

Introduction to Random Matrix Theory

I] General Overview and applications of RMT

II] Ensembles of RMT

III] Coulomb gas approach

IV] Local statistics (bulk & edge)

I] Intro.

M : $N \times N$ matrix w. random entries M_{jk}
Real & Complex

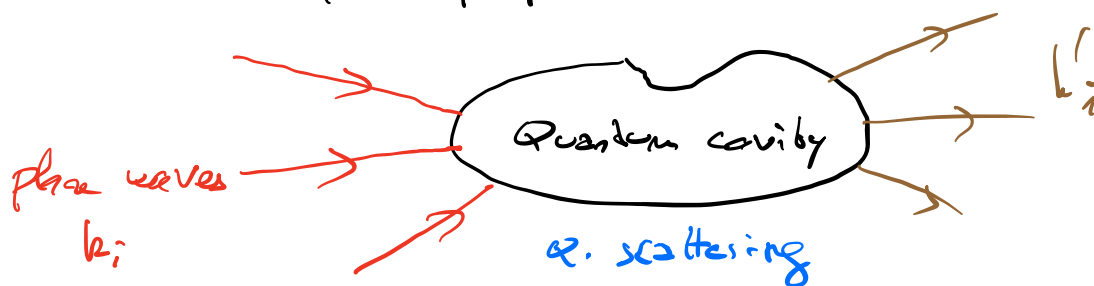
$\mathcal{P}$: stat. of eigenvalues and eigenvectors
 $N \rightarrow \infty$

In physics, introduced by Wigner ('50)

↳ study of nuclear physics to model
the energy levels of big nuclei.

Since then many applications:

* Mesoscopic physics:



$$\underline{k}' = S \underline{k}$$

random matrix

* Random graph : adjacency matrix

* nber theory (Riemann Zeta - Function)

* combinatorics

* stochastic / Fluctuating interfaces
Kardar-Parisi-Zhang equation

↳ related to directed polymers + disorder

In fact RMT was introduced by J. Wishart
(1928)

↳ correlation in time series

$$\vec{X}_t = \begin{pmatrix} x_{1,t} \\ x_{2,t} \\ \vdots \\ x_{n,t} \end{pmatrix}$$

stock price, daily temp.

$$\hat{C}_{jk} = \frac{1}{T} \sum_{t=1}^T x_{jt} x_{kt}$$



random correlated matrix

→ applications in data analysis.

2) Ensembles of RMT

→ Matrices w. real spectrum real symmetric
complex Hermitian

→ Two main categories of rand. matrices:

a) Ens. with independent entries

"Wigner matrices"

$$M_{jk} = x_{jk} + i y_{jk}$$

Joint probab. density function

$$\mathcal{P}(M) = \mathcal{P}(\{M_{jk}\})$$

$M_{ij} = M_{ji} \quad i \neq j$

$M \equiv \begin{pmatrix} \text{shaded upper triangle} \end{pmatrix}$
fixed by symmetry

$$\mathcal{P}(M) = \prod_{i=1}^N f_i(x_{ii}) \prod_{j < k} f_{jk}^{(R)}(x_{jk}) f_{jk}^{(I)}(y_{jk})$$

$$\text{Nber of degrees of freedom} = 1 + 2 + 3 + \dots + N \\ = \frac{N(N+1)}{2}$$

b) Rotationally invariant ensembles

M : real symmetric $\Rightarrow \exists (O, \Lambda)$

$$\Lambda = \begin{pmatrix} \lambda_1 & & 0 \\ & \ddots & \\ 0 & & \lambda_N \end{pmatrix}, \lambda_i = \text{eigenvalues of } M$$

$$w. \quad O O^T = \mathbb{1} \text{ such that } M = O \Lambda O^{-1}$$

Rotationally inv.: $\mathcal{S}(M) = \mathcal{S}(O M O^{-1})$

$\hookrightarrow$ the eigenvectors do not play a very important role.

$$\Leftrightarrow \mathcal{S}(M) \propto e^{-\text{Tr } V(M)}$$

$$P(\lambda_1, \dots, \lambda_N) \propto \prod_{i < j} |\lambda_i - \lambda_j|^\beta e^{-\sum_i V(\lambda_i)}$$

$\uparrow$
 some function

Porter-Rosenzweig thm ('60): the only ensembles that belong to a) and b)

are the Gaussian ens. :

$$\mathcal{P}(M) = \frac{1}{Z_N} e^{-a \text{Tr}(M^2) - b \text{Tr}(M)}$$

. show that this is type a) [Wigner]

$$\begin{aligned} \mathcal{P}(M) &= \frac{1}{Z_N} e^{-a \sum_{j,h} M_{jh} \underbrace{M_{hj}}_{=M_{jh}, j \neq h} - b \sum_j M_{jj}} \\ &= \frac{1}{Z_N} e^{-a \sum_j M_{jj}^2 - a \sum_{j \neq h} M_{jh}^2 - b \sum_j M_{jj}} \\ &= \frac{1}{Z_N} e^{-a \sum_j M_{jj}^2 - b \sum_j M_{jj} - 2a \sum_{j < h} M_{jh}^2} \\ &= \frac{1}{Z_N} \prod_j e^{-a M_{jj}^2 - b M_{jj}} \prod_{i < j} e^{-2a M_{ij}^2} \end{aligned}$$

↳ belongs to a)

NB: for $b=0$, the variance of the diag.

terms is twice the one of the off-diag.
terms.

