

ALESSANDRO BACCHETTA, PAVIA U. AND INFN

STATUS OF QUARK TMDS

A NEW ERA OF DISCOVERY

THE 2023 LONG RANGE PLAN FOR NUCLEAR SCIENCE

2023 | VERSION 1.2



A NEW ERA OF DISCOVERY

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RECOMMENDATION 3
We recommend the expeditious completion of the EIC as the highest priority for facility construction.

A NEW ERA OF DISCOVERY

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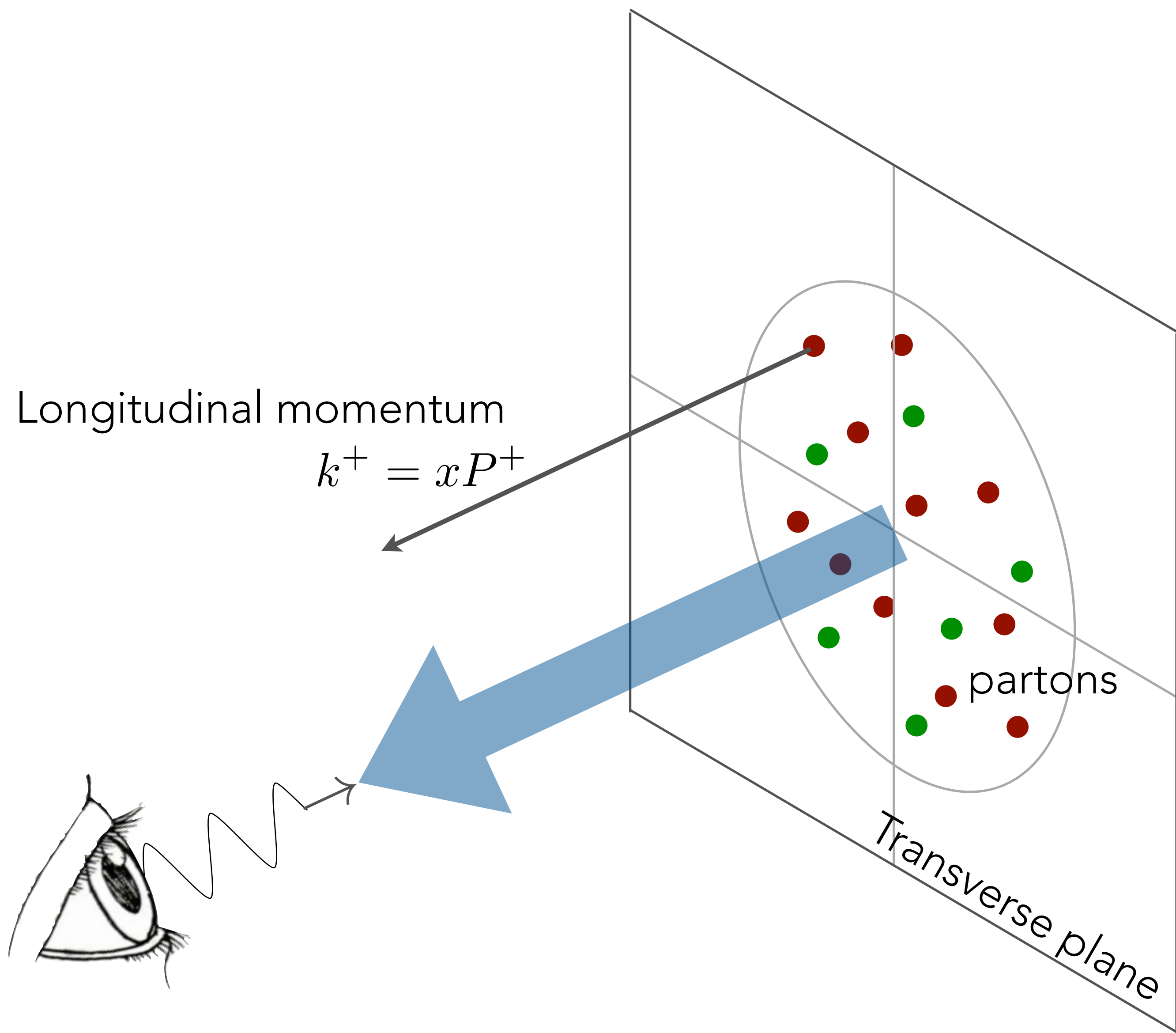


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“

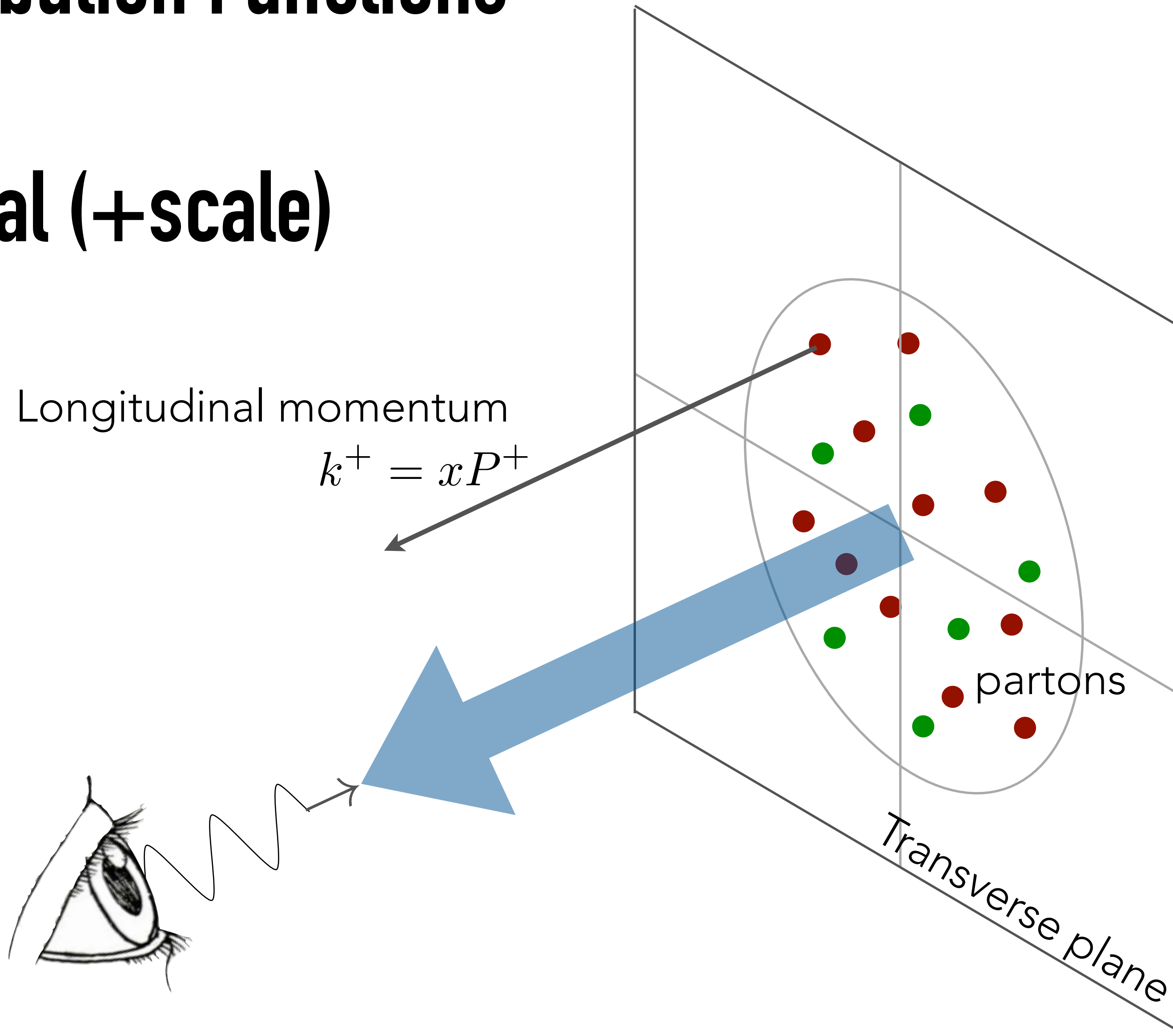
The EIC is a powerful discovery machine, a precision microscope capable of taking three-dimensional pictures of nuclear matter at femtometer scales.

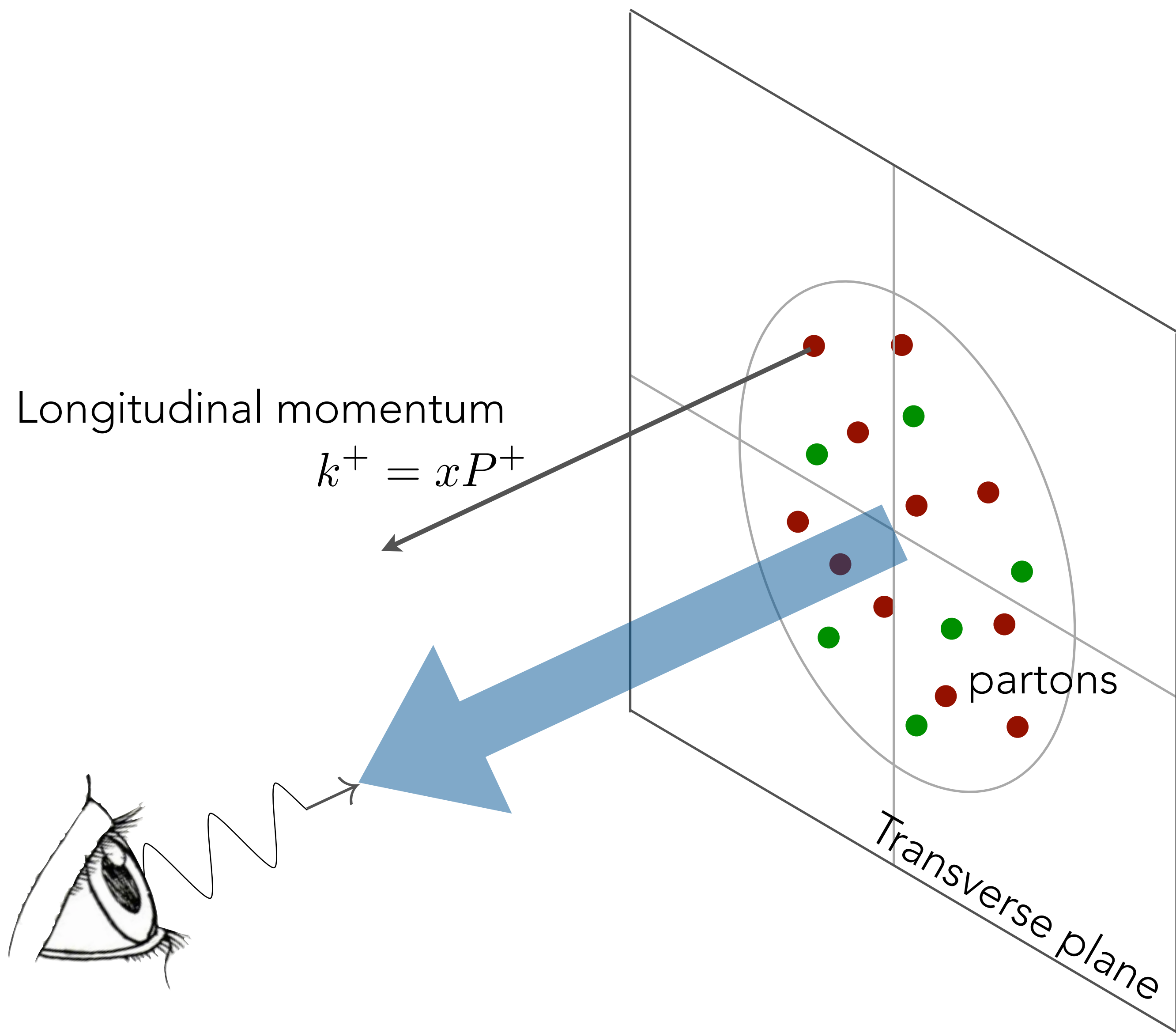


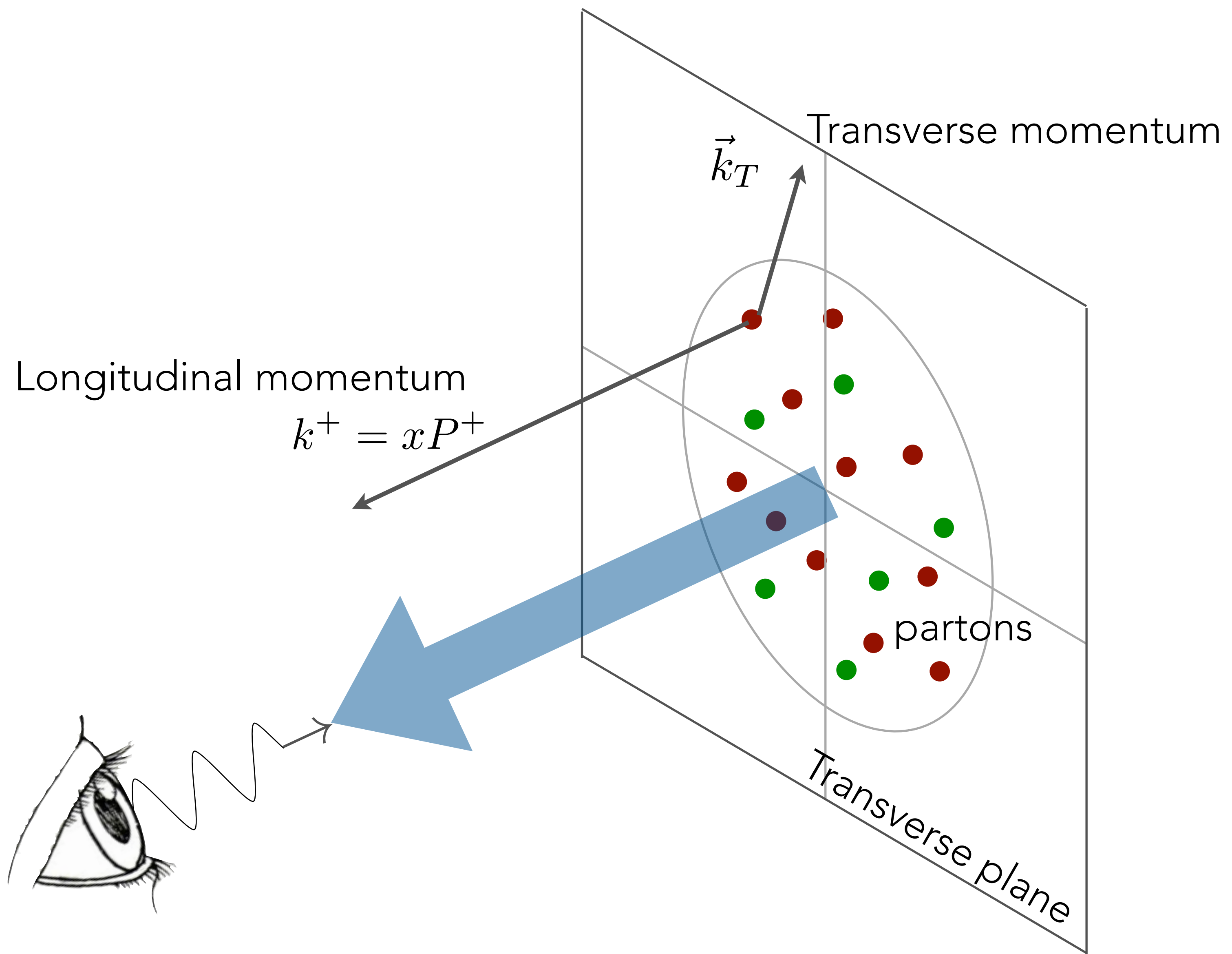
Parton Distribution Functions

$$f(x)$$

1 dimensional (+scale)



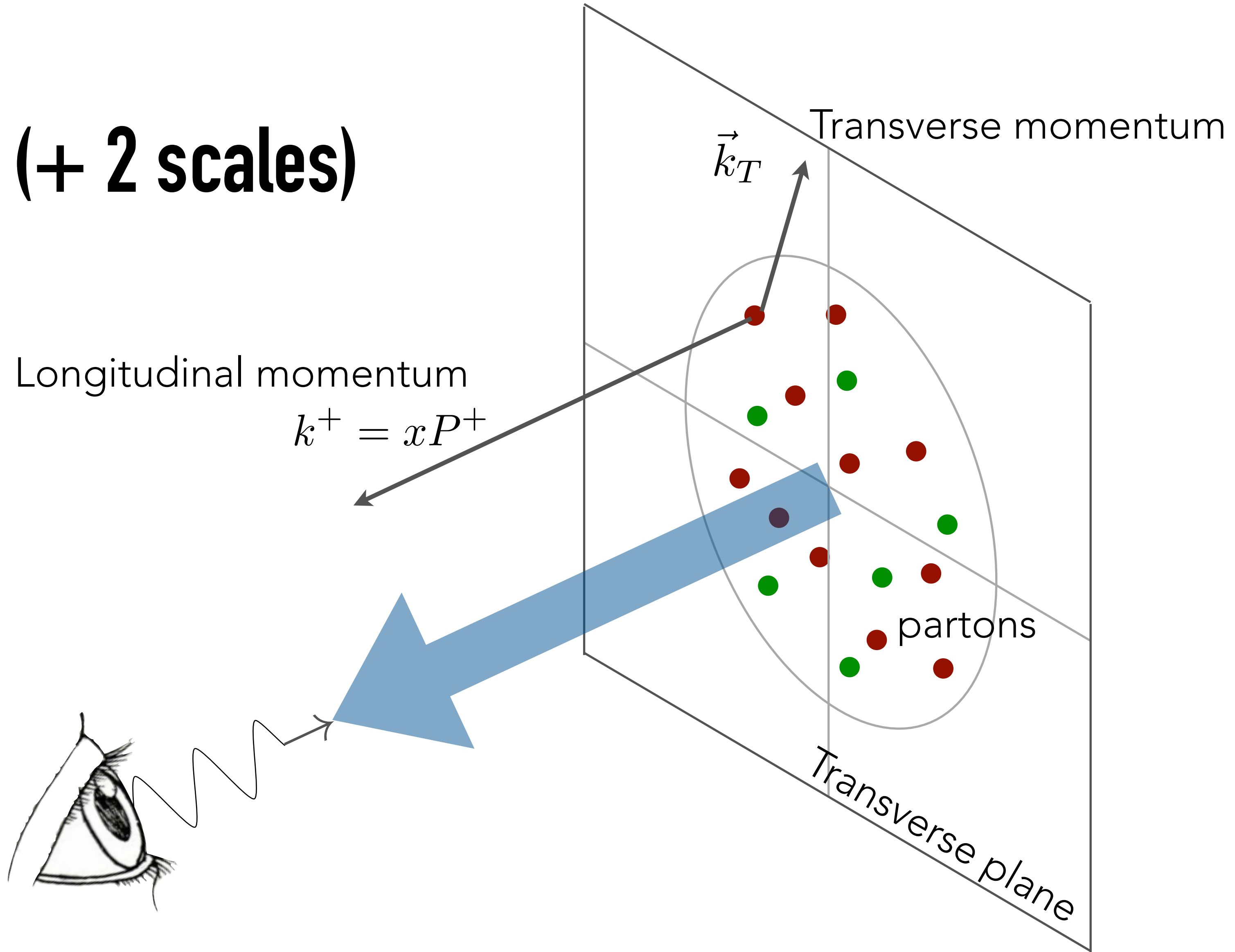


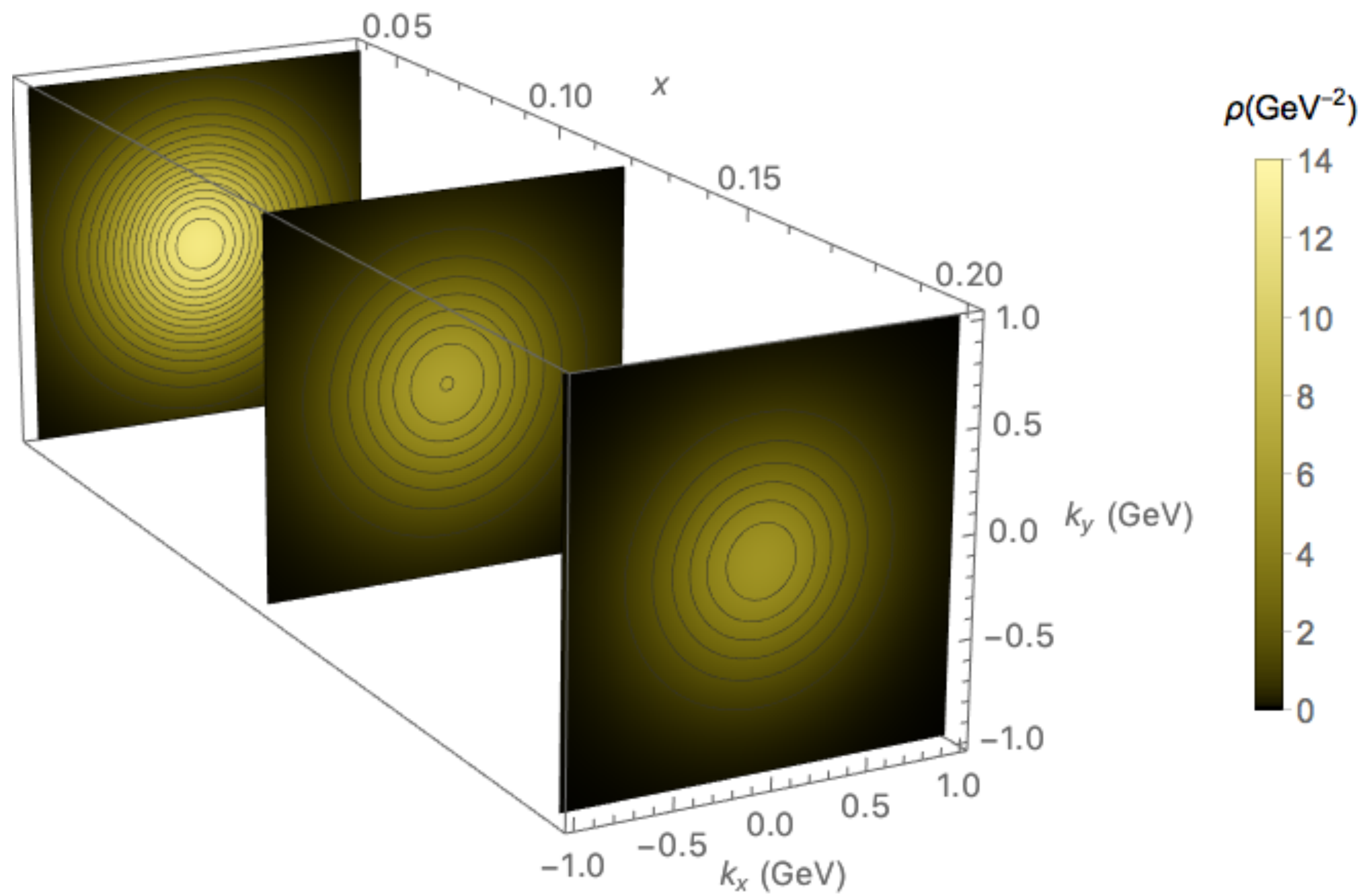


Transverse-Momentum Distributions

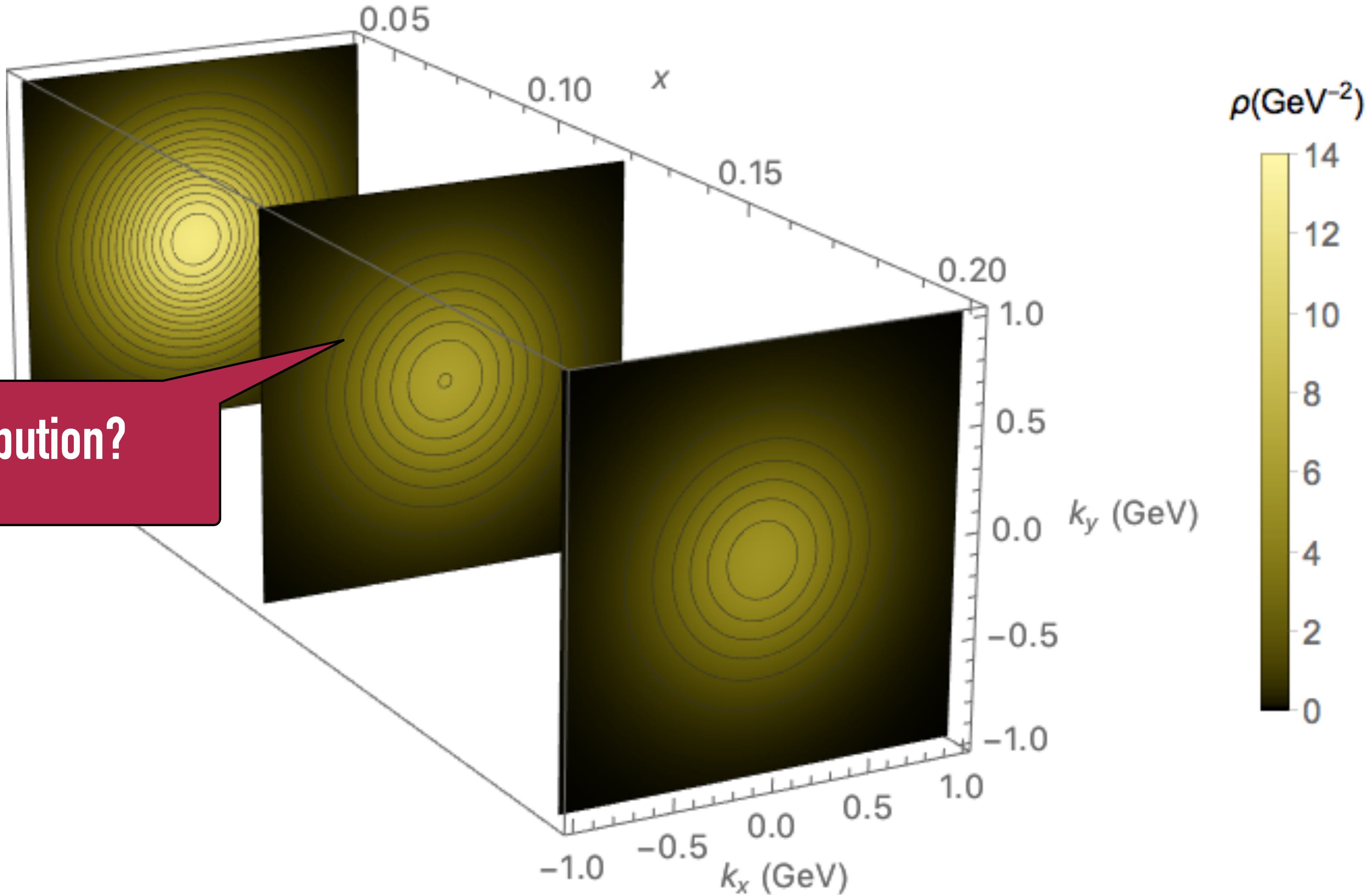
$$f(x, \vec{k}_T)$$

3 dimensional (+ 2 scales)



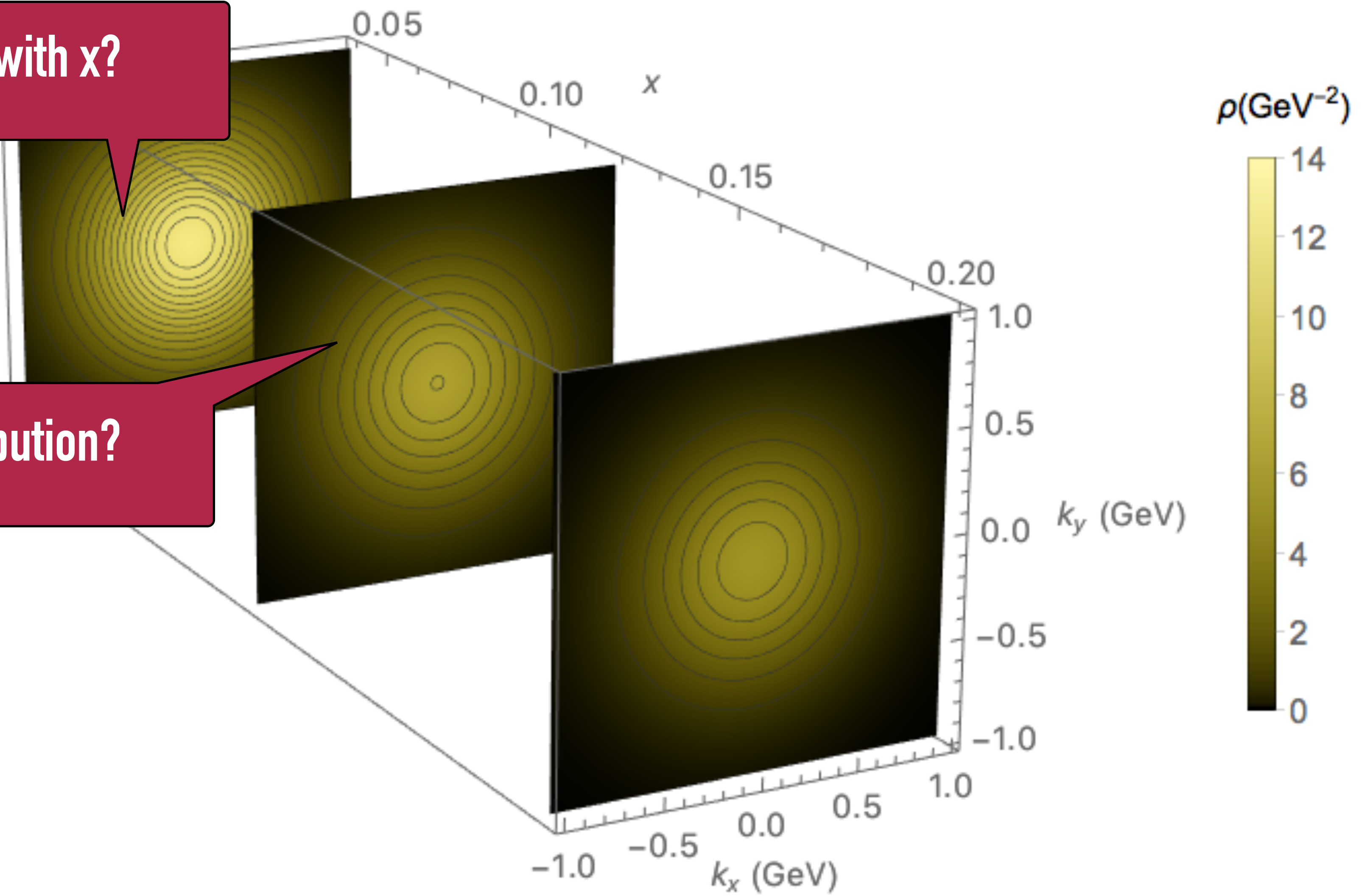


How "wide" is the distribution?



How does it change with x ?

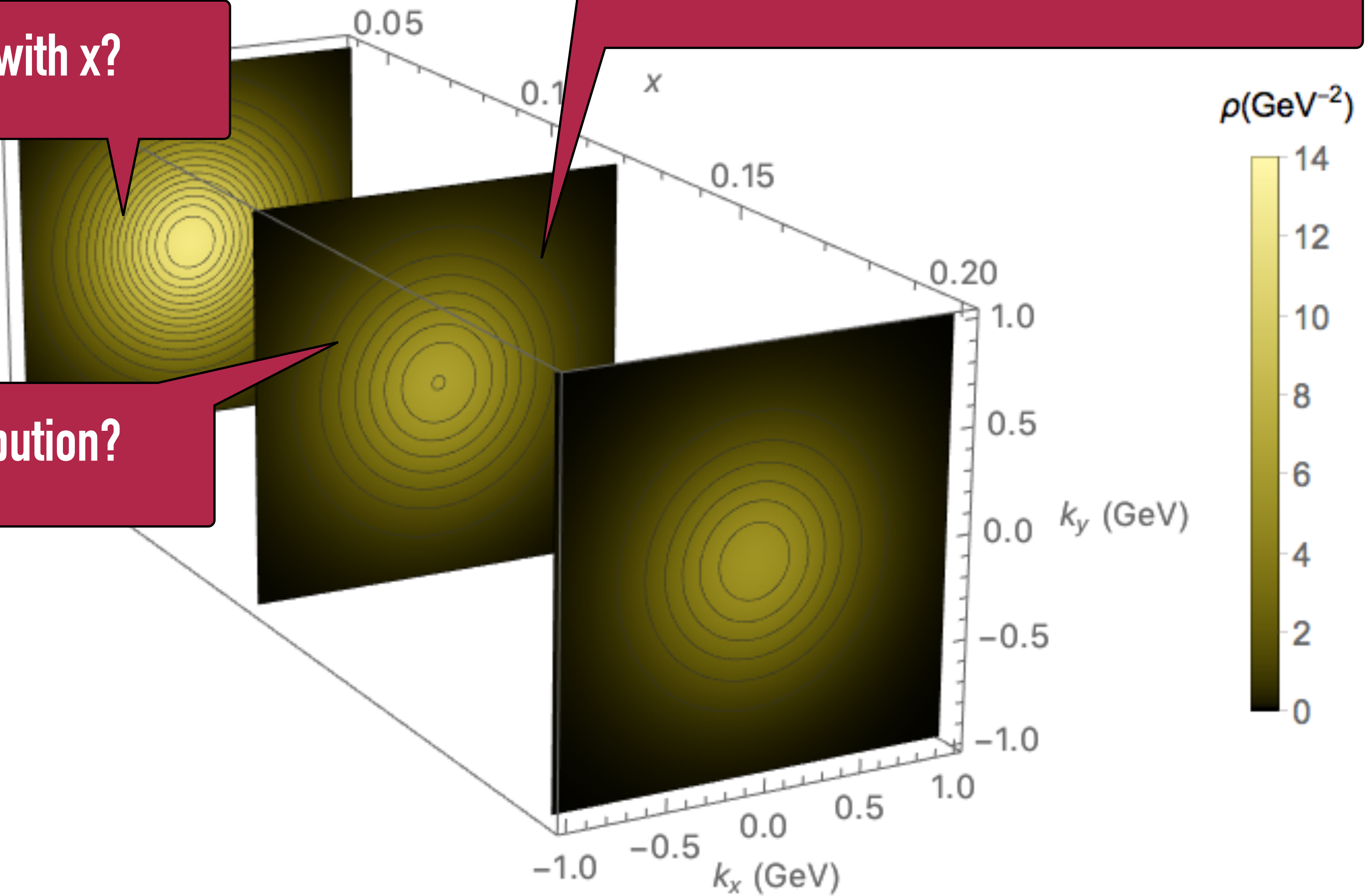
How "wide" is the distribution?



How does it change with x ?

Is there a difference between flavors?

How "wide" is the distribution?

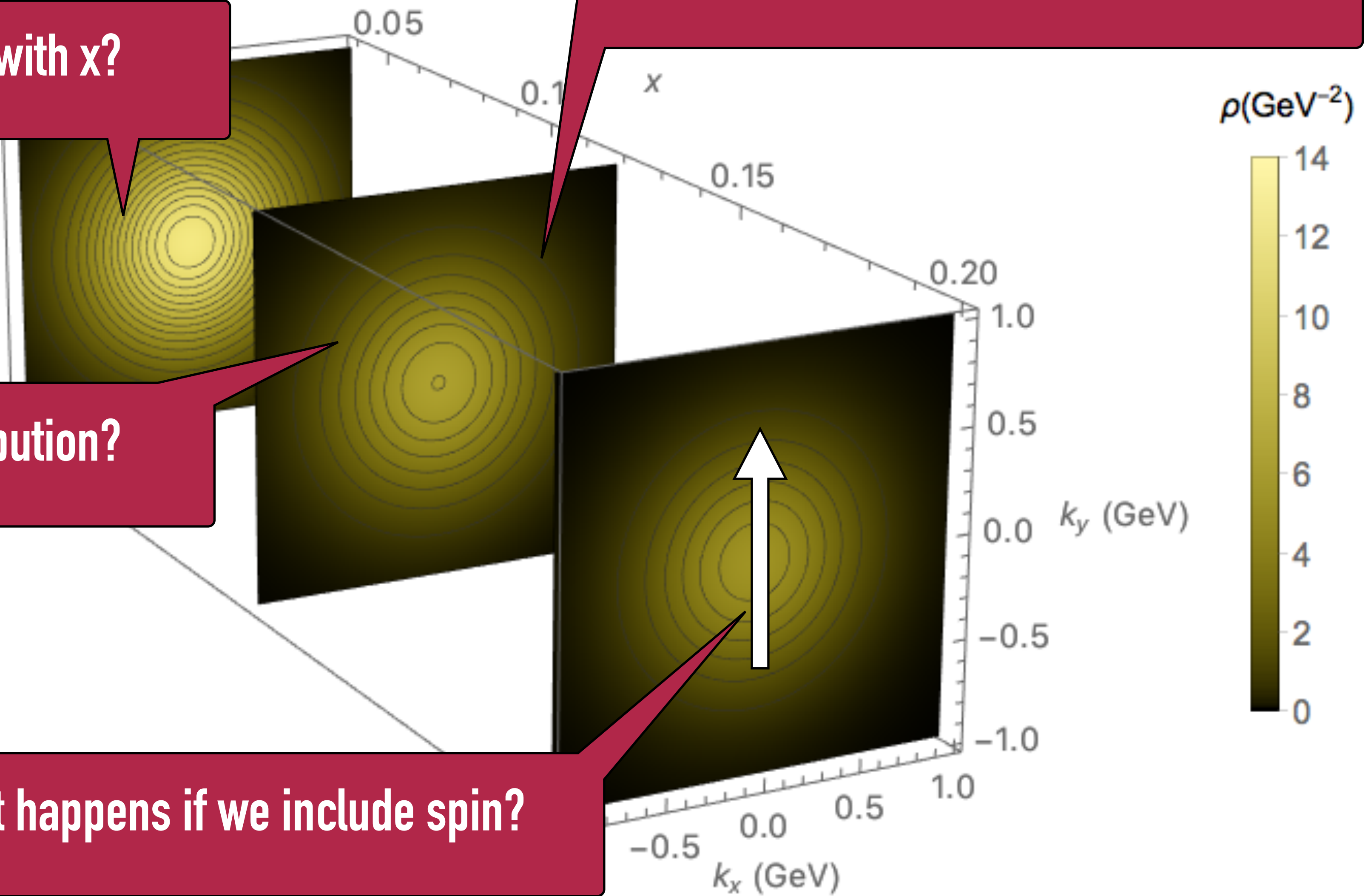


How does it change with x ?

Is there a difference between flavors?

How "wide" is the distribution?

What happens if we include spin?



[TMD collaboration, "TMD Handbook," arXiv:2304.03302](https://arxiv.org/abs/2304.03302)



Preprints: JLAB-THY-23-3780, LA-UR-21-20798, MIT-CTP/5386

TMD Handbook

Renaud Boussarie¹, Matthias Burkardt², Martha Constantinou³, William Detmold⁴, Markus Ebert^{4,5},
Michael Engelhardt², Sean Fleming⁶, Leonard Gamberg⁷, Xiangdong Ji⁸, Zhong-Bo Kang⁹,
Christopher Lee¹⁰, Keh-Fei Liu¹¹, Simonetta Liuti¹², Thomas Mehen¹³, Andreas Metz³, John Negele⁴,
Daniel Pitonyak¹⁴, Alexei Prokudin^{7,16}, Jian-Wei Qiu^{16,17}, Abha Rajan^{12,18}, Marc Schlegel^{2,19},
Phiala Shanahan⁴, Peter Schweitzer²⁰, Iain W. Stewart⁴, Andrey Tarasov^{21,22}, Raju Venugopalan¹⁸,
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471 PAGES

[Mulders-Tangerman, NPB 461 \(96\)](#)

[Boer-Mulders, PRD 57 \(98\)](#)

quark pol.

	U	L	T
nucleon pol.	U		f_1
	L	g_{1L}	h_{1L}^\perp
	T	f_{1T}^\perp	g_{1T}

TMDs in **black** survive integration over transverse momentum

TMDs in **red** are time-reversal odd

[Mulders-Tangerman, NPB 461 \(96\)](#)

[Boer-Mulders, PRD 57 \(98\)](#)

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nucleon pol.

Sivers

TMDs in **black** survive integration over transverse momentum

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TMD TABLES: QUARK, LEADING TWIST

[Mulders-Tangerman, NPB 461 \(96\)](#)

[Boer-Mulders, PRD 57 \(98\)](#)

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nucleon pol.

