

# Quantum Hamiltonian Engineering via Parametric Drives

Michael Hatridge

Department of Physics and Astronomy, University of Pittsburgh

ICTS NHP 2018



# What a quantum machine looks like



Classical control lines

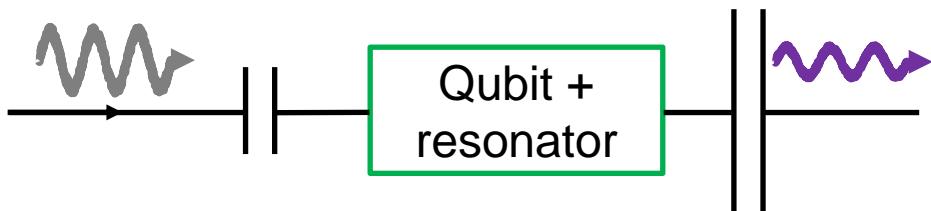
Dilution unit (cryogenic cooling)

50 or so readout channels

50 or so quantum bits

IBM

# Quantum signals from qubits



We might care to

- amplify this pulse
- route it to another qubit/cavity
- Interfere it with another pulse
- ....

**Key limitation: quantum efficiency**

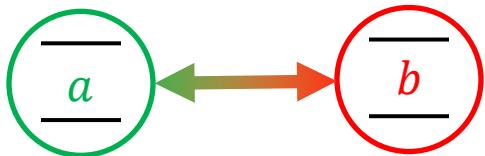
**Solution: parametric amplification**

## Outline

- **Parametric amplification**
- **Effect of 4<sup>th</sup> order terms**
- **Multiple parametric drives**

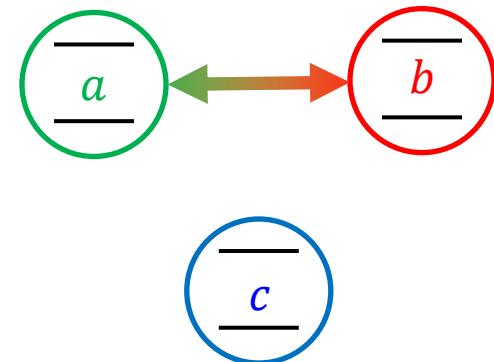
# Parametrically driven couplings

direct exchange



$$\frac{\mathcal{H}_{int}}{\hbar} = g(a b^\dagger + a^\dagger b)$$

parametrically driven  
exchange (Conv)



- if  $\omega_a - \omega_b \gg \kappa_{a,b}$  this term dies due to energy conservation (RWA)
- interaction also turns off slowly vs. detuning, limiting the on/off ratio

$$\frac{\mathcal{H}_{int}}{\hbar} = g(a b^\dagger c^\dagger + a^\dagger b c)$$

- if  $\omega_c \neq \omega_a - \omega_b$  we can drive the  $c$ -mode ‘stiffly’

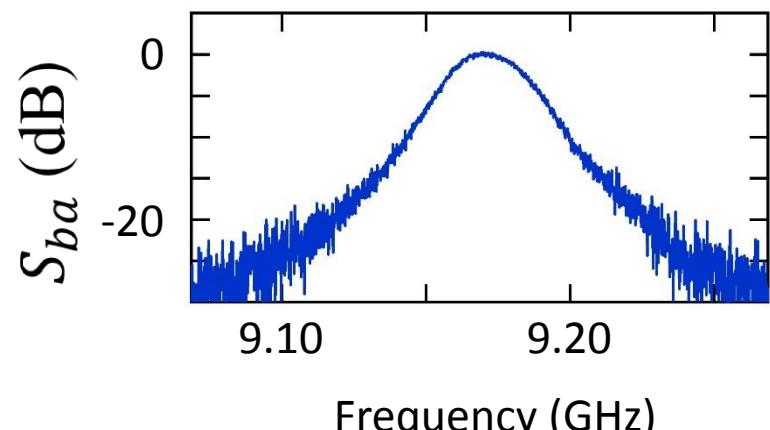
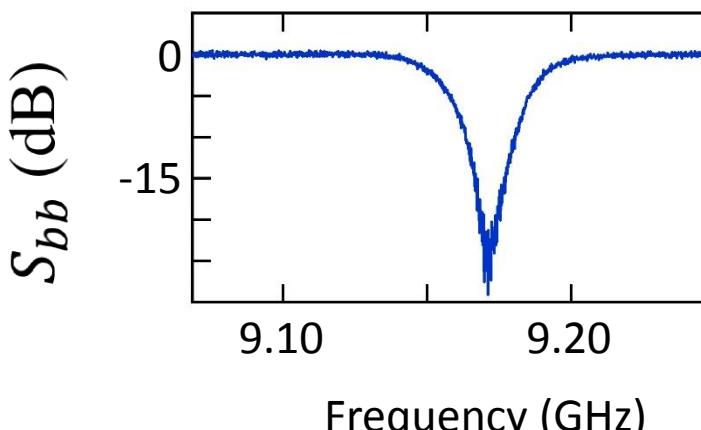
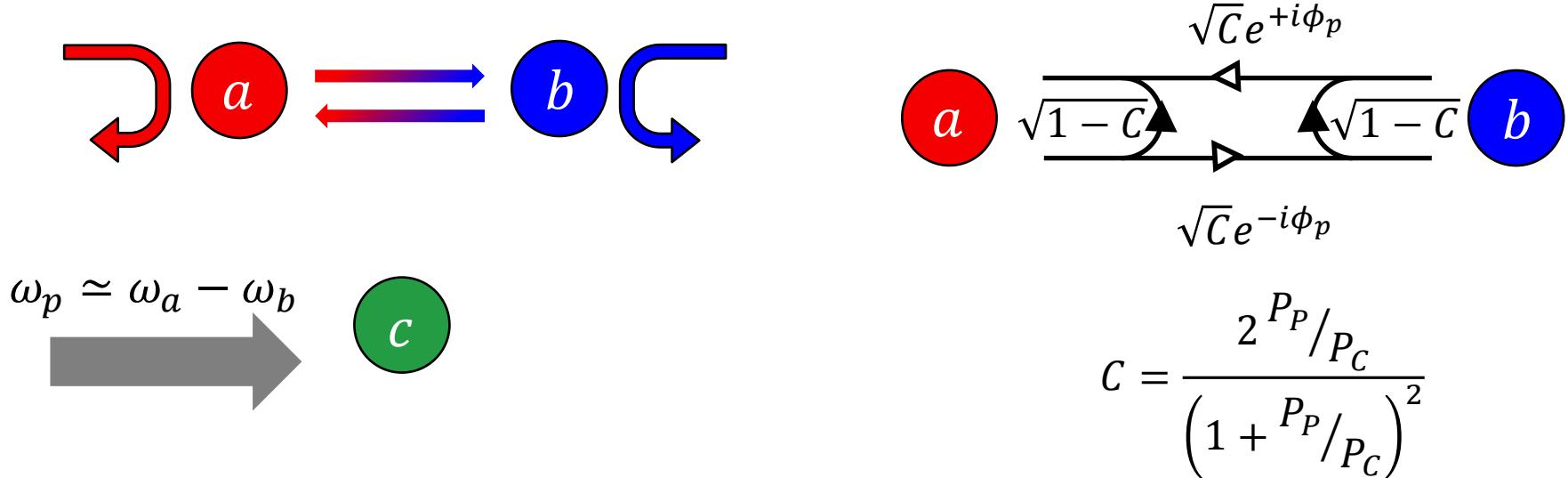
$$\frac{\mathcal{H}_{int}}{\hbar} = g a b^\dagger + g^* a^\dagger b$$

- Parametric drive fully controls strength and phase of interaction

# Photon conversion (Conv)

$$\omega_p = \omega_a - \omega_b \neq \omega_c$$

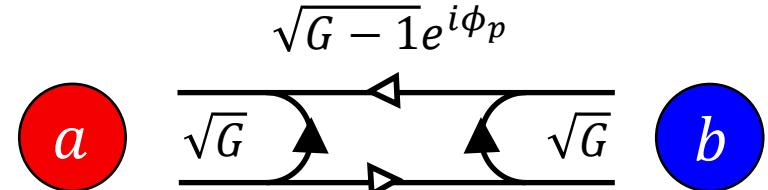
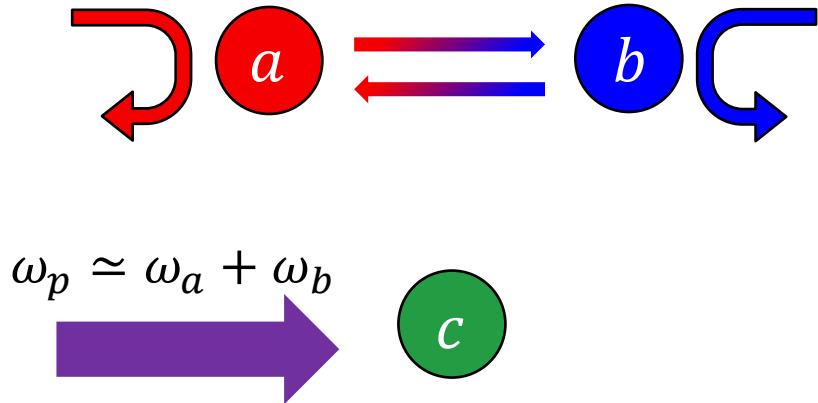
$$H_G = \hbar g(a^\dagger b e^{i\phi_p} + a b^\dagger e^{-i\phi_p})$$



