Quantum variance for automorphic forms

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(Based on joint work with Stephen Lester)

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L-functions, Circle-Method and Applications (HYBRID)
@ ICTS

Contents

Quantum chaos: general theory

Classical dynamics:

- M = a (compact) Riemann surface.
- $T^1 M$ = its unit tangent bundle.
- $\Phi^t : T^1 M \to T^1 M$ the geodesic flow.

Quantum dynamics:

- $\{\psi_j\}$ = an orthonormal basis of $L^2(M)$ consisting of the eigenfunctions of the Laplacian ($\Delta\psi_j = \lambda_j \psi_j$, with $\lambda_j = 1/4 + t_j^2$).
- Weyl's law: $\#\{t_j \leq T\} \sim c_M \cdot T^2$.

Random wave conjecture: Michael Berry (1977) suggested eigenfunctions for <u>chaotic</u> systems are modeled by random waves in the semiclassical limit, that is, its distribution should be <u>Gaussian</u>.

- The semiclassical limit $\Longleftrightarrow$ eigenvalues tend to infinity.
- Chaotic = ergodic + exponential divergence of orbits ...

Fluctuations of matrix elements

- $a \in C^\infty(M)$ an observable.
- $\mathrm{Op}(a) : L^2(M) \to L^2(M)$ a quantization (the multiplication operator).
- $\{\psi_j\}$ an orthonormal basis of the eigenfunctions of the Laplacian.
- Matrix elements: $\langle \mathrm{Op}(a)\psi_j, \psi_j \rangle = \langle a, |\psi_j|^2 \rangle = \int_M a|\psi_j|^2$.

Mean value of matrix elements = classical average (local Weyl law):

$$\frac{1}{N(T)} \sum_{t_j \asymp T} \langle \mathrm{Op}(a)\psi_j, \psi_j \rangle \sim \int_M a, \quad \text{where } N(T) = \sum_{t_j \asymp T} 1.$$

Quantum ergodicity theorem (Schnirelman 1974; Colin de Verdiere 1985; Zelditch 1987)

If the geodesic flow of M is ergodic, then the variance of the matrix elements vanishes, i.e.,

$$\mathrm{Var}(T) := \frac{1}{N(T)} \sum_{t_j \asymp T} \left| \langle \mathrm{Op}(a)\psi_j, \psi_j \rangle - \int_M a \right|^2 \to 0, \quad T \to \infty.$$

Quantum fluctuation conjecture

Quantum fluctuation conjecture (Feingold–Peres 1986, Eckhart–Fishman–Keating–Agam–Main–Müller 1995)

Normalize $\int_M a = 0$. If the geodesic flow is chaotic then generically

1) we have

$$\mathrm{Var}(T) \sim \mathrm{V}_{\mathrm{cl}}(a)/T^{\dim(M)-1},$$

where $\mathrm{V}_{\mathrm{cl}}(a) := \int_{\mathbb{R}} \langle a \circ \Phi^t, a\rangle_M \mathrm{d}t$ is the average auto-correlation of classical observable a.

2) The normalized matrix elements

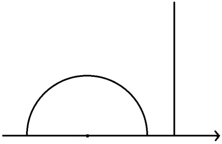
$$F_j(a) := t_j^{\frac{\dim(M)-1}{2}} \langle \mathrm{Op}(a)\psi_j, \psi_j\rangle$$

have a Gaussian distribution with mean zero and variance $\mathrm{V}_{\mathrm{cl}}(a)$.

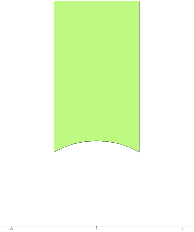
Hyperbolic surfaces

- $\mathbb{H} = \{z = x + iy : y > 0\}$ the upper half-plane with measure $\mathrm{d}\mu z = \mathrm{d}x\mathrm{d}y/y^2$.
- $\Delta = -y^2 \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right)$ the Laplace operator.
- Gauss curvature on $\mathbb{H}$ is negative $(= -1)$.
- Geodesics are semicircles subtended on $y = 0$ and vertical lines.