Sivers

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Transversity

- ▶ Very good knowledge of x dependence of f_1 and g_{1L}

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- ▶ Good knowledge of the k_T dependence of f_1 (also for pions)

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- ▶ Very good knowledge of x dependence of f_1 and g_{1L}
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- ▶ Fair knowledge of Sivers and transversity (mainly x dependence)

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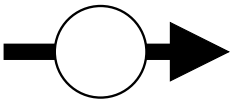

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- ▶ Some hints about all others

INTERPRETATION OF POLARIZED TMDS

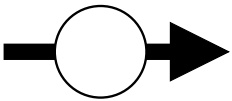

  nucleon with transverse or longitudinal spin

  parton with transverse or longitudinal spin

 parton transverse momentum

INTERPRETATION OF POLARIZED TMDS

Proton goes out of the screen/
photon goes into the screen

  nucleon with transverse or longitudinal spin

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Proton goes out of the screen/
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$$f_1 = \text{circle with red parton}$$

$$g_1 = \text{circle with black dot and red dot} - \text{circle with black dot and red cross}$$

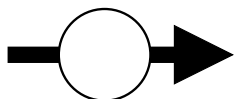
Helicity

$$h_1 = \text{circle with red parton and black arrow} - \text{circle with red parton and black arrow pointing left}$$

Transversity

INTERPRETATION OF POLARIZED TMDs

Proton goes out of the screen/
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  nucleon with transverse or longitudinal spin

  parton with transverse or longitudinal spin

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Helicity

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Transversity

$$f_{1T}^\perp = \text{circle with blue arrow pointing down} - \text{circle with blue arrow pointing up}$$

Sivers

$$h_{1T}^\perp = \text{circle with blue arrow pointing down and red arrow pointing right} - \text{circle with blue arrow pointing up and red arrow pointing right}$$

Boer-Mulders

$$g_{1T} = \text{circle with red dot and blue arrow pointing right} - \text{circle with red dot and blue arrow pointing left}$$

worm-gear

$$h_{1L}^\perp = \text{circle with black dot and red arrow pointing right} - \text{circle with black dot and red arrow pointing left}$$

worm-gear

$$h_{1T}^\perp = \text{circle with red dot and blue arrow pointing right} - \text{circle with red dot and blue arrow pointing left}$$

pretzelosity

AVAILABLE EXTRACTIONS (NEWEST ONLY)

Unpol. TMD	MAP 22 arXiv:2206.07598 , ART23 2305.07473
Helicity	
Transversity	arXiv:1505.05589 , arXiv:1612.06413 , arXiv:2205.00999
Sivers	MAP20 arXiv:2004.14278 , arXiv:2009.10710 , arXiv:2103.03270 , arXiv:2205.00999 , arXiv:2304.14328
Boer-Mulders	arXiv:0912.2031 , arXiv:1502.04214 , arXiv:2004.02117
Worm-gear g1T	arXiv:2110.10253
Worm-gear h1L	
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Worm-gear h1L	
Pretzelosity	arXiv:1411.0580

Not mentioned: pion TMDs, TMD fragmentation functions, nuclear TMDs

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[Mulders-Tangerman, NPB 461 \(96\)](#)

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[Bacchetta, Mulders, Pijlman, hep-ph/0405154](#)

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- ▶ Some knowledge of g_T x-dependence

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- ▶ Lots of progress from the theory side
- ▶ Some knowledge of g_T x-dependence
- ▶ First hints about e x-dependence

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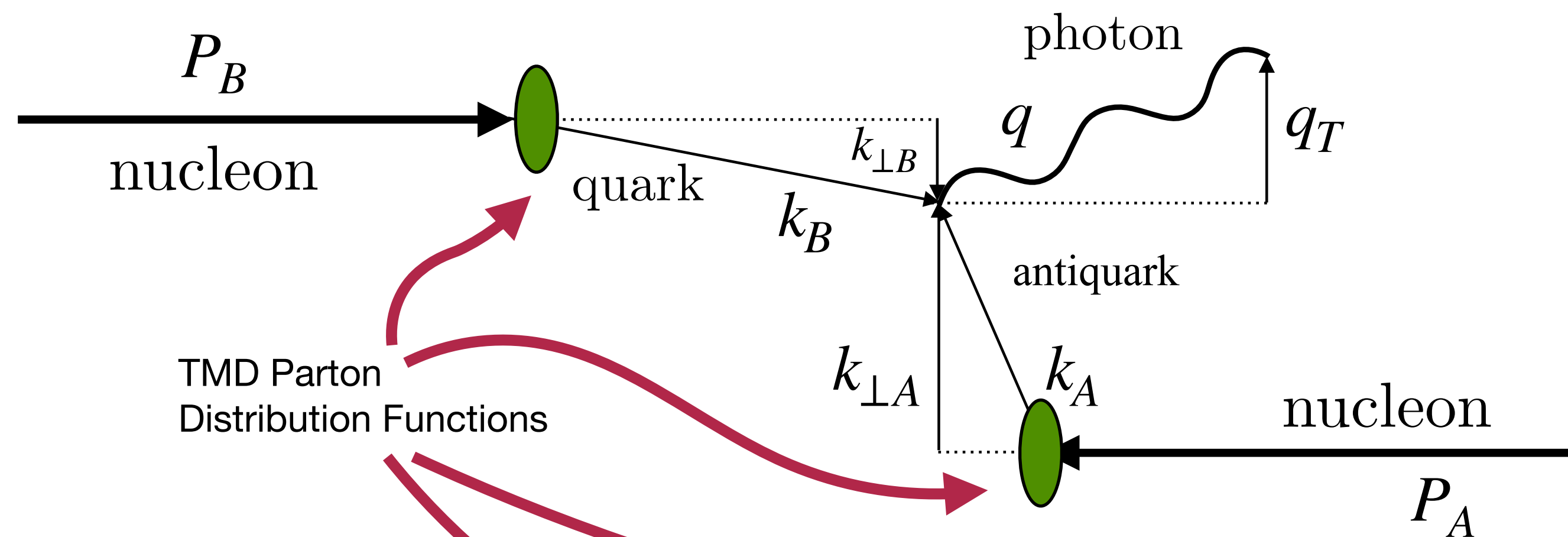
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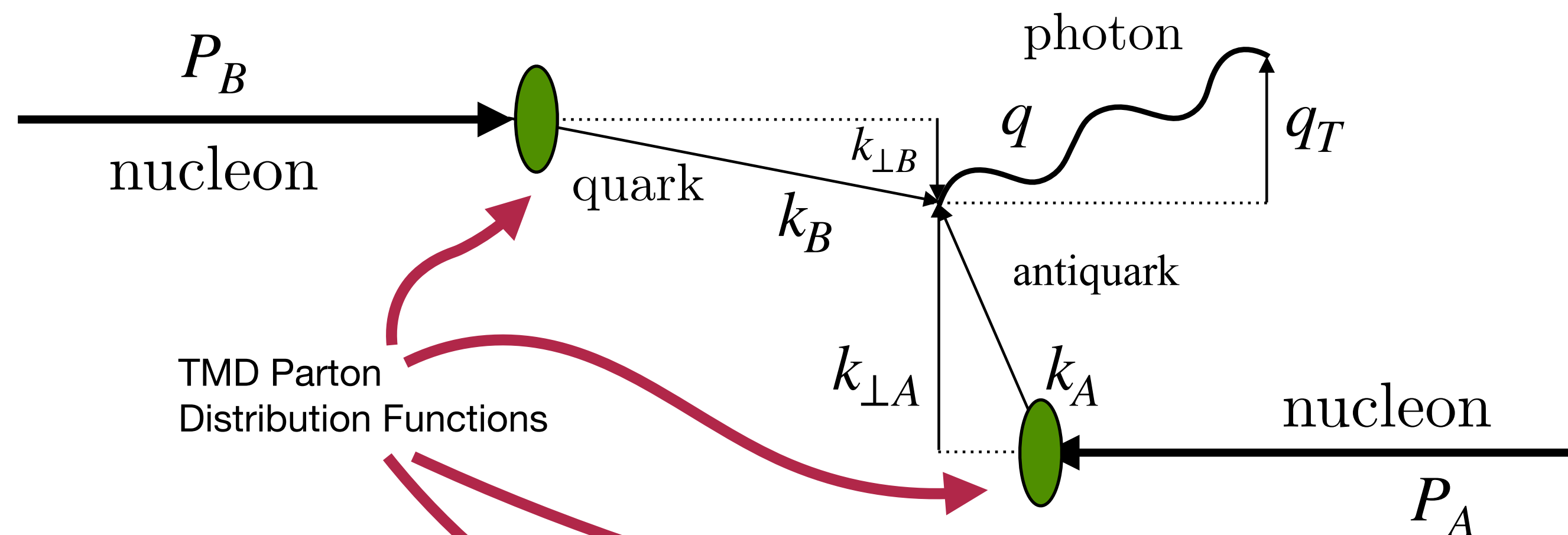
- ▶ Lots of progress from the theory side
- ▶ Some knowledge of g_T x-dependence
- ▶ First hints about e x-dependence
- ▶ All others unknown

UNPOLARIZED TMD



$$F_{UU}^1(x_A, x_B, \mathbf{q}_T^2, Q^2)$$

$$\approx \sum_q \mathcal{H}_{UU}^{1q}(Q^2, \mu^2) \int d^2\mathbf{k}_{\perp A} d^2\mathbf{k}_{\perp B} f_1^q(x_A, \mathbf{k}_{\perp A}^2; \mu^2) f_1^{\bar{q}}(x_B, \mathbf{k}_{\perp B}^2; \mu^2) \delta^{(2)}(\mathbf{k}_{\perp A} - \mathbf{q}_T + \mathbf{k}_{\perp B})$$

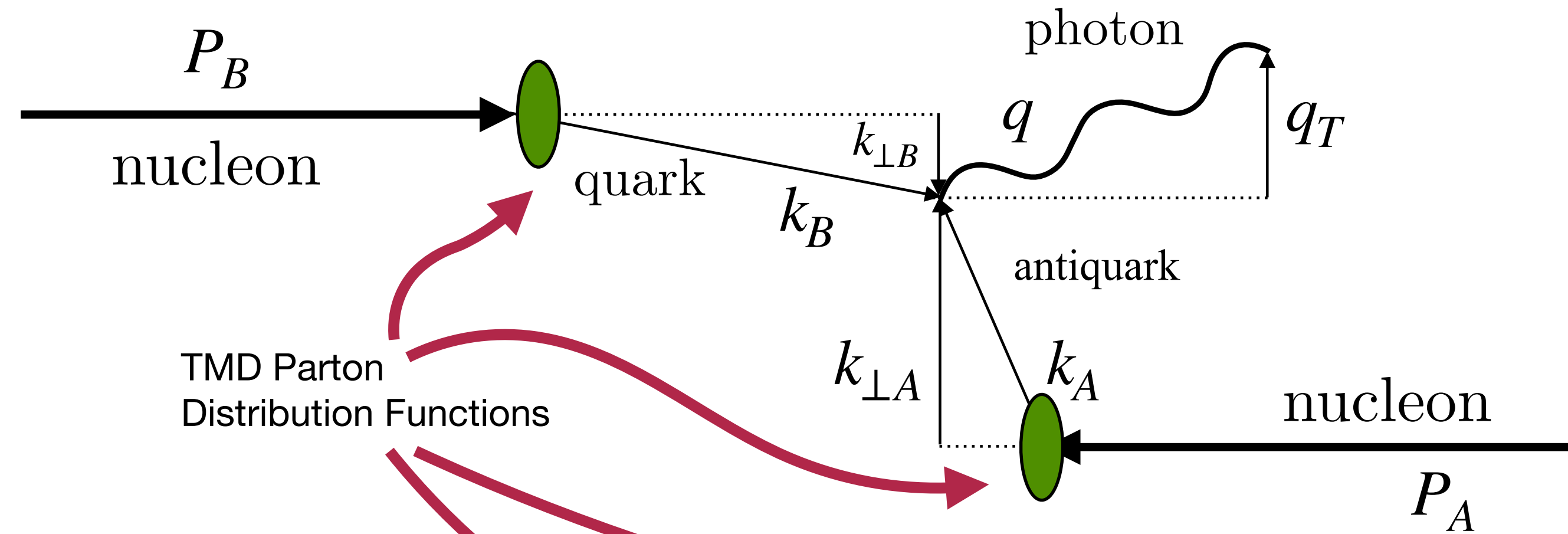


$$F_{UU}^1(x_A, x_B, \mathbf{q}_T^2, Q^2)$$

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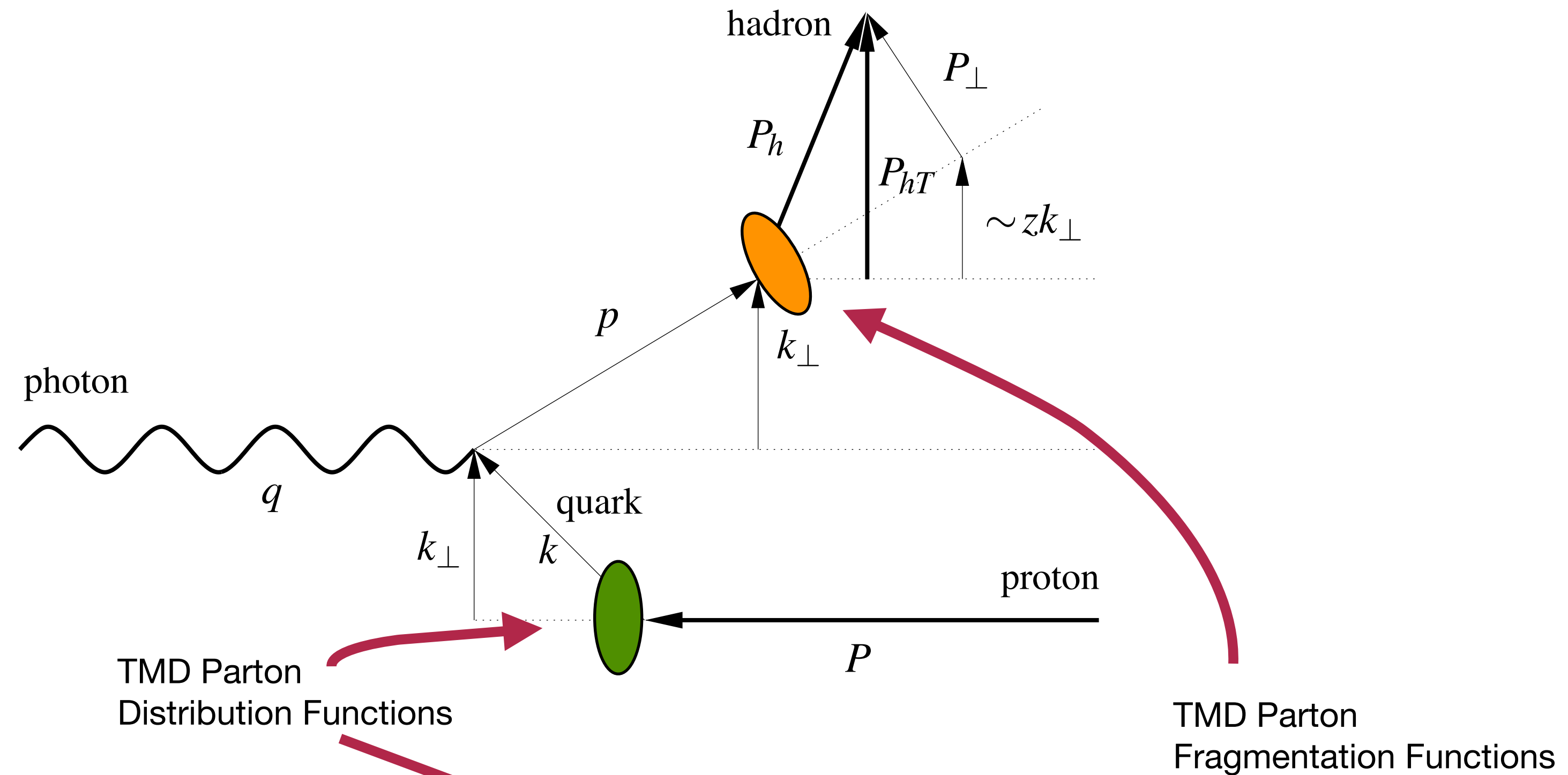
$$= \sum_q \mathcal{H}_{UU}^{1q}(Q^2, \mu^2) \int db_T b_T J_0(b_T |\mathbf{q}_T|) \hat{f}_1^q(x_A, b_T^2; \mu^2) \hat{f}_1^{\bar{q}}(x_B, b_T^2; \mu^2)$$

The analysis is usually done in Fourier-transformed space



$$\begin{aligned}
 F_{UU}^1(x_A, x_B, \mathbf{q}_T^2, Q^2) & \\
 & \approx \sum_q \mathcal{H}_{UU}^{1q}(Q^2, \mu^2) \int d^2\mathbf{k}_{\perp A} d^2\mathbf{k}_{\perp B} f_1^q(x_A, \mathbf{k}_{\perp A}^2; \mu^2) f_1^{\bar{q}}(x_B, \mathbf{k}_{\perp B}^2; \mu^2) \delta^{(2)}(\mathbf{k}_{\perp A} - \mathbf{q}_T + \mathbf{k}_{\perp B}) \\
 & = \sum_q \mathcal{H}_{UU}^{1q}(Q^2, \mu^2) \int db_T b_T J_0(b_T |\mathbf{q}_T|) \hat{f}_1^q(x_A, b_T^2; \mu^2) \hat{f}_1^{\bar{q}}(x_B, b_T^2; \mu^2)
 \end{aligned}$$

The analysis is usually done in Fourier-transformed space
 TMDs formally depend on two scales, but we set them equal.



$$\begin{aligned}
 F_{UU,T}(x, z, \mathbf{P}_{hT}^2, Q^2) &= x \sum_q \mathcal{H}_{UU,T}^q(Q^2, \mu^2) \int d^2 \mathbf{k}_\perp d^2 \mathbf{P}_\perp f_1^a(x, \mathbf{k}_\perp^2; \mu^2) D_1^{a \rightarrow h}(z, \mathbf{P}_\perp^2; \mu^2) \delta(z \mathbf{k}_\perp - \mathbf{P}_{hT} + \mathbf{P}_\perp) \\
 &= x \sum_a \mathcal{H}_{UU,T}^a(Q^2, \mu^2) \int db_T b_T J_0(b_T |\mathbf{P}_{h\perp}|) \hat{f}_1^q(x, z^2 b_\perp^2; \mu^2) \hat{D}_1^{a \rightarrow h}(z, b_\perp^2; \mu^2)
 \end{aligned}$$

$$\hat{f}_1^a(x, |\mathbf{b}_T|; \mu, \zeta) = \int d^2\mathbf{k}_\perp e^{i\mathbf{b}_T \cdot \mathbf{k}_\perp} f_1^a(x, \mathbf{k}_\perp^2; \mu, \zeta)$$

[see, e.g., Collins, "Foundations of Perturbative QCD" \(11\)](#)
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$$\hat{f}_1^a(x, b_T^2; \mu_f, \zeta_f) = [C \otimes f_1](x, \mu b_*) e^{\int_{\mu b_*}^{\mu_f} \frac{d\mu}{\mu} (\gamma_F - \gamma_K \ln \frac{\sqrt{\zeta_f}}{\mu})} \left(\frac{\sqrt{\zeta_f}}{\mu b_*} \right)^{K_{\text{resum}} + g_K}$$

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$$\hat{f}_1^a(x, |\mathbf{b}_T|; \mu, \zeta) = \int d^2\mathbf{k}_\perp e^{i\mathbf{b}_T \cdot \mathbf{k}_\perp} f_1^a(x, \mathbf{k}_\perp^2; \mu, \zeta)$$

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$$\mu_{b_*} = \frac{2e^{-\gamma_E}}{b_T}$$

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perturbative Sudakov form factor

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collinear PDF

Collins-Soper kernel (perturbative and nonperturbative)

$$\mu_b = \frac{2e^{-\gamma_E}}{b_T}$$

matching coefficients (perturbative)

[see, e.g., Collins, "Foundations of Perturbative QCD" \(11\)](#)
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$$\hat{f}_1^a(x, |\mathbf{b}_T|; \mu, \zeta) = \int d^2\mathbf{k}_\perp e^{i\mathbf{b}_T \cdot \mathbf{k}_\perp} f_1^a(x, \mathbf{k}_\perp^2; \mu, \zeta)$$

perturbative Sudakov form factor

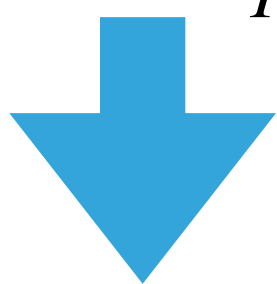
$$\hat{f}_1^a(x, b_T^2; \mu_f, \zeta_f) = [C \otimes f_1](x, \mu_{b_*}) e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu} (\gamma_F - \gamma_K \ln \frac{\sqrt{\zeta_f}}{\mu})} \left(\frac{\sqrt{\zeta_f}}{\mu_{b_*}} \right)^{K_{\text{resum}} + g_K}$$

collinear PDF

Collins-Soper kernel (perturbative and nonperturbative)

matching coefficients (perturbative)

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$$\mu_{b_*} = \frac{2e^{-\gamma_E}}{\bar{b}_*}$$

[see, e.g., Collins, "Foundations of Perturbative QCD" \(11\)](#)
[TMD collaboration, "TMD Handbook," arXiv:2304.03302](#)

$$\hat{f}_1^a(x, |\mathbf{b}_T|; \mu, \zeta) = \int d^2\mathbf{k}_\perp e^{i\mathbf{b}_T \cdot \mathbf{k}_\perp} f_1^a(x, \mathbf{k}_\perp^2; \mu, \zeta)$$

perturbative Sudakov form factor

$$\hat{f}_1^a(x, b_T^2; \mu_f, \zeta_f) = [C \otimes f_1](x, \mu_{b_*}) e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu} (\gamma_F - \gamma_K \ln \frac{\sqrt{\zeta_f}}{\mu})} \left(\frac{\sqrt{\zeta_f}}{\mu_{b_*}} \right)^{K_{\text{resum}} + g_K} f_{1NP}(x, b_T^2; \zeta_f, Q_0)$$

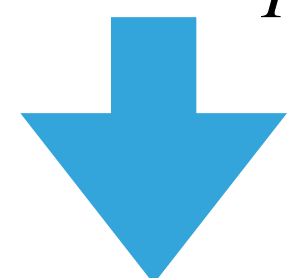
collinear PDF

Collins-Soper kernel (perturbative and nonperturbative)

nonperturbative part of TMD

matching coefficients (perturbative)

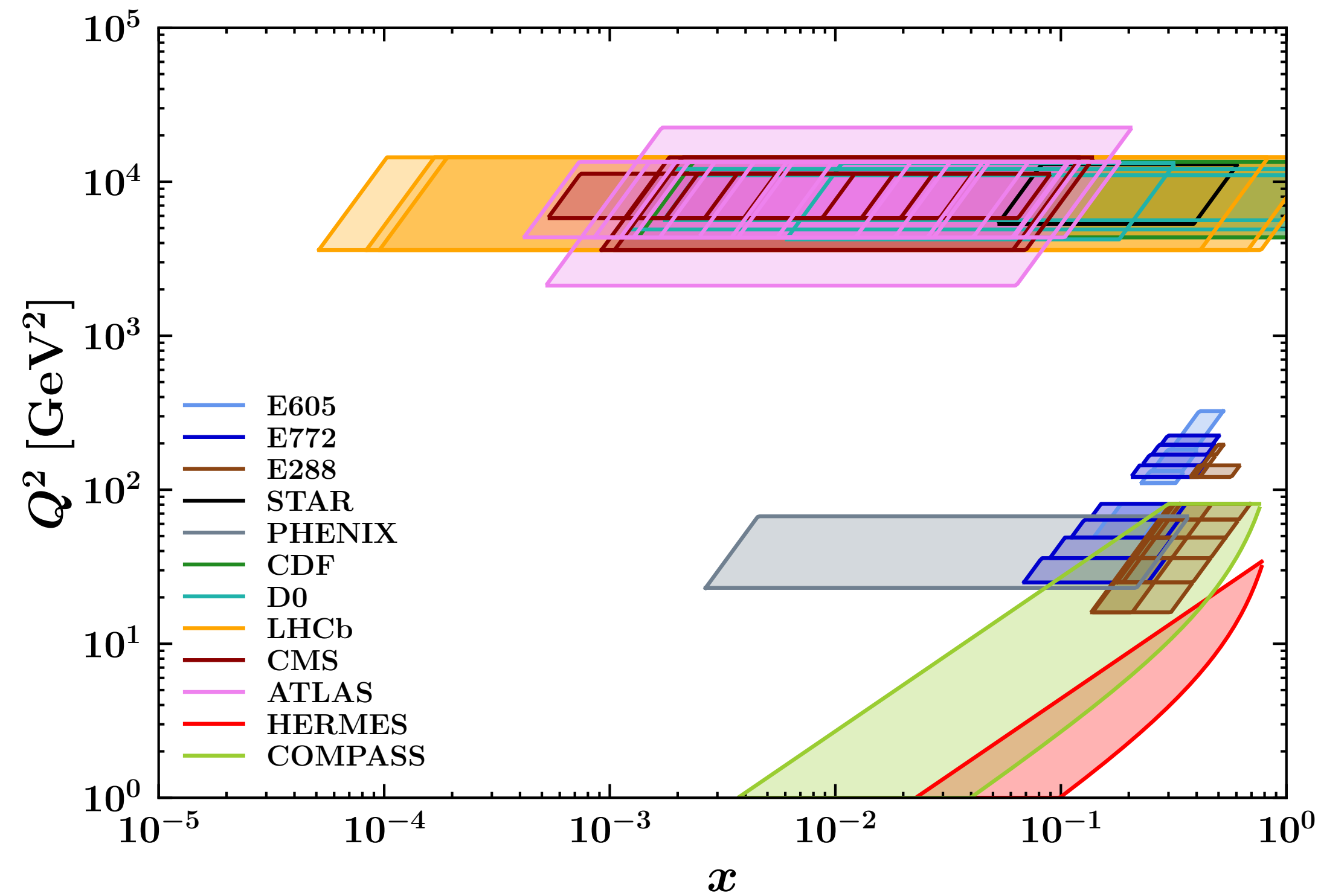
$$\mu_b = \frac{2e^{-\gamma_E}}{b_T}$$



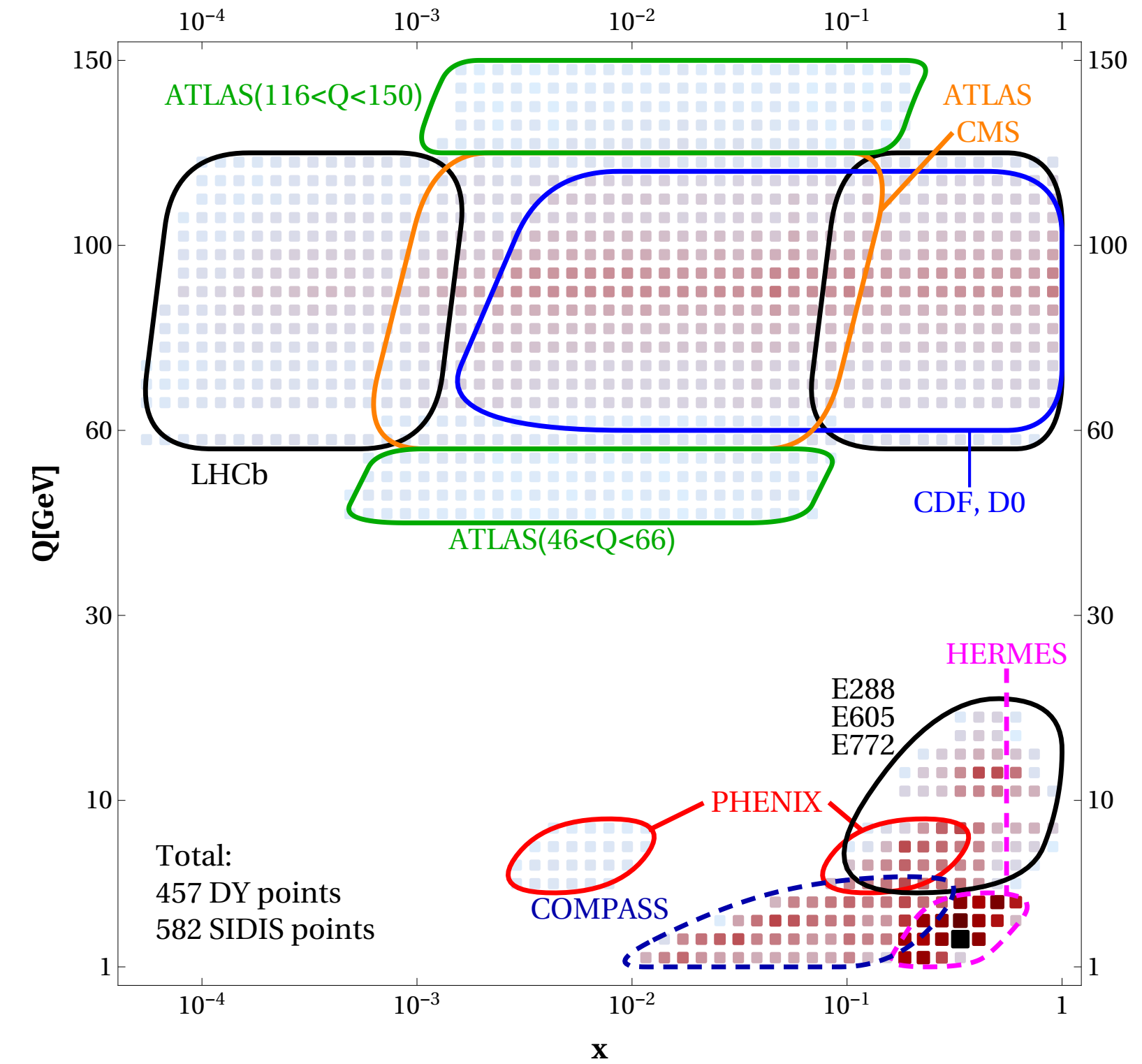
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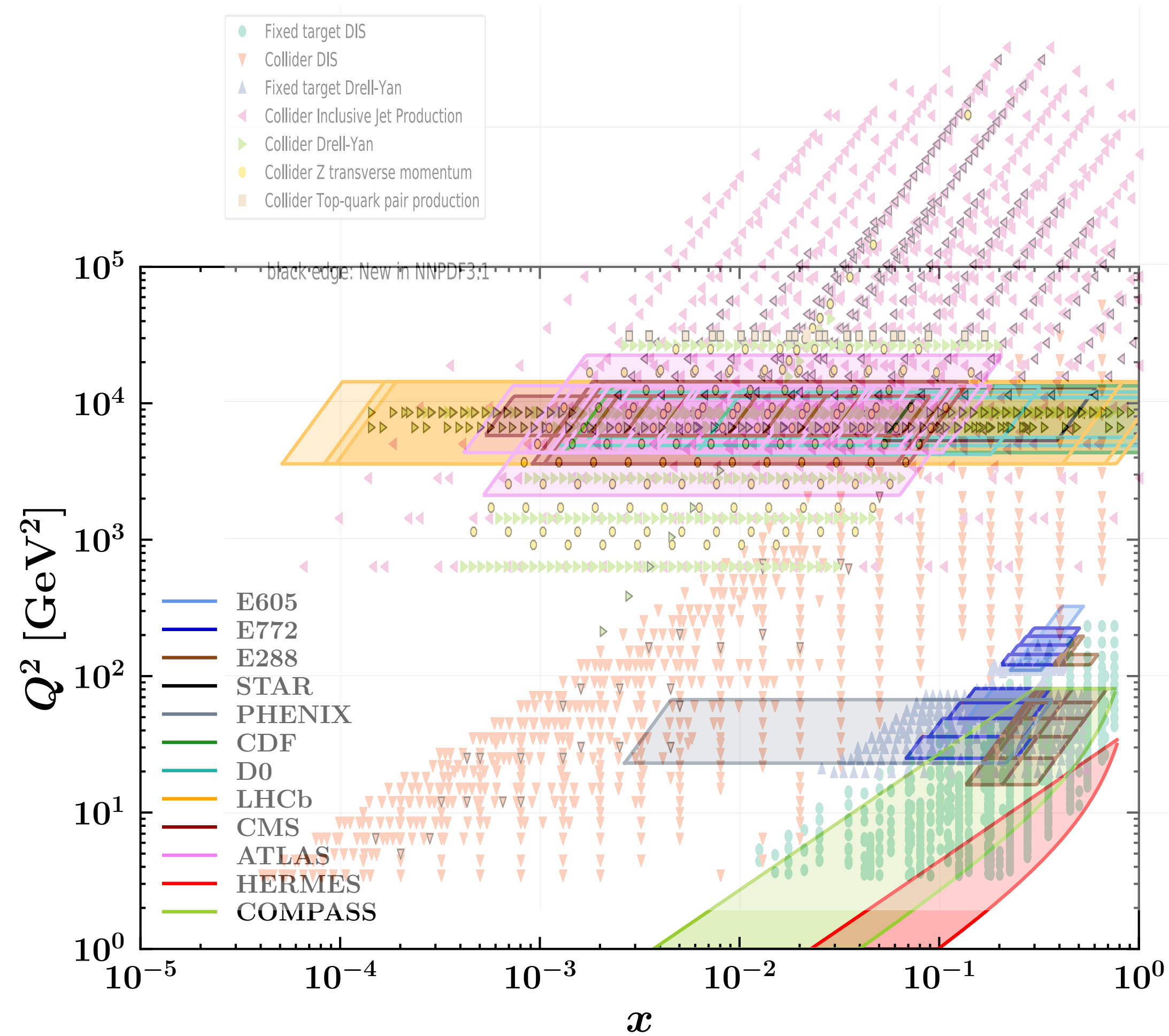
	Accuracy	SIDIS HERMES	SIDIS COMPASS	DY fixed target	DY collider	N of points	χ^2/N_{points}
Pavia 2017 arXiv:1703.10157	NLL	✓	✓	✓	✓	8059	1.55
SV 2019 arXiv:1912.06532	N ³ LL ⁻	✓	✓	✓	✓	1039	1.06
MAP22 arXiv:2206.07598	N ³ LL ⁻	✓	✓	✓	✓	2031	1.06



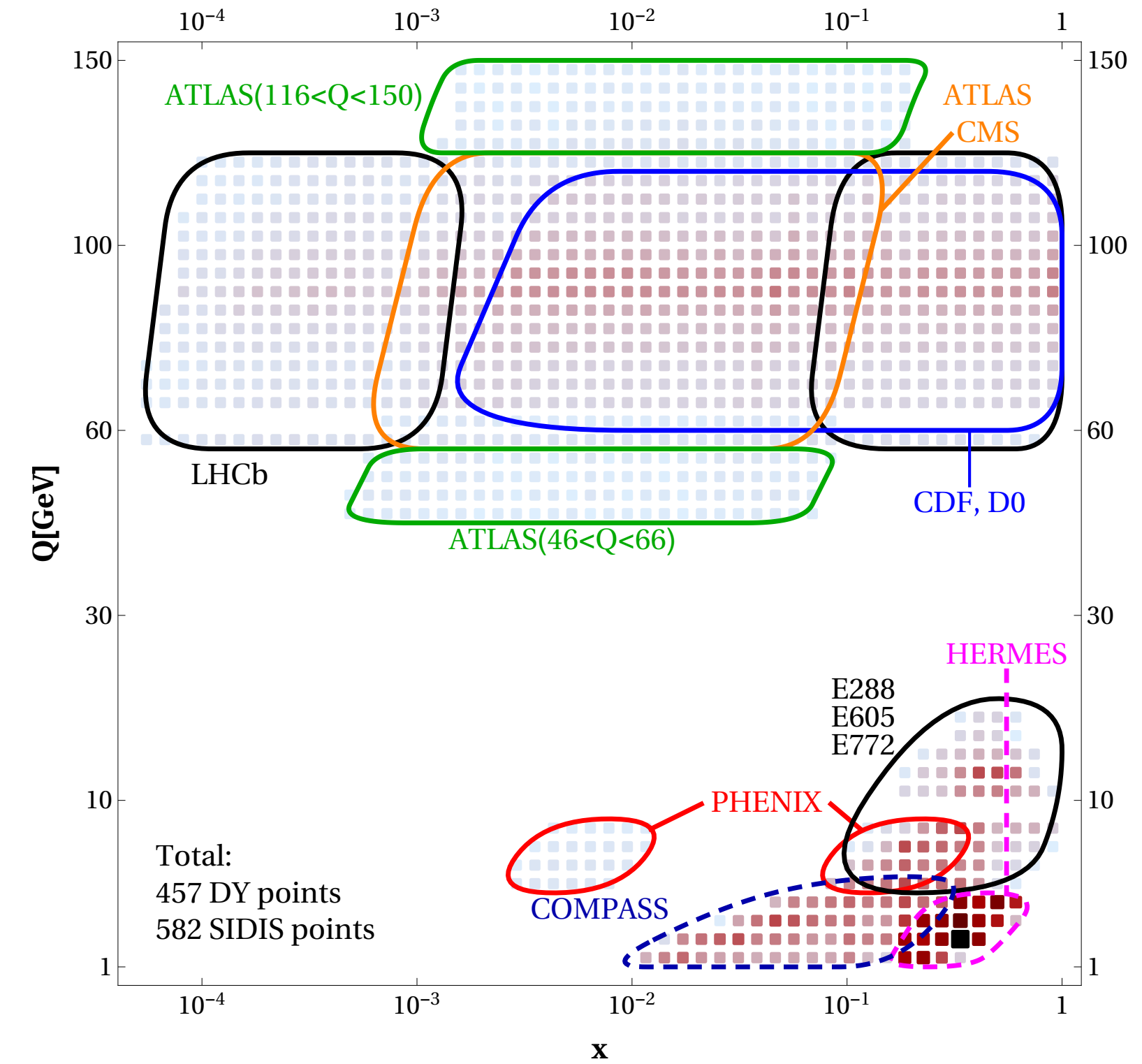
[MAP Collaboration](#)
[Bacchetta, Bertone, Bissolotti, Bozzi, Cerutti,](#)
[Piacenza, Radici, Signori, arXiv:2206.07598](#)



[Scimemi, Vladimirov,](#)
[arXiv:1912.06532](#)



[MAP Collaboration](#)
[Bacchetta, Bertone, Bissolotti, Bozzi, Cerutti, Piacenza, Radici, Signori, arXiv:2206.07598](#)



[Scimemi, Vladimirov, arXiv:1912.06532](#)

Data set	N_{dat}	χ_D^2/N_{dat}	$\chi_\lambda^2/N_{\text{dat}}$	χ_0^2/N_{dat}
Tevatron total	71	0.87	0.06	0.93
LHCb total	21	1.15	0.3	1.45
ATLAS total	72	4.56	0.48	5.05
CMS total	78	0.53	0.02	0.55
PHENIX 200	2	2.21	0.88	3.08
STAR 510	7	1.05	0.10	1.15
DY collider total	251	1.86	0.2	2.06
DY fixed-target total	233	0.85	0.4	1.24
HERMES total	344	0.48	0.23	0.71
COMPASS total	1203	0.62	0.3	0.92
SIDIS total	1547	0.59	0.28	0.87
Total	2031	0.77	0.29	1.06

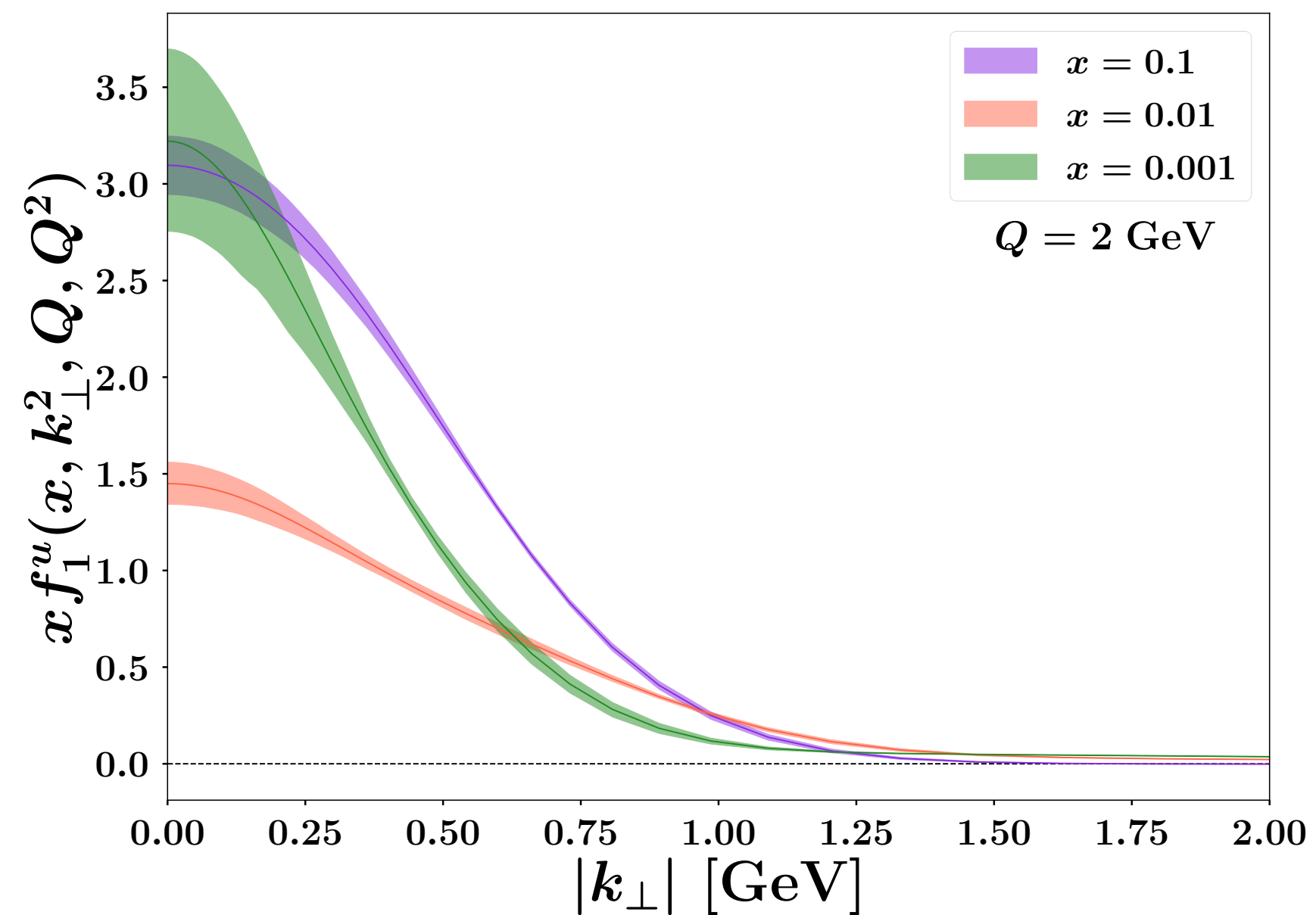


FIG. 13: The TMD PDF of the up quark in a proton at $\mu = \sqrt{\zeta} = Q = 2$ GeV (left panel) and 10 GeV (right panel) as a function of the partonic transverse momentum $|k_\perp|$ for $x = 0.001, 0.01$ and 0.1 . The uncertainty bands represent the 68% CL.

