

ICTS, Non-Hermitian physics 23/03/21

Stochastic description
of noisy
open quantum systems

Tony Jin

University of Geneva

Open systems



Open systems



- In general very difficult
ex: Kondo problem

Open systems



- In general very difficult
ex: Kondo problem
- Simplifying hypothesis
↳ Markov property

GKSL equation (Gorini, Kosakowski, Sudanshan, Lindblad)

Map s.t. : * completely positive * trace-preserving
* Markov property

GKSL equation (Gorini, Kossakowski, Sudanshan, Lindblad)

Map s.t. : * completely positive * trace-preserving
* Markov property

\hookrightarrow := Quantum channel \mathbb{I}

GKSL equation (Gorini, Kossakowski, Sudarshan, Lindblad)

Map s.t. : * completely positive * trace-preserving
* Markov property

\hookrightarrow := Quantum channel Φ

$\mathcal{L}(\rho)$ generator of the map : $e^{\mathcal{L}t}[\rho] := \Phi[\rho]$

$$\frac{d}{dt} \rho = -i[H, \rho] + \sum_i \gamma_i (L_i \rho L_i^\dagger - \frac{1}{2} \{L_i^\dagger L_i, \rho\})$$

\hookrightarrow "Non-Hermitian" complex spectrum

Lifting the Lindblad equation

Stinespring's theorem:

Always possible to enlarge the space enough to get Φ as a reduced unitary evolution

$$\Phi(\rho) = \text{Tr}_B (U \rho \otimes \rho_E U^\dagger)$$

↖ "Noise"

Lifting the Lindblad equation

Stinespring's theorem:

Always possible to enlarge the space enough to get Φ as a reduced unitary evolution

$$\Phi(\rho) = \text{Tr}_B (U \rho \otimes \rho_E U^\dagger)$$

average \searrow

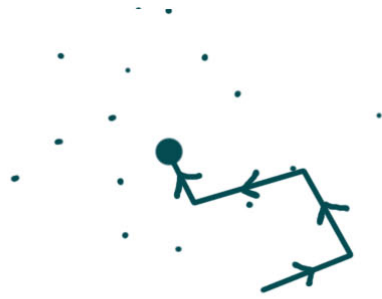
\nwarrow "Noise"

Stochastic unitary \rightarrow Quantum channel

Stochastic Hamiltonian \rightarrow Lindblad

Classical

Brownian
motion



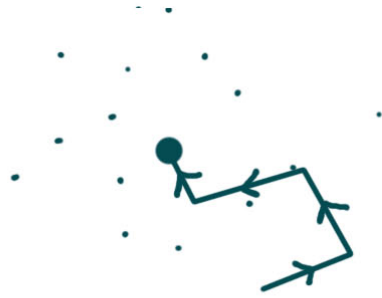
Langevin eq:

$$m dv = -\alpha v dt + d\beta t$$

$\frac{d\beta t}{dt}$ White noise

Classical

Brownian
motion



Langevin eq:

$$m dv = -\alpha v dt + d\beta t$$

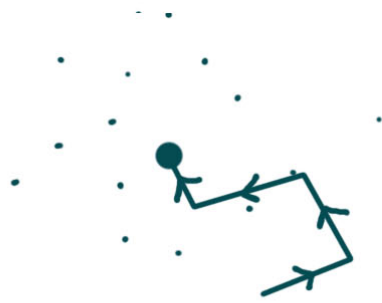
Damping eq: ↓ average

$$m dv = -\alpha v dt$$

$\frac{d\beta t}{dt}$ White noise

Classical

Brownian motion



Langevin eq:

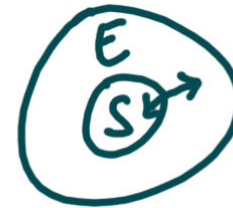
$$m dv = -\alpha v dt + d\beta t$$

Damping eq: ↓ average

$$m dv = -\alpha v dt$$

$\frac{d\beta t}{dt}$ White noise

Quantum

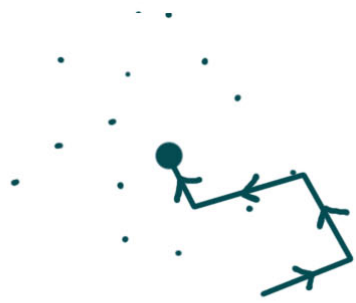


Stochastic Hamiltonian:

$$| \psi_{t+dt} \rangle = e^{-i dH t} | \psi_t \rangle$$

Classical

Brownian motion



Langevin eq:

