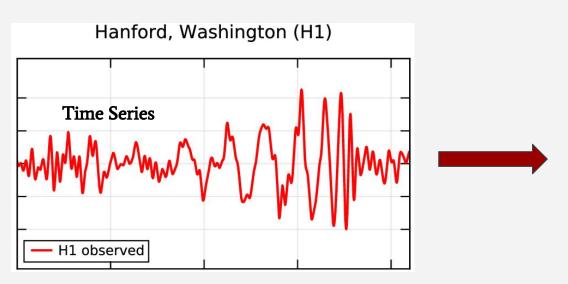
Prospects for inferring the population properties of compact binaries

Aditya Vijaykumar

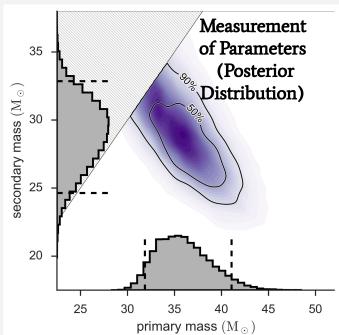
The Future of Gravitational Wave Astronomy 28th October 2025



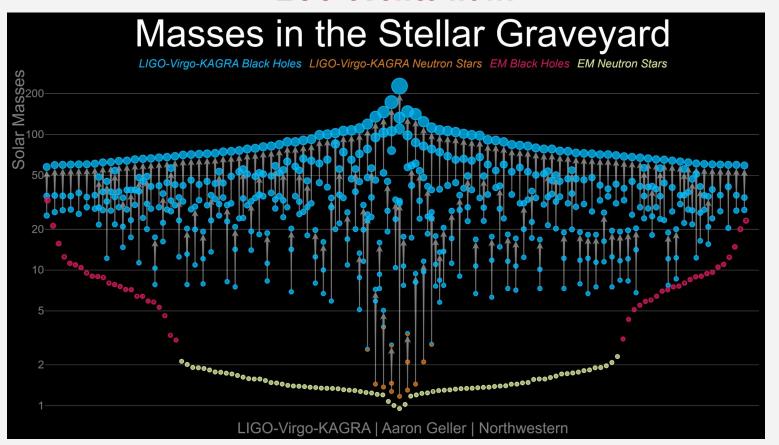
GW150914



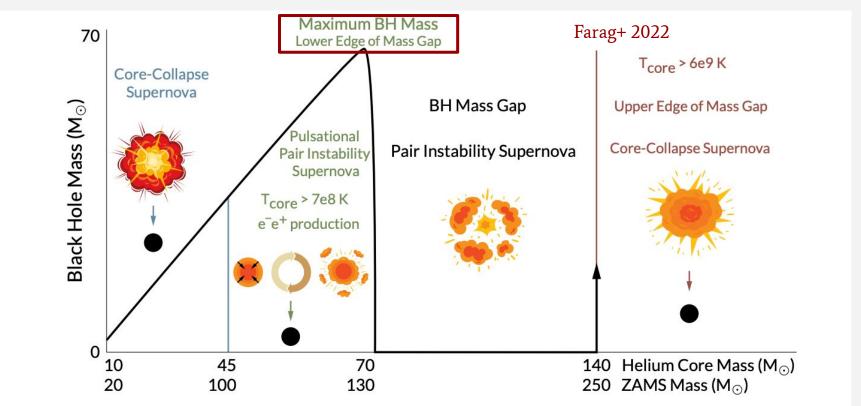
Abbott+ 2016, arXiv:1602.03837, arXiv:1602.03840



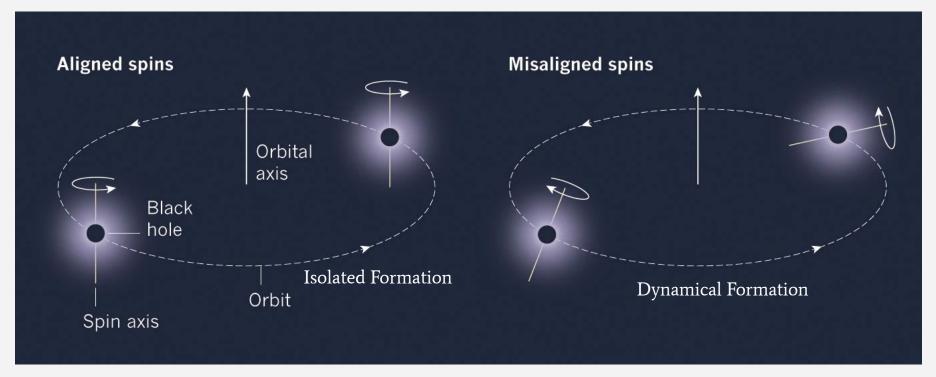
~200 events now!