# Phase preserving gain (Gain)

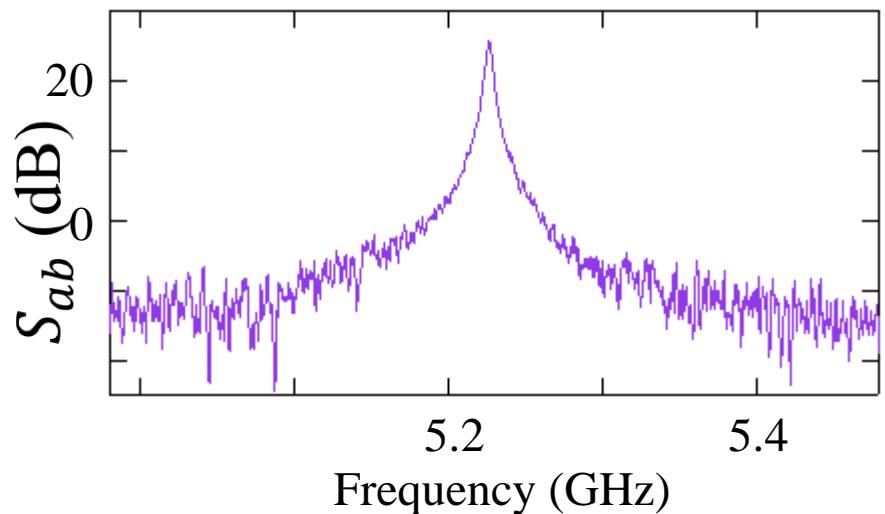
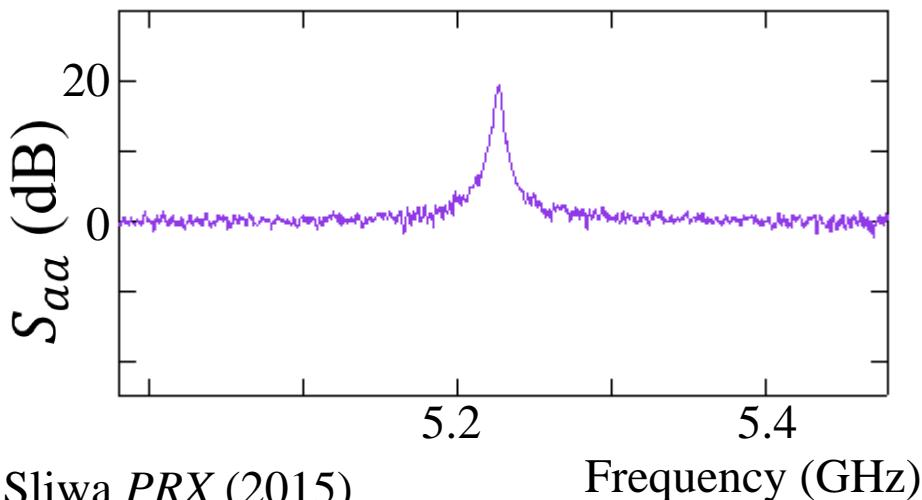
$$\omega_p = \omega_a + \omega_b \neq \omega_c$$

$$H_G = \hbar g(a^\dagger b^\dagger e^{i\phi_p} + ab e^{-i\phi_p})$$

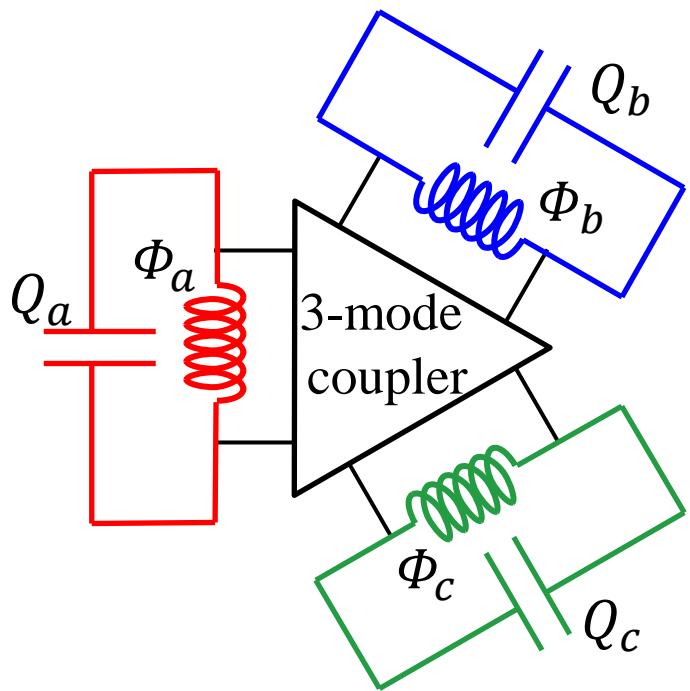


$$\sqrt{G-1}e^{-i\phi_p}$$

$$G = \frac{\left(1 + P_P/P_C\right)^2}{\left(1 - P_P/P_C\right)^2}$$



# Parametric coupling overview



$$H_{couple} \propto \Phi_a \Phi_b \Phi_c$$

Re-write in terms of  $a, b, c$

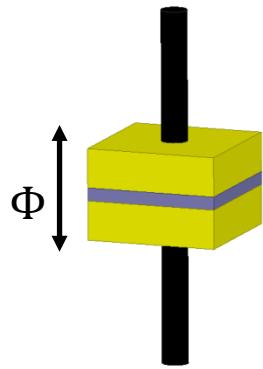
$$\begin{aligned} H_{couple} &= \hbar g_3 (a + a^\dagger)(b + b^\dagger)(c + c^\dagger) \\ &= \hbar g_3 (abc^\dagger + a^\dagger b^\dagger c + ab^\dagger c + a^\dagger bc^\dagger + \dots) \end{aligned}$$

- If frequencies  $\omega_{a,b,c}$  all very different, all terms die in the rotating wave approx.
- Drive one mode ( $c$ ) at  $\omega_p = \omega_a + \omega_b \neq \omega_c$
- Stiff pump:  $c \rightarrow \langle c \rangle = |c| e^{i\phi_p}$

$$H_G = \hbar g (a^\dagger b^\dagger e^{i\phi_p} + a b e^{-i\phi_p})$$

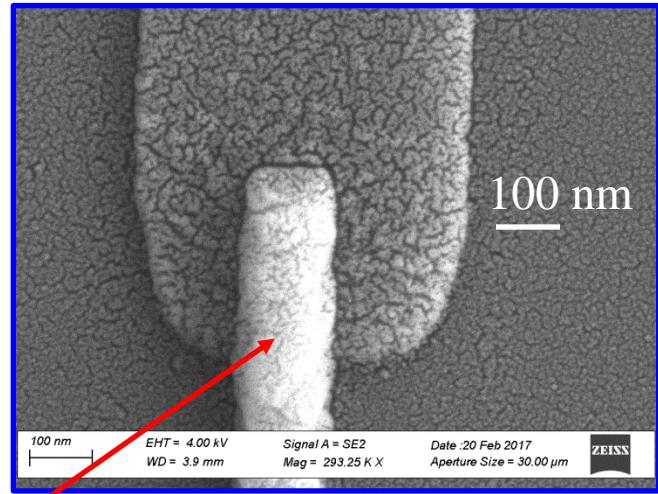
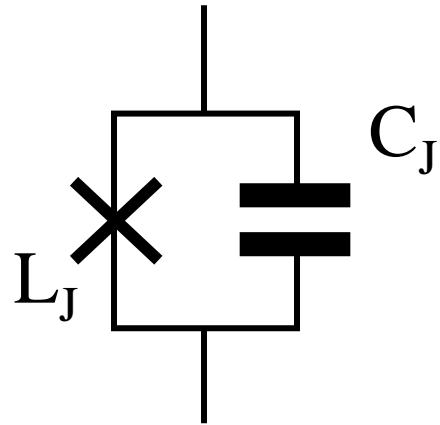
Physical Implementations: Josephson junctions, opto-mechanics, diodes, optical fibers....

# The Josephson tunnel junction

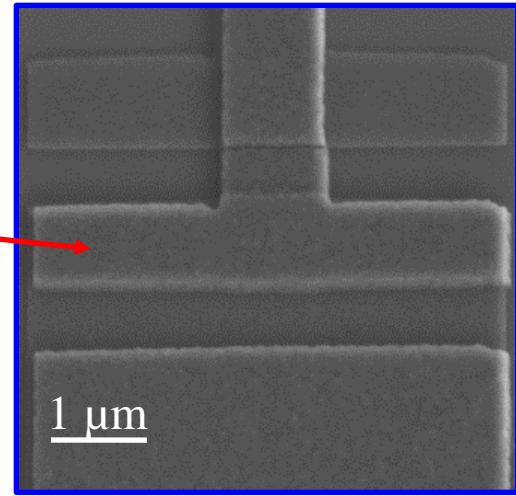


SUPERCONDUCTING  
TUNNEL JUNCTION

$$I = I_0 \sin\left(\frac{2\pi}{\Phi_0}\Phi\right)$$

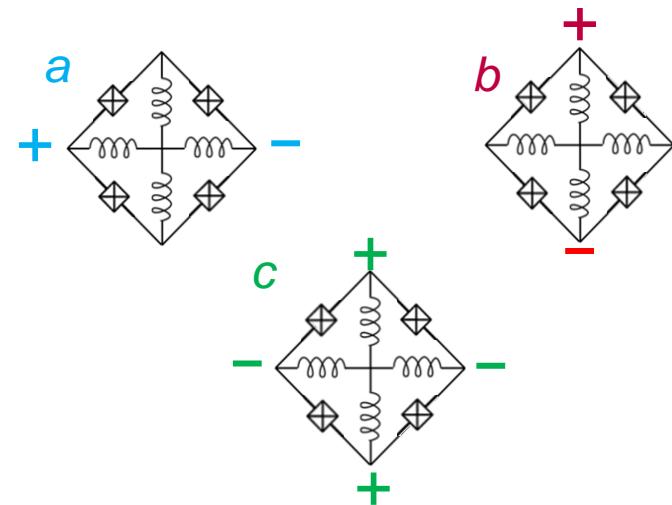
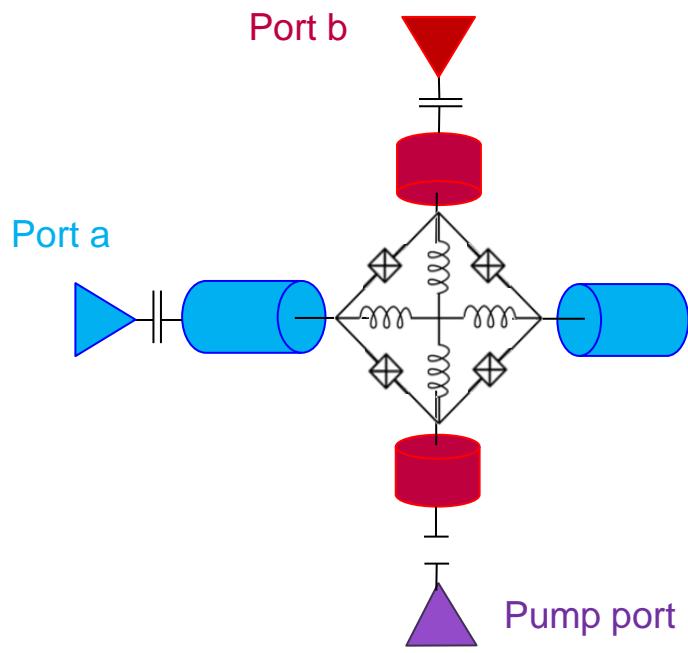


