

EFT coefficient	Lower bound	Upper bound
$\tilde{g}_3$	-10.346	3
$\tilde{g}_4$	0	0.5
$\tilde{g}_5$	-4.096	2.5
$\tilde{g}_6$	0	0.25
$\tilde{g}'_6$	-12.83	3
$\tilde{g}_7$	-1.548	1.75
$\tilde{g}_8$	0	0.125
$\tilde{g}'_8$	-10.03	4
$\tilde{g}_9$	-0.524	1.125
$\tilde{g}'_9$	-13.60	3
$\tilde{g}_{10}$	0	0.0625
$\tilde{g}'_{10}$	-6.32	3.75

$$g_3 \geq 0$$

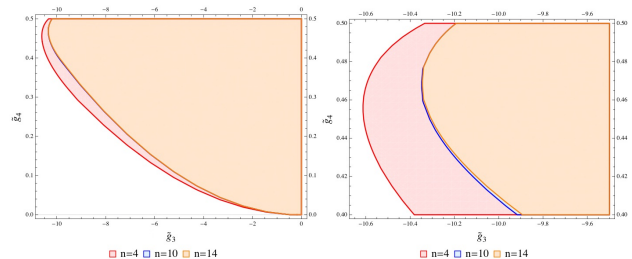
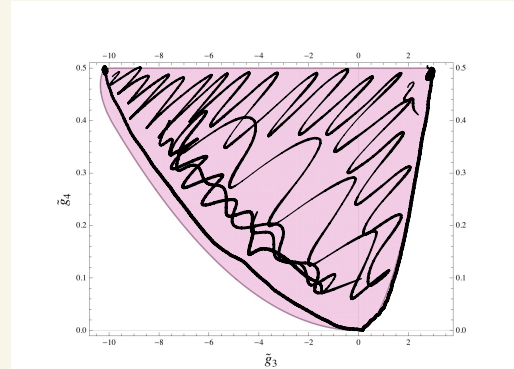
$$\tilde{g}_3 = g_3^2 / g_2$$

"Higher-derivative couplings don't grow faster than geometric"

$$\frac{1}{m^2 s} \rightarrow \frac{1}{m^2} + \frac{s}{m^4} + \frac{s^2}{m^6} + \dots$$

even if UV theory is strongly coupled, and has states of all spins.

Causality $\rightarrow$ EFT-like power counting.



(a) Convergence is rapidly achieved by adding null (b) Close-up near the left kink at $(-10.19, 0.5)$ constraints.

Figure 10: Convergence of the $(\tilde{g}_3, \tilde{g}_4)$ allowed region (shaded) with the Mandelstam degree n of crossing equations. The positive section of the horizontal axis is omitted since that region converges immediately.

Two dual approaches:

i) Rule out

ii) Rule in (exhibit some $2 \rightarrow 2$ S-matrix that satisfies our axioms)

Ruling 'in':

$$\mathcal{M}_{\text{spin-0}} = \frac{1}{m^2 - s} + \frac{1}{m^2 - t} + \frac{1}{m^2 - u},$$

$$\mathcal{M}_{\text{stu-pole}} = \frac{m^4}{(m^2 - s)(m^2 - t)(m^2 - u)} - \gamma(d) \mathcal{M}_{\text{spin-0}}$$

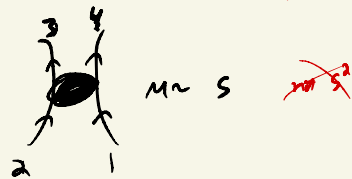
satisfies all our axioms!

cautions:

- High-dim ops in SM $\rightarrow (HL)^2 \Rightarrow \nu$ masses dim 5.

$\rightarrow \eta (H^\dagger d_u H)^2, (\bar{\psi} \psi)^2, \dots$ dim 6

these have spin 1, and are not ruled by spin ≥ 1 dispersion.



