# Microscopic Legendre Transform and Canonical Ensemble

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#### Overview

Legendre Transform

Canonical ensemble

Microscopic Legendre Transform

Equilibrium statistical mechanics with MLT

Comparison with Jaynes' approach

## What is a Legendre Transform?

- ➤ A Legendre transform switches from one set of variables to another—especially useful when the new variables are more natural to work with.
- ► Common in physics when changing from quantities like velocity to momentum, or entropy to temperature.

#### Mathematical Definition

The Legendre-Fenchel transform:

$$f^*(p) = \inf_{x} (px - f(x)).$$

If f(x) is concave and differentiable:

$$f^*(p) = px - f(x)$$
, where  $p = \frac{df}{dx}$ 

► It replaces variable x with  $p = \frac{df}{dx}$ Known simply as Legendre transform in physics.

Legendre-Fenchel transforms in a nutshell, by Hugo Touchette

## Legendre Transform in Thermodynamics

- ▶ Internal energy: U(S, V)
- ► Temperature:  $T = \left(\frac{\partial U}{\partial S}\right)_V$
- ▶ Helmholtz free energy:  $S \rightarrow T$ .

$$F(T,V) = U(S,V) - TS$$

Legendre Transform of entropy:

$$\beta F(\beta, V) = \beta U - S(U, V).$$

Inverse Legendre Transform:

$$S(U, V) = \beta U - \beta F(\beta, V).$$



#### Canonical ensemble

The system is in thermal equilibrium with the heat reservoir:

$$p_i^* = \frac{e^{-\beta \varepsilon_i}}{\sum_{i=1}^W e^{-\beta \varepsilon_i}},$$

The equilibrium free energy:

$$\beta F(\beta) = -\ln \sum_{i=1}^{W} e^{-\beta \varepsilon_i}.$$

#### Canonical ensemble

Mean energy:

$$U^* = \frac{\partial}{\partial \beta}(\beta F) = \sum_{i=1}^W p_i^* \varepsilon_i$$

Entropy:

$$S(U^*) = \beta U^* - \beta F(\beta)$$
$$= -\sum_{i=1}^{W} p_i^* \ln p_i^* \equiv S^*$$

Thermodynamic entropy is equal to the Shannon entropy of the equilibrium distribution.

## Jaynes' Maximum Entropy Principle

#### Statement

Given a set of constraints, the most unbiased estimate of a probability distribution is the one that maximizes the Shannon entropy:

$$S=-\sum_i p_i \ln p_i$$

subject to:

$$\sum_i p_i = 1$$
 (normalization)

$$\sum_{i} p_{i} f_{i} = \langle f \rangle \quad \text{(constraint on mean value)}$$

#### Mathematical Solution

Use the method of Lagrange multipliers to incorporate constraints:

$$\mathcal{L} = -\sum_{i} p_{i} \ln p_{i} - \alpha \left( \sum_{i} p_{i} - 1 \right) - \lambda \left( \sum_{i} p_{i} f_{i} - \langle f \rangle \right)$$

▶ Solving this  $(\delta \mathcal{L} = 0)$  gives:

$$p_i = \frac{e^{-\lambda f_i}}{Z}, \quad Z = \sum_i e^{-\lambda f_i}$$

## The Maxent principle

Constraints: k = 1, ..., m

$$\langle f_k(x)\rangle = \sum_{i=1}^n p_i f_k(x_i) = F_k$$

Solution:

$$p_i = \exp[-\lambda_0 - \sum_{j=1}^m \lambda_j f_j(x_i)]$$

$$\lambda_0 = \ln Z = \ln \sum_{i=1}^n \exp[-\sum_{j=1}^m \lambda_j f_j(x_i)].$$

$$F_k = -\frac{\partial}{\partial \lambda_k} \ln Z.$$

$$S(F_1, ..., F_m) = \sum_{i=1}^m \lambda_k F_k + \ln Z(\lambda_1, ..., \lambda_m).$$

### Example: Canonical Ensemble

- ▶ Maximize entropy with given average energy  $(\sum_i p_i E_i = U)$ :
- Variational condition:

$$\sum_{i=1}^{W} (\beta \varepsilon_i + 1 + \ln p_i + \alpha) dp_i = 0.$$

Resulting distribution:

$$p_i = \frac{e^{-\beta E_i}}{Z}, \quad Z = \sum_i e^{-\beta E_i}$$

► Legendre Transform structure:

$$S(U) = \beta U + \ln Z(\beta).$$

### Interpretation

#### Canonical entropy ←⇒ Shannon entropy

- Equilibrium free energy is related to thermodynamic entropy via Legendre transform
- Thermodynamic entropy may be related to Shannon entropy via Jaynes' principle.