Al/AlO<sub>x</sub>/Al  
tunnel junction  
 $T \sim 20$  mK

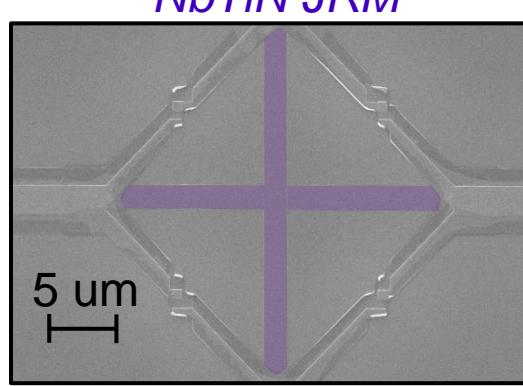


$$H = \frac{Q^2}{2C} - E_J \cos\left(\frac{2\pi}{\Phi_0}\Phi\right) = \hbar\omega_0 b^\dagger b - \lambda(b^\dagger b)^2 + \dots$$

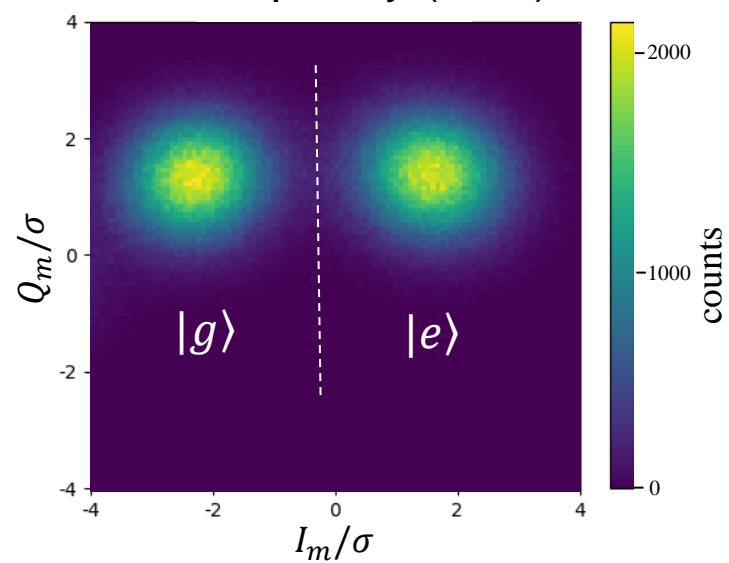
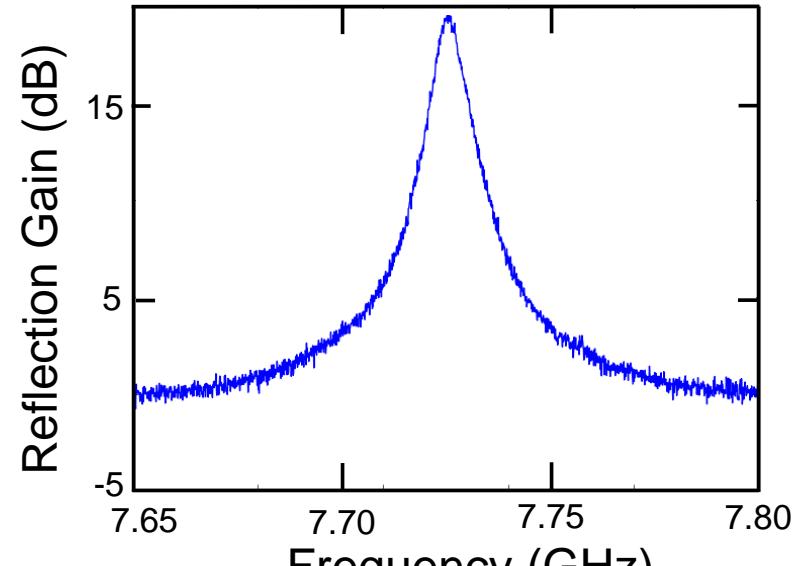
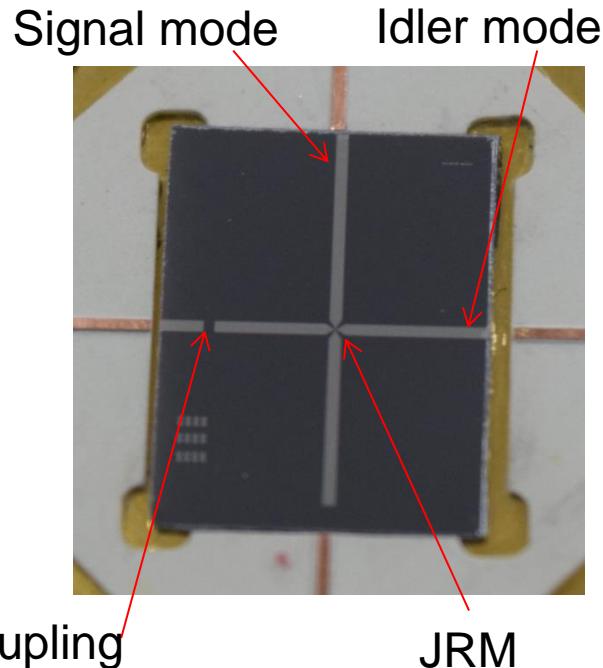
# Josephson parametric converter (JPC)



$$\frac{H_{Coupling}}{\hbar} = g (a + a^\dagger)(b + b^\dagger)(c + c^\dagger) - \sum_{m=a}^c \sum_{n=m}^c K_{mn} a_m^\dagger a_m a_n^\dagger a_n + O(a_i^5) + \dots$$



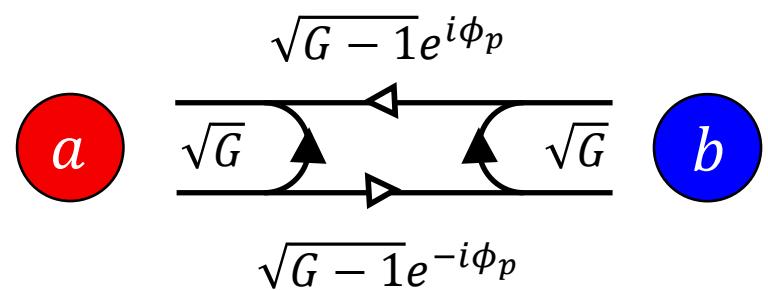
# The 8-junction Josephson Parametric Converter



Nearly quantum limited ( $\eta \sim 0.5 - 0.6$ )!

# Amplifier Limitations

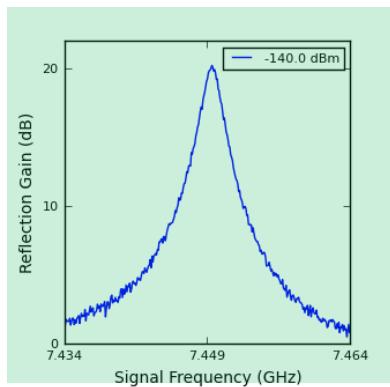
1. Operates in reflection



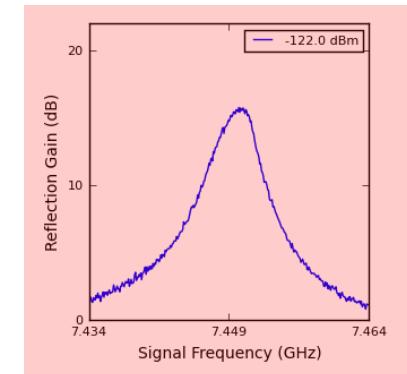
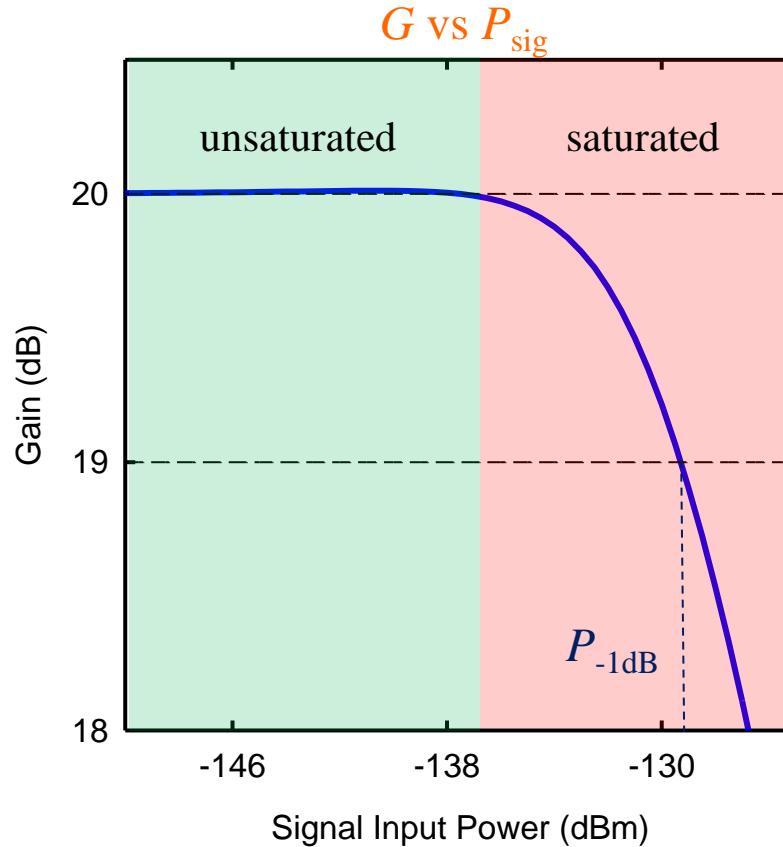
2. Has narrow bandwidth