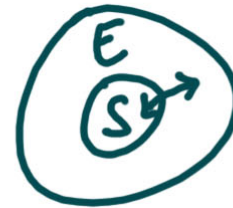
$$m dv = -\alpha v dt + d\beta t$$

Damping eq: ↓ average

$$m dv = -\alpha v dt$$

$\frac{d\beta t}{dt}$ White noise

Quantum



Stochastic Hamiltonian:

$$| \psi_{t+dt} \rangle = e^{-i dH t} | \psi_t \rangle$$

Lindblad eq: ↓ average

$$\frac{d}{dt} \rho = \mathcal{L}(\rho)$$

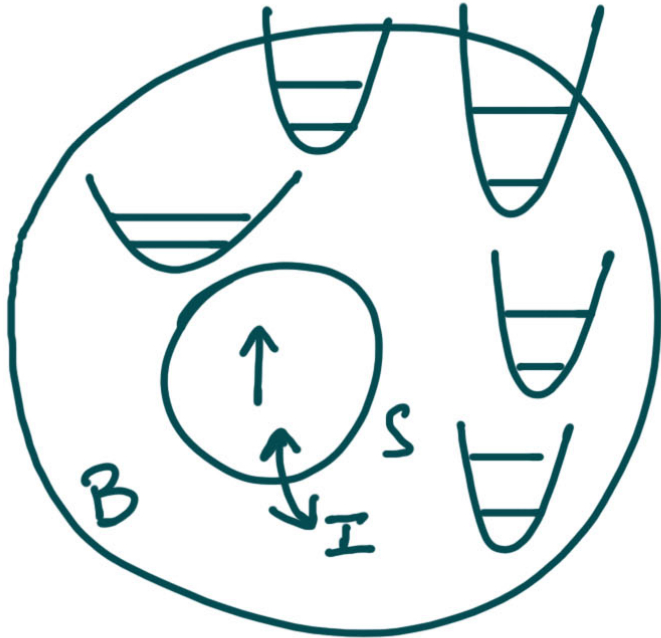
↑
Lift

Disclaimer: Similar tools but
 \neq weak measurements

Why should you care?

- Useful to have unitary evolution
- Rich physics beyond the mean
- Stochastic process in the Hilbert space

Spin-boson (Caldeira-Leggett)



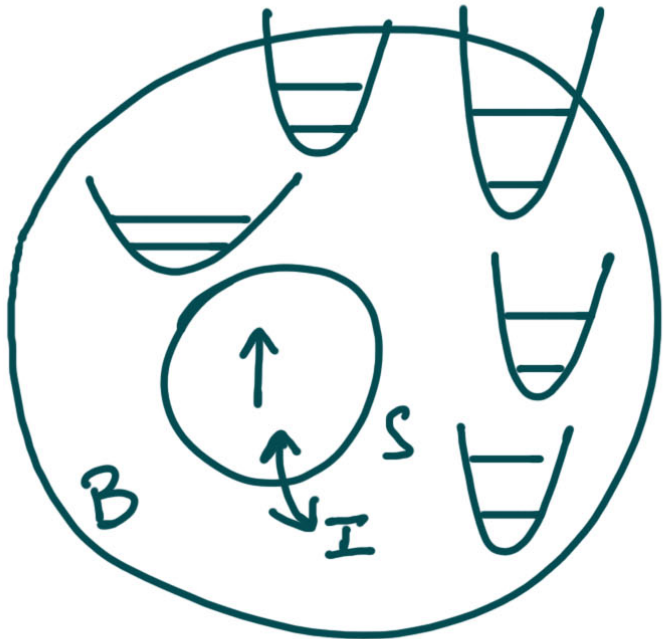
$$H = H_S + H_B + H_I$$

$$H_S = 0$$

$$H_B = \int \frac{dk}{2\pi} \epsilon_k a_k^\dagger a_k$$

$$H_I = \int \frac{dk}{2\pi} \chi_k (a_k^\dagger \sigma_S^- + a_k \sigma_S^+)$$

Spin-boson (Caldeira-Leggett)



