

Introduction to flavour experiments

In memoriam of Sheldon Stone (Feb. 14, 1946 – Oct. 6, 2021)



<https://cerncourier.com/a/sheldon-stone-1946-2021/>

ICTS 2022 Bengaluru (India), April 2022

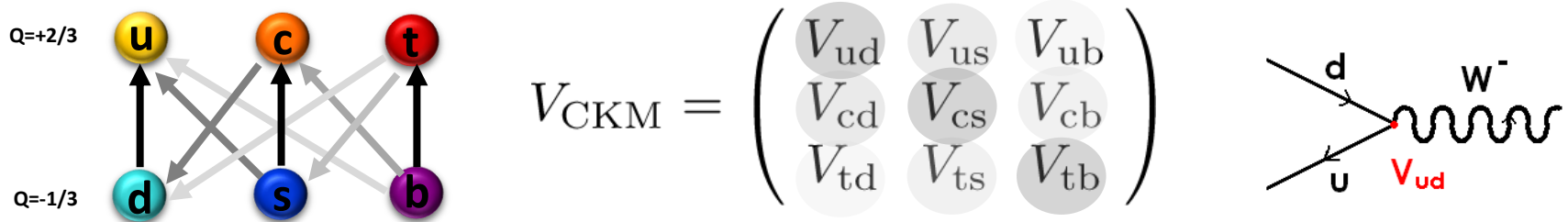
Arantza.Oyanguren@ific.uv.es

Outline

- Lesson 1: Introduction to flavour physics
- Lesson 2: The CKM matrix
- Lesson 3: Rare decays of heavy hadrons
- Lesson 4: Mixing and CP violation

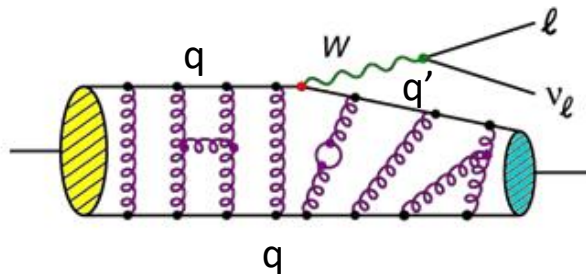
The CKM matrix

- In the Standard Model of Particle Physics, transitions between different quarks are governed by the CKM mechanism:



- The amplitude of a hadron decay process can be described using Effective Field Theories: Operator Product Expansion (OPE)

$$A(M \rightarrow F) = \langle F | \mathcal{H}_{eff} | M \rangle = \frac{G_F}{\sqrt{2}} \sum_i V_{CKM}^i C_i(\mu) \langle F | O_i(\mu) | M \rangle$$



CKM couplings	Wilson Coefficients ($\mu = \text{scale}$)	Hadronic Matrix Elements
------------------	--	-----------------------------

The CKM matrix

The CKM matrix V_{CKM} describes rotation for quarks between the weak eigenstates (d',s',b') and mass eigenstates (d,s,b)

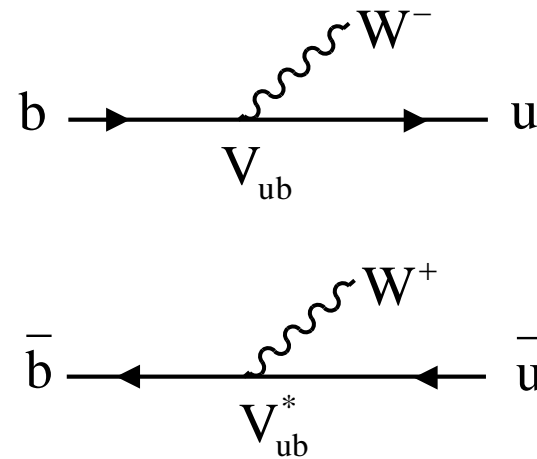
Quarks

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

Antiquarks

$$\begin{pmatrix} \bar{d}' \\ \bar{s}' \\ \bar{b}' \end{pmatrix} = \begin{pmatrix} V_{ud}^* & V_{us}^* & V_{ub}^* \\ V_{cd}^* & V_{cs}^* & V_{cb}^* \\ V_{td}^* & V_{ts}^* & V_{tb}^* \end{pmatrix} \begin{pmatrix} \bar{d} \\ \bar{s} \\ \bar{b} \end{pmatrix}$$

V_{ij} governs the transition from quark j to quark i



CP violation due to complex phases of CKM matrix elements

The CKM matrix

- The CKM matrix is complex and unitary $\hat{V}_{CKM}^+ \hat{V}_{CKM} = 1$
- 4 independent parameters
 - **Fundamental constants** of the Standard Model
 - Must be determined from **experiment**
- **Standard parametrization (PDG):**
 - 3 angles: θ_{12} , θ_{23} , θ_{13} and 1 phase δ

$$V_{CKM} = R_{23} \times R_{13} \times R_{12}$$

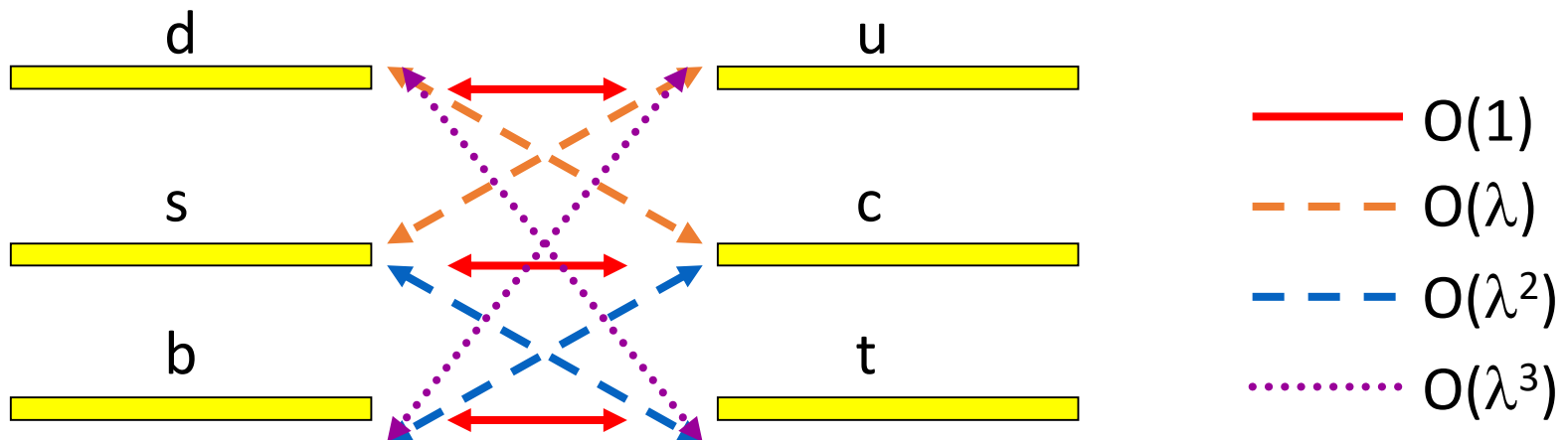
$$R_{12} = \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \quad R_{23} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \quad R_{13} = \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix}$$

$$s_{ij} = \sin \theta_{ij} \quad c_{ij} = \cos \theta_{ij}$$

The CKM matrix

- **Wolfenstein parameterization:** $s_{12} \sim 0.2, s_{23} \sim 0.04, s_{23} \sim 0.004$
- Perturbative, reflects the hierarchy of the matrix elements in terms of λ
 - $\lambda = \sin \theta_{12} \approx 0.23$ (Cabibbo angle)
- The four parameters are defined as:

$$\lambda = s_{12} \quad A = \frac{s_{23}}{s_{12}^2} \quad \rho = \frac{s_{13} \cos \delta}{s_{12} s_{23}} \quad \eta = \frac{s_{13} \sin \delta}{s_{12} s_{23}}$$



The CKM matrix

Wolfenstein parameterization at $O(\lambda^3)$:

$$V_{CKM} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} = \begin{pmatrix} 1 - \frac{\lambda^2}{2} & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \frac{\lambda^2}{2} & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix} + O(\lambda^4)$$

(next-to leading order corrections in λ may be important when increasing precision)

$$V_{CKM} = \begin{pmatrix} 1 - \frac{\lambda^2}{2} - \frac{\lambda^4}{8} & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda + A^2\lambda^5\left(\frac{1}{2} - \rho - i\eta\right) & 1 - \frac{\lambda^2}{2} - \frac{\lambda^4}{8}(1 + 4A^2) & A\lambda^2 \\ A\lambda^3(1 - \bar{\rho} - i\bar{\eta}) & -A\lambda^2 + A\lambda^4(1/2 - \rho - i\eta) & 1 - A^2\lambda^4/2 \end{pmatrix} + O(\lambda^6)$$

$$(\bar{\rho}, \bar{\eta}) \equiv (1 - \lambda^2/2)(\rho, \eta)$$

The CKM matrix

- CP Violation in the Standard Model:
 - Requirements for CP violation

$$\left(m_t^2 - m_c^2\right)\left(m_t^2 - m_u^2\right)\left(m_c^2 - m_u^2\right) \\ \times \left(m_b^2 - m_s^2\right)\left(m_b^2 - m_d^2\right)\left(m_s^2 - m_d^2\right) \times J_{CP} \neq 0$$

$$J_{CP} = \left| \text{Im} \left\{ V_{i\alpha} V_{j\beta} V_{i\beta}^* V_{j\alpha}^* \right\} \right| \quad (i \neq j, \alpha \neq \beta)$$

- Jarlskog invariant:

$$J_{CP} = s_{12}s_{13}s_{23}c_{12}c_{23}c_{13} \sin \delta = \lambda^6 A^2 \eta = O(10^{-5})$$

→ CP violation is small in the Standard Model

(and cannot explain the observed baryon asymmetry of the Universe)