(even if we assumed convergence for $J=1$ sum rules,

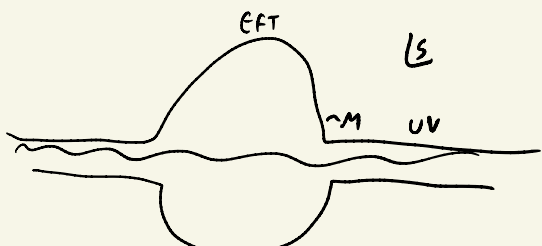
they are not "positive" $\oint \frac{ds}{s^2} u \Rightarrow g_1 = \langle \frac{|c^{\phi\phi}|^2 - |c^{\phi\bar{\phi}}|^2}{m^2} \rangle$

- unitarity $\Rightarrow 0 \leq \text{Im } a_\ell(s) \leq 2$

$$|1 + i a_\ell| \leq 1 \Leftrightarrow \begin{pmatrix} \text{Im } a_\ell & \text{Re } a_\ell \\ \text{Re } a_\ell & 2 - \text{Im } a_\ell \end{pmatrix} \succeq 0$$

Guerrini + Sever 2106.10257

Chiang + Huang + Li + Rodina + Weng 2105.02862, 2201.07177



Tolley + Wong + Zhou 2011.02400
SCH + van Duong 2011.02357

Scalar EFT $M_4 \supset g_2 (s^2 + t^2 + u^2) + g_3 stu + \dots$

$$g_2 = \langle \frac{1}{m^4} \rangle, \quad g_3 = \langle \frac{3 - 2i(i+1)}{m^6} \rangle, \quad 0 = \langle \frac{5i(i+1)(i(i+1)-8)}{m^8} \rangle$$

$$\langle \dots \rangle = \sum_{J=0, \text{ even}} \left(\sum_{m^2} \frac{d m^2}{2\pi} \frac{\text{Im } a_J(m)}{m^2} \right) \langle \dots \rangle$$

Q: given g_2 , what is $\min(g_3)$?

Anatomy of a bootstrap problem:

- positive unknowns
- some obs. to max/min
- crossing equation

$$(CFT) \quad F_0^2 \quad \gamma \left(\gamma \left(\gamma - \gamma \right) \right) = 0$$

$$\text{Ad-hoc Cauchy-Schwarz} \Rightarrow \langle \frac{5i(i+1)}{m^6} \rangle \leq \frac{8}{m^2} g_2 \Rightarrow \underline{M^2 g_3} \geq -16 \underline{g_2}$$

Systematic strategy:

take linear combination:

$$F[J, m] = \alpha \frac{1}{m^4} + \beta \frac{3 - 2i(i+1)}{m^6} + \gamma \frac{5i(i+1)(i(i+1)-8)}{m^8} M^4$$

s.t. $F[J, m] \geq 0$ for all possible heavy states ($J \geq 0$ even, $m \geq m$)

$$\langle F[J, m] \rangle \geq 0$$

$$\Rightarrow \alpha g_2 + g_3 + \beta \cdot 0 \geq 0$$

$$M^2 g_3 \geq -\alpha g_2 \Rightarrow \text{maximize } -\alpha \text{ over all such positive functionals.}$$

Good news: efficient algorithms (numerically) solve such problems!

i) simplest strategy: sample N points on a fine grid of (J, m)

$$N \downarrow \left(\begin{pmatrix} \vdots \\ \vdots \\ \vdots \end{pmatrix} \right) \cdot \begin{pmatrix} \alpha \\ 1 \\ g_3 \end{pmatrix} \succeq \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \end{pmatrix} \quad N \downarrow$$

Mathematica

LinearOptimization[] can solve this.

ii) impose $\forall$ mass above, for some finite list of spins.

thm: a poly(x) is positive for $x \in \mathbb{R}$ iff $\text{poly}(x) = \sum (p_i(x))^2$
 $x \geq 0$ iff $\text{poly}(x) = \sum (c_i)^2 + x \sum (c_i')^2$.

$$\frac{1}{m^2} \rightarrow \sum_{x=1}^{\infty} \quad m \geq 1 \Rightarrow x \geq 0$$

a poly of degree $2d$ is a sum of squares iff

$$\text{iff } \text{poly}(x) = (1 \ x \ \dots \ x^d) \cdot M \cdot \begin{pmatrix} 1 \\ x \\ \vdots \\ x^d \end{pmatrix} \quad \text{where } M \succeq 0$$

$\Rightarrow$ by finding positive matrices M , we solve our constraints ($\forall m \geq m$)

"Semi-definite programming" SDPB.

once a "numerically positive" functional is found,

positivity is easy to verify

EFTs of gravity

massless spin 2 particles $d=4$ $M \ll M_{Pl}$

Loantz, causality

$u \sim \langle 14 \rangle [23]^4 \left(\frac{8\pi 6}{5t^0} + 8\pi 6 \frac{g_3^2 s u}{t} + g_4 + \dots \right)$

$0 = \oint F(s, p^2 - s) ds$

$-B_2|_{low} = B_2|_{high}$

SCH + Mazac + Rastelli + Simmons-Ruffin
2102.08951

SCH + Li + Parra-Martinez + Simmons-Ruffin
2201.06602

contact terms don't contribute!