$$2\pi B \simeq \frac{\sqrt{\kappa_a \kappa_b}}{\sqrt{G}}$$

# Limitation 3: Gain Saturation



$$G_0 = \left( \frac{1 + n_p/n_p^c}{1 - n_p/n_p^c} \right)^2$$



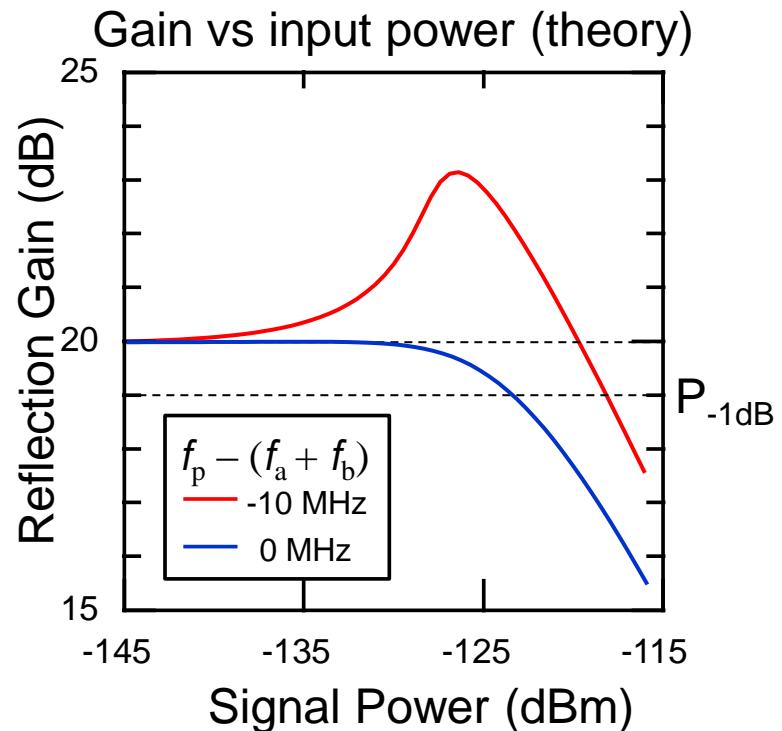
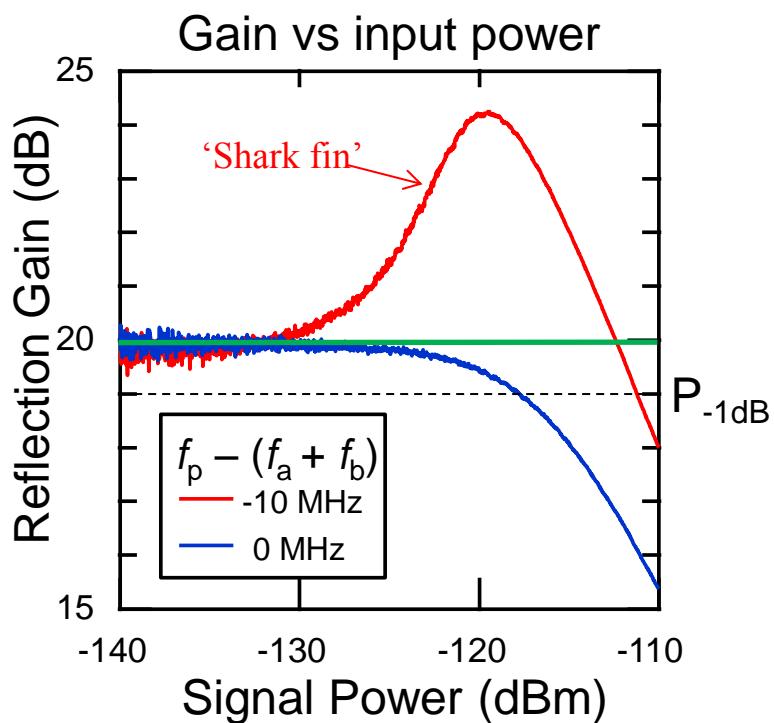
- Pump depletion (PD) theory:
- $P_{sig}$  ↑,  $n_p$  ↓,  $G$  ↓
  - $n_p$  ↑,  $P_{-1dB}$  ↑

Problem: We don't understand well all causes

# Cancelling Higher Order Effects

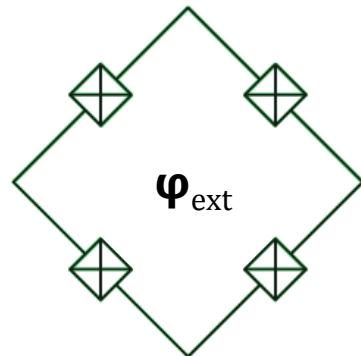
# Kerr terms effect performance

$$H_{coupling} = g(ab + a^\dagger b^\dagger) \\ + \kappa_{aa}(a^\dagger a)^2 + \kappa_{bb}(b^\dagger b)^2 + \kappa_{ab}a^\dagger ab^\dagger b$$

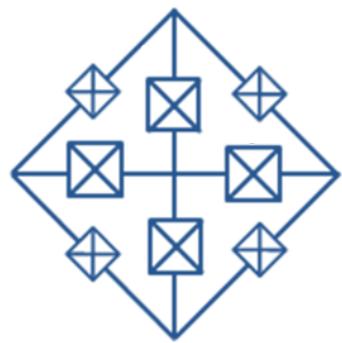


# Using ‘shunted’ JRMs to achieve cancellation

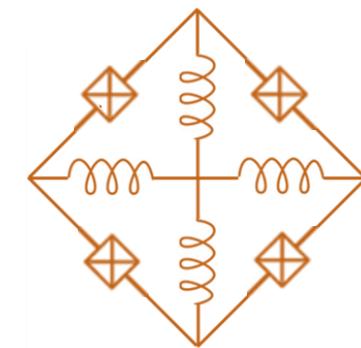
4-Junction JRM



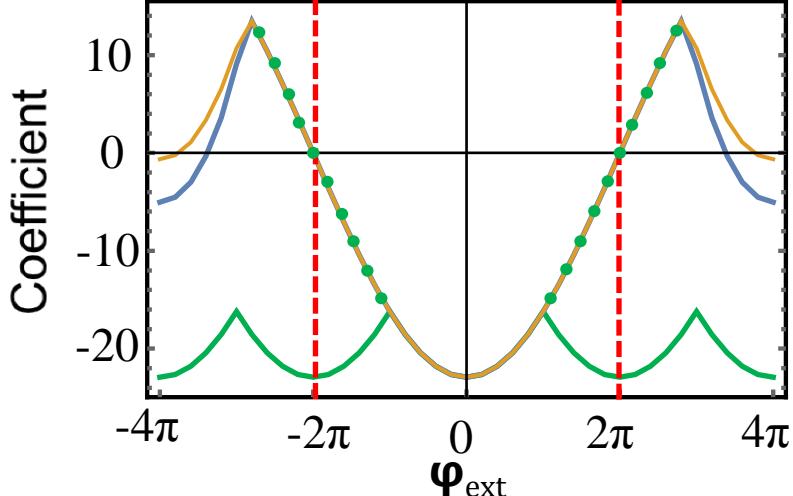
8-Junction JRM



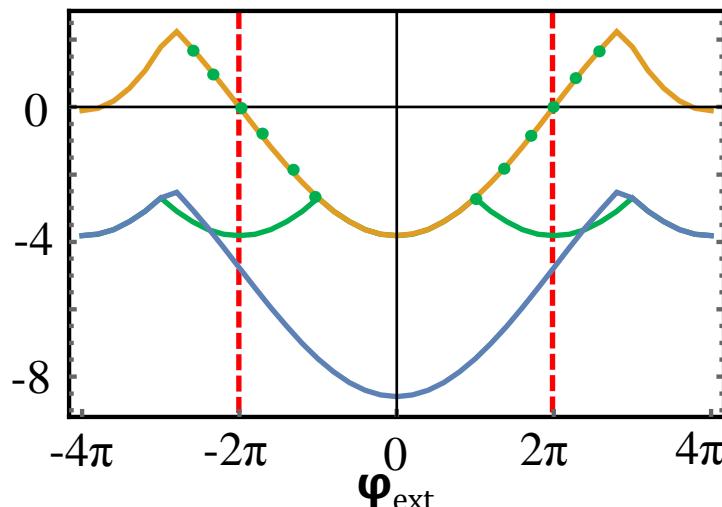
4-Junction JRM + Linear Inductance (WJRM)



Cross Kerr

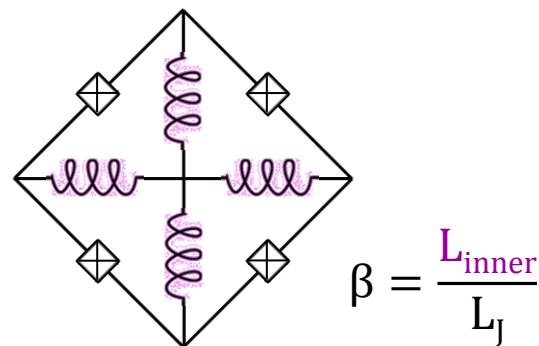


Self Kerr

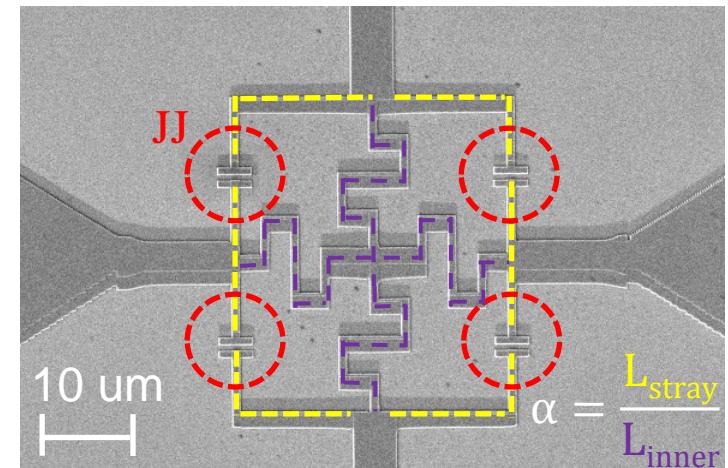


# How strongly to shunt?

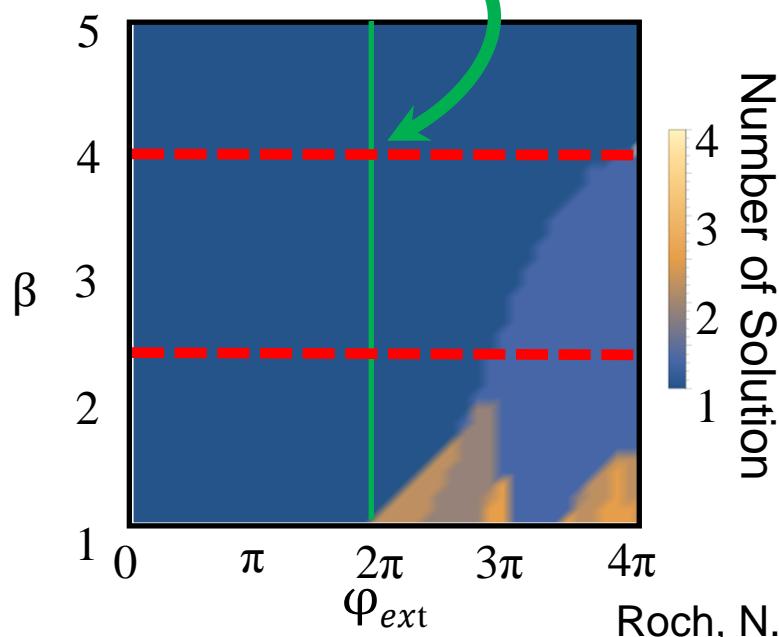
4 Junction JRM + Linear Inductances



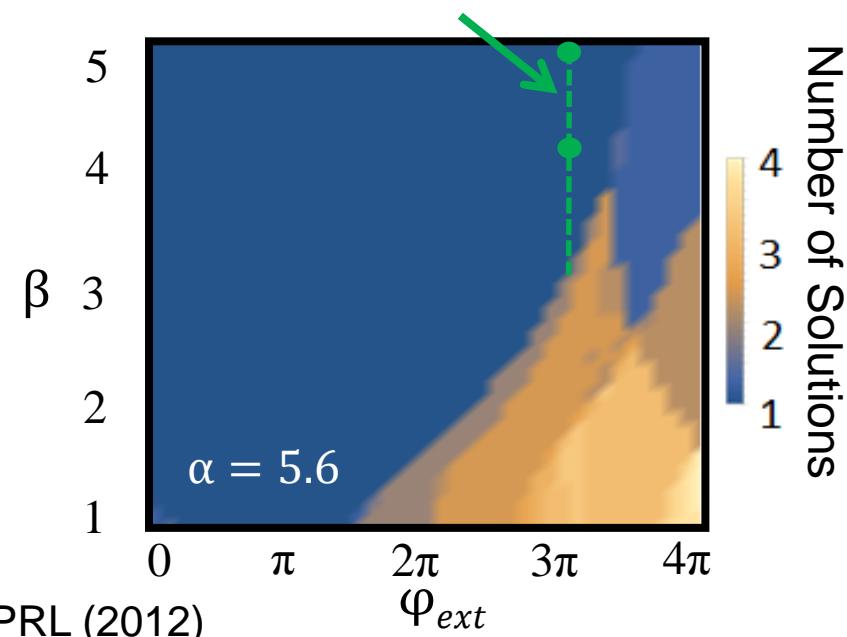
Stray Inductance (yellow)