The CKM matrix

- PDG 2021: [https://pdg.lbl.gov/2021/reviews/contents_sports.html]

$$0.97401 \pm 0.00011$$

superallowed $0^+ \rightarrow 0^+$ β decays

$$0.22650 \pm 0.00048$$

semileptonic / leptonic kaon decays
hadronic tau decays

$$(3.61 \pm 0.11) \times 10^{-3}$$

semileptonic / leptonic B decays

$$0.22636 \pm 0.00048$$

semileptonic charm decays
charm production in neutrino beams

$$0.97320 \pm 0.00011$$

semileptonic / leptonic charm decays

$$(40.53 \pm 0.83) \times 10^{-3}$$

semileptonic B decays

$$(8.54 \pm 0.23) \times 10^{-3}$$

B_d oscillations

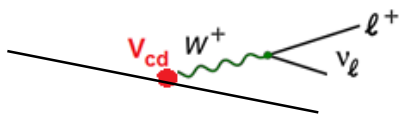
$$(39.78 \pm 0.82) \times 10^{-3}$$

B_s oscillations

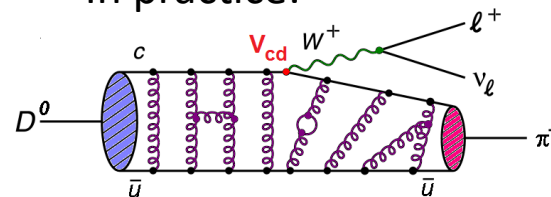
$$0.999172 \pm 0.000035$$

single top production

In theory:



In practice:



→ Need theory to describe QCD effects (lattice QCD)

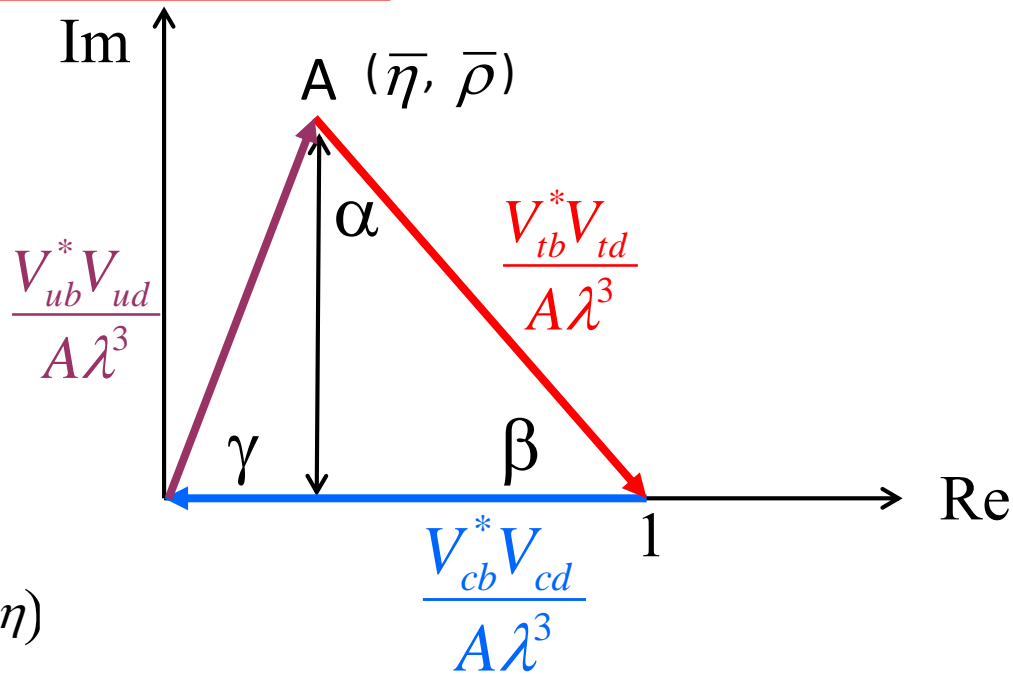
Q : How well is satisfied unitarity from the measured CKMs?

The CKM matrix

The Unitarity Triangle → CKM is unitary:

$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} V_{ud} & V_{cd} & V_{td} \\ V_{us} & V_{cs} & V_{ts} \\ V_{ub} & V_{cb} & V_{tb} \end{pmatrix}^* = \mathbf{I}$$

$$V_{ud} V_{ub}^* + V_{cd} V_{cb}^* + V_{td} V_{tb}^* = 0$$



$$(\bar{\rho}, \bar{\eta}) \equiv (1 - \lambda^2/2)(\rho, \eta)$$

$$\gamma \equiv \arg \left[-\frac{V_{ub}^* V_{ud}}{V_{cb}^* V_{cd}} \right] = \tan^{-1} \frac{\bar{\eta}}{\bar{\rho}} \sim 70^\circ$$

$$\beta \equiv \arg \left[-\frac{V_{cb}^* V_{cd}}{V_{tb}^* V_{td}} \right] = \tan^{-1} \frac{\bar{\eta}}{1 - \bar{\rho}} \sim 21^\circ$$

$$\alpha \equiv \pi - \beta - \gamma$$

The CKM matrix

The Unitarity Triangle

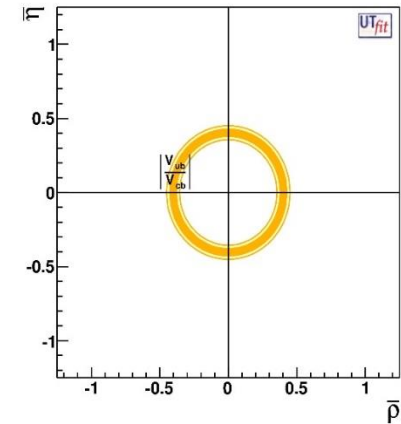
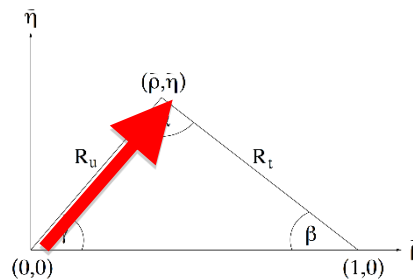
The idea: try to measure as many flavour observables as possible

overconstrain the unitarity triangle

Ex: Measuring the $b \rightarrow u \ell \nu$ vs the $b \rightarrow c \ell \nu$ transition

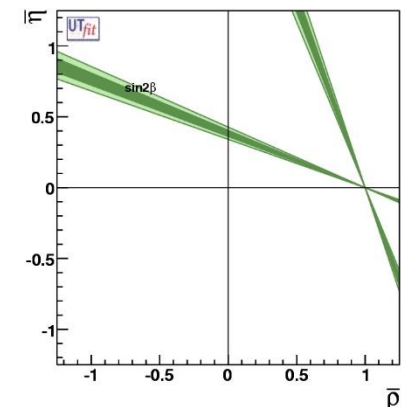
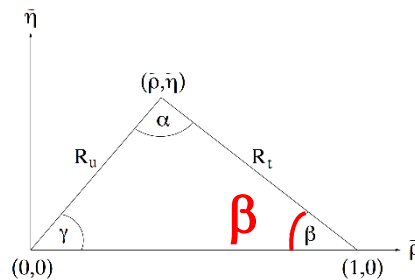
$$\frac{\Gamma(b \rightarrow u \ell \nu)}{\Gamma(b \rightarrow c \ell \nu)} \sim \left| \frac{V_{ub}}{V_{cb}} \right|^2 \sim \frac{1}{50}$$

$$\left| \frac{V_{ub}}{V_{cb}} \right| = \frac{\lambda}{1 - \lambda^2/2} \sqrt{\bar{\rho}^2 + \bar{\eta}^2}$$



Ex: Measuring time-dependent asymmetries in $b \rightarrow c\bar{c} s$ decays
(effect from interference of mixing and decay)

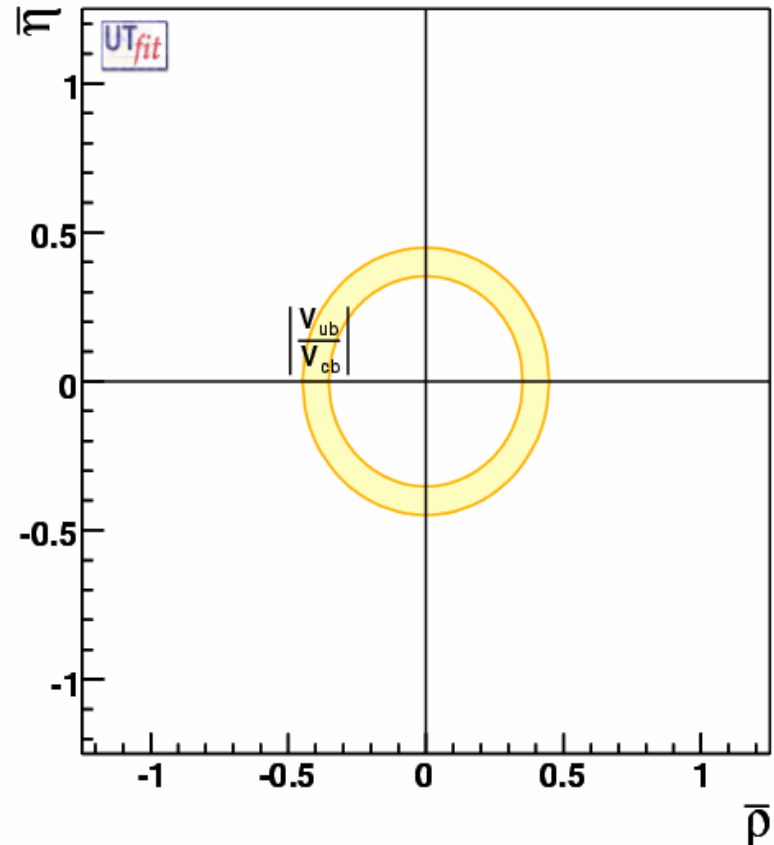
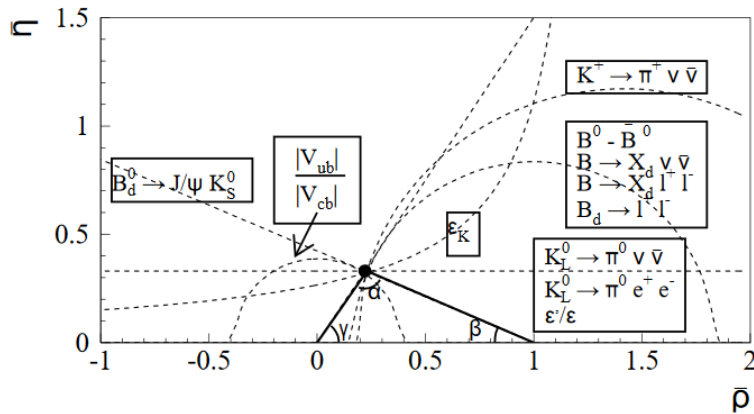
$$\beta = \tan^{-1} \frac{\bar{\eta}}{1 - \bar{\rho}}$$