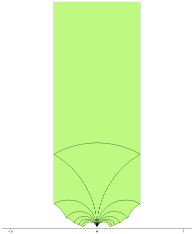
- $\Gamma = \mathrm{SL}_2(\mathbb{Z})$ and $\Gamma_0(D) = \{\gamma = \left(\begin{smallmatrix} a & b \\ c & d \end{smallmatrix}\right) \in \mathrm{SL}_2(\mathbb{Z}) : c \equiv 0 \pmod{D}\}$.
  $\Gamma \curvearrowright \mathbb{H}$ as fractional linear transforms.
- $\mathbb{X} = \Gamma \backslash \mathbb{H}$ or $\Gamma_0(D) \backslash \mathbb{H}$ a hyperbolic surface.
- Gauss curvature on $\mathbb{X}$ is negative $\Rightarrow$ The geodesic flow on $T^1\mathbb{X}$ is chaotic.



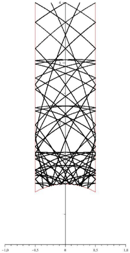
Geodesics in $\mathbb{H}$



$\mathrm{SL}_2(\mathbb{Z}) \backslash \mathbb{H}$



$\Gamma_0(13) \backslash \mathbb{H}$



A projected geodesic

Hecke–Maass forms

A *cuspidal Hecke–Maass newform* ϕ of level D with a nebentypus character χ of modulus D satisfies the automorphy condition

$$\phi(\gamma z) = \chi(d)\phi(z), \quad \gamma = \left(\begin{smallmatrix} a & b \\ c & d \end{smallmatrix}\right) \in \Gamma_0(D), \quad z \in \mathbb{H},$$

and is an eigenfunction of the Laplace operator Δ with eigenvalue λ_ϕ, and of the Hecke operators. Define the spectral parameter t_ϕ by $\lambda_\phi = 1/4 + t_\phi^2$.

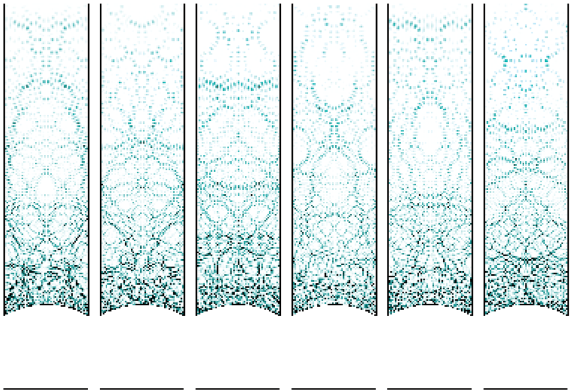
Weyl's law (Selberg):

$$\#\{\phi : |t_\phi| \leq T\} \sim \frac{\mathrm{vol}(\mathbb{X})}{4\pi} T^2.$$

Here we have $\mathrm{vol}(\mathbb{X}) = \frac{\pi}{3} D \prod_{p|D} (1 + p^{-1})$.

Let $\{u_j\}_{j=1}^{\infty}$ be an orthonormal basis of the cuspidal Hecke–Maass forms, which corresponding to the discrete spectrum.

Value distribution



This depicts the densities of a sequence of Maass forms on the hyperbolic surface.

Eisenstein series for $\mathrm{SL}_2(\mathbb{Z})$

- Eisenstein series:

$$E(z,s) := \sum_{\gamma\in\Gamma_\infty\backslash\,\mathrm{PSL}_2(\mathbb{Z})} (\operatorname{Im}\gamma z)^s = \frac{1}{2}\sum_{\substack{(c,d)\in\mathbb{Z}^2\\(c,d)=1}} \frac{y^s}{|cz+d|^{2s}}, \quad \mathrm{Re}(s)>1.$$

Here $\Gamma_\infty = \{\left(\begin{smallmatrix}1 & n\\ & 1\end{smallmatrix}\right) : n\in\mathbb{Z}\}$.

