COSMOLOGICAL BACKGROUNDS of GRAVITATIONAL WAVES



SCHEDULE

PHYSICS OF THE EARLY UNIVERSE (HYBRID) 3 - 12 January 2022

Note that all dates and times are in Indian Standard Time (IST). IST is behind Japan Standard Time by 3:30 hours, ahead of Central European Winter Time by 4:30 hours, Greenwich Mean time by 5:30 hours and Eastern Standard Time by 10:30 hours.

Introductory remarks by the Organisers on 3rd January at 11:15 hrs

Welcome remarks by Centre Director, Prof. Rajesh Gopakumar on 4th January at 11:30 hrs

	09:30-11:00	11:30-13:00	14:30-16:00	16:30-18:00	18:30-20:00
Monday, January 3		Introductory remarks (11:15 hrs)		Inflation I	CPT T
		CPT I			
Tuesday, January 4		Welcome remarks CPT II	Bounces I	Inflation II	Inflation T
Wednesday, January 5		CPT III	Bounces II	GW I	Inflation III
Thursday, January 6	PPEU I	CPT IV	Bounces III	GW II	Imlation IV
Friday, January 7	PPEU II	CWD I	PPEU T	GW III	GW T
Monday, January 10	PPEU III	CWD II	DEMG I	DM I	Bounces T
Tuesday, January 11	PPEU IV	CWD III	DEMG II	DM II	DEMG T
Wednesday, January 12	CWD T	CWD IV	DEMG III	DM III	DM T
					Closing remarks (20:00 hrs)

GW

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CPT

- Cosmological perturbation theory - by David Wands

Inflation

- The inflationary paradigm - by William Kinney

Bounces

- The bouncing scenario - by Patrick Peter

PPEU

- Particle physics in the early universe - by Masahide Yamaguchi

CWD

- Comparison of models of the early universe with the cosmological data - by Shiv Sethi

GWs

- Generation and imprints of primordial gravitational waves - by Daniel Figueroa

DM

- Dark matter - by Katelin Schutz

DEMG

- Dark energy and modified gravity - by Alessandra Silvestri

GW

PHYSICS OF THE EARLY UNIVERSE, ICTS Bangalore, India (Jan. 3rd - 12th 2022)



- Generation and imprints of primordial gravitational waves - by Daniel Figueroa



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GWs



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COSMOLOGICAL BACKGROUNDS of GRAVITATIONAL WAVES

DANIEL G. FIGUEROA IFIC, Valencia, Spain

PHYSICS OF THE EARLY UNIVERSE, ICTS Bangalore, India (Jan. 3rd - 12th 2022)

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GWs



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COSMOLOGICAL BACKGROUNDS



'PRIMORDIAL' BACKGROUNDS



- Generation and imprints of primordial gravitational waves - by Daniel Figueroa





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COSMOLOGICAL BACKGROUNDS

'PRIMORDIAL' BACKGROUNDS

GW = Gravitational Waves



- Generation and imprints of primordial gravitational waves - by Daniel Figueroa



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COSMOLOGICAL BACKGROUNDS



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GW: PRIMORDIAL, ergo COSMOLOGICAL



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COSMOLOGICAL BACKGROUNDS



'PRIMORDIAL' BACKGROUNDS

GW: PRIMORDIAL, ergo COSMOLOGICAL Early Universe



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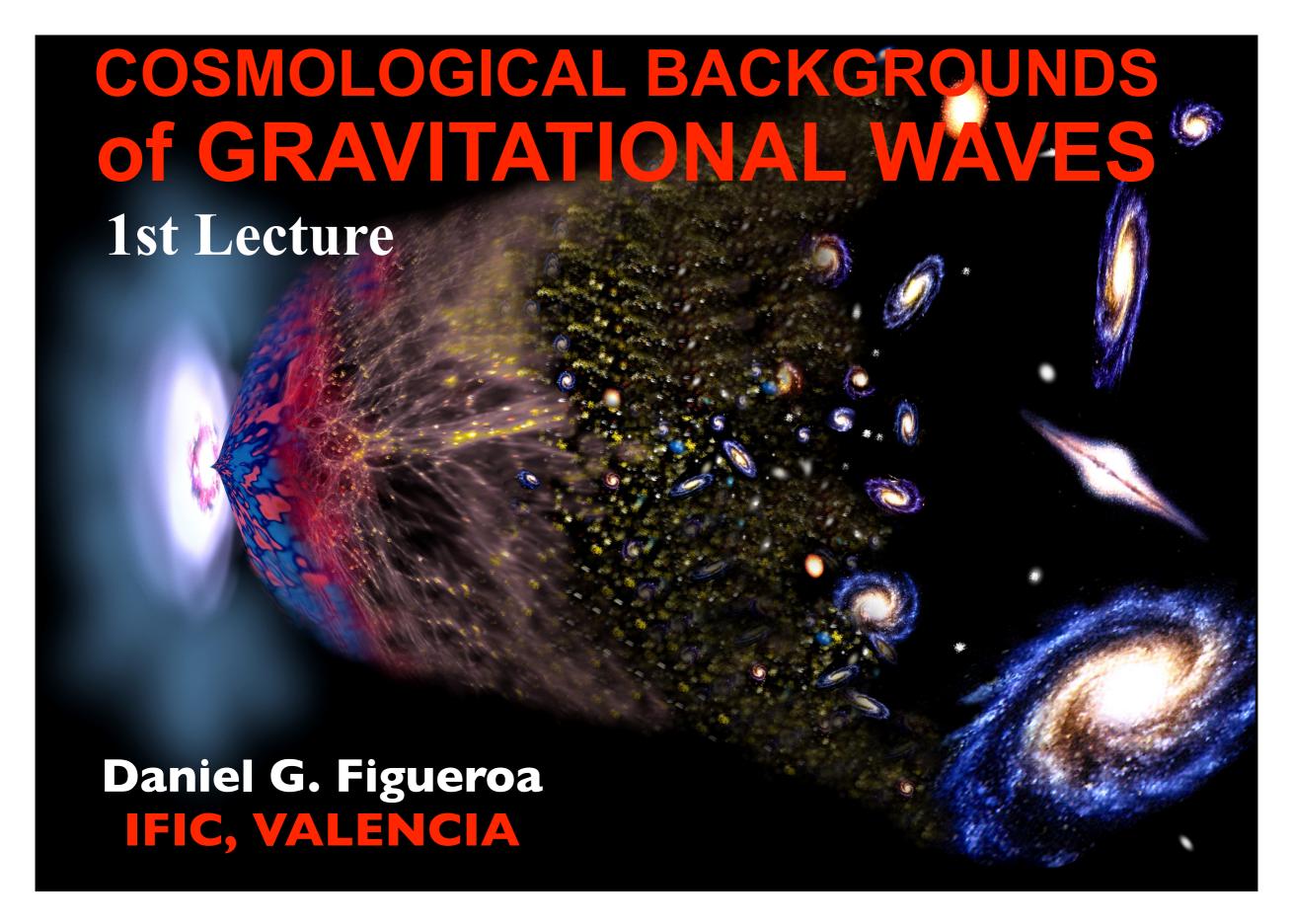
COSMOLOGICAL BACKGROUNDS



'PRIMORDIAL' BACKGROUNDS

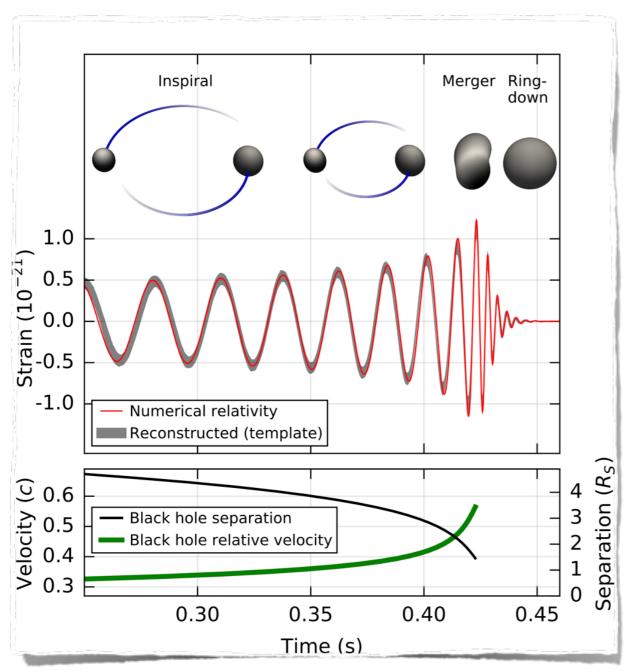
GW: PRIMORDIAL, ergo COSMOLOGICAL

STOCHASTIC



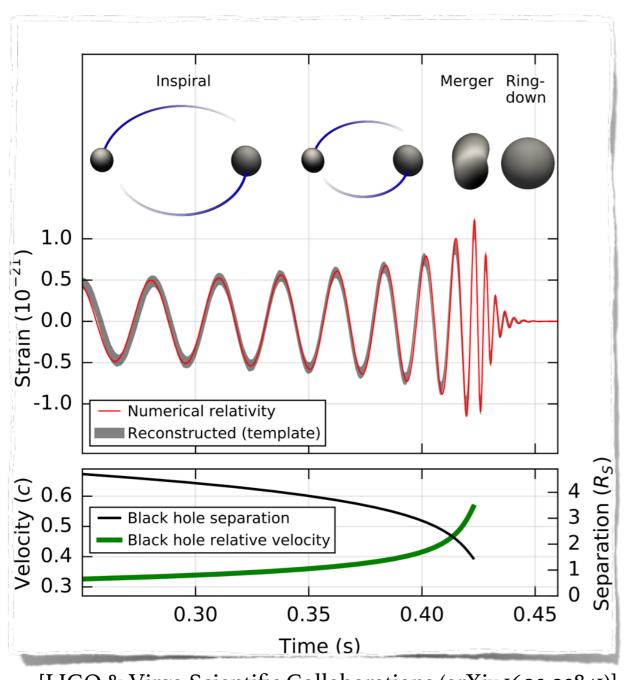
MOTIVATION

(cosmologist biased)



[LIGO & Virgo Scientific Collaborations (arXiv:1602.03841)]

Gravitational Waves (GWs) detected! [by LIGO/VIRGO]





[LIGO & Virgo Scientific Collaborations (arXiv:1602.03841)]

Einstein 1916 ... LIGO/VIRGO 2015-2021

* O(10) Solar mass Black Holes (BH) exist

* We can test the BH's paradigm and Neutron Star physics

* We can further test General Relativity (GR) [so far no deviation]

* We can observe the Universe through GWs



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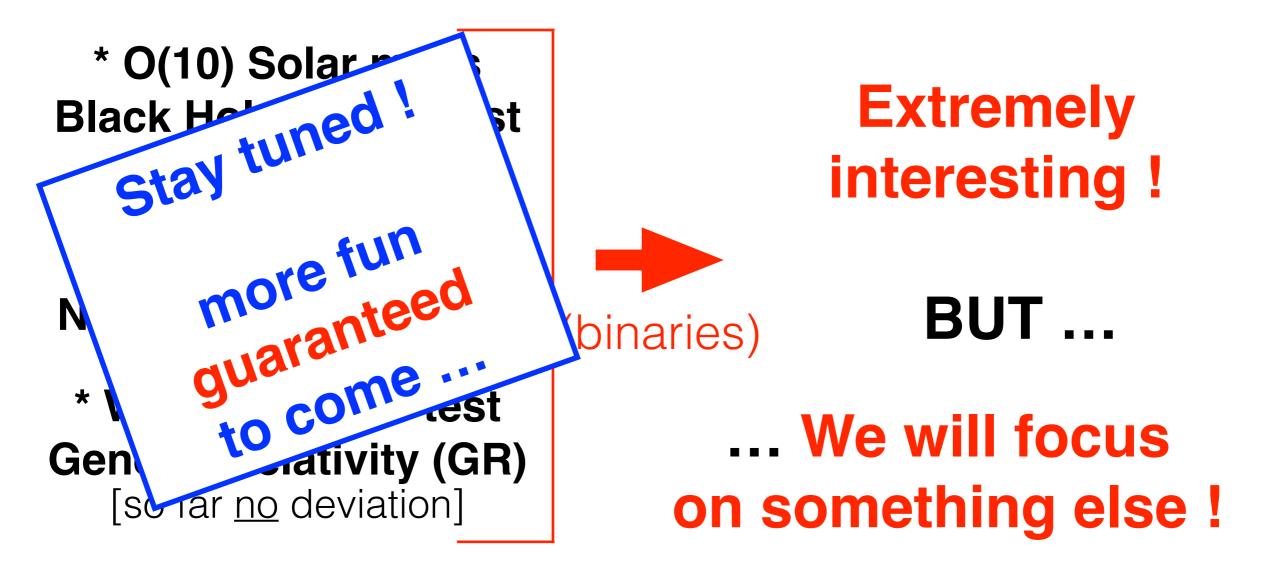
* We can observe the Universe through GWs

Extremely interesting!



BUT ...

... We will focus on something else!



* We can observe the Universe through GWs

* O(10) Solar mass Black Holes (BH) exist

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BUT ...

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* Cosmology with GWs

* Cosmology with GWs

* Late Universe:

* Early Universe:

* Cosmology with GWs

* Late Universe:

Standard sirens: distances in cosmology; Measuring H0 and EoS dark energy; cosmological parameters; modify gravity, lensing, ...



* Cosmology with GWs

* Late Universe:

* Early Universe: High Energy Particle Physics

* Cosmology with GWs

* Late Universe:

* Early Universe: High Energy Particle Physics

* Cosmology with GWs

* Late Universe:

* Early Universe: High Energy Particle Physics

Can we really probe High Energy Physics using Gravitational Waves (GWs)? How?

Why?

One and ONLY One reason ...

WEAKNESS of GRAVITY:

ADVANTAGE: GW DECOUPLE upon Production

DISADVANTAGE: DIFFICULT DETECTION

WEAKNESS of GRAVITY:

ADVANTAGE: GW DECOUPLE upon Production

DISADVANTAGE: DIFFICULT DETECTION

 \bigcirc **ADVANTAGE**: GW \rightarrow Probe for Early Universe

```
\rightarrow \left\{ \begin{array}{l} \textbf{Decouple} \rightarrow \underline{Spectral\ Form\ Retained} \\ \textbf{Specific\ HEP} \Leftrightarrow \underline{Specific\ GW} \end{array} \right.
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WEAKNESS of GRAVITY:

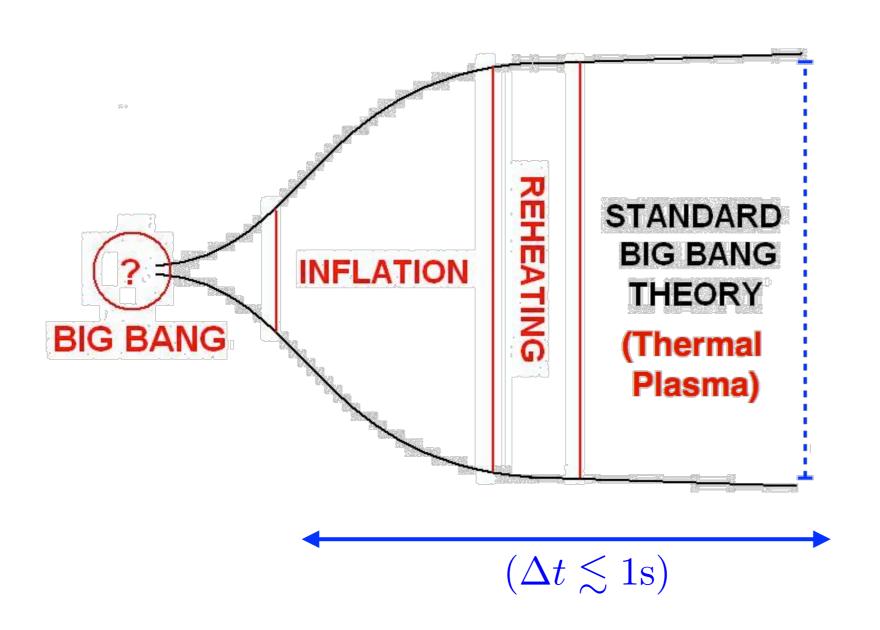
ADVANTAGE: GW DECOUPLE upon Production

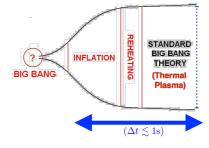
DISADVANTAGE: DIFFICULT DETECTION

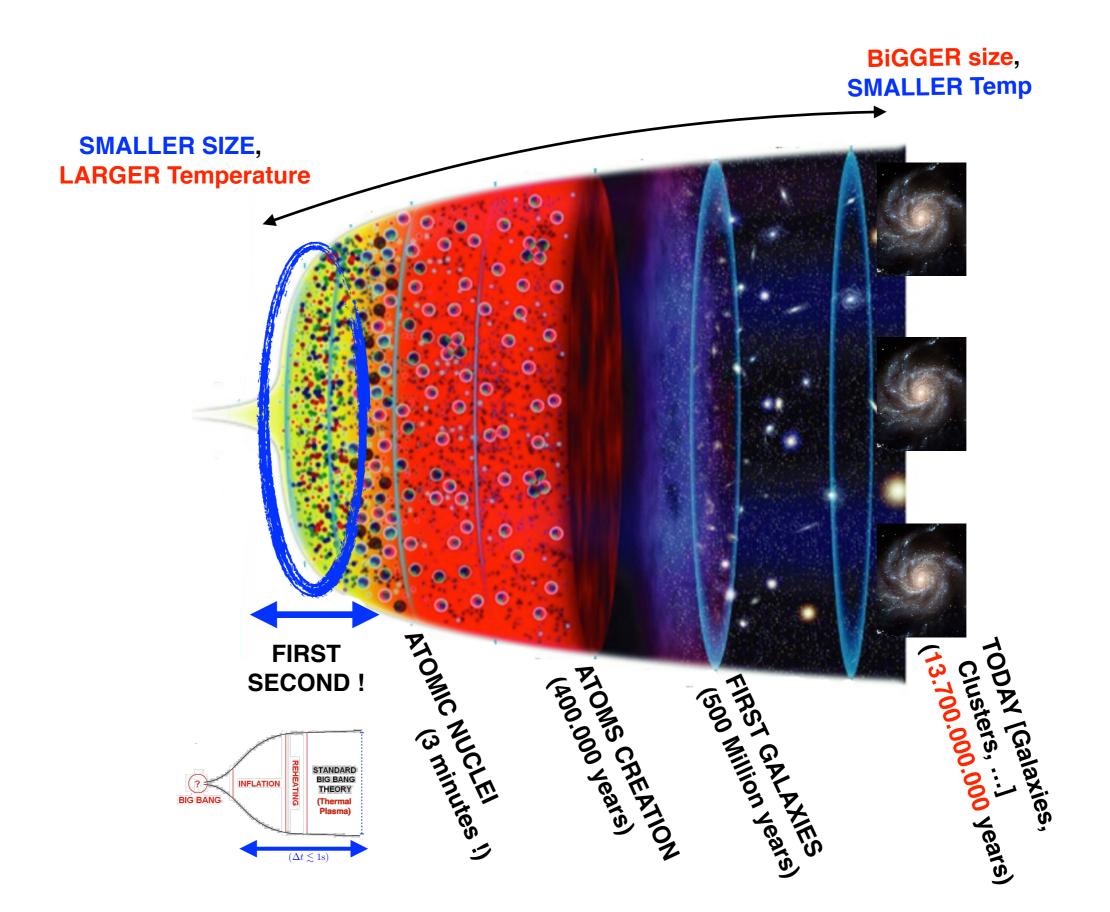
2 ADVANTAGE: GW \rightarrow Probe for Early Universe

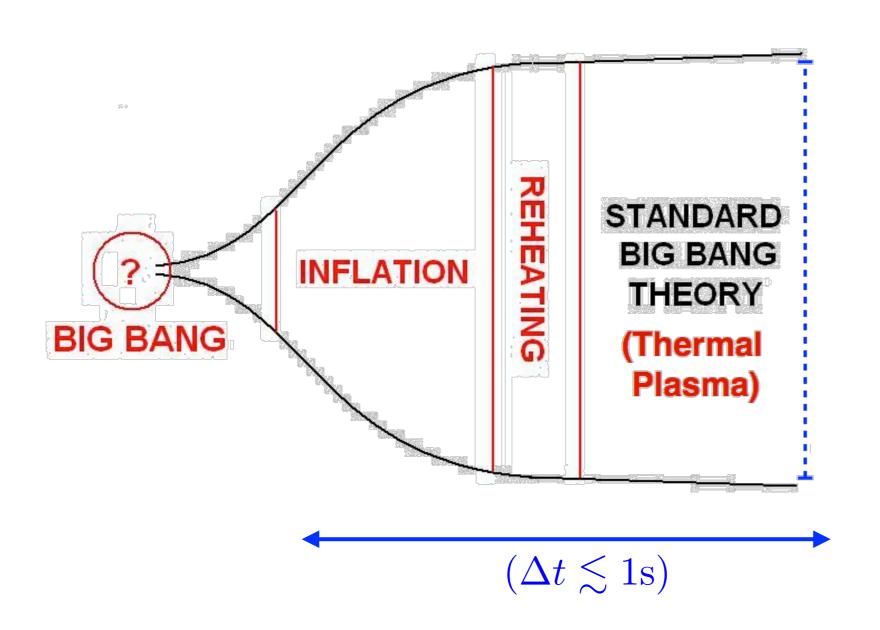
```
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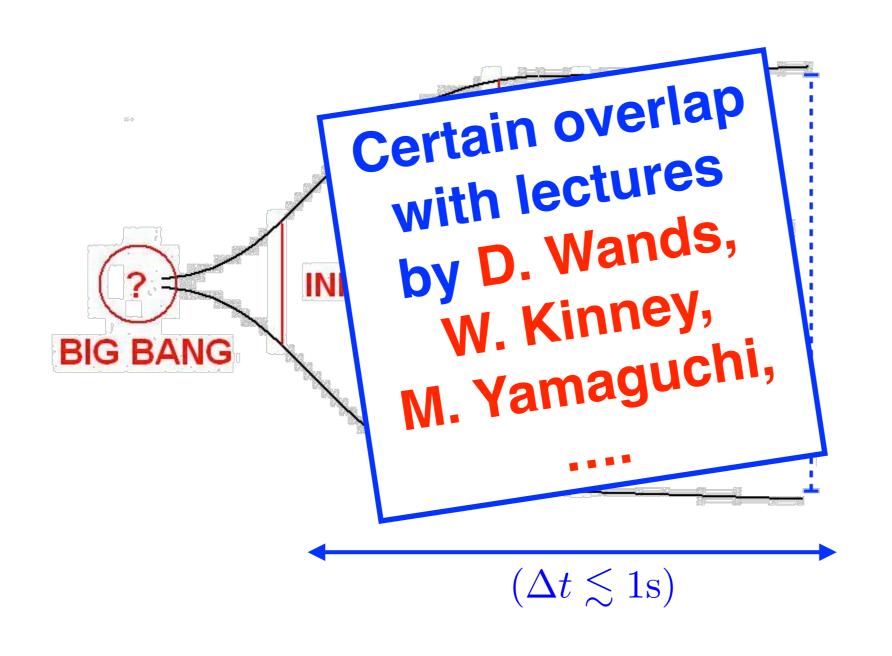
What processes of the early Universe?