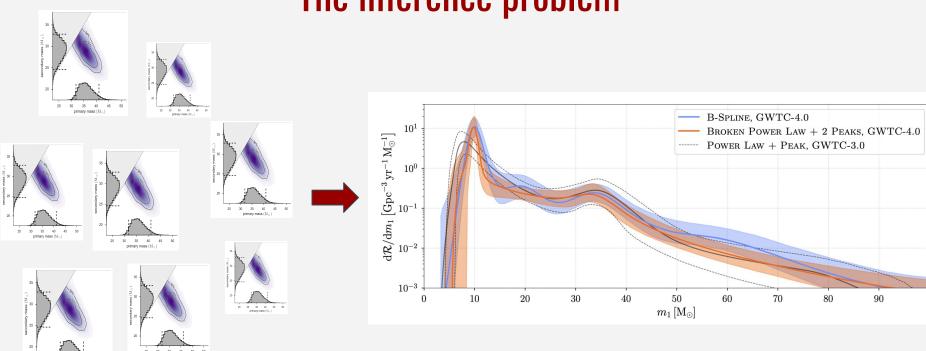
Population-level properties encode astrophysics



Population-level properties encode astrophysics



The inference problem



Many uncertain measurements from different sources

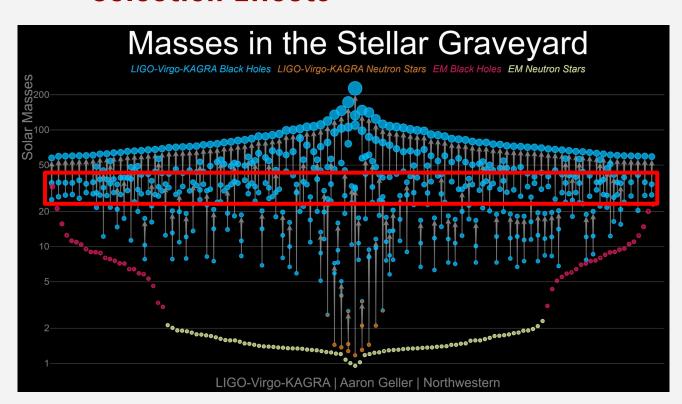
Measurements of population properties [LVK+, arXiv:2508.08083]

Selection Effects

 $A_{
m GW} \propto rac{{\cal M}^{5/6}}{D_L}$

Inherently easier to detect massive binaries, and binaries that are close by.

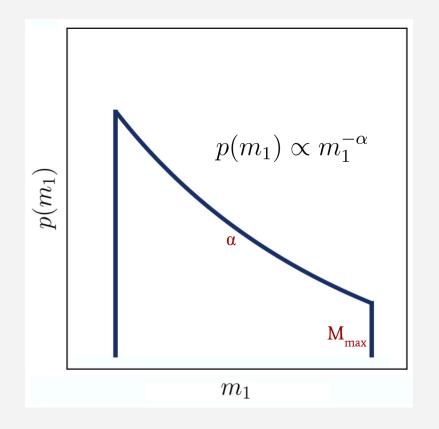
...but we can correct for this!



Population Model

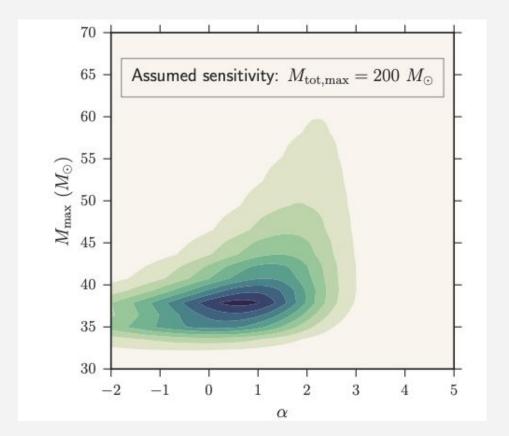
Simple Model: The mass function of stars is a power law with exponent 2.3 with a M_{max} ~300 Msun

What are the corresponding numbers for black holes?