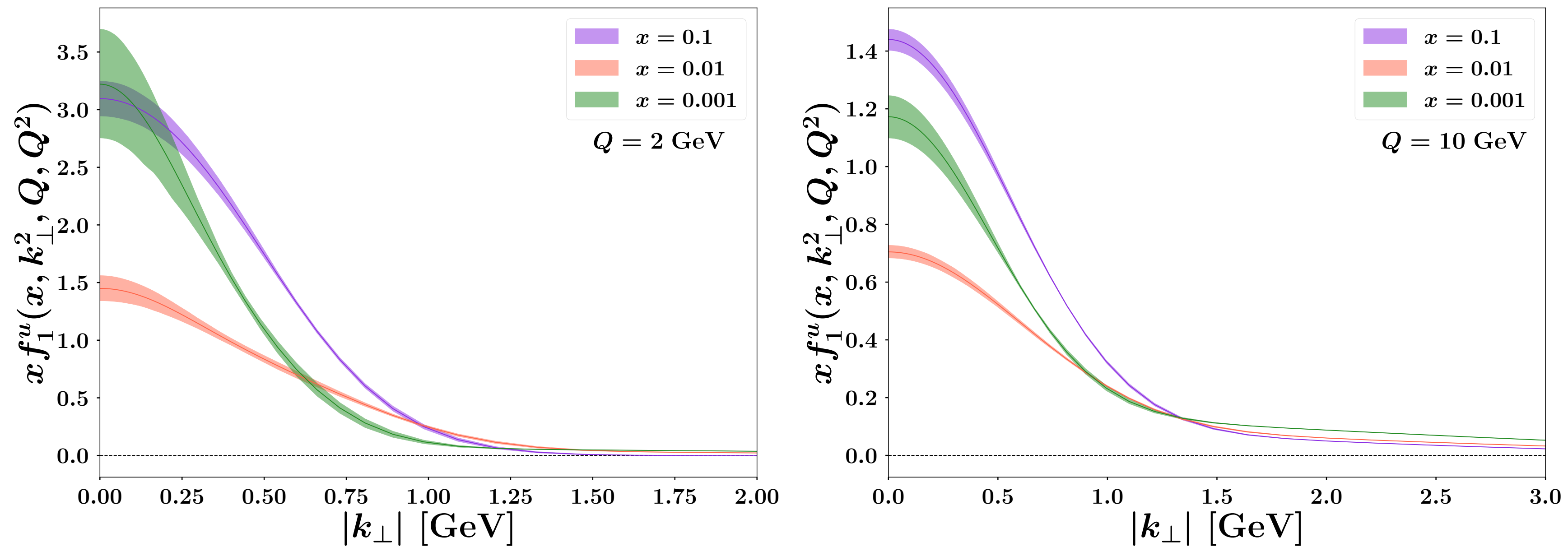
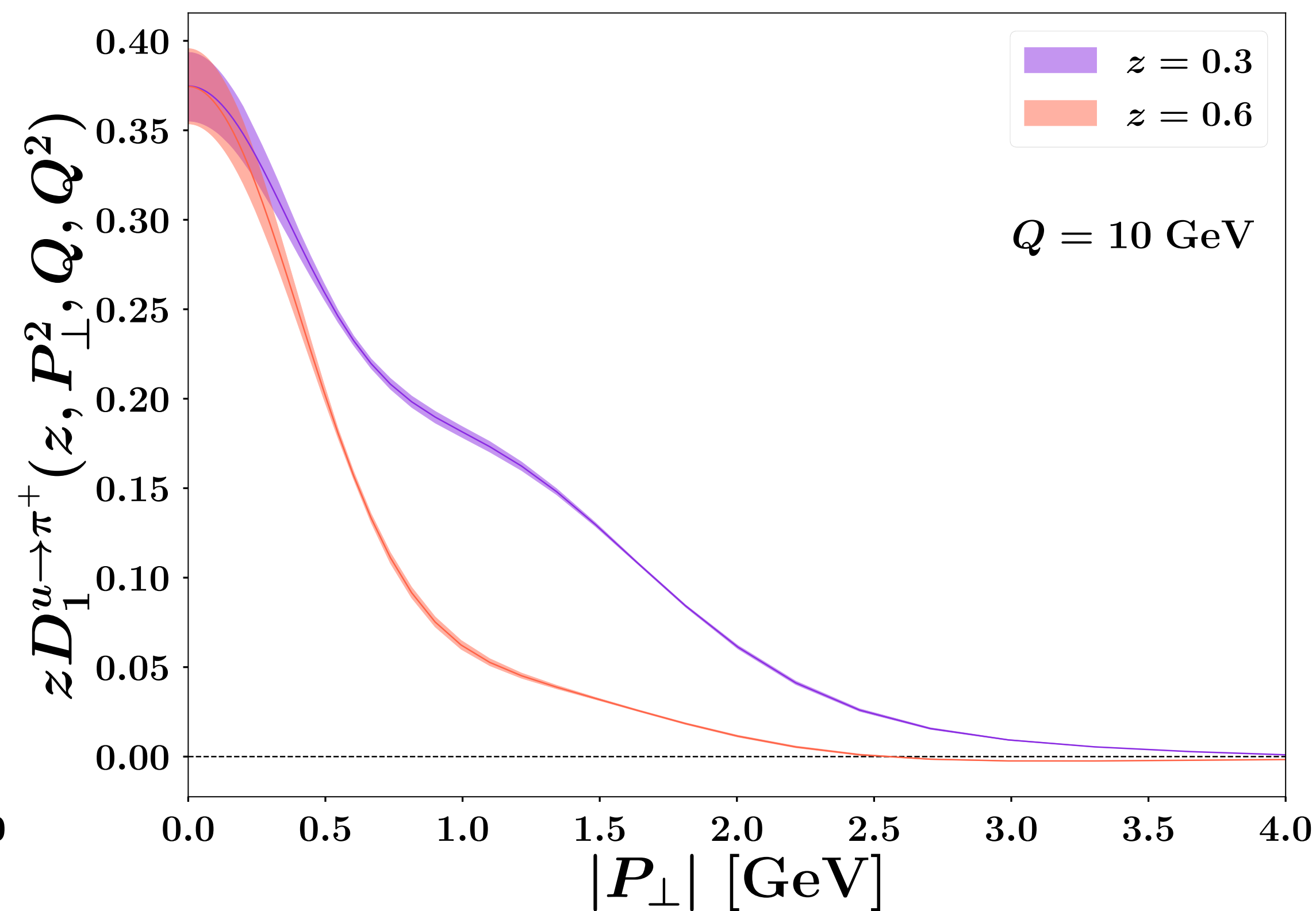
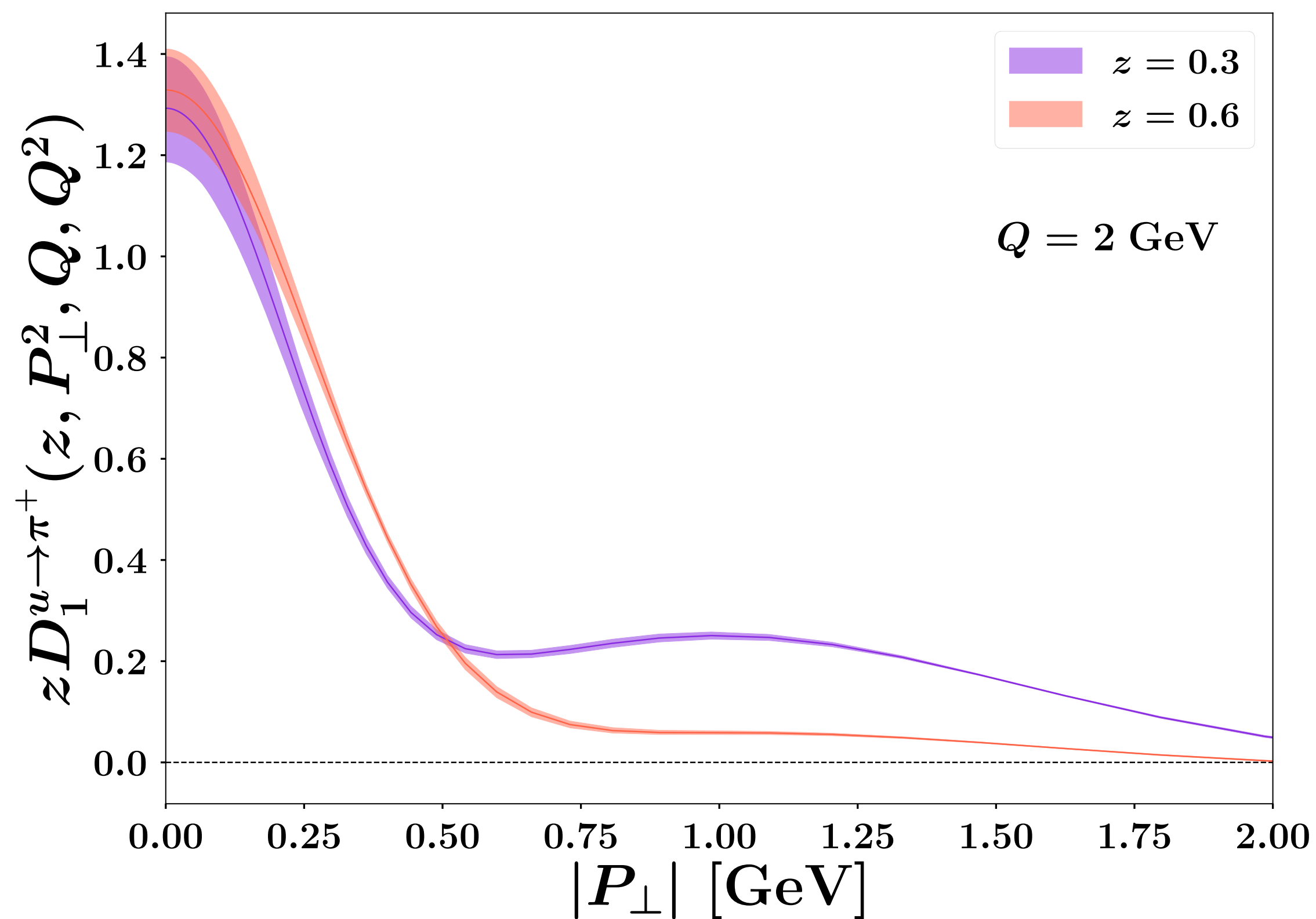
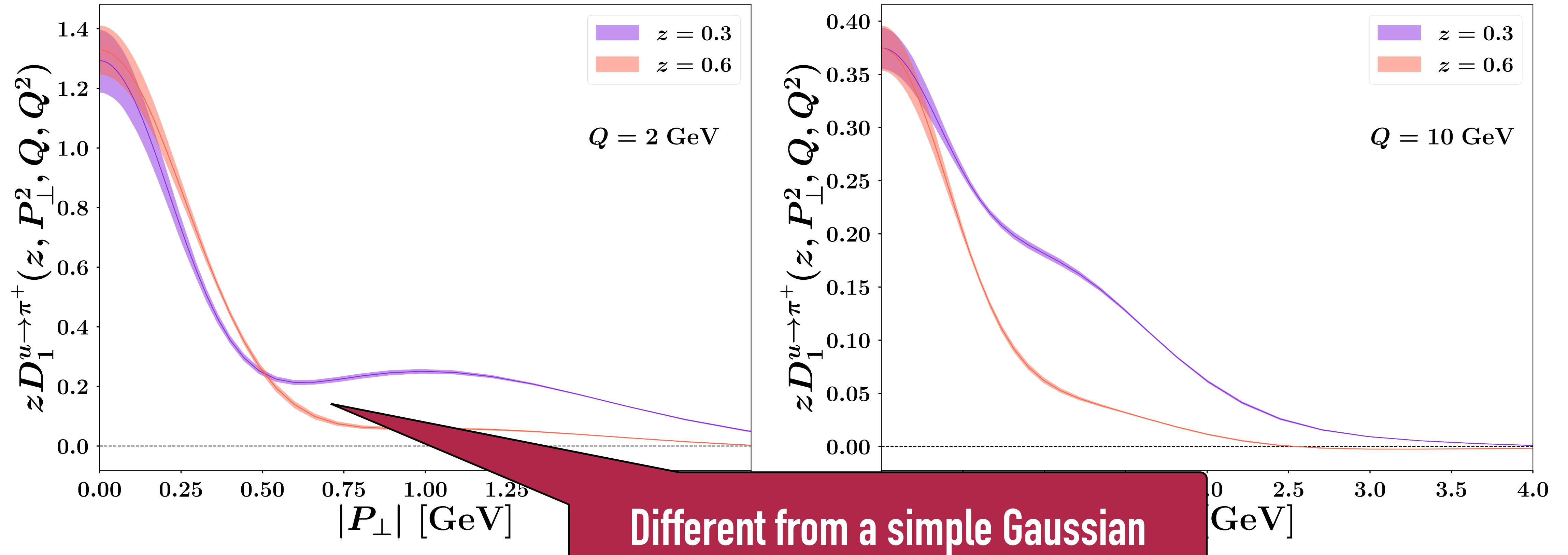


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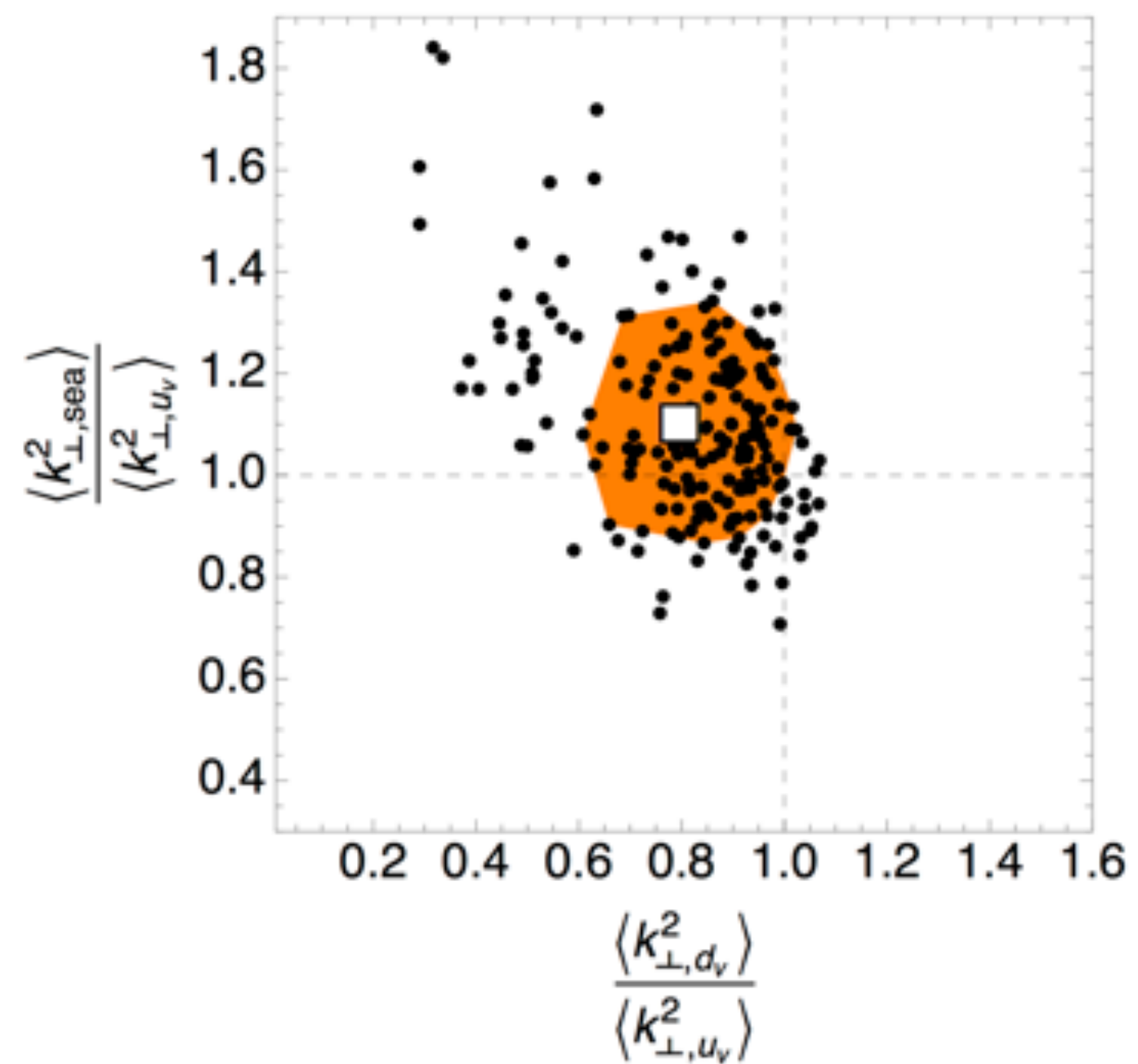
RESULTING TMD FRAGMENTATION FUNCTIONS

MAP Collaboration, arXiv:2206.07598



[Signori, Bacchetta, Radici, Schnell JHEP 1311 \(13\)](#)

Ratio of width of sea /
width of up valence

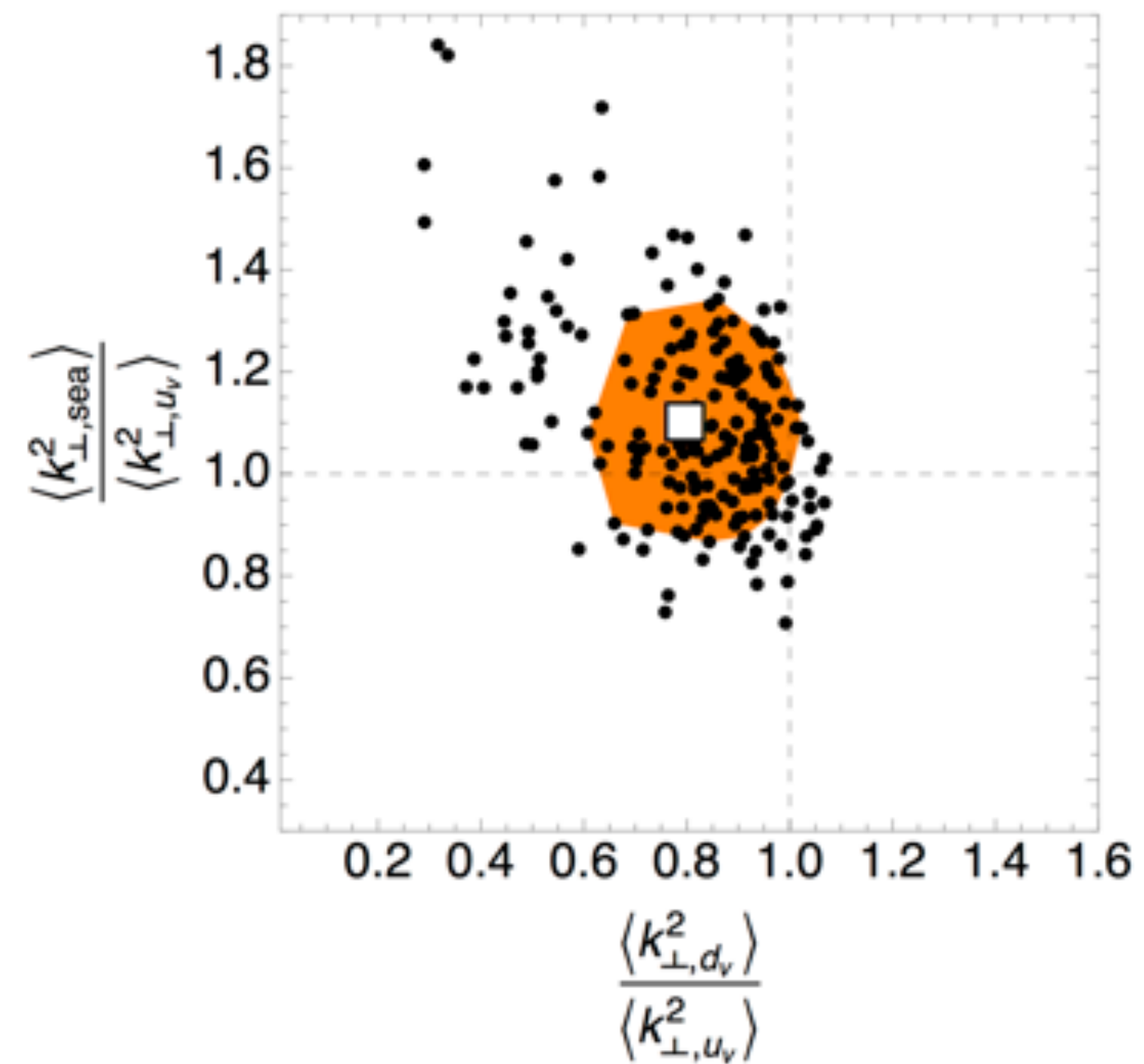


Ratio width of down valence/
width of up valence

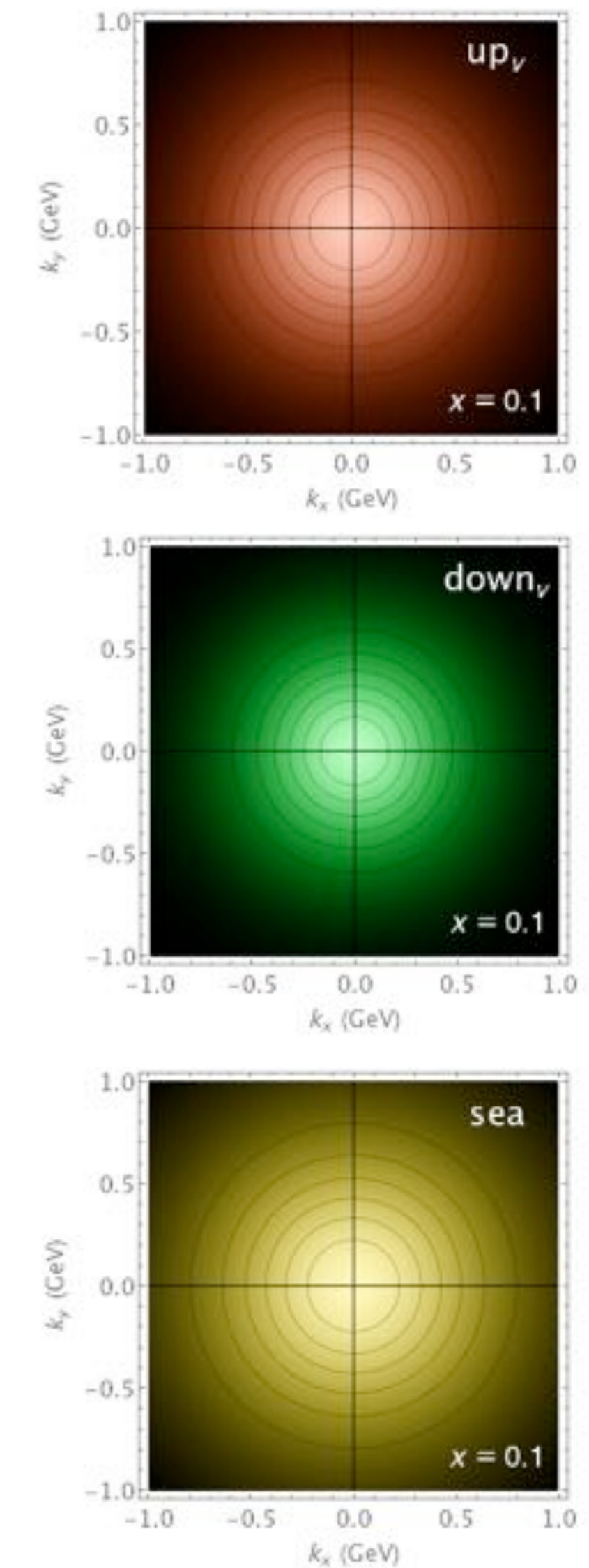
FLAVOR DEPENDENCE OF TMDS

[Signori, Bacchetta, Radici, Schnell JHEP 1311 \(13\)](#)

Ratio of width of sea /
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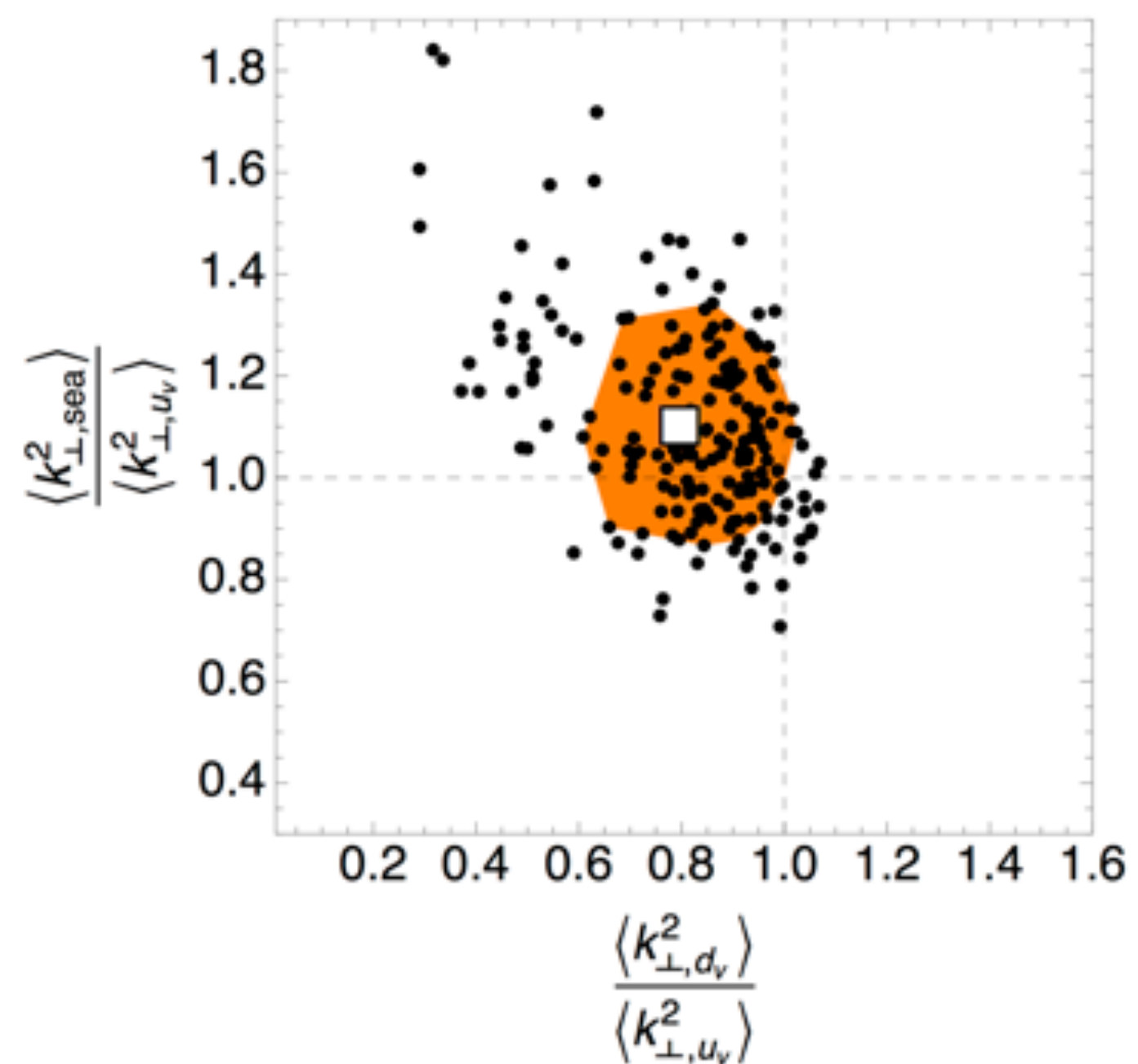
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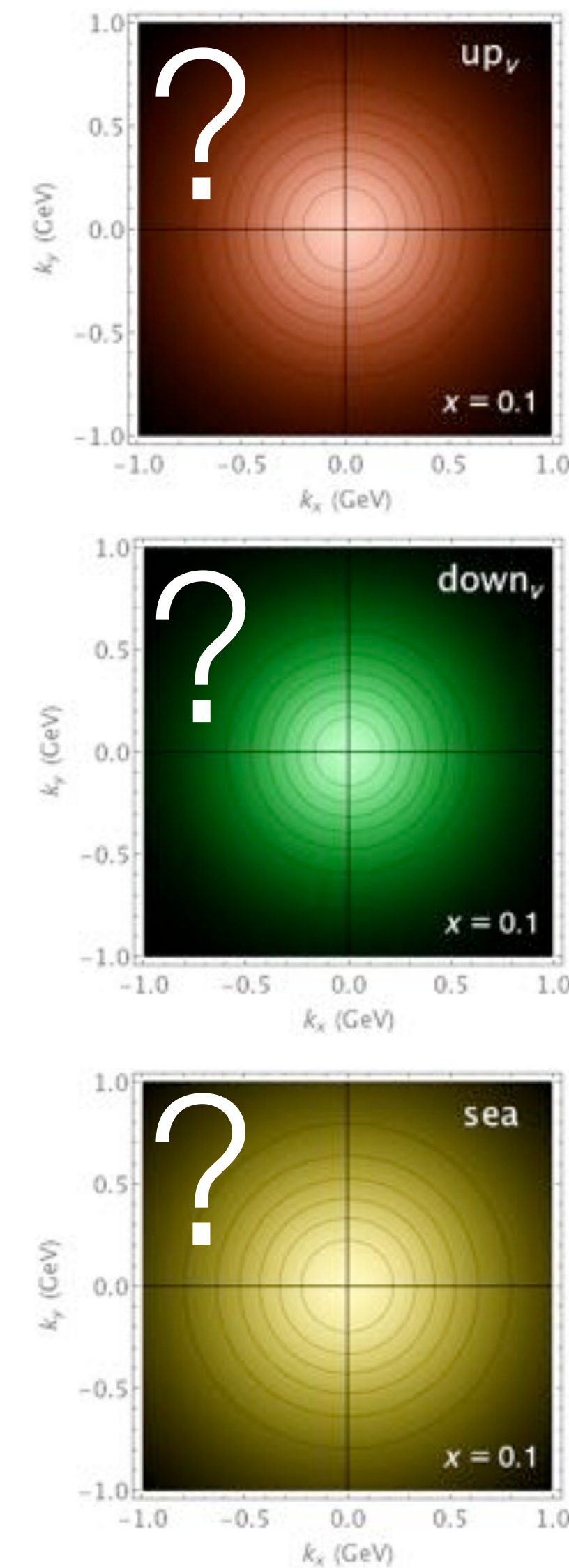
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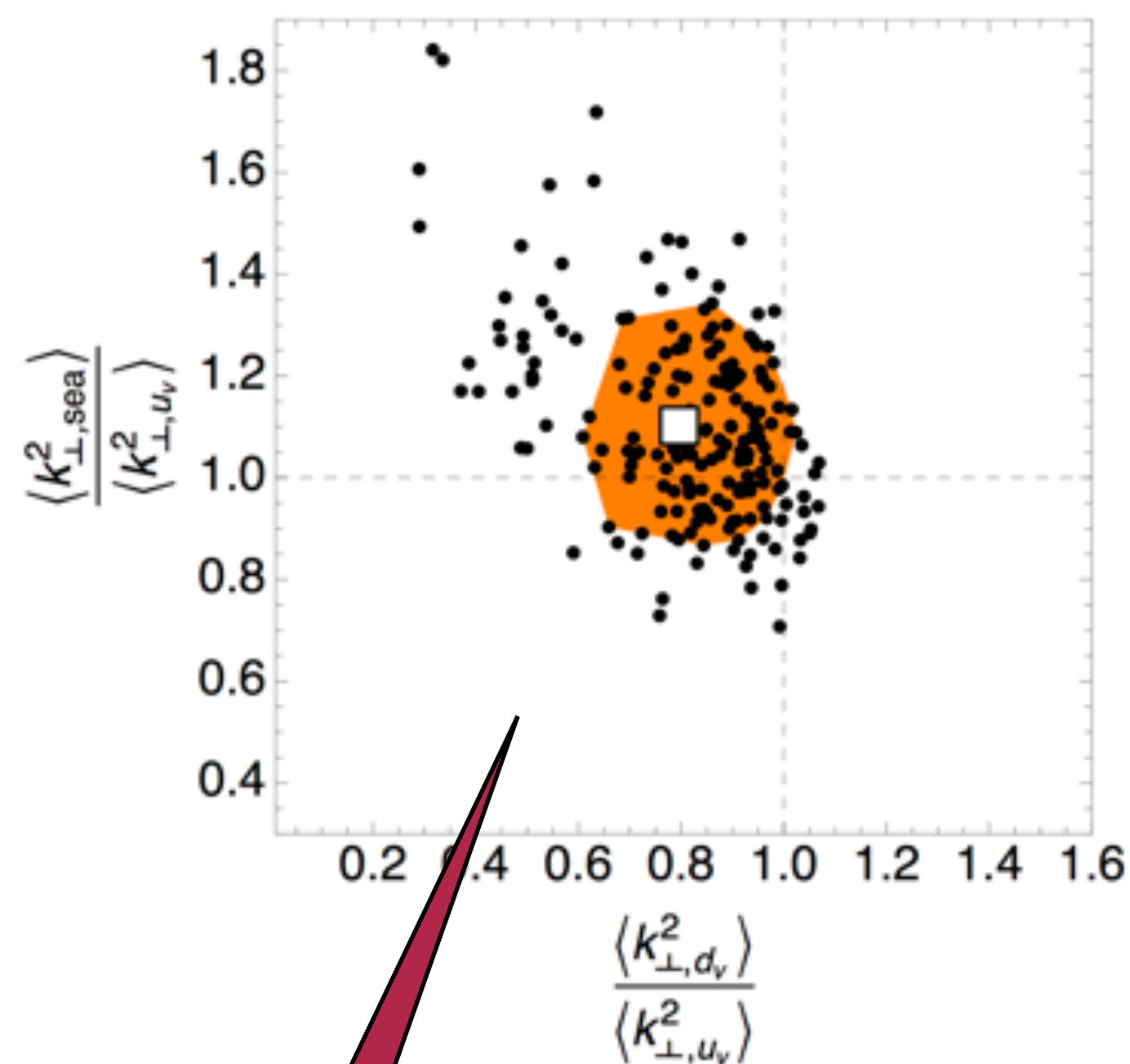
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FLAVOR DEPENDENCE OF TMDS

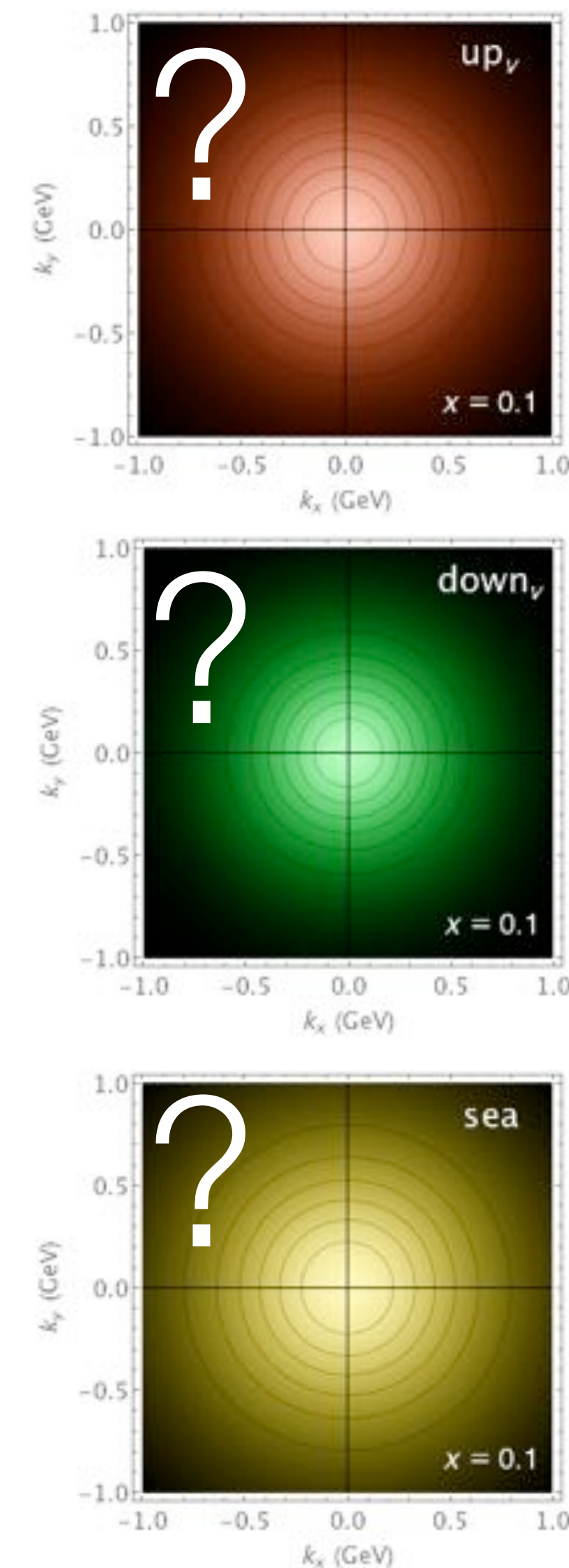
Signori, Bacchetta, Radici, Schnell JHEP 1311 (13)

Ratio of width of sea /
width of up valence



Ratio width of down valence/
width of up valence

**First investigation 10 years ago:
there is room for flavor dependence,
but we don't control it well**



ART23

[Moos, Scimemi, Vladimirov, Zurita, 2305.07473](#)

N⁴LL- accuracy

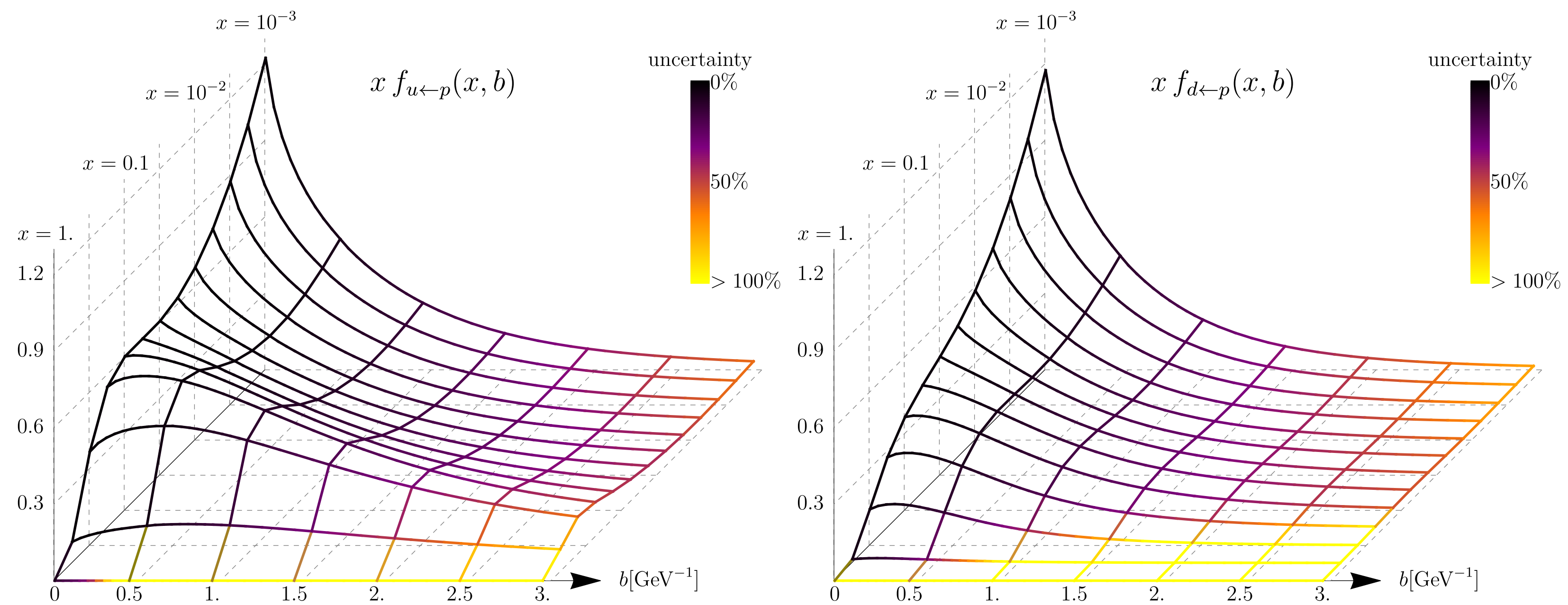
Drell-Yan only

ART23

[Moos, Scimemi, Vladimirov, Zurita, 2305.07473](#)