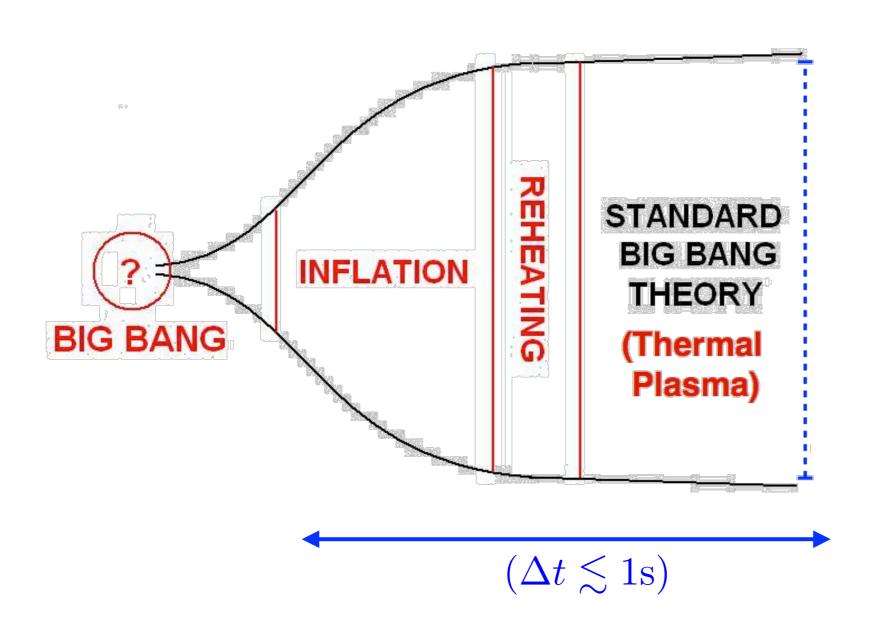


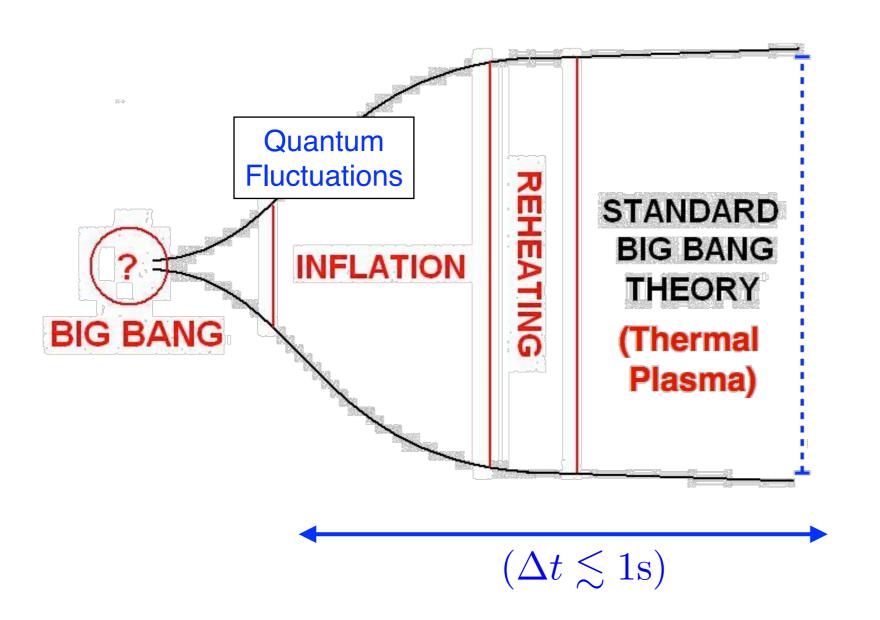


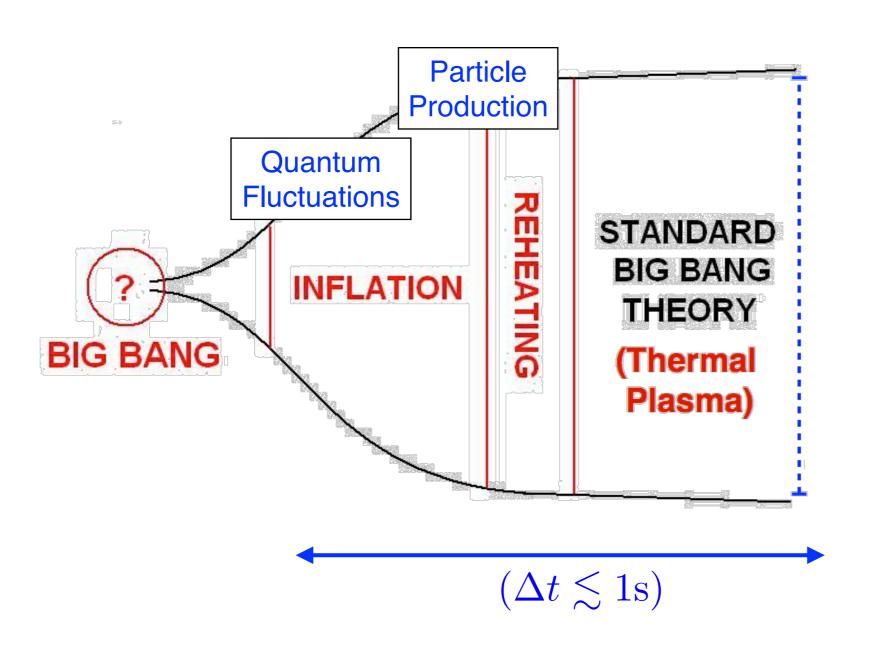


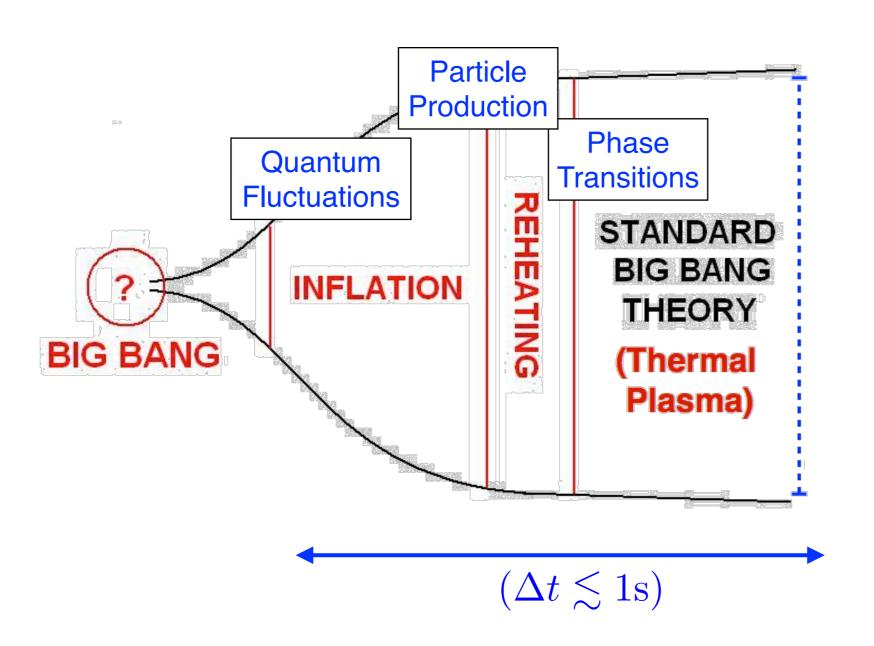


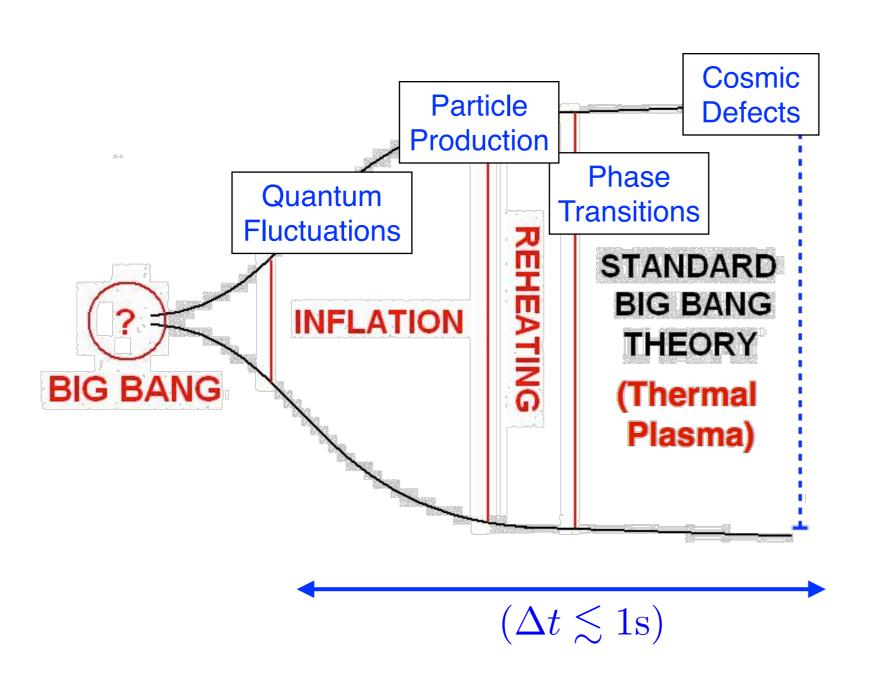


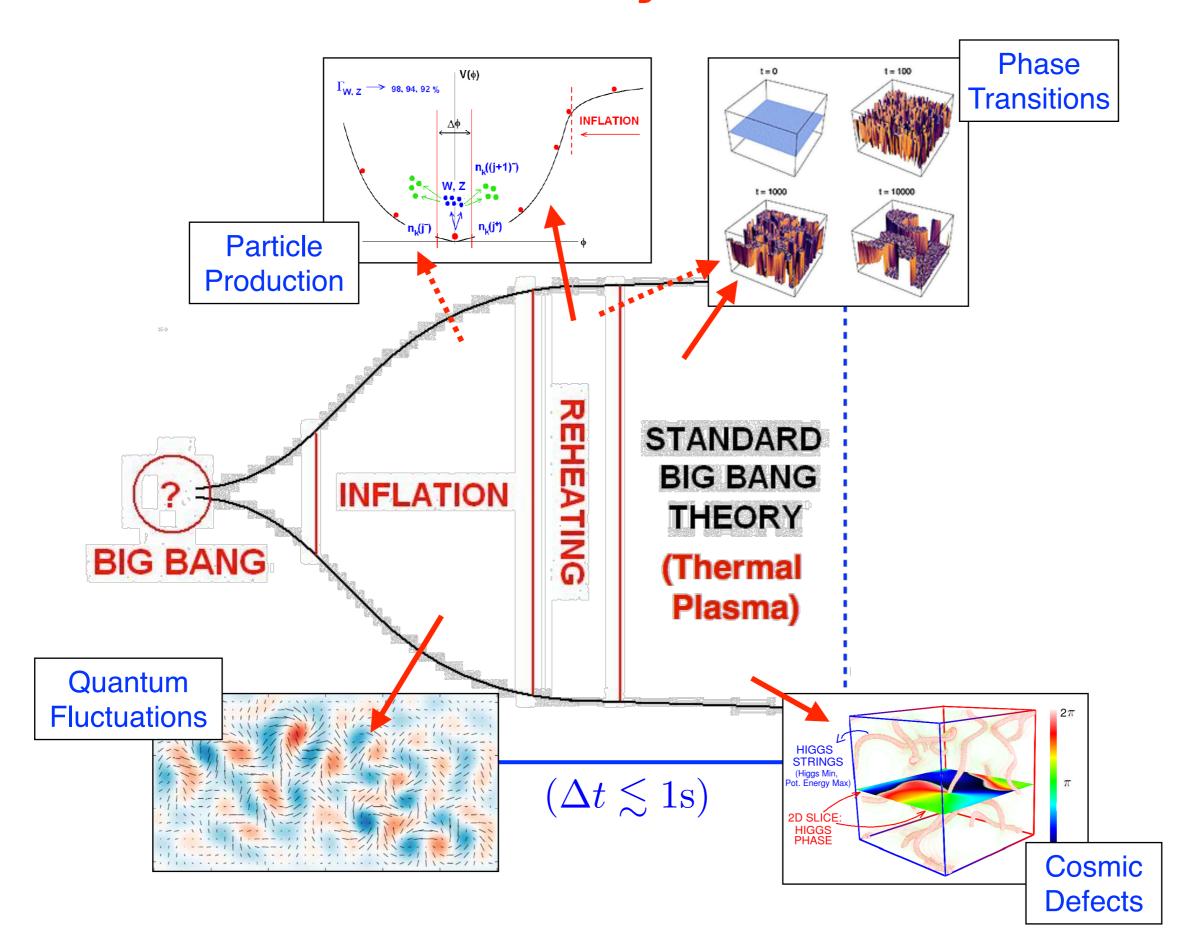


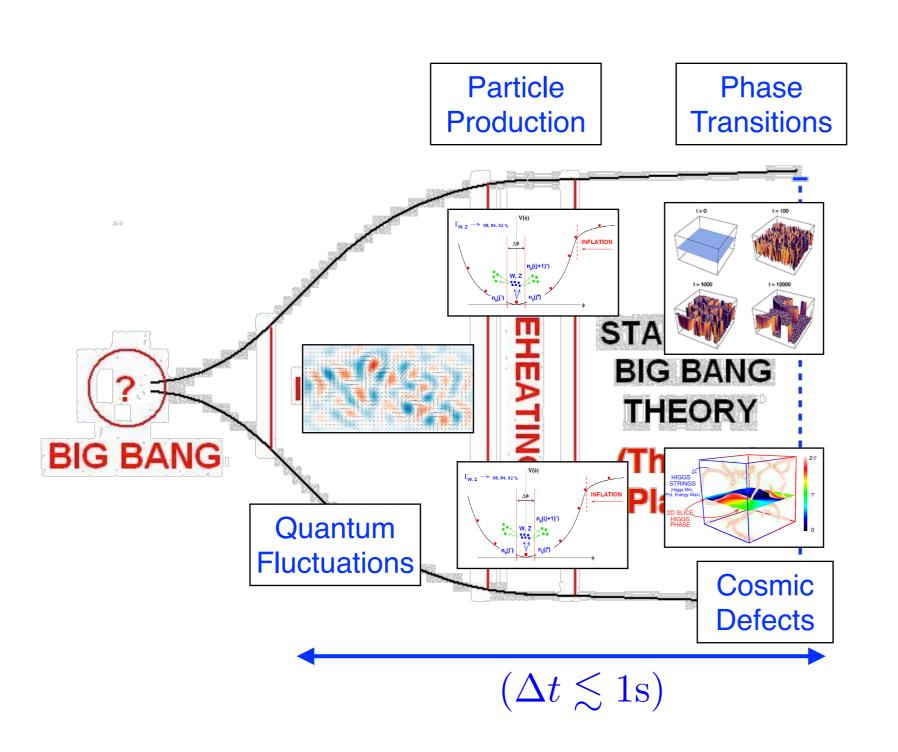


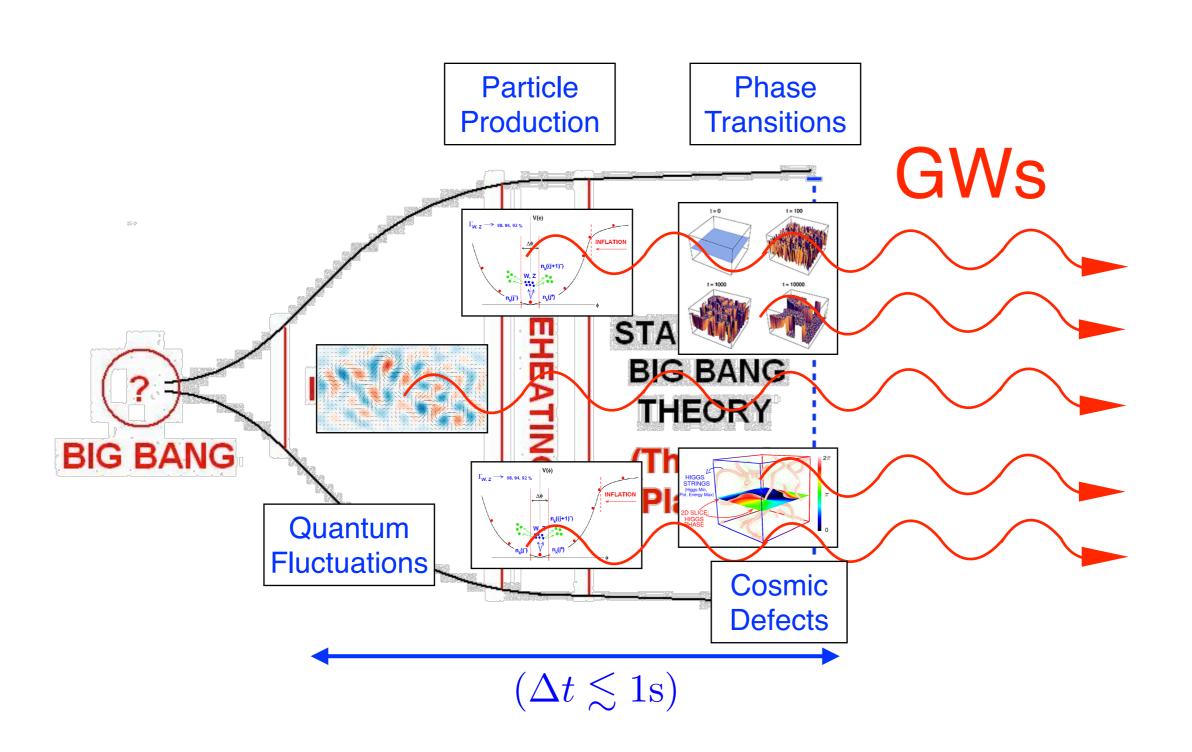


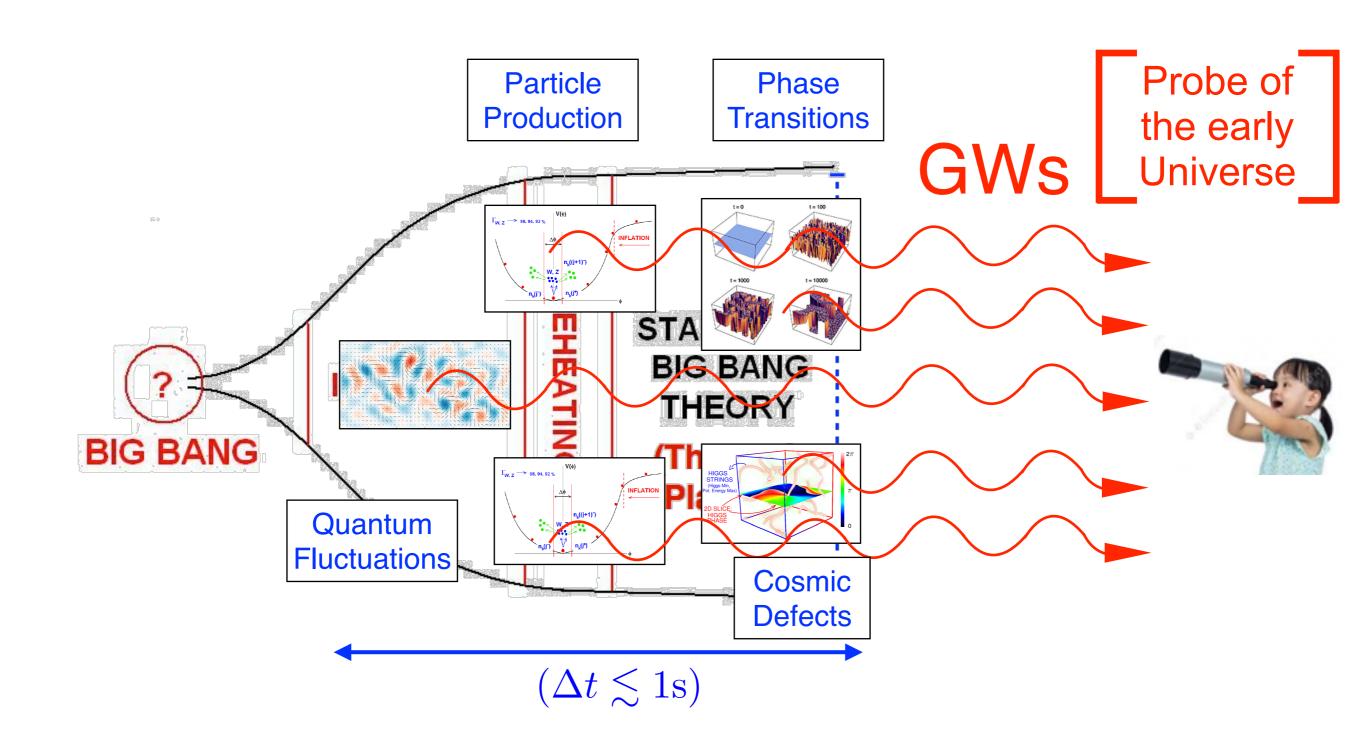




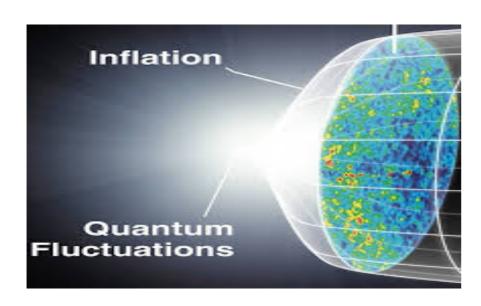






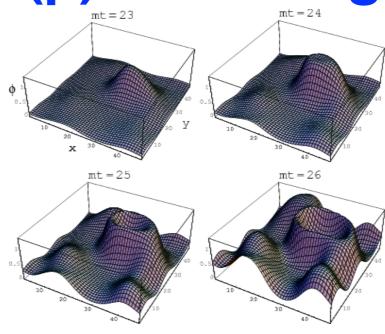


Inflationary Period



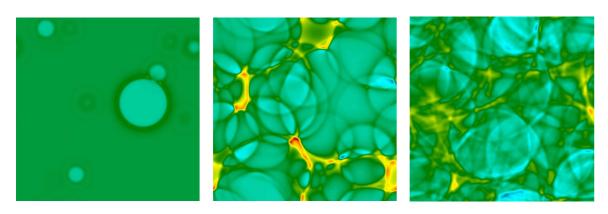
(Image: Google Search)

(p)Reheating



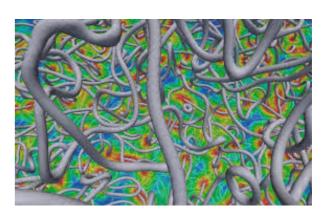
(Fig. credit: Phys.Rev. D67 103501)

Phase Transitions



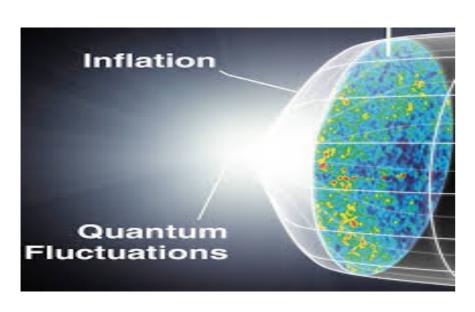
(Image: PRL 112 (2014) 041301)

Cosmic Defects



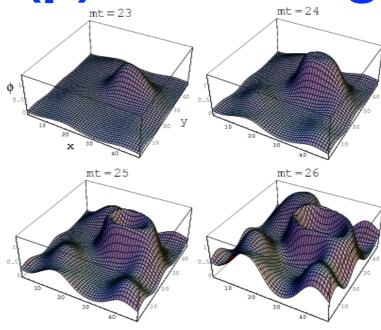
(Image: Daverio et al, 2013)

Inflationary Period



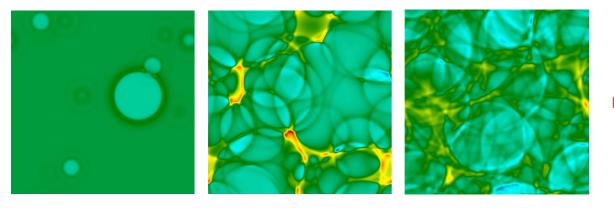
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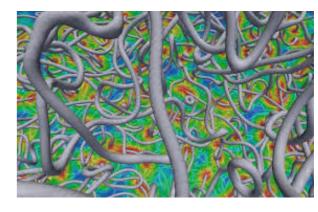
(Fig. credit: Phys.Rev. D67 103501)

Phase Transitions



(Image: PRL 112 (2014) 041301)

Cosmic Defects



(Image: Daverio et al, 2013)

OUTLINE

- 1) GWs from Inflation
- 2) GWs from Preheating
- 3) GWs from Phase Transitions
- 4) GWs from Cosmic Defects

OUTLINE (~ 4.5 h)

Early

2) GWs from Preheating —— ~ 1 h

1) GWs from Inflation —— ~ 1 h

3) GWs from Phase Transitions ——— ~ 1 h

4) GWs from Cosmic Defects —— ~ 1 h

Universe

OUTLINE (~ 4.5 h)

- 0) Gravitational Waves definition
- 1) GWs from Inflation
- **Early Universe**
- 2) GWs from Preheating
- 3) GWs from Phase Transitions
- 4) GWs from Cosmic Defects

OUTLINE (~ 4.5 h)

- 1) Gravitational Waves definition
- 2) GWs from Inflation
- **Early Universe**
- 3) GWs from Preheating
- 4) GWs from Phase Transitions
- 5) GWs from Cosmic Defects

OUTLINE (~ 4.5 h)

1) Gravitational Waves definition

1st lecture (~ 1.5 h)

- 2) GWs from Inflation
- 3) GWs from Preheating
- 4) GWs from Phase Transitions
- 5) GWs from Cosmic Defects

OUTLINE (~ 4.5 h)

1) Gravitational Waves definition

1st lecture (~ 1.5 h)

- 2) GWs from Inflation
- 3) GWs from Preheating

2nd lecture (~ 1.5 h)

- 4) GWs from Phase Transitions
- 5) GWs from Cosmic Defects

OUTLINE (~ 4.5 h)

1) Gravitational Waves definition

1st lecture (~ 1.5 h)

Early 3) GWs from Preheating

2) GWs from Inflation

2nd lecture (~ 1.5 h)

4) GWs from Phase Transitions

ansitions

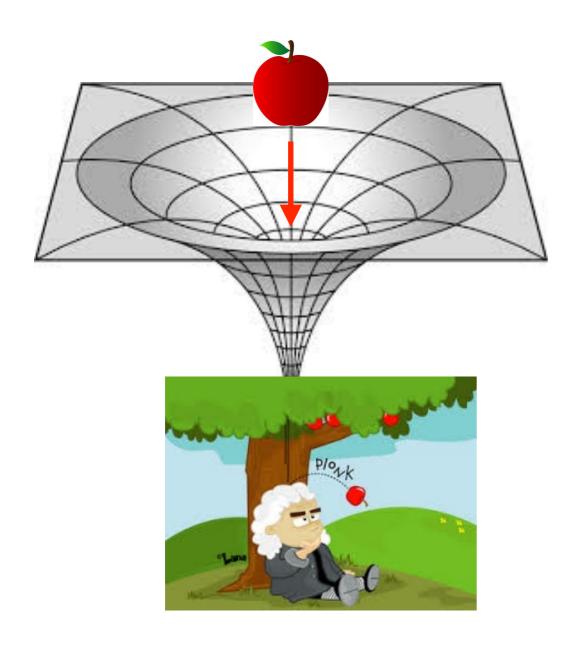
5) GWs from Cosmic Defects

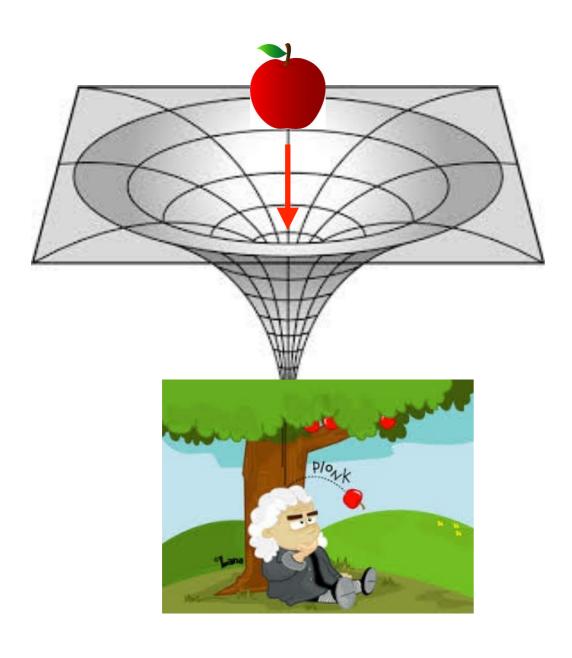
3rd lecture (~ 1.5 h)



Let's Start!