$$H = H_S + H_B + H_I$$

$$H_S = 0$$

$$H_B = \int \frac{dk}{2\pi} \epsilon_k a_k^\dagger a_k$$

$$H_I = \int \frac{dk}{2\pi} z_k (a_k^\dagger \sigma_S^- + a_k \sigma_S^+)$$

$$|\psi_{t+dt}\rangle = e^{-idHt} |\psi_t\rangle \downarrow \text{Markov hypothesis}$$

$dHt = \sigma_S^+ d\bar{W}_t + \sigma_S^- dW_t$	$dW_t = dB_t^1 + i dB_t^2$ $d\bar{W}_t = dB_t^1 - i dB_t^2$
---	--

$$dH_t = \sigma_s^+ d\bar{W}_t + \sigma_s^- dW_t$$

$$|\psi_{t+dt}\rangle = e^{-idH_t} |\psi_t\rangle$$

↓ average

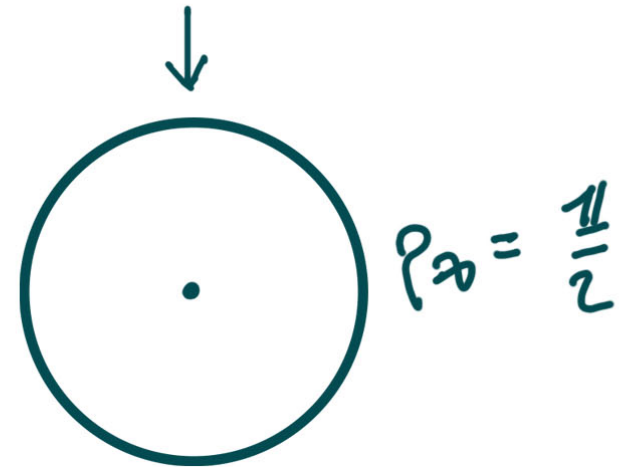
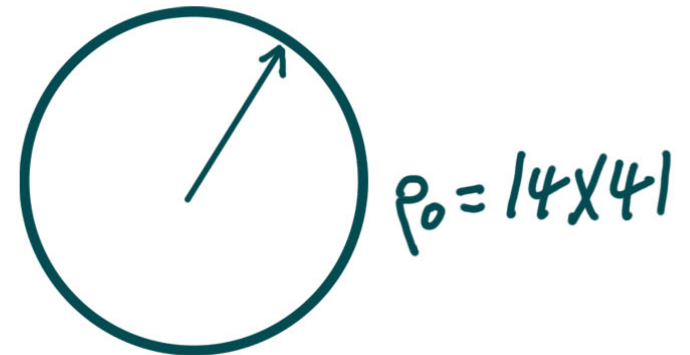
$$\begin{aligned} \frac{d\rho}{dt} = & \sigma^+ \rho_t \sigma^- - \frac{1}{2} \{ \sigma^+ \sigma^+, \rho_t \} \\ & + \sigma^- \rho_t \sigma^+ - \frac{1}{2} \{ \sigma^+ \sigma^-, \rho_t \} \end{aligned}$$