Abbott et al 2020 (GWTC-2 Population paper)

Population Model

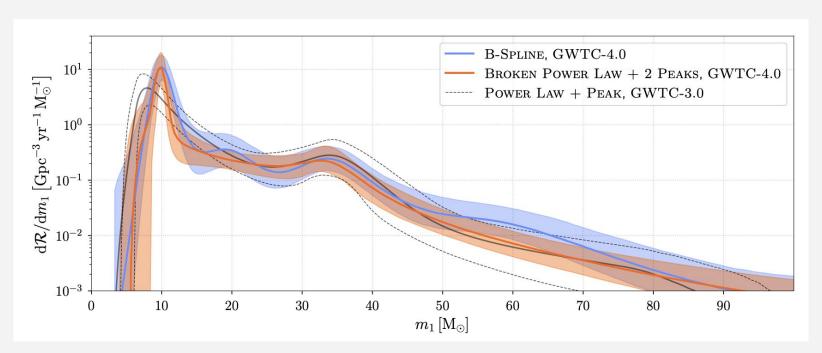


 $\Lambda = {\alpha, M_{max}}$

Fishbach and Holz 2017

Used only the first 6 reported GW detections

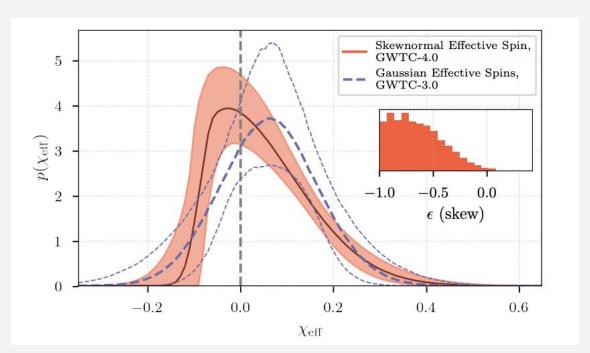
Mass Distribution



With more data, we now know that a simple power law is not a good fit to the data.

LVK+, arXiv:2508.08083

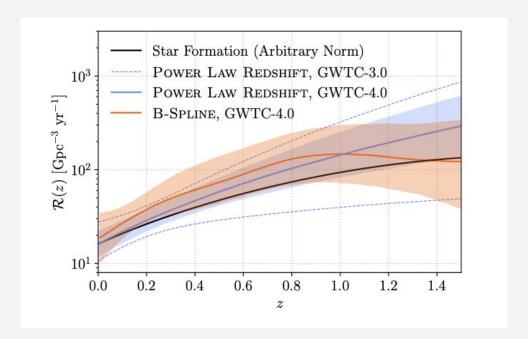
Effective Spin Distribution



Distribution of effective spins skewed towards positive values, but significant support for negative spins.

LVK+, arXiv:2508.08083

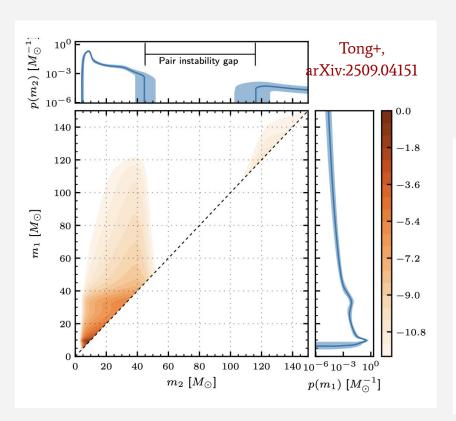
Redshift Distribution



The rate of mergers per unit volume increases with redshift, consistent with star formation rate density.

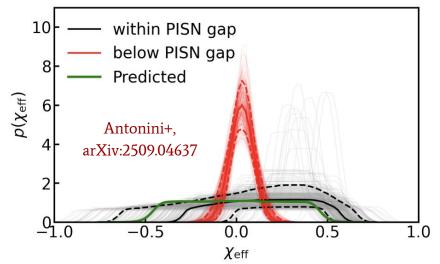
LVK+, arXiv:2508.08083

Pair instability feature in the data?



The *secondary* mass distribution breaks after about 45 Msun, but there are primary masses above that value.

Hierarchical mergers?



Challenges: What features should we look for in the data?

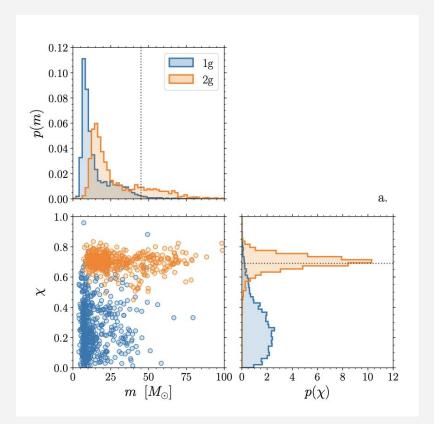
Difficult to disentangle formation channels: e.g., most formation channels can produce 30 Msun BHs.

Need to look for robust features.