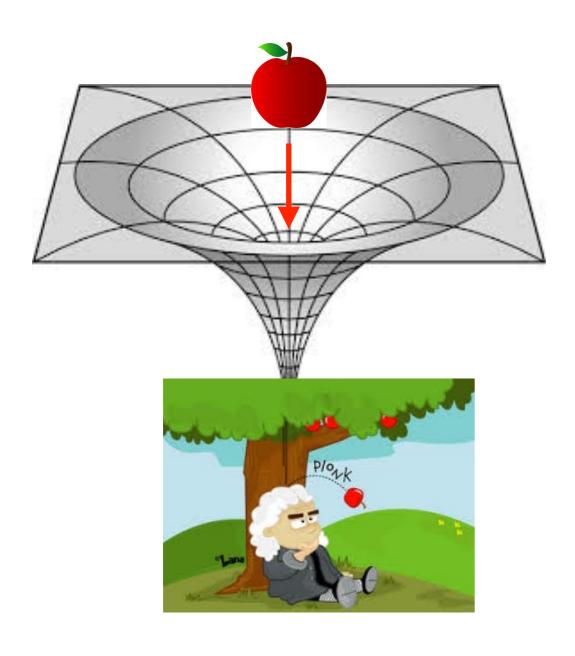
A PRIMER ON GRAVITATIONAL WAVES





$$G_{\mu
u} = rac{1}{m_p^2} T_{\mu
u}$$
 geometry matter

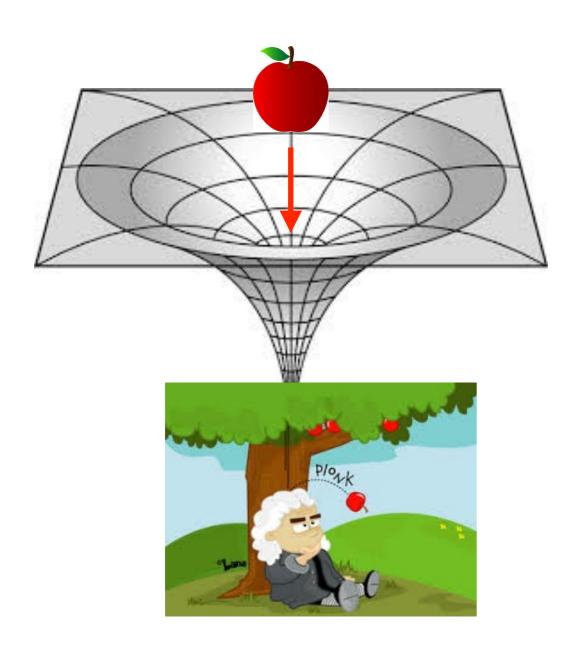
$$\left[m_p = (8\pi G)^{-1/2} = 2.44 \cdot 10^{18} \,\text{GeV}\right] \begin{array}{c} \text{Reduced} \\ \text{Planck mass} \end{array}$$



$$G_{\mu\nu} = \tfrac{1}{m_p^2} T_{\mu\nu}$$
 geometry matter

$$\left[m_p = (8\pi G)^{-1/2} = 2.44 \cdot 10^{18} \,\text{GeV}\right] \begin{array}{c} \text{Reduced} \\ \text{Planck mass} \end{array}$$

$$ds^2 = g_{\mu\nu}(x)dx^{\mu}dx^{\nu}$$

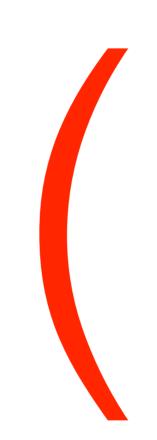


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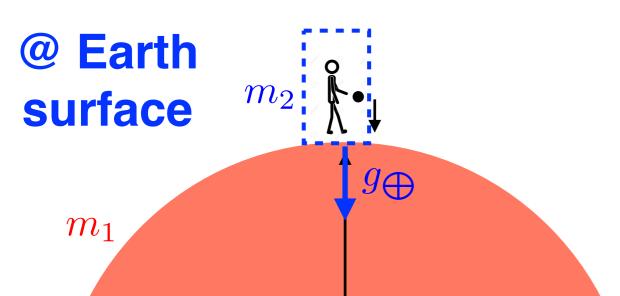
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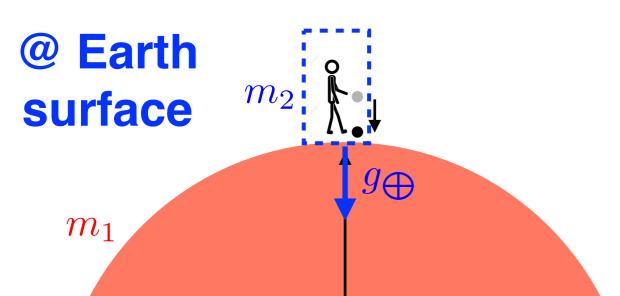
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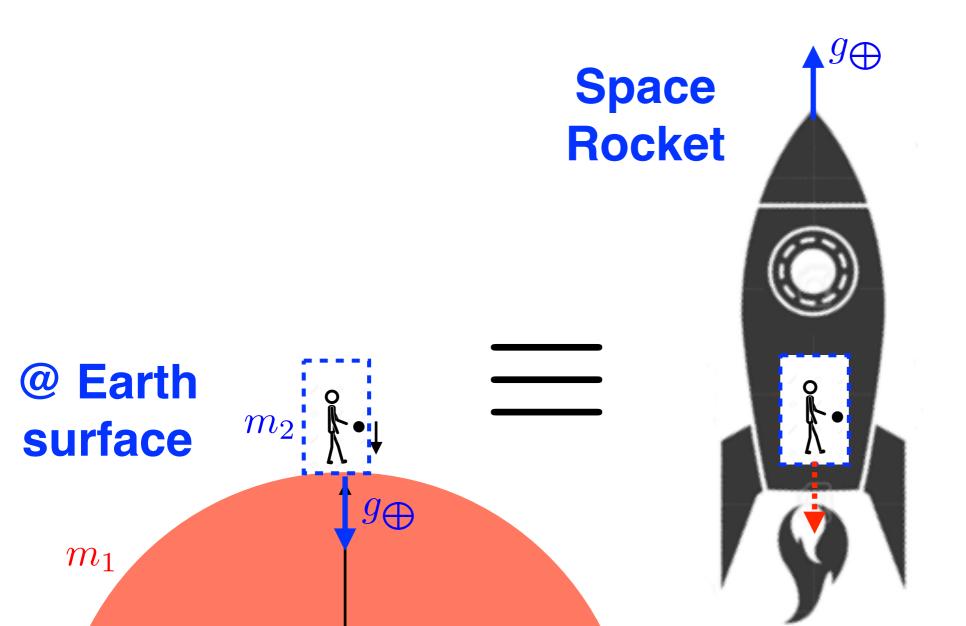
DIFF:
$$x^{\mu} \rightarrow x'^{\mu}(x)_{\text{etry}}$$

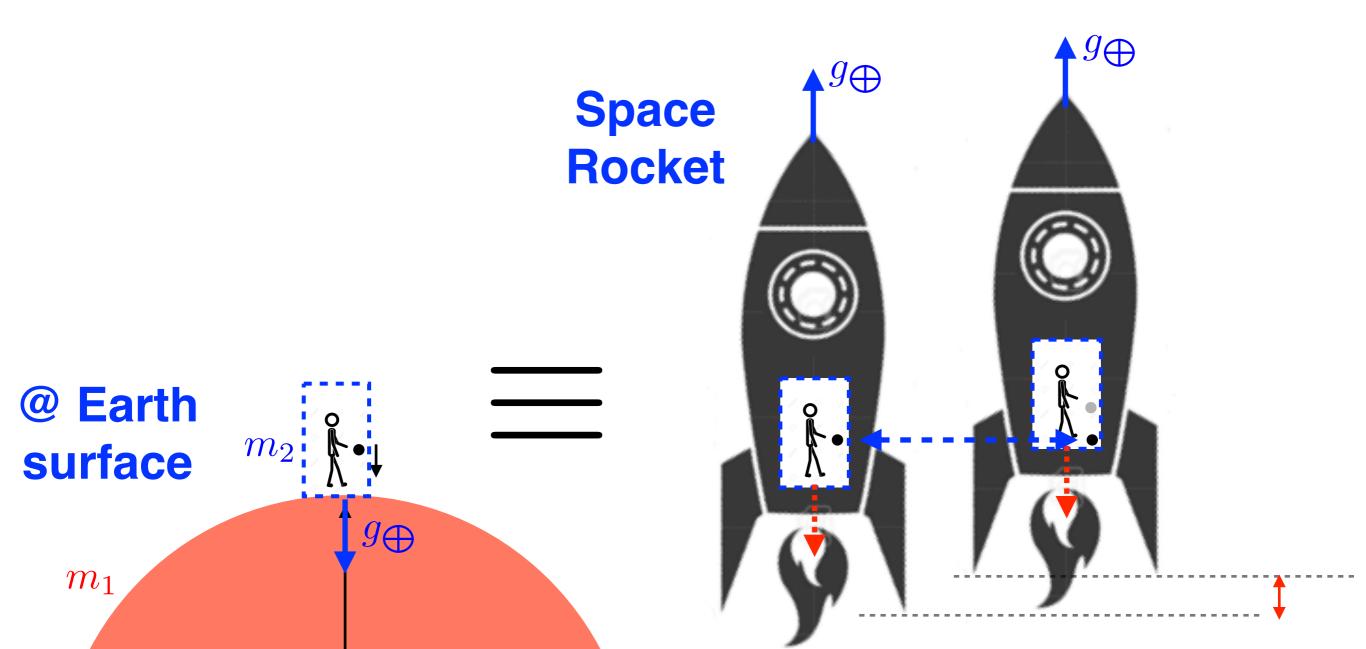


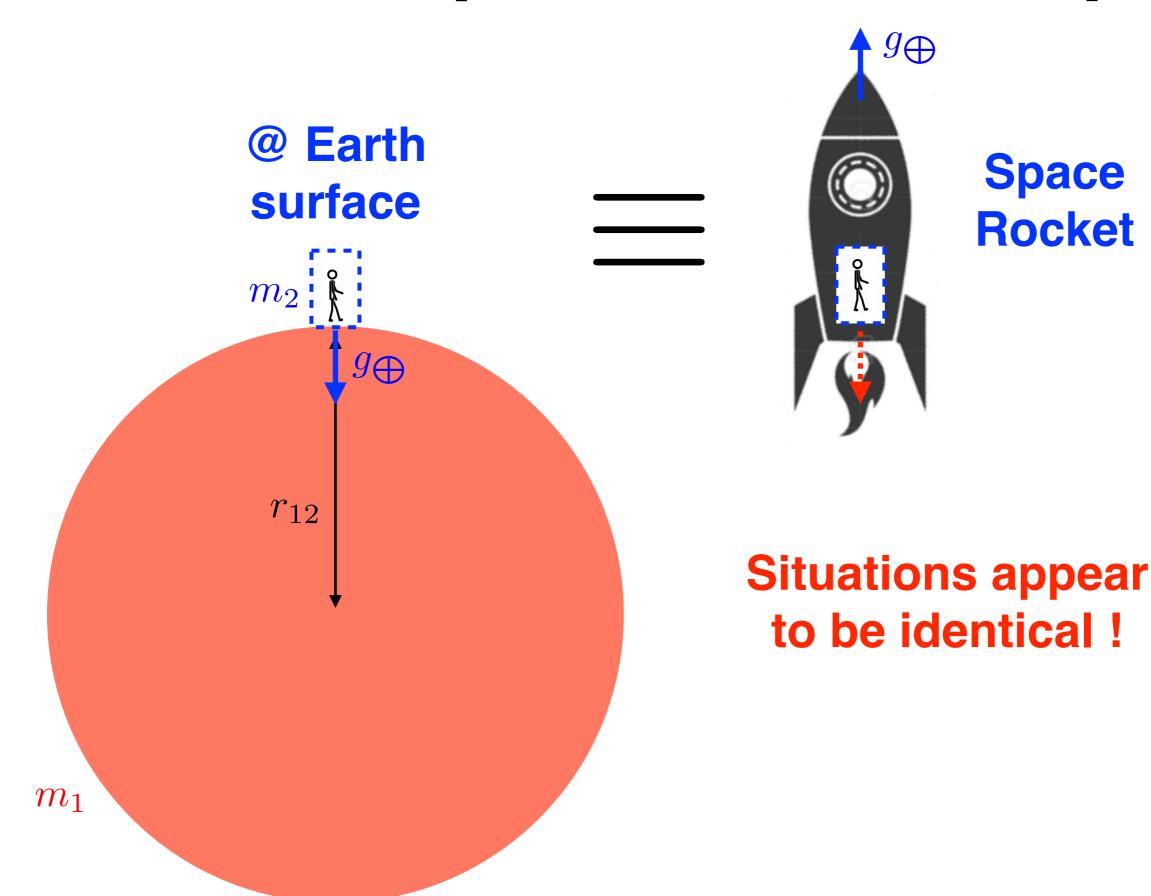
A PRIMER ON GENERAL RELATIVITY

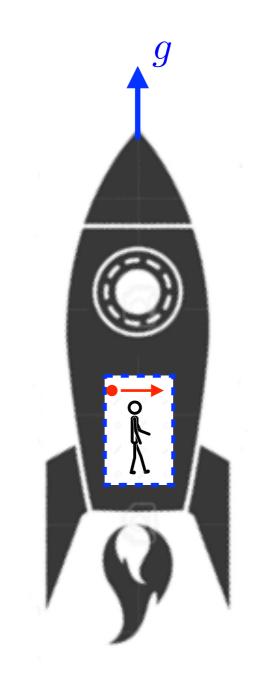




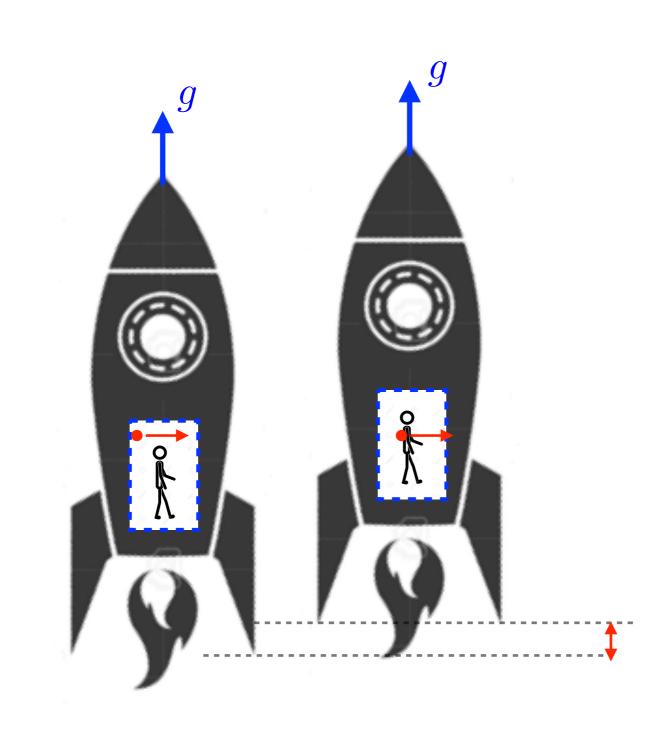




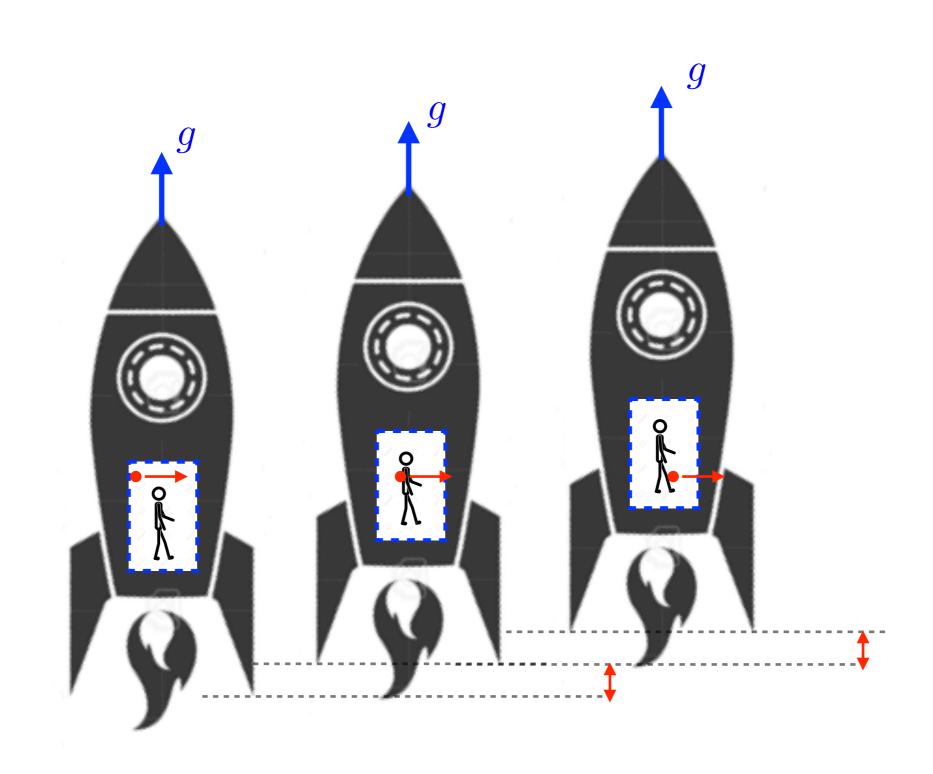




Laser beam

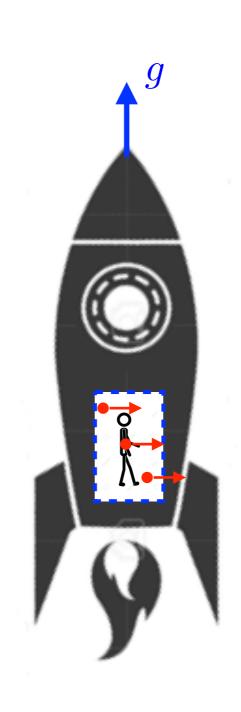


Laser beam

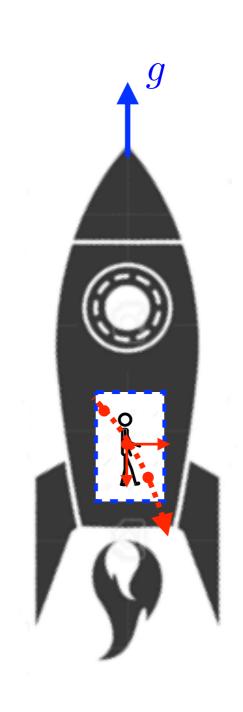


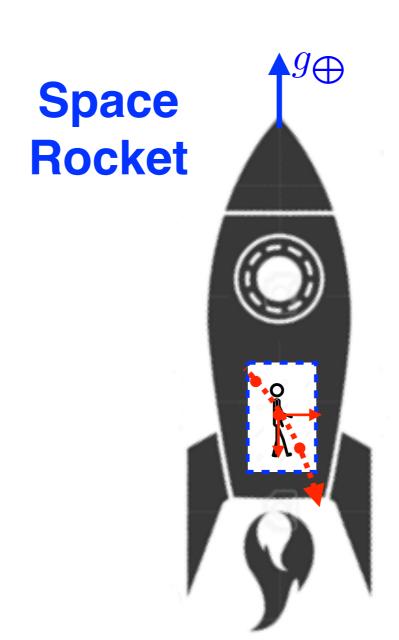
Laser beam

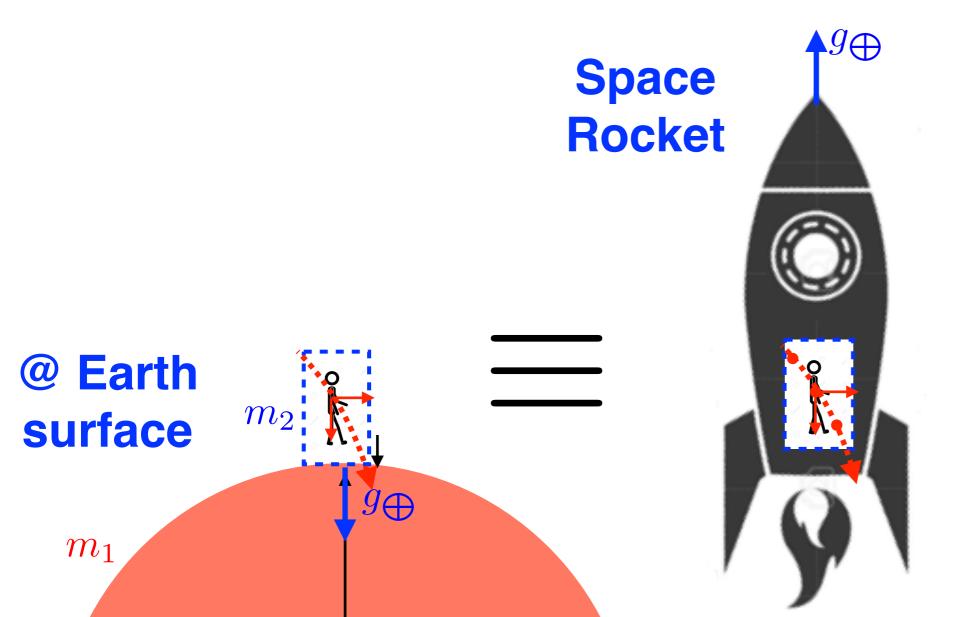
Laser beam



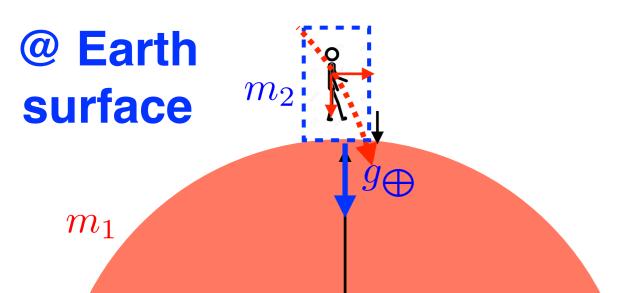
Laser beam





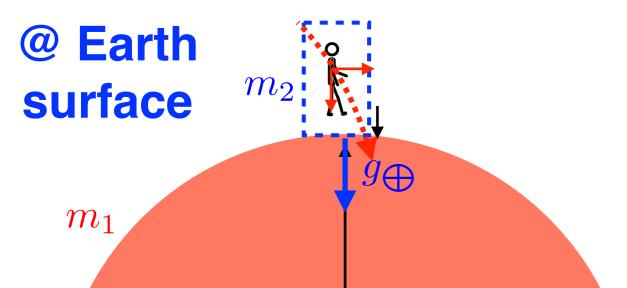


Gravitational field must bend light!



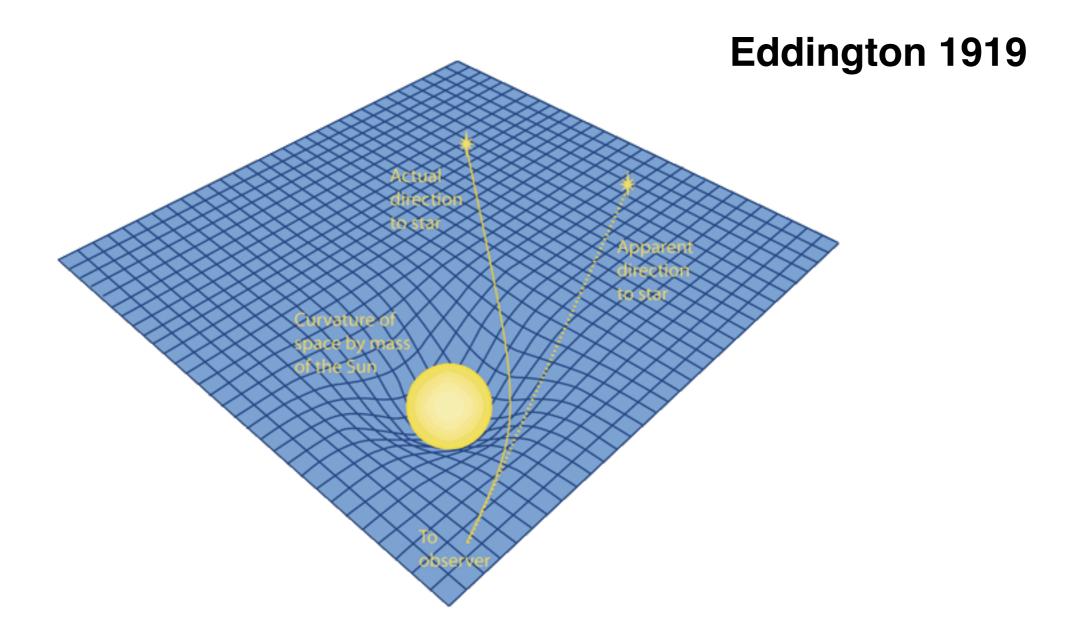
(Earth gravity too weak to observe the effect, but ...)

Gravitational field must bend light!



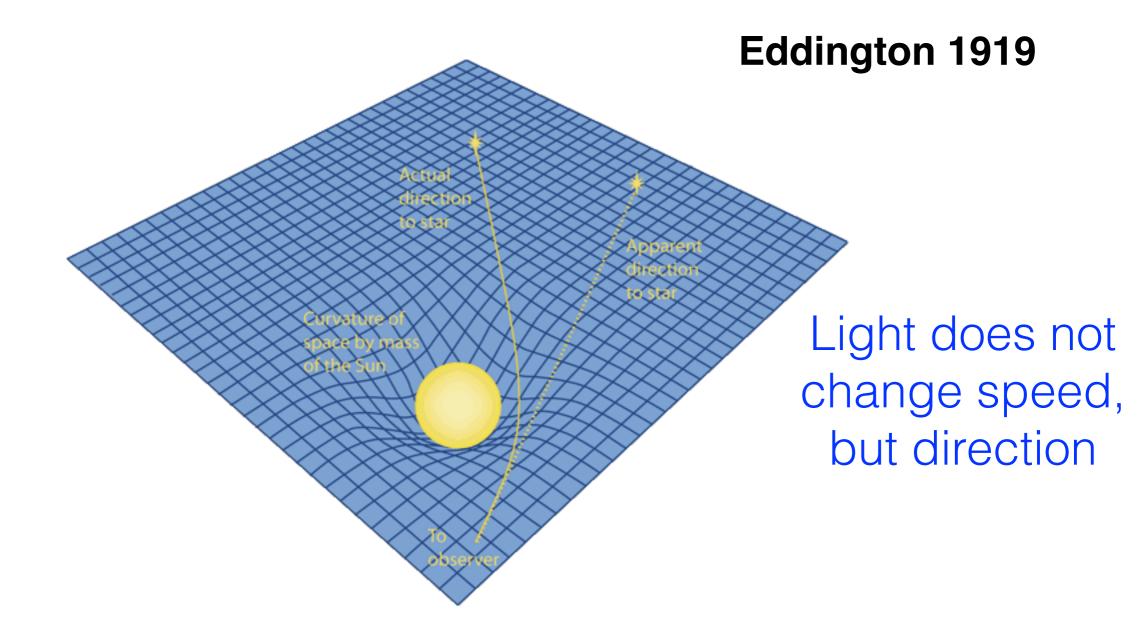
(Earth gravity too weak to observe the effect, but ...) →

the effect, but ...) e.g. the sun does a better job!