- Fourier expansion:

$$E(z,s) = y^s + \phi(s)y^{1-s} + \frac{2\sqrt{y}}{\pi^{-s}\Gamma(s)\zeta(2s)}\sum_{n\neq 0}\eta_{s-1/2}(n)K_{s-1/2}(2\pi|n|y)e(nx),$$

with $\phi(s) = \sqrt{\pi}\frac{\Gamma(s-1/2)\zeta(2s-1)}{\Gamma(s)\zeta(2s)}$, $\eta_s(n) = \sum_{ab=|n|,\ a>0}(a/b)^s$,

$$K_s(y) = \tfrac{1}{2}\int_0^\infty e^{-(u+1/u)y/2}u^{s-1}\mathrm{d}u, \text{ and } e(x) = e^{2\pi i x}.$$

- The Selberg spectral decomposition:

$$L^2(\Gamma\backslash\mathbb{H}) = \mathbb{C}\oplus L^2_{\mathrm{cusp}}(\Gamma\backslash\mathbb{H})\oplus L^2_{\mathrm{cont}}(\Gamma\backslash\mathbb{H}).$$

That is,

$$f(z) = \langle f,1\rangle\frac{3}{\pi} + \sum_{j\geq 1}\langle f,u_j\rangle u_j(z) + \frac{1}{4\pi i}\int_{(1/2)}\langle f,E(*,s)\rangle E(z,s)\mathrm{d}s.$$

Quantum Unique Ergodicity

For a test function $\psi : \mathbb{X} \to \mathbb{C}$, define

$$\mu_j(\psi) := \langle \psi, |u_j|^2 \rangle = \int_{\mathbb{X}} \psi(z) |u_j(z)|^2 \frac{\mathrm{d}x\mathrm{d}y}{y^2}, \quad \mu_t(\psi) := \langle \psi, |E(*, 1/2 + it)|^2 \rangle.$$

Quantum Unique Ergodicity (Rudnick–Sarnak conjecture 1994):

$$\mu_j(\psi) \sim \frac{3}{\pi} \int_{\mathbb{X}} \psi(z) \frac{\mathrm{d}x\mathrm{d}y}{y^2}, \quad \text{as } j \to \infty.$$

- Luo–Sarnak 1995:
$$\mu_t(\psi) \sim \frac{6}{\pi} \log t \int_{\mathbb{X}} \psi(z) \frac{\mathrm{d}x\mathrm{d}y}{y^2}, \quad \text{as } t \to \infty.$$
- Sarnak 2001 & Liu–Ye 2002: QUE holds for dihedral Maass forms.
- Lindenstrauss 2006 & Soundararajan 2010: QUE holds for Hecke–Maass cusp forms.
- Holowinsky and Soundararajan 2010: QUE holds for holomorphic Hecke eigenforms.

Quantum variance for cusp forms

Luo–Sarnak 1995, Zhao 2010, Sarnak–Zhao 2019, and Nelson 2016-2019 computed the quantum variance for the discrete spectrum. E.g.

Theorem (Luo–Sarnak 2004; Zhao 2010)

Define the quantum variance for cusp forms by

$$Q_C(\phi,\psi) := \lim_{T\to\infty} \frac{1}{T} \sum_{t_j \sim T} \mu_j(\phi)\overline{\mu_j(\psi)},$$

for fixed $\phi, \psi \in \{u_j\}$. Then we have

$$Q_C(\phi,\psi) = \begin{cases} c(\phi)L(1/2,\phi)\mathrm{V}_{\mathrm{cl}}(\phi), & \text{if } \phi = \psi \text{ is even,} \\ 0, & \text{otherwise,} \end{cases}$$

with the classical variance

$$\mathrm{V}_{\mathrm{cl}}(\phi) = \frac{\left|\Gamma(\frac{1}{4} + \frac{it_\phi}{2})\right|^4}{2\pi|\Gamma(\frac{1}{2} + it_\phi)|^2}.$$

Tools: Poincaé series, trace formulas, Hecke operators, ...