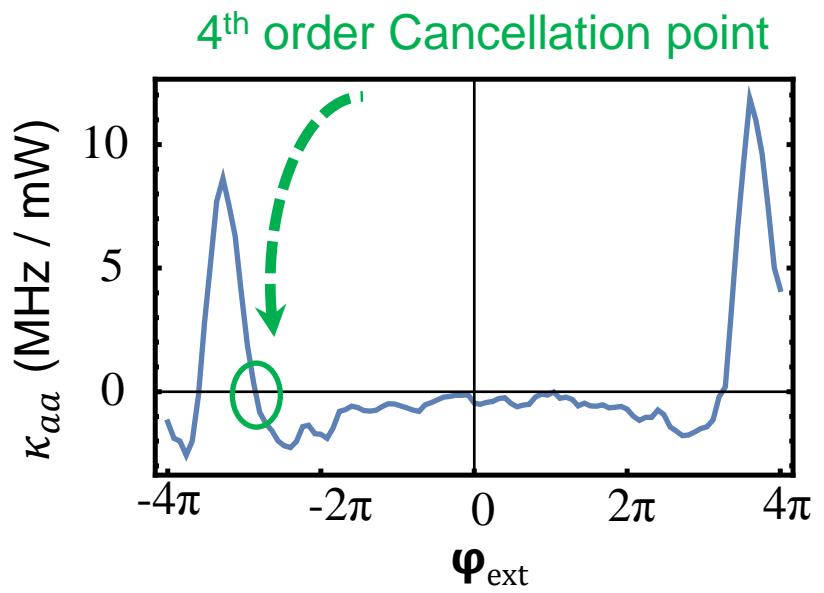
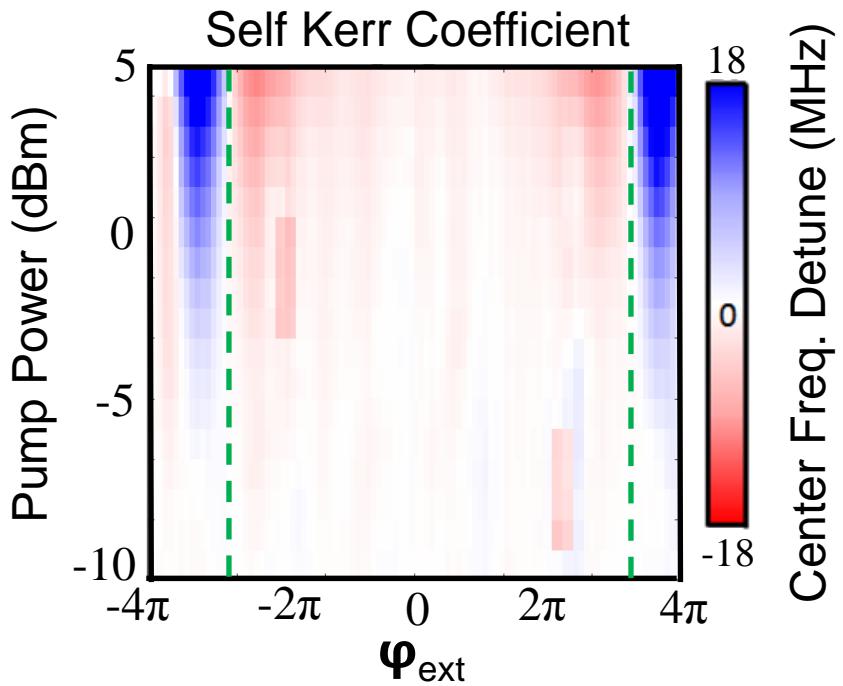
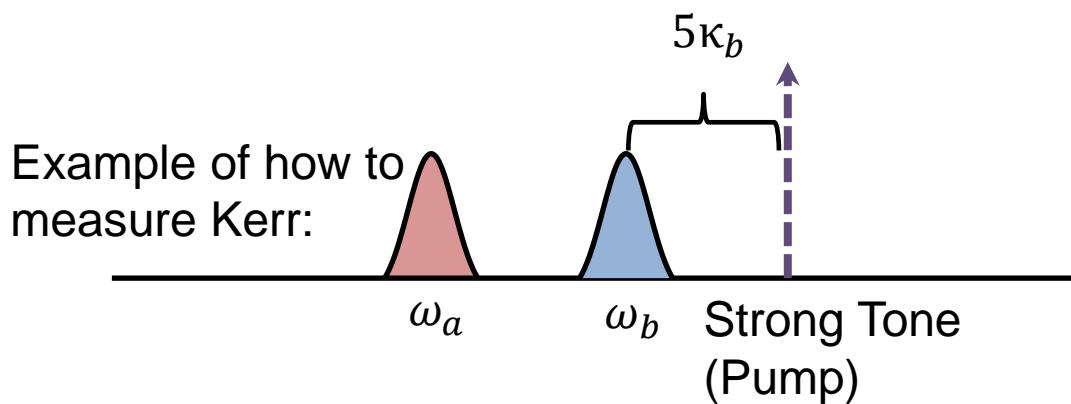
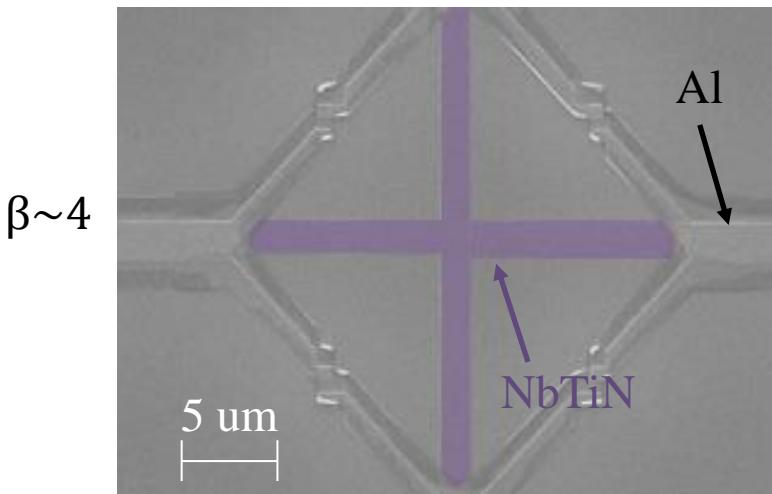
4<sup>th</sup> order Cancellation points



Cancellation Points

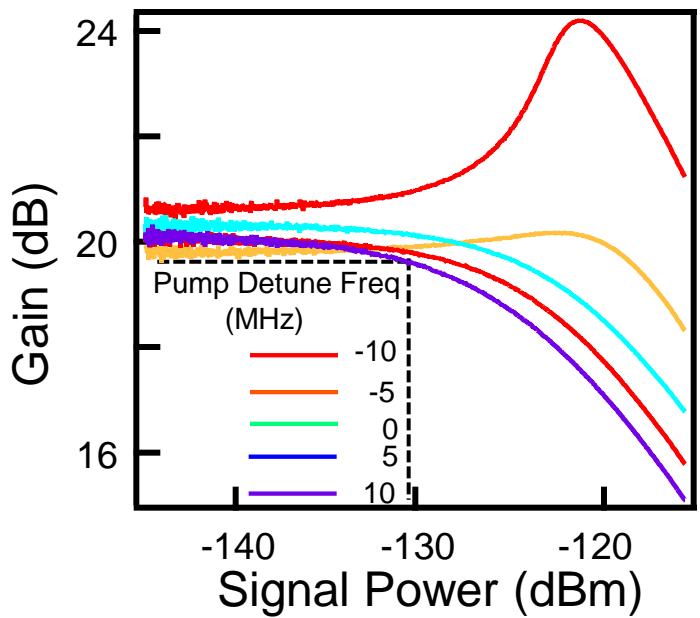
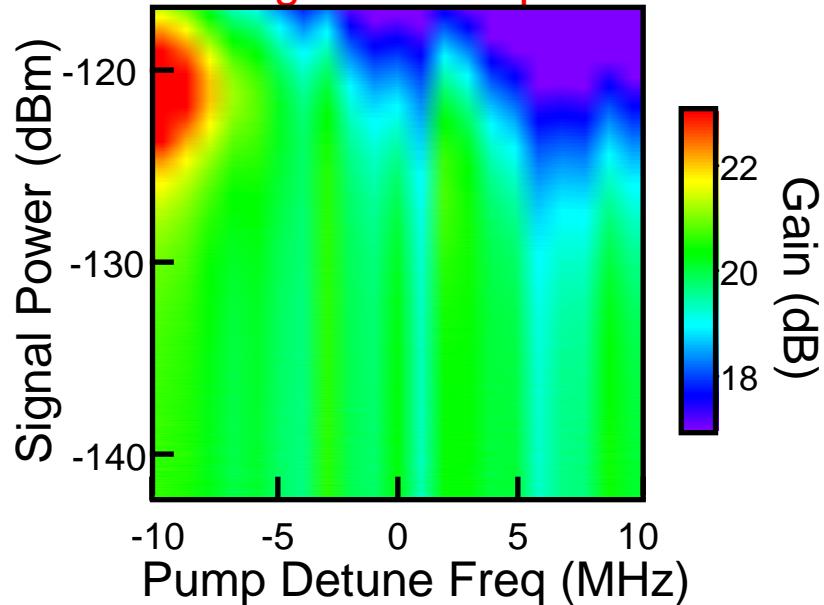