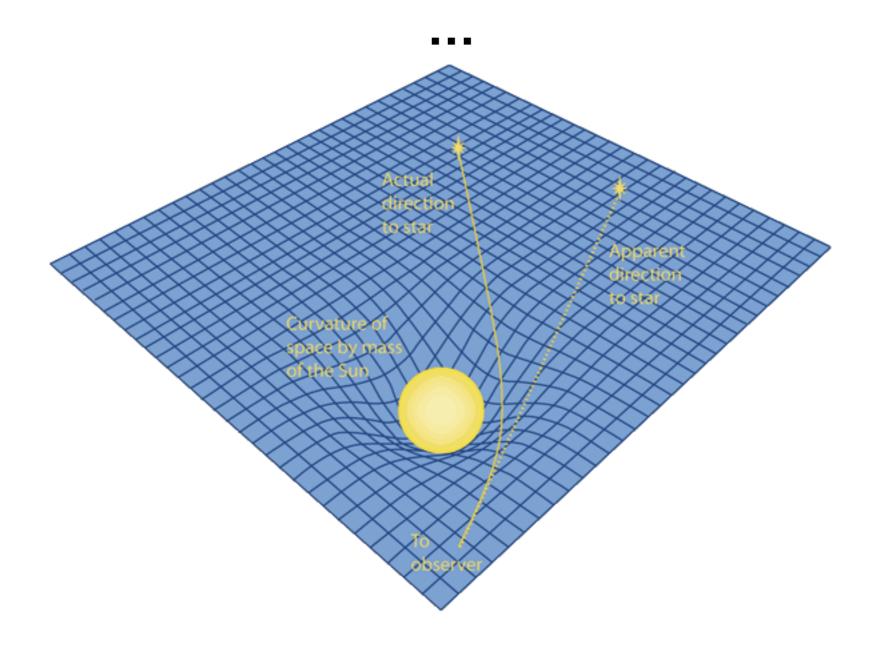
(Earth gravity too weak to observe the effect, but ...) →

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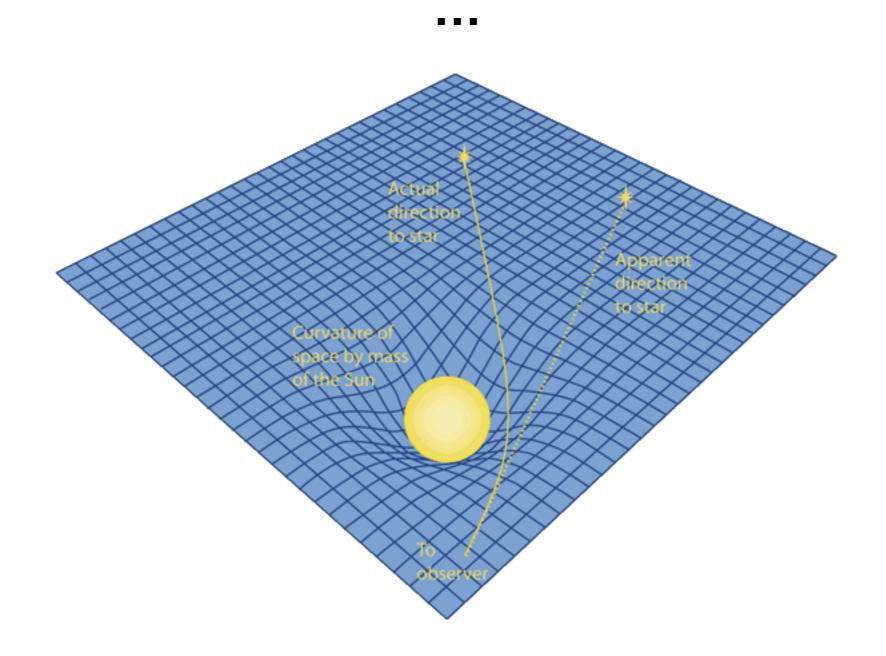
Einstein understood like this...

light bending, light red/blue-shifting, gravitational time dilation,

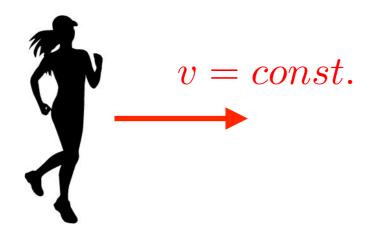


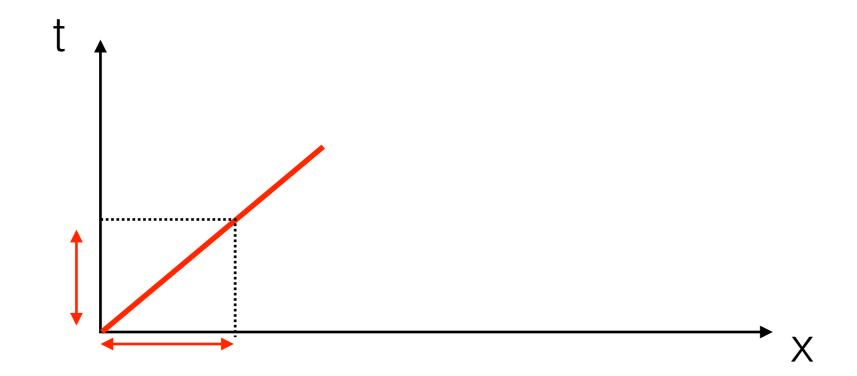
Einstein understood like this...

a mathematical formulation was needed!

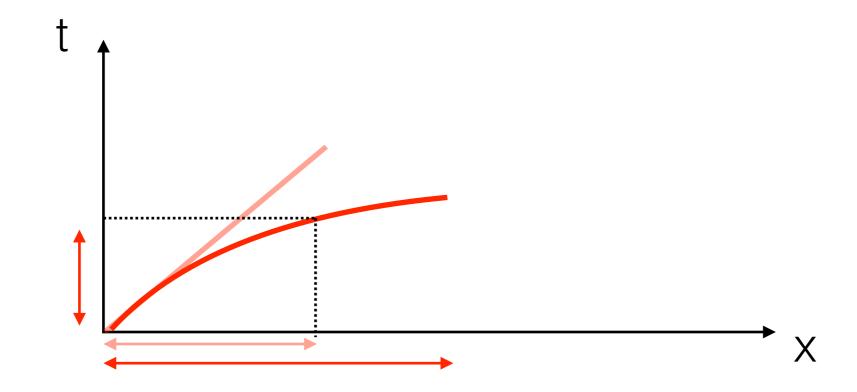


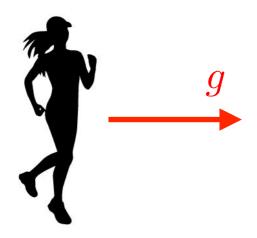
Mathematical formulation of General Relativity (GR)



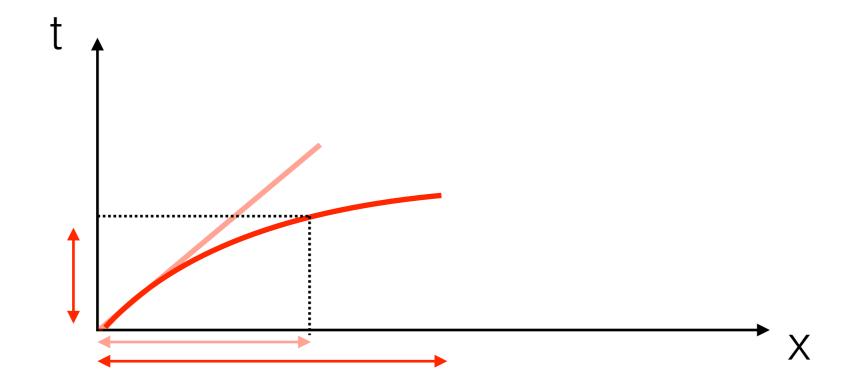






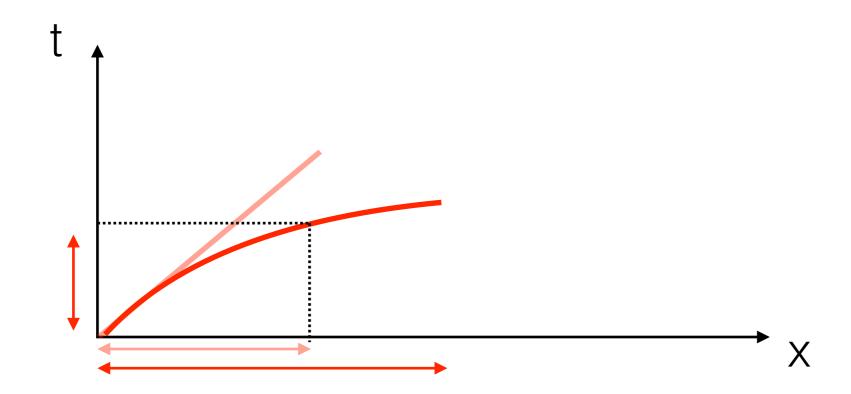


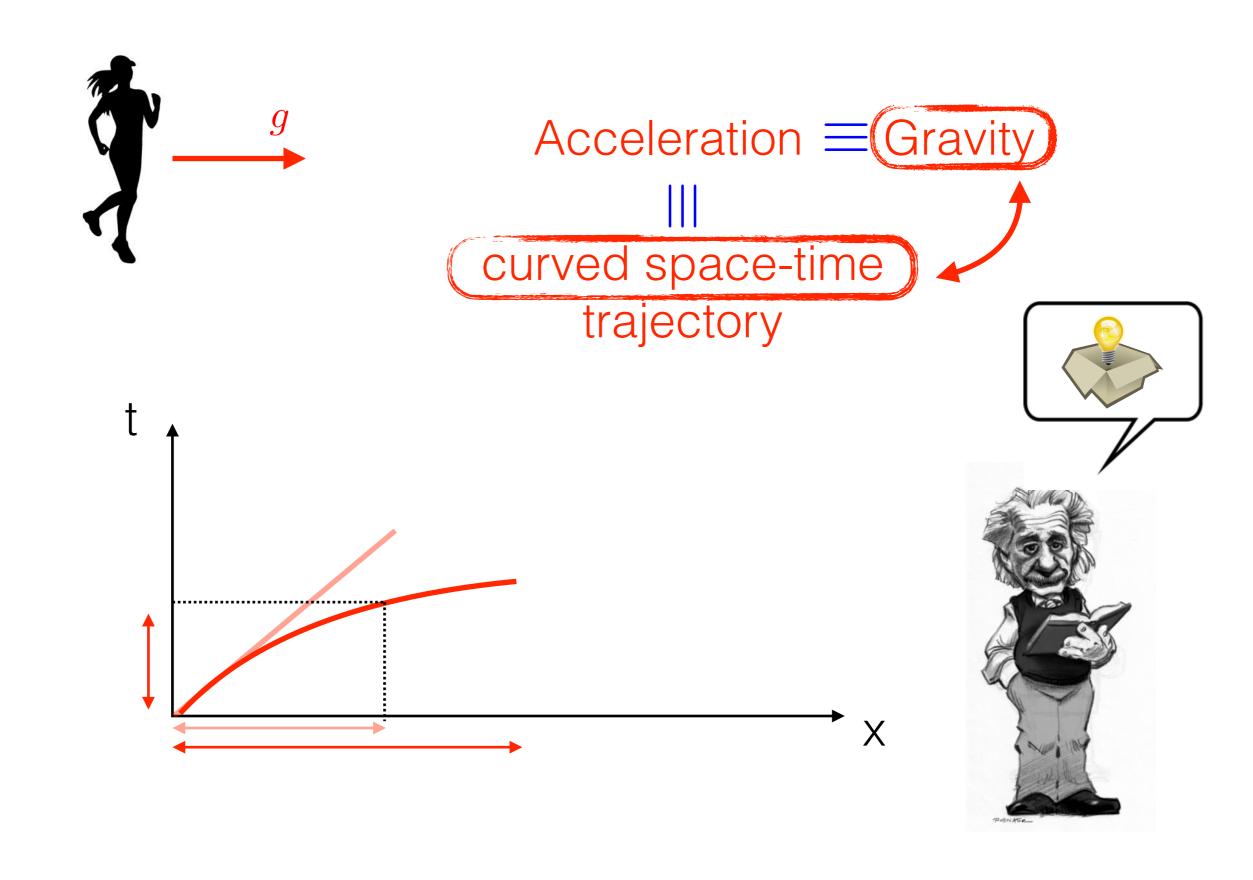
Acceleration = Gravity



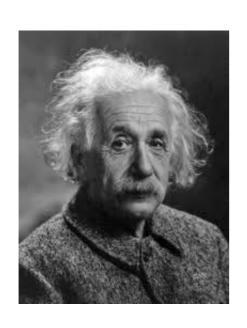


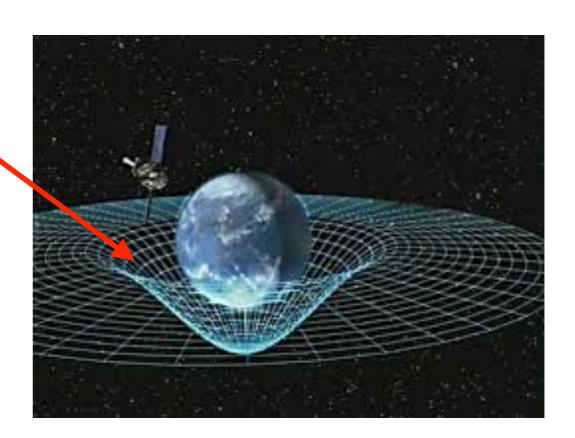
Acceleration = Gravity
|||
curved space-time
trajectory



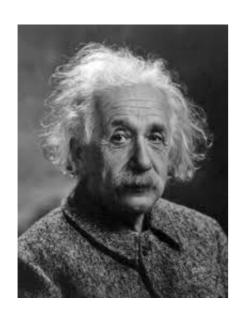


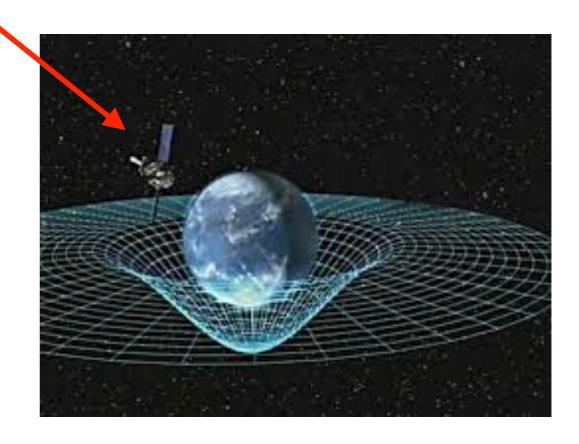
Presence of Matter (Energy/p) dictates 'Space-Time' Geometry

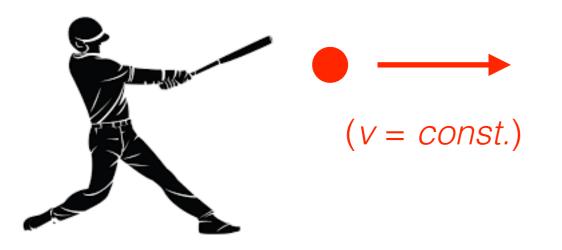


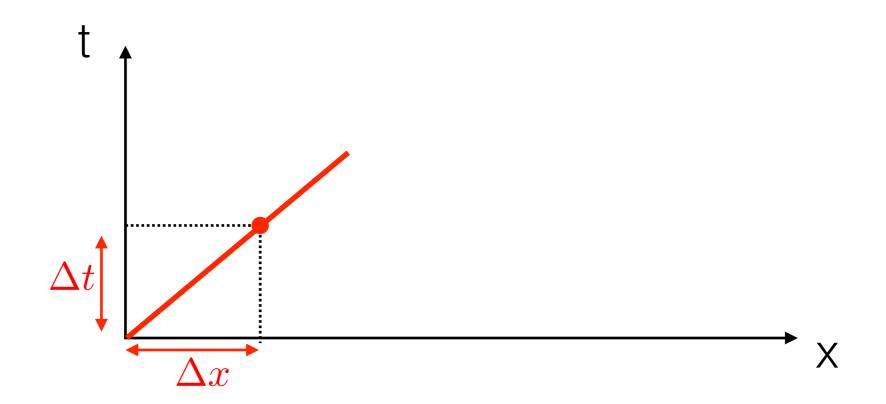


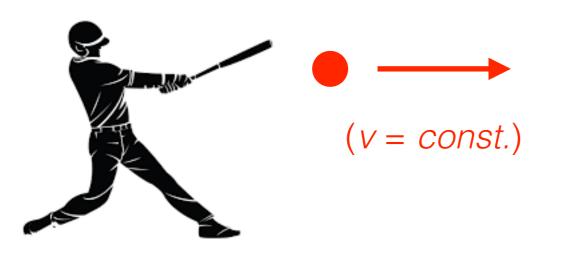
'Space-Time' Geometry dictates Movement of Matter

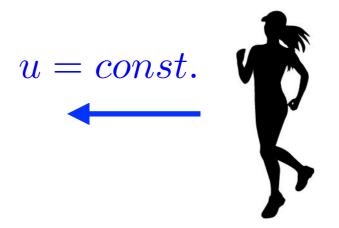


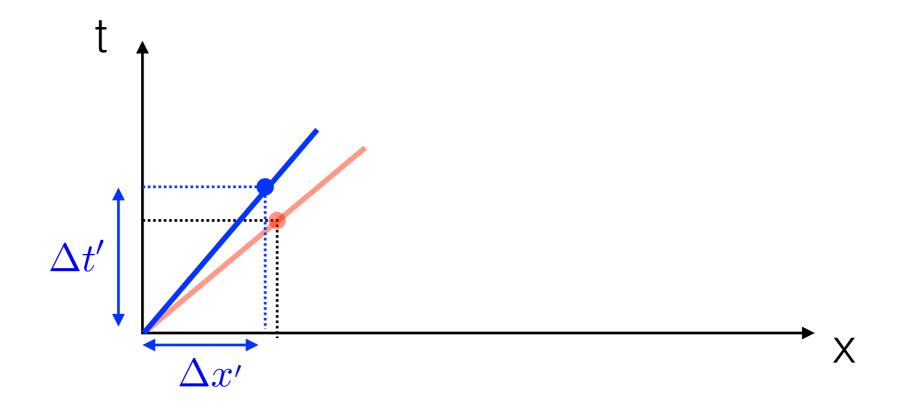


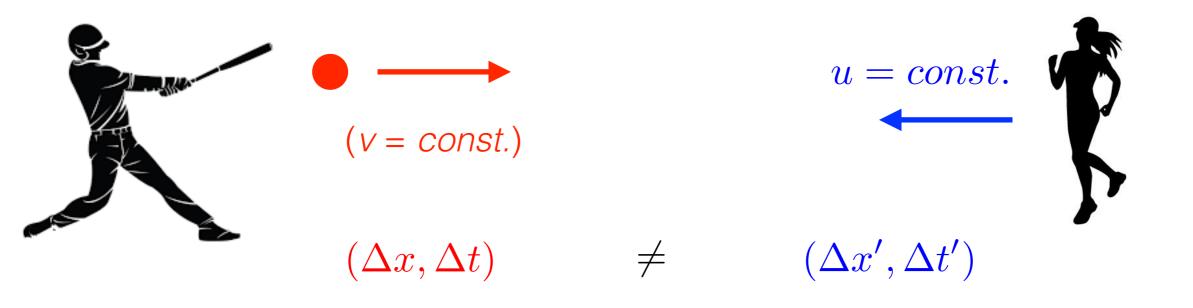


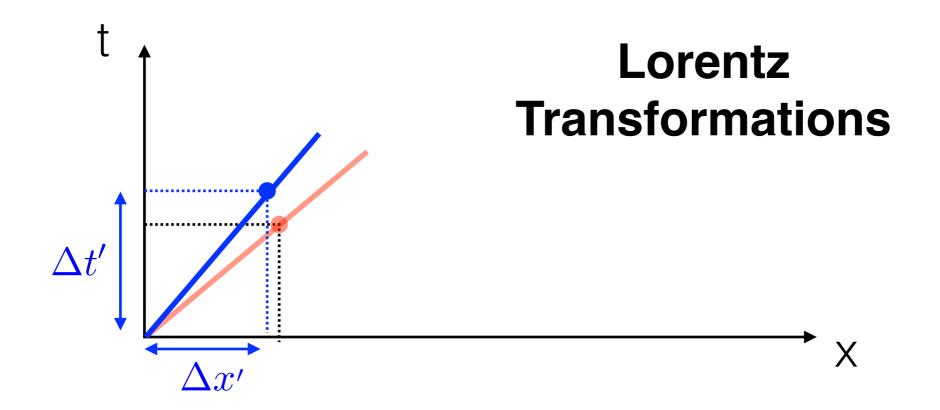


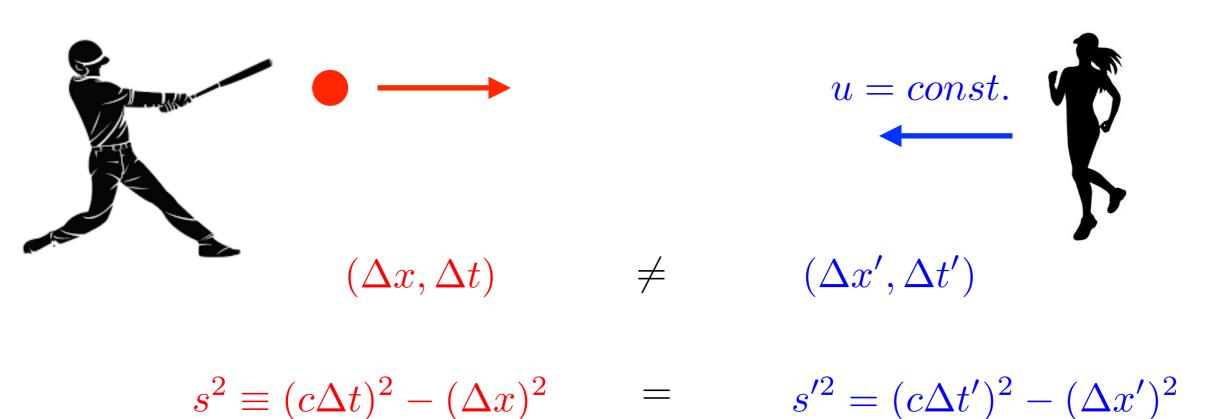














$$s^2 \equiv (c\Delta t)^2 - (\Delta x)^2 = s'^2 = (c\Delta t')^2 - (\Delta x')^2$$

$$ds^2 = c^2 dt^2 - \sum_j dx_j dx^j$$

Special Relativity

Space-time interval invariant

$$ds^2 = -c^2 dt^2 + \sum_j dx_j dx^j$$

Space-time invariant interval (**Special Relativity**)

$$ds^2 = -c^2 dt^2 + \sum_j dx_j dx^j \longrightarrow ds^2 = \eta_{\mu\nu} dx^\mu dx^\nu$$

Einstein convention

Summation over repeated indices

Space-time invariant interval (**Special Relativity**)

$$\eta \equiv \begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

$$ds^{2} = -c^{2}dt^{2} + \sum_{j} dx_{j} dx^{j} \longrightarrow ds^{2} = \eta_{\mu\nu} dx^{\mu} dx^{\nu}$$

Einstein convention

Summation over repeated indices

Space-time invariant interval (**Special Relativity**)

Minkowski Metric $\eta \equiv diag(-, +, +, +)$

$$ds^{2} = -c^{2}dt^{2} + \sum_{j} dx_{j} dx^{j} \longrightarrow ds^{2} = \eta_{\mu\nu} dx^{\mu} dx^{\nu}$$

Einstein convention
Summation

over repeated indices

Space-time invariant interval (**Special Relativity**)

Minkowski Metric $\eta \equiv \operatorname{diag}(-, +, +, +)$

$$ds^2 = g_{\mu\nu}(x)dx^{\mu}dx^{\nu}$$

$$ds^{2} = -c^{2}dt^{2} + \sum_{j} dx_{j} dx^{j} \longrightarrow ds^{2} = \eta_{\mu\nu} dx^{\mu} dx^{\nu}$$

Einstein convention
Summation
over repeated
indices

Space-time invariant interval (**Special Relativity**)

Minkowski Metric $\eta \equiv \operatorname{diag}(-, +, +, +)$

Space-time invariant interval (**General Relativity**)

$$g_{\mu\nu}(x)dx^{\mu}dx^{\nu} = g'_{\alpha\beta}(x')dx'^{\alpha}dx'^{\beta}$$

$$ds^2 = g_{\mu\nu}(x)dx^{\mu}dx^{\nu}$$



$$ds^{2} = -c^{2}dt^{2} + \sum_{j} dx_{j} dx^{j} \longrightarrow ds^{2} = \eta_{\mu\nu} dx^{\mu} dx^{\nu}$$

Einstein convention
Summation
over repeated
indices

Space-time invariant interval (**Special Relativity**)

Minkowski Metric $\eta \equiv \operatorname{diag}(-, +, +, +)$

Space-time invariant interval (**General Relativity**)

$$ds^2 = g_{\mu\nu}(x)dx^{\mu}dx^{\nu}$$

$$g_{\mu\nu}(x)dx^{\mu}dx^{\nu} = g'_{\alpha\beta}(x')dx'^{\alpha}dx'^{\beta}$$
$$g_{\mu\nu}(x) = \frac{\partial x'^{\alpha}}{\partial x^{\mu}} \frac{\partial x'^{\beta}}{\partial x^{\nu}} g'_{\alpha\beta}(x')$$

General Relativity: Generalisation of Special Relativity

* Equivalence Principle

Geodesic motion

$$\frac{d^2x^{\mu}}{ds^2} + \Gamma^{\mu}_{\alpha\beta}[g_{**}] \frac{dx^{\alpha}}{ds} \frac{dx^{\beta}}{ds} = 0$$

$$\Gamma^{\mu}_{\alpha\beta} = \frac{1}{2} g^{\mu\lambda} \left(g_{\lambda\alpha,\beta} + g_{\lambda\beta,\alpha} - g_{\alpha\beta,\lambda} \right)$$

Christoffel Symbol

General Relativity: Generalisation of Special Relativity

* Equivalence Principle

Geodesic motion

$$\frac{d^2x^{\mu}}{ds^2} + \Gamma^{\mu}_{\alpha\beta}[g_{**}] \frac{dx^{\alpha}}{ds} \frac{dx^{\beta}}{ds} = 0$$

* Principle of Relativity

$$x'^{\mu} = x'^{\mu}(\{x^{\alpha}\}) \qquad \text{change of }$$

Arbitrary coordinates

$$g_{\mu\nu}(x)dx^{\mu}dx^{\nu} = g'_{\alpha\beta}(x')dx'^{\alpha}dx'^{\beta}$$

$$g_{\mu\nu}(x) = \frac{\partial x'^{\alpha}}{\partial x^{\mu}} \frac{\partial x'^{\beta}}{\partial x^{\nu}} g'_{\alpha\beta}(x')$$
;