N⁴LL⁻ accuracy

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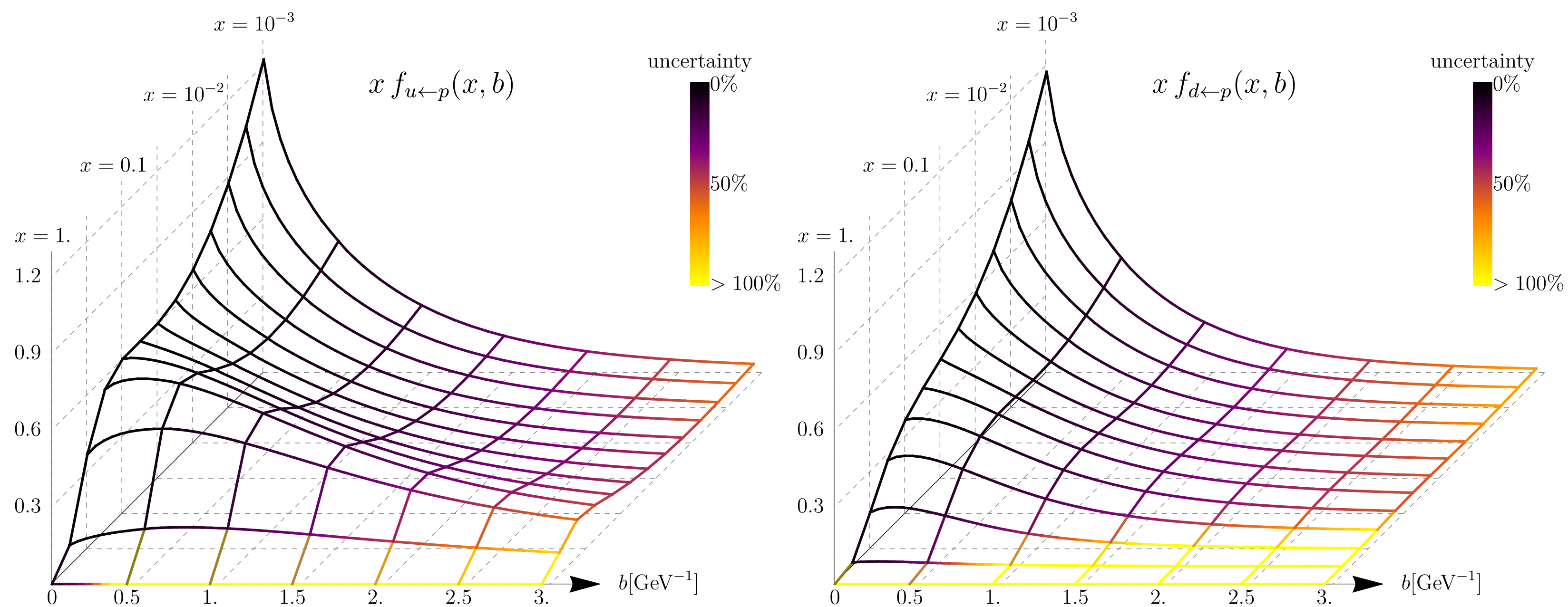
ART23

[Moos, Scimemi, Vladimirov, Zurita, 2305.07473](#)

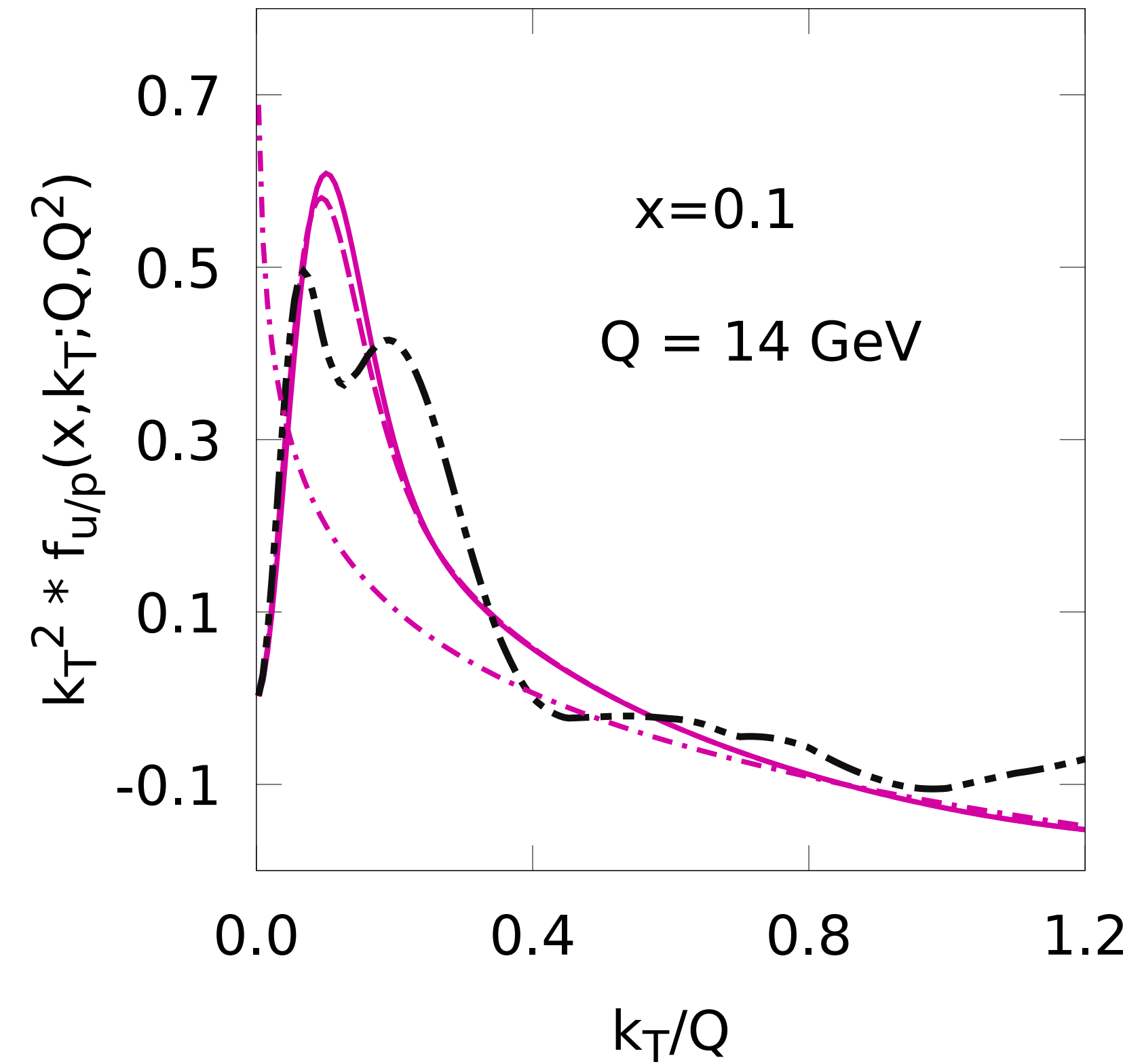
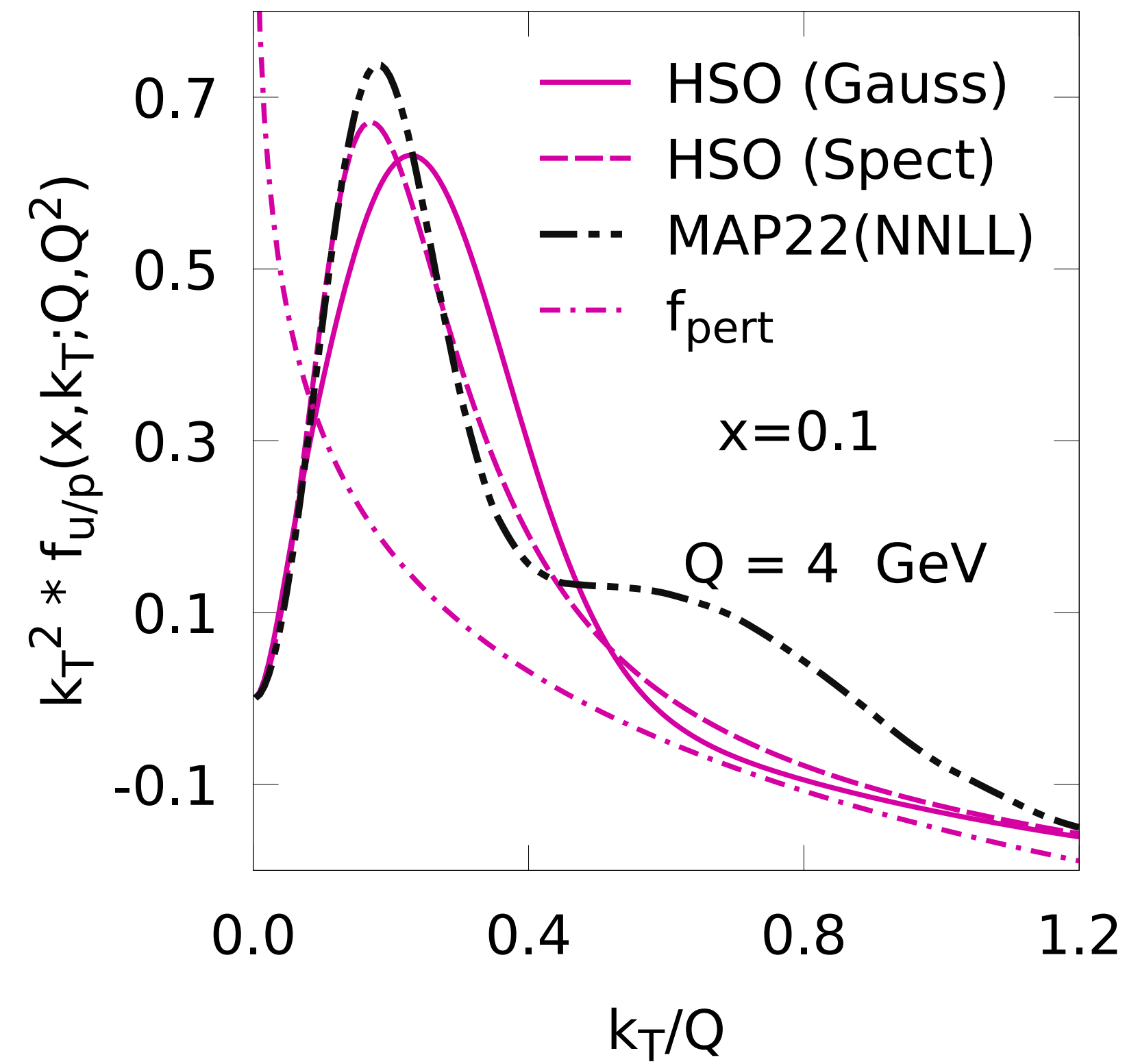
N⁴LL⁻ accuracy

Drell-Yan only

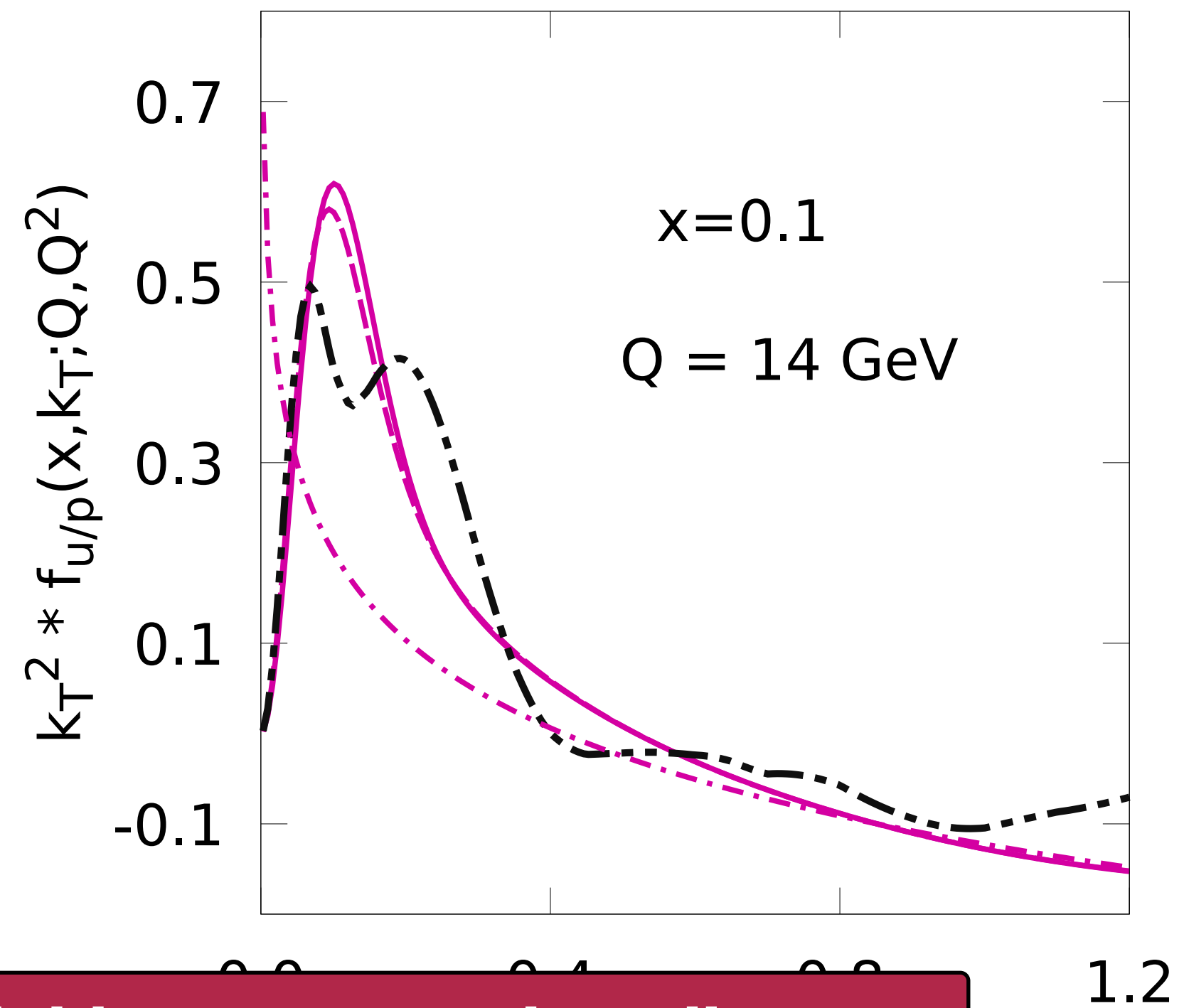
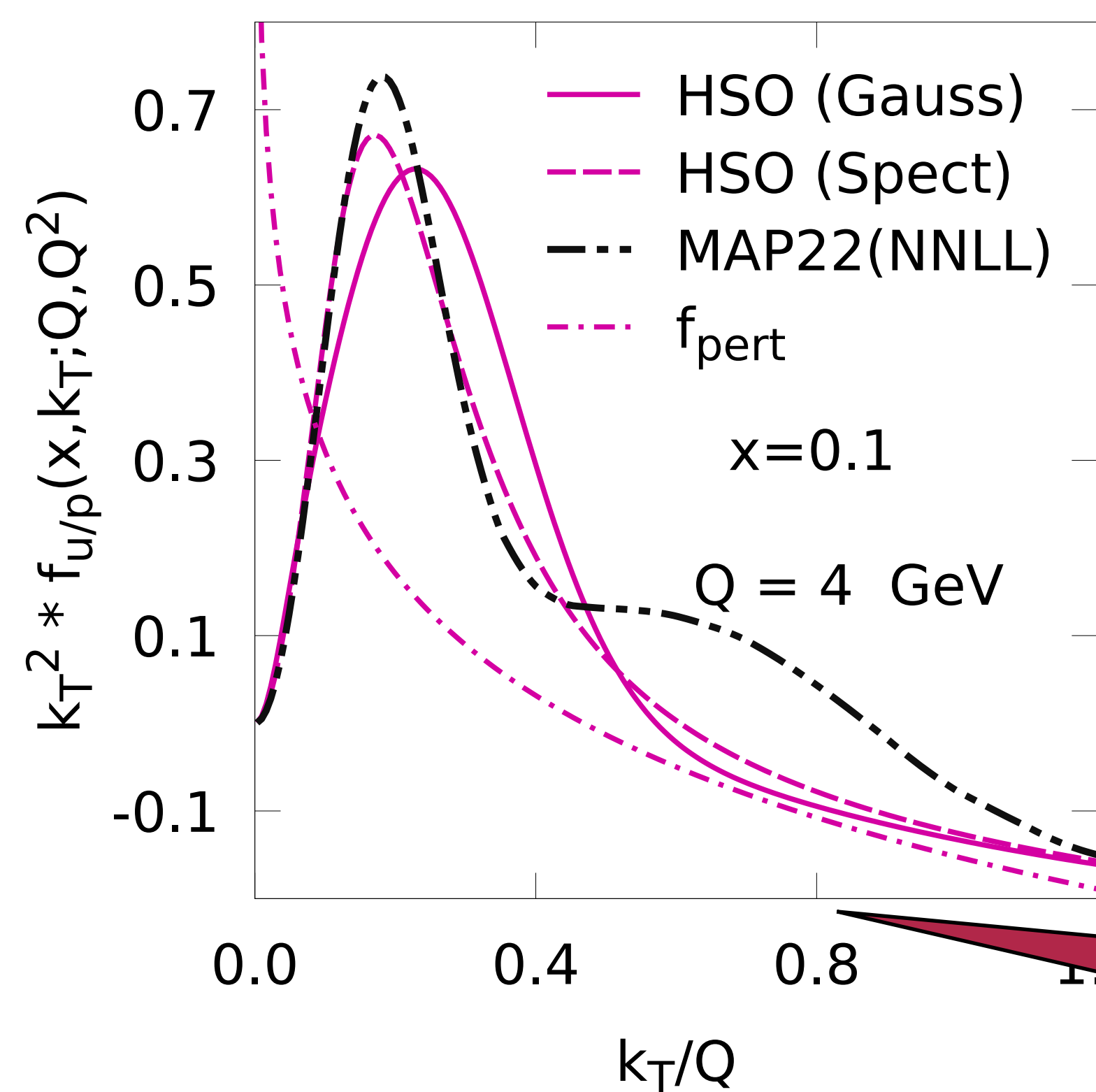
Different up and down TMDs



[Aslan, Boglione, Gonzalez-Hernandez, Rainaldi, Rogers, Simonelli, 2401.14266](#)

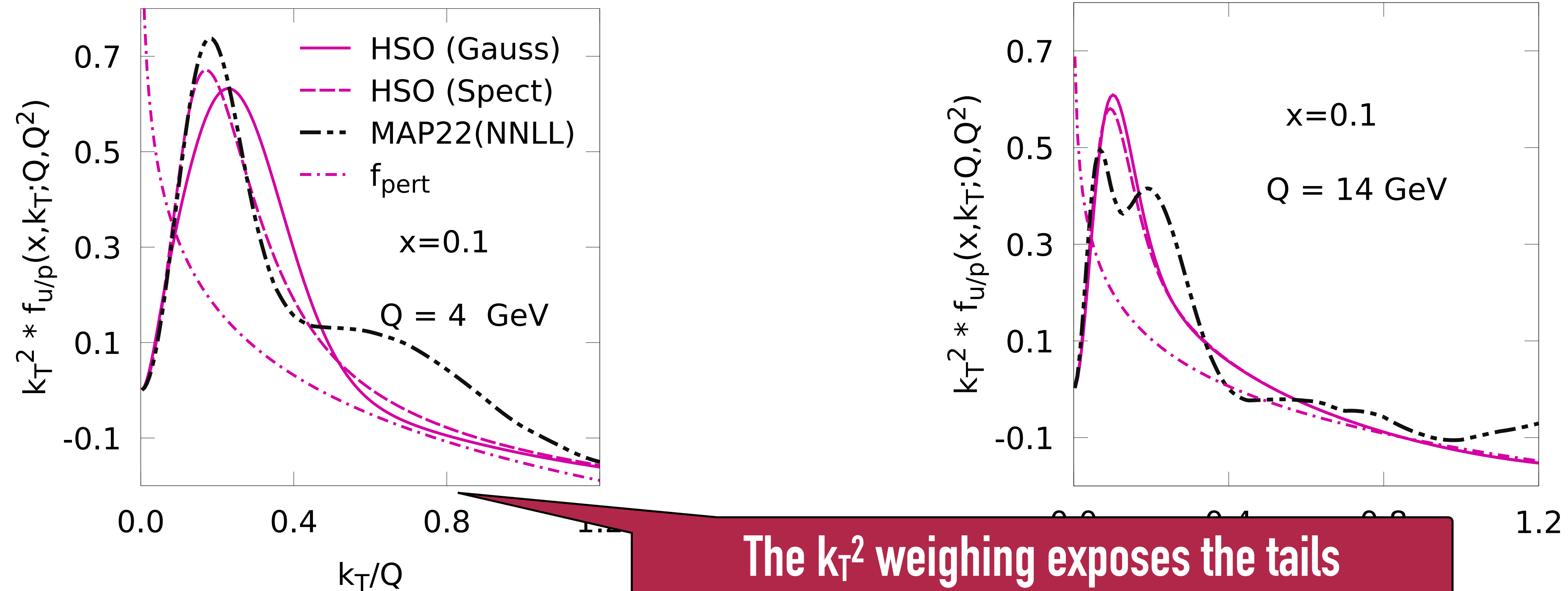


[Aslan, Boglione, Gonzalez-Hernandez, Rainaldi, Rogers, Simonelli, 2401.14266](#)



The k_T^2 weighing exposes the tails
(multiplied by a factor of 10 in this case)

[Aslan, Boglione, Gonzalez-Hernandez, Rainaldi, Rogers, Simonelli, 2401.14266](#)



The k_T^2 weighing exposes the tails
(multiplied by a factor of 10 in this case)

The paper emphasizes the relevance of prescription choices and simultaneous TMD-PDF fit, but does not provide a fit to extended data sets.

PDFS

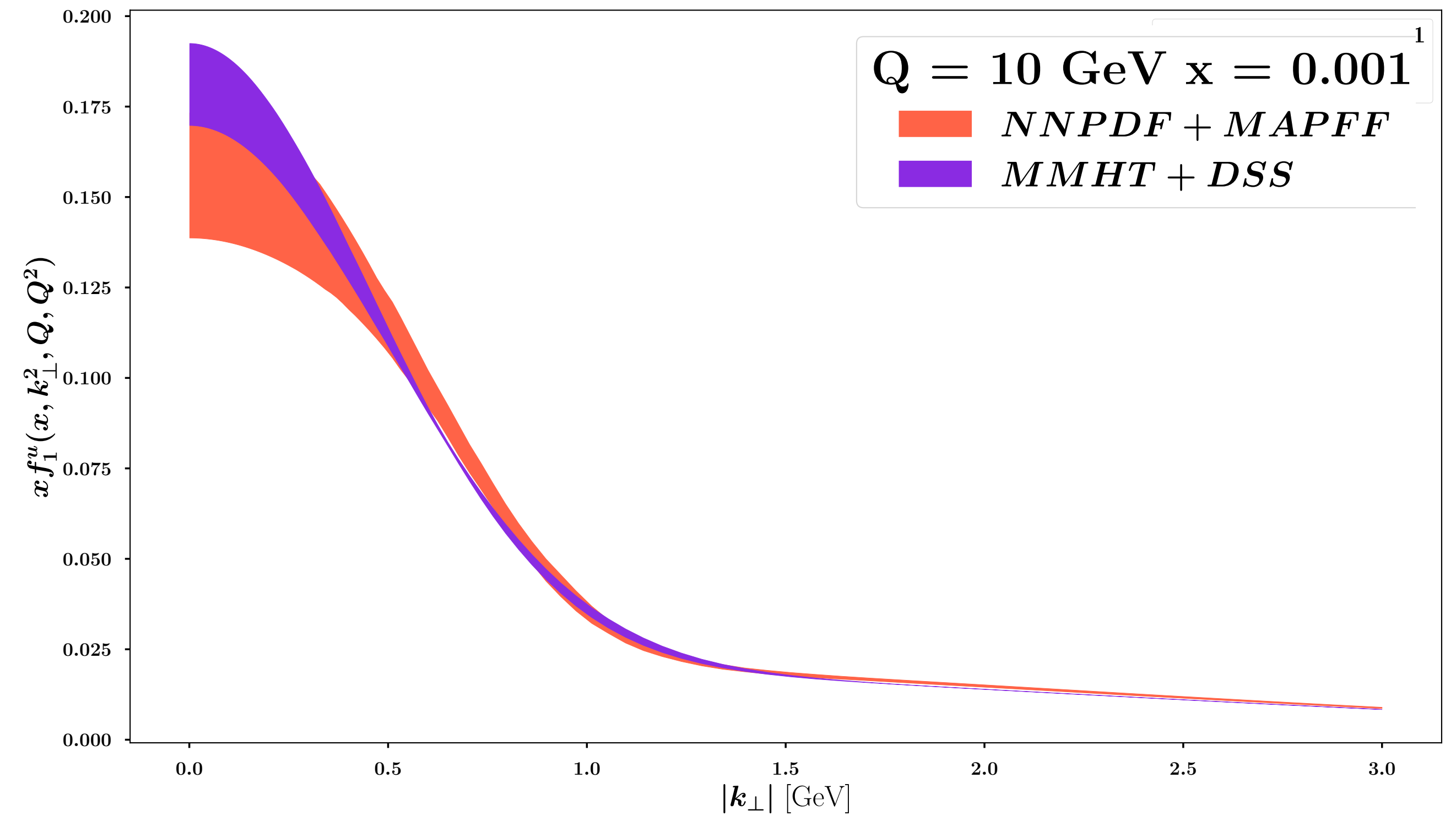
MMHT2014

Hessian set



NNPDF3.1

Monte Carlo set



PDFS

MMHT2014

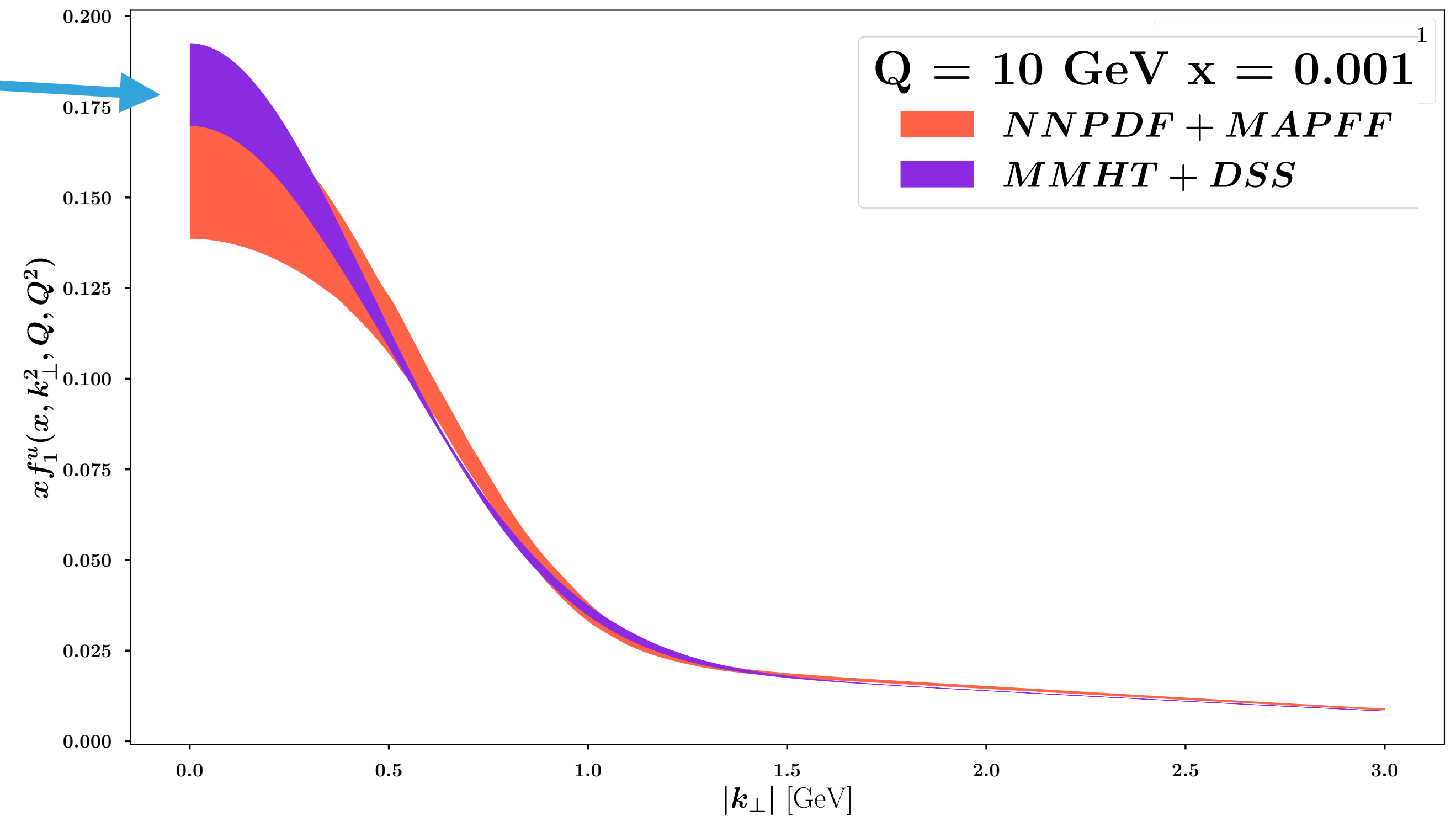
Hessian set



NNPDF3.1

Monte Carlo set

MAP22 fit



PDFS

MMHT2014

Hessian set

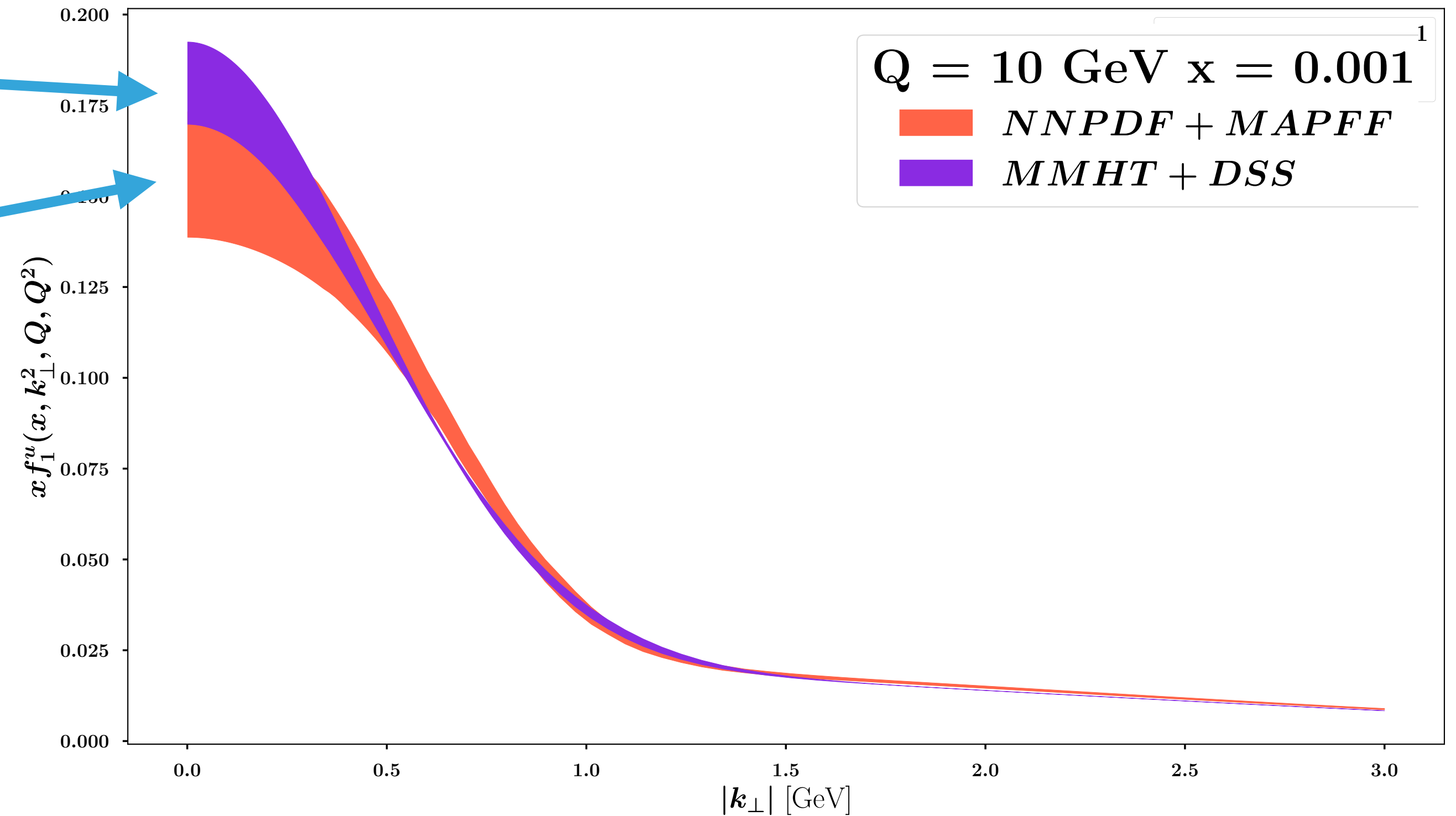
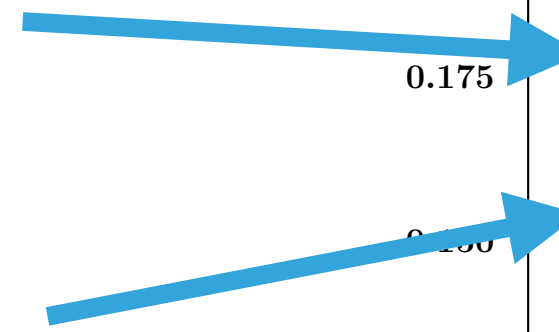


NNPDF3.1

Monte Carlo set

MAP22 fit

Fit with NNPDF set



PDFS

MMHT2014

Hessian set

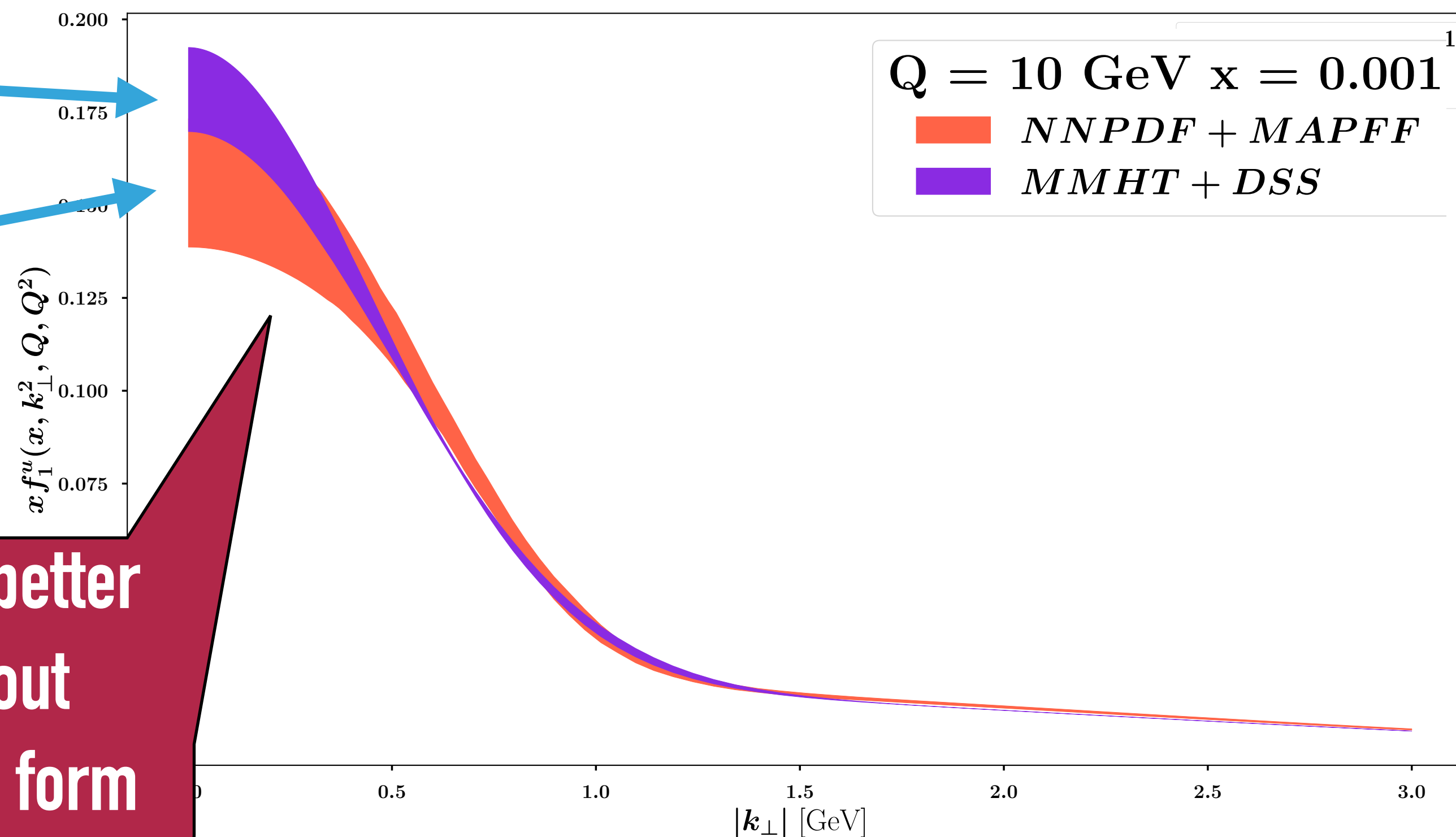


NNPDF3.1

Monte Carlo set

MAP22 fit

Fit with NNPDF set



The use of a Monte Carlo PDF set gives a better idea of the full TMD uncertainties, without dramatically changing the TMD functional form

Nonpert. TMD components of FF
equal for pions and kaons

Nonpert. TMD components of FF
equal for pions and kaons

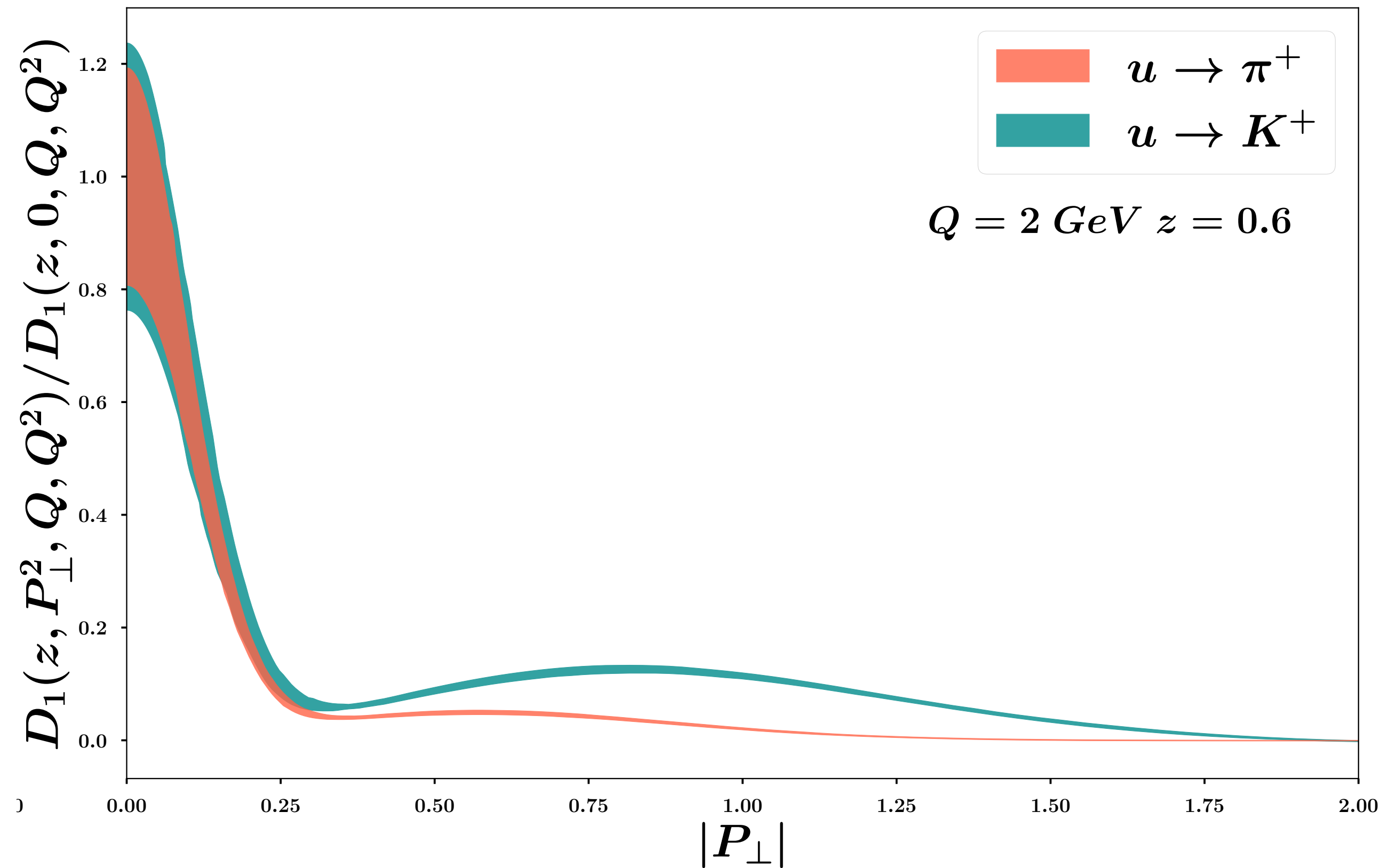


Distinction between pions and kaons

Nonpert. TMD components of FF
equal for pions and kaons



Distinction between pions and kaons



LESSONS LEARNED

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- ▶ Simple Gaussians or bell-like shapes are not sufficient to describe data