Quantum variance for Eisenstein series

The continuous spectrum for $\mathrm{SL}_2(\mathbb{Z})\backslash\mathbb{H}$ is parametrized by Eisenstein series $E(z,s)$. Recall that $\mu_t(\psi) = \langle \psi, |E(*, 1/2+it)|^2\rangle$.

Theorem (H. 2021)

Define the quantum variance for Eisenstein series by

$$Q_E(\phi,\psi) := \lim_{T\to\infty} \frac{1}{\log T} \int_T^{2T} \mu_t(\phi)\overline{\mu_t(\psi)}\mathrm{d}t,$$

for $\phi, \psi \in \{u_j\}$. Then we have

$$Q_E(\phi,\psi) = \begin{cases} C(\phi)L(1/2,\phi)^2\mathrm{V}_{\mathrm{cl}}(\phi), & \text{if } \phi=\psi \text{ is even,} \\ 0, & \text{otherwise,} \end{cases}$$

where $C(\phi)$ is an explicit constant depending on ϕ.

Remark:

- In fact, we proved asymptotic formula with quantitative error terms.
- Rudnick–Soundararajan (2005) showed higher moments blow up (no CLT).

Hecke Grössencharacters

$D > 0$ squarefree and $D \equiv 1 \pmod 4$.

$F = \mathbb{Q}(\sqrt{D})$ be a fixed real quadratic fields with discriminant D.

For simplicity, we assume that F has the narrow class number 1, and D is a product of two distinct primes congruent to $3 \pmod 4$. For example $D = 21$.

$\omega_D = \frac{1+\sqrt{D}}{2}$.
$\epsilon_D > 1$ the fundamental unit of F.
$\mathcal{O}_F = \mathbb{Z}[\omega_D]$ the ring of integers of F.
$U_F = \{\pm 1\} \times \epsilon_D^{\mathbb{Z}}$ the group of units.

For integer $k \neq 0$, we have the **Hecke Grössencharacter** Ξ_k of F defined by

$$\Xi_k((\alpha)) := \left|\frac{\alpha}{\tilde{\alpha}}\right|^{\frac{\pi i k}{\log \epsilon_D}} \quad \text{for ideal } (\alpha) \subset \mathcal{O}_F \text{ with generator } \alpha,$$

where $\tilde{\alpha}$ is the conjugate of α under the nontrivial automorphism of F.

Dihedral Maass forms

Let $\mathcal{B}_0^*(D, \chi_D)$ denote the set of L^2-normalized newforms of weight 0 for $\Gamma_0(D)$, with nebentypus character χ_D (the Kronecker symbol). Maass showed that the theta-like series associated to Ξ_k by

$$\phi_k(z) := \rho_k(1)\, y^{1/2} \sum_{\substack{\mathfrak{a} \subset \mathcal{O}_F \\ \mathfrak{a} \neq \{0\}}} \Xi_k(\mathfrak{a}) K_{it_k}(2\pi \,\mathrm{N}(\mathfrak{a}) y)\big(e(\mathrm{N}(\mathfrak{a})x) + e(-\mathrm{N}(\mathfrak{a})x)\big) \in \mathcal{B}_0^*(D, \chi_D),$$