In the following: set $b = 0$

For N real sym.: Gaussian Orthogonal Ens. ^{GOE}

N complex Hermitian: Gaussian Unitary Ens. ^{GUE}

c) Joint law of eigenvalues for GOE/GUE

GOE: $M = M^t$ and real.

$$M = \underbrace{O}_{\text{eigenvectors}} \underbrace{\Lambda}_{\text{eigenvalues}} \underbrace{O^{-1}}_{\text{eigenvectors}}, \quad \Lambda = \begin{pmatrix} \lambda_1 & & 0 \\ & \ddots & \\ 0 & & \lambda_N \end{pmatrix}$$
$$OO^t = \mathbb{1}$$

Change of variables: $M \equiv \{M_{jk}\} \rightarrow \{\Lambda, O\}$

Rh: nber of degrees of freedom:

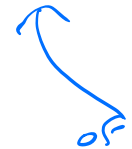
$$\frac{N(N+1)}{2} = N + \frac{N(N-1)}{2}$$

$$\mathcal{P}(M) \longrightarrow P(\{\lambda_1, \dots, \lambda_N\}, O)$$
$$e^{-\alpha \text{Tr}(M^2)}$$

Related by a Jacobian:

$$P(\{\lambda_1, \dots, \lambda_N\}, \underline{0}) = \mathcal{Z}(N) |\det J|$$

$$J = \left\{ \frac{\partial \pi_{i0}}{\partial \lambda_k}, \frac{\partial \pi_{i0}}{\partial \omega_{kl}} \right\}$$


 size $\frac{N(N+1)}{2} \times \frac{N(N+1)}{2}$

$$(GOE) \quad |\det J| = \prod_{j < k} (\lambda_j - \lambda_k)$$

indep. of $\underline{0}$

$$\Rightarrow P(\{\lambda_1, \dots, \lambda_N\}, \underline{0}) = \frac{1}{Z_N} e^{-\alpha \sum_{i=1}^N \lambda_i^2} \underbrace{\prod_{j < k} (\lambda_j - \lambda_k)^\beta}_{\text{level repulsion}}$$

for GOE $\beta = 1$

GOE $\beta = 2$

• Vandermonde determinant: $x_1, x_2, \dots, x_N$

$$V_N = \begin{pmatrix} 1 & 1 & \dots & 1 \\ x_1 & x_2 & & x_N \\ x_1^2 & x_2^2 & & x_N^2 \\ \vdots & \vdots & & \vdots \\ x_1^{N-1} & x_2^{N-1} & & x_N^{N-1} \end{pmatrix}$$

$$\underline{N=2}: \begin{pmatrix} 1 & 1 \\ x_1 & x_2 \end{pmatrix} = V_2 \Rightarrow \det V_2 = x_2 - x_1.$$

$$\det V_N = \prod_{i < j} (x_j - x_i)$$

Here we choose:

$$P(\lambda_1, \dots, \lambda_N) = B_N e^{-\beta \frac{N}{2} \sum_{i=1}^N \lambda_i^2} \prod_{i < j} |\lambda_i - \lambda_j|^\beta$$

III. 1 Coulomb gas approach (Dyson '62)

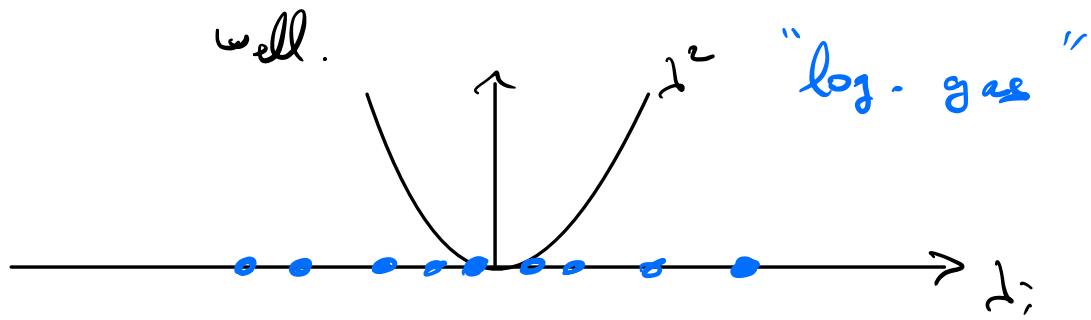
$$\text{Rewriting: } \prod_{i < j} |\lambda_i - \lambda_j|^\beta = e^{\beta \sum_{i < j} \ln |\lambda_i - \lambda_j|} = e^{\frac{\beta}{2} \sum_{i \neq j} \ln |\lambda_i - \lambda_j|}$$

$$\Rightarrow P(\lambda_1, \dots, \lambda_N) = B_N e^{-\beta E(\{\lambda_i\})}$$

$$E(\{\lambda_i\}) = \frac{N}{2} \sum_{i=1}^N \lambda_i^2 - \frac{1}{2} \sum_{i \neq j} \ln |\lambda_i - \lambda_j|$$

repulsive

$\Rightarrow \lambda_i \equiv$ positions of charged particles
interacting via the 2d Coulomb interaction
confined on a line within a harmonic



$\Rightarrow$ competition between confinement and
the repulsive interactions