The CKM matrix

The Unitarity Triangle

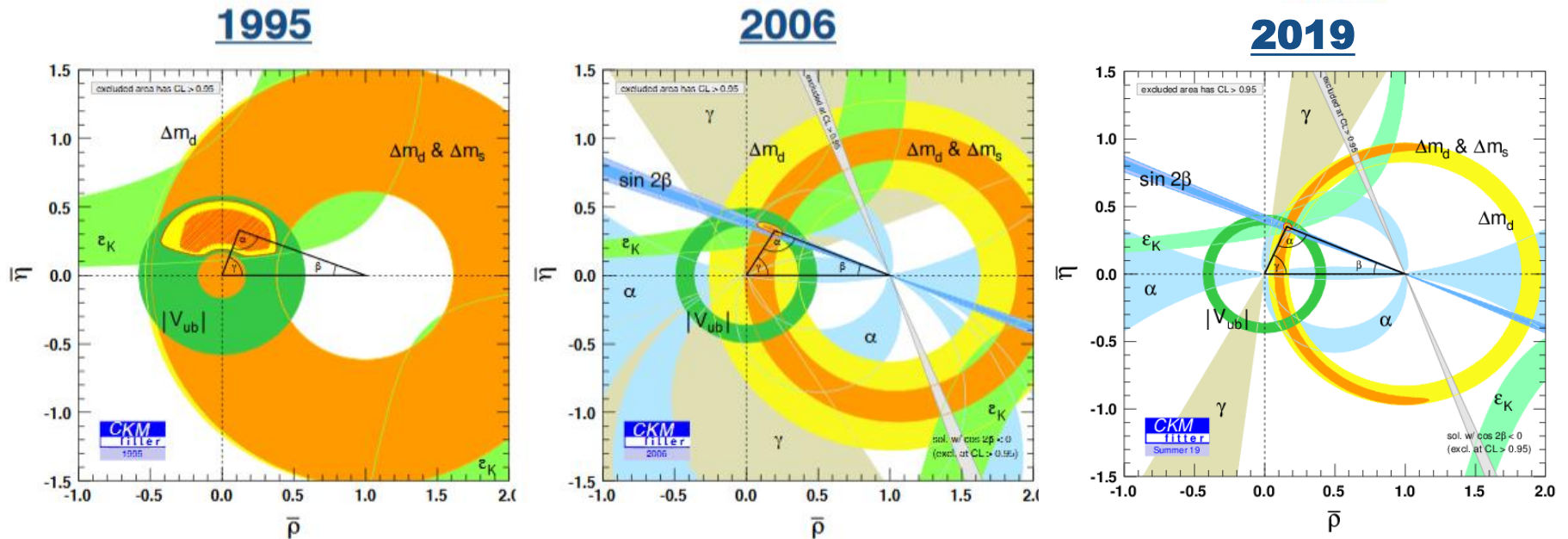
The idea: try to measure as many flavour observables as possible
overconstrain the unitarity triangle



- If all measurements meet in the same apex → good understanding of the flavour structure of the SM
- If not → New Physics !

The CKM matrix

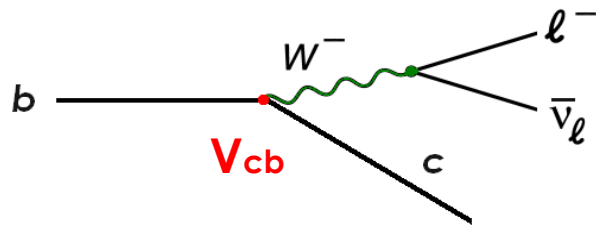
The Unitarity Triangle



- Very precise measurements of CKM elements (fundamental parameters!)
- Angles and sides have to be measured in many different ways to look for inconsistencies \rightarrow quantum effects from new particles
- Tree level processes (new physics is less expected) should be compared to loop processes sensitive to new particles
- If everything is consistent with increasing precision the new physics scale has to be higher

The CKM matrix

Measuring CKM elements from semileptonic decays (an example)
(sides of the Unitarity Triangle)



$$V_{\text{CKM}} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

Decay width, partial width and lifetime:

$$\Gamma = \sum_f \Gamma_f = \frac{1}{\tau}$$

Decay length:

$$\langle \Delta \ell \rangle = \beta \gamma c \tau$$

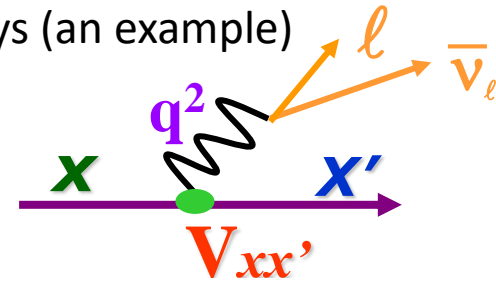
Branching fraction:

$$1 = \sum_{f=1}^m \frac{\Gamma_f}{\Gamma} = \sum_{f=1}^m B_f$$

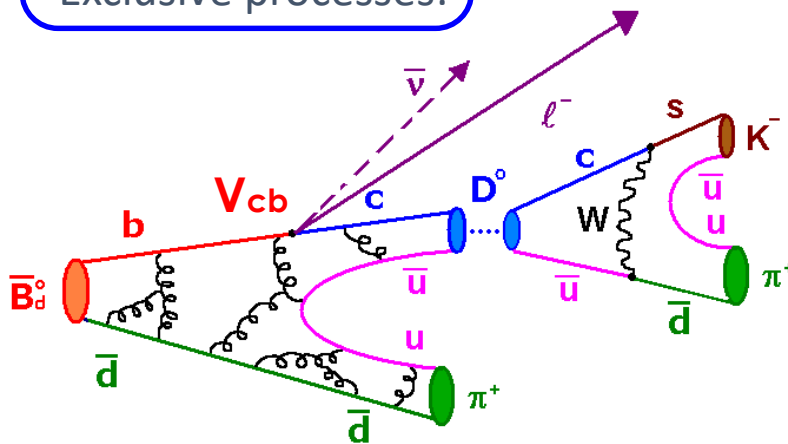
The CKM matrix

Measuring CKM elements from semileptonic decays (an example)
(sides of the Unitarity Triangle)

$$\Gamma_{sl} = |V_{xx'}|^2 f(\text{theory})$$



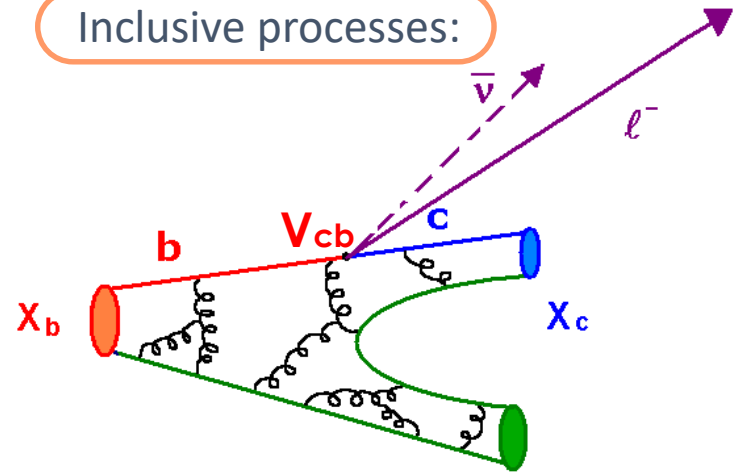
Exclusive processes:



$$\Gamma_{sl}^{D^*} = \frac{\mathcal{B}(\bar{B} \rightarrow D^* l^- \bar{\nu}_l)}{\tau_{\bar{B}}}$$

Theory: HQET, LQCD, LCSR
Experiment: low backgrounds (mainly D**)

Inclusive processes:



$$\Gamma_{sl}^{tot} = \frac{\mathcal{B}(b \rightarrow c l^- \bar{\nu}_l)}{\tau_b}$$

Theory: OPE
Experiment: higher backgrounds 15

The CKM matrix

Measurement of the $|V_{cb}|$ element by MARK II at PEP e+e- collider at 29 GeV (California, 1983)