$\left(\frac{8\pi 6}{p^2} + 8\pi 6 g_3^2 p^6 = \sum_j \left(\frac{d_j^2}{m_j^6} \left[\frac{p^2}{2} \left(1 - \frac{p^2}{m_j^2} \right) + p^2 \right]^2 \right) \right)$

Task: Find a positive functional $\int_0^m dp \psi(p) \beta_0(p^2)$ s.t. RHS is pos, and LHS is $8\pi 6$ -const.

$\Rightarrow \psi^{motic}$ proves that gravity is attractive.

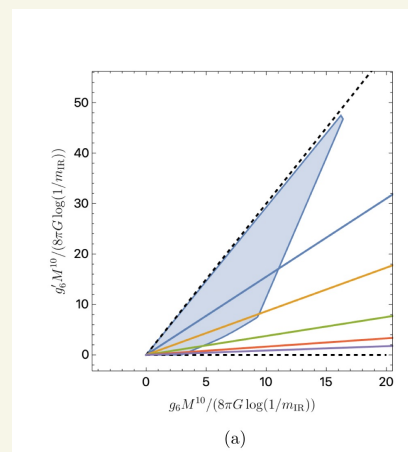
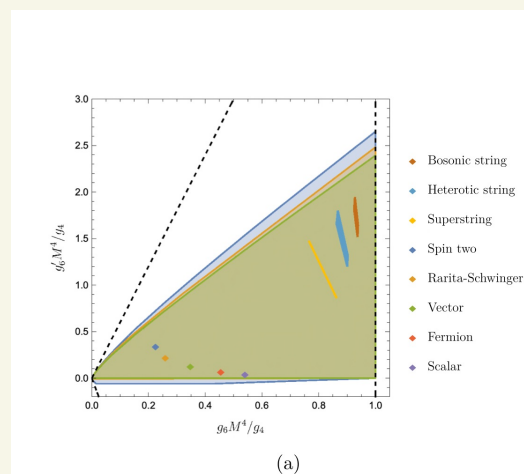
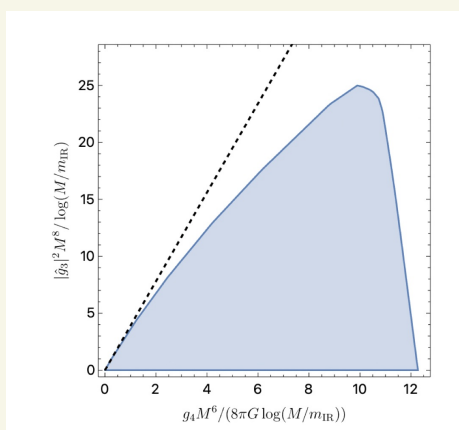
guarantees bounds: $6 - \#g_4 = \psi^{motic} - \psi_{g_4} \geq 0$

$\frac{g_4}{8\pi 6} \leq \# \frac{1}{M^6} \log(M/m_{IR})$

(in $d=4$: if positive at all $b \Rightarrow$ can't vanish at $p=0 \Rightarrow$ log-IR divergent on graviton pole)

interpretation:

$\mathcal{L} = \frac{1}{16\pi 6} \int d^4x \left(R + g_3 \frac{Riem^2}{2!} + \frac{g_4 Riem^4}{8\pi 6} + \dots \right)$



g_3 heavy

$M \ll M_{Pl}$

can't exceed gravitational strength

$\sqrt{8\pi 6} \cdot \alpha(l)$

if $g_3 \sim 10^4$ were measured, new higher-spin states must exist at $M \leq 10$

(Ex: $\sim (10 \text{ km})^4 \Rightarrow M \leq 10^{11} \text{ eV}$)

$f(R) = R + (R)^2 + V(R)$

$\uparrow$ unconstrained by us.

With gravity:

- how far from GR can it be? max(g_4) given mass M of higher-spin states
- how close to GR can it be? need Planck-scale effects?

100 SUSY

Guerrieri + Penedones + Vieira 2102.02847

Theory that minimizes $Riem^4$ is a string theory!