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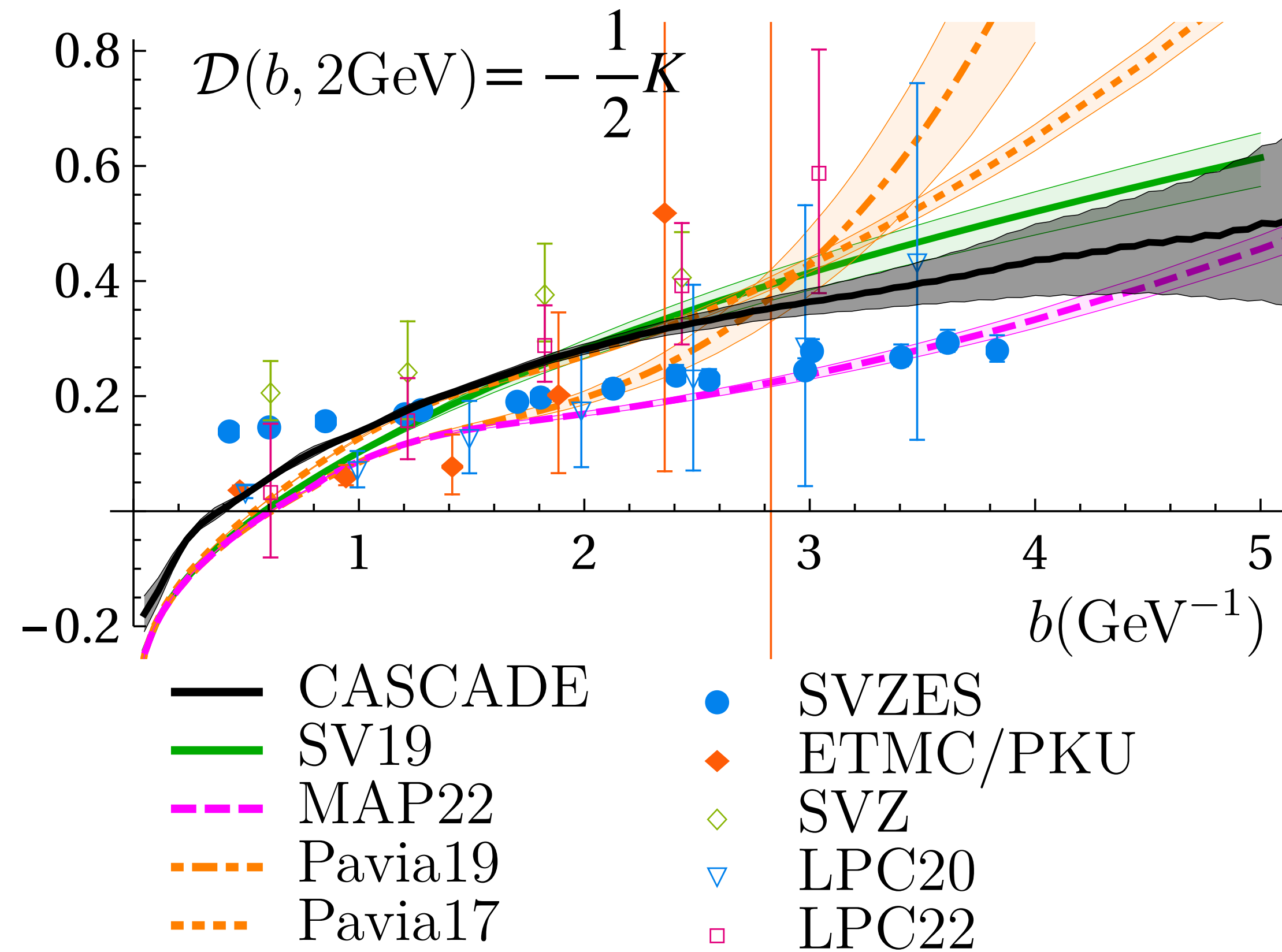
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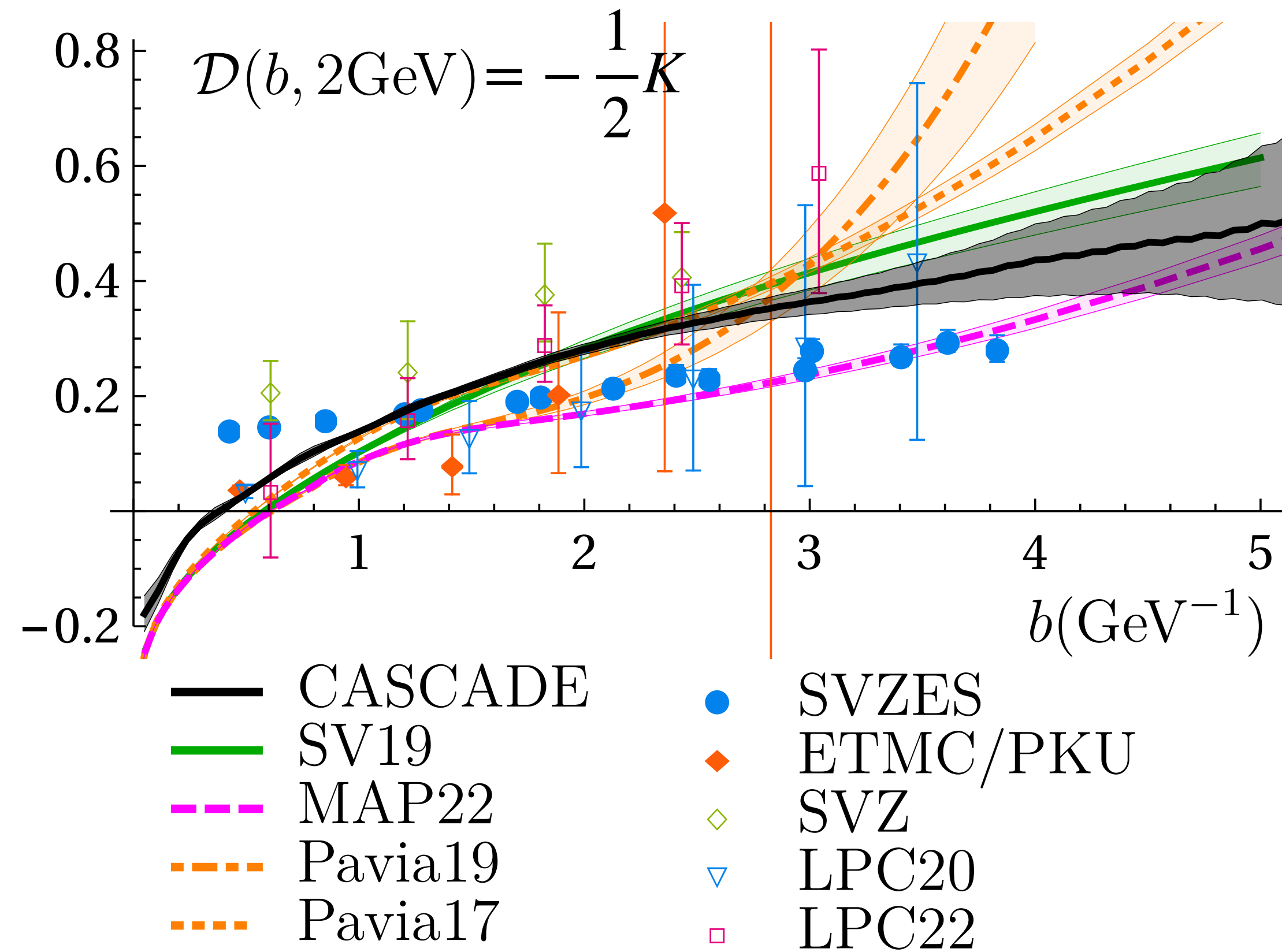
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- ▶ The TMDs are probably different for different quark flavors

Bermudez Martinez, Vladimirov, arXiv:2206.01105

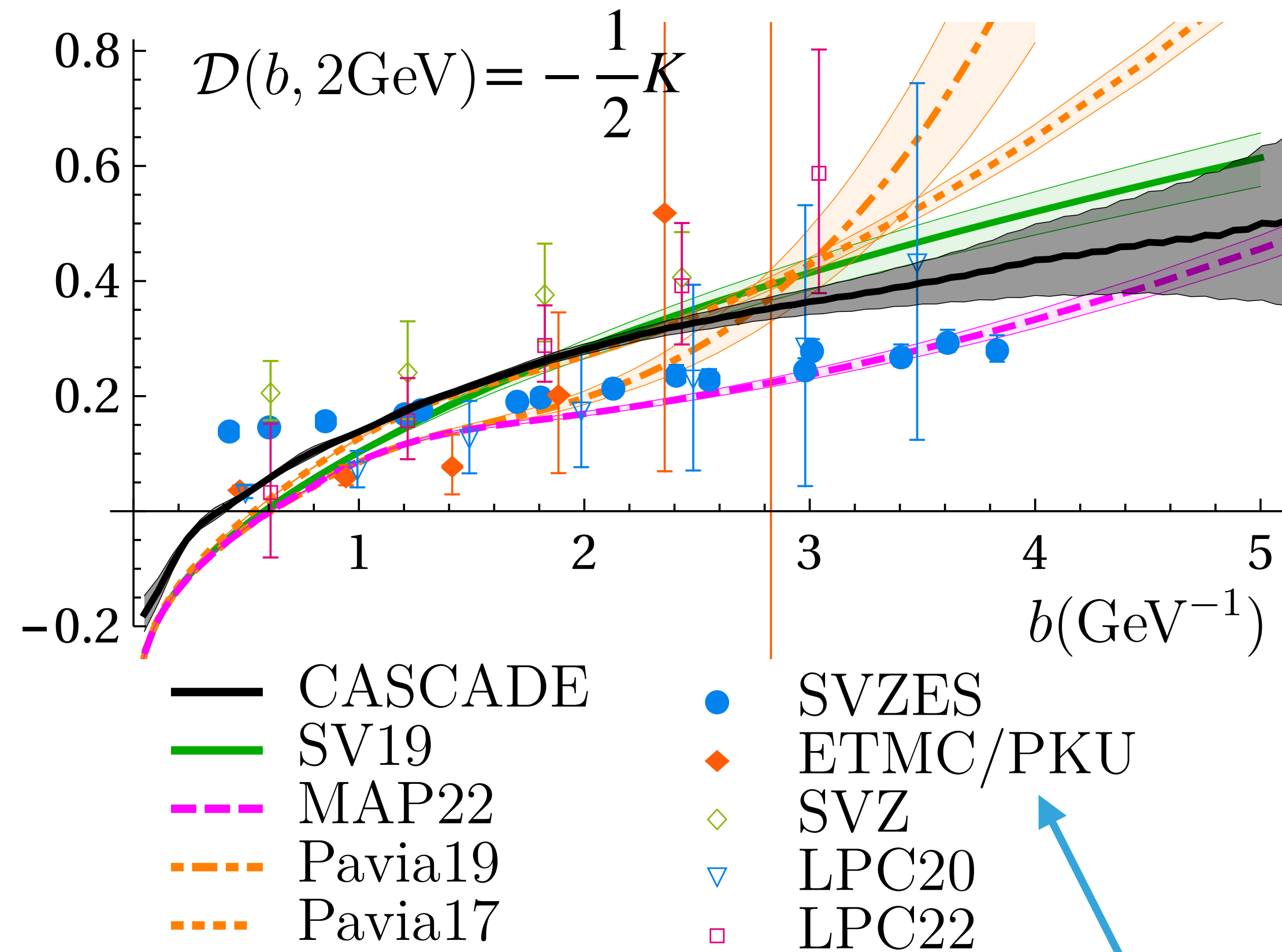


[Bermudez Martinez, Vladimirov, arXiv:2206.01105](#)



TMD phenomenology

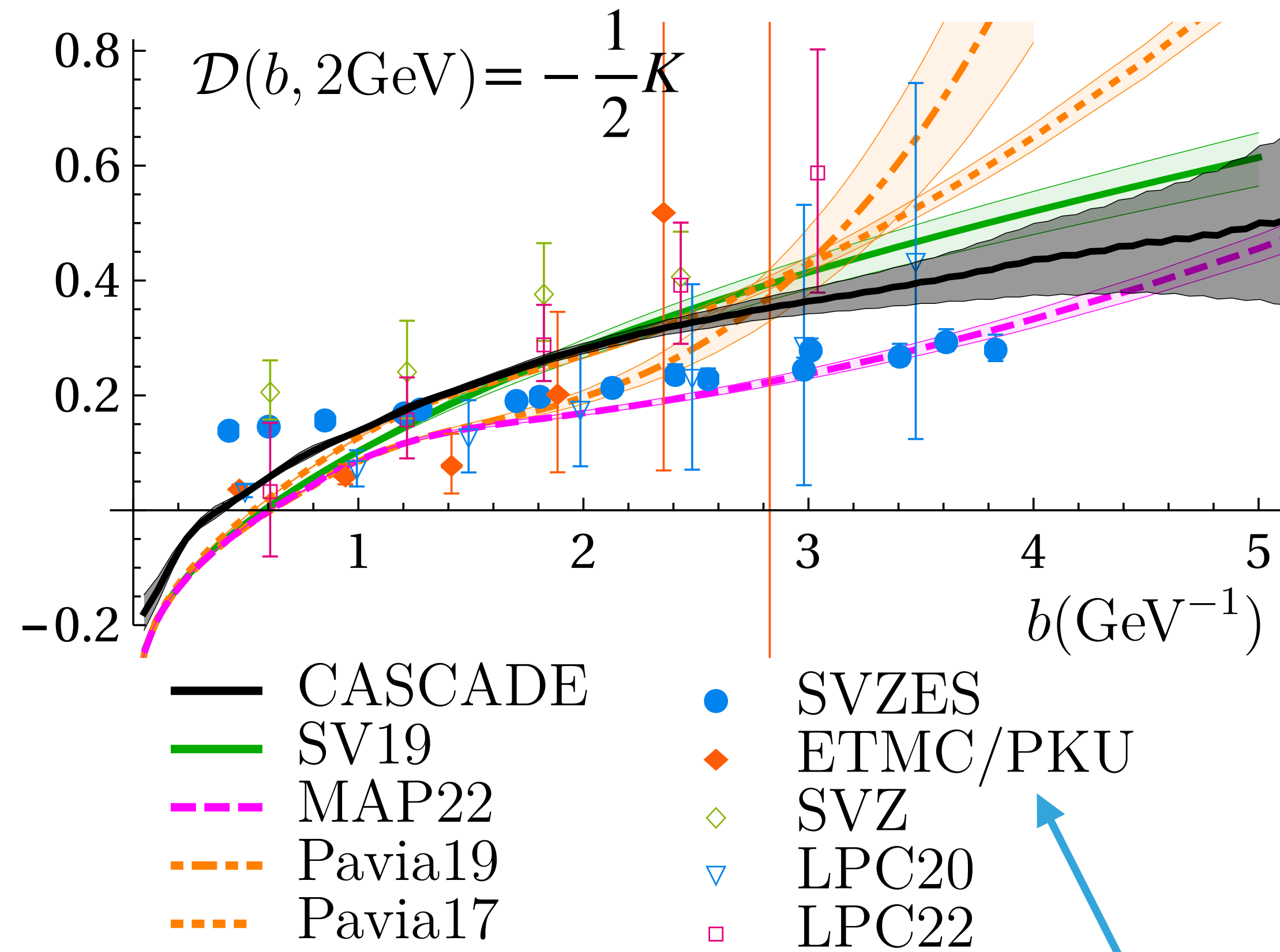
Bermudez Martinez, Vladimirov, arXiv:2206.01105



TMD phenomenology

Lattice QCD

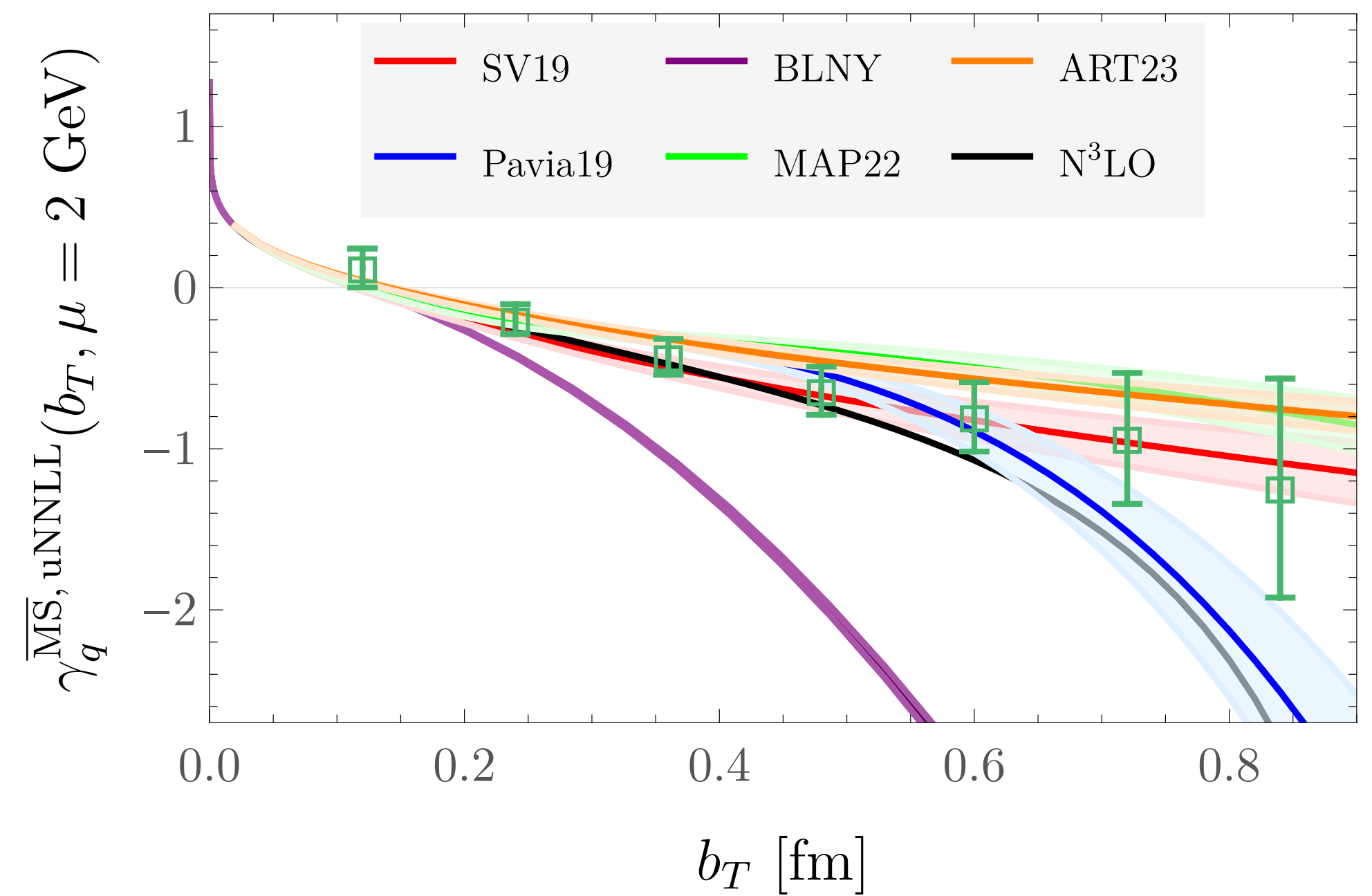
[Bermudez Martinez, Vladimirov, arXiv:2206.01105](#)

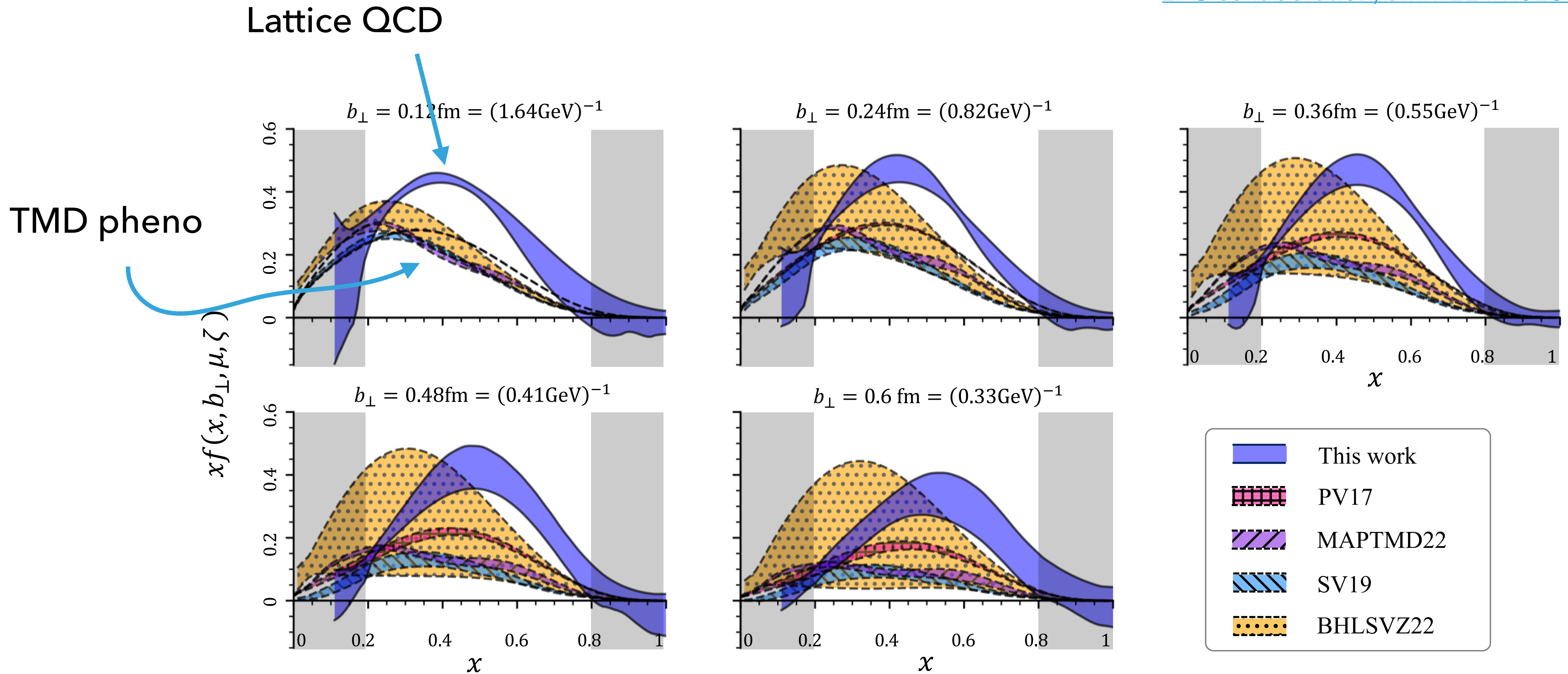


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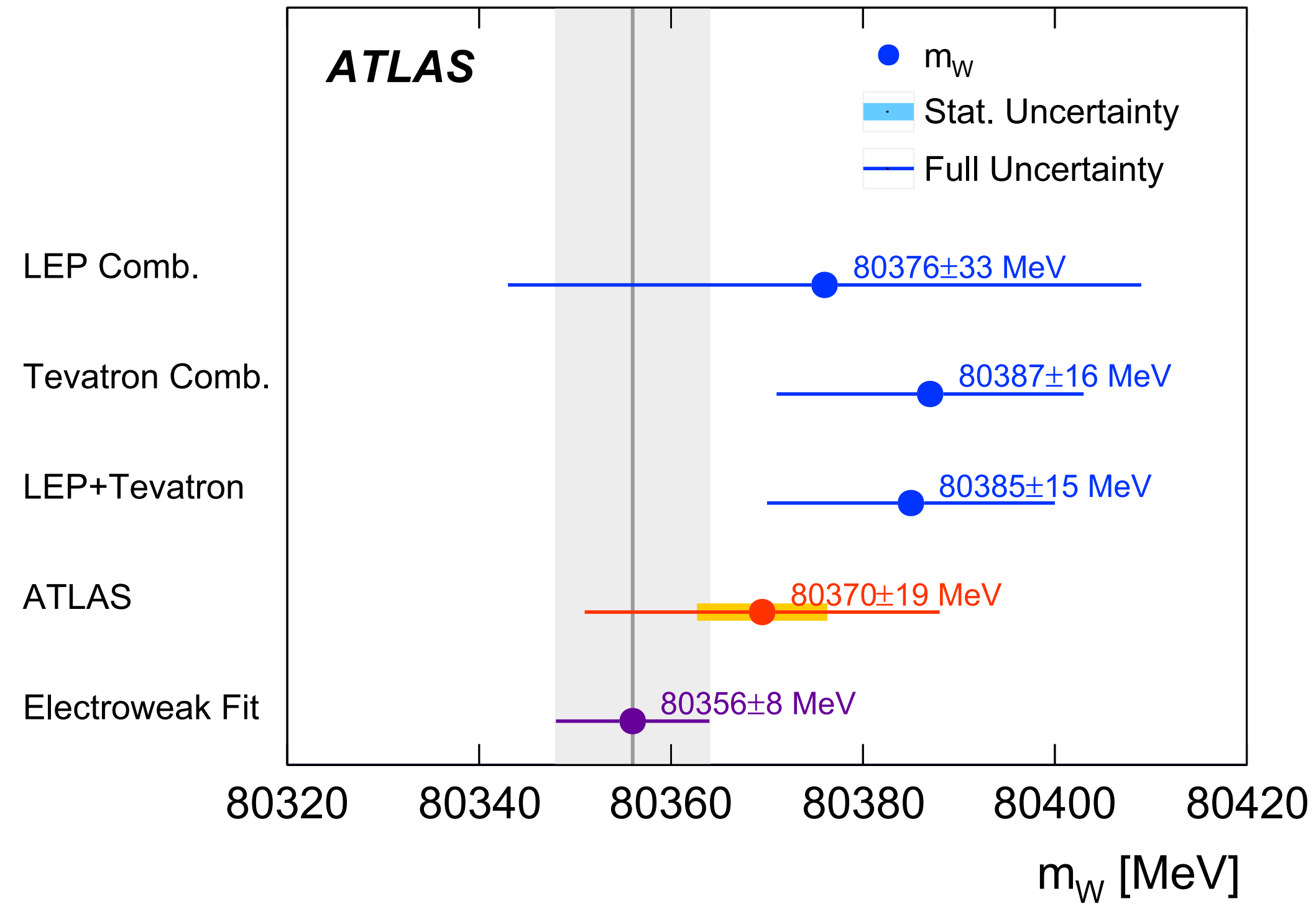
Lattice QCD

[Avkhadiev, Shanahan, Wagman, Zhao, arXiv:2307.12359](#)





[ATLAS Collab. arXiv:1701.07240](https://arxiv.org/abs/1701.07240)



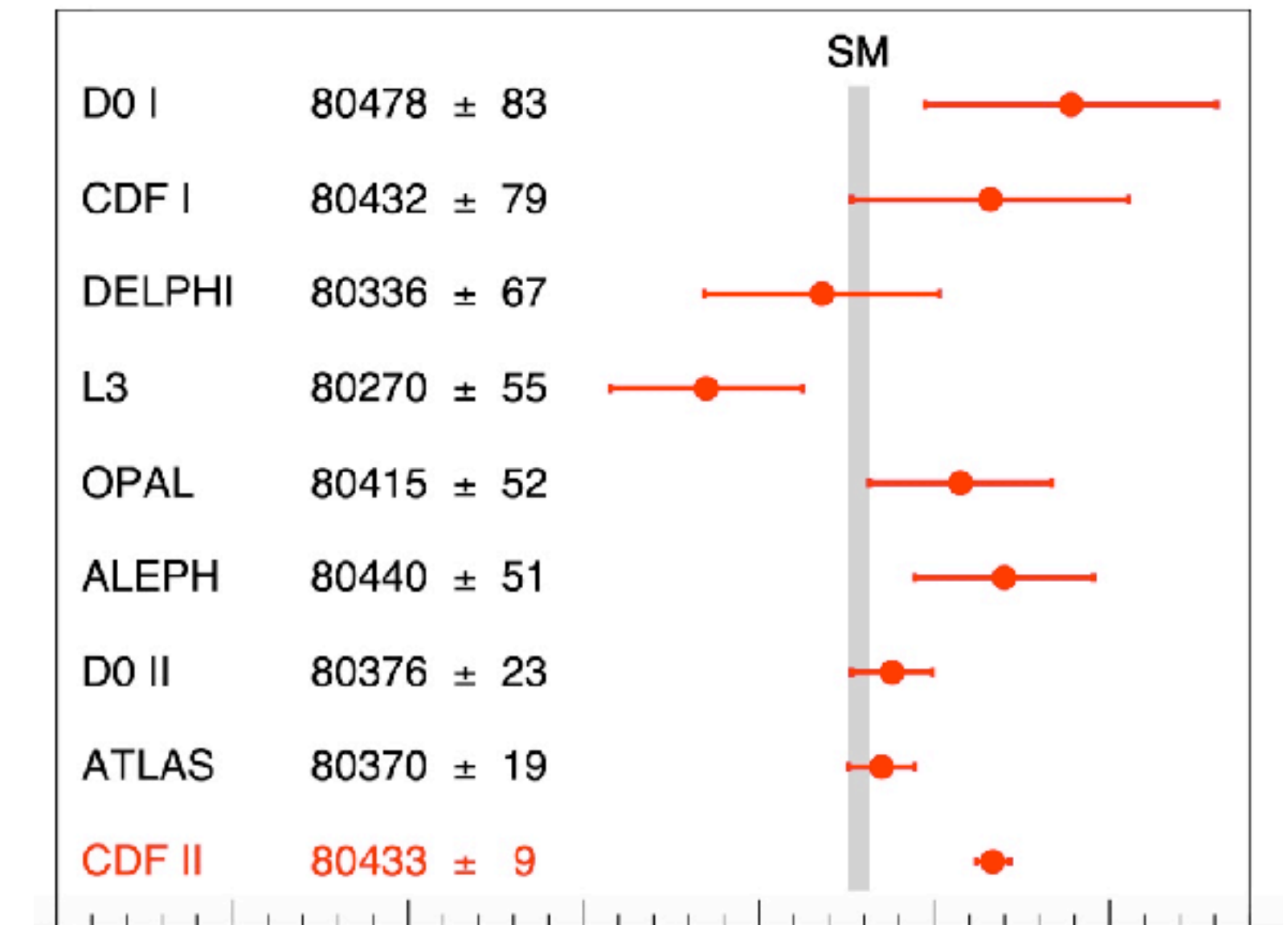
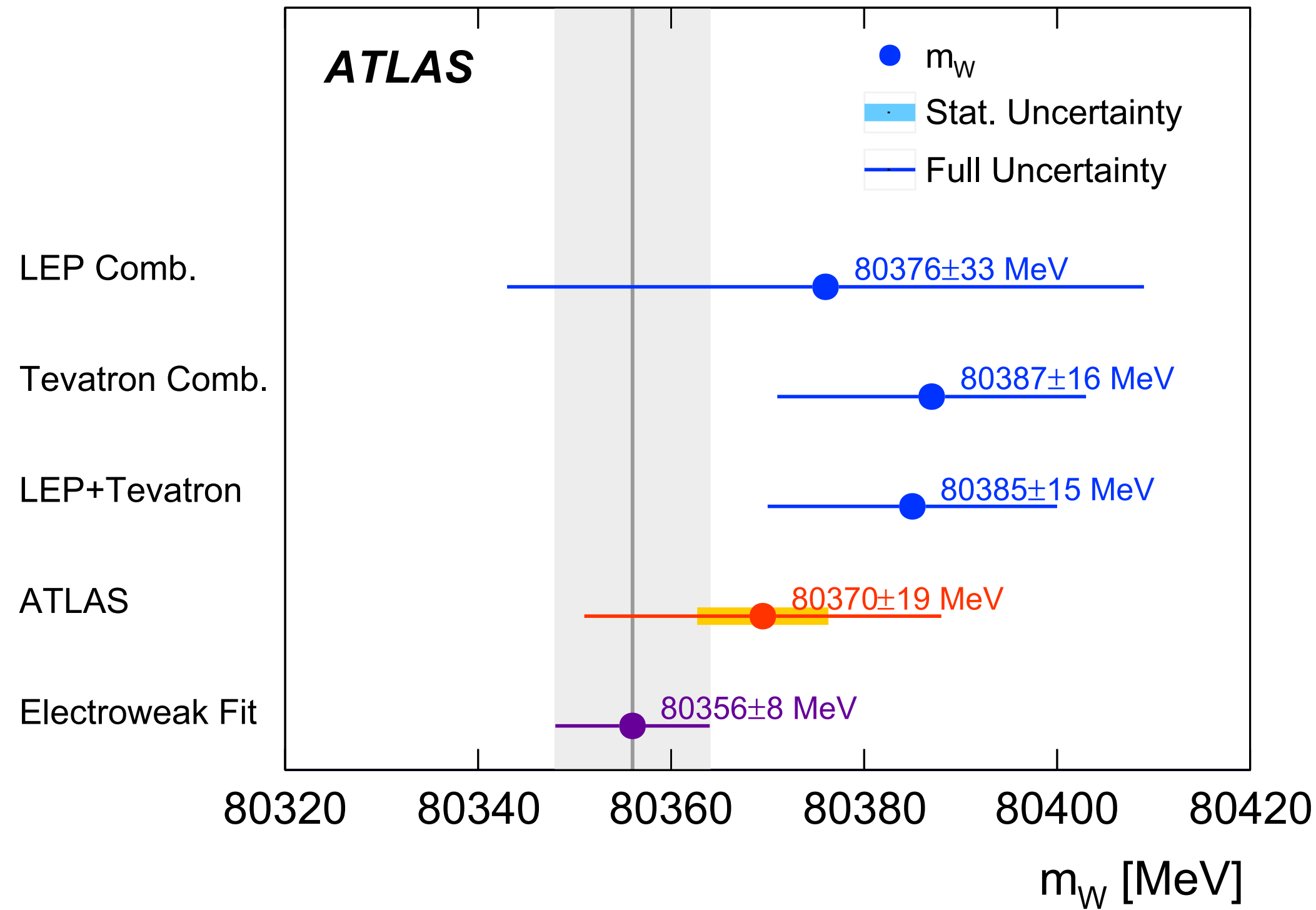
$$m_W = 80370 \pm 7 \text{ (stat.)} \pm 11 \text{ (exp. syst.)} \pm 14 \text{ (mod. syst.) MeV}$$

$$= 80370 \pm 19 \text{ MeV,}$$

$$m_{W^+} - m_{W^-} = -29 \pm 28 \text{ MeV.}$$

[ATLAS Collab. arXiv:1701.07240](https://arxiv.org/abs/1701.07240)

[CDF Collab.. Science 2022](https://arxiv.org/abs/2106.11977)



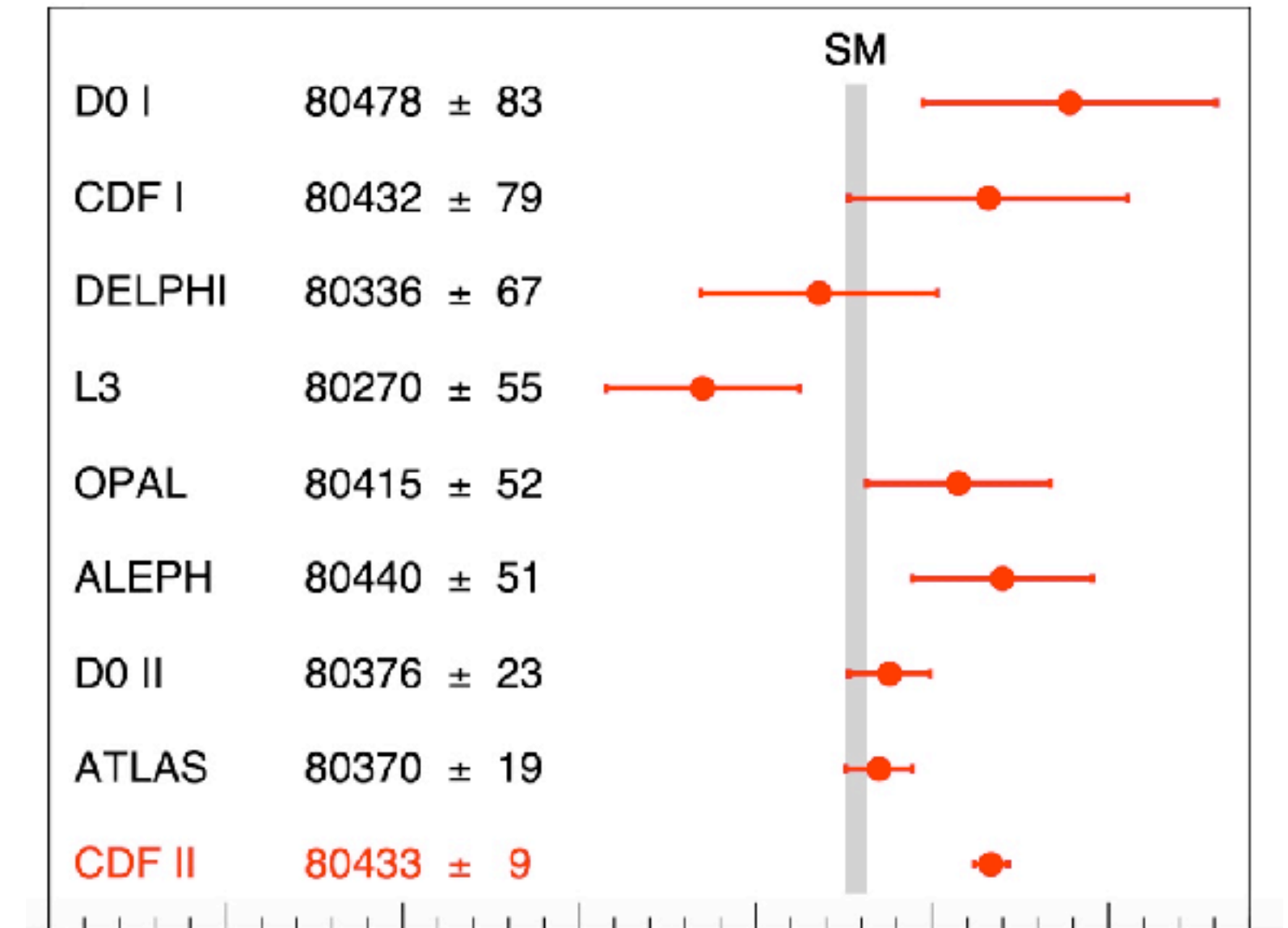
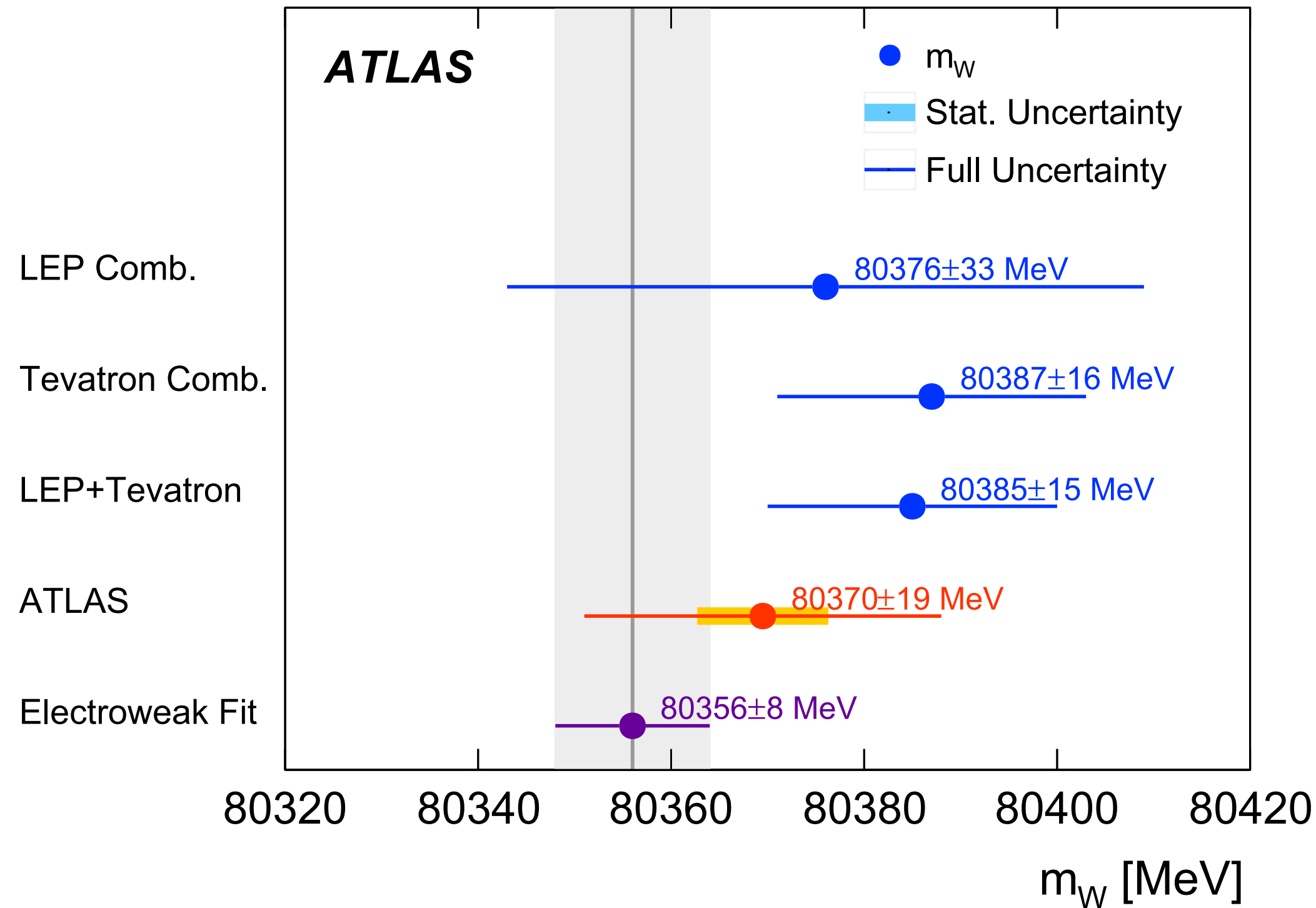
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All analyses assume that TMDs are not flavor dependent. What happens if they are?