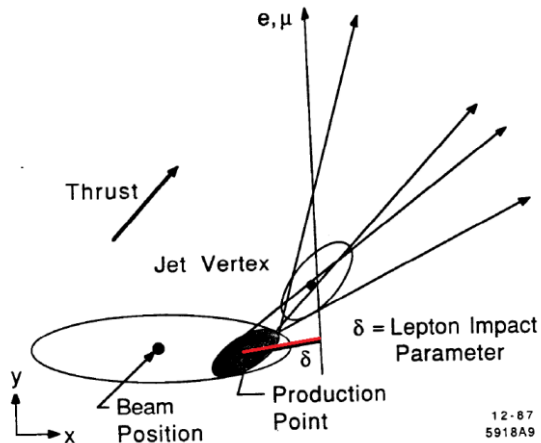
Lifetime of b-hadrons

Phys. Rev. Lett. 51 (1983) 1316

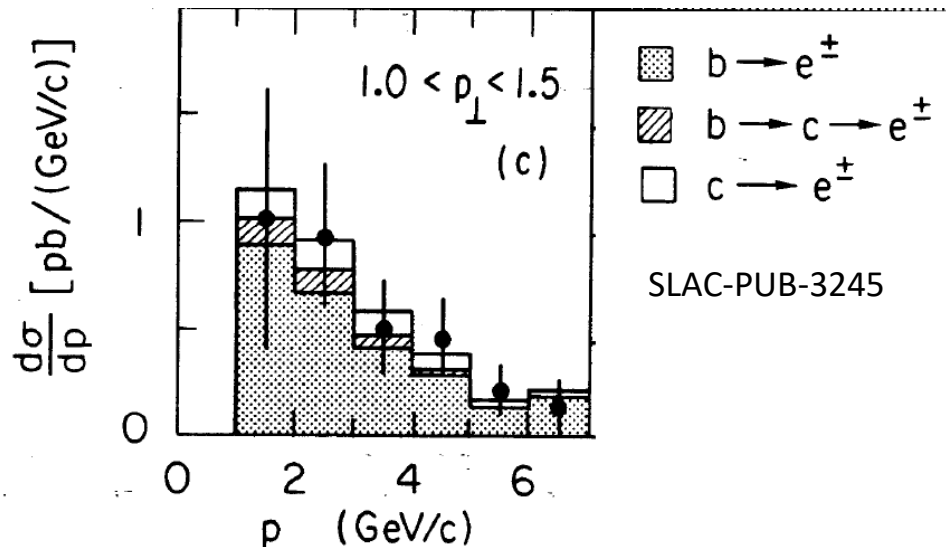
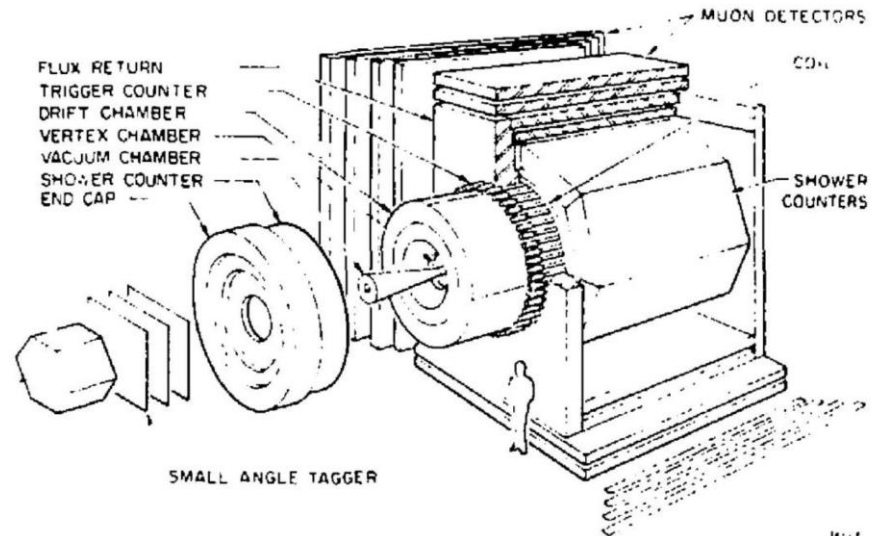
$$\frac{1}{\tau_b} \equiv \Gamma_{tot} = \frac{\Gamma(b \rightarrow x l \nu)}{\text{BR}(b \rightarrow x l \nu)}$$

$$\Gamma(b \rightarrow x l \nu) = \frac{G_F^2 m_b^5}{192 \pi^3} (|V_{ub}|^2 + 0.49 |V_{cb}|^2) (\text{GeV})$$

Impact parameter:

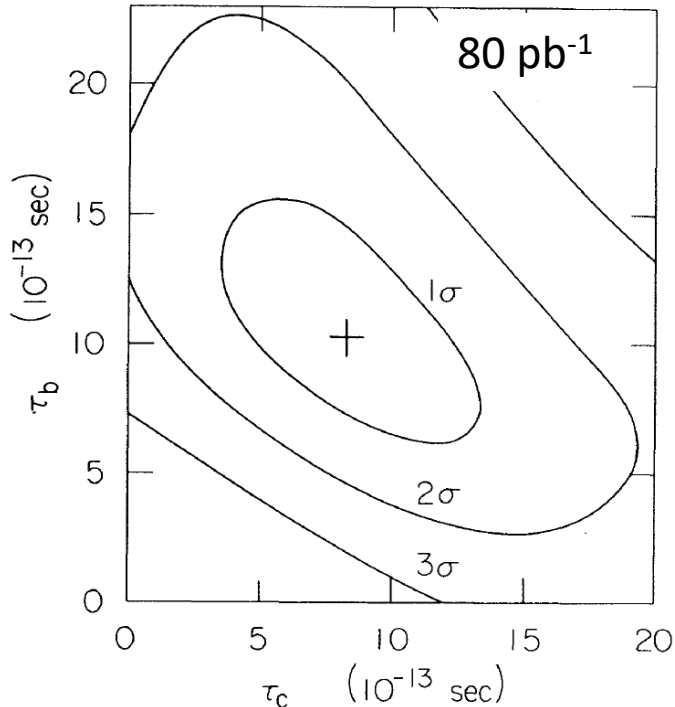


12-87
5918A9



The CKM matrix

Phys. Rev. Lett. 51 (1983) 1316



$$\tau_b = (12.0^{+4.5}_{-3.6} \pm 3.0) \times 10^{-13} \text{ sec}$$

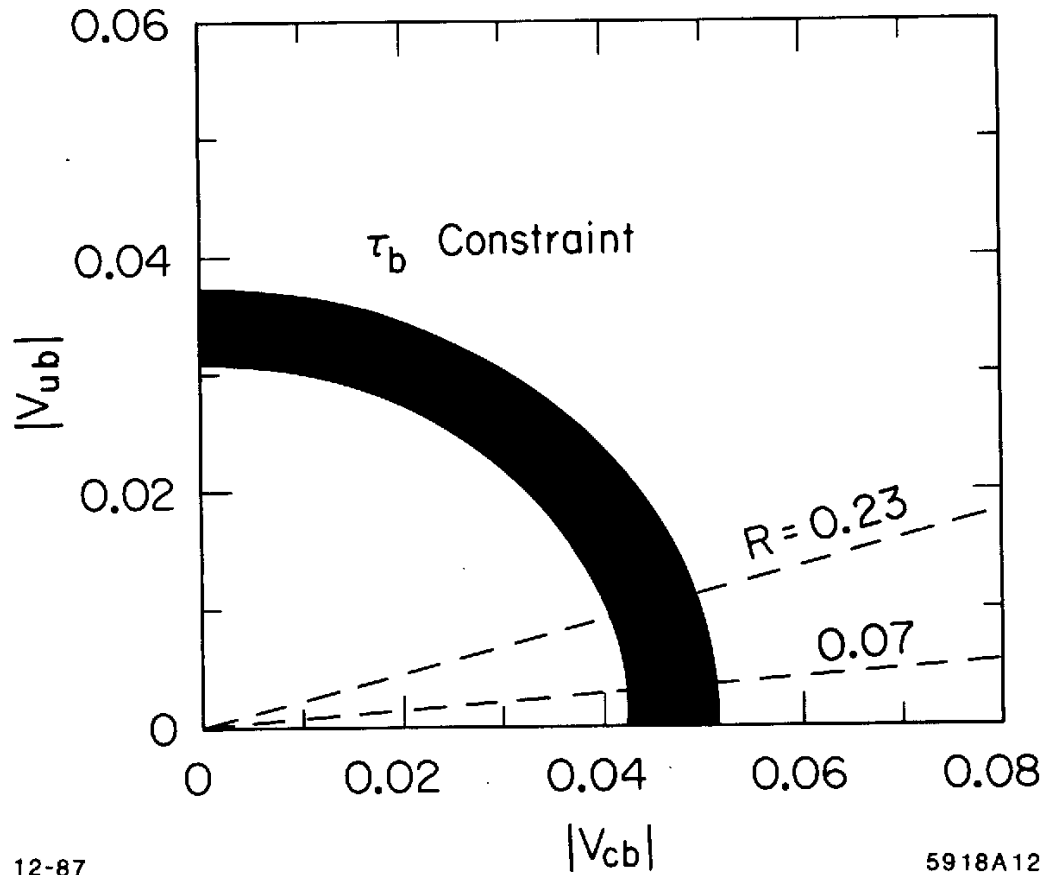
The B lifetime is long \rightarrow the mixing between the third generation of quarks and the lighter quarks is much weaker than the mixing between the two first generations.

the U_{bu} term in the expression for the lifetime is negligible. If we put $m_b = 5 \text{ GeV}/c^2$, we find $|U_{bc}| = 0.053^{+0.010}_{-0.009}$, where the error is statistical only. This value is appreciably smaller than that of the analogous matrix element which describes strange-particle decay, $U_{su} = 0.22$, the sine of the Cabibbo angle.

$\sim 20\%$ error

The CKM matrix

Status by December 1987...



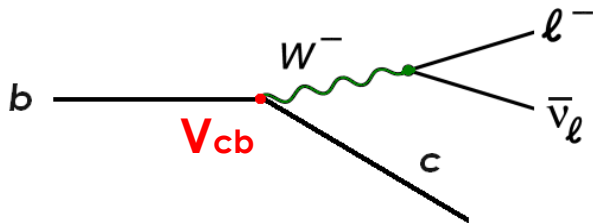
12-87

5918A12

Constraints on the KM matrix elements for B decay. The dark band corresponds to the constraint imposed by the B lifetime measurements. The dotted lines reflect the constraints imposed by the ratio $R = |V_{ub}|/|V_{cb}|$. SLAC-PUB-4503

The CKM matrix

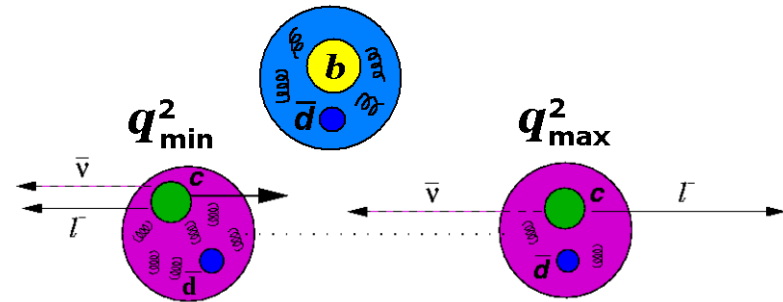
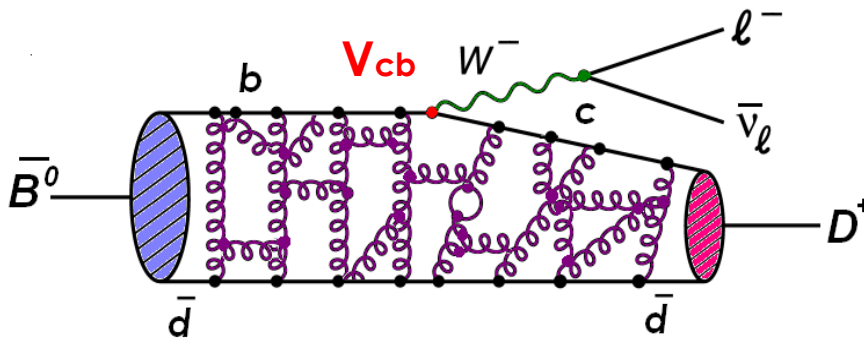
Measuring CKM elements from semileptonic decays (an example)
(sides of the Unitarity Triangle)