General Relativity: Generalisation of Special Relativity

* Equivalence Principle

Geodesic motion

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 change of

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$$g_{\mu\nu}(x)dx^{\mu}dx^{\nu} = g'_{\alpha\beta}(x')dx'^{\alpha}dx'^{\beta}$$

$$g_{\mu\nu}(x) = \frac{\partial x'^{\alpha}}{\partial x^{\mu}} \frac{\partial x'^{\beta}}{\partial x^{\nu}} g'_{\alpha\beta}(x')$$

$$; G_{\mu\nu}[g_{**}] \equiv R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \frac{1}{m_p^2} T_{\mu\nu}$$

$$\text{space-time geometry}$$

$$\text{matter (energy/p)}$$

$$x'^{\mu} = x'^{\mu}(\{x^{\alpha}\}) \quad \text{Arbitrary} \\ \text{coordinates} \qquad ; \qquad g_{\mu\nu}(x)dx^{\mu}dx^{\nu} = g'_{\alpha\beta}(x')dx'^{\alpha}dx'^{\beta} \\ g_{\mu\nu}(x) = \frac{\partial x'^{\alpha}}{\partial x^{\mu}} \frac{\partial x'^{\beta}}{\partial x^{\nu}} g'_{\alpha\beta}(x')$$

$$x'^{\mu} = x'^{\mu}(\{x^{\alpha}\}) \quad \text{Arbitrary} \\ \text{coordinates} \qquad ; \qquad g_{\mu\nu}(x)dx^{\mu}dx^{\nu} = g'_{\alpha\beta}(x')dx'^{\alpha}dx'^{\beta} \\ \vdots \\ g_{\mu\nu}(x) = \frac{\partial x'^{\alpha}}{\partial x^{\mu}} \frac{\partial x'^{\beta}}{\partial x^{\nu}} g'_{\alpha\beta}(x')$$

$$G_{\mu\nu}[g_{**}] \equiv R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \frac{1}{m_p^2}T_{\mu\nu} \quad ; \quad \frac{8\pi G}{c^4} \equiv \frac{1}{m_p^2} \quad ; \quad m_p = 2.44 \cdot 10^{18} \; \mathrm{GeV}$$
 space-time matter (energy/p)

$$x'^{\mu} = x'^{\mu}(\{x^{\alpha}\}) \quad \text{Arbitrary} \\ \text{coordinates} \qquad ; \qquad g_{\mu\nu}(x) dx^{\mu} dx^{\nu} = g'_{\alpha\beta}(x') dx'^{\alpha} dx'^{\beta} \\ \text{coordinates} \qquad ; \qquad g_{\mu\nu}(x) = \frac{\partial x'^{\alpha}}{\partial x^{\mu}} \frac{\partial x'^{\beta}}{\partial x^{\nu}} g'_{\alpha\beta}(x')$$

$$G_{\mu\nu}[g_{**}] \equiv R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \frac{1}{m_p^2}T_{\mu\nu} \quad ; \quad \frac{8\pi G}{c^4} \equiv \frac{1}{m_p^2} \quad ; \quad m_p = 2.44 \cdot 10^{18} \; \mathrm{GeV}$$
 space-time geometry matter (energy/p)

$$R_{\alpha\beta} = \Gamma^{\mu}_{\ \alpha\beta,\mu} - \Gamma^{\mu}_{\ \alpha\mu,\beta} + \Gamma^{\mu}_{\ \lambda\mu} \Gamma^{\lambda}_{\ \alpha\beta} - \Gamma^{\mu}_{\ \lambda\beta} \Gamma^{\lambda}_{\ \alpha\mu} \qquad \text{Ricci tensor}$$

$$\Gamma^{\mu}_{\ \alpha\beta} = \frac{1}{2} g^{\mu\lambda} \left(g_{\lambda\alpha,\beta} + g_{\lambda\beta,\alpha} - g_{\alpha\beta,\lambda} \right) \sim (metric)^2 \qquad \text{Christoffel Symbol}$$

General Relativity (GR)

$$G_{\mu
u} = rac{1}{m_p^2} T_{\mu
u}$$
 geometry matter

metric
$$G_{\mu\nu} \equiv \mathcal{D}[g_{\alpha\beta}] = m_p^{-2} T_{\mu\nu}({\rm Matt,Rad,Top.Defects,DarkEnergy,...})$$
 source 2nd order, non-Linear

General Relativity (GR)

$$G_{\mu
u} = rac{1}{m_p^2} T_{\mu
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 geometry matter

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 source 2nd order, non-Linear



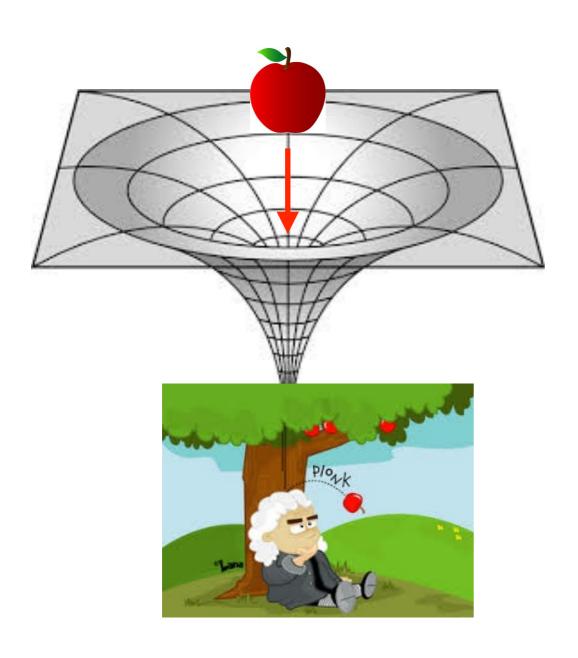
Extremely difficult to solve!

END of digression on GENERAL RELATIVITY



Let's continue with PRIMER ON GRAVITATIONAL WAVES

General Relativity (GR)



$$G_{\mu\nu} = \tfrac{1}{m_p^2} T_{\mu\nu}$$
 geometry matter

$$\left[m_p = (8\pi G)^{-1/2} = 2.44 \cdot 10^{18} \,\text{GeV}\right] \begin{array}{c} \text{Reduced} \\ \text{Planck mass} \end{array}$$

$$ds^2 = g_{\mu\nu}(x)dx^{\mu}dx^{\nu}$$

DIFF:
$$x^{\mu} \rightarrow x'^{\mu}(x)_{\text{etry}}$$

General Relativity (GR)

$$G_{\mu
u} = rac{1}{m_p^2} T_{\mu
u}$$
 geometry matter

metric
$$G_{\mu\nu} \equiv \mathcal{D}[g_{\alpha\beta}] = m_p^{-2} T_{\mu\nu}({\rm Matt,Rad,Top.Defects,DarkEnergy,...})$$
 source 2nd order, non-Linear

General Relativity (GR)

$$G_{\mu
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u}$$
 geometry matter

metric
$$G_{\mu\nu} \equiv \mathcal{D}[g_{\alpha\beta}] = m_p^{-2} T_{\mu\nu}({\rm Matt,Rad,Top.Defects,DarkEnergy},...)$$
 source 2nd order, non-Linear

How do we define GWs?

General Relativity (GR)

$$G_{\mu
u} = rac{1}{m_p^2} T_{\mu
u}$$
 geometry matter

metric
$$G_{\mu\nu} \equiv \mathcal{D}[g_{\alpha\beta}] = m_p^{-2} T_{\mu\nu}(\mathrm{Matt,Rad,Top.Defects,DarkEnergy,...})$$
 source expand in perturbations

$$g_{\alpha\beta} = \bar{g}_{\alpha\beta} + \delta g_{\alpha\beta}$$

How do we define GWs?

General Relativity (GR)

$$G_{\mu
u} = rac{1}{m_p^2} T_{\mu
u}$$
 geometry matter

metric
$$G_{\mu\nu} \equiv \mathcal{D}[g_{\alpha\beta}] = m_p^{-2} T_{\mu\nu}({\rm Matt,Rad,Top.Defects,DarkEnergy,...})$$
 source of GWs

expand in perturbations

$$g_{\alpha\beta} = \bar{g}_{\alpha\beta} + \delta g_{\alpha\beta}$$

How do we define GWs?

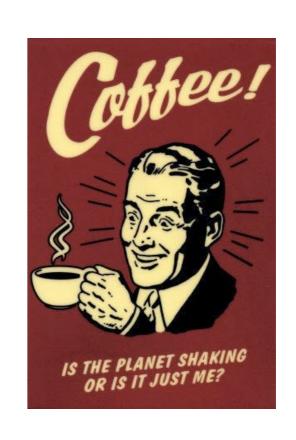
$$g_{\alpha\beta} = \bar{g}_{\alpha\beta} + \delta g_{\alpha\beta}$$

Let's continue this approach...

$$g_{\alpha\beta} = \bar{g}_{\alpha\beta} + \delta g_{\alpha\beta}$$

Let's continue this approach...

But not before you load yourselves with coffee ('cause you are gonna need it)



Definition of GWs 1st approach

Minkowski
$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}(x) \ (|h_{\mu\nu}| \ll 1)$$



Minkowski
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Minkowski
$$g_{\mu\nu}=\eta_{\mu\nu}+h_{\mu\nu}(x) \ \ \begin{array}{l} \text{fixed} \\ \text{frame} \\ (\ |h_{\mu\nu}|\ll 1\) \end{array}$$

Minkowski
$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}(x) \ \ \text{fixed} \ \ (|h_{\mu\nu}| \ll 1)$$

DIFF:
$$x^{\mu} \rightarrow x'^{\mu}(x)$$
symmetry?

Minkowski
$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}(x) \ \ \text{fixed} \ \ (|h_{\mu\nu}| \ll 1)$$

DIFF:
$$x^{\mu} \rightarrow x'^{\mu}(x)$$

Minkowski
$$g_{\mu\nu}=\eta_{\mu\nu}+h_{\mu\nu}(x) \ \ {\it fixed} \ \ (|h_{\mu\nu}|\ll 1)$$

DIFF:
$$x^{\mu} \rightarrow x'^{\mu}(x)$$

$$x^{\mu} \rightarrow x'^{\mu} = x^{\mu} + \xi^{\mu}(x)$$
 residual $(|\partial_{\mu}\xi_{\nu}(x)| \lesssim |h_{\mu\nu}|)$ symm.

$$h_{\mu\nu}(x) \to h'_{\mu\nu}(x') = h_{\mu\nu}(x) - \partial_{(\mu}\xi_{\nu)}$$

Minkowski
$$g_{\mu\nu}=\eta_{\mu\nu}+h_{\mu\nu}(x) \ \ \begin{array}{l} \text{fixed} \\ \text{frame} \\ (\ |h_{\mu\nu}|\ll 1\) \end{array}$$

DIFF:
$$x^{\mu} \not\rightarrow x'^{\mu}(x)$$

$$x^{\mu} \rightarrow x'^{\mu} = x^{\mu} + \xi^{\mu}(x)$$
 residual $(|\partial_{\mu}\xi_{\nu}(x)| \lesssim |h_{\mu\nu}|)$ symm.

$$h_{\mu\nu}(x) \to h'_{\mu\nu}(x') = h_{\mu\nu}(x) - \partial_{(\mu}\xi_{\nu)}$$

Notation:
$$\begin{cases} \partial_{(\mu}\xi_{\nu)} \equiv \partial_{\mu}\xi_{\nu} + \partial_{\nu}\xi_{\mu} \\ \partial_{[\mu}\xi_{\nu]} \equiv \partial_{\mu}\xi_{\nu} - \partial_{\nu}\xi_{\mu} \end{cases}$$

Minkowski
$$g_{\mu\nu}=\eta_{\mu\nu}+h_{\mu\nu}(x) \ \ {\it fixed} \ \ (|h_{\mu\nu}|\ll 1)$$

DIFF:
$$x^{\mu} \rightarrow x'^{\mu}(x)$$

$$x^{\mu} \rightarrow x'^{\mu} = x^{\mu} + \xi^{\mu}(x)$$
 residual $(|\partial_{\mu}\xi_{\nu}(x)| \lesssim |h_{\mu\nu}|)$ symm.

$$h_{\mu\nu}(x) \to h'_{\mu\nu}(x') = h_{\mu\nu}(x) - \partial_{(\mu}\xi_{\nu)}$$

1st approach to GWs

Minkowski
$$g_{\mu\nu}=\eta_{\mu\nu}+h_{\mu\nu}(x) \ \ \begin{array}{l} \text{fixed} \\ \text{frame} \\ (\,|h_{\mu\nu}|\ll 1\,) \end{array}$$

Let's expand Einstein Equations!

1st approach to GWs

Minkowski
$$g_{\mu\nu}=\eta_{\mu\nu}+h_{\mu\nu}(x) \ \ {\it fixed} \ \ (|h_{\mu\nu}|\ll 1\)$$

Trace-reversed

$$\bar{h}_{\mu\nu} \equiv h_{\mu\nu} - \frac{1}{2}h\eta_{\mu\nu}$$

1st approach to GWs

Minkowski
$$g_{\mu\nu}= \begin{matrix} \uparrow \\ \eta_{\mu\nu} + h_{\mu\nu}(x) \end{matrix} \quad \begin{array}{l} \text{fixed} \\ \text{frame} \end{matrix} \quad (|h_{\mu\nu}| \ll 1) \end{matrix}$$

Trace-reversed

$$\bar{h}_{\mu\nu} \equiv h_{\mu\nu} - \frac{1}{2}h\eta_{\mu\nu} \qquad \qquad \qquad \partial^{\alpha}\partial_{\alpha}\bar{h}_{\mu\nu} + \eta_{\mu\nu}\partial^{\alpha}\partial^{\beta}\bar{h}_{\alpha\beta} - \partial^{\alpha}\partial_{(\mu}\bar{h}_{\alpha\nu)} = -\frac{2}{m_{p}^{2}}T_{\mu\nu}$$

1st approach to GWs

Minkowski
$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}(x) \begin{tabular}{l} {\bf fixed} \\ {\bf frame} \\ (\ |h_{\mu\nu}| \ll 1 \) \end{tabular}$$

Trace-reversed

$$\bar{h}_{\mu\nu} \equiv h_{\mu\nu} - \frac{1}{2}h\eta_{\mu\nu} \longrightarrow \partial^{\alpha}\partial_{\alpha}\bar{h}_{\mu\nu} + \eta_{\mu\nu}\partial^{\alpha}\partial^{\beta}\bar{h}_{\alpha\beta} - \partial^{\alpha}\partial_{(\mu}\bar{h}_{\alpha\nu)} = -\frac{2}{m_{p}^{2}}T_{\mu\nu}$$

1st approach to GWs

Minkowski
$$g_{\mu\nu}= \begin{matrix} \uparrow \\ \eta_{\mu\nu} + h_{\mu\nu}(x) \end{matrix} \quad \begin{array}{l} \text{fixed} \\ \text{frame} \end{matrix} \quad (|h_{\mu\nu}| \ll 1) \end{matrix}$$

Trace-reversed

$$\bar{h}_{\mu\nu} \equiv h_{\mu\nu} - \frac{1}{2}h\eta_{\mu\nu} \longrightarrow \partial^{\alpha}\partial_{\alpha}\bar{h}_{\mu\nu} + \eta_{\mu\nu}\partial^{\alpha}\partial^{\beta}\bar{h}_{\alpha\beta} - \partial^{\alpha}\partial_{(\mu}\bar{h}_{\alpha\nu)} = -\frac{2}{m_p^2}T_{\mu\nu}$$

residual
$$x^{\mu} \rightarrow x'^{\mu} = x^{\mu} + \xi^{\mu}(x)$$
 symm. $(|\partial_{\mu}\xi_{\nu}(x)| \lesssim |h_{\mu\nu}|)$

1st approach to GWs

Minkowski
$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}(x) \begin{tabular}{l} {\bf fixed} \\ {\bf frame} \\ (\ |h_{\mu\nu}| \ll 1 \) \end{tabular}$$

Trace-reversed

$$\bar{h}_{\mu\nu} \equiv h_{\mu\nu} - \frac{1}{2}h\eta_{\mu\nu} \longrightarrow \partial^{\alpha}\partial_{\alpha}\bar{h}_{\mu\nu} + \eta_{\mu\nu}\partial^{\alpha}\partial^{\beta}\bar{h}_{\alpha\beta} - \partial^{\alpha}\partial_{(\mu}\bar{h}_{\alpha\nu)} = -\frac{2}{m_p^2}T_{\mu\nu}$$

residual

$$\partial^{\nu}\bar{h}_{\mu\nu} = 0$$

Lorentz gauge

1st approach to GWs

Minkowski
$$g_{\mu\nu}= \begin{matrix} \uparrow \\ \eta_{\mu\nu} + h_{\mu\nu}(x) \end{matrix} \quad \begin{array}{l} \text{fixed} \\ \text{frame} \end{matrix} \quad (|h_{\mu\nu}| \ll 1) \end{matrix}$$

Trace-reversed

$$\bar{h}_{\mu\nu} \equiv h_{\mu\nu} - \frac{1}{2}h\eta_{\mu\nu}$$

$$\partial^{\alpha}\partial_{\alpha}\bar{h}_{\mu\nu} + \eta_{\mu\nu}\partial^{\alpha}\partial^{\beta}\bar{h}_{\alpha\beta} - \partial^{\alpha}\partial_{(\mu}\bar{h}_{\alpha\nu)} = -\frac{2}{m_{p}^{2}}T_{\mu\nu}$$

 $\mathcal{O}(h_{**})$ Einstein tensor expanded

residual
$$\partial^{\nu} \bar{h}_{\mu\nu} = 0$$

Lorentz gauge

1st approach to GWs

Minkowski
$$g_{\mu\nu}=\eta_{\mu\nu}+h_{\mu\nu}(x) \ \ \begin{array}{c} \text{fixed} \\ \text{frame} \\ (\ |h_{\mu\nu}|\ll 1\) \end{array}$$

Trace-reversed

$$\bar{h}_{\mu\nu} \equiv h_{\mu\nu} - \frac{1}{2}h\eta_{\mu\nu} \qquad \qquad \partial$$

$$\partial^{\alpha}\partial_{\alpha}\bar{h}_{\mu\nu} + \eta_{\mu\nu}\partial^{\alpha}\partial^{\beta}\bar{h}_{\alpha\beta} - \partial^{\alpha}\partial_{(\mu}\bar{h}_{\alpha\nu)} = -\frac{2}{m_{p}^{2}}T_{\mu\nu}$$

 $\mathcal{O}(h_{**})$ Einstein tensor expanded

residual
$$\partial^{\nu} \bar{h}_{\mu\nu} = 0$$

Lorentz gauge

$$x^{\mu} \to x'^{\mu} = x^{\mu} + \xi^{\mu}(x)$$

$$h_{\mu\nu}(x) \to h'_{\mu\nu}(x') = h_{\mu\nu}(x) - \partial_{[\mu}\xi_{\nu]}$$

1st approach to GWs

Minkowski
$$g_{\mu\nu}=\eta_{\mu\nu}+h_{\mu\nu}(x) \ \ {\it fixed} \ \ (|h_{\mu\nu}|\ll 1\)$$

Trace-reversed

$$\bar{h}_{\mu\nu} \equiv h_{\mu\nu} - \frac{1}{2}h\eta_{\mu\nu} \qquad \qquad \partial^{\alpha}\partial_{\alpha}\bar{h}_{\mu\nu} + \eta_{\mu\nu}\partial^{\alpha}\partial^{\beta}\bar{h}_{\alpha\beta} - \partial^{\alpha}\partial_{(\mu}\bar{h}_{\alpha\nu)} = -\frac{2}{m_{p}^{2}}T_{\mu\nu}$$

 $\mathcal{O}(h_{**})$ Einstein tensor expanded

residual
$$\partial^{\nu} \bar{h}_{\mu\nu} = 0$$

Lorentz gauge

$$\partial^{\prime\mu}\bar{h}_{\mu\nu}'(x') = \underbrace{\partial^{\mu}\bar{h}_{\mu\nu}(x)}_{\equiv f(x) \neq 0} - \Box\xi_{\nu}$$

1st approach to GWs

Minkowski
$$g_{\mu\nu}=\eta_{\mu\nu}+h_{\mu\nu}(x) \ \ \begin{array}{l} \text{fixed} \\ \text{frame} \\ (\,|h_{\mu\nu}|\ll 1\,\,) \end{array}$$

Trace-reversed

$$\bar{h}_{\mu\nu} \equiv h_{\mu\nu} - \frac{1}{2}h\eta_{\mu\nu}$$

$$\partial^{\alpha}\partial_{\alpha}\bar{h}_{\mu\nu} + \eta_{\mu\nu}\partial^{\alpha}\partial^{\beta}\bar{h}_{\alpha\beta} - \partial^{\alpha}\partial_{(\mu}\bar{h}_{\alpha\nu)} = -\frac{2}{m_{p}^{2}}T_{\mu\nu}$$

 $\mathcal{O}(h_{**})$ Einstein tensor expanded

residual
$$\partial^{\nu} \bar{h}_{\mu\nu} = 0$$

Lorentz gauge

$$\partial'^{\mu} \bar{h}'_{\mu\nu}(x') = \underbrace{\partial^{\mu} \bar{h}_{\mu\nu}(x)}_{=} - \Box \xi_{\nu} = 0 \quad \iff \quad \Box \xi_{\nu} = f(x)$$
$$\equiv f(x) \neq 0$$

1st approach to GWs

Minkowski
$$g_{\mu\nu}= \begin{matrix} \uparrow \\ \eta_{\mu\nu} + h_{\mu\nu}(x) \end{matrix} \quad \begin{array}{l} \text{fixed} \\ \text{frame} \end{matrix} \quad (|h_{\mu\nu}| \ll 1) \end{matrix}$$

Trace-reversed

$$\bar{h}_{\mu\nu} \equiv h_{\mu\nu} - \frac{1}{2}h\eta_{\mu\nu}$$

$$\partial^{\alpha}\partial_{\alpha}\bar{h}_{\mu\nu} + \eta_{\mu\nu}\partial^{\alpha}\partial^{\beta}\bar{h}_{\alpha\beta} - \partial^{\alpha}\partial_{(\mu}\bar{h}_{\alpha\nu)} = -\frac{2}{m_{p}^{2}}T_{\mu\nu}$$

 $\mathcal{O}(h_{**})$ Einstein tensor expanded

residual
$$\partial^{\nu} \bar{h}_{\mu\nu} = 0$$

Lorentz gauge

$$\frac{\partial'^{\mu}\bar{h}'_{\mu\nu}(x')}{\equiv f(x) \neq 0} = \underbrace{\partial^{\mu}\bar{h}_{\mu\nu}(x)} - \Box\xi_{\nu} = 0 \iff \Box\xi_{\nu} = f(x)$$
(solution always!)

1st approach to GWs

Minkowski
$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}(x) \begin{tabular}{l} {\bf fixed} \\ {\bf frame} \\ (\ |h_{\mu\nu}| \ll 1 \) \end{tabular}$$

Trace-reversed

$$\bar{h}_{\mu\nu} \equiv h_{\mu\nu} - \frac{1}{2}h\eta_{\mu\nu} \longrightarrow \partial^{\alpha}\partial_{\alpha}\bar{h}_{\mu\nu} + \eta_{\mu\nu}\partial^{\alpha}\partial^{\beta}\bar{h}_{\alpha\beta} - \partial^{\alpha}\partial_{(\mu}\bar{h}_{\alpha\nu)} = -\frac{2}{m_p^2}T_{\mu\nu}$$

residual

$$\partial^{\nu}\bar{h}_{\mu\nu} = 0$$

Lorentz gauge

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$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}(x) \begin{tabular}{l} \textbf{fixed} \\ \textbf{frame} \\ (|h_{\mu\nu}| \ll 1) \end{tabular}$$

Trace-reversed

$$\bar{h}_{\mu\nu} \equiv h_{\mu\nu} - \frac{1}{2}h\eta_{\mu\nu}$$

$$= 0$$

$$\partial^{\alpha}\partial_{\alpha}\bar{h}_{\mu\nu} + \underbrace{\eta_{\mu\nu}\partial^{\alpha}\partial^{\beta}\bar{h}_{\alpha\beta}}_{=0} - \underbrace{\partial^{\alpha}\partial_{(\mu}\bar{h}_{\alpha\nu)}}_{=0} = -\frac{2}{m_{p}^{2}}T_{\mu\nu}$$

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$$\bar{h}_{\mu\nu} \equiv h_{\mu\nu} - \frac{1}{2}h\eta_{\mu\nu} \qquad \qquad \partial^{\alpha}\partial_{\alpha}\bar{h}_{\mu\nu} + \eta_{\mu\nu}\partial^{\alpha}\partial^{\beta}\bar{h}_{\alpha\beta} - \partial^{\alpha}\partial_{(\mu}\bar{h}_{\alpha\nu)} = -\frac{2}{m_{p}^{2}}T_{\mu\nu}$$

$$\partial^{\nu} \bar{h}_{\mu\nu} = 0$$

Lorentz gauge

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$$(10 - 4 = 6 \text{ d.o.f.})$$

1st approach to GWs

Minkowski
$$g_{\mu\nu}=\eta_{\mu\nu}+h_{\mu\nu}(x) \ \ \begin{array}{l} \text{fixed} \\ \text{frame} \end{array} \\ (\ |h_{\mu\nu}|\ll 1\) \end{array}$$

Is that all?

1st approach to GWs

Minkowski
$$g_{\mu\nu}=\eta_{\mu\nu}+h_{\mu\nu}(x) \ \ {\it fixed} \ \ (|h_{\mu\nu}|\ll 1\)$$

Is that all? Not really ...