where $z = x + iy \in \mathbb{H}$, $t_k := t_{\phi_k} = \pi k / \log \epsilon_D$ and ϕ_k has Laplace eigenvalue $1/4 + t_k^2$. Here $\mathrm{N}(\mathfrak{a}) = \#\mathcal{O}_F/\mathfrak{a}$ is the norm of a nonzero ideal $\mathfrak{a} \subset \mathcal{O}_F$, $K_s(z)$ is the modified Bessel function, and $\rho_k(1)$ is the positive real number such that ϕ_k is L^2-normalized, i.e.,

$$\|\phi_k\|_2^2 = \langle \phi_k, \phi_k \rangle_D = \int_{\Gamma_0(D)\backslash\mathbb{H}} |\phi_k(z)|^2 \frac{\mathrm{d}x\mathrm{d}y}{y^2} = 1.$$

Weyl's law: Let $t_k := \frac{\pi k}{\log \epsilon_D}$, then

$$\#\{\phi_k : 0 < t_k \leq T\} \sim \frac{\log \epsilon_D}{\pi} T.$$

Quantum variance for dihedral Maass forms

Define

$$\mu_k(\psi) := \langle \psi, |\phi_k|^2 \rangle.$$

Define the quantum covariance for dihedral Maass forms by

$$Q(\psi_1, \psi_2; K; \Phi) := \sum_{k \in \mathbb{Z}} \mu_k(\psi_1) \overline{\mu_k(\psi_2)} \Phi\left(\frac{k}{K}\right)$$

for $\psi_1, \psi_2 \in L^2_{\mathrm{cusp}}(\mathbb{X})$.

Define the harmonic weighted quantum covariance by

$$Q^{\mathrm{h}}(\psi_1, \psi_2; K; \Phi) := \sum_{k \in \mathbb{Z}} L(1, \phi_{2k})^2 \mu_k(\psi_1) \overline{\mu_k(\psi_2)} \Phi\left(\frac{k}{K}\right),$$

where $L(s, \phi_{2k})$ is the L-function of ϕ_{2k}.

Quantum variance for dihedral Maass forms

Theorem (H.–Lester 2020)

Let ψ be an even Hecke–Maass cuspidal newform on $\Gamma_0(D)$. Then as $K \to \infty$ we have that

$$Q^{\mathrm{h}}(\psi, \psi; K; \Phi) = \widetilde{\Phi}(0) A^{\mathrm{h}}(\psi)\, L(\tfrac{1}{2}, \psi) L(\tfrac{1}{2}, \psi \times \chi_D) \mathrm{V}_{\mathrm{cl}}(\psi) + o(1),$$

where

$$A^{\mathrm{h}}(\psi) = \frac{\pi \log \epsilon_D}{2D^2 \zeta_D(2) L(1, \chi_D)} \left(1 + \frac{\lambda_\psi(p_1)}{\sqrt{p_1}} + \frac{\lambda_\psi(p_2)}{\sqrt{p_2}} + \frac{\lambda_\psi(D)}{\sqrt{D}}\right).$$

Assume the Generalized Ramanujan Conjecture (GRC). Then as $K \to \infty$ we have that

$$Q(\psi, \psi; K; \Phi) = \widetilde{\Phi}(0) A(\psi)\, L(\tfrac{1}{2}, \psi) L(\tfrac{1}{2}, \psi \times \chi_D) V(\psi) + o(1),$$

where $A(\psi) = A^{\mathrm{h}}(\psi) C'_{D,\psi}$, with an explicit $C'_{D,\psi}$.

Application: If $A^{\mathrm{h}}(\psi)\, L(\frac{1}{2}, \psi) L(\frac{1}{2}, \psi \times \chi_D) \neq 0$, then $\mu_k(\psi) = \Omega(k^{-1/2-\varepsilon})$.
Conjecture: $\mu_k(\psi) = O(k^{-1/2+\varepsilon})$. GRH implies this conjecture.

