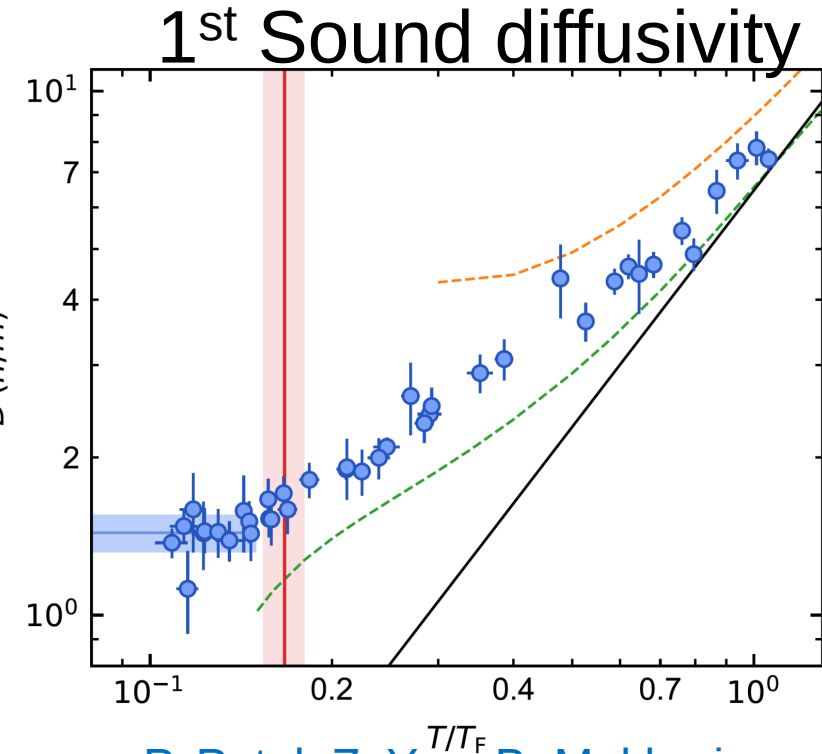


A.T. Sommer, M.J.H. Ku, G. Roati,
MZ, Nature 472, 201 (2011)

Kinematic viscosity (= momentum diffusivity)

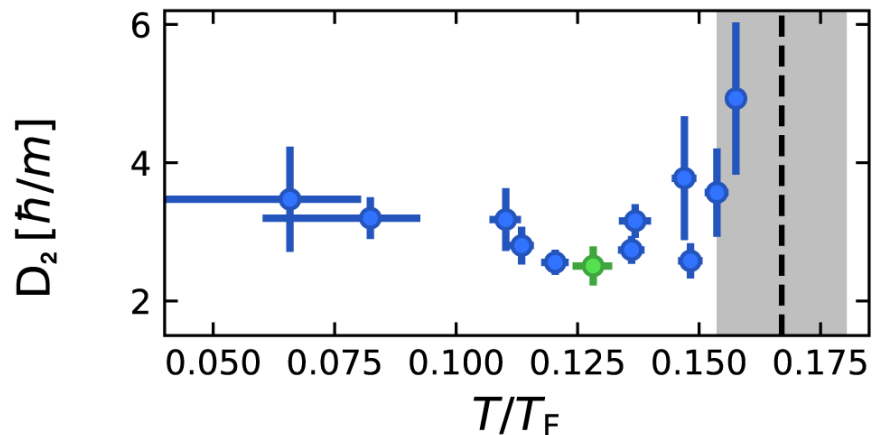


Cao et al., John Thomas group,
Science 331, 58 (2011)



P. Patel, Z. Yan, B. Mukherjee,
R. Fletcher, J. Struck, MZ,
Science 370, 1222 (2020)

2nd Sound diffusivity



Diffusivity (charge, spin, momentum, thermal)

$$D \propto vl \propto \frac{\hbar}{m}$$

Transport in Strongly Correlated Quantum Gases

Unifying themes of strongly interacting Fermi systems:

Loss of quasi-particle description: $\tau^{-1} \approx E_F / \hbar$

Quantum-Limited diffusivities:

$$D \approx \frac{\hbar}{m}$$

Quantum gas simulators are poised to further elucidate
Interplay between charge, spin, heat and Hall transport

- Universal Sound Diffusion in a Strongly Interacting Fermi Gas
Parth B. Patel, Zhenjie Yan, Biswaroop Mukherjee, Richard J. Fletcher, Julian Struck, MZ
Science, 370, 1222 (2020)
- Spin Transport in a Mott Insulator of Ultracold Fermions
Matthew Nichols, Lawrence Cheuk, Melih Okan, Thomas Hartke, Enrique Mendez, T.
Senthil, Ehsan Khatami, Hao Zhang, MZ, Science 363, 383 (2019)
- Doublon-Hole Correlations and Fluctuation Thermometry in a Fermi-Hubbard Gas
T. Hartke, B. Oreg, N. Jia, MZ, PRL 125, 113601 (2020)
- Bose polarons near quantum criticality
Zoe Z. Yan, Yiqi Ni, Carsten Robens, MZ, Science, 368, 190-194 (2020)
- Geometric squeezing into the lowest Landau level
Richard J. Fletcher, Airlia Shaffer, Cedric C. Wilson, Parth B. Patel, Zhenjie Yan, Valentin
Crépel, Biswaroop Mukherjee, MZ, arXiv:1911.12347 (Science 2021, to be published)

Bosons and Fermions



Bosons and Fermions

BEC 1/Fermi3

Fermionic Superfluids

Biswaroop Mukherjee
Parth Patel
Zhenjie Yan
Dr. Richard Fletcher
Visiting Professors:
Zoran Hadzibabic
Chris Vale
Former Members:
Ariel Sommer
Mark Ku
Dr. Tarik Yefsah
(→ CNRS, ENS Paris)
Dr. Julian Struck
(→ ENS Paris)

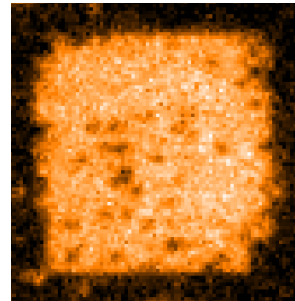
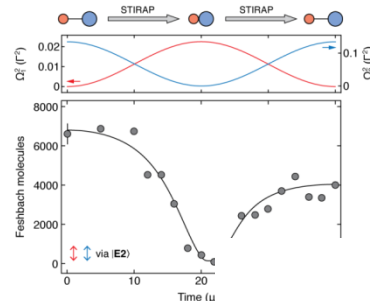


Fermi 1

NaK Dipolar Molecules

Yiqi Ni
Alex Yu Chuang
Eric Wolf
Dr. Carsten Robens

Former Members:
Zoe Yan
Jee Woo Park (→ SNU Korea)
Dr. Sebastian Will
(→ Prof., Columbia U.)
Cheng-Hsun Wu (→ Google X)
Dr. Huanqian Loh (→ Singap.)



Fermi 2

Fermi Gas Microscope

Thomas Hartke
Botond Oreg
Dr. Ningyuan Jia

Former Members:
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(→ Prof.@Princeton)
Melih Okan
Matthew Nichols
Dr. Hao Zhang
Katherine Lawrence
Vinay Ramasesh
Thomas Gersdorf
Dr. Thomas Lompe
(→ Hamburg)
Dr. Waseem Bakr
(→ Prof.@Princeton)

Fermi3

Rotating Superfluids

Airlia Shaffer-Moag
Cedric Wilson
Biswaroop Mukherjee
Parth Patel
Zhenjie Yan
Dr. Richard Fletcher
Valentin Crepel