# Kerr measurement of NbTiN shunted JRM

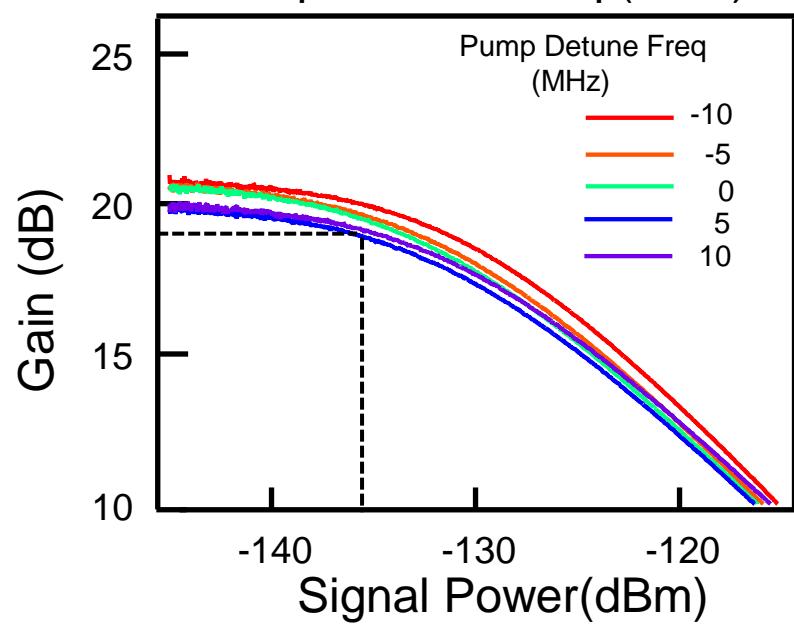
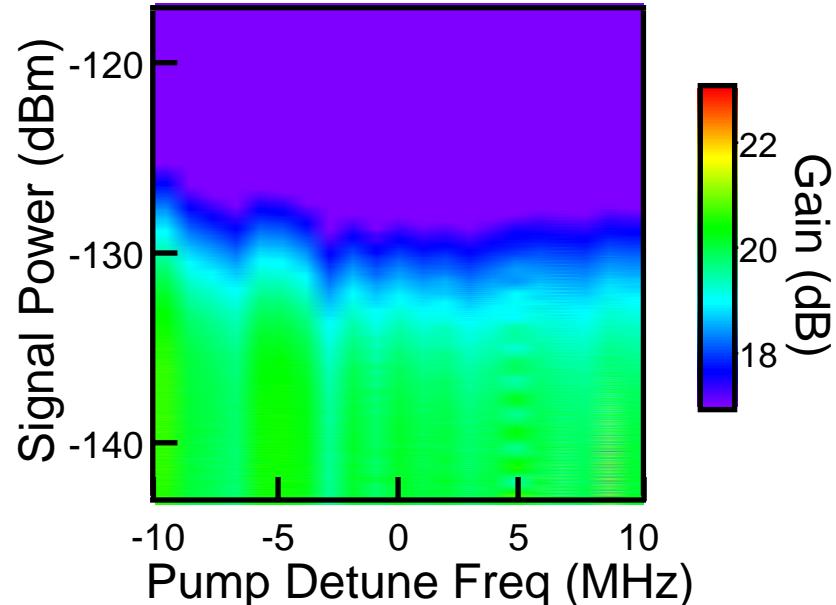


# Cancellation Point Comparison

Negative Kerr point

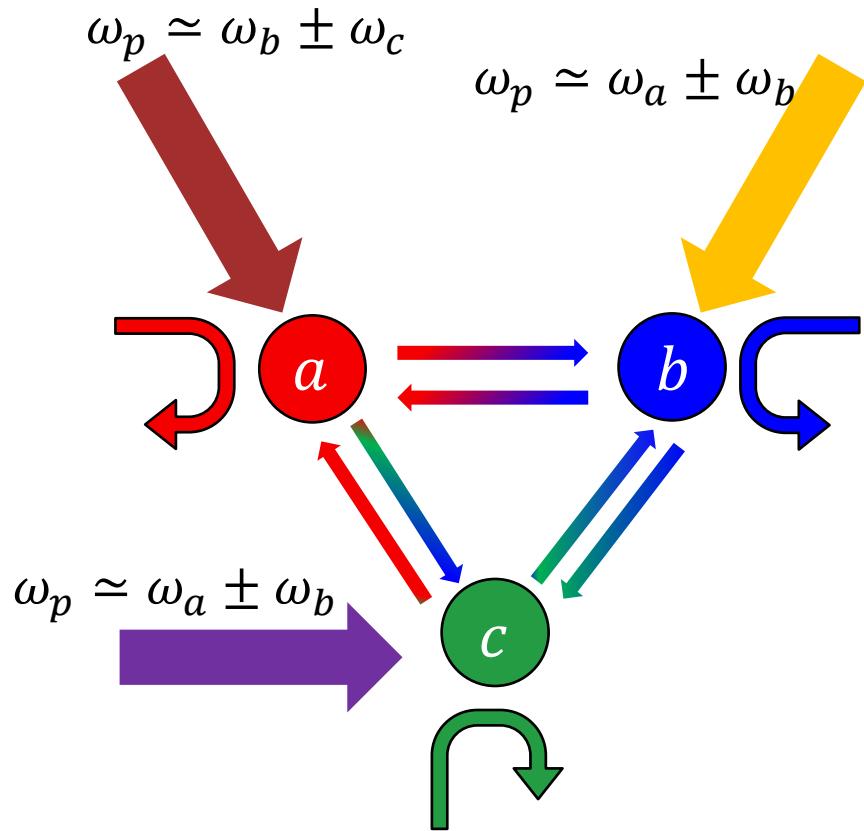


4<sup>th</sup> order cancellation point



# **Multiple Parametric Couplings**

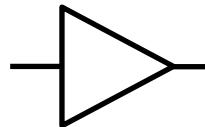
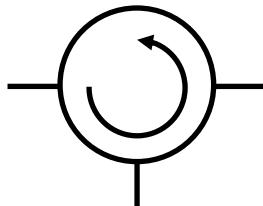
# Adding interference



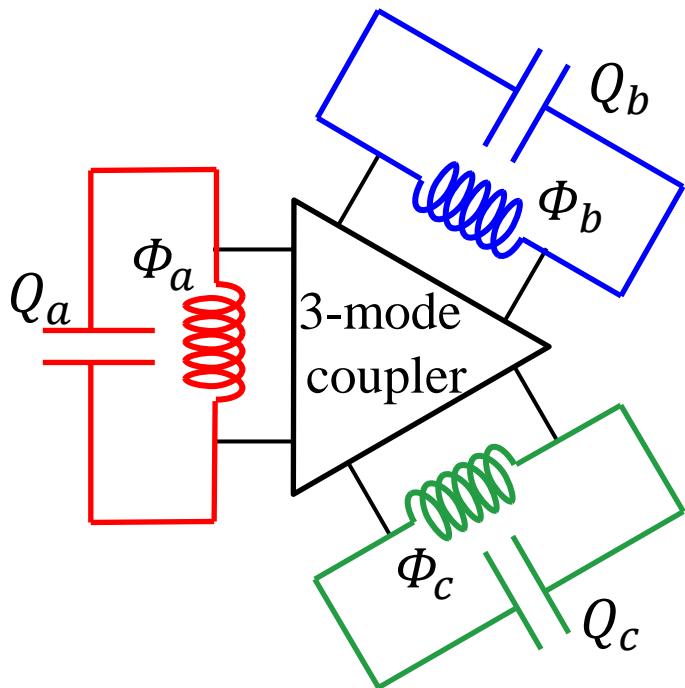
- Pumps span frequencies and imprint phases on each leg
- Controls:
  - Coupling type ( $\omega_p$ )
  - Coupling strength ( $P_p$ )
  - Coupling phase ( $\phi_p$ )
- Frequency (non) conjugation
  - Gain:  $(\omega_a + \Delta) \rightarrow (\omega_b - \Delta)$
  - Conv:  $(\omega_a + \Delta) \rightarrow (\omega_b + \Delta)$

# Engineering quantum information processors

- Passive/reciprocal circuit elements (hybrids, directional couplers) are straightforward
- Gainful and non-reciprocal devices are much harder



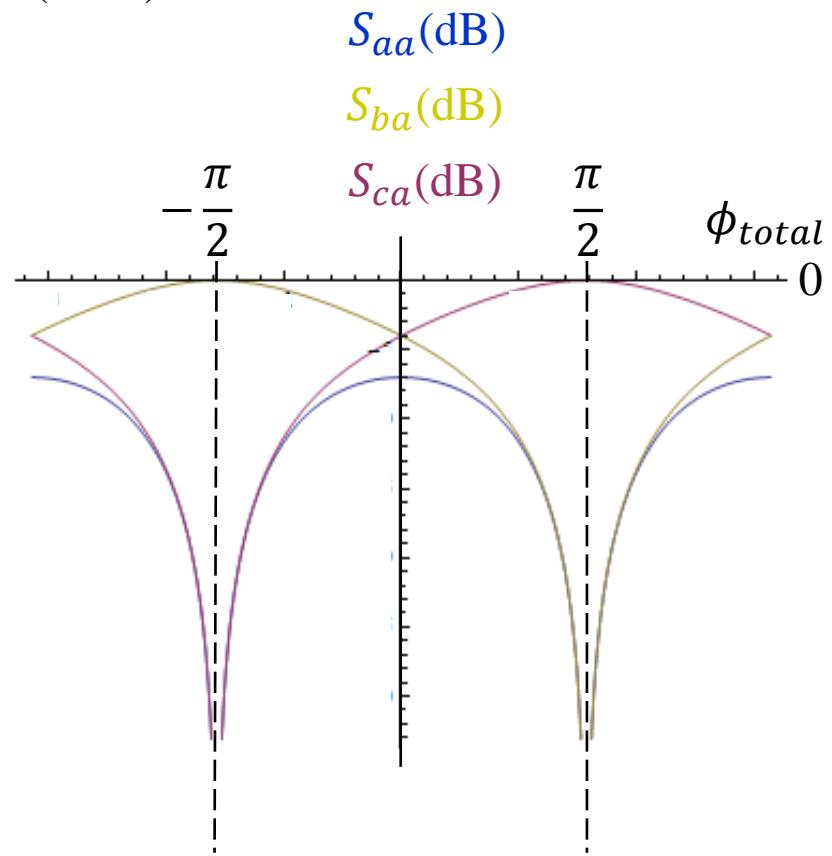
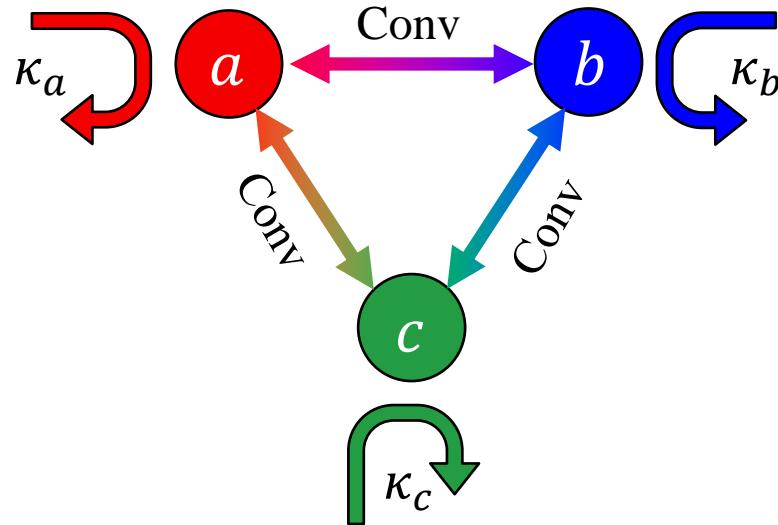
- These *quantum* devices must be lossless, and operate on *propagating* states of microwave light