Can the equilibrium free energy be directly related to Shannon entropy via an optimization procedure?

## Microscopic Legendre Transform $(\mathscr{L}_{\mathscr{M}})$

$$\beta F(\beta, V, N) = \min_{U} \{\beta U - S(U, V, N)\}.$$
$$\beta U = \beta \sum_{i=1}^{W} p_{i} \varepsilon_{i} \to \sum_{i=1}^{W} p_{i}.(\beta \varepsilon_{i})$$

$$\beta F(\beta \varepsilon_1, ..., \beta \varepsilon_W) = \min_{p_1, ..., p_W} \left\{ \sum_{i=1}^W p_i.(\beta \varepsilon_i) - \tilde{\mathcal{S}}(p_1, ..., p_W) \right\}$$
$$\tilde{\mathcal{S}}(p_1, ..., p_W) = -\sum_{i=1}^W p_i \ln p_i - \alpha \left( \sum_{i=1}^W p_i - 1 \right)$$

Optimum condition:

$$\left. \frac{\partial \tilde{\mathcal{S}}}{\partial p_i} \right|_{p_i = p_i^*} = \beta \varepsilon_i.$$

## Exact differential of entropy

$$\tilde{S} = -\sum_{i=1}^{W} p_i \ln p_i - \alpha \left( \sum_{i=1}^{W} p_i - 1 \right)$$
$$d\tilde{S} = \sum_{i=1}^{W} \frac{\partial \tilde{S}}{\partial p_i} dp_i,$$

where

$$\frac{\partial \tilde{\mathcal{S}}}{\partial p_i} = -(1 + \ln p_i + \alpha).$$

At the optimal point  $(p_i = p_i^*)$ :

$$\frac{\partial \tilde{\mathcal{S}}}{\partial p_i} = \beta \varepsilon_i.$$

$$\beta \varepsilon_i + 1 + \ln p_i^* + \alpha = 0 \implies p_i^* = e^{-(1+\alpha)} \cdot e^{-\beta \varepsilon_i}$$

#### Definition of heat

$$dQ = \sum_{i=1}^{W} \varepsilon_i dp_i$$

$$d\tilde{S} = \sum_{i=1}^{W} \frac{\partial \tilde{S}}{\partial p_i} dp_i.$$

At equilibrium,

$$\frac{\partial \tilde{\mathcal{S}}}{\partial p_i} = \beta \varepsilon_i.$$

$$\therefore d\tilde{\mathcal{S}}^* = \beta \sum_{i=1}^W \varepsilon_i dp_i^*.$$

## MLT and Jaynes' approach

Consider a system in equilibrium with the heat reservoir.

A reversible process,

$$d\tilde{S} + dS_{\rm R} = 0.$$

$$d\tilde{S} = -\sum_{i=1}^{W} (1 + \ln p_i + \alpha) dp_i.$$
 
$$d\tilde{Q} = \sum_{i=1}^{W} \varepsilon_i dp_i.$$

Since this heat is exchanged with the heat reservoir, so the change in the entropy of the latter is  $dS_{\rm R} = -\beta \ dQ$ .

$$\sum_{i=1}^{W} (\beta \varepsilon_i + 1 + \ln p_i + \alpha) dp_i = 0.$$

#### Conclusion

- The microscopic Legendre transform offers a novel perspective on the relationship between Shannon entropy and Helmholtz free energy.
- It provides an alternative derivation of the canonical ensemble and underscores the connection between information-theoretic and thermodynamic entropy.
- ► Future applications: Non-equilibrium processes. Generalized entropies.

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RESEARCH

Microscopic Legendre Transform, Canonical Ensemble and Jaynes' Maximum Entropy Principle

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# Thank You for your attention