Hierarchical merger spins are a robust feature, peaking around 0.7.

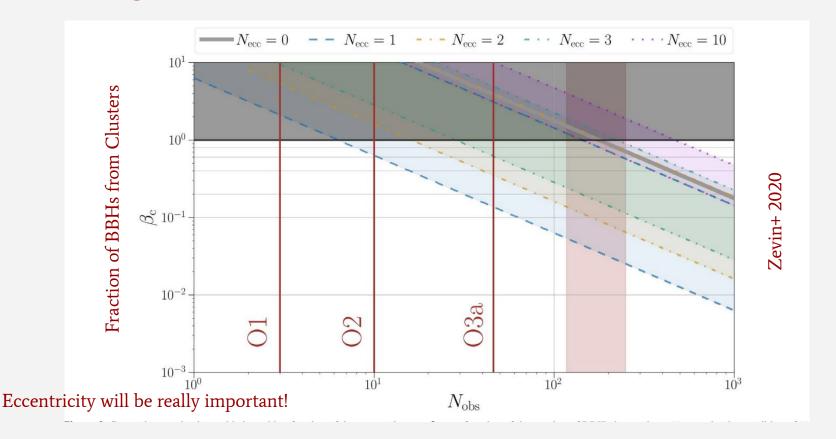
More candidate features:

- 1. Do higher mass black holes have larger spins?
- Does the 1G mass distribution shift to larger values as a function of redshift?
 a etc.

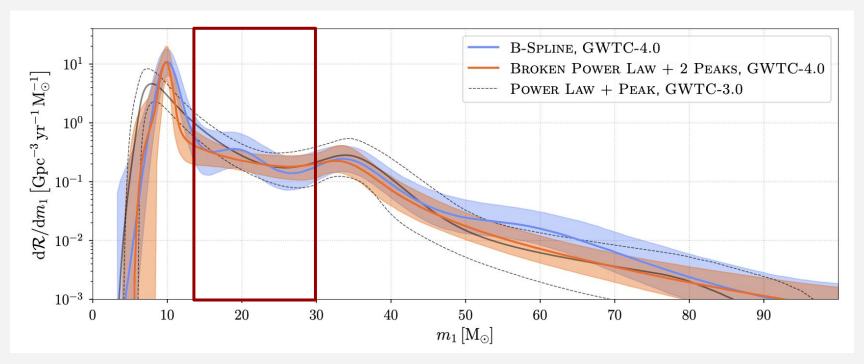


Gerosa & Fishbach 2021

Challenges: What features should we look for in the data?



Challenges: How should we look for features in the data?



Flexible vs parametric models: Bias variance trade-off! Current flexible models are expensive in multiple dimensions. New methods?

Challenges: How should we do selection effects?

Compact Binary Coalescence Sensitivity Estimates with Injection Campaigns during the LIGO-Virgo-KAGRA Collaborations' Fourth Observing Run

Reed Essick,^{1, 2, 3, a} Michael W. Coughlin,^{4, 5} Michael Zevin,^{6, 7, 8} Deep Chatterjee,⁹ Teagan A.

We describe the effort to characterize gravitational-wave searches and detector sensitivity to different types of compact binary coalescences during the LIGO-Virgo-KAGRA Collaborations' fourth observing run. We discuss the design requirements and example use cases for this data product, constructed from $> 4.33 \times 10^8$ injections during O4a alone. We also identify subtle effects with high confidence, like diurnal duty cycles within detectors. This paper accompanies a public data release of the curated injection set [1, 2], and the appendixes give detailed examples of how to use the publicly available data.

Selection effects are *expensive*, and requirements will only increase with time.

Summary

Thinking beyond single-event parameters and viewing events as a population can help
 us constrain astrophysical properties

- Many challenges:
 - Modeling the population
 - Selection effects
 - Looking for robust features
- Synergies with other observational channels?
 - o LISA?
 - o Gaia BHs?
 - X-ray Binaries?

Thanks for Listening!

Get in Touch!

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Email: aditya@utoronto.ca

Extra Slides

Ingredients

$$p\left(\Lambda \mid \{d\}\right) \propto p\left(\{d\} \mid \Lambda\right) p(\Lambda)$$

Population parameters, could be:

Population Likelihood

Prior

- 1. Slope of the mass distribution
- 2. Minimum and maximum masses
- 3. Location of e.g. "peaks" in the mass distribution
- 4. Mean and variance of spin magnitudes
- 5. ...etc.