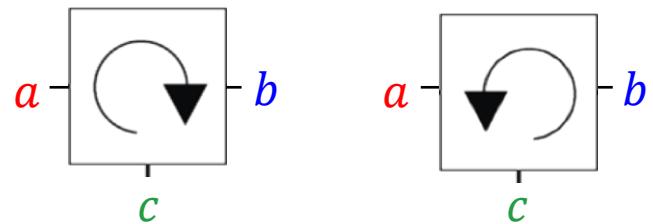
- idea is to engineer these devices from the Hamiltonian up using parametric drives
- key element non-linear Josephson junctions

# 3 port circulator

Sliwa *Phys. Rev. X* (2015)



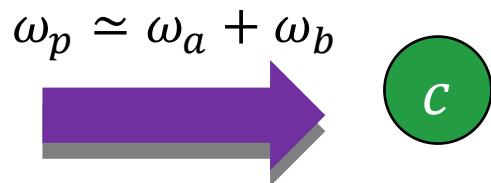
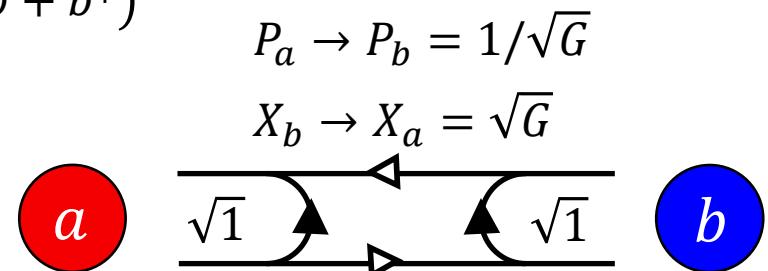
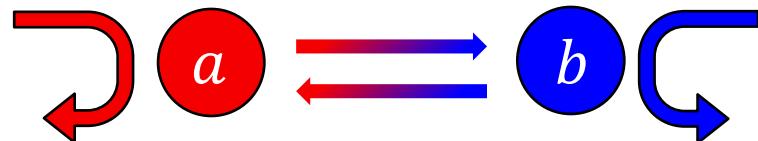
- balance three hopping term amplitudes with dissipation rates
- balance phases so that loop phase  $\phi_{total} = \phi_{ab} + \phi_{bc} - \phi_{ca} = \pm\pi/2$



# Bi-direction phase sensitive gain (XX)

$$\omega_{p1,p2} = \omega_a \pm \omega_b \neq \omega_c$$

$$H_G = \hbar g(a + a^\dagger)(b + b^\dagger)$$

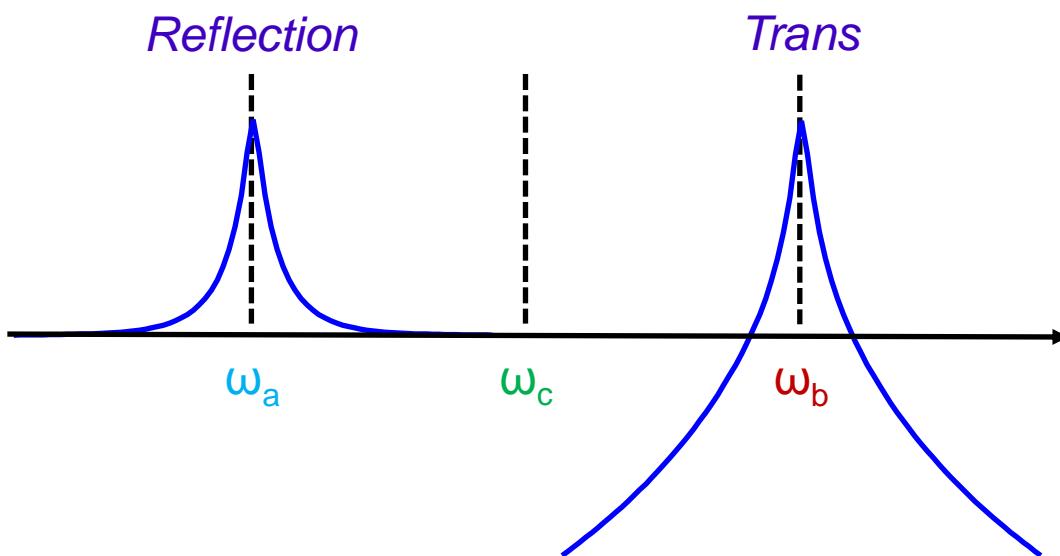
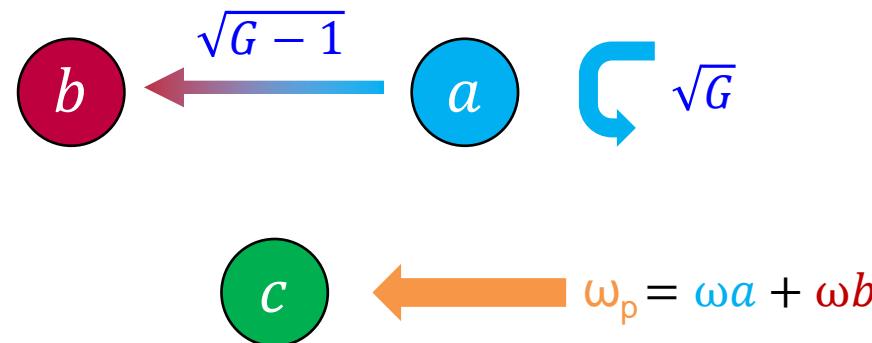


$$\omega_p \simeq \omega_a - \omega_b$$

$$G = P_P / P_C$$

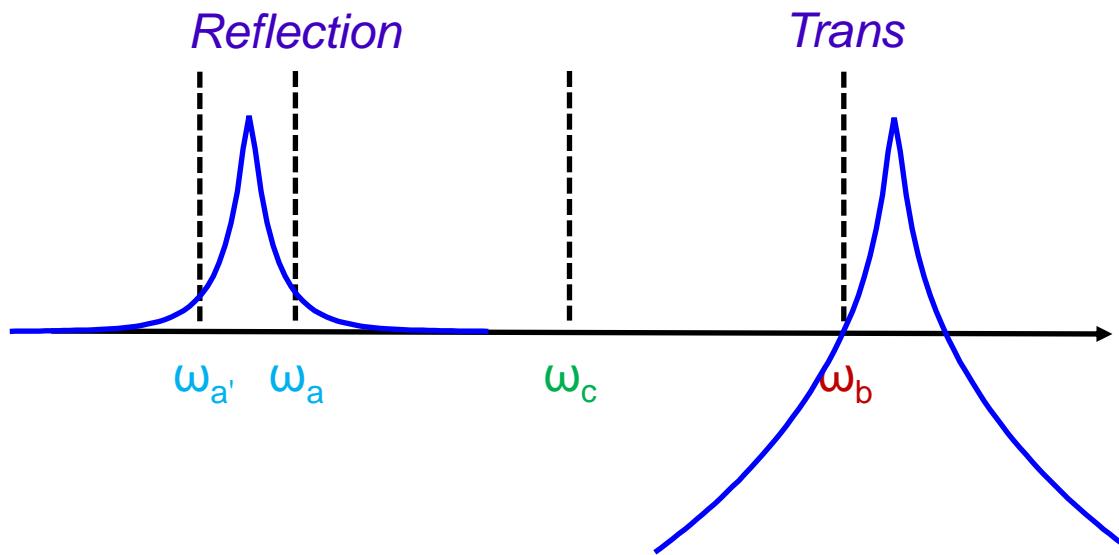
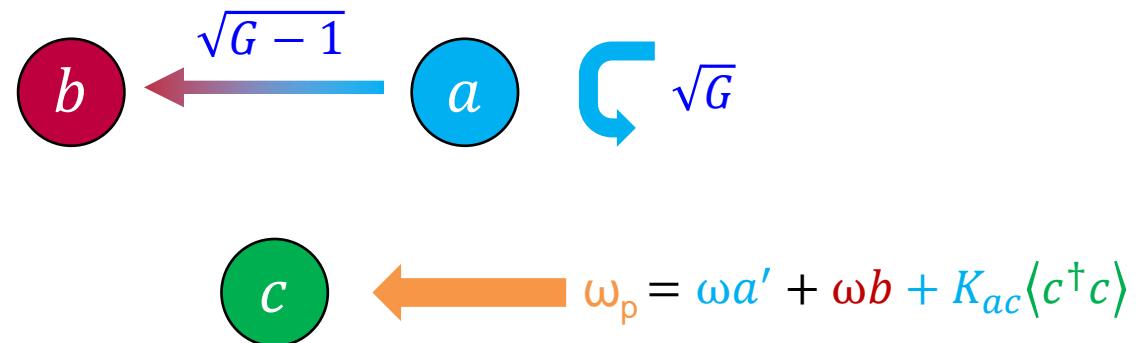
- Enhanced bandwidth  $BW \simeq \frac{1}{2}\kappa/2\pi$  and saturation power!
- Have to match mode bandwidths
- Have to pump HARD!