$$V_{\text{CKM}} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

Quarks are confined inside the hadrons

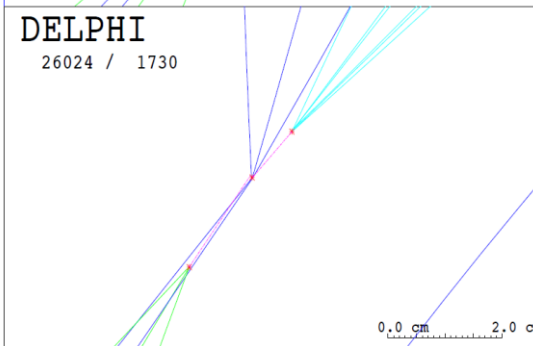
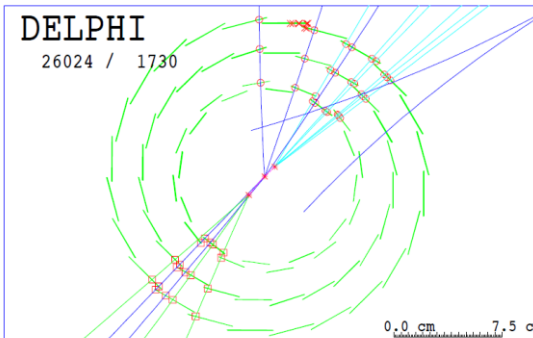
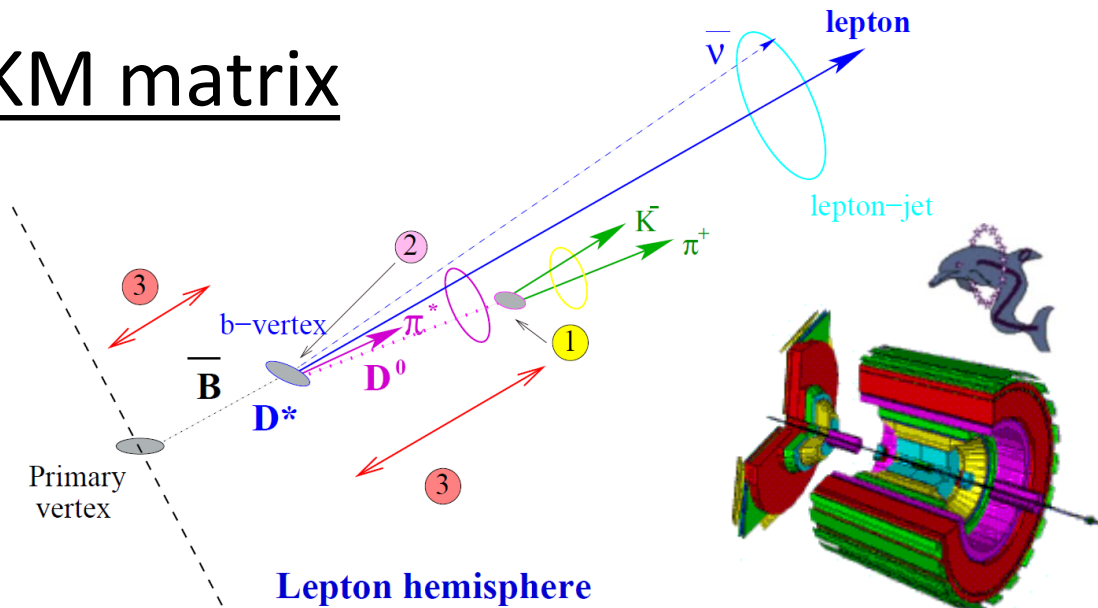
$$q^2 = (p_{\ell^-} + p_{\bar{\nu}_\ell})^2 = (p_{\bar{B}} - p_{D^*})^2$$



$$\frac{d\Gamma(B \rightarrow D^* l \bar{\nu})}{dq^2} = \frac{G_F^2}{48\pi^3} |V_{cb}|^2 F(q^2)^2 \mathcal{K}(q^2)$$

The CKM matrix

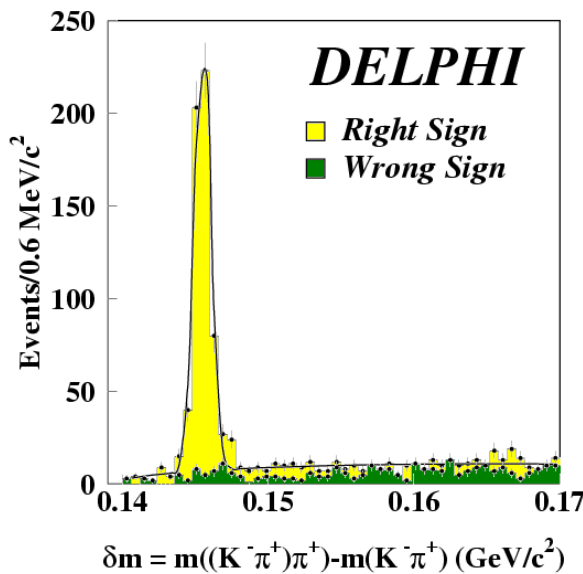
20 years after MARK II,
and before the B factories:
 V_{cb} from Z decays at LEP



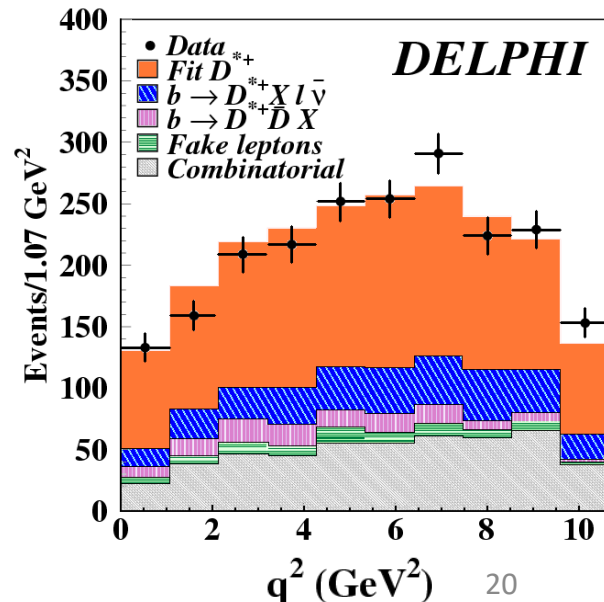
$D^{*+} - \ell^-$ candidates



Eur. Phys. J. C33 (2004) 213



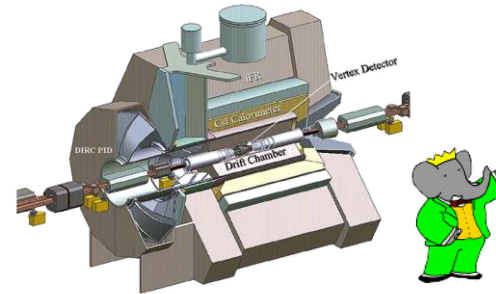
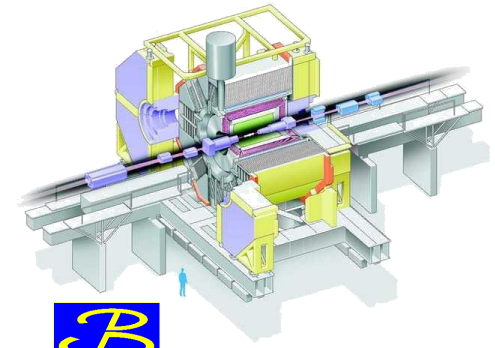
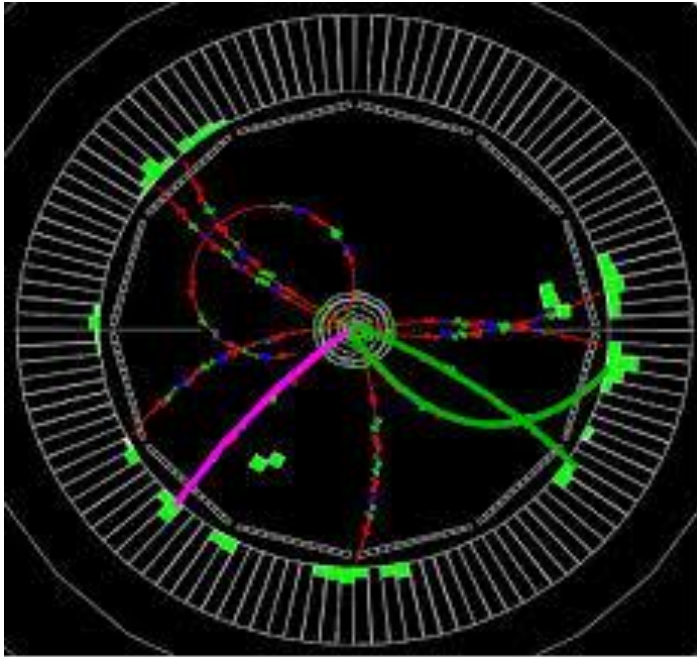
$\delta|V_{cb}|/|V_{cb}| \sim 5\%$



$$q_{\text{max}}^2 = (m_B - m_{D^*})^2 \approx 10.69 \text{ GeV}^2$$

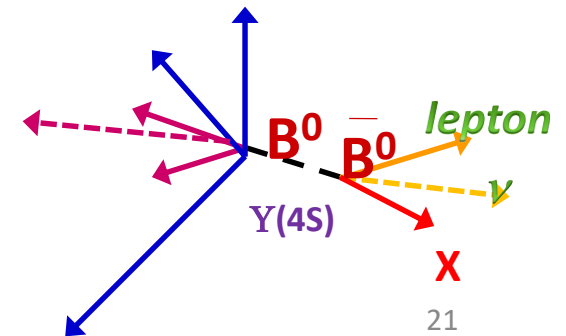
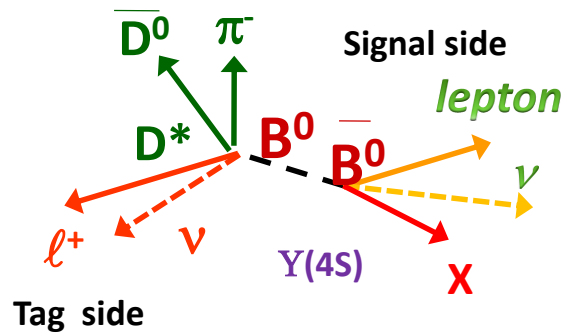
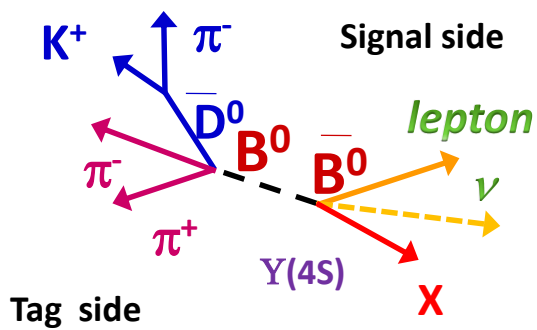
The CKM matrix

