<u>A wave-wave interaction mechanism</u> for near-inertial waves

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Wind-to-near-inertial-motions



M.H. Alford (2003)

Wind work on ocean inertial motion

Wind stress from reanalysis and ocean surface near-inertial motion from slab model (Pollard & Millard 1970)

Enhanced near-inertial wind work during storms

Total work done of the same order as tidal dissipation by conversion to internal tides

MIXED LAYER NEAR-INERTIAL OSCILLATIONS – GENERATION BY WIND STRESS

Sudden wind stress in the x direction (westerly wind) – Gill (1984)

$$\frac{\partial (U_E + iV_E)}{\partial t} + if(U_E + iV_E)) = \frac{\tau_x^S}{\rho}$$
$$U_E + iV_E = -i(\tau_x^S/\rho f)(1 - exp(-ift))$$

Steady state plus inertial oscillation

Pollard & Millard (1970) slab mixed layer model – add damping term

$$\frac{dZ}{dt} + (r_d + if)Z = \frac{T}{H}$$

$$Z = U_E + iV_E \qquad T = \rho^{-1}(\tau_x^S + i\tau_y^S)$$

$$H - mixed \ layer \ thickness$$

 $r_d^{-1} - decay time - arbitrarily chosen to fit observations$



NEAR-INERTIAL CURRENTS - GENERATION & DISSIPATION



Obvious downward propagation in early August (relatively large MLD); Fresh water influx in September – small MLD & strong stratification below – not so obvious downward propagation.

<u>Decay mechanisms for mixed layer near-</u> inertial currents

- Radiation of downward-propagating NIOs Gill (1984), D'Asaro (1989), Young & Ben Jelloul (1997)
- Nonlinear Interactions transferring energy to other frequencies Henyey et al. (1986)
- Turbulent Dissipation Hebert & Moum (1993)
- High-frequency internal waves Polton et al. (2008)

✓ Relative importance of each mechanism is unclear. In this talk, we will propose a potential nonlinear interaction mechanism

Internal wave modes – weakly nonlinear solutions

$$\begin{split} \frac{\partial^2}{\partial t^2} (\nabla^2 \psi) + f^2 \frac{\partial^2 \psi}{\partial z^2} &= \frac{g}{\rho^*} \frac{\partial}{\partial x} [J(\psi, \rho)] - \frac{\partial}{\partial t} [J(\psi, \nabla^2 \psi)] + f \frac{\partial}{\partial z} [J(\psi, v)] \\ \frac{\partial \rho}{\partial t} &= -J(\psi, \rho), \\ \frac{\partial v}{\partial t} + J(\psi, v) &= f \frac{\partial \psi}{\partial z}, \end{split} \quad \begin{split} & \frac{Fully \text{ nonlinear, } 2D, \text{ inviscid,}}{Boussinesq equations} \end{split}$$

$$(\psi, v, \rho) = (\psi_0, v_0, \rho_0) + \epsilon(\psi_1, v_1, \rho_1) + \epsilon^2(\psi_2, v_2, \rho_2) + \dots,$$



Weakly nonlinear solutions

 $\psi_2 = \sum_{m=-\infty}^{\infty} \sum_{n=-\infty}^{\infty} \left[h_{mn}(z) \cos\left((k_m + k_n)x - 2\omega t + \alpha_m + \alpha_n\right) + g_{mn}(z) \cos\left((k_m - k_n)x + \alpha_m - \alpha_n\right) \right]$

Weakly nonlinear solution diverges if

Uniform Stratification

$$\frac{\omega^2}{N_0^2} = \frac{(m+n)^2 - (m-n)^2 (4 - f^2/\omega^2)/(1 - f^2/\omega^2)}{4(m+n)^2 - (m-n)^2 (4 - f^2/\omega^2)/(1 - f^2/\omega^2)}$$

 $(m/3) < n < 3m, \ m \neq n$

Non-dimensional parameters

1. ω/N_0

2. f/ω

3. *m/n*

> High-mode interactions at nearinertial frequencies result in strong low-mode superharmonics





> Nonuniform stratification increases the number of triadic resonances

> Appearance of several resonances at the near-inertial frequencies

Nonuniform stratifications – Distribution of modes

Presence of several modes – triadic resonance at near-inertial frequency highly likely

m²s⁻²/cph

m²s⁻²/cph

2015 WHOI Mooring –
 Thanks to *Dipanjan Chaudhuri* for the figure

Energy at 2f may be present at subsurface depths also



Ongoing work

Rate of energy transfer to superharmonics (with Dheeraj Varma, IIT Madras)

> Numerical Simulations (with Vamsi Chalamalla, IIT Delhi)

Field Data (with Dipanjan Chaudhuri, IISc)

Dheeraj Varma & Manikandan Mathur, J. Fluid Mech. 2017