# Gain / Conversion



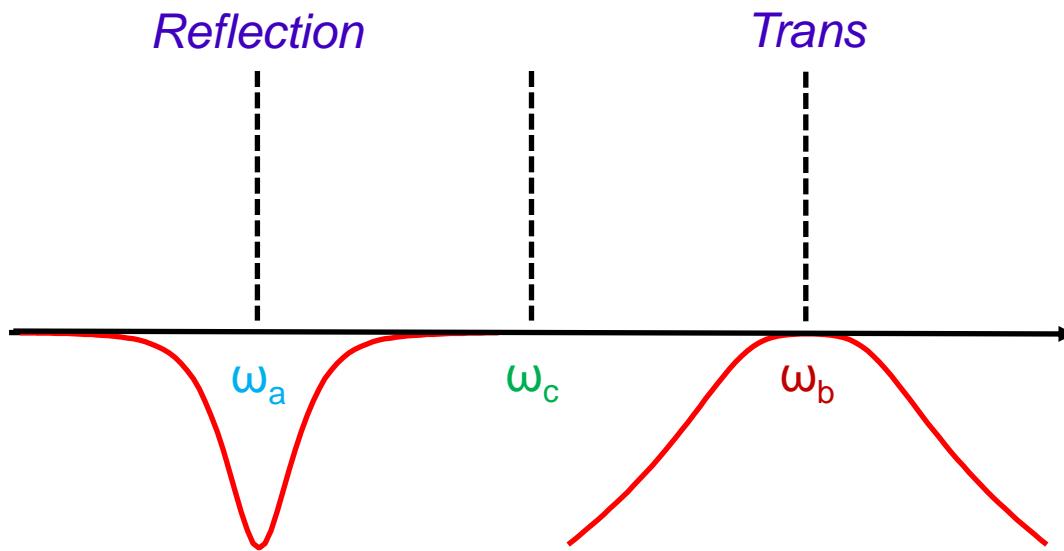
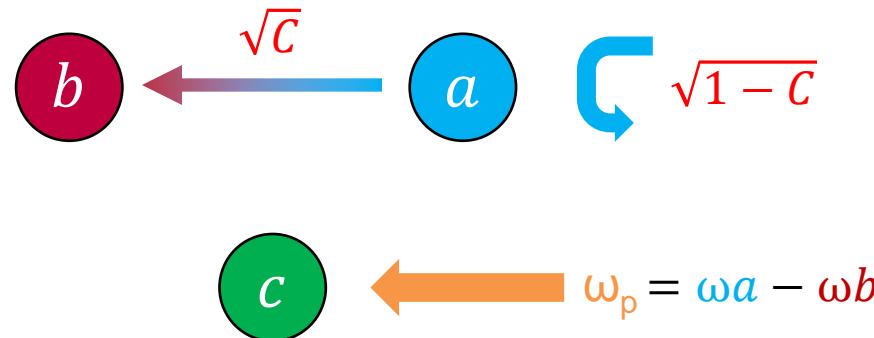
$$\frac{H_{\text{Coupling}}}{\hbar} = g_G(a^\dagger b^\dagger + a b)$$

# Gain / Conversion



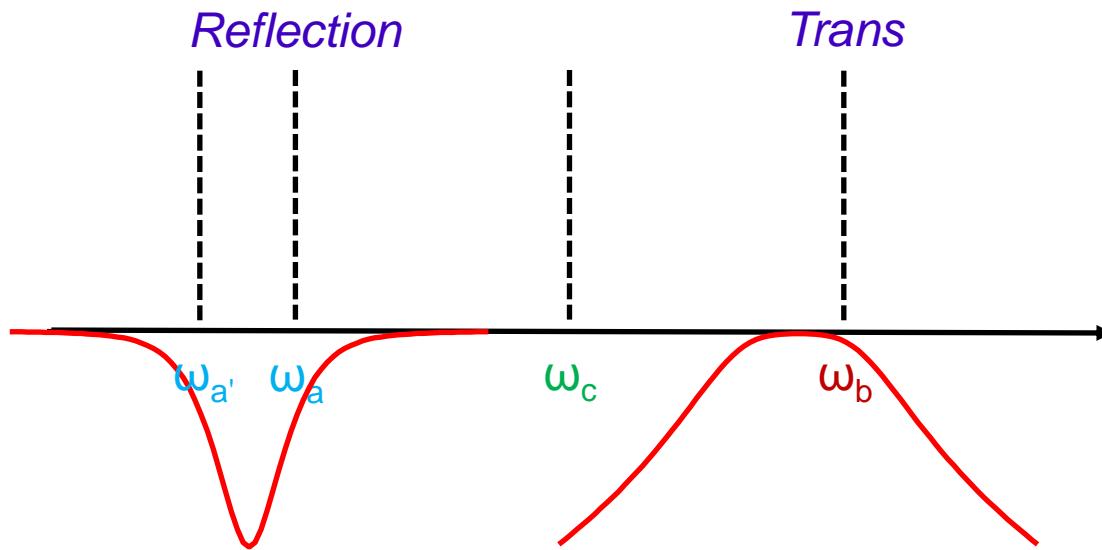
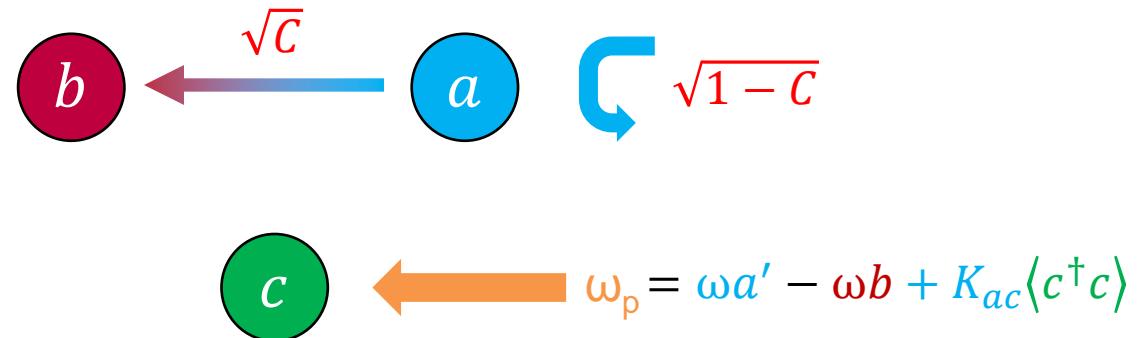
$$\frac{H_{\text{Coupling}}}{\hbar} = g_G (a^\dagger b^\dagger + a b) - K_{ac} a^\dagger a \langle c^\dagger c \rangle$$

# Gain / Conversion



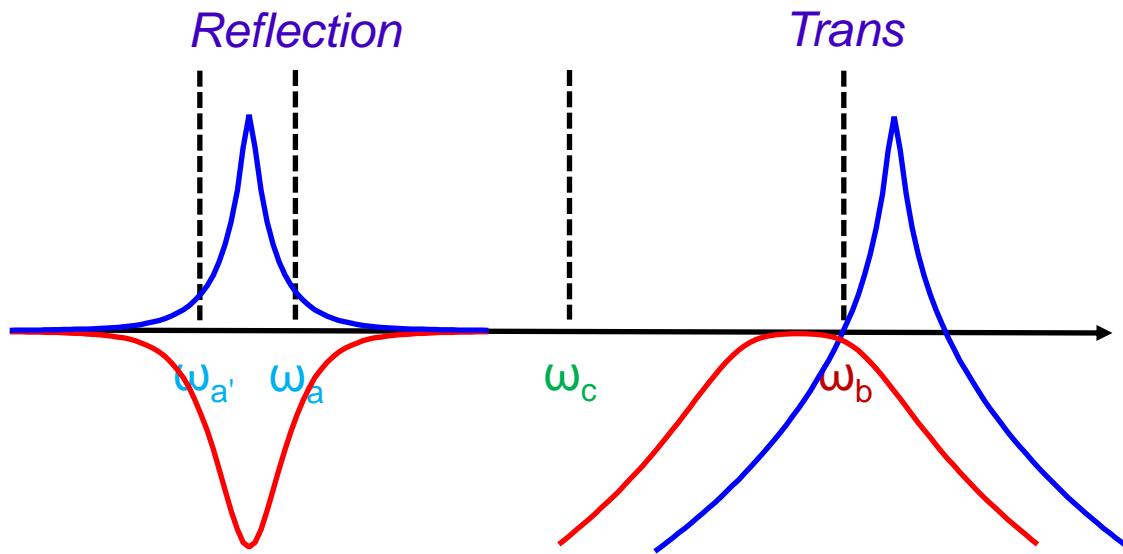
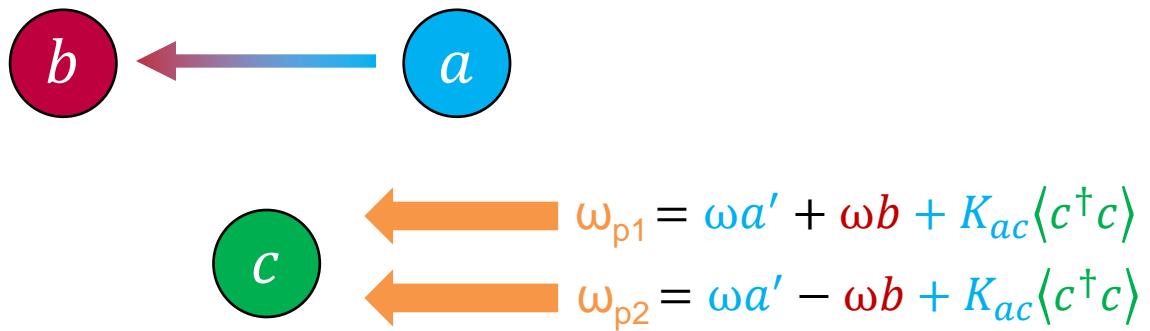
$$\frac{H_{\text{Coupling}}}{\hbar} = g_c(a^\dagger b + a b^\dagger)$$

# Gain / Conversion



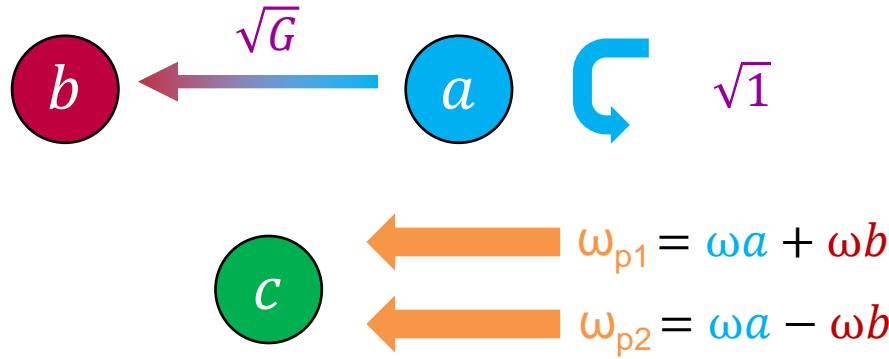
$$\frac{H_{\text{Coupling}}}{\hbar} = g_c (a^\dagger b + a b^\dagger) - K_{ac} a^\dagger a \langle c^\dagger c \rangle$$

# G-C amplification w/ 4<sup>th</sup> order