Measurements from BELLE and BaBar:



Hadronic or Semileptonic Tag (High purity)

Untagged (High efficiency)



The CKM matrix

$$\frac{d\Gamma(B \rightarrow D^{(*)} \ell \nu)}{dw} = \frac{G_F^2}{48\pi^3 \hbar} \mathcal{F}^2(w) \mathcal{K}(w) |V_{cb}|^2$$

Form factor
(shape & normalization)

Kinematic factor

$$w = V_B \cdot V_{D^{(*)}} = \frac{(M_B^2 + M_{D^{(*)}}^2 - q^2)}{(2M_B M_{D^{(*)}})}$$

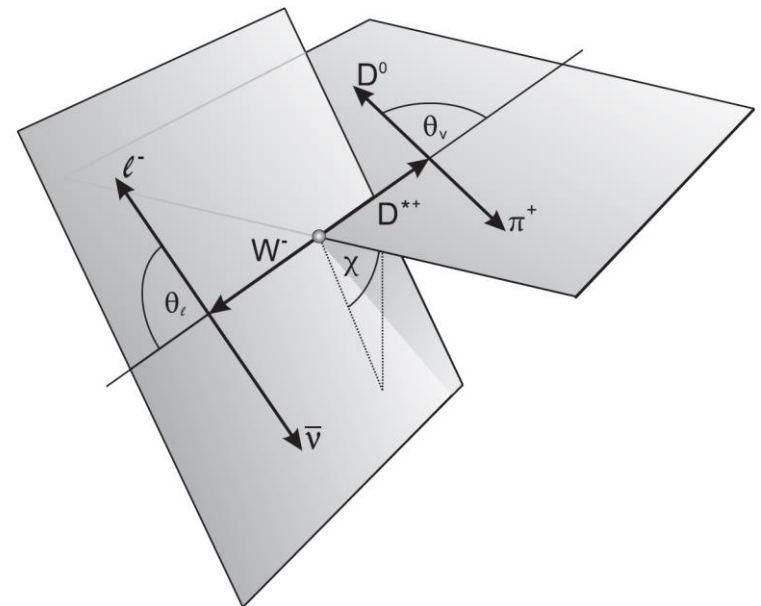
(w ranges from 1 to 1.5)

Normalized by HQET ($m_Q \rightarrow \infty$) at q^2_{\max} ($w=1$) $\equiv F_{D^*}(1)=1$

LQCD $\rightarrow F_{D^*}(1) = 0.921 \pm 0.024$

Belle and BaBar have a lot of data! ($\sim 700 \text{ fb}^{-1}$)
Possible to fit also the angular distribution of the decay particles:

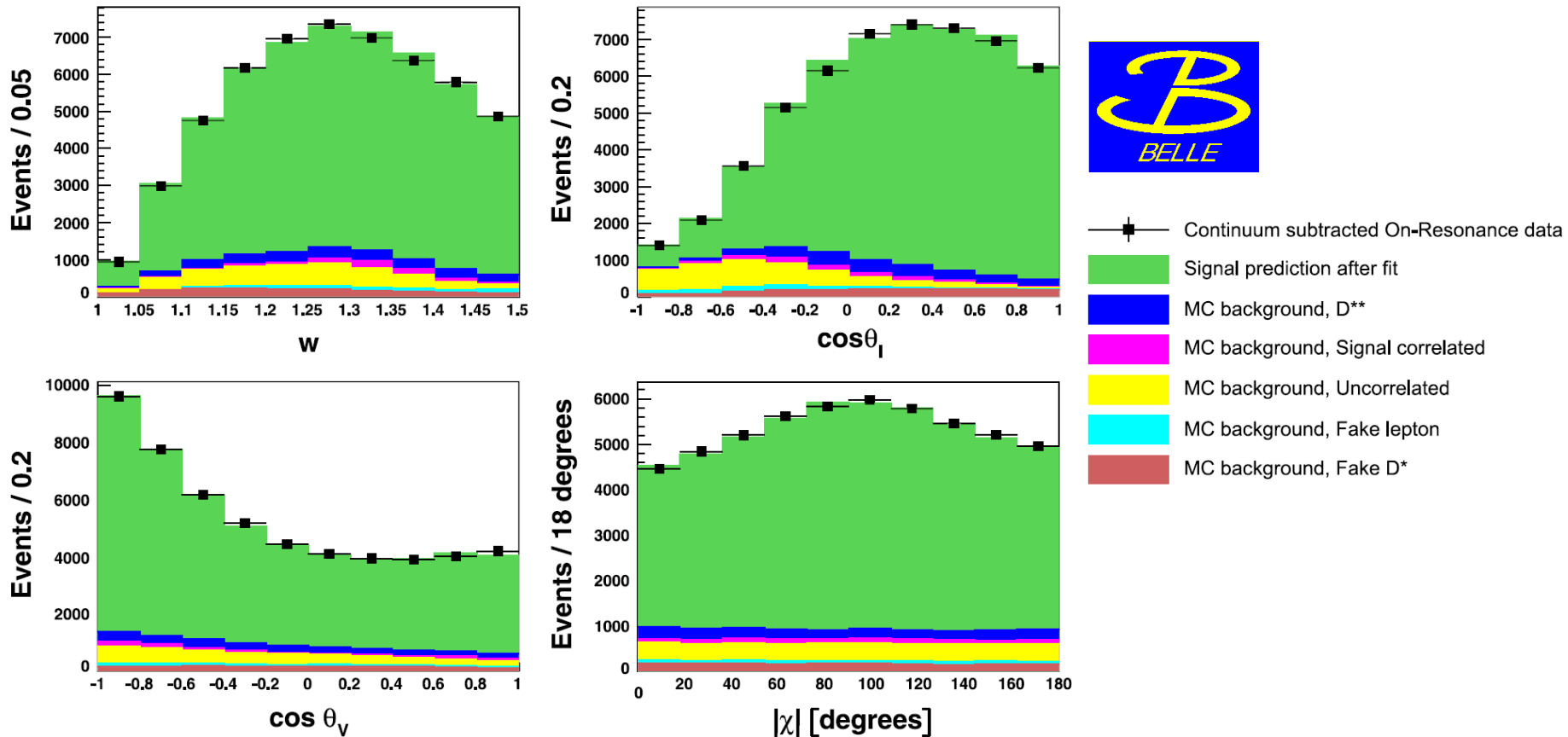
$$\frac{d^4\Gamma(B^0 \rightarrow D^{*-} \ell^+ \nu_\ell)}{dw d(\cos\theta_\ell) d(\cos\theta_V) d\chi}$$



$$J(B^0) = 0, \quad J(D^*) = 1, \quad J(W) = 1$$

The CKM matrix

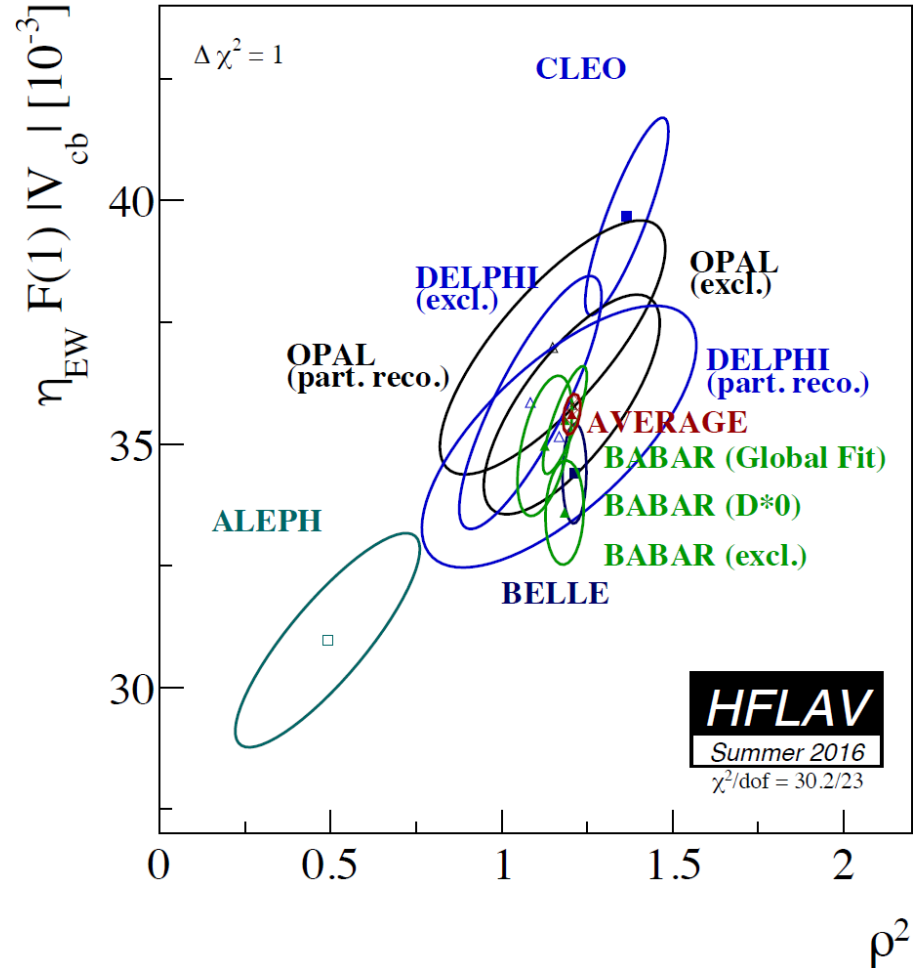
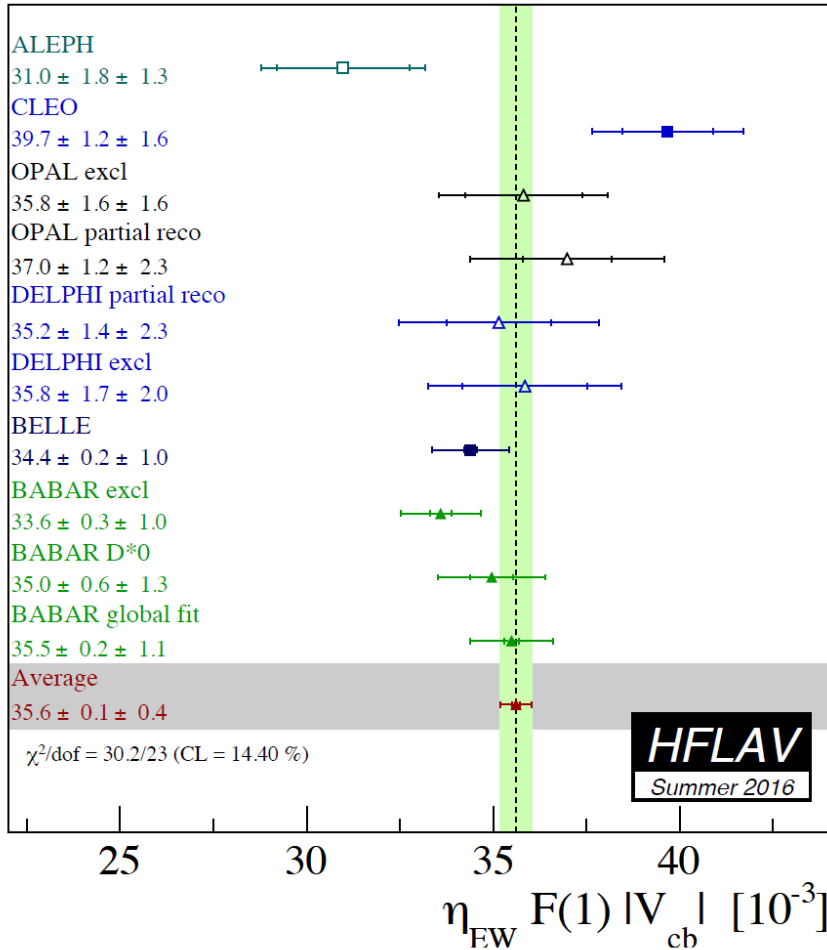
120000 $B \rightarrow D^* \ell \nu$! (muons and electrons), very few backgrounds