$$dH_t = \sigma_s^+ d\bar{W}_t + \sigma_s^- dW_t$$

$$|\psi_{t+dt}\rangle = e^{-idH_t} |\psi_t\rangle$$

↓ average

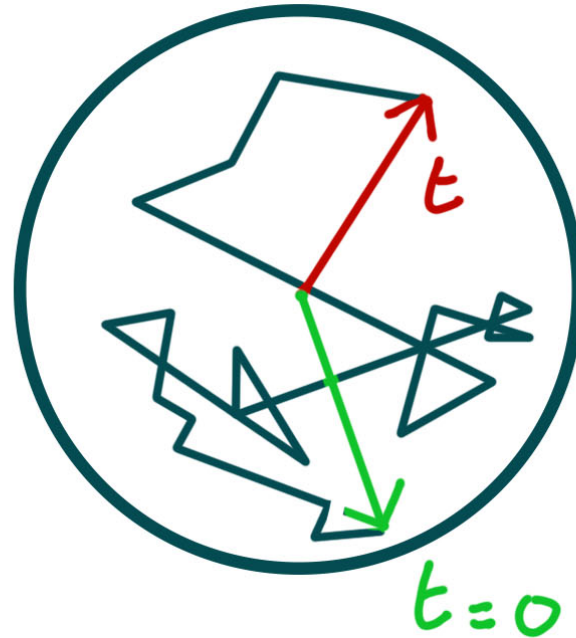
$$\frac{d\rho}{dt} = \sigma^+ \rho_t \sigma^- - \frac{1}{2} \{ \sigma^+ \sigma^+, \rho_t \} \\ + \sigma^- \rho_t \sigma^+ - \frac{1}{2} \{ \sigma^+ \sigma^-, \rho_t \}$$



Decoherence towards the
centre of the Bloch sphere.

Stochastic picture

Ergodic random walk on the Bloch sphere



No decoherence at the stochastic level

$$\text{Ex: } \langle \text{tr}(\rho S^j)^2 \rangle = \frac{1}{3} \quad j = x, y, z$$

Stationary distribution

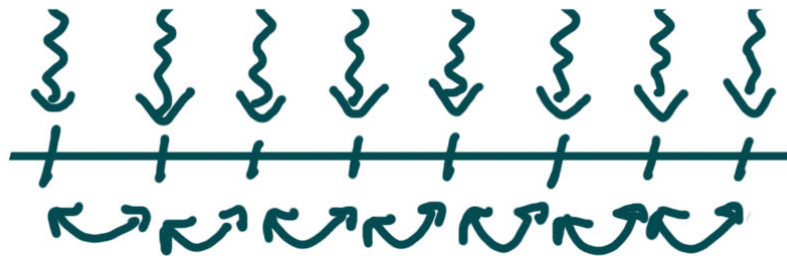
$$Z(A) := \langle \exp(\text{tr} A \rho) \rangle$$

$$= \int_{U \in \text{SU}(2)} d\eta(U) e^{\text{tr}(AU \rho_0 U^\dagger)}$$

↳ The momenta probe the structure of the invariants of $\text{SU}(2)$

Many-body systems

Stochastic dephasing model

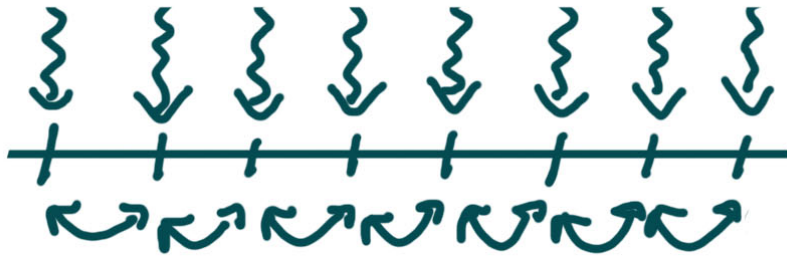


Quantum exclusion processes.



Many-body systems

Stochastic dephasing model



Quantum exclusion processes.



- Diffusive transport
- Rich structure beyond the mean
- First step towards quantum MFT?

References

• Quantum noise

«Path integral approach to quantum Brownian motion»

Caldeira-Legett, 1983

«Quantum Ito's formula and stochastic evolutions»

Hudson, Parthasarathy, 1984

• Dephasing model

«Dephasing-induced diffusive transport in the anisotropic Heisenberg model»

Znidaric, 2009

«Stochastic dissipative quantum spin chains (I): Quantum fluctuating discrete hydrodynamics»

Bauer, Bernard, Jin, 2017

• Quantum exclusion processes

«Equilibrium fluctuations in maximally noisy extended quantum systems»

Bauer, Bernard, Jin, 2018

«Open Quantum Symmetric Simple Exclusion Process»

Bernard, Jin, 2019

«From stochastic spin chains to quantum Kardar-Parisi-Zhang dynamics»

Jin, Krajenbrink, Bernard, 2020

Thank you!