Quantum covariance for dihedral Maass forms

Theorem (H.–Lester 2020)

Assume Generalized Riemann Hypothesis (GRH). Let ψ_1, ψ_2 be two orthogonal even Hecke–Maass cuspidal newforms. Then we have as $K \to \infty$ that

$$Q(\psi_1, \psi_2; K; \Phi) \longrightarrow 0.$$

In particular, the quadratic form $Q = \lim_{K \to \infty} Q(\cdot, \cdot; K; \Phi)$ is diagonalized by the orthonormal basis of Hecke–Maass cuspidal newforms on $\mathcal{B}_0^*(D)$.

Using the method of Rudnick–Soundararajan for lower bounds for moments of L-functions one can show the moments of $\mu_k(\psi)$ blow up.

Rankin–Selberg and Watson–Ichino

Let ϕ, u_j be Hecke–Maass forms of level 1, ϕ_k be a dihedral Maass form of level D, and ψ a Hecke–Maass newform of level D with trivial nebentypus.

- Rankin–Selberg method:

$$|\mu_t(\phi)|^2 = \frac{|\Lambda(1/2+2it,\phi)|^2\Lambda(1/2,\phi)^2}{2|\xi(1+2it)|^4\Lambda(1,\operatorname{sym}^2\phi)},$$

  where Λ means the corresponding completed L-functions and $\xi(s) = \pi^{-s/2}\Gamma(s/2)\zeta(s)$.

- Watson–Ichino formula:

$$|\mu_j(\phi)|^2 = \frac{\Lambda(1/2,\operatorname{sym}^2 u_j\times\phi)\Lambda(1/2,\phi)}{8\Lambda(1,\operatorname{sym}^2 u_j)^2\Lambda(1,\operatorname{sym}^2\phi)},$$

$$|\mu_k(\psi)|^2 = \frac{1}{8\sqrt{D}}\frac{\Lambda(\frac{1}{2},\psi)\Lambda(\frac{1}{2},\psi\times\chi_D)\Lambda(\frac{1}{2},\psi\times\phi_{2k})}{\Lambda(1,\operatorname{sym}^2\psi)\Lambda(1,\chi_D)^2\Lambda(1,\phi_{2k})^2}.$$

Quantum variances to moments of L-functions

By the Rankin–Selberg method, for quantum variance of Eisenstein series, we need to estimate the following moment of L-functions:

$$\frac{1}{T\log T}\int_0^\infty w\Big(\frac{t}{T}\Big)\frac{L(\frac{1}{2}-2it,\phi)L(\frac{1}{2}+2it,\psi)}{|\zeta(1+2it)|^4}\mathrm{d}t,$$

for some nice weight function w.

Lemma

Let $t \asymp T$ be large enough. Then for any $T^\varepsilon \ll x \ll T^B$, we have

$$\frac{1}{\zeta(1+2it)^2}=\sum_{k\le x^{1+\varepsilon}}\frac{\alpha(k)}{k^{1+2it}}e^{-k/x}+O(e^{-(\log T)^{1/5}}),$$

where $\alpha(k)=\sum_{mn=k}\mu(m)\mu(n)$.

Shifted Convolution Sums

Theorem (Harcos 2003)

Let f be a smooth function on $(\mathbb{R}_{>0})^2$ satisfying

$$x^i y^j f^{(i,j)}(x,y) \ll_{i,j} \left(1+\frac{x}{X}\right)^{-1}\left(1+\frac{y}{Y}\right)^{-1} P^{i+j},$$

with some $P, X, Y \geq 1$ for all $i, j \geq 0$. Let $\lambda_\phi(m)$ (resp. $\lambda_\psi(n)$) be the normalized Fourier coefficients of a holomorphic or Maass cusp form ϕ (resp. ψ) of arbitrary level and nebentypus. Define

$$D_f(k,\ell;h) = \sum_{km \pm \ell n = h} \lambda_\phi(m)\lambda_\psi(n) f(km, \ell n),$$

where k, ℓ, h are positive integers. Then for coprime k and ℓ, we have

$$D_f(k,\ell;h) \ll P^{11/10}(k\ell)^{-1/10}(X+Y)^{1/10}(XY)^{2/5+\varepsilon},$$

where the implied constant depends only on ε and the forms ϕ, ψ.