Minkowski
$$g_{\mu\nu}=\eta_{\mu\nu}+h_{\mu\nu}(x) \ \ {\it fixed} \ \ (|h_{\mu\nu}|\ll 1)$$

$$x'^{\mu} = x^{\mu} + \xi^{\mu}(x)$$
 with $\partial_{\alpha}\partial^{\alpha}\xi_{\mu} = 0$ (further residual gauge)

Minkowski
$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}(x) \ \ {\it fixed} \ \ (|h_{\mu\nu}| \ll 1)$$

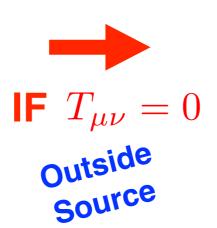
$$x'^{\mu}=x^{\mu}+\xi^{\mu}(x)$$
 with $\partial_{\alpha}\partial^{\alpha}\xi_{\mu}=0$ (further residual gauge) $(\partial^{\mu}\bar{h}_{\mu\nu}=0 \ o \ \partial'^{\mu}\bar{h}'_{\mu\nu}=0)$

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$$g_{\mu\nu}=\eta_{\mu\nu}+h_{\mu\nu}(x) \ \ {\it fixed} \ \ (|h_{\mu\nu}|\ll 1)$$

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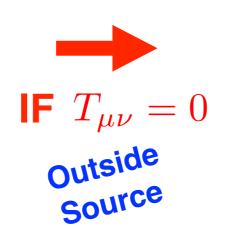
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1st approach to GWs

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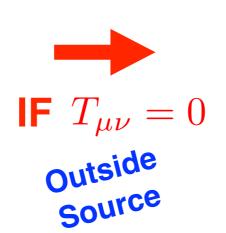


$$h^{0\mu} = 0, \quad h_i^i = 0, \quad \partial_j h_{ij} = 0$$

(transverse traceless gauge)

Minkowski
$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}(x) \begin{tabular}{l} fixed \\ frame \\ (|h_{\mu\nu}| \ll 1) \end{tabular}$$

$$x'^{\mu}=x^{\mu}+\xi^{\mu}(x)$$
 with $\partial_{\alpha}\partial^{\alpha}\xi_{\mu}=0$ (further residual gauge)



$$h^{0\mu}=0\,,\quad h^i_i=0\,,\quad \partial_j h_{ij}=0$$
 $h^{0\mu}=0$ $h^{0\mu}=0$ $h^{ij}=0$ h^{i

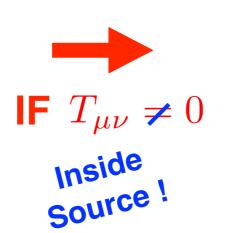
Minkowski
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Minkowski
$$g_{\mu\nu}= \begin{matrix} \uparrow \\ \eta_{\mu\nu} + h_{\mu\nu}(x) \end{matrix} \quad \begin{array}{l} \text{fixed} \\ \text{frame} \end{matrix} \\ (\ |h_{\mu\nu}| \ll 1 \) \end{array}$$

$$x'^{\mu} = x^{\mu} + \xi^{\mu}(x)$$
 with $\partial_{\alpha}\partial^{\alpha}\xi_{\mu} = 0$ (further residual gauge)



1st approach to GWs

Minkowski
$$g_{\mu\nu}=\eta_{\mu\nu}+h_{\mu\nu}(x) \ \ \begin{array}{l} \text{fixed} \\ \text{frame} \\ (\,|h_{\mu\nu}|\ll 1\,) \end{array}$$

$$x'^{\mu}=x^{\mu}+\xi^{\mu}(x)$$
 with $\partial_{\alpha}\partial^{\alpha}\xi_{\mu}=0$ (further residual gauge)



Cannot make $h_{*0}=0$

$$\partial_{lpha}\partial^{lpha}ar{h}_{\mu
u}=-rac{2}{m_{p}^{2}}T_{\mu
u}$$
 (6 - 4 = 2 d.o.f.)

Yet there are still only 2 <u>radiative</u> dof!

(TT gauge: 6 - 4 = 2 d.o.f.)

1st approach to GWs

$$h^{0\mu} = 0, \quad h_i^i = 0, \quad \partial_j h_{ij} = 0$$

Outside

(TT gauge: 6 - 4 = 2 d.o.f.)

1st approach to GWs

$$h^{0\mu} = 0, \quad h_i^i = 0, \quad \partial_j h_{ij} = 0$$

$$\partial_{\mu}\partial^{\mu}h_{ij} = 0$$

Wave Eq. → Gravitational Waves!

(TT gauge: 6 - 4 = 2 d.o.f.)

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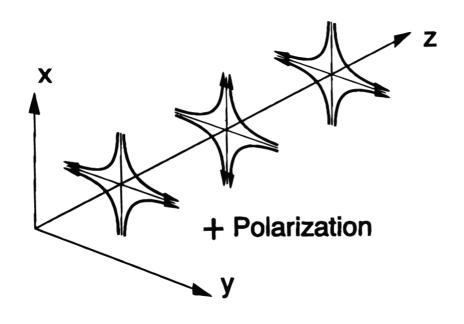
$$h^{0\mu} = 0, \quad h_i^i = 0, \quad \partial_j h_{ij} = 0$$

 $\partial_{\mu}\partial^{\mu}h_{ij} = 0$

Wave Eq. → Gravitational Waves!

can GW be 'gauged away'? No!

direction of propagation



Transverse (& Traceless)

Outside

Source

(TT gauge: 6 - 4 = 2 d.o.f.)

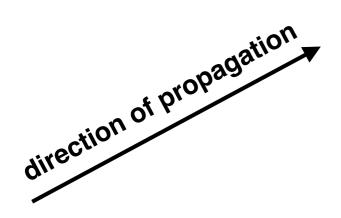
1st approach to GWs

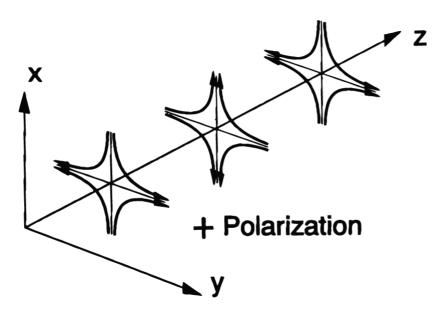
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Wave Eq. → Gravitational Waves!

can GW be 'gauged away'? No!





2 *dof* = 2 polarizations

Outside

Source

(TT gauge: 6 - 4 = 2 d.o.f.)

1st approach to GWs

(TT gauge: 6 - 4 = 2 d.o.f.)
$$h^{0\mu} = 0 , \quad h^i_i = 0 , \quad \partial_j h_{ij} = 0$$
 Outside Source

$$\partial_{\mu}\partial^{\mu}h_{ij} = 0$$

 $\partial_{\mu}\partial^{\mu}h_{ij}=0$ Wave Eq. \longrightarrow Gravitational Waves!

2
$$dof$$
 = 2 polarizations $h_{ab}(t,\mathbf{x}) = \int_{-\infty}^{\infty} df \int d\hat{n} \, h_{ab}(f,\hat{n}) e^{-2\pi i f(t-\hat{n}\mathbf{x})}$ (plane wave) transverse plane

(TT gauge: 6 - 4 = 2 d.o.f.)

1st approach to GWs

$$h^{0\mu}=0\,, \qquad h^i_i=0\,, \qquad \partial_j h_{ij}=0$$
 Outside Source

$$\partial_{\mu}\partial^{\mu}h_{ij} = 0$$

 $\partial_{\mu}\partial^{\mu}h_{ij}=0$ Wave Eq. \longrightarrow Gravitational Waves!

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$$h_{ab}(f,\hat{n}) = \sum_{A=+,\mathbf{x}} h_A(f,\hat{n}) \epsilon_{ab}^{(A)}(\hat{n}) \ = \left(\begin{array}{ccc} h_+ & h_x & 0 \\ h_x & -h_+ & 0 \\ 0 & 0 & 0 \end{array} \right) \quad \begin{array}{c} \text{Transverse-} \\ \text{Traceless} \\ \text{(2 dof)} \end{array}$$

(TT gauge: 6 - 4 = 2 d.o.f.)

1st approach to GWs

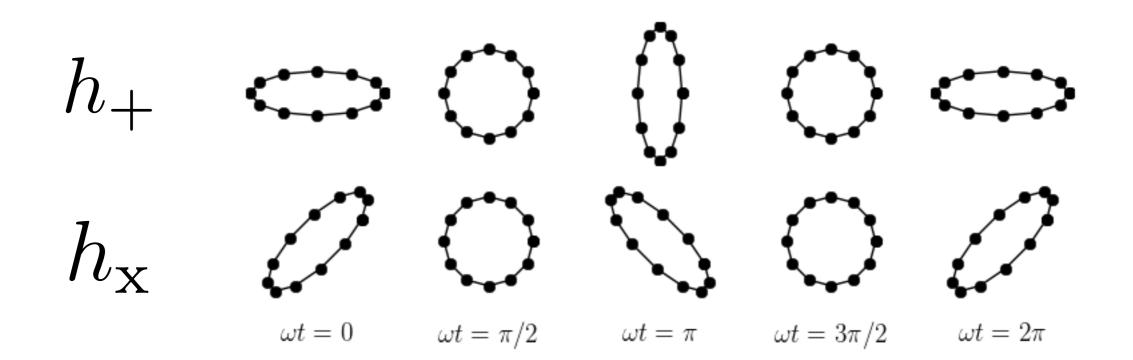
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Wave Eq. → Gravitational Waves!



Definition of GWs 2nd approach

2nd approach to GWs

(gauge invariant def.)

Minkowski
$$g_{\mu\nu} = \eta_{\mu\nu} + \delta g_{\mu\nu} \qquad (|\delta g_{\mu\nu}| \ll 1)$$

2nd approach to GWs

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$$g_{\mu\nu} = \eta_{\mu\nu} + \delta g_{\mu\nu} \qquad (|\delta g_{\mu\nu}| \ll 1)$$

 $\delta g_{00} = -2\phi,$

$$\delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i),$$

$$\delta g_{ij} = \delta g_{ji} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij},$$

(svt decomposition)

s: scalar

v: vector

t: tensor

2nd approach to GWs

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Recall lectures by David Wands

2nd approach to GWs

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s: scalar

v: vector

t: tensor

 $T_{00} = \rho$,

(svt decomposition)

$$T_{0i} = T_{i0} = \partial_i u + u_i,$$

$$T_{ij} = T_{ji} = p \,\delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) \sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}.$$

(svt metric perturbations)

$$\delta g_{00} = -2\phi,$$

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,

(svt E/p-tensor components)

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 $\delta g_{00}=-2\phi,$ (svt metric perturbations)

$$\delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i),$$

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	$\delta g_{\mu u}$	$T_{\mu u}$
Scalar(s) Vector(s) Tensor(s) $\in \Re^3$	ϕ, B, ψ, E S_i, F_i h_{ij}	$ ho, u, p, \sigma$ u_i, v_i Π_{ij}

 $\delta g_{00} = -2\phi,$ (svt metric perturbations) $\delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i),$

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	$\delta g_{\mu u}$	$T_{\mu u}$
Scalar(s)	ϕ, B, ψ, E	ρ , u , p , σ
$Vector(s) $ $\in \Re^3$	S_i, F_i	u_i, v_i
$ \left\{ \begin{array}{l} \text{Vector(s)} \\ \text{Tensor(s)} \end{array} \right\} \in \Re^3 $	h_{ij}	Π_{ij}

 $\delta g_{00} = -2\phi$, (svt metric perturbations)

$$\delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i),$$

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Scalar(s)	ϕ , B , ψ , E	ρ, u, p, σ
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	$\delta g_{\mu u}$	$T_{\mu u}$
Scalar(s) Vector(s) Tensor(s) $\in \Re^3$	ϕ, B, ψ, E $\longrightarrow S_i, F_i \longleftarrow$ h_{ij}	$ ho, u, p, \sigma$ u_i, v_i Π_{ij}

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	$\delta g_{\mu u}$	$T_{\mu u}$
Scalar(s)	ϕ , B , ψ , E	ρ , u , p , σ
$\{Vector(s)\}\in\Re^3$	S_i, F_i	$\longrightarrow u_i, v_i \longleftarrow$
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	$\delta g_{\mu u}$	$T_{\mu u}$
Scalar(s) Vector(s) Tensor(s) $\in \Re^3$	ϕ, B, ψ, E S_i, F_i h_{ij}	ρ, u, p, σ u_i, v_i Π_{ij}

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16 degrees of freedom

 $T_{00} = \rho$,

(svt E/p-tensor components)

$$T_{0i}=T_{i0}=\partial_i u+u_i,$$

 $T_{ij} = T_{ji} = p \, \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) \sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}.$

	$\delta g_{\mu u}$	$T_{\mu u}$
Scalar(s) Vector(s) Tensor(s) $\in \Re^3$	ϕ, B, ψ, E S_i, F_i h_{ij}	$ ho, u, p, \sigma$ u_i, v_i Π_{ij}

(svt metric perturbations)

$$\delta g_{00} = -2\phi,$$

$$\delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i),$$

$$\delta g_{ij} = \delta g_{ji} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij},$$

16 degrees of freedom

$$T_{00} = \rho$$
,

(svt E/p-tensor components)

$$T_{0i} = T_{i0} = \partial_i u + u_i,$$

16 degrees of freedom

$$T_{ij} = T_{ji} = p \,\delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) \sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}.$$

In order NOT to over-count degrees of freedom

 $\delta g_{00} = -2\phi$,

(svt metric perturbations)

$$\delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i)$$

$$\delta g_{ij} = \delta g_{ji} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij},$$

16 degrees of freedom

 $T_{00} = \rho$,

(svt E/p-tensor components)

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16 degrees of freedom

$$T_{ij} = T_{ji} = p \,\delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) \sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}.$$

In order NOT to over-count degrees of freedom

$$\partial_i S_i = 0$$
 (1 constraint), $\partial_i F_i = 0$ (1 constraint), $\partial_i h_{ij} = 0$ (3 constraints), $h_{ii} = 0$ (1 constraint)

Metric perturbations

 $\delta g_{00} = -2\phi$,

(svt metric perturbations)

$$\delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i)$$

$$\delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i),$$

$$\delta g_{ij} = \delta g_{ji} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij},$$

16 degrees of freedom

 $T_{00} = \rho$

(svt E/p-tensor components)

$$T_{0i} = T_{i0} = \partial_i u + u_i,$$

16 degrees of freedom

$$T_{ij} = T_{ji} = p \, \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) \sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}.$$

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Metric perturbations

$$\partial_i u_i = 0$$
 (1 constraint), $\partial_i v_i = 0$ (1 constraint), $\partial_i \Pi_{ij} = 0$ (3 constraints), $\Pi_{ii} = 0$ (1 constraint),

Energy/Momentum tensor

 $\delta g_{00} = -2\phi,$

(svt metric perturbations)

$$\delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i),$$

$$\delta g_{ij} = \delta g_{ji} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij},$$

16 degrees of freedom

 $T_{00} = \rho$,

(svt E/p-tensor components)

$$T_{0i} = T_{i0} = \partial_i u + u_i,$$

16 degrees of freedom

$$T_{ij} = T_{ji} = p \,\delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) \sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}.$$

In order NOT to over-count degrees of freedom

$$\partial_i S_i = 0$$
 (1 constraint), $\partial_i F_i = 0$ (1 constraint),

$$\partial_i h_{ij} = 0$$
 (3 constraints), $h_{ii} = 0$ (1 constraint)

6 constraints for metric perturbations

$$\partial_i u_i = 0$$
 (1 constraint), $\partial_i v_i = 0$ (1 constraint),

$$\partial_i \Pi_{ij} = 0$$
 (3 constraints), $\Pi_{ii} = 0$ (1 constraint),

6 constraints for E/p tensor components

 $\delta g_{00} = -2\phi,$

(svt metric perturbations)

$$S_{\alpha} = S_{\alpha} = (3 D + C)$$

$$\delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i),$$

$$\delta g_{ij} = \delta g_{ji} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij},$$

10

16 degrees of freedom

$$T_{00} = \rho$$
,

(svt E/p-tensor components)

$$T_{0i}=T_{i0}=\partial_i u+u_i,$$

$$T_{ij} = T_{ji} = p \, \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) \sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}.$$

16 degrees of freedom

In order NOT to over-count degrees of freedom

$$\partial_i S_i = 0$$
 (1 constraint), $\partial_i F_i = 0$ (1 constraint),

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 (3 constraints), $\Pi_{ii} = 0$ (1 constraint),

6 constraints for E/p tensor components

 $\delta g_{00} = -2\phi$,

(svt metric perturbations)

$$\delta g_{00} = -2\phi$$

$$\delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i),$$

$$\delta g_{ij} = \delta g_{ji} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij},$$

10 degrees of freedom

 $T_{00} = \rho$

(svt E/p-tensor components)

$$T_{0i} = T_{i0} = \partial_i u + u_i,$$

10 degrees of freedom

$$T_{ij} = T_{ji} = p \,\delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) \sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}.$$

In order NOT to over-count degrees of freedom

$$\partial_i S_i = 0$$
 (1 constraint), $\partial_i F_i = 0$ (1 constraint),

$$\partial_i h_{ij} = 0$$
 (3 constraints), $h_{ii} = 0$ (1 constraint)

(6) constraints for metric perturbations

$$\partial_i u_i = 0$$
 (1 constraint), $\partial_i v_i = 0$ (1 constraint),

$$\partial_i \Pi_{ij} = 0$$
 (3 constraints), $\Pi_{ii} = 0$ (1 constraint),

(6) constraints for E/p tensor components

(svt metric perturbations)

$$\delta g_{00} = -2\phi$$
,

$$\delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i),$$

$$\delta g_{ij} = \delta g_{ji} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij},$$

10 degrees of freedom

$$T_{00} = \rho$$
,

(svt E/p-tensor components)

$$T_{0i}=T_{i0}=\partial_i u+u_i,$$

10 degrees of freedom

$$T_{ij} = T_{ji} = p \,\delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) \sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}.$$

 $\delta g_{00} = -2\phi,$

(svt metric perturbations)

$$\delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i),$$

10 degrees of freedom

$$\delta g_{ij} = \delta g_{ji} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij},$$

 $T_{00} = \rho$,

(svt E/p-tensor components)

$$T_{0i}=T_{i0}=\partial_i u+u_i,$$

10 degrees of freedom

$$T_{ij} = T_{ji} = p \,\delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) \sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}.$$

Physical Constraints

$$\partial^{\mu}T_{\mu\nu}=0$$

 $\delta g_{00} = -2\phi$,

(svt metric perturbations)

$$\delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i)$$

$$\delta g_{ij} = \delta g_{ji} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij},$$

10 degrees of freedom

 $T_{00} = \rho$

(svt E/p-tensor components)

$$T_{0i} = T_{i0} = \partial_i u + u_i$$

$$T_{ij} = T_{ji} = p \,\delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) \sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}.$$

10 degrees of freedom

Physical Constraints

$$\partial^{\mu}T_{\mu\nu}=0$$

$$\partial^{\mu}T_{\mu\nu} = 0 \implies \begin{cases} \nabla^{2}u = \dot{\rho} \text{ (1 constraint),} \\ \nabla^{2}\sigma = \frac{3}{2}(\dot{u} - p) \text{ (1 constraint),} \end{cases}$$
 4 constraints (due to E/p conservation)
$$\nabla^{2}v_{i} = \dot{u}_{i} \text{ (2 constraints).}$$

$$\nabla^2 v_i = \dot{u}_i$$
 (2 constraints)

 $\delta g_{00} = -2\phi$,

(svt metric perturbations)

$$\delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i),$$

$$\delta g_{ij} = \delta g_{ji} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij},$$

10 degrees of freedom

 $T_{00} = \rho$

(svt E/p-tensor components)

$$T_{0i} = T_{i0} = \partial_i u + u_i,$$

$$T_{ij} = T_{ji} = p \,\delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) \sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}.$$

10 degrees of freedom

Physical Constraints

$$\partial^{\mu}T_{\mu\nu}=0$$

$$\nabla^2 u = \dot{\rho} \ (1 \text{ constraint}),$$

$$\nabla^2 \sigma = \frac{3}{2}(\dot{u} - p) \quad (1 \text{ constraint}),$$

$$\nabla^2 v_i = \dot{u}_i \ (2 \text{ constraints})$$

 $\partial^{\mu}T_{\mu\nu} = 0 \implies \begin{cases} \nabla^{2}u = \dot{\rho} \text{ (1 constraint),} \\ \nabla^{2}\sigma = \frac{3}{2}(\dot{u} - p) \text{ (1 constraint),} \end{cases}$ (due to E/p conservation) $\nabla^{2}v_{i} = \dot{u}_{i} \text{ (2 constraints).}$

 $\delta g_{00} = -2\phi,$

(svt metric perturbations)

$$\delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i),$$

$$\delta g_{ij} = \delta g_{ji} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij},$$

10 degrees of freedom

 $T_{00} = \rho$,

(svt E/p-tensor components)

$$T_{0i} = T_{i0} = \partial_i u + u_i,$$

6 degrees of freedom

$$T_{ij} = T_{ji} = p \,\delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) \sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}.$$

$$\partial^{\mu}T_{\mu\nu}=0$$

 $\delta g_{00} = -2\phi,$

(svt metric perturbations)

$$\delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i),$$

10 degrees of freedom

$$\delta g_{ij} = \delta g_{ji} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij},$$

 $T_{00}=\rho$,

(svt E/p-tensor components)

$$T_{0i} = T_{i0} = \partial_i u + u_i,$$

6 degrees of freedom

$$T_{ij} = T_{ji} = p \,\delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) \sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}.$$

Physical Constraints

$$\partial^{\mu}G_{\mu\nu}=0$$
 \Longrightarrow $\left[\ldots\right]$

 $\delta g_{00} = -2\phi,$

(svt metric perturbations)

$$\delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i),$$

$$\delta g_{ij} = \delta g_{ji} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij},$$

10 degrees of freedom

 $T_{00} = \rho$,

(svt E/p-tensor components)

$$T_{0i} = T_{i0} = \partial_i u + u_i,$$

$$T_{ij} = T_{ji} = p \,\delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) \sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}.$$

6 degrees of freedom

Physical Symmetry

$$\int x_{\mu} \longrightarrow x_{\mu} + \xi_{\mu}$$

$$\delta g_{\mu\nu} \longrightarrow \delta g_{\mu\nu} - \partial_{\mu} \xi_{\nu} - \partial_{\nu} \xi_{\mu}$$

 2ϕ , (svt metric perturbations)

$$\delta g_{00} = -2\phi,$$

$$\delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i),$$

$$\delta g_{ij} = \delta g_{ji} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij},$$

10 degrees of freedom

$$T_{00} = \rho$$
,

(svt E/p-tensor components)

$$T_{0i} = T_{i0} = \partial_i u + u_i,$$

6 degrees of freedom

$$T_{ij} = T_{ji} = p \,\delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) \sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}.$$

Physical Symmetry

(4 *d.o.f.* spurious)

$$x_{\mu} \longrightarrow x_{\mu} + \xi_{\mu}$$

$$\delta g_{\mu\nu} \longrightarrow \delta g_{\mu\nu} - \partial_{\mu} \xi_{\nu} - \partial_{\nu} \xi_{\mu}$$

$$\xi_{\mu} = (\xi_0, \xi_i) \equiv (d_0, \partial_i d + d_i)$$

with $\partial_i d_i = 0$,

 $\delta g_{00} = -2\phi$

(svt metric perturbations)

$$\delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i),$$

$$\delta g_{ij} = \delta g_{ji} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij},$$

10 degrees of freedom

 $T_{00} = \rho$,

(svt E/p-tensor components)

$$T_{0i} = T_{i0} = \partial_i u + u_i,$$

 $T_{ij} = T_{ji} = p \,\delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) \sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}.$

Physical Symmetry

(4 d.o.f. spurious)

$$x_{\mu} \longrightarrow x_{\mu} + \xi_{\mu}$$

$$\delta g_{\mu\nu} \longrightarrow \delta g_{\mu\nu} - \partial_{\mu} \xi_{\nu} - \partial_{\nu} \xi_{\mu}$$

$$\xi_{\mu} = (\xi_0, \xi_i) \equiv (d_0, \partial_i d + d_i)$$

with $\partial_i d_i = 0$,

$$S_i \longrightarrow S_i - \dot{d}_i, \quad F_i \longrightarrow \quad F_i - 2d_i,$$

$$h_{ij} \longrightarrow h_{ij}.$$

 $\delta g_{00} = -2\phi$

(svt metric perturbations)

$$\delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i),$$

$$\delta g_{ij} = \delta g_{ji} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij},$$

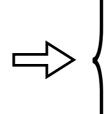
10 degrees of freedom

 $T_{00} = \rho$,

(svt E/p-tensor components)

$$T_{0i} = T_{i0} = \partial_i u + u_i,$$

$$T_{ij} = T_{ji} = p \, \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) \sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}.$$

Physical Symmetry 

$$\phi \longrightarrow \phi - \dot{d}_0, \quad B \longrightarrow B - d_0 - \dot{d},$$
 $\psi \longrightarrow \psi + \frac{1}{3} \nabla^2 d, \quad E \longrightarrow E - 2d,$

$$S_i \longrightarrow S_i - \dot{d}_i, \quad F_i \longrightarrow \quad F_i - 2d_i$$

$$h_{ij} \longrightarrow h_{ij}$$

 $\delta g_{00} = -2\phi,$

(svt metric perturbations)

$$\delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i)$$

$$\delta g_{ij} = \delta g_{ji} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij},$$

6 degrees of freedom

 $T_{00} = \rho$,

(svt E/p-tensor components)

$$T_{0i} = T_{i0} = \partial_i u + u_i,$$

6 degrees of freedom

$$T_{ij} = T_{ji} = p \,\delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) \sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}.$$

Physical Symmetry (4 *d.o.f.*

spurious)

$$\begin{cases} \phi \longrightarrow \phi - \dot{d}_0, & B \longrightarrow B - d_0 - \dot{d}, \\ \psi \longrightarrow \psi + \frac{1}{3} \nabla^2 d, & E \longrightarrow E - 2d, \end{cases}$$

$$S_i \longrightarrow S_i - \dot{d}_i, \quad F_i \longrightarrow F_i - 2d_i,$$

$$h_{ij} \longrightarrow h_{ij}.$$

 $\delta g_{00} = -2\phi$,

(svt metric perturbations)

$$\delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i)$$

$$\delta g_{ij} = \delta g_{ji} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij},$$

6 degrees of freedom

 $T_{00} = \rho$,

(svt E/p-tensor components)

$$T_{0i} = T_{i0} = \partial_i u + u_i$$

$$(S_{ii}\nabla^2)\sigma + \partial_iv_i + \partial_iv_i + \Pi_{ii}$$

6 degrees of freedom

 $T_{ij} = T_{ji} = p \,\delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) \sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}.$

Gauge Invariant!

$$\begin{cases} \phi \longrightarrow \phi - \dot{d}_0, & B \longrightarrow B - d_0 - \dot{d}, \\ \psi \longrightarrow \psi + \frac{1}{3} \nabla^2 d, & E \longrightarrow E - 2d, \\ \text{Symmetry} \end{cases}$$

$$\begin{cases} S_i \longrightarrow S_i - \dot{d}_i, & F_i \longrightarrow F_i - 2d_i, \end{cases}$$



$$\Theta \equiv -2\psi - \frac{1}{3}\nabla^2 E,$$

 $\Phi \equiv -\phi + \dot{B} - \frac{1}{2}\ddot{E},$

$$\Sigma_i \equiv S_i - \frac{1}{2}\dot{F}_i,$$

with $\partial_i \Sigma_i = 0$

(4 d.o.f. spurious)