[Bacchetta, Bozzi, Radici, Ritzmann, Signori, arXiv:1807.02101](#)

Try some judicious choices of flavour dependent widths and check

[Bacchetta, Bozzi, Radici, Ritzmann, Signori, arXiv:1807.02101](#)

Try some judicious choices of flavour dependent widths and check

Set	u_v	d_v	u_s	d_s	s
1	0.34	0.26	0.46	0.59	0.32
2	0.34	0.46	0.56	0.32	0.51
3	0.55	0.34	0.33	0.55	0.30
4	0.53	0.49	0.37	0.22	0.52
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narrow, medium, large
narrow, large, narrow
large, narrow, large
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medium, narrow, large

[Bacchetta, Bozzi, Radici, Ritzmann, Signori, arXiv:1807.02101](#)

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	ΔM_{W+}		ΔM_{W-}	
Set	m_T	$p_{T\ell}$	m_T	$p_{T\ell}$
1	0	-1	-2	3
2	0	-6	-2	0
3	-1	9	-2	-4
4	0	0	-2	-4
5	0	4	-1	-3

[Bacchetta, Bozzi, Radici, Ritzmann, Signori, arXiv:1807.02101](#)

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Not taking into account the flavor dependence of TMDs can lead to errors in the determination of the W mass, of the order of a few MeVs

[Bacchetta, Bozzi, Radici, Ritzmann, Signori, arXiv:1807.02101](#)

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 narrow, large, narrow
 large, narrow, large
 large, medium, narrow
 medium, narrow, large

They all describe the Z spectrum very well

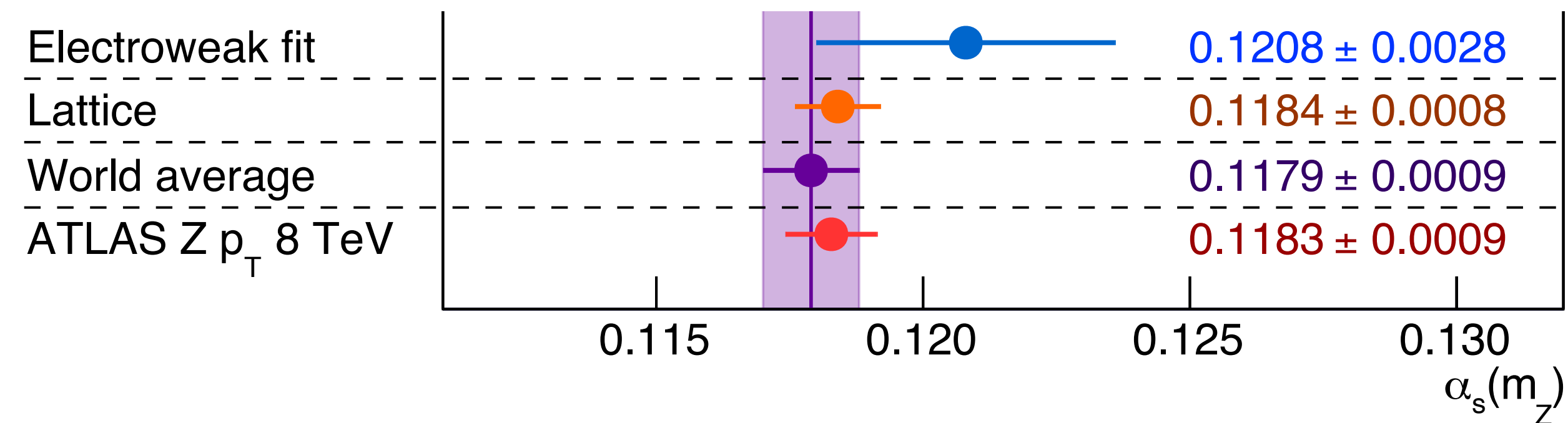


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The coupling constant of the strong force is determined from the transverse-momentum distribution of Z bosons produced in 8 TeV proton–proton collisions at the LHC and recorded by the ATLAS experiment.

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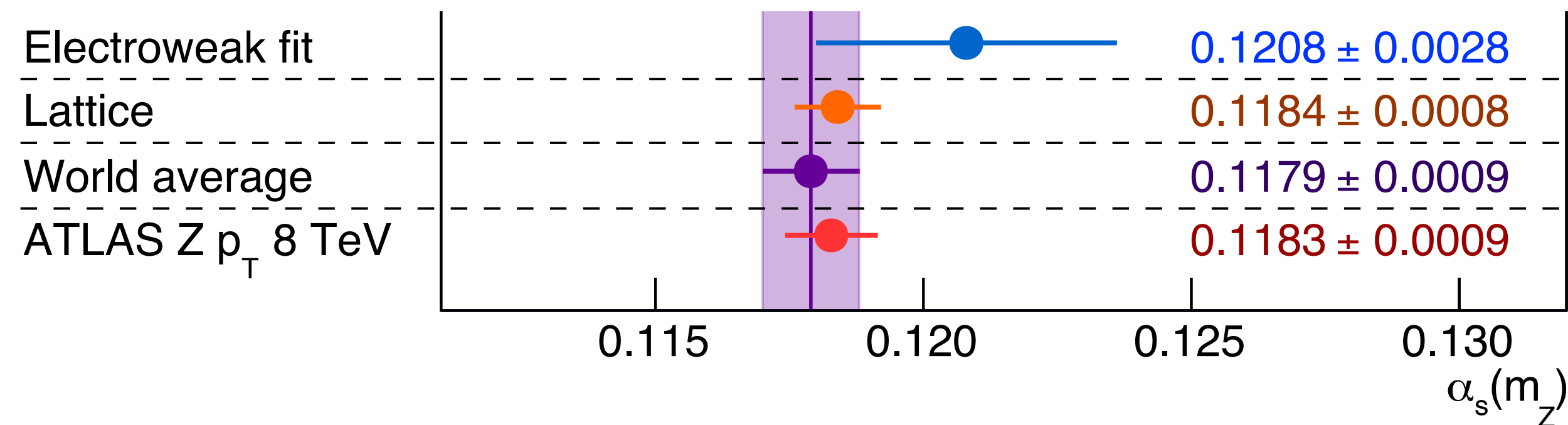
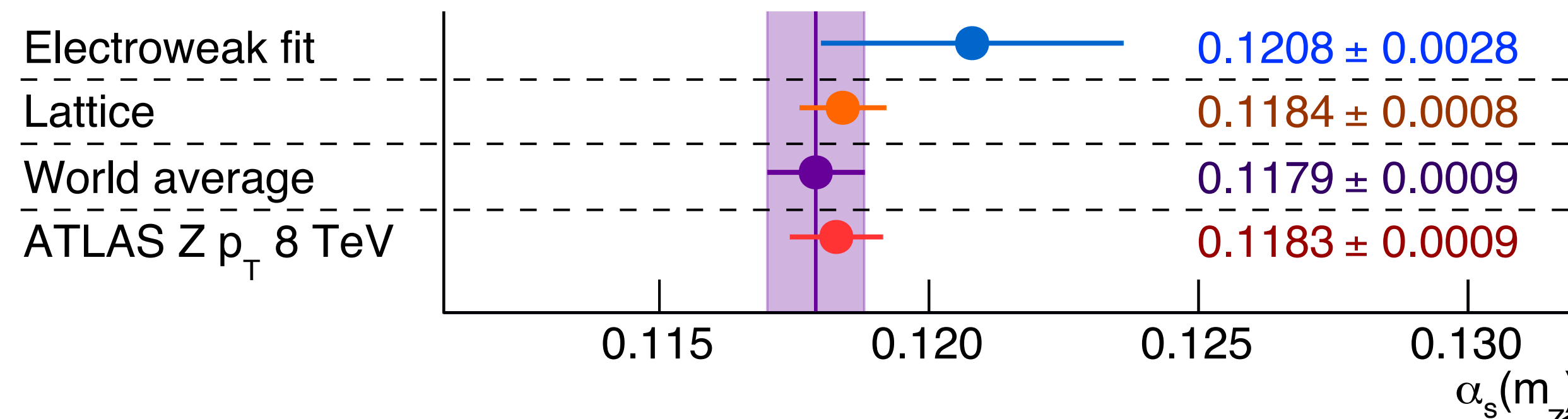


Table 1: Summary of the uncertainties in the determination of $\alpha_s(m_Z)$, in units of 10^{-3} .

Experimental uncertainty	± 0.44
PDF uncertainty	± 0.51
Scale variation uncertainties	± 0.42
Matching to fixed order	0 -0.08
Non-perturbative model	$+0.12$ -0.20
Flavour model	$+0.40$ -0.29
QED ISR	± 0.14
N^4LL approximation	± 0.04
Total	$+0.91$ -0.88

The coupling constant of the strong force is determined from the transverse-momentum distribution of Z bosons produced in 8 TeV proton–proton collisions at the LHC and recorded by the ATLAS experiment.

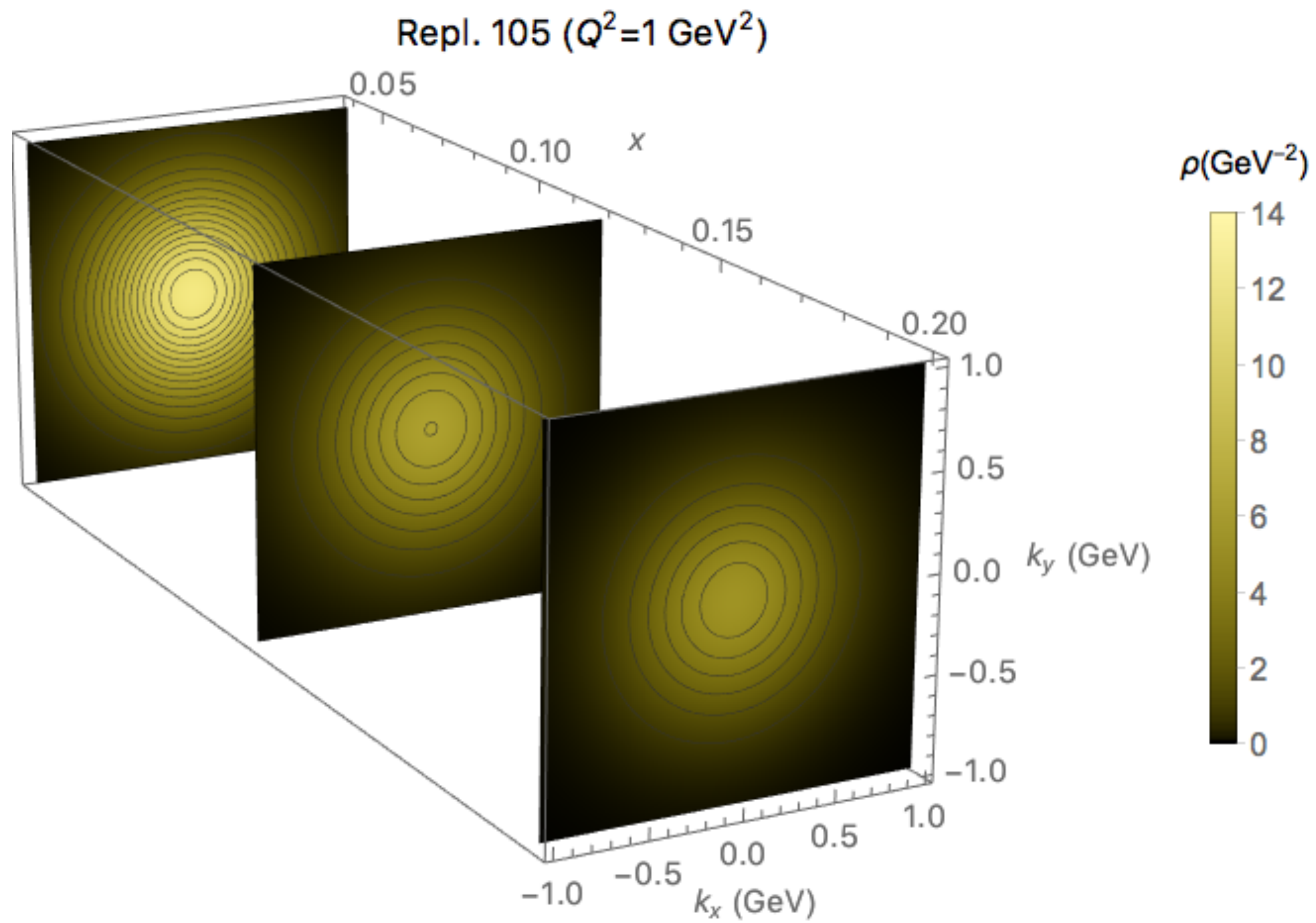


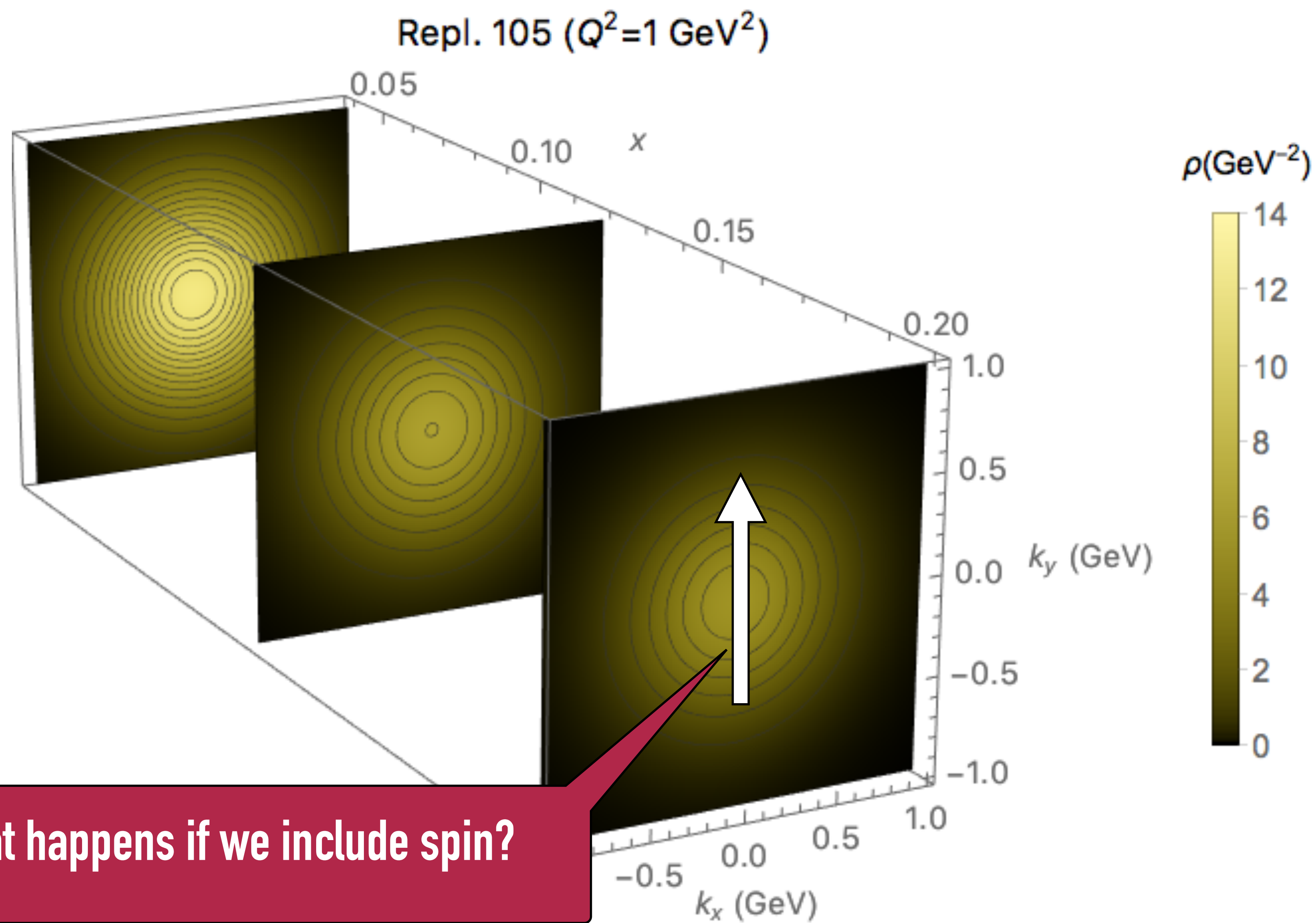
Estimate that does not take into account recent TMD results

Table 1: Summary of the uncertainties in the determination of $\alpha_s(m_Z)$, in units of 10^{-3} .

Experimental uncertainty	± 0.44	
PDF uncertainty	± 0.51	
Scale variation uncertainties	± 0.42	
Matching to fixed order	0	-0.08
Non-perturbative model	+0.12	-0.20
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SIVERS TMD





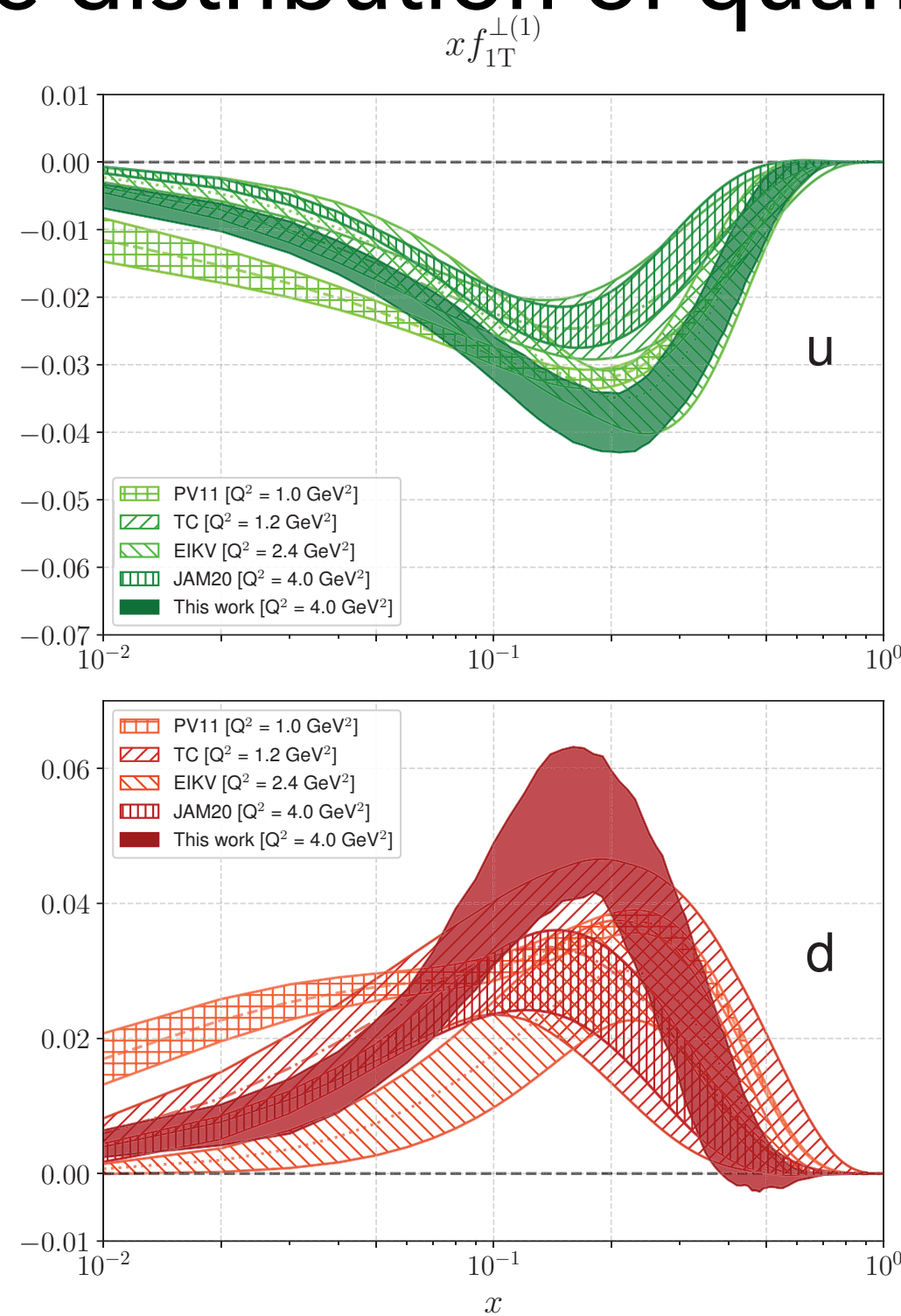
What happens if we include spin?

$$\rho_{N\uparrow}^q(x, k_x, k_y; Q^2) = f_1^q(x, k_T^2; Q^2) - \frac{k_x}{M} f_{1T}^{\perp q}(x, k_T^2; Q^2)$$

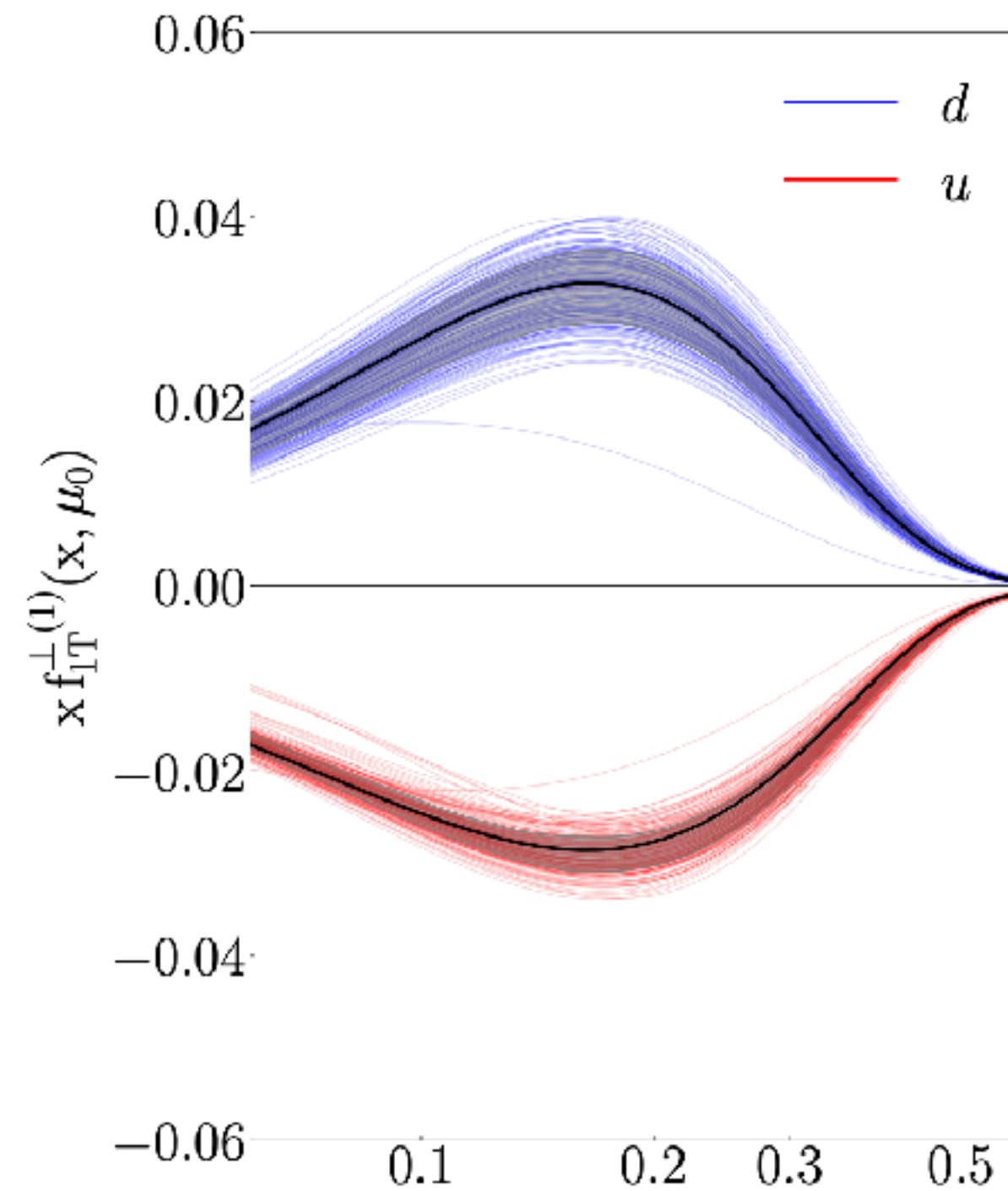
In a nucleon polarized in the +y direction,
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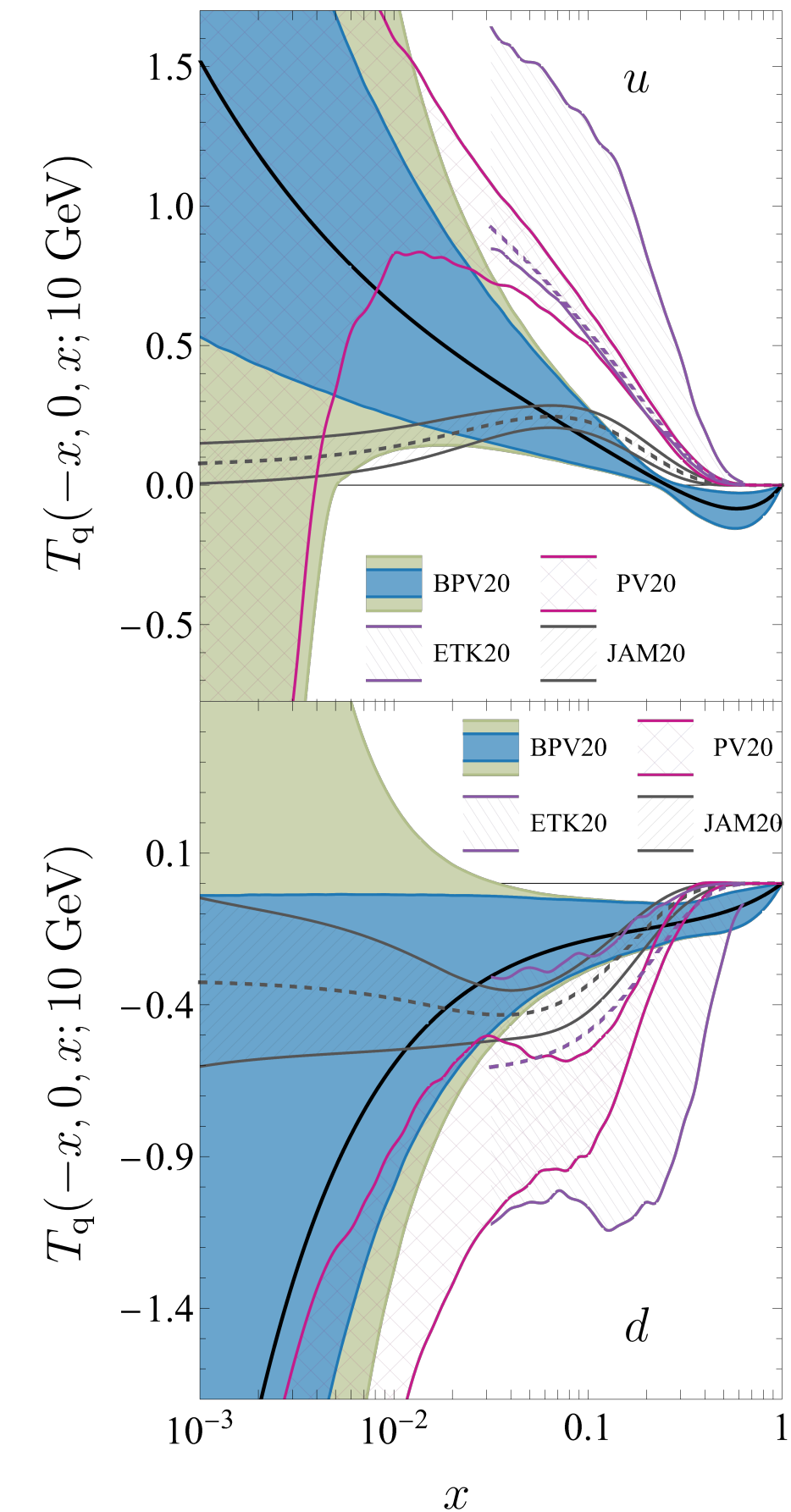
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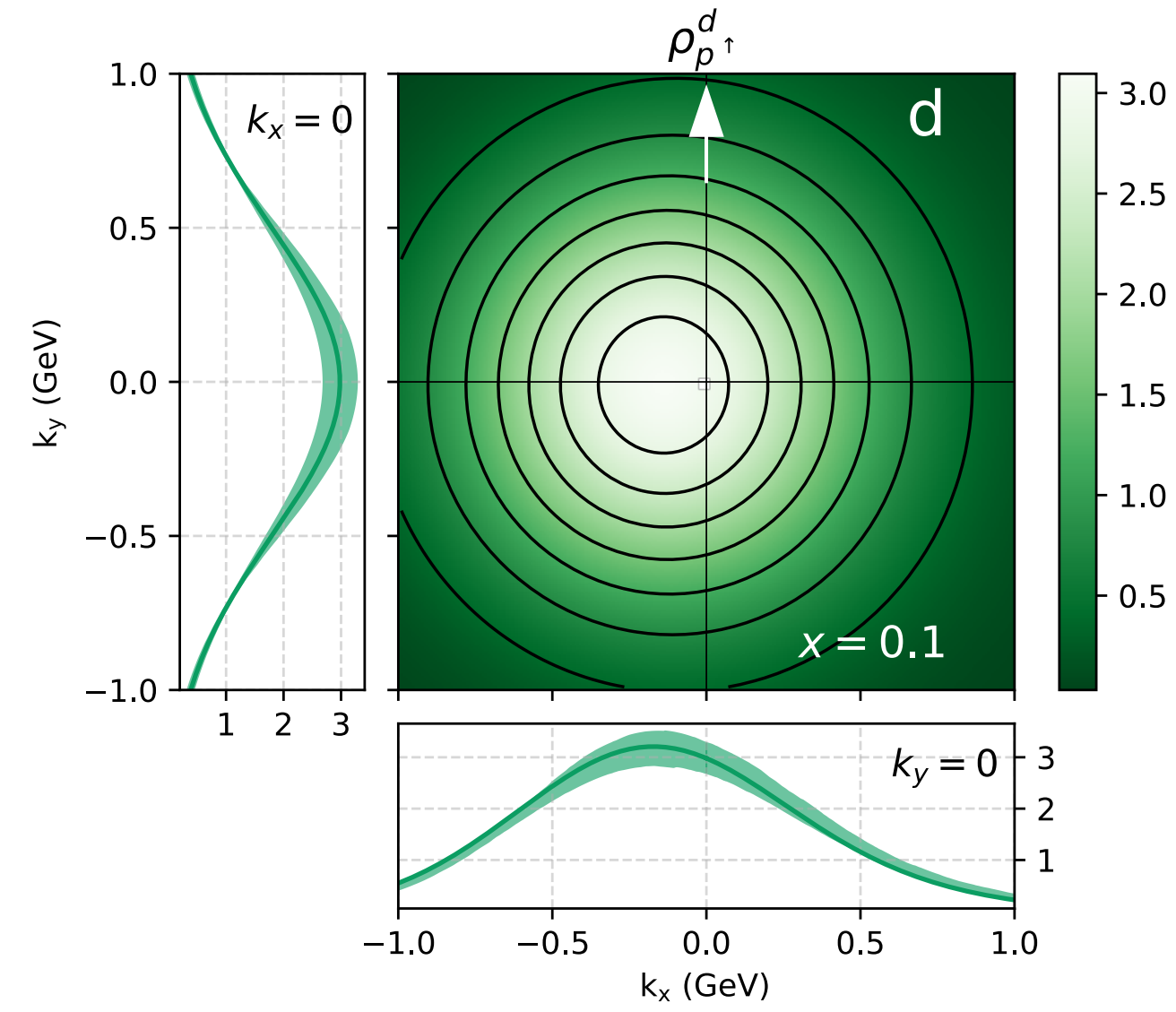
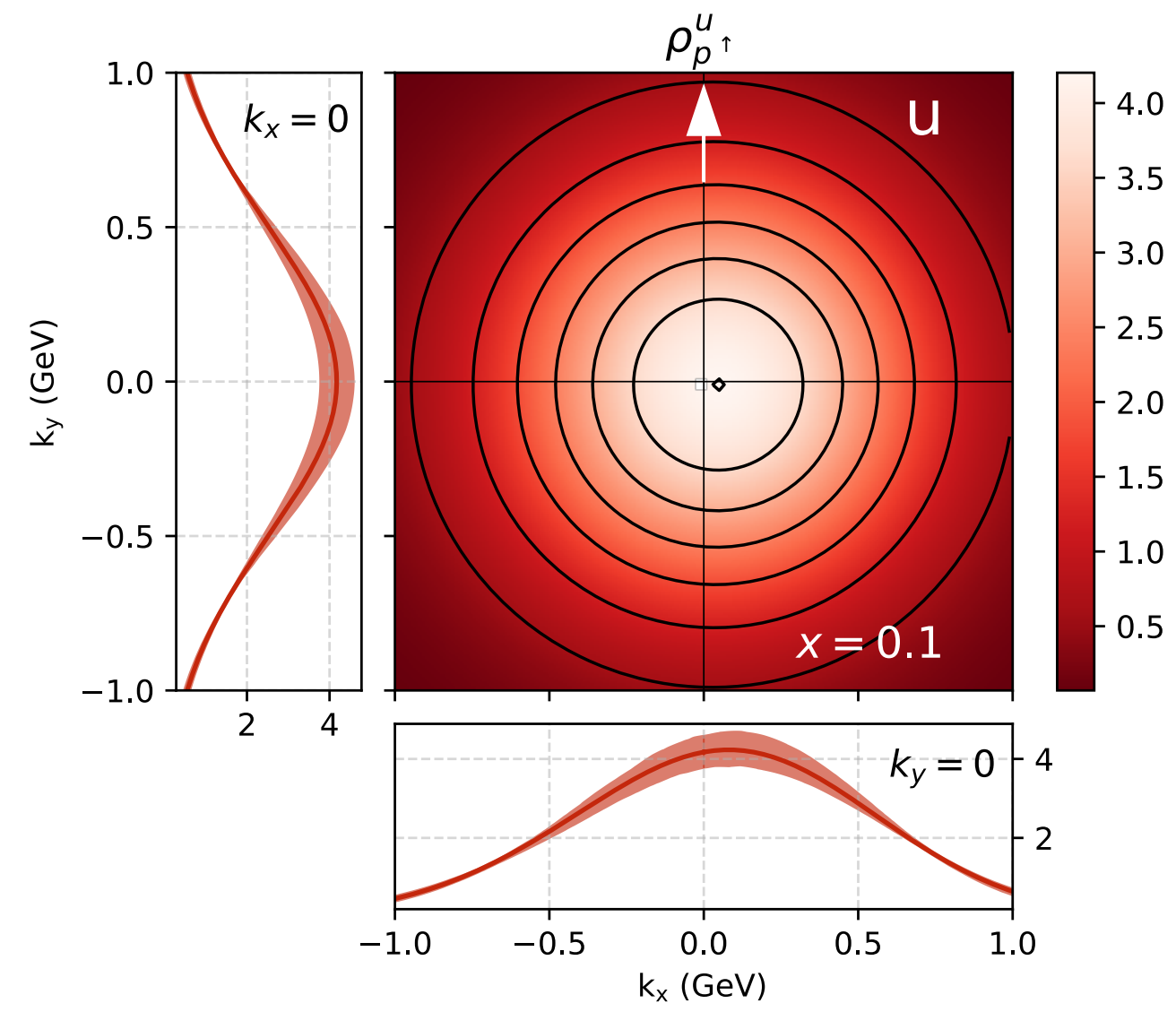
[Bacchetta, Delcarro, Pisano, Radici, arXiv:2004.14278](#)



[Echevarria, Kang, Terry, arXiv:2009.10710](#)