A 4 dimensional fit allows the determination of the form factor shape and $|V_{cb}|$

The CKM matrix

After BELLE and BaBar:



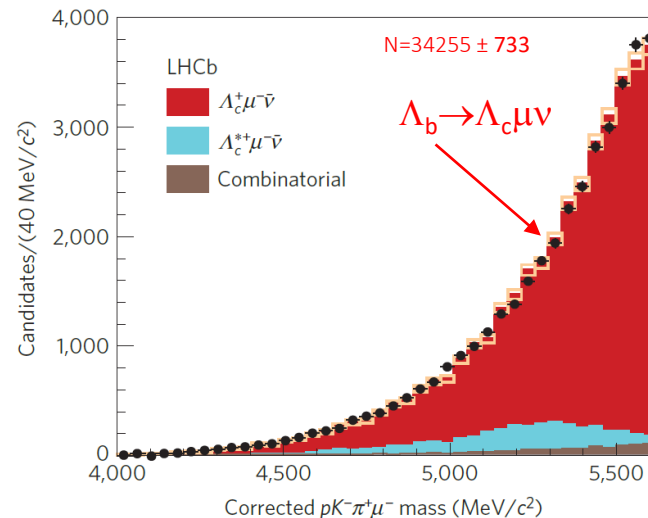
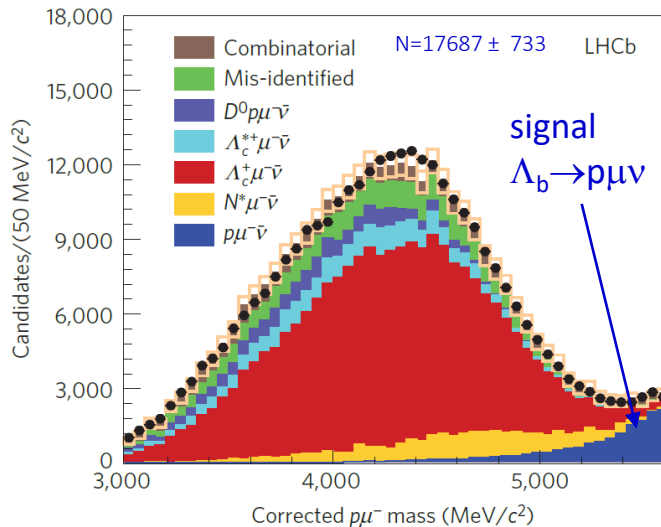
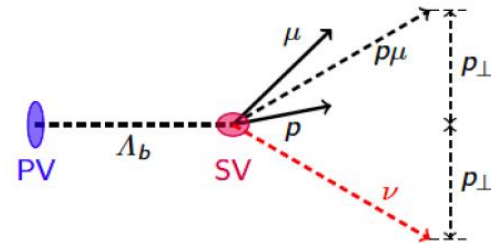
$\delta|V_{cb}|/|V_{cb}| \sim 1\% !$

The CKM matrix

At LHCb semileptonic decays are very challenging due to the missing neutrino:
Using semileptonic decays of b-baryons:

$$\frac{\mathcal{B}(\Lambda_b^0 \rightarrow p\mu^-\bar{\nu}_\mu)}{\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+\mu^-\bar{\nu}_\mu)} = \frac{|V_{ub}|^2}{|V_{cb}|^2} \times \text{Ratio of form factors} \quad \leftarrow \text{(5\% accuracy from LQCD)}$$

- Use information from displaced vertex
- Select high q^2 region (theory more precise)
- Corrected mass: $m_{\text{corr}} = \sqrt{m_{h\mu}^2 + p_\perp^2} + p_\perp$



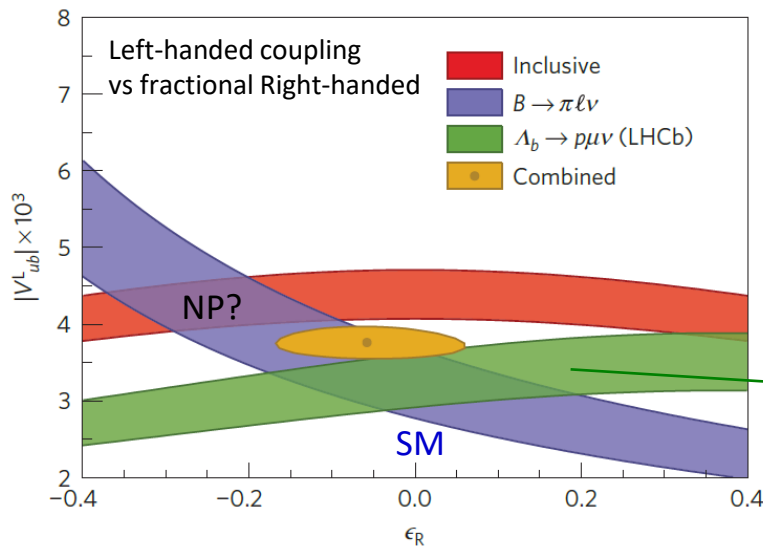
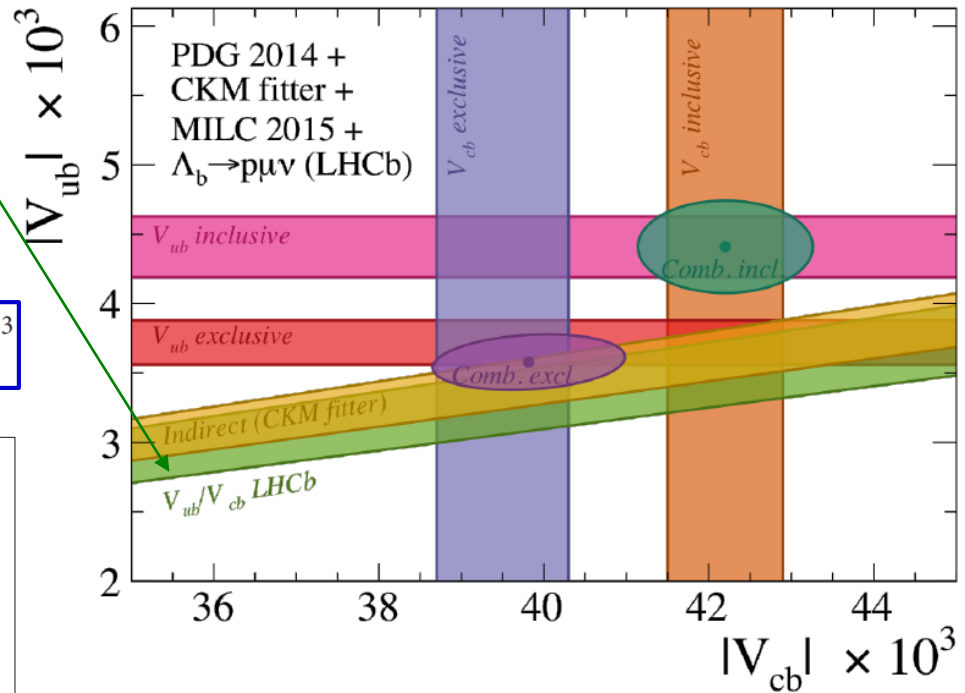
The CKM matrix

[Nature Physics 10 (2015) 1038]

$$\frac{|V_{ub}|}{|V_{cb}|} = 0.083 \pm 0.004 \pm 0.004$$

Using the world average from exclusive V_{cb} :