 $\delta g_{00} = -2\phi$,

(svt metric perturbations)

$$\delta g_{0i} = \delta g_{i0} = (\partial_i B + S_i),$$

$$\delta g_{ij} = \delta g_{ji} = -2\psi \delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) E + \partial_i F_j + \partial_j F_i + h_{ij},$$

6 degrees of freedom

 $T_{00} = \rho$

(svt E/p-tensor components)

$$T_{0i} = T_{i0} = \partial_i u + u_i,$$

$$-\partial_i v_i + \Pi_{ii}.$$

6 degrees of freedom

 $T_{ij} = T_{ji} = p \,\delta_{ij} + (\partial_i \partial_j - \frac{1}{3} \delta_{ij} \nabla^2) \sigma + \partial_i v_j + \partial_j v_i + \Pi_{ij}.$

Gauge Invariant!

$$\begin{cases} \phi \longrightarrow \phi - \dot{d}_0, & B \longrightarrow B - d_0 - \dot{d}, \\ \psi \longrightarrow \psi + \frac{1}{3} \nabla^2 d, & E \longrightarrow E - 2d, \end{cases}$$
 Physical

$$S_i \longrightarrow S_i - \dot{d}_i, \quad F_i \longrightarrow F_i - 2d_i,$$

$$h_{ij} \longrightarrow h_{ij}$$
. (2

$$\Theta \equiv -2\psi - \frac{1}{3}\nabla^2 E, \quad (1)$$

 $\Phi \equiv -\phi + \dot{B} - \frac{1}{2} \ddot{E}, \quad (1)$

$$\Sigma_i \equiv S_i - \frac{1}{2}\dot{F}_i, \qquad (2)$$

with $\partial_i \Sigma_i = 0$

Symmetry (4 d.o.f. spurious)

Gauge Invariant!

$$\Phi \equiv -\phi + \dot{B} - \frac{1}{2} \ddot{E}, \qquad (1)$$

$$\Theta \equiv -2\psi - \frac{1}{3}\nabla^2 E, \qquad (1)$$

$$\Phi \equiv -\phi + \dot{B} - rac{1}{2}\ddot{E},$$
 (1)
 $\Theta \equiv -2\psi - rac{1}{3}
abla^2 E,$ (1)
 $\Sigma_i \equiv S_i - rac{1}{2}\dot{F}_i,$ ($\partial_i\Sigma_i = 0$) (2)
 $h_{ij} \equiv h_{ij},$ ($h_{ii} = \partial_i h_{ij} = 0$) (2)

$$h_{ij} \equiv h_{ij} , \qquad (h_{ii} = \partial_i h_{ij} = 0) \qquad (2)$$

6 gauge invariant degrees of freedom

Gauge Invariant!

$$\Phi \equiv -\phi + \dot{B} - \frac{1}{2}\ddot{E}, \qquad (1)$$

$$\Theta \equiv -2\psi - \frac{1}{3}\nabla^2 E, \qquad (1)$$

$$\Phi \equiv -\phi + \dot{B} - rac{1}{2}\ddot{E},$$
 (1)
 $\Theta \equiv -2\psi - rac{1}{3}
abla^2 E,$ (1)
 $\Sigma_i \equiv S_i - rac{1}{2}\dot{F}_i,$ ($\partial_i\Sigma_i = 0$) (2)
 $h_{ij} \equiv h_{ij},$ ($h_{ii} = \partial_i h_{ij} = 0$) (2)

$$h_{ij} \equiv h_{ij} , \qquad (h_{ii} = \partial_i h_{ij} = 0) \qquad (2)$$

6 gauge invariant degrees of freedom



Gauge Invariant Einstein Tensor

$$G_{00} = -\nabla^2 \Theta,$$

$$G_{0i} = -rac{1}{2}
abla^2\Sigma_i - \partial_i\dot{\Theta},$$

$$G_{00} = -
abla^2 \Theta,$$

$$G_{0i} = -\frac{1}{2}
abla^2 \Sigma_i - \partial_i \dot{\Theta},$$

$$G_{ij} = -\frac{1}{2}
abla h_{ij} - \partial_{(i} \dot{\Sigma}_{j)} - \frac{1}{2} \partial_i \partial_j (2\Phi + \Theta) + \delta_{ij} \left[\frac{1}{2}
abla^2 (2\Phi + \Theta) - \ddot{\Theta} \right].$$

Gauge Invariant!

$$\Phi \equiv -\phi + \dot{B} - \frac{1}{2} \ddot{E},$$
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6 gauge invariant degrees of freedom



Gauge Invariant (perturbed) Einstein Eqs.

$$abla^2 \Theta = -\frac{1}{m_p^2} \rho, \qquad
abla^2 \Phi = \frac{1}{2m_p^2} \left(\rho + 3p - 3\dot{u} \right)$$

$$abla^2 \Sigma_i = -rac{2}{m_p^2} S_i, \qquad \Box h_{ij} = -rac{2}{m_p^2} \Pi_{ij}.$$

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Gravitational Waves (GWs) are TT *d.o.f.* metric perturbations, independently of system of reference

Definition of GWs 3rd approach

3rd approach to GWs

(for a curved space-time)

$$g_{\mu\nu}(x) = \bar{g}_{\mu\nu}(x) + \delta g_{\mu\nu}(x) , \quad |\delta g_{\mu\nu}| \ll 1$$

(separation not well defined)

3rd approach to GWs

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More subtle problem! Solution: Separation of scales!

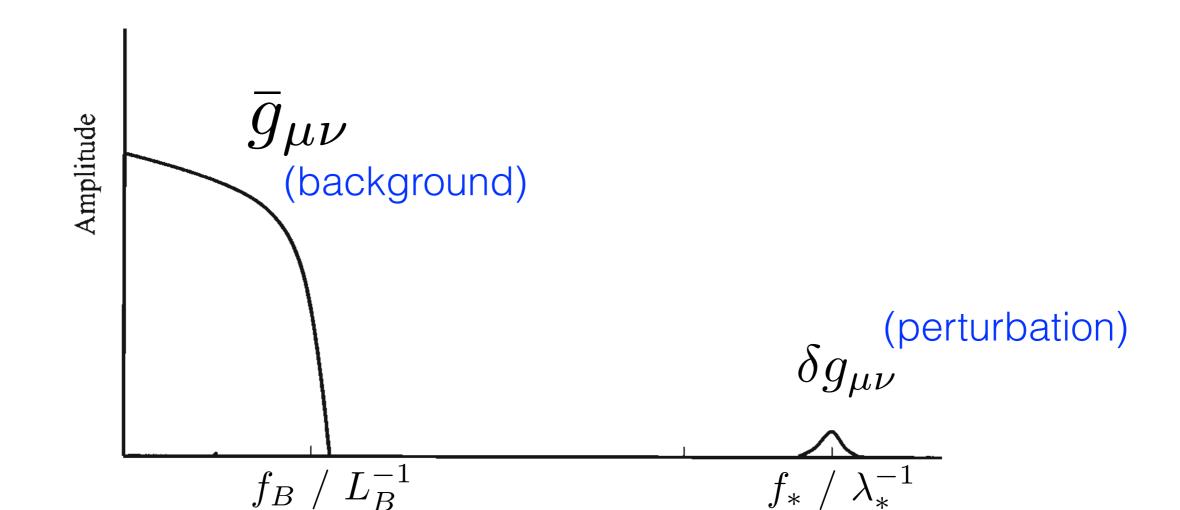
Maggiore's 1st Maggiore's 1st Book on GWs

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Low Freq. / Long Scale: $\bar{R}_{\mu\nu} = -[R^{(2)}_{\mu\nu}]^{\rm Low} + \frac{1}{m_p^2} \left(T_{\mu\nu} - \frac{1}{2} g_{\mu\nu} T \right)^{\rm Low}$

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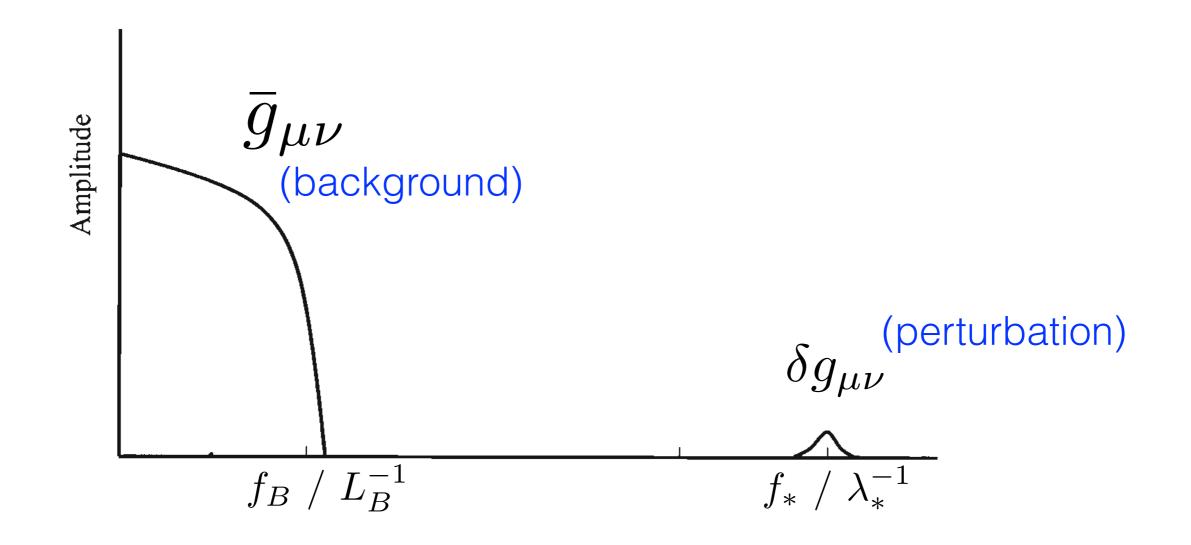
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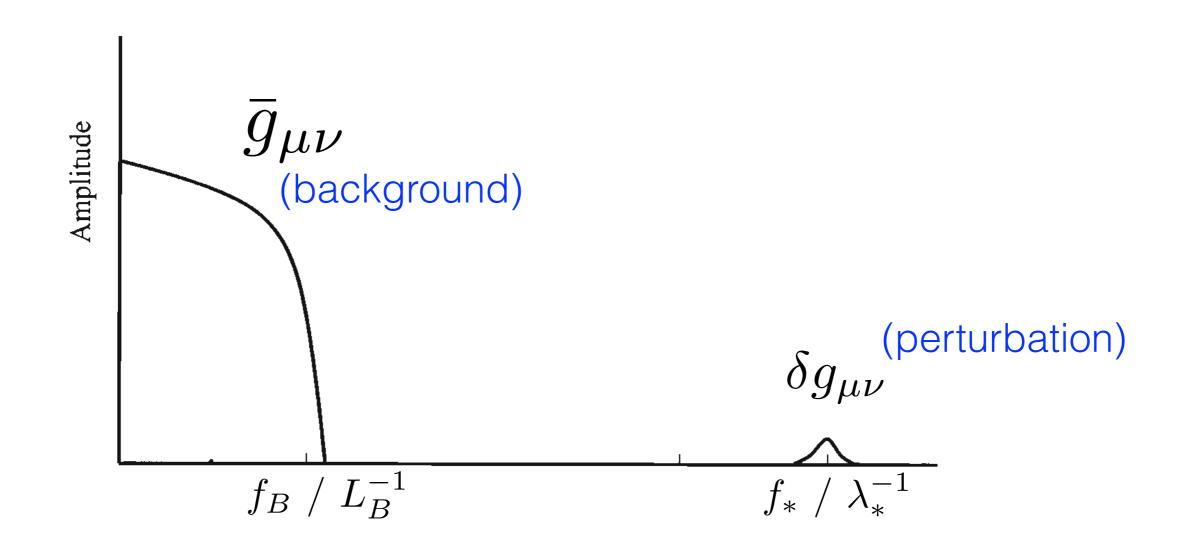
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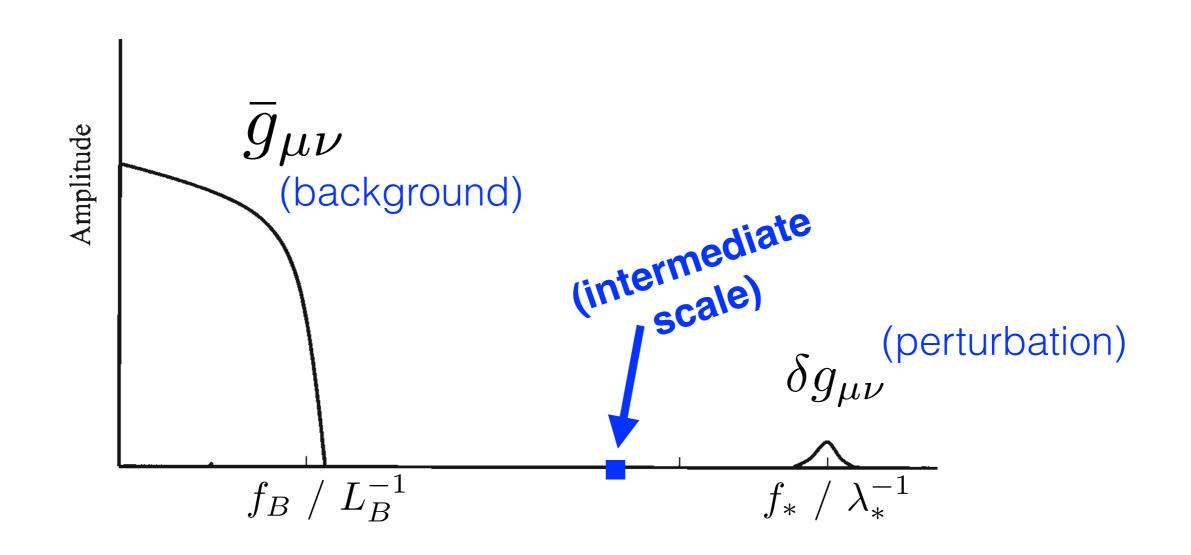
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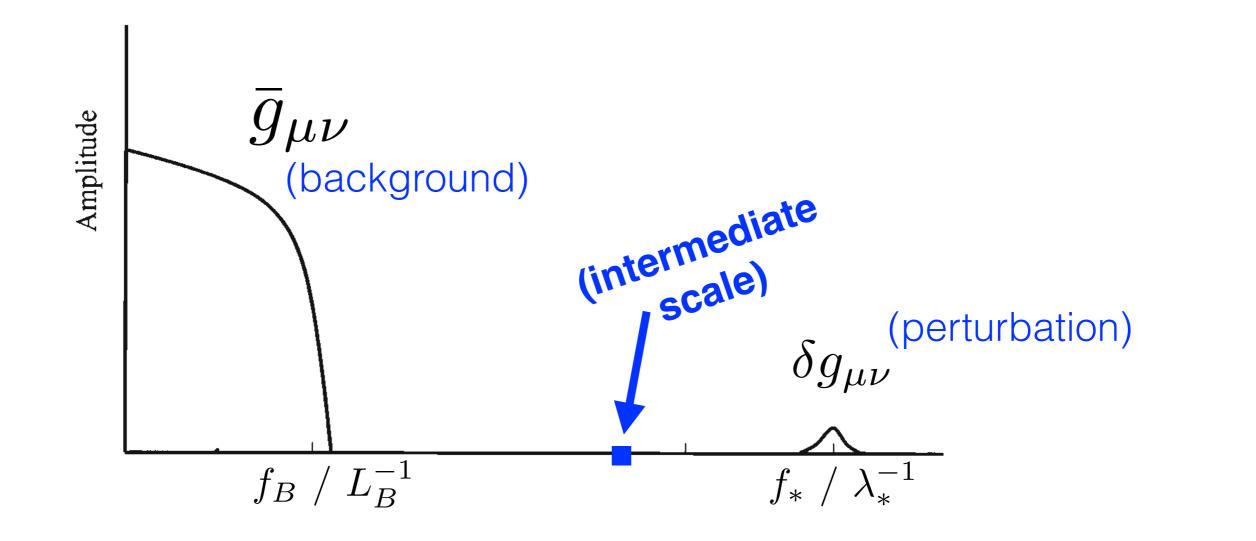
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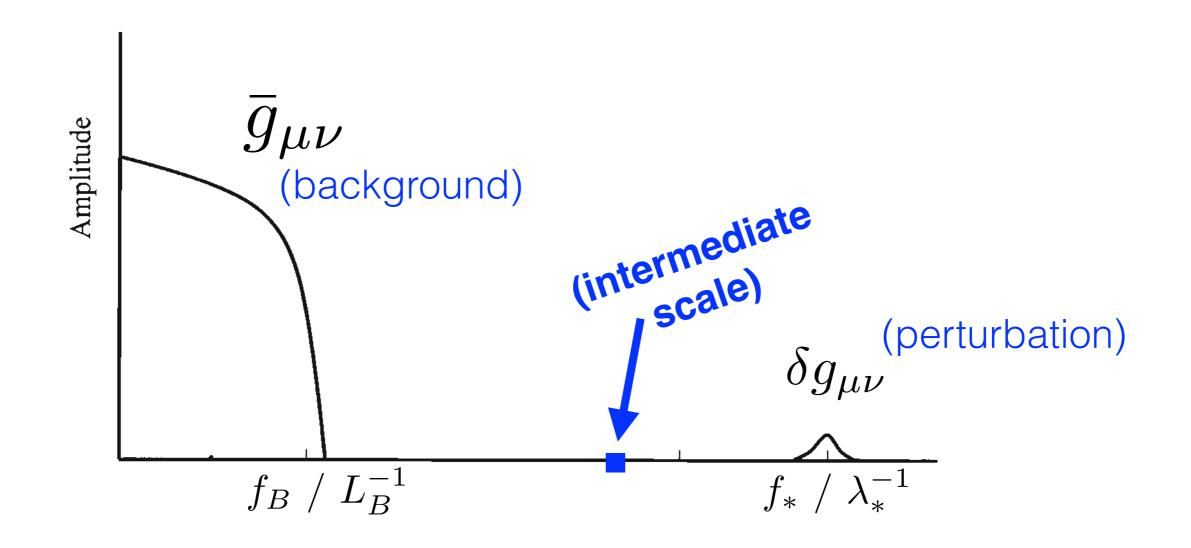


Low Freq. / Long Scale:
$$\bar{R}_{\mu\nu} = -\langle R_{\mu\nu}^{(2)} \rangle + \frac{1}{m_p^2} \langle T_{\mu\nu} - \frac{1}{2} g_{\mu\nu} T \rangle \xrightarrow{\text{average}} t_{\mu\nu} = -\frac{1}{m_p^2} \langle R_{\mu\nu}^{(2)} - \frac{1}{2} \bar{g}_{\mu\nu} R^{(2)} \rangle \qquad \langle T_{\mu\nu} - \frac{1}{2} g_{\mu\nu} T \rangle = \bar{T}^{\mu\nu} - \frac{1}{2} \bar{g}_{\mu\nu} \bar{T}$$



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GW energy-momentum tensor

Low Freq. / Long Scale:
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It can be shown that only TT dof contribute to < ... >

$$t_{\mu\nu} = \frac{m_p^2}{4} \langle \partial_{\mu} \delta g_{ij}^{\text{TT}} \partial_{\nu} \delta g_{ij}^{\text{TT}} \rangle \left| \frac{dE}{dAdt} = \frac{m_p^2}{4} \langle \dot{h}_{ij}^{\text{TT}} \dot{h}_{ij}^{\text{TT}} \rangle \right|$$

GW energy-momentum tensor

GW power/area radiated

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It can be shown that only TT dof contribute to < ... >

$$t_{\mu\nu} = \frac{m_p^2}{4} \langle \partial_{\mu} \delta g_{ij}^{\text{TT}} \partial_{\nu} \delta g_{ij}^{\text{TT}} \rangle \left| \begin{array}{c} \bullet \\ \delta g_{ij} \equiv h_{ij} \end{array} \right| \rho_{\text{GW}} \equiv \frac{m_p^2}{4} \langle \dot{h}_{ij}^{\text{TT}} \dot{h}_{ij}^{\text{TT}} \rangle$$

GW energy-momentum tensor

GW energy density

What about the High Freq. / Short Scale?
$$R_{\mu\nu}^{(1)} = -[R_{\mu\nu}^{(2)}]^{\rm High} + \frac{1}{m_p^2} \left(T_{\mu\nu} - \frac{1}{2}g_{\mu\nu}T\right)^{\rm High}$$

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$$\frac{|R_{\mu}^{(2)}|^{\mathrm{High}}}{|R_{\mu}^{(1)}|} \sim \mathcal{O}\left(\frac{\lambda_*}{L_B}\right) \longrightarrow |R_{\mu}^{(2)}|^{\mathrm{High}} \text{ negligible}$$

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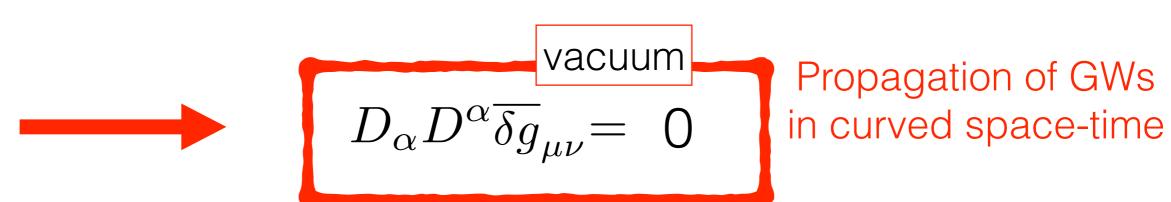
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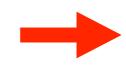


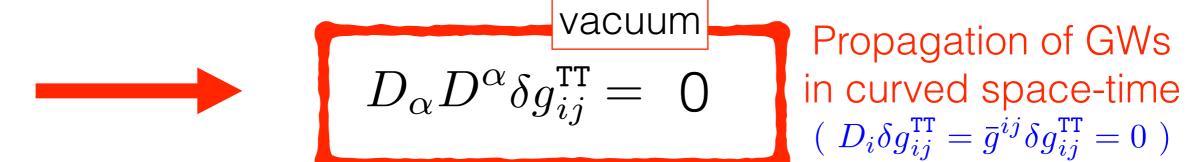
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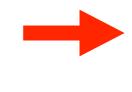


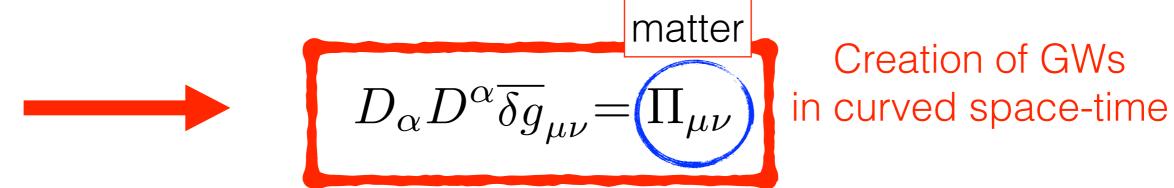
Propagation of GWs $(D_i \delta g_{ij}^{\mathsf{TT}} = \bar{g}^{ij} \delta g_{ij}^{\mathsf{TT}} = 0)$

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$$\begin{split} R_{\mu\nu}^{(1)} &= \bar{g}^{\alpha\beta} \left(D_{\alpha} D_{(\mu} \overline{\delta g_{\nu)\beta}} - D_{\mu} D_{\nu} \overline{\delta g_{\alpha\beta}} - D_{\alpha} D_{\beta} \overline{\delta g_{\mu\nu}} \right) \\ D_{\mu} \overline{\delta g}_{\mu\nu} &= 0 \quad (\, \overline{\delta g}_{\mu\nu} \equiv \delta g_{\mu\nu} - \frac{\mathrm{i}}{2} \bar{g}_{\mu\nu} \bar{g}^{\alpha\beta} \delta g_{\alpha\beta} \,) \quad \text{Lorentz} \\ \mathrm{gauge} \end{split}$$





Creation of GWs

What about the

What about the High Freq. / Short Scale?
$$R_{\mu\nu}^{(1)}=-[R_{\mu\nu}^{(2)}]^{\rm High}\left(+\frac{1}{m_p^2}\left(T_{\mu\nu}-\frac{1}{2}g_{\mu\nu}T\right)^{\rm High}\right)$$

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Creation of GWs in curved space-time

TT dof = truly radiative! [no gauge choice]

Definition of GWs

* 1st approach: Lin Grav in Minkowski



* 2nd approach: SVT decomp. <



Some perspective

GW Propagation/Creation in Cosmology

FLRW:
$$ds^2 = a^2(-d\eta^2 + (\delta_{ij} + h_{ij})dx^i dx^j)$$
, $TT: \begin{cases} h_{ii} = 0 \\ h_{ij},_j = 0 \end{cases}$ (conformal time)

GW Propagation/Creation in Cosmology

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 TT:
$$\begin{cases} h_{ii} = 0 \\ h_{ij},_j = 0 \end{cases}$$
 (conformal time)

Creation/Propagation GWs in FLRW

Eom: $h_{ij}^{\prime\prime} + 2\mathcal{H}h_{ij}^{\prime} - \nabla^2 h_{ij} = 16\pi G \Pi_{ij}^{\mathrm{TT}}.$