Tool: The Jutila circle method, Voronoi summation formula, ...

Proof ingredients: Moments of L-functions

Proposition (H.–Lester 2020)

Let ψ be an even Hecke–Maass cuspidal newform on $\Gamma_0(D)$ with trivial nebentypus and $\eta_\psi(D)$ denote the W_D-eigenvalue of ψ, where W_D is the Atkin-Lehner operator. Suppose $\eta_\psi(D) = 1$. Let w be a Schwartz function with compact support in $[\frac{1}{2}, 2]$ such that $w^{(j)}(x) \ll P^j$, where $P \geq 1$ is a large parameter. Then there exists $A_0 > 0$ such that

$$\sum_{k\in\mathbb{Z}} L(\tfrac{1}{2}, \psi \times \phi_{2k})\, w\left(\frac{k}{K}\right) = \tilde{w}(1) \cdot C_{D,\psi} \cdot K + O(P^{A_0} \cdot K^{\frac{1}{2}+\vartheta+\varepsilon}),$$

where the implied constant depends at most on ψ, D. Here ϑ is the bound toward the Ramanujan–Selberg conjecture and

$$C_{D,\psi} = 2 \cdot \frac{L(1,\chi_D)}{\zeta_D(2)} L(1, \operatorname{sym}^2 \psi) \left(1 + \frac{\lambda_\psi(p_1)}{\sqrt{p_1}} + \frac{\lambda_\psi(p_2)}{\sqrt{p_2}} + \frac{\lambda_\psi(D)}{\sqrt{D}}\right)$$

where $\zeta_D(s) = \zeta(s) \prod_{p|D}(1 - p^{-s})$. Recall that $\tilde{w}(s) := \int_0^\infty w(x) x^{s-1}\, dx$ is the Mellin transform of w.

Proof ingredients: Twisted moments of L-functions

Proposition (H.–Lester 2020)

Assume GRC. Suppose $\eta_\psi(D) = 1$. Then there exists $A_0 > 0$ such that for $n \in \mathbb{N}$

$$\sum_{k \in \mathbb{Z}} L(\tfrac{1}{2}, \psi \times \phi_{2k}) \cdot \lambda_{2k}(n)\, w\left(\frac{k}{K}\right) = \tilde{w}(1) \cdot C_{D,\psi} \cdot h\left(\frac{n}{(n, D)}\right) \cdot K + O((Pn)^{A_0} \cdot K^{\frac{1}{2}+\vartheta+\varepsilon}),$$

for certain multiplicative function h.

Proposition (H.–Lester 2020)

Assume GRC. Suppose $\eta_\psi(D) = 1$. Also, suppose $P \leq K^\delta$ for some $\delta > 0$ sufficiently small. Then for any $A \geq 1$ we have that

$$\sum_{k \in \mathbb{Z}} \frac{L(\frac{1}{2}, \psi \times \phi_{2k})}{L(1, \phi_{2k})^2} w\left(\frac{k}{K}\right) = \tilde{w}(1) \cdot C'_{D,\psi} \cdot C_{D,\psi} \cdot K + O\left(\frac{K}{(\log K)^A}\right)$$

where $C'_{D,\psi}$ is an explicit constant depending on D and ψ.

Proof ingredients: Non-split Sums

Let ψ be a Hecke-Maass form of level N and trivial nebentypus.
Define

$$\mathcal{S} := \sum_{n\geq 1} \lambda_\psi(an^2+bn+c) W\left(\frac{an^2+bn+c}{Y}\right)$$

where $a, b, c \in \mathbb{Z}$ $0 < |a| \ll Q$, $b \ll QR$, $c \ll QR^2$ and $\Delta := b^2 - 4ac > 0$, $W(x)$ is a smooth function with compact support in $[1, 2]$ and $W^{(j)}(x) \ll P^j$. Here the parameters P, Q, R satisfy $P, Q, R \leq Y^\delta$ for some $\delta > 0$ sufficiently small. WLOG we may assume $a > 0$.