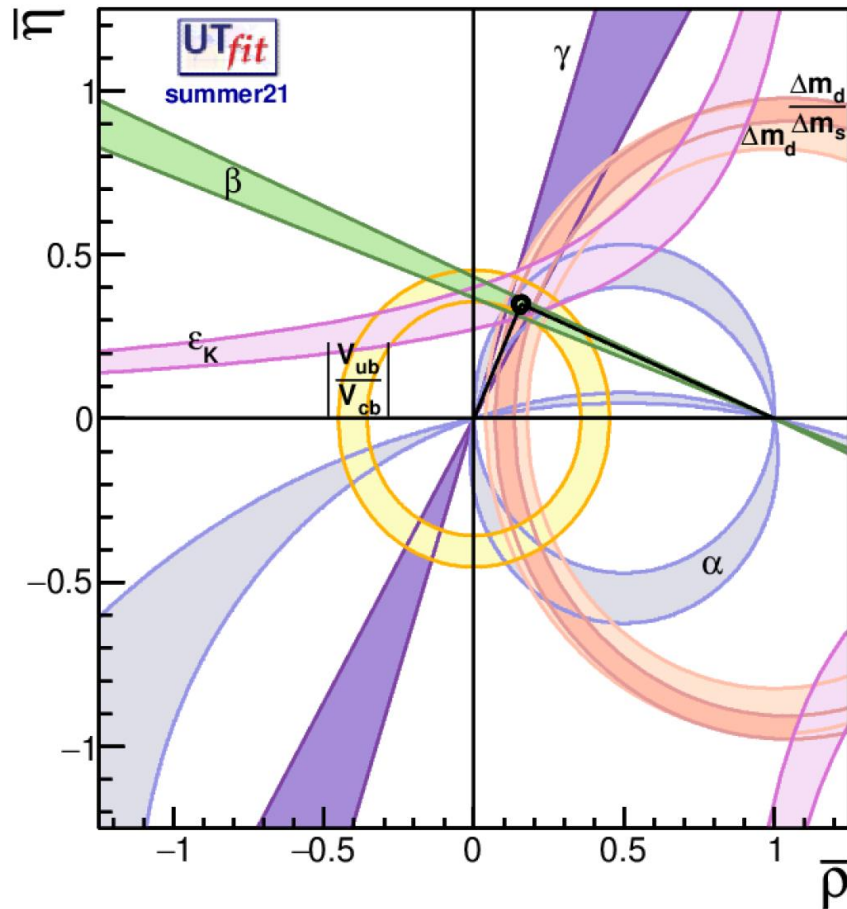
$$|V_{ub}| = (3.27 \pm 0.15 \pm 0.16 \pm 0.06) \times 10^{-3}$$



Disfavours New Physics models with Right-handed currents

The CKM matrix

Present status:



- Very high level of precision (few %)
- No inconsistencies: validation of Standard Model in the flavour sector
- Understanding from QCD is crucial

$$\bar{\rho} = 0.157 \pm 0.012$$
$$\bar{\eta} = 0.350 \pm 0.010$$

<http://ckmfitter.in2p3.fr/>

<http://www.utfit.org/UTfit/>

The CKM matrix

- Nevertheless, some anomalies are found at tree level in B semileptonic decays
 → Ratio of semi-tauonic and semi-muonic branching fractions:

$$\mathcal{R}(D^*) = \frac{\mathcal{B}(\bar{B}^0 \rightarrow D^{*+} \tau^- \bar{\nu}_\tau)}{\mathcal{B}(\bar{B}^0 \rightarrow D^{*+} \mu^- \bar{\nu}_\mu)}$$

SM predictions very precise :

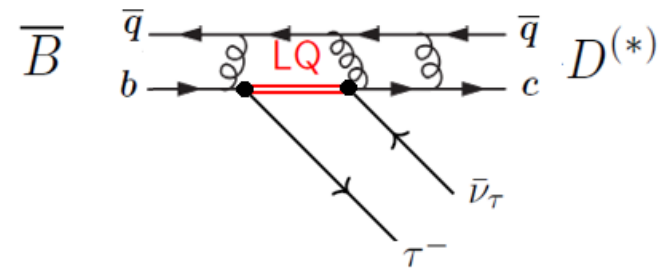
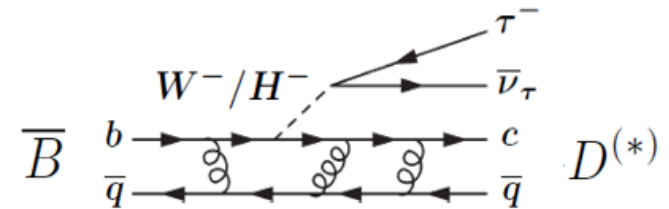
(V_{cb} and form factors (partially) cancel)

$$\mathcal{R}(D)_{SM} = 0.299 \pm 0.003$$

$$\mathcal{R}(D^*)_{SM} = 0.252 \pm 0.003$$

▶ Test of lepton universality

▶ Sensitive to charged Higgs bosons and leptoquarks



The CKM matrix

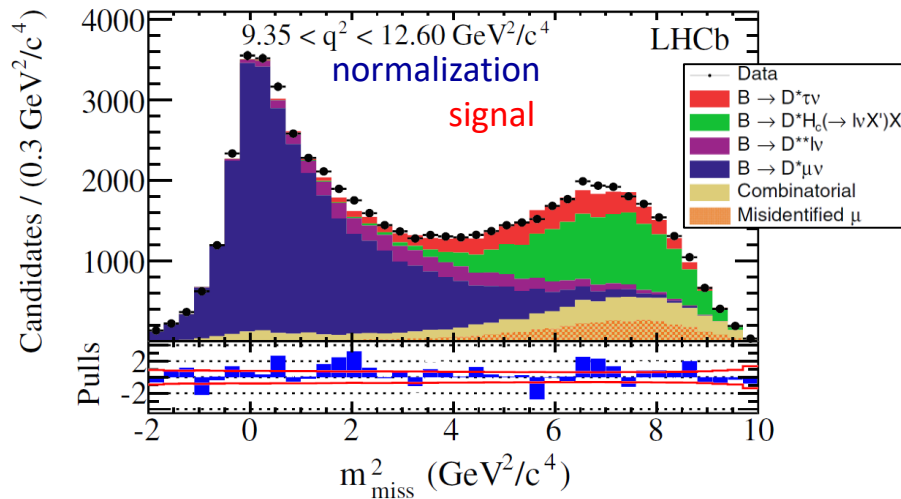
BaBar measured an excess of $B^0 \rightarrow D^{(*)} \tau^- \bar{\nu}_\tau$ (**3 σ away from SM!**) [Nature 546 (2017) 227]

LHCb: $R(D^*)$ $\left\{ \begin{array}{l} \blacksquare B^0 \rightarrow D^{*+} \tau^- \bar{\nu}_\tau, \text{ with } \tau^- \rightarrow \mu^- \bar{\nu}_\mu \nu_\tau \text{ [PRL 115 (2015) 111803]} \\ \blacksquare B^0 \rightarrow D^{*-} \tau^+ \nu_\tau, \text{ with } \tau^+ \rightarrow \pi^+ \pi^- \pi^+ (\pi^0) \bar{\nu}_\tau \text{ [PRL 120 (2018) 171802]} \end{array} \right.$

■ Using $\tau^- \rightarrow \mu^- \bar{\nu}_\mu \nu_\tau$

Information from the missing mass squared $m_{\text{miss}}^2 = (P_B - P_{D^*} - P_\mu)^2$ and muon energy

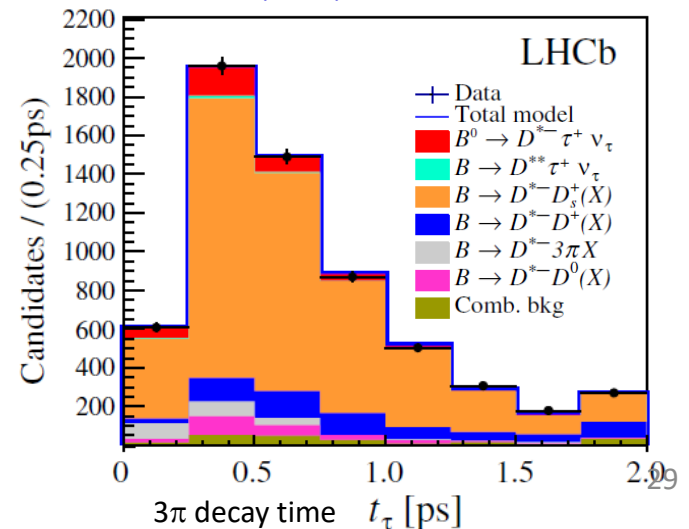
[PRL 115 (2015) 111803]



■ Using $\tau^+ \rightarrow \pi^+ \pi^- \pi^+ \bar{\nu}_\tau$

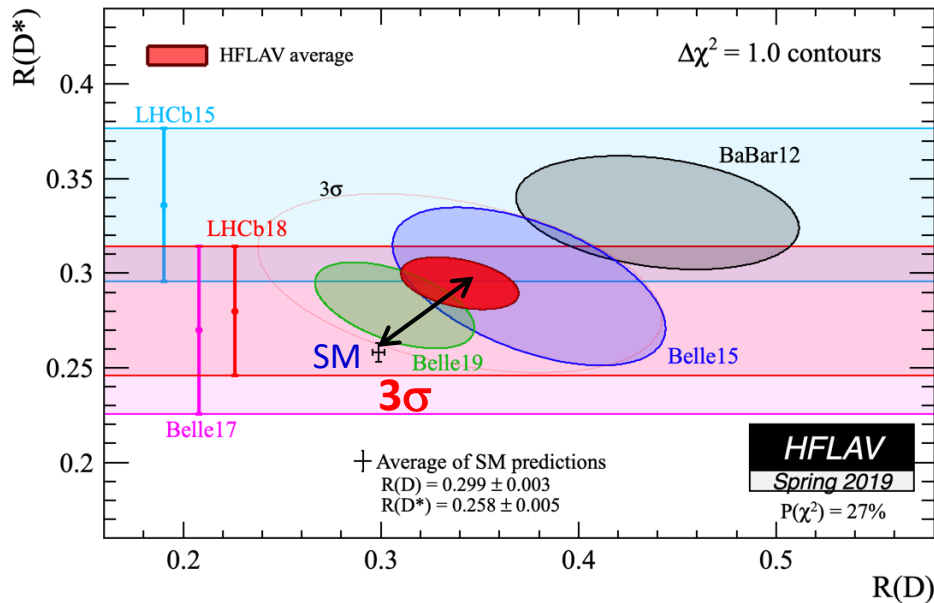
Information from the position of the pions. Normalized to $B^0 \rightarrow D^{*-} \pi^+ \pi^- \pi^+$

[PRL 120 (2018) 121801]



The CKM matrix

- Present global picture of R_D and R_{D^*} :



→ Average: 3σ deviation from SM