Anisotropic Stress

$$\Pi_{ij} = T_{ij} - \langle T_{ij} \rangle_{\mathsf{FLRW}}$$

GW Propagation/Creation in Cosmology

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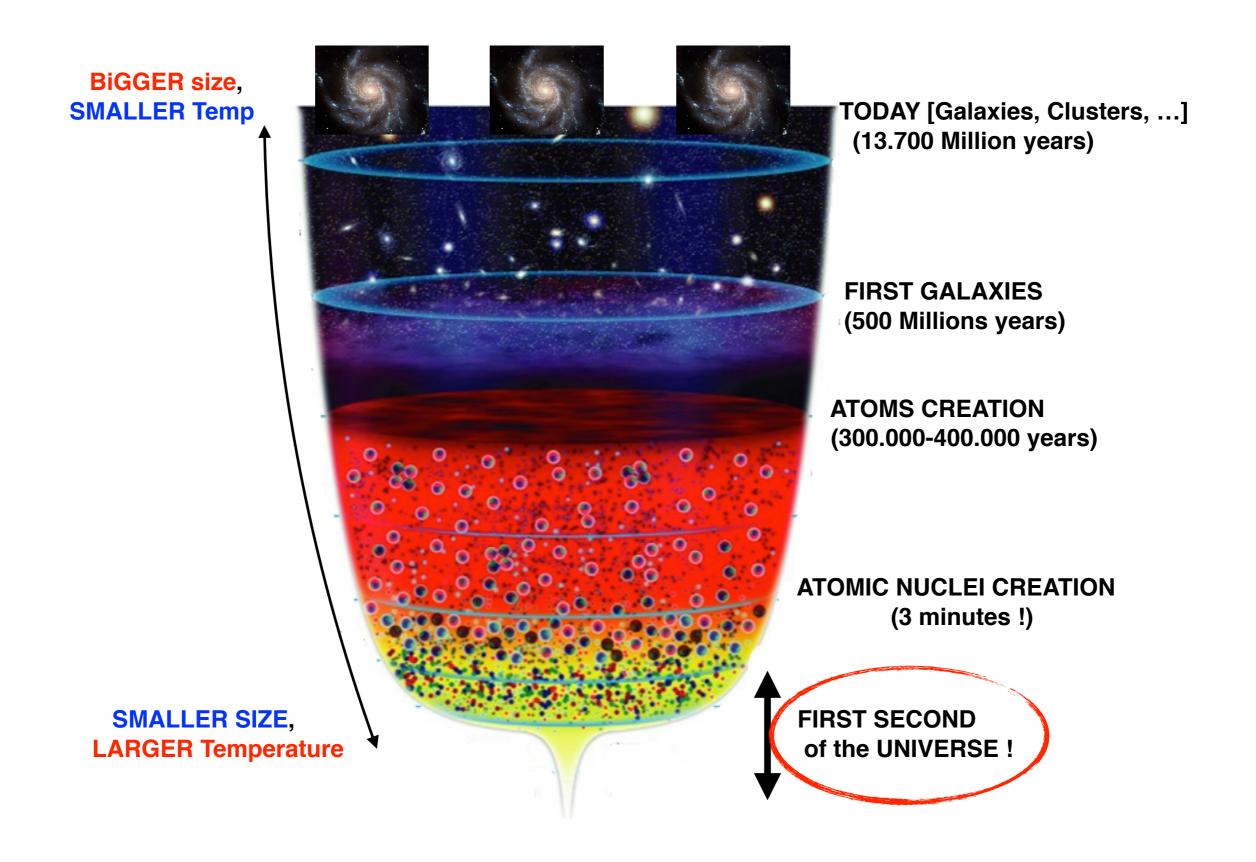
Eom:
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Anisotropic Stress

$$\Pi_{ij} = T_{ij} - \left\langle T_{ij} \right\rangle_{\mathsf{FLRW}}$$

GW Source(s): (SCALARS , VECTOR , FERMIONS)
$$\Pi_{ij}^{TT} \propto \{\partial_i \chi^a \partial_j \chi^a\}^{TT}, \{E_i E_j + B_i B_j\}^{TT}, \{\bar{\psi} \gamma_i D_j \psi\}^{TT}$$

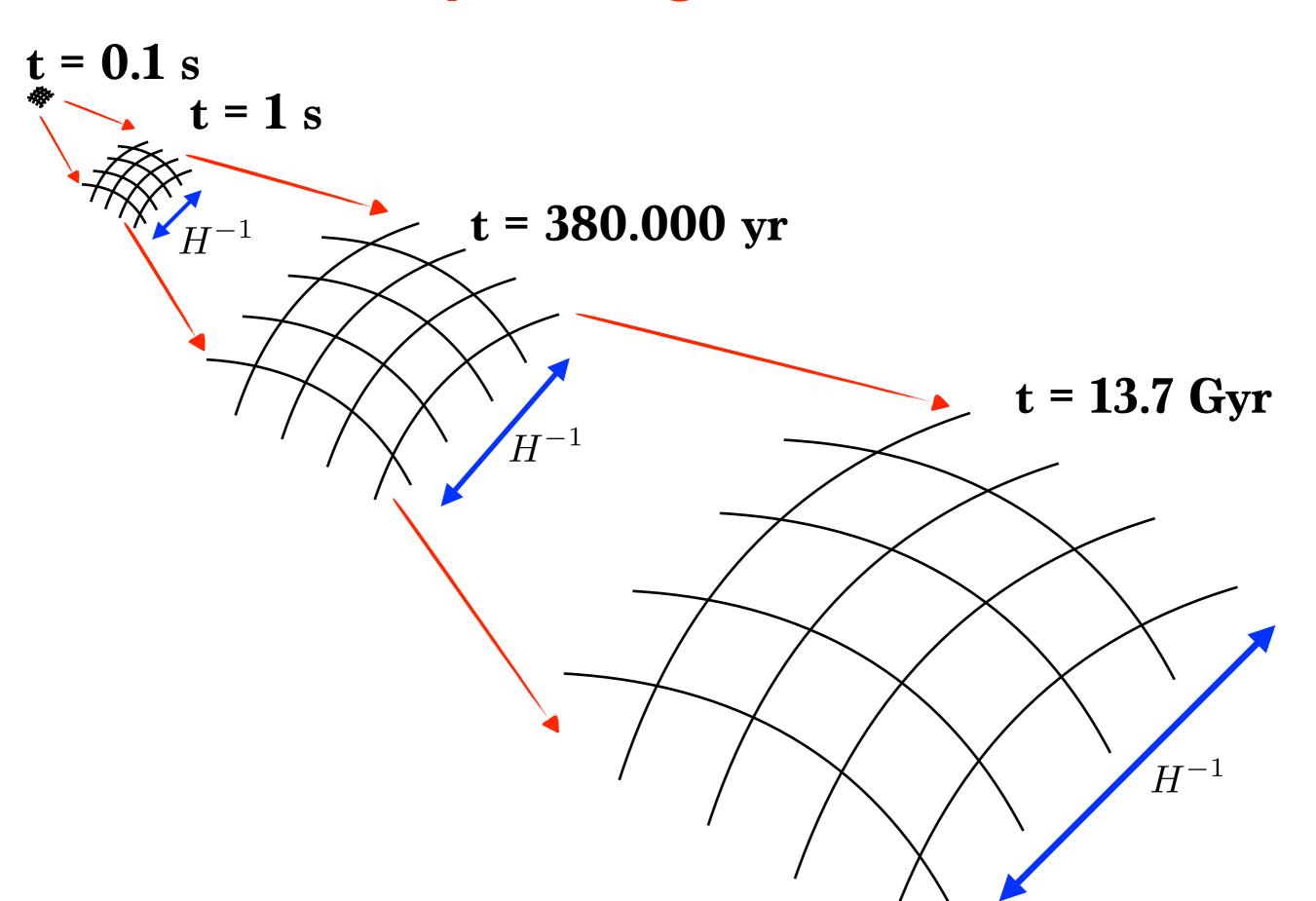
Cosmic History



To Be Continued ...

BACK SLIDES

FLRW



H & I

$$T^{\mu}_{\nu} \equiv \operatorname{diag}(-\rho, p, p, p)$$



$$m_p^2 G^{\mu}_{\nu} \left[g_{**}^{(FRW)} \right] = T^{\mu}_{\nu}$$

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Friedmann Equations

$$\frac{1}{a}\frac{d^2a}{dt^2} = -\frac{\rho}{6m_p^2}(1+3w)$$

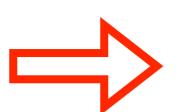
$$\left(w \equiv \frac{p}{\rho}\right)$$
 Equation of State (EoS)

H & I

$$T^{\mu}_{\nu} \equiv \operatorname{diag}(-\rho, p, p, p)$$



$$m_p^2 G^{\mu}_{\nu} \left[g_{**}^{(FRW)} \right] = T^{\mu}_{\nu}$$



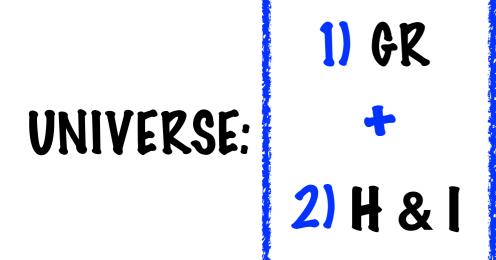
Friedmann Equations

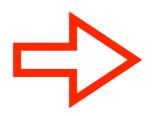
$$H^2 \equiv \left(\frac{1}{a}\frac{da}{dt}\right)^2 = \frac{\rho}{3m_p^2} - \frac{K}{a^2}$$

$$\left(w \equiv \frac{p}{\rho}\right)$$
 Equation of State (EoS)

(I)+(II)
$$\frac{1}{\rho}\frac{d\rho}{dt} = -\frac{3}{a}\frac{da}{dt}(1+w)$$
 (III)

Friedmann Equations





$$\frac{1}{a}\frac{d^{2}a}{dt^{2}} = -\frac{\rho}{6m_{p}^{2}}(1+3w)$$
 (I)
$$H^{2} \equiv \left(\frac{1}{a}\frac{da}{dt}\right)^{2} = \frac{\rho}{3m_{p}^{2}} - \frac{K}{a^{2}}$$
 (II)

$$H^2 \equiv \left(\frac{1}{a}\frac{da}{dt}\right)^2 = \frac{\rho}{3m_p^2} - \frac{K}{a^2}$$

$$\frac{1}{\rho}\frac{d\rho}{dt} = -\frac{3}{a}\frac{da}{dt}(1+w) \quad \text{(III)}$$

$$\left(w \equiv \frac{p}{\rho}\right)$$
 Equation of State (EoS)

(II)
$$H^2 \equiv \left(\frac{1}{a}\frac{da}{dt}\right)^2 = \frac{\rho}{3m_p^2} - \frac{K}{a^2} \longrightarrow \begin{bmatrix} \rho_c \equiv 3m_p^2 H^2 \end{bmatrix}$$

Critical density $(\rho = \rho_c \Leftrightarrow K = 0)$

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Critical density $(\rho = \rho_c \Leftrightarrow K = 0)$

$$\rho = \sum_{i} \rho_{i} ; \quad \Omega_{i} \equiv \frac{\rho_{i}}{\rho_{c}} \quad \Longrightarrow \quad \Omega \equiv \frac{\rho}{\rho_{c}} = \sum_{i} \Omega_{i} \quad \Longrightarrow \quad \Omega - 1 \equiv \frac{k}{a^{2}H^{2}}$$

Cosmic Sum

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Cosmic Sum

$$\begin{cases} \Omega > 1 \Rightarrow \operatorname{Close}(k > 0) \\ \Omega = 1 \Rightarrow \operatorname{Flat}(k = 0) \\ \Omega < 1 \Rightarrow \operatorname{Open}(k < 0) \end{cases}$$
 one-to-one correlation

(II)
$$H^2 \equiv \left(\frac{1}{a}\frac{da}{dt}\right)^2 = \frac{\rho}{3m_p^2} - \frac{K}{a^2}$$
 \longrightarrow $\rho_c \equiv 3m_p^2H^2$ Critical density $(\rho = \rho_c \Leftrightarrow K = 0)$

$$\rho = \sum_{i} \rho_{i} \; ; \; \Omega_{i} \equiv \frac{\rho_{i}}{\rho_{c}} \; \Longrightarrow \; \Omega \equiv \frac{\rho}{\rho_{c}} = \sum_{i} \Omega_{i} \; \Longrightarrow \; \Omega - 1 \equiv \frac{k}{a^{2}H^{2}}$$

Cosmic Sum

(III)
$$\frac{1}{\rho} \frac{d\rho}{dt} = -\frac{3}{a} \frac{da}{dt} (1+w) \implies \rho \propto e^{-3\int \frac{da}{a}(1+w)} = \begin{cases}
1/a^3, & \text{Mat.}(w=0) \\
1/a^4, & \text{Rad.}(w=1/3) \\
& \text{const.}, & \text{C.C.}(w=-1)
\end{cases}$$

(II)
$$H^2 \equiv \left(\frac{1}{a}\frac{da}{dt}\right)^2 = \frac{\rho}{3m_p^2} - \frac{K}{a^2} \longrightarrow \frac{\rho_c \equiv 3m_p^2 H^2}{\text{Critical density}} \quad (\rho = \rho_c \Leftrightarrow K = 0)$$

$$\rho = \sum_{i} \rho_{i} \; ; \quad \Omega_{i} \equiv \frac{\rho_{i}}{\rho_{c}} \quad \Longrightarrow \quad \Omega \equiv \frac{\rho}{\rho_{c}} = \sum_{i} \Omega_{i} \quad \Longrightarrow \quad \Omega - 1 \equiv \frac{k}{a^{2}H^{2}}$$

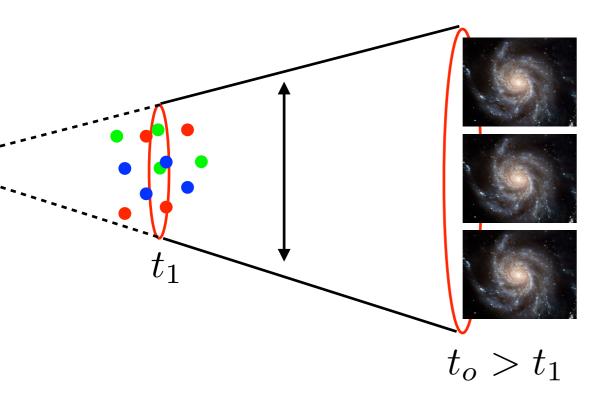
Cosmic Sum

(III) + (II):

$$H^{2}(a) = H_{o}^{2} \left\{ \Omega_{R}^{(o)} \left(\frac{a_{o}}{a} \right)^{4} + \Omega_{M}^{(o)} \left(\frac{a_{o}}{a} \right)^{3} + \Omega_{k}^{(o)} \left(\frac{a_{o}}{a} \right)^{2} + \Omega_{DE}^{(o)} e^{-3 \int \frac{da}{a} (1+w)} \right\}$$

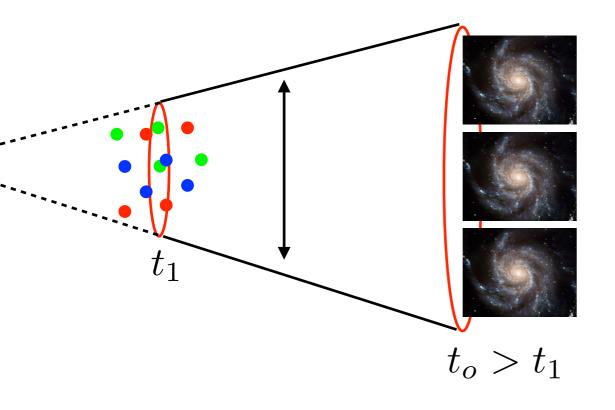
$$\equiv H_{o}^{2} E^{2}(a)$$

$$E(a) \equiv \sqrt{\Omega_{\rm R}^{(o)} \left(\frac{a_o}{a}\right)^4 + \Omega_{\rm M}^{(o)} \left(\frac{a_o}{a}\right)^3 + \Omega_{\rm k}^{(o)} \left(\frac{a_o}{a}\right)^2 + \Omega_{\rm DE}^{(o)} e^{-3\int \frac{da}{a}(1+w)} \qquad \Omega_{\rm k}^{(o)} \equiv -\frac{k}{a_o^2 H_o^2}}$$



Past: particle ensemble

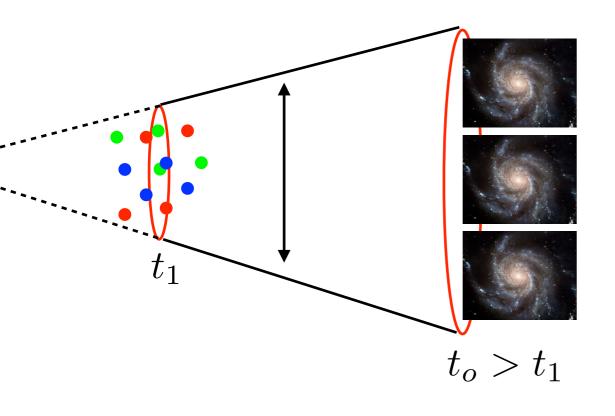
Statistical Mechanics



Past: particle ensemble

Statistical Mechanics

(III)
$$\frac{d\rho}{dt} + 3H(\rho + p) = 0 \longrightarrow \frac{dU}{dt} + p\frac{dV}{dt} = 0, \qquad \begin{cases} U = a^3\rho, \\ V = a^3 \end{cases}$$

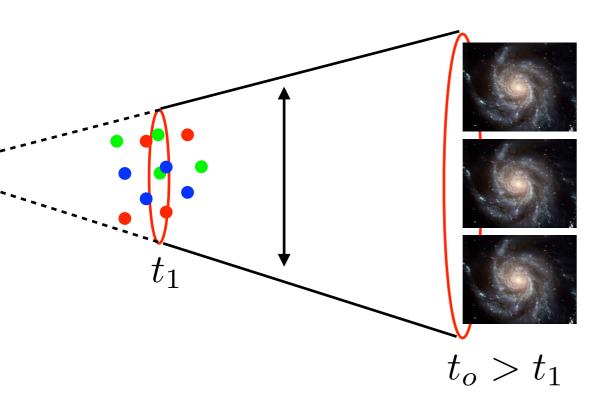


Past: particle ensemble

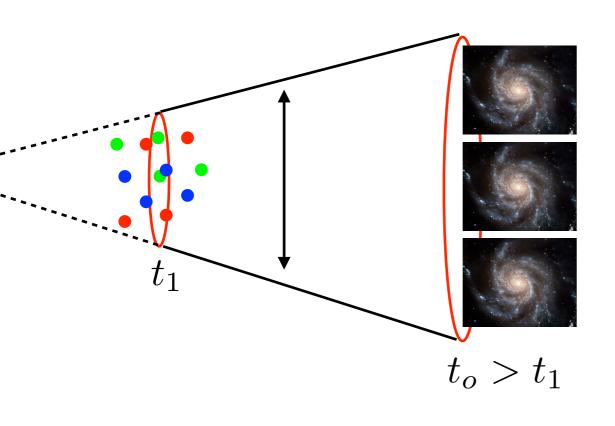
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$$\begin{cases} \frac{dU}{dt} + p\frac{dV}{dt} = T\frac{dS}{dt}, & \longrightarrow \text{ Thermal Eq.} \\ \frac{dS}{dt} = 0, & \longrightarrow \text{ Adiabatic Exp.} \end{cases}$$

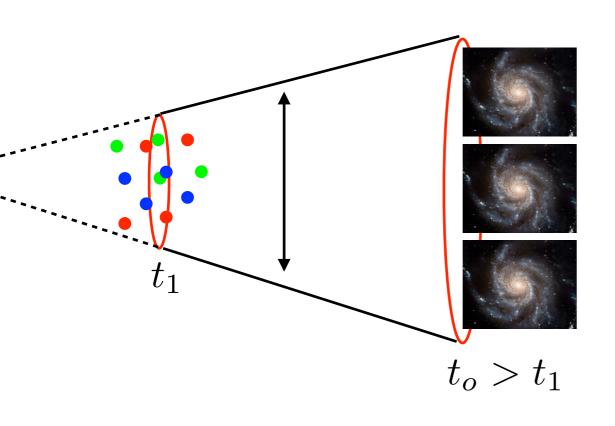


Thermal Eq. (densities)
$$\begin{cases} n = g_* \int d\vec{p} \, f(\vec{p}) \,, & \text{number} \\ \rho = g_* \int d\vec{p} \, E(\vec{p}) f(\vec{p}) \,, & \text{energy} \\ p = g_* \int d\vec{p} \, \frac{|\vec{p}|^2}{3E(\vec{p})} f(\vec{p}) \,, & \text{pressure} \\ dof & \text{Dispersion} & \text{Statistical} \\ relation & \text{Distribution} \end{cases}$$



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Bose-Einstein / Fermi-Dirac: $f(\vec{p}) = \left(e^{E(\vec{p})/T} \pm 1\right)^{-1}, \begin{cases} F(+) & \text{[fermions]} \\ B(-) & \text{[bosons]} \end{cases}$

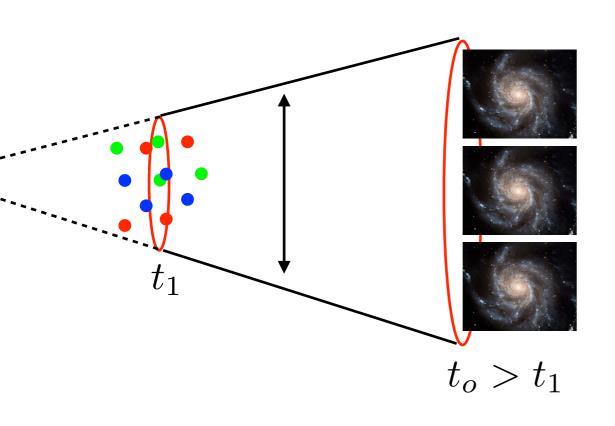


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.....

$$\frac{\rho_R \propto 1/a^4}{\rho_M \propto 1/a^3} \propto 1/a \,, \quad \Rightarrow \quad z \ge z_{\rm EQ} \, \left(t \le t_{\rm EQ} \right), \quad \rho_R > \rho_M$$

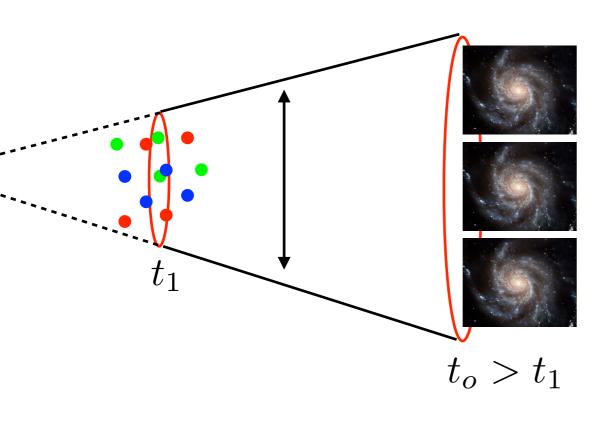


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Past: Radiation Domination !
$$1 + z_{\rm EQ} = \Omega_{\rm M}^{(o)}/\Omega_{\rm Rad}^{(o)} \sim 3400$$

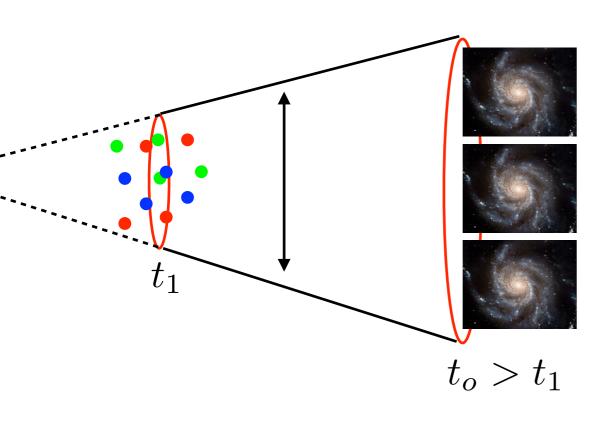


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Past: Radiation Domination (RD)

$$\rho_R^{(i)} = f_i g_*^{(i)} \frac{\pi^2}{30} T_i^4, \quad f_i = \begin{cases} 1, & B \\ \frac{7}{8}, & F \end{cases}$$

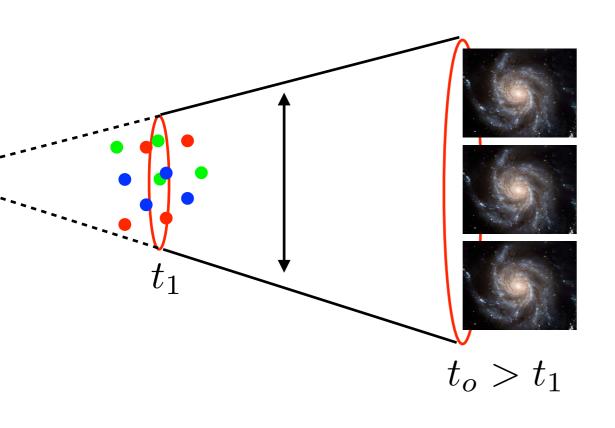


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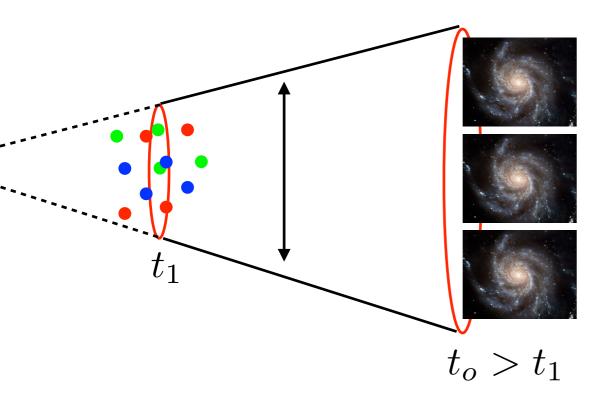
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$$\rho_R = \sum_i \rho_R^{(i)} \equiv g_*(T) \frac{\pi^2}{30} T^4$$

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$$g_*(T) \equiv \sum_{i} g_{*,i}^{(B)} \left(\frac{T_i}{T}\right)^4 + \frac{7}{8} \sum_{i} g_{*,i}^{(F)} \left(\frac{T_i}{T}\right)^4$$

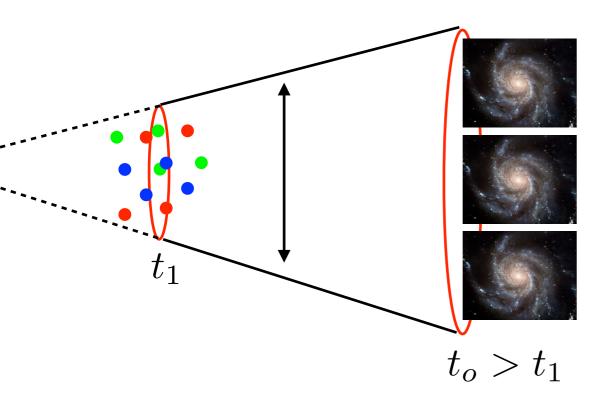


Adiabatic Exp:

$$S = \frac{a^3(\rho + p)}{T} = const.$$

$$a^3 T^3 g_*^{(s)}(T) = const.$$

$$g_*^{(s)}(T) \equiv \sum_i g_{*,i}^{(B)} \left(\frac{T_i}{T}\right)^3 + \frac{7}{8} \sum_i g_{*,i}^{(F)} \left(\frac{T_i}{T}\right)^3$$



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$$S = \frac{a^3(\rho + p)}{T} = const.$$

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When do $g_*(T), g_*^{(s)}(T)$ change ? $\begin{cases} 1 \text{ Species Decoupling }, T \to T_i, \\ 2 \text{ Mass threshold }, T < 2m_i, \end{cases}$