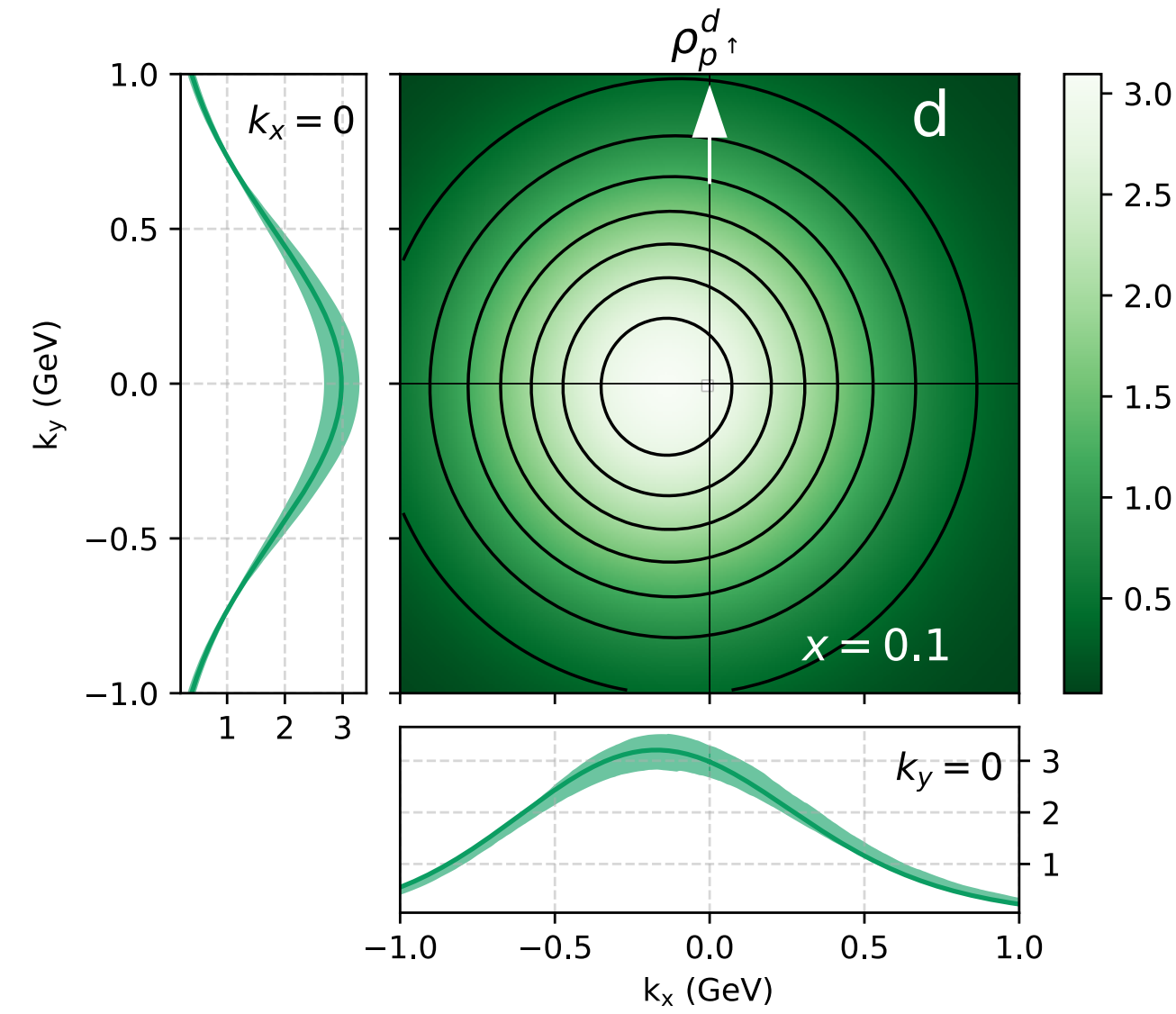
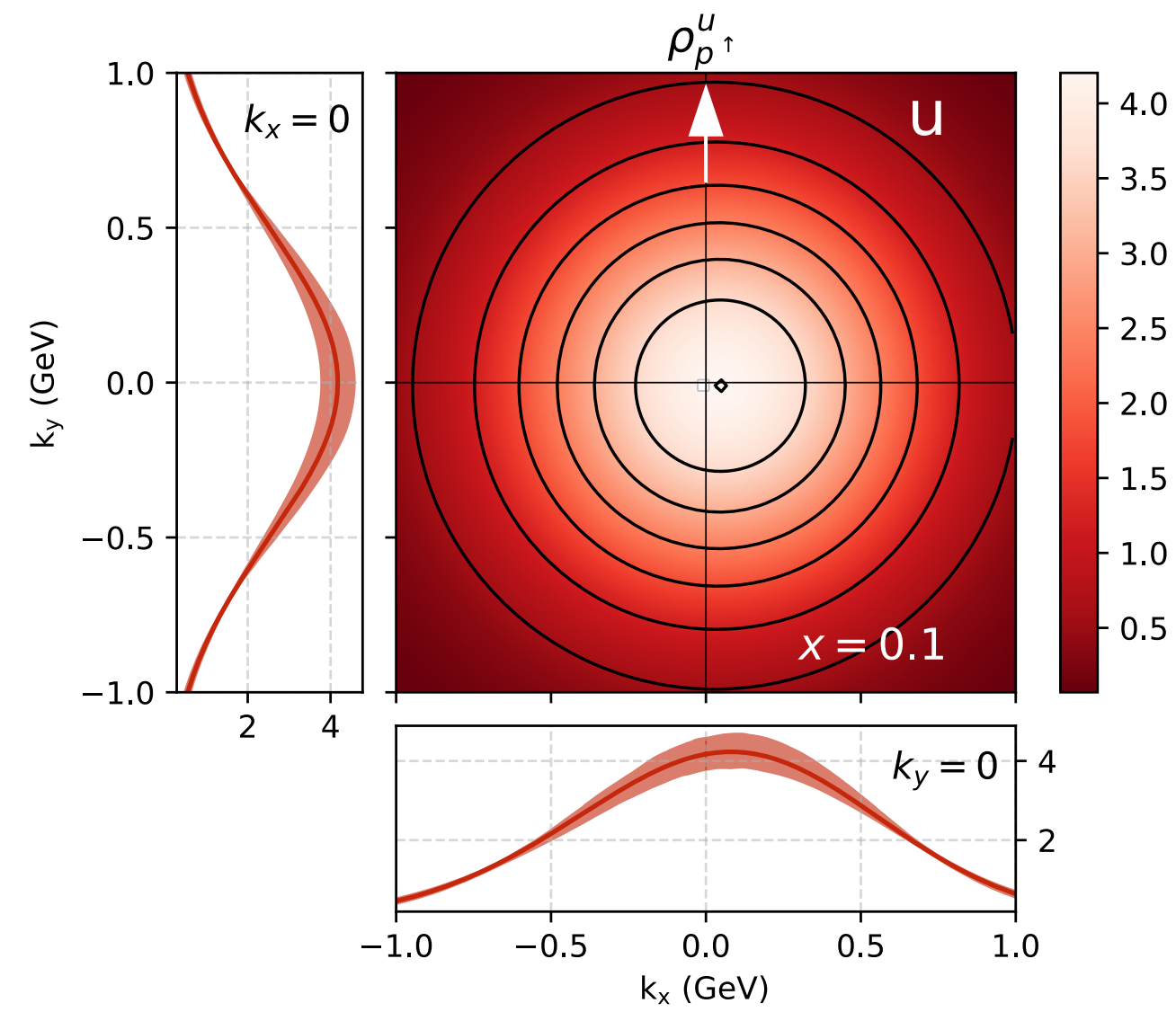


[Bury, Prokudin, Vladimirov, arXiv:2103.03270](#)



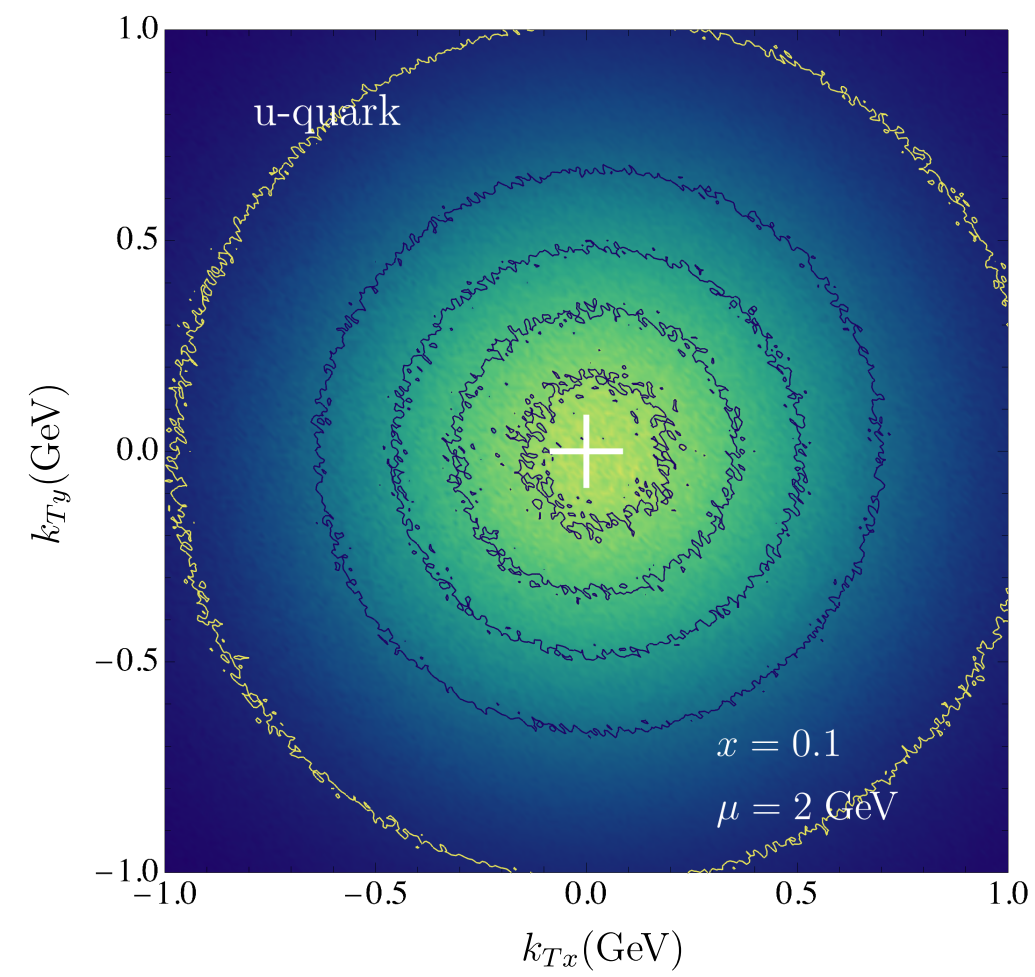
$Q = 2\text{GeV}$

[Bacchetta, Delcarro,
Pisano, Radici,
arXiv:2004.14278](#)

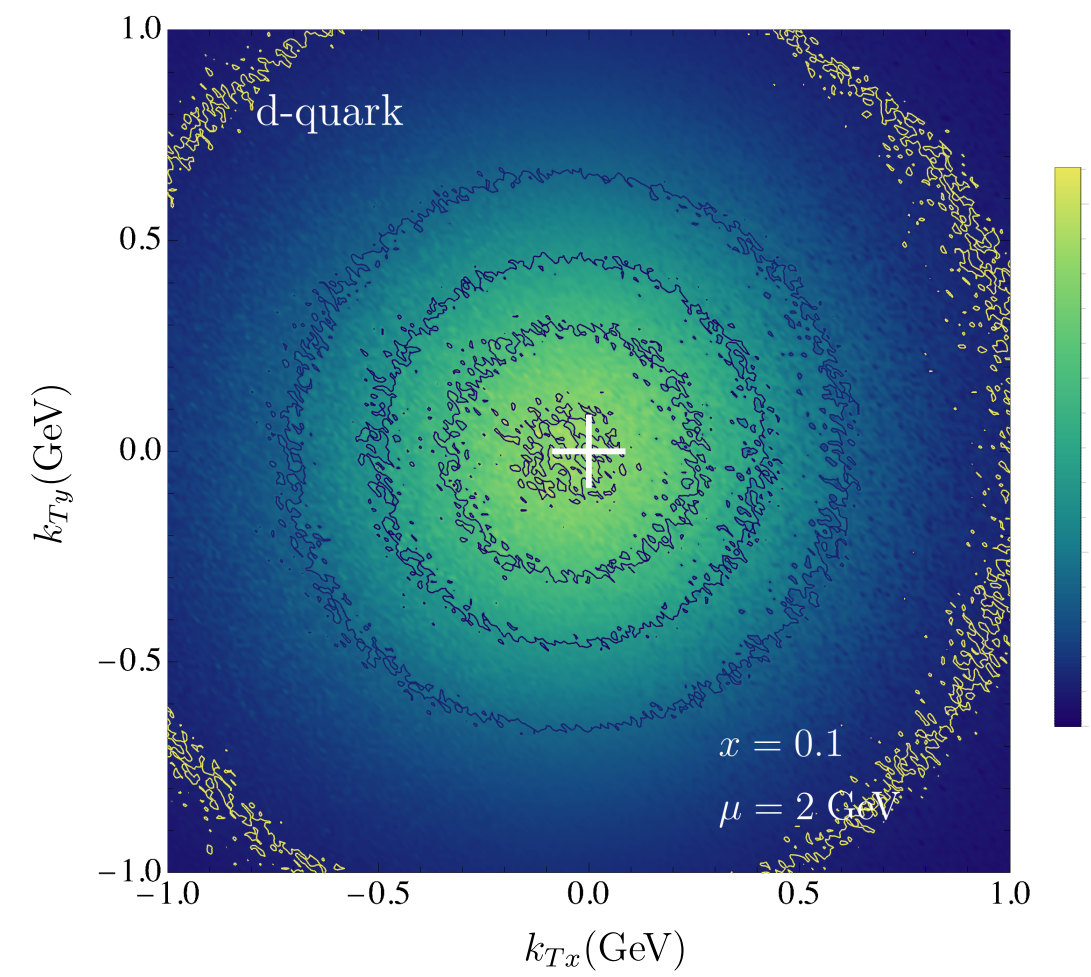


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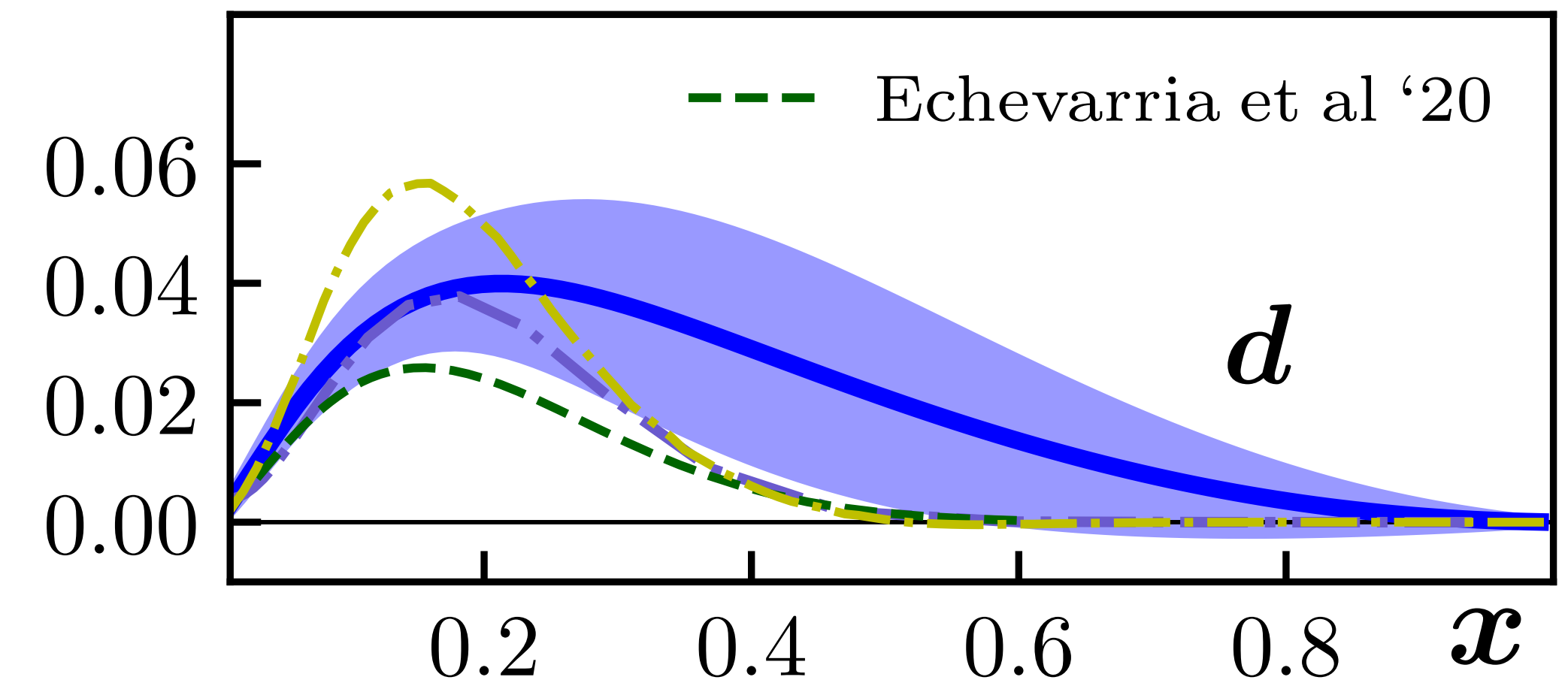
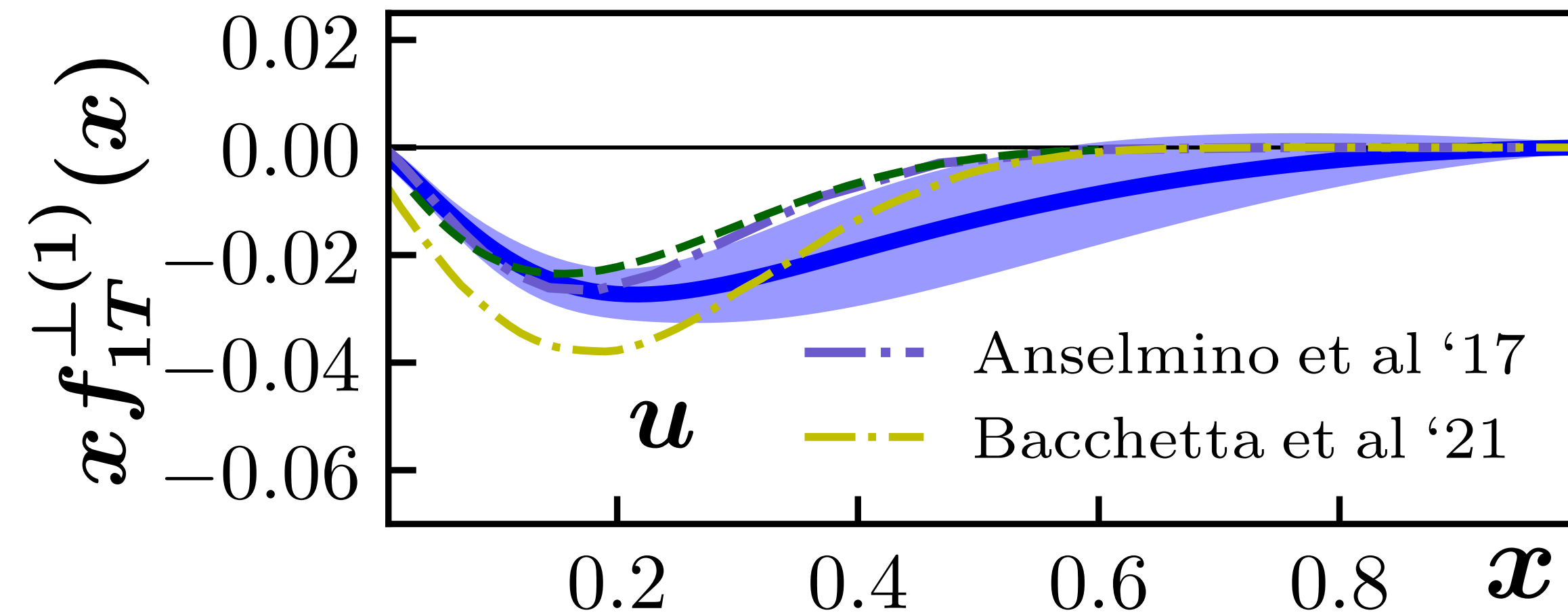


(a)

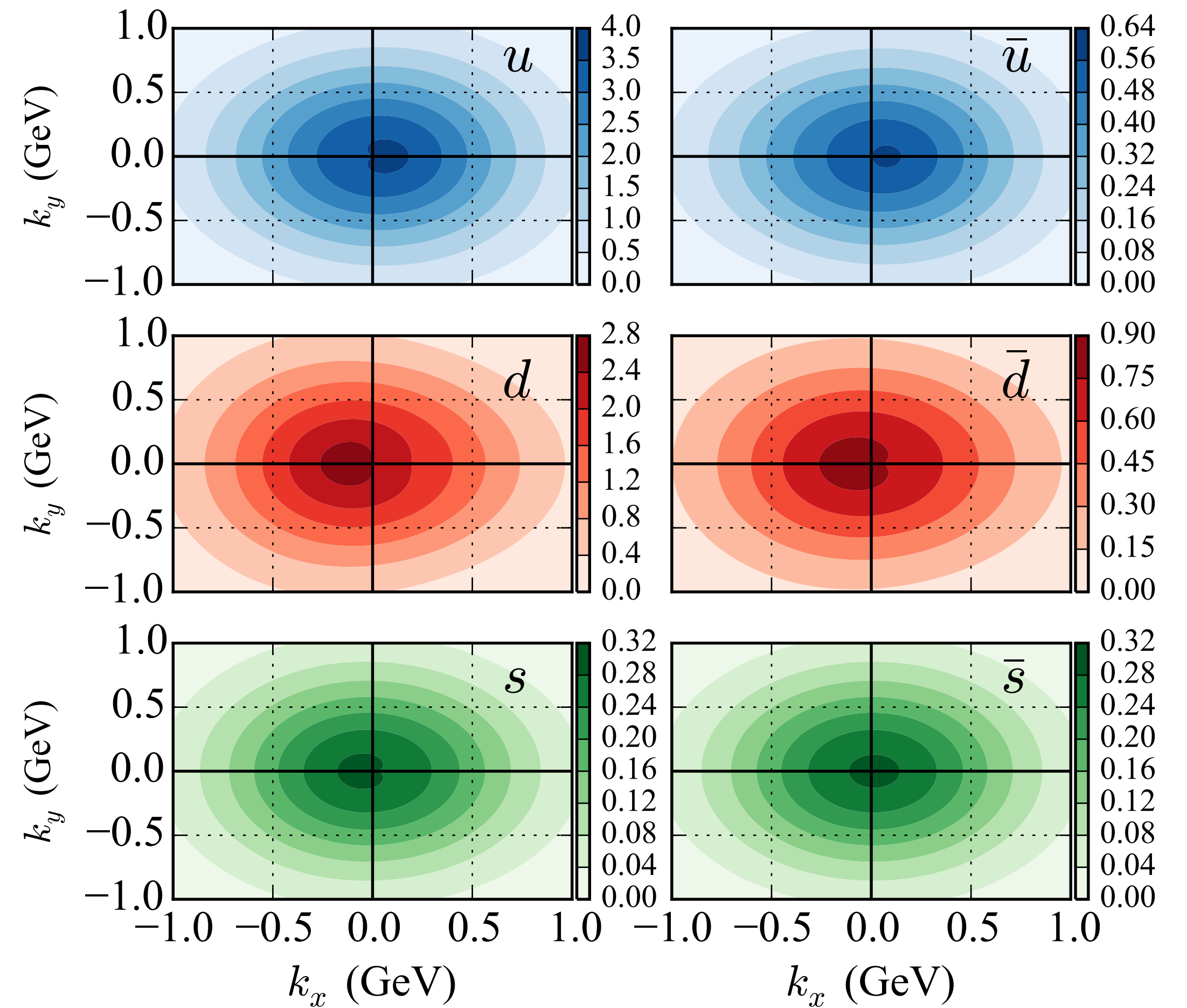
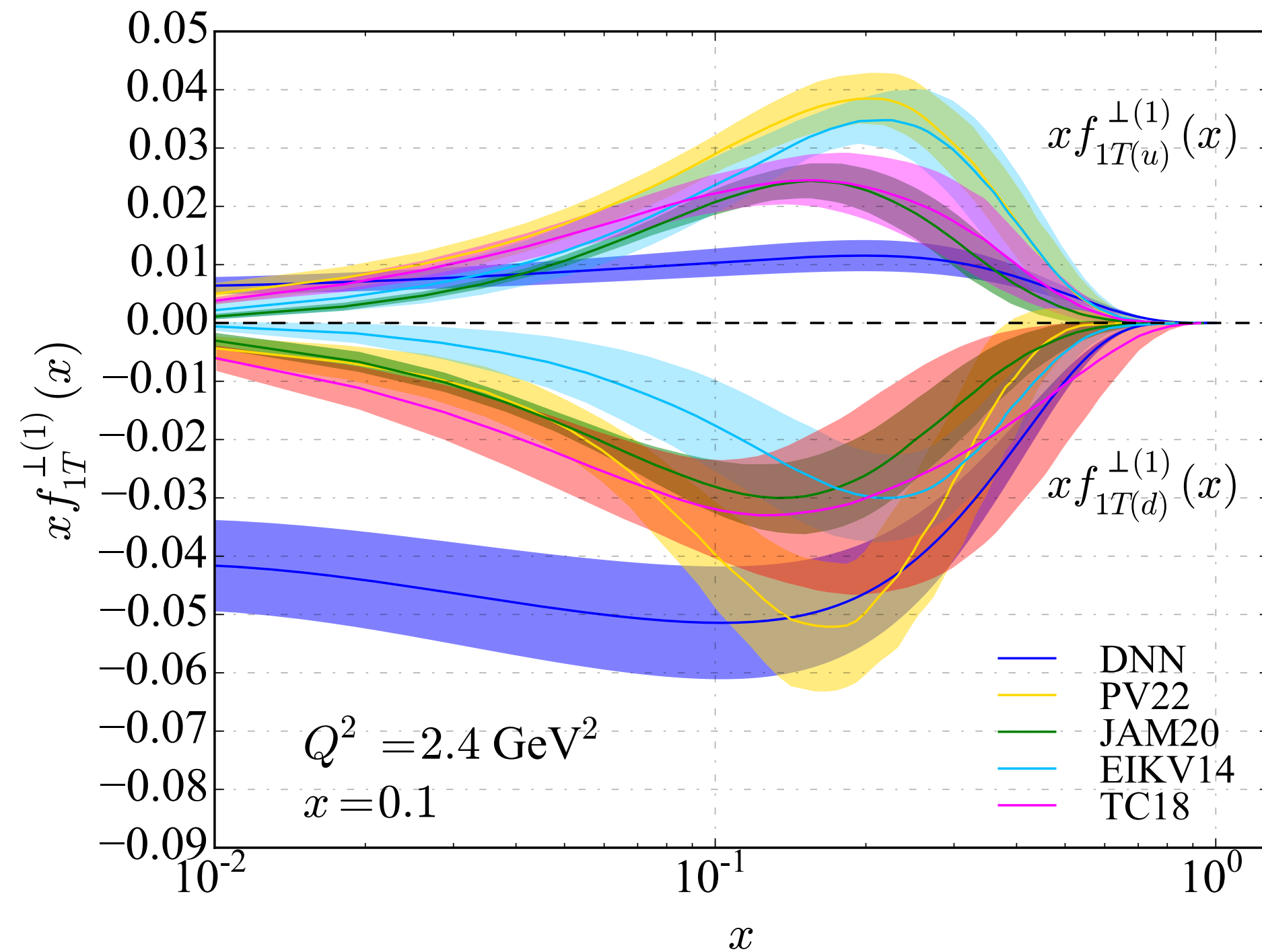


(b)

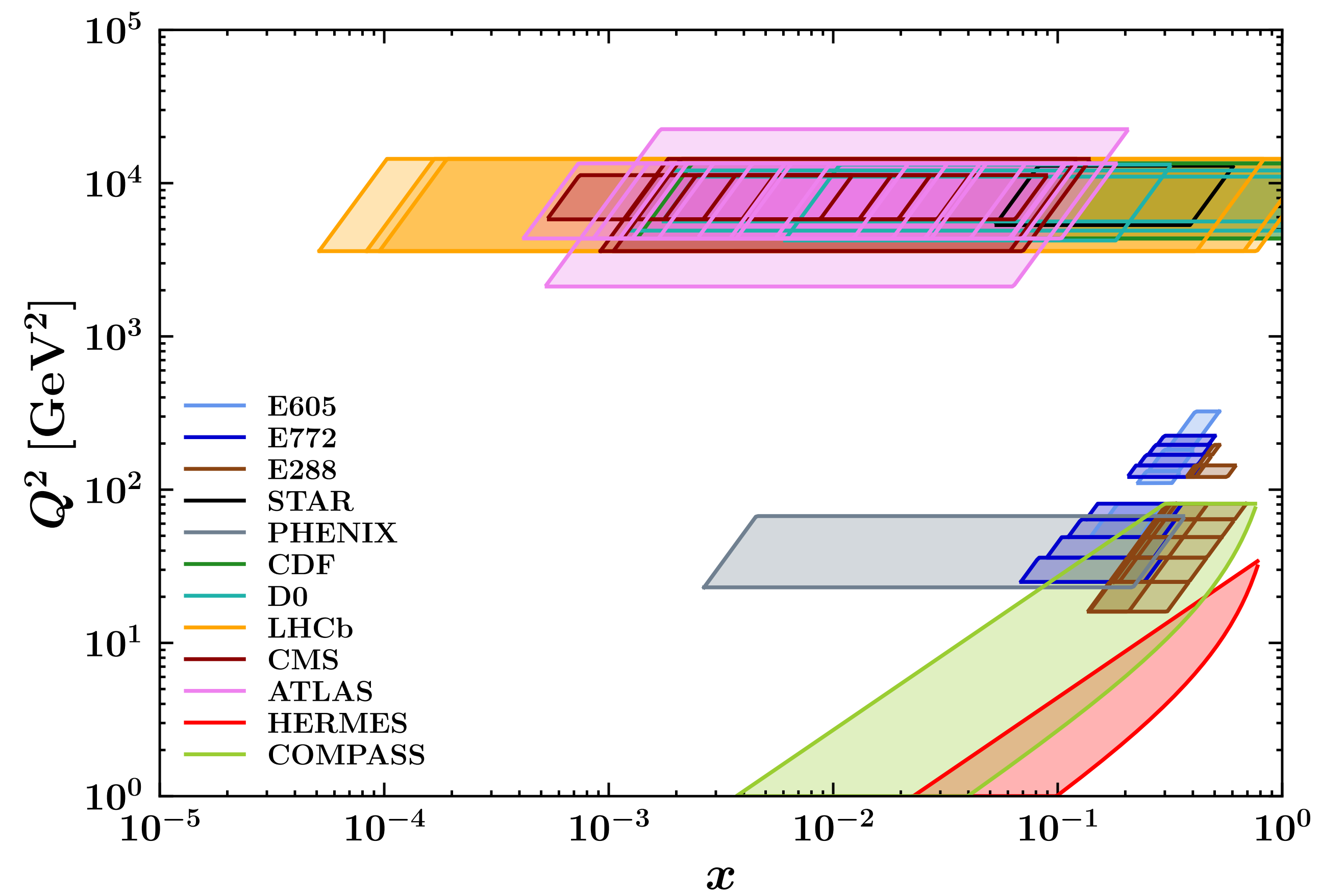
[Bury, Prokudin, Vladimirov, arXiv:2103.03270](#)

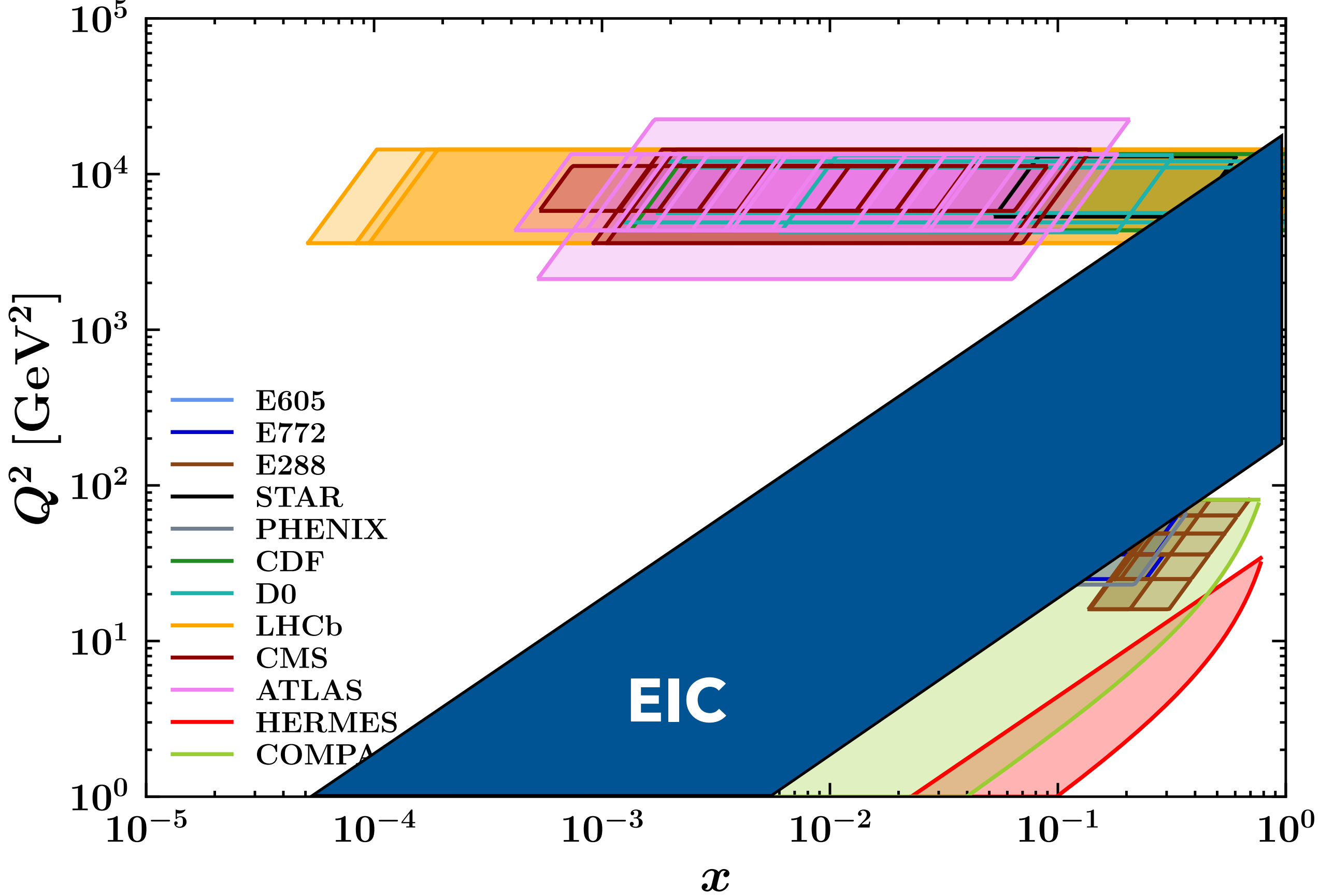


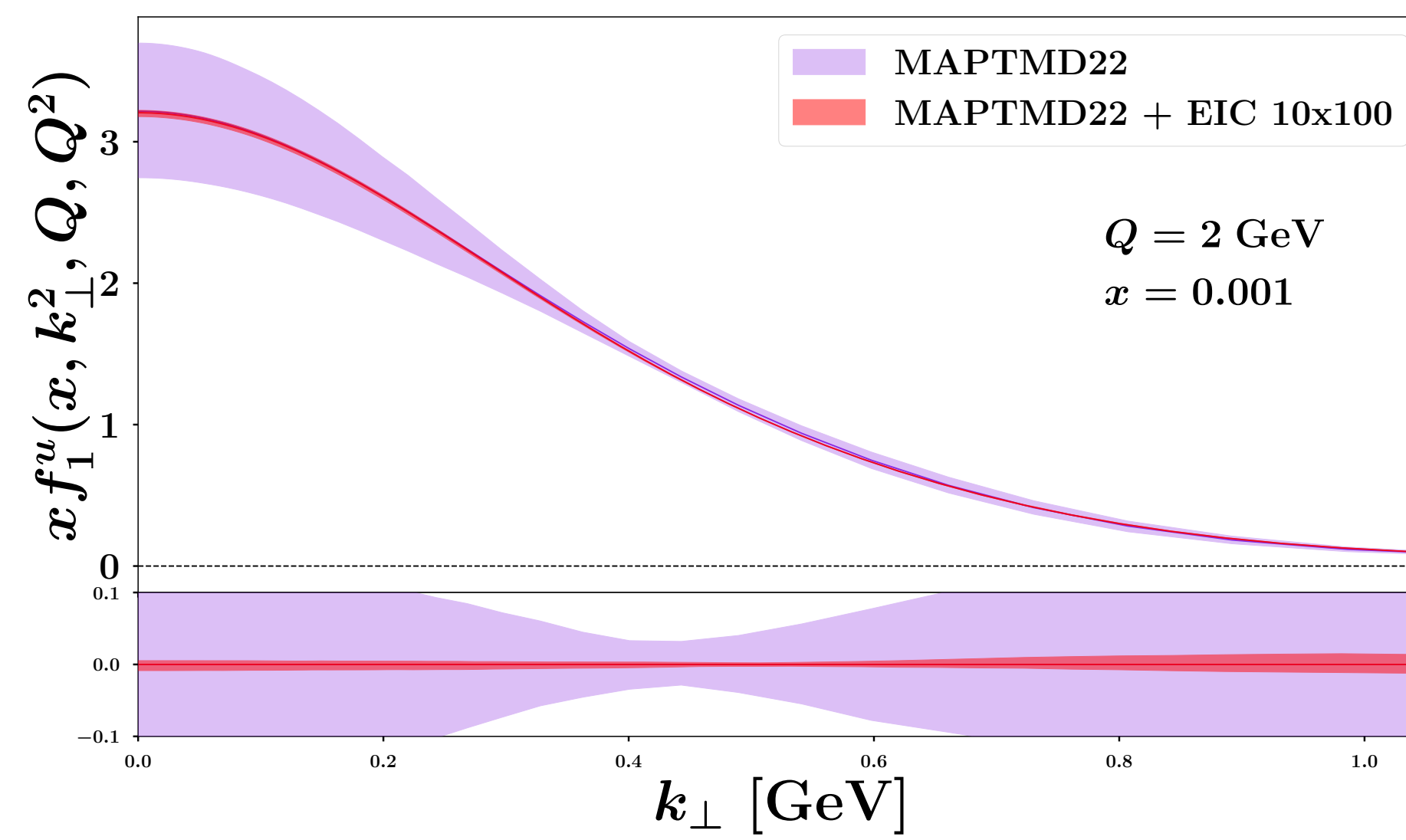
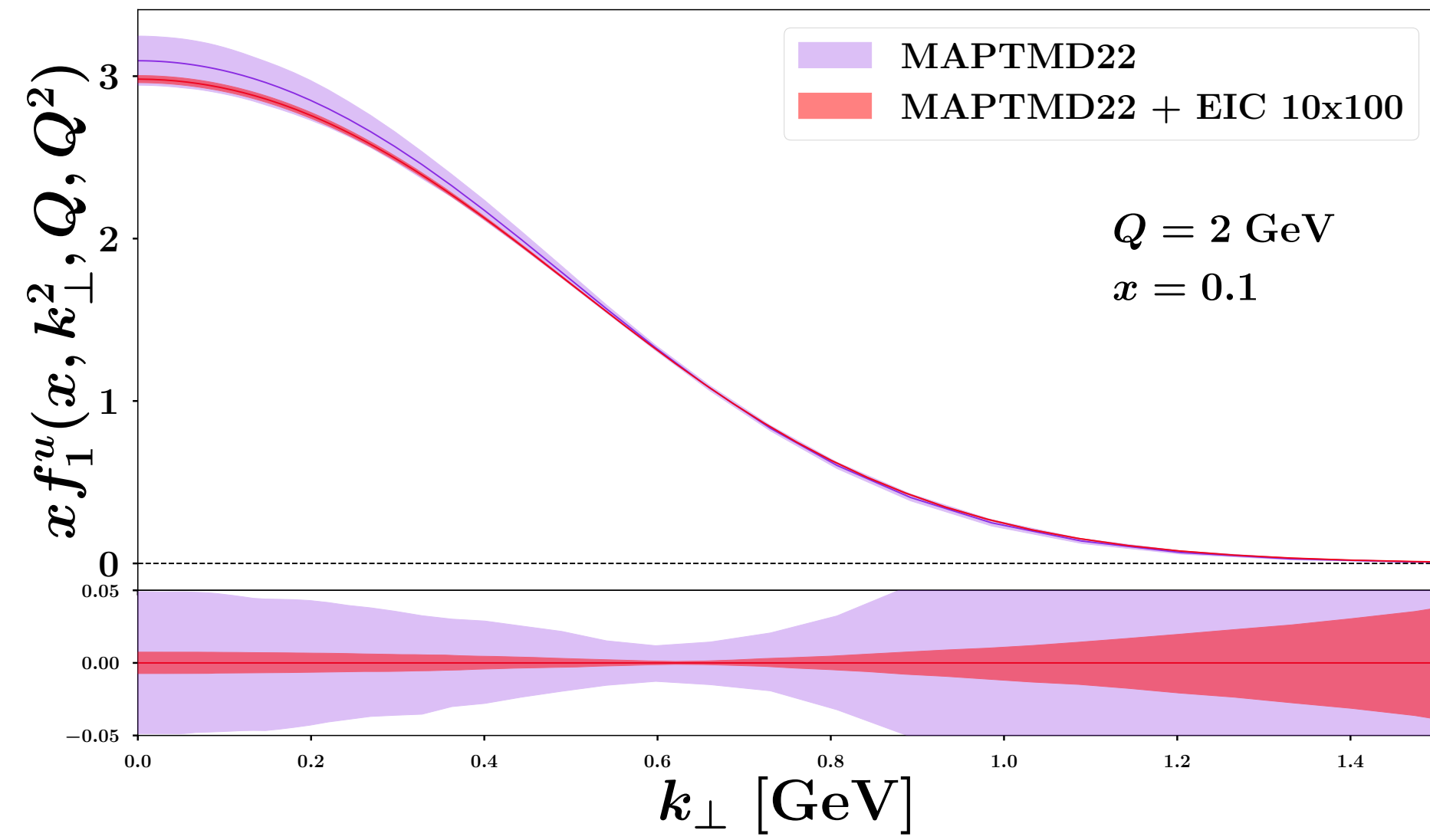
Interesting work from the point of view of simultaneous use of several measurements, but still limited from other perspectives (lack of TMD evolution and knowledge of the unpolarized function)



Interesting work from the point of view of the use of Neural Networks, but still limited from other perspectives (lack of TMD evolution and knowledge of the unpolarized function)







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- ▶ Progress is ongoing concerning higher-twist and gluon TMDs (see Cristian Pisano's talk)
- ▶ Extractions of unpolarized TMDs are reaching a good level of sophistication, but there are still several open questions
- ▶ For other TMDs, the study has barely started

BACKUP

MAP22

$$f_{1NP}(x, b_T^2) \propto \text{F.T. of } \left(e^{-\frac{k_T^2}{g^1}} + \lambda^2 k_T^2 e^{-\frac{k_T^2}{g^1 B}} + \lambda_2^2 e^{-\frac{k_T^2}{g^1 C}} \right)$$

$$g_K(b_T^2) = -\frac{g_2^2}{2} b_T^2$$

$$g_1(x) = N_1 \frac{(1-x)^\alpha x^\sigma}{(1-\hat{x})^\alpha \hat{x}^\sigma}$$

11 parameters for TMD PDF
+ 1 for NP evolution +9 for FF
= 21 free parameters

SV19

$$f_{NP}(x, b) = \exp \left(-\frac{\lambda_1(1-x) + \lambda_2 x + x(1-x)\lambda_5}{\sqrt{1 + \lambda_3 x^{\lambda_4} b^2}} b^2 \right)$$

$$g_K(b_T^2) = -2c_0 \frac{b_T^2}{1 + b_T^2/B_{NP}^2}$$

5 parameters for TMD PDF
+ 2 for NP evolution +5 for FF
= 11 free parameters

ART23

$$f_{NP}^f(x, b) = \frac{1}{\cosh \left(\left(\lambda_1^f(1-x) + \lambda_2^f x \right) b \right)}$$