Compare to the shifted convolution sums:

$$\sum_{n\geq 1} \lambda_\psi(n)\lambda_\psi(n+h) V\left(\frac{n}{N}\right).$$

Hooley 1963, Sarnak 1984, Blomer 2008, Templier 2011, Templier–Tsimerman 2013, ect, considered:

$$\sum_{n\geq 1} \lambda_\psi(n^2+d) W\left(\frac{n^2+d}{Y}\right).$$

Proof ingredients: Non-split Sums

Theorem (H.–Lester 2020)

Let ψ be a Hecke-Maass form of level N and trivial nebentypus. Each of the following holds:

(i) There exists $C_{\psi,a,b,c}$ and $B > 0$ such that

$$\mathcal{S} = C_{\psi,a,b,c}\tilde{W}(1/2)Y^{1/2} + O\left(Y^{1/4+\vartheta/2+\varepsilon}(PQR)^B\right),$$

where the implied constant depends on at most ψ and ε.

(ii) If $N = p_1p_2$ where $p_1 \equiv p_2 \equiv 3 \pmod 4$ are distinct primes, $a \mid N$ and $a \mid b$ then we have that $C_{\psi,a,b,c} = 0$.

(iii) Assume GRC. Then we have that $C_{\psi,a,b,c} = 0$.

In fact, $C_{\psi,a,b,c}$ is equal to

$$\frac{\pi^{1/4}\Delta^{1/4}}{2^{1/2}\varrho_\psi\varphi(a')}\sum_{\chi\ (a')}\frac{\pi^{\nu_\chi}d^{\nu_\chi}\bar{\chi}(b')}{\Gamma(1/2+\nu_\chi/2+it_\psi)\Gamma(1/2+\nu_\chi/2-it_\psi)}\sum_{f_j\in\mathbf{H}_\kappa(M,\chi_\nu,i/4)}\overline{\rho_j(\Delta)}\langle f_j,\overline{\Psi}\theta_{\chi,d^2}\rangle,$$

and $d = (2a,b)$, $a' = 2a/d$, $b' = b/d$, $\Psi(z) = \psi(4az) \in L^2_{\text{cusp}}(\Gamma_0(4aN)\backslash\mathbb{H})$.

Proof ingredients: Conditional upper bounds

Theorem (Soundararajan 2009)

Let $M_k(T) := \int_0^T |\zeta(1/2+it)|^{2k}\mathrm{d}t$. Assume RH. Then we have

$$T(\log T)^{k^2} \ll_k M_k(T) \ll_{k,\varepsilon} T(\log T)^{k^2+\varepsilon}.$$

Proposition (H.–Lester 2020)

Assume GRH. Let $n \geq 1$. Also, let $\psi_1, \ldots, \psi_n$ be pairwise orthogonal Hecke–Maass cuspidal newforms on $\Gamma_0(D)$ with trivial nebentypus. Then for any real numbers $\ell_1, \cdots, \ell_n > 0$ we have that

$$\sum_{K<k\leq 2K} L(\tfrac{1}{2}, \psi_1 \times \phi_{2k})^{\ell_1} \cdots L(\tfrac{1}{2}, \psi_n \times \phi_{2k})^{\ell_n} \ll K \cdot (\log K)^{\frac{\ell_1(\ell_1-1)}{2}+\cdots+\frac{\ell_n(\ell_n-1)}{2}+\varepsilon}.$$

$n=2,\ \ell_1=\ell_2=1/2 \quad \Rightarrow \quad$ Quantum covariance for dMF vanishes.

Thank you for your attention!