Single-particle and collective excitations in transitional nuclei

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Plan of the talk :

- Physics interest in transitional nuclei
- Role of intruder orbital and triaxiality
- Experiments with INGA High spin band structure in Ce and Tl nuclei
- Development of collectivity towards neutron rich nuclei



Shapes and excitation in transitional nuclei

- Interplay of single particle and collective excitation
- Band structures based on multi-quasiparticle excitation
- Role of
 - intruder shape driving orbitals
 - triaxial degrees of freedom
- New modes of excitation
 - Magnetic rotation
 - Chiral bands



Rotation in weakly deformed system : MR band



Novel way of generating angular momentum in a weakly deformed system with low quadrupole collectivity.

Tilted Axis Cranking (TAC)

S. Frauendorf NPA677, 115 (2000).

Very regular band structure ($\Delta I = 1$) with a few valance particles (holes) in high j orbital



Non-axial deformation $\rightarrow \gamma$ -soft or rigid triaxial ?



N.V. Zamfir and R.F. Casten PLB 260, 265 (1991)



Non-axial deformation \rightarrow triaxiality

Experimental manifestation

Signature inversion

explained as a manifestation of the drift of the rotation axis in the intrinsic frame when a triaxial nucleus rotates. - degree of triaxiality in a deformed basis

phase of the staggering of S(I) = E(I) - E(I-1) reverses

can as well be consistent with an axially symmetric shape

Wobbling excitation

TSD bands in ^{163,165}Lu

Chiral bands

based on the specific geometry of three components of the total angular momentum



R. Bengtsson et al., Winter meeting on Nucl. Phys. Bromino (1982)

I. Hamamoto, PLB235, 221 (1990)

Nuclear chirality

A Triaxial Nucleus becomes Chiral if it rotates about an axis that lies outside the three planes spanned by the principal axes of it's triaxial ellipsoidal shape.

 Two identical bands almost same Energy and Parity
 De-excitation of the partner bands in a very similar way



S. Frauendorf, J. Meng, NPA 617, 131 (1997) S. Frauendorf, Rev. Mod. Phys 73, 463 (2001)

 $\frac{A \sim 130}{\text{Odd} - \text{Odd} (\text{ph}_{11/2}\text{nh}_{11/2}^{-1})}$ $^{134}\text{Pr}, ^{136}\text{Pm}, ^{130}\text{Cs}, ^{128}\text{Cs}$ $\text{Odd} - A (\text{p}(\text{h}_{11/2})^2\text{nh}_{11/2}^{-1})$ ^{135}Nd

 $\frac{A \sim 105}{Odd - Odd (pg_{9/2}^{-1}nh_{11/2})}$ $^{106}Ag, \,^{106}Rh, \,^{104}Rh, \,^{102}Rh$ $Odd - A (pg_{9/2}^{-1}n(h_{11/2})^2)$ $^{107}Ag, \,^{105}Rh, \,^{103}Rh$

<u>A~190</u> Odd-Odd (ph_{9/2}ni_{13/2}) ¹⁸⁸Ir (?) ¹⁹⁸Tl

T. Koike, NPA 834, 36c (2010)



Nuclei with few particles or holes around Z/N = 82





Transitional nuclei in A~130 region

- Multi-quasiparticle excitations with valance protons and neutrons

 Shape driving effects, γ softness
- Presence of high-j h_{11/2} orbital for both protons and neutrons
- > Prolate driving $\pi h_{11/2}$ & Oblate driving $\nu h_{11/2}$
 - Structutre associated with $\pi h_{11/2} \otimes \nu h_{11/2}$ configuration in odd-odd nuclei







Properties of band structure in ¹³⁸Ce

0.8 mg/cm² ¹³⁰Te on 4.8 g/cm² Au backing





18 Clover Ge Indian National Gamma Array © IUAC, New Delhi



Level structure of ¹³⁸Ce



Multi-quasiparticle band in ¹³⁸Ce



Heavier rare earth nuclei in A~190-200

> Role of low- $\Omega \pi h_{9/2}$ orbital to generate the band structure - similar as the $h_{11/2}$ states for lighter rare-earth nuclei.

Rotational bands built on the intruder states i_{13/2} at higher spin.

The ground state of odd-A Tl nuclei are $1/2^+$ \rightarrow occupation of the unpaired proton in the $3s_{1/2}$ orbital below the Z = 82.

> $\pi h_{9/2}$ orbital above the Z = 82 shell closure, is accessible by the odd-proton in Tl nuclei for oblate deformation.

For heavy Tl isotopes, the neutron numbers favor the oblate deformation as they occupy the upper part of the i_{13/2} orbital.





Core (Hg) systematics

Even Hg isotopes





Possible chirality in the doubly-odd ¹⁹⁸Tl

E. A. Lawrie *et al.*, PRC 78, 021305(R) (2008)





High spin structure of ²⁰⁰TI, ²⁰¹TI



15 Clover Ge detectors INGA @TIFR





Band structure of ²⁰⁰Tl





 $\pi h_{9/2} \otimes \nu i_{13/2}$

A.J. Kreiner et al PRC 23, 748 (1981) Two minima with similar deformation of $\beta_2 = 0.12$

at near oblate shape ($\gamma = -64^{\circ}$) \rightarrow an oblate deformation for the yrast band

at triaxial shape (γ = 48°) \rightarrow possibility to observe Chiral structure



²⁰¹TI : states above h_{9/2} isomer





What is happening away from stability ?





New region of deformation near N=28

Breakdown of N=28 shell gap at Z=14



\rightarrow high collectivity

- \rightarrow strong deformation (β =0.45)
- \rightarrow no magicity



γ-soft

D. Sohler *et al.* PRC 66, 054302 (2002)

Shape mixing in low energy states

Coexistence of 1p-1h and 2p-2h Configuration

D. Santiago-Gonzalez et al PRC, 83, 061305(R) (2011)

Erosion of N=28 shell closure at Z=16



Development of collectivity near N=28

spherical

onset of deformation

shape coexistence

well deformed (axial)





Strong presence of N_{ph} excitations from vf_{7/2} orbital
 Role of neutron proton correlation on collective motion
 Strong quadrupole interaction between Protons in sd and neutrons in pf
 Rapid development of collectivity







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Signature of triaxiality in ⁴⁸Ar





Properties of calculated states





Summary and conclusion

Transitional nuclei in A~130 and A~200

 Role of shape driving intruder orbitals triaxial degrees of freedom

✓ Magnetic Rotation and multi-quasiparticle bands in ¹³⁸Ce

High spin spectroscopy of ^{200,201}TI

Development of collectivity near N=28 for neutron rich nuclei





Collaboration

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