Population Likelihood

• The arrival of GWs at our detectors is a Poisson process, so we might think

$$p(N_{\text{ev}}|N_{\text{exp}}) = \frac{N_{\text{exp}}^{N_{\text{ev}}}e^{-N_{\text{exp}}}}{N_{\text{ev}}!}$$

 N_{ev} is the number of events observed, and N_{exp} is the number of events expected given a population model parameterized by Λ .

Population Likelihood

• The arrival of GWs at our detectors is a Poisson process, so we might think

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 N_{ev} is the number of events observed, and N_{exp} is the number of events expected given a population model parameterized by Λ .

- ...But it is not a homogeneous Poisson process: N_{exp} depends on the population model and the detector sensitivity to each member of the population
 - e.g. some population model might predict 5 Msun and 50 Msun BHs at the same rate, but selection effects mean we'll detect more 50 Msun BHs.

Population Likelihood

Assuming a inhomogeneous Poisson process, the population likelihood can be written as

$$p(\{d\}|\{\theta\},\Lambda) \propto e^{-N_{\rm exp}(\Lambda)} \prod_i^{N_{\rm ev}} \int {\rm d}\theta_i \ p(d_i|\theta_i) \left[\mathcal{R} \ p(\theta_i|\Lambda)\right]$$
 For a full derivation, see
$$\begin{array}{c} {\rm Single-event} \\ {\rm Vitale\ et\ al\ 2020} \end{array}$$
 Population model

$$N_{\mathrm{exp}}(\Lambda) = \int \mathrm{d}\theta \, \mathrm{d}t \, \left[\mathcal{R} \, p(\theta|\Lambda) \right] \, p(\det|\theta)$$

Probability of detecting a signal with parameters θ . Is either 1 or 0 .

$$p(\{d\}|\{\theta\},\Lambda) \propto e^{-N_{\exp}(\Lambda)} \prod_{i}^{N_{\text{ev}}} \int d\theta_i \ p(d_i|\theta_i) \left[\mathcal{R} \ p(\theta_i|\Lambda) \right]$$
$$N_{\exp}(\Lambda) = \int d\theta \ dt \left[\mathcal{R} \ p(\theta|\Lambda) \right] \ p(\det|\theta)$$

Both these expressions involve integrals in high dimensions. Solution: Convert integrals into Monte Carlo sums.

$$\int dx f(x) p(x) = \mathbb{E}_p [f(x)] = \frac{1}{N_{\text{samps}}} \sum_{k=1}^{N_{\text{samps}}} f(x_k)|_{x_k \sim p(x)}$$

Integrals like these are expectation values of a function f(x) under some probability distribution p(x).

Can be approximated by drawing samples from p(x), and calculating an average of the f(x) over those samples

$$N_{\rm exp}(\Lambda) = \int d\theta \, dt \, \left[\mathcal{R} \, p(\det|\theta) \, \frac{p(\theta|\Lambda)}{p_{\rm draw}(\theta|\Lambda_{\rm draw})} \right] \, p_{\rm draw}(\theta|\Lambda_{\rm draw})$$

$$N_{\rm exp}(\Lambda) = \frac{1}{N_{\rm inj}} \sum_{k=1}^{N_{\rm found}} \frac{p(\theta_k | \Lambda)}{p_{\rm draw}(\theta_k | \Lambda_{\rm draw})}$$

Choose a fiducial draw population distribution → simulate and inject event parameters → check if they are detected.

Approximate N_{exp} as a sum over "found" injections.

$$\int d\theta_i \ p(d_i|\theta_i) \left[p(\theta_i|\Lambda) \right] = \int d\theta_i \ p(\theta_i|d_i) \left[\frac{p(\theta_i|\Lambda)}{\pi_{\text{PE}}(\theta_i)} \right] = \frac{1}{N_{\text{samps}}^i} \sum_{k=1}^{N_{\text{samps}}^i} \left[\frac{p(\theta_i^k|\Lambda)}{\pi_{\text{PE}}(\theta_i^k)} \right]$$

Similarly, the integral over individual event likelihoods can be approximated using single event posterior samples.

We now have everything to calculate the likelihood!

$$p(\{d\}|\{\theta\},\Lambda) \propto e^{-N_{\rm exp}(\Lambda)} \prod_i^{N_{\rm ev}} \int {\rm d}\theta_i \ p(d_i|\theta_i) \left[\mathcal{R} \ p(\theta_i|\Lambda)\right]$$
 Approximate using simulated "injections" Approximate using single-event PE samples

However, Monte Carlo sums have uncertainties, which become important especially when the population model has sharp features [Essick and Farr 2022, Talbot and Golomb 2023].

Typical strategy: set likelihood to 0 when variance in the Monte Carlo Estimator is above a threshold.