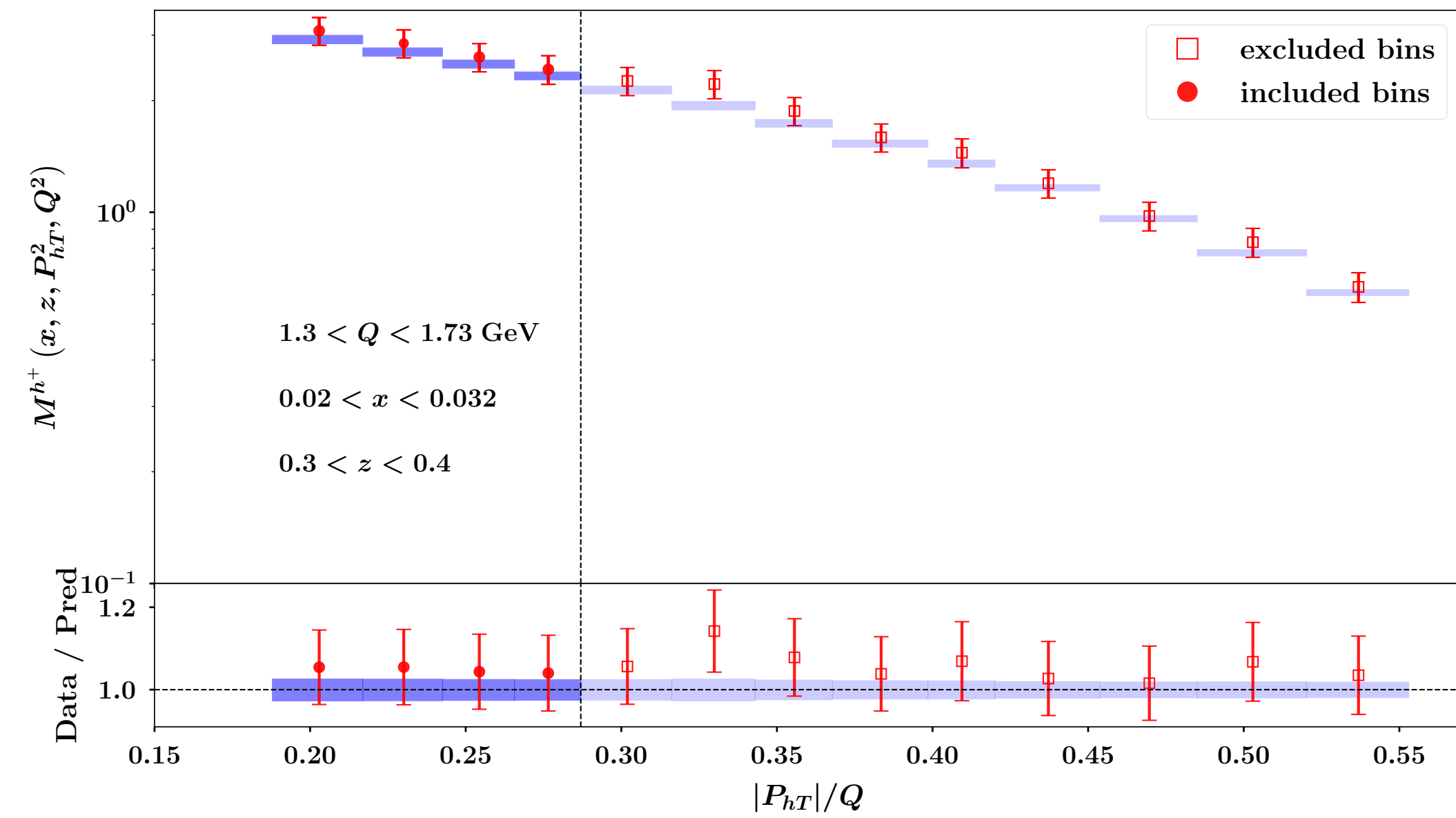
$$\mathcal{D}_{NP}(b) = bb^* \left[c_0 + c_1 \ln \left(\frac{b^*}{B_{NP}} \right) \right]$$

2401.14266

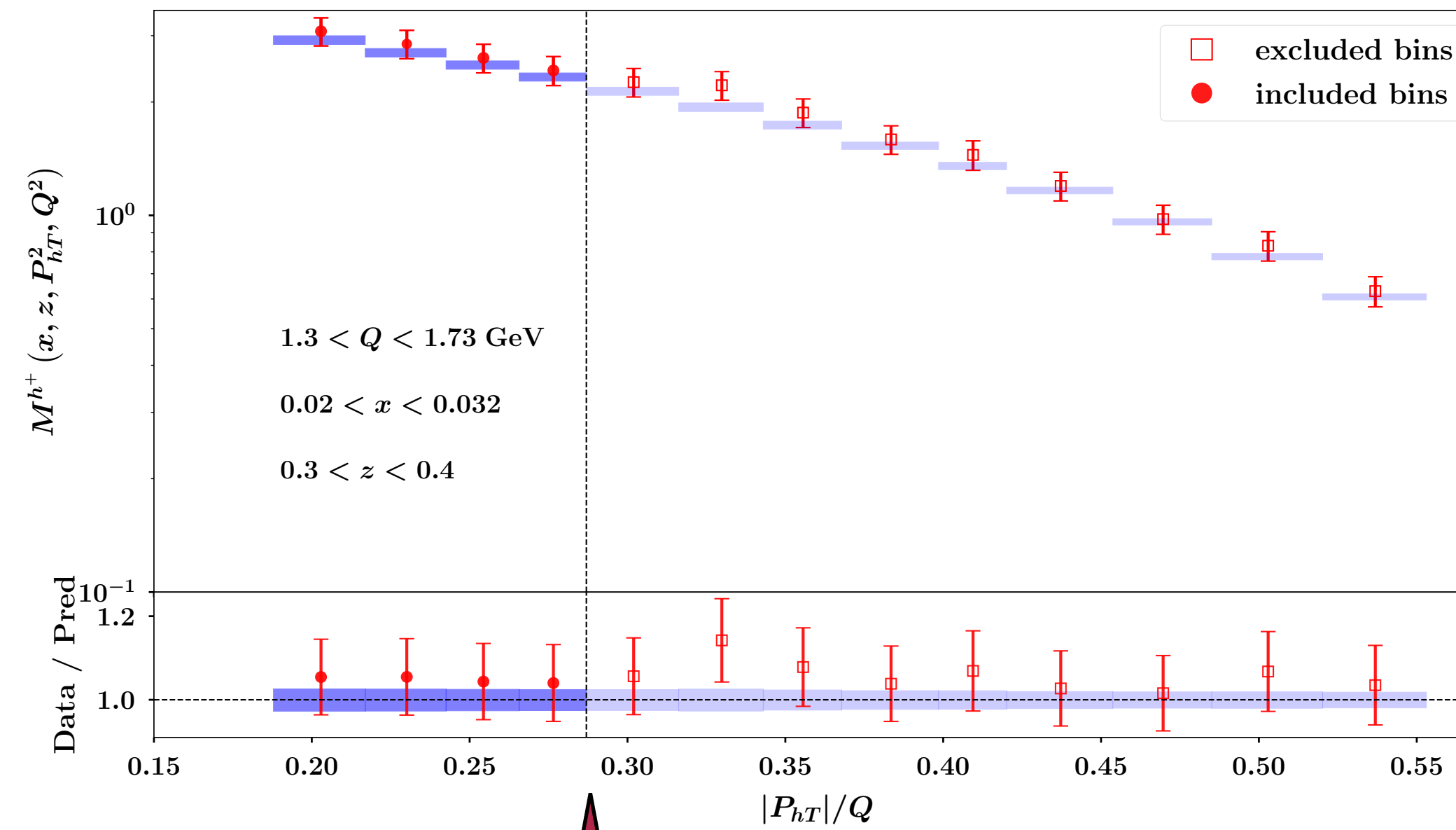
$$f_{\text{core},i/p}^{\text{Spect}}(x, \mathbf{k}_T; Q_0^2) = \frac{1}{\pi} \frac{6L^6}{L^2 + 2(m_q + x M_p)^2} \frac{k_T^2 + (m_q + x M_p)^2}{(k_T^2 + L^2)^4}$$

10 parameters for TMD PDF
+ 3 for NP evolution
= 13 free parameters

$$|q_T| = |P_{hT}|/z \ll Q$$

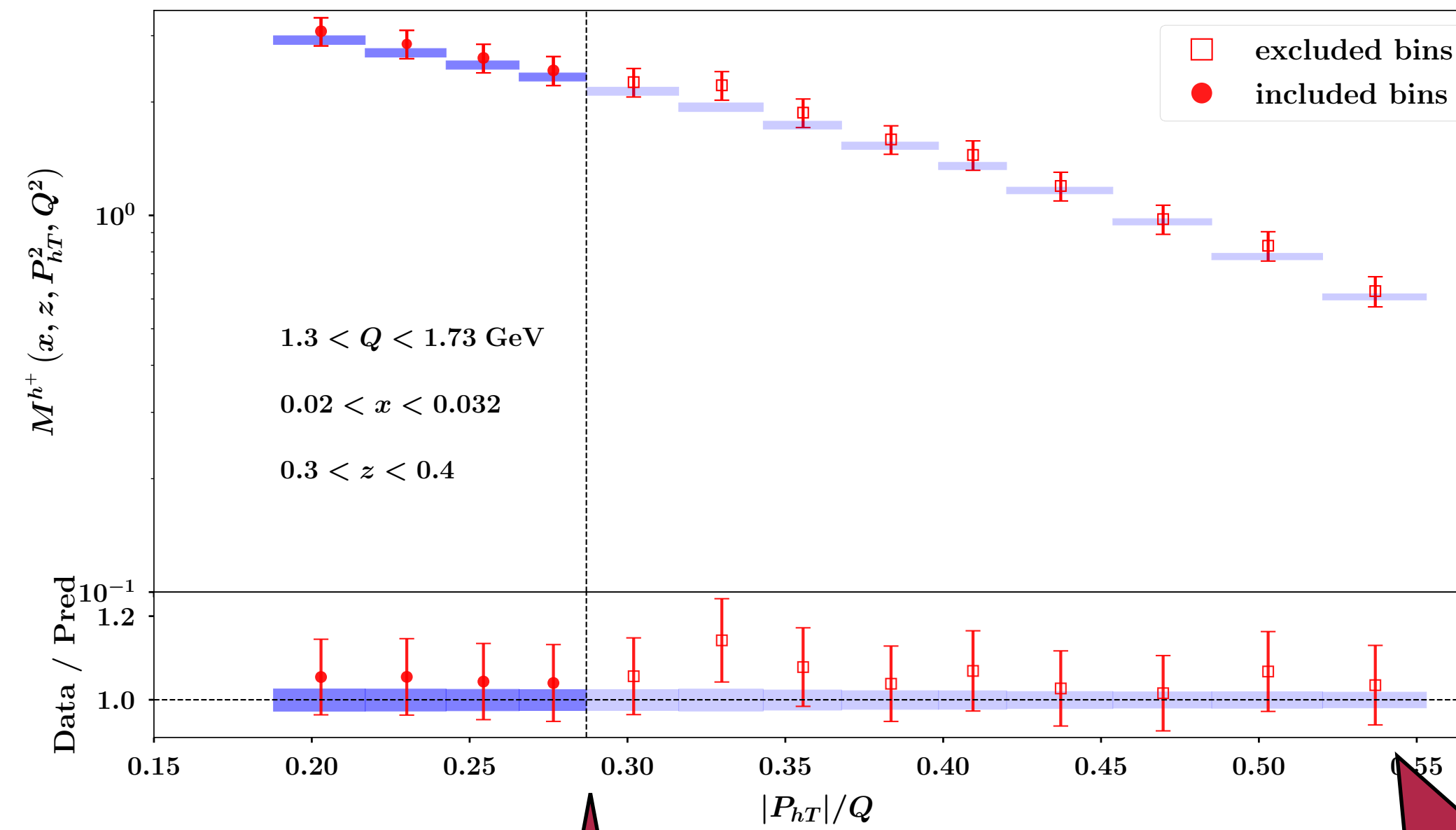


$$|q_T| = |P_{hT}|/z \ll Q$$



MAP22 cut

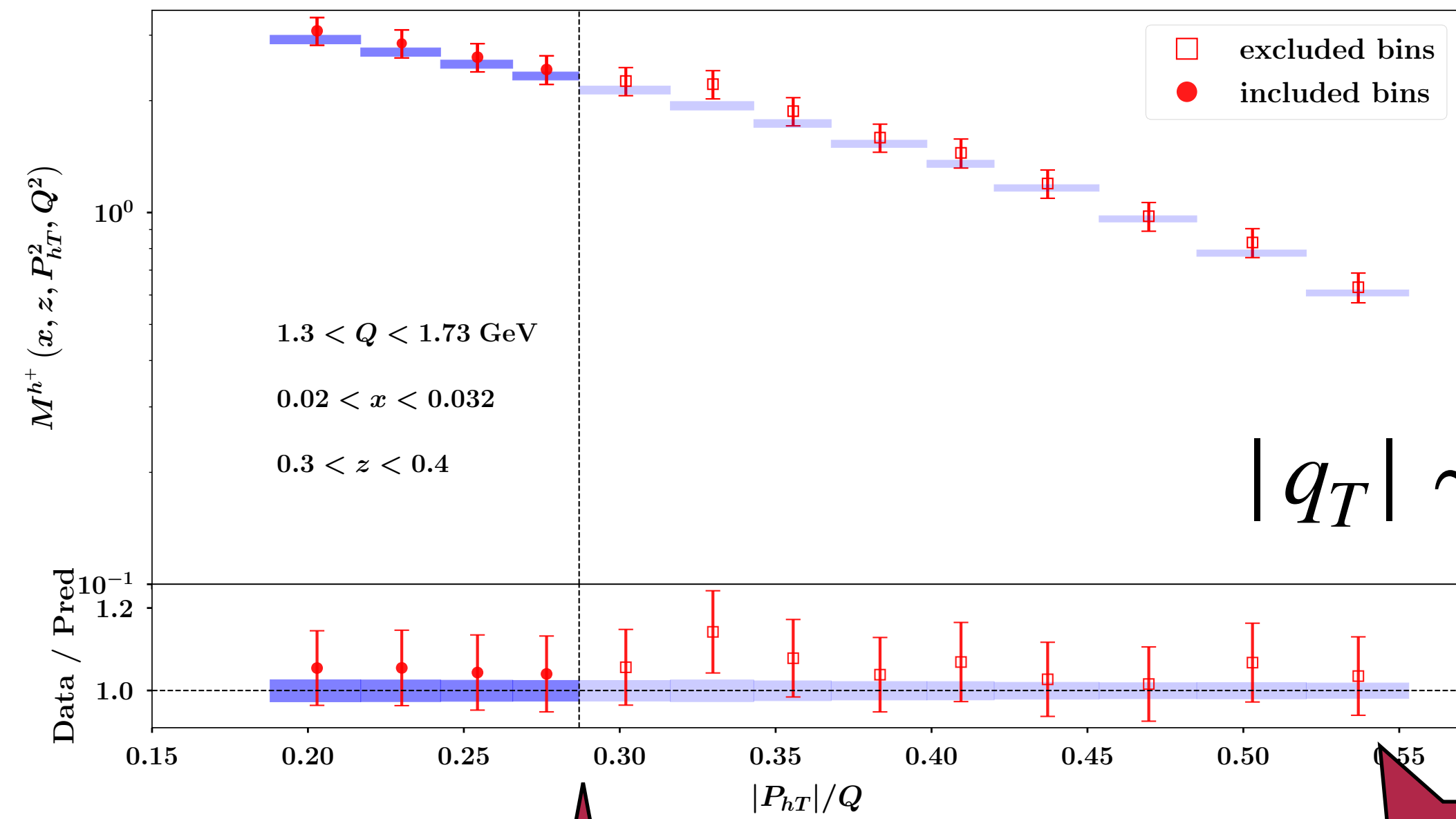
$$|q_T| = |P_{hT}|/z \ll Q$$



MAP22 cut

MAP22
extrapolation

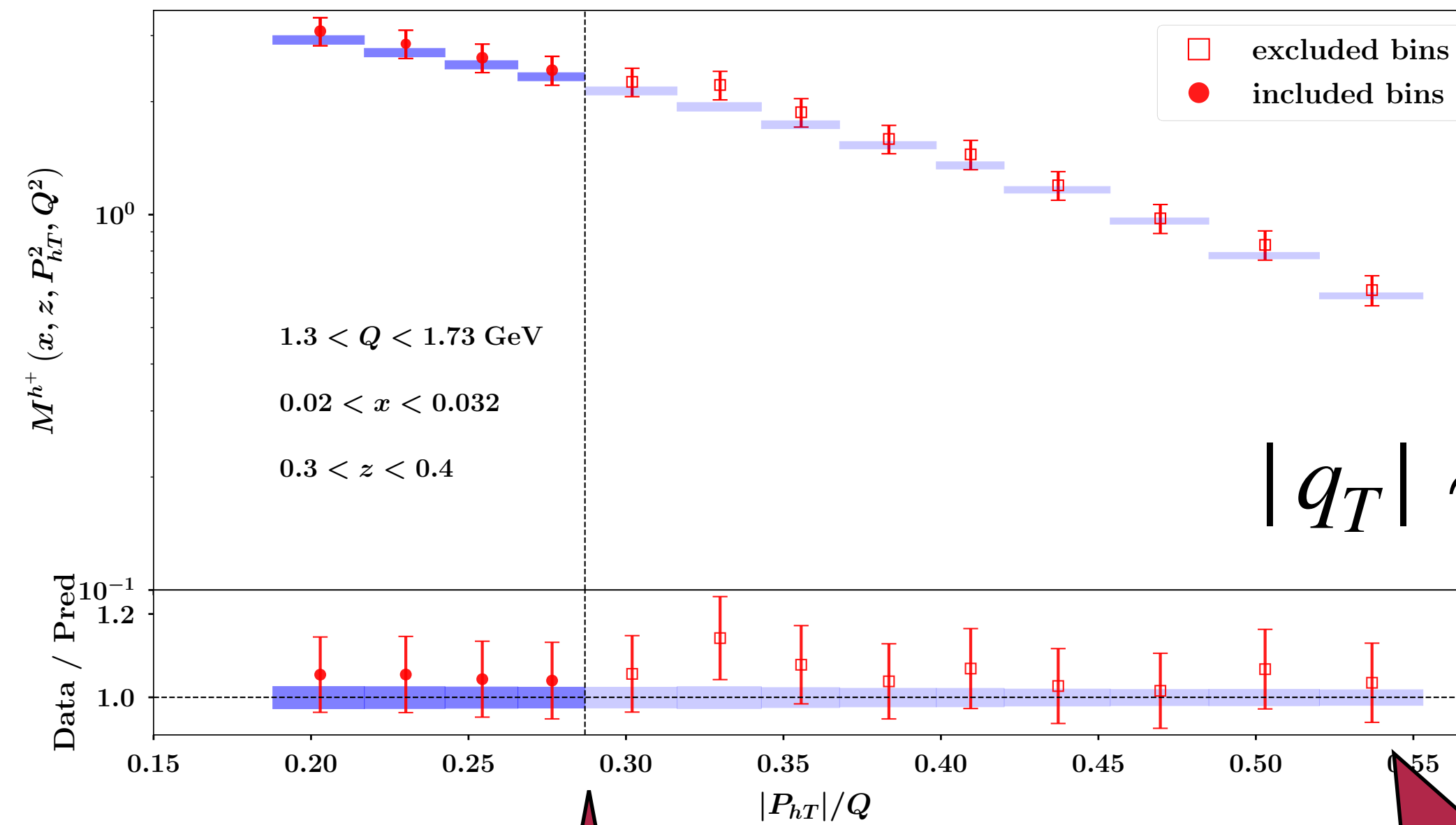
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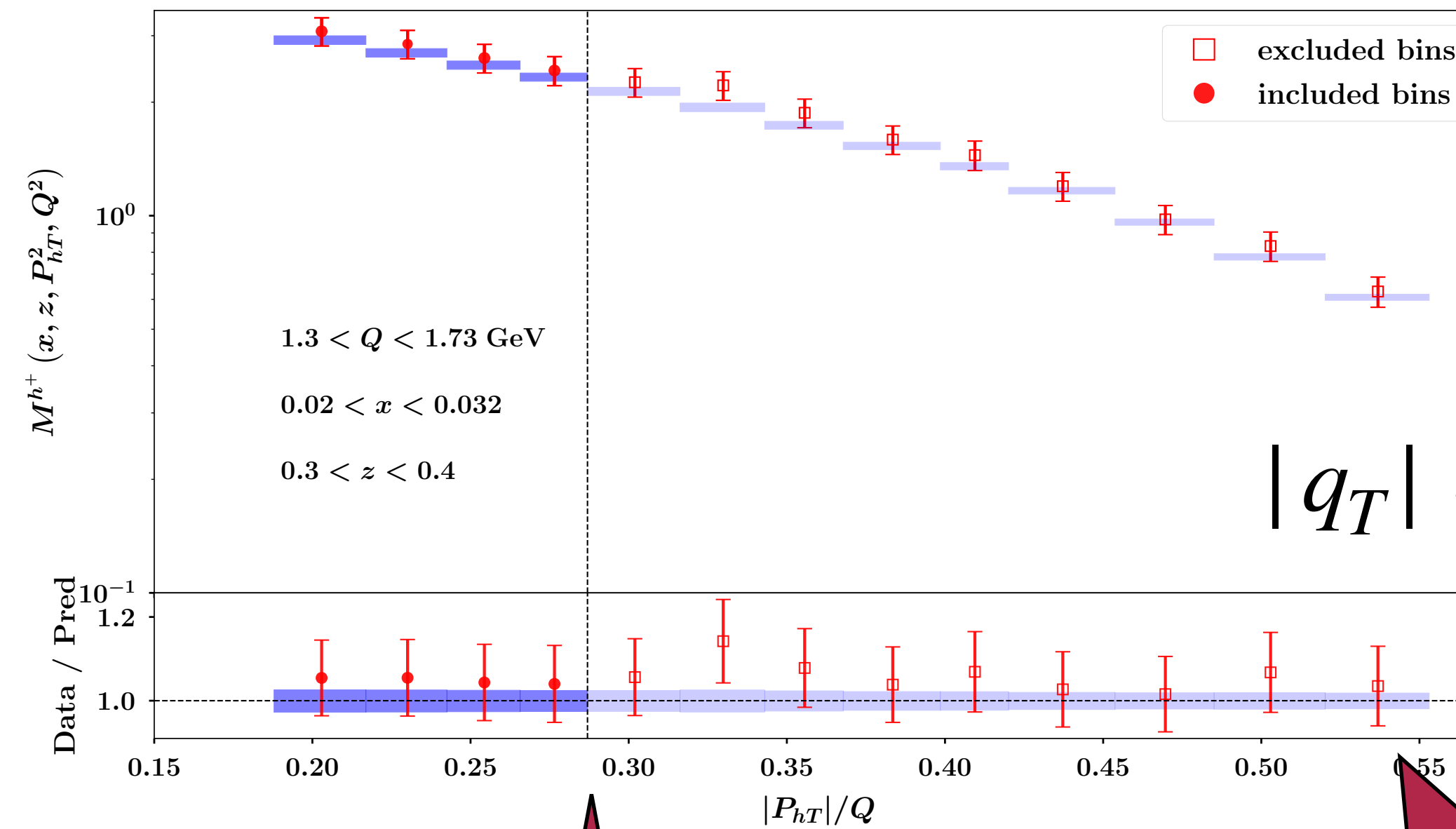


MAP22 cut

MAP22 extrapolation

The MAP22 cut is already considered to be "generous", but the physics seems to be the same for a much wider transverse momentum

$$|q_T| = |P_{hT}|/z \ll Q$$



MAP22 cut

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The MAP22 cut is already considered to be "generous", but the physics seems to be the same for a much wider transverse momentum

$$\log(Q^2 b_T^2) \rightarrow \log(Q^2 b_T^2 + 1)$$

*[see, e.g., Bozzi, Catani, De Florian, Grazzini
hep-ph/0302104](#)*

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[hep-ph/0302104](https://arxiv.org/abs/hep-ph/0302104)

$$b_*(b_c(b_T)) = \sqrt{\frac{b_T^2 + b_0^2/(C_5^2 Q^2)}{1 + b_T^2/b_{\max}^2 + b_0^2/(C_5^2 Q^2 b_{\max}^2)}} \quad b_{\min} \equiv b_*(b_c(0)) = \frac{b_0}{C_5 Q} \sqrt{\frac{1}{1 + b_0^2/(C_5^2 Q^2 b_{\max}^2)}}$$

Collins et al.
[arXiv:1605.00671](https://arxiv.org/abs/1605.00671)

$$b_* \equiv \frac{b_T}{\sqrt{1 + b_T^2/b_{\max}^2}}$$

Collins, Soper, Sterman, NPB250 (85)

$$\mu_0 = 1 \text{ GeV}$$

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$$\mu_b = 2e^{-\gamma_E}/b_*$$

$$\bar{b}_* \equiv b_{\text{max}} \left(\frac{1 - e^{-b_T^4/b_{\text{max}}^4}}{1 - e^{-b_T^4/b_{\text{min}}^4}} \right)^{1/4}$$

$$b_{\text{max}} = 2e^{-\gamma_E}$$

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These are all choices that should be at some point checked/challenged

$$\hat{f}_1^q(x, b_T; \mu^2) = \sum_i (C_{qi} \otimes f_1^i)(x, b_*; \mu_b) e^{\tilde{S}(b_*; \mu_b, \mu)} e^{g_K(b_T) \ln \frac{\mu}{\mu_0}} \hat{f}_{\text{NP}}^q(x, b_T)$$

$$\mu_0 = 1 \text{ GeV}$$

$$b_* \equiv \frac{b_T}{\sqrt{1 + b_T^2/b_{\text{max}}^2}} \quad \text{Collins, Soper, Sterman, NPB250 (85)}$$

$$\mu_b = 2e^{-\gamma_E}/b_* \quad \bar{b}_* \equiv b_{\text{max}} \left(\frac{1 - e^{-b_T^4/b_{\text{max}}^4}}{1 - e^{-b_T^4/b_{\text{min}}^4}} \right)^{1/4} \quad b_{\text{max}} = 2e^{-\gamma_E}$$

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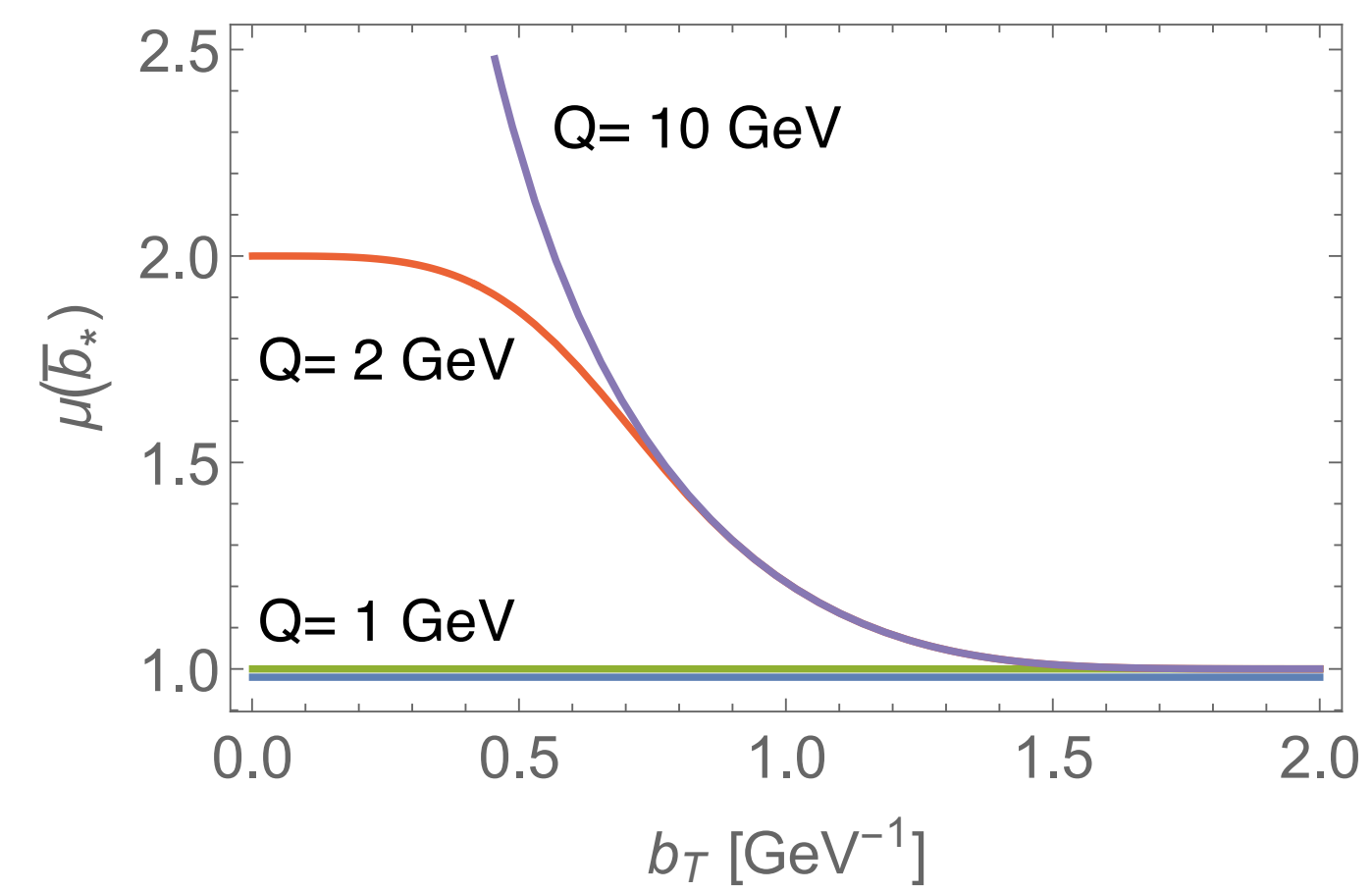
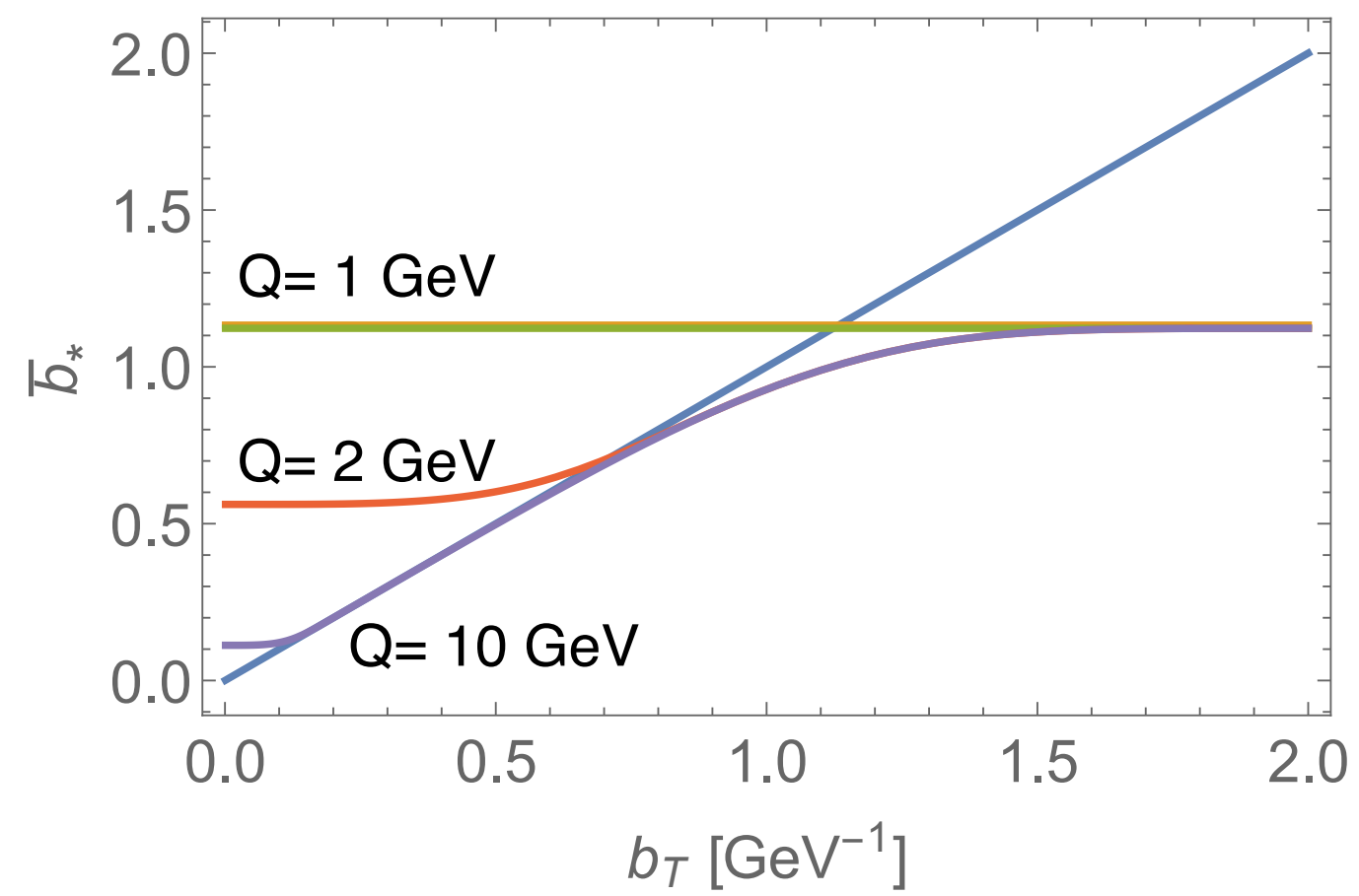
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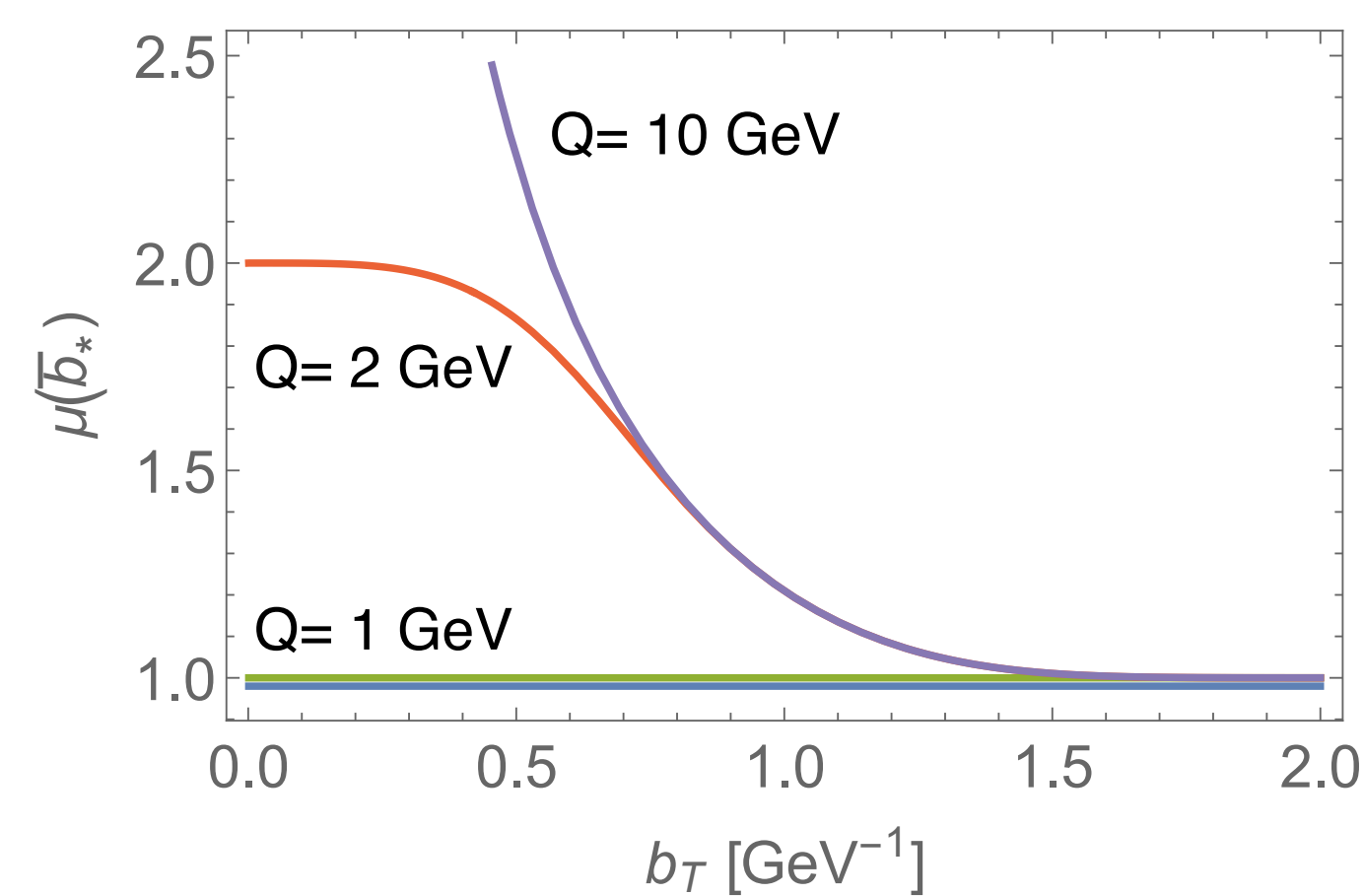
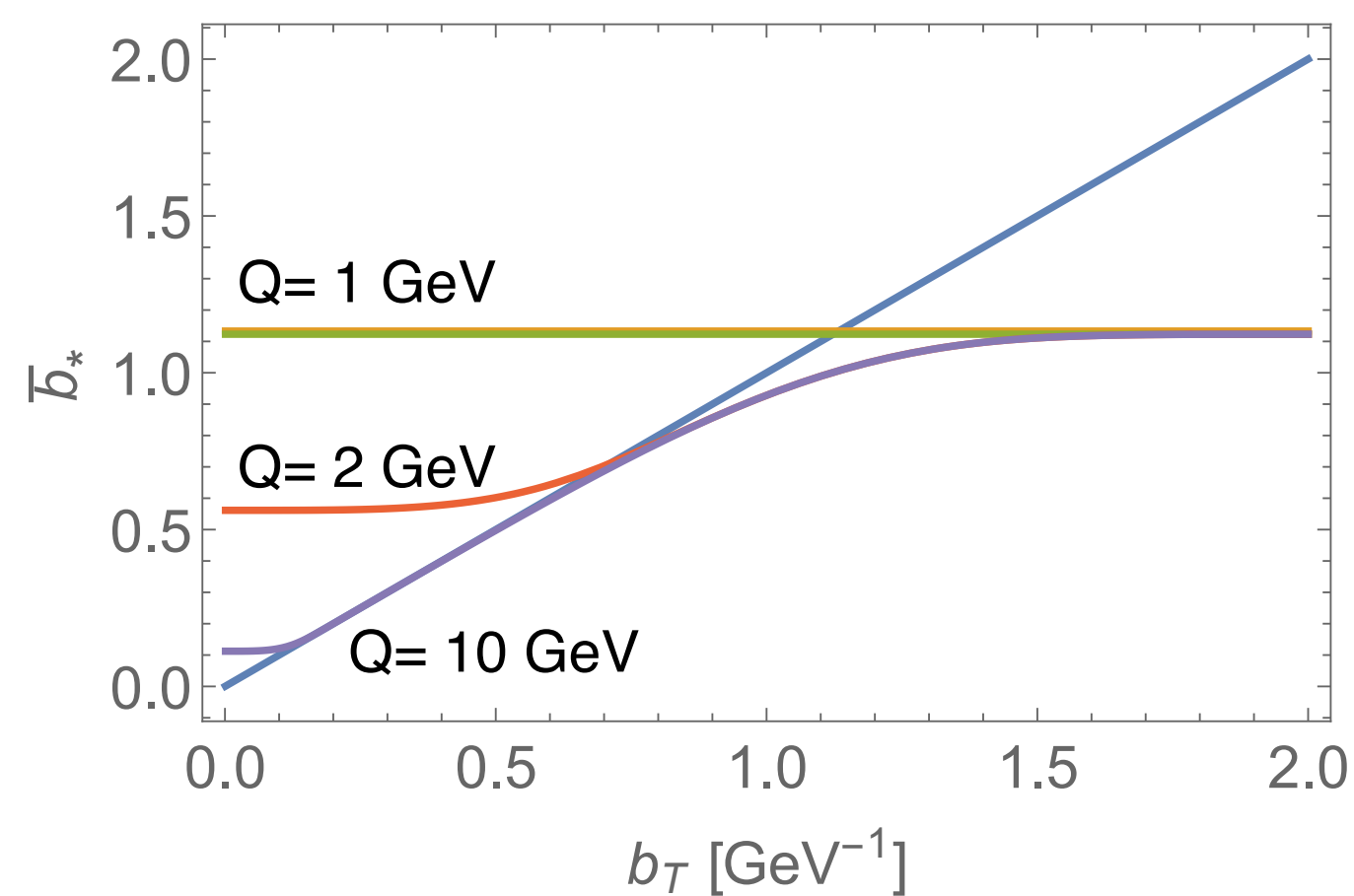


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No significant effect at high Q , but large effect at low Q
(inhibits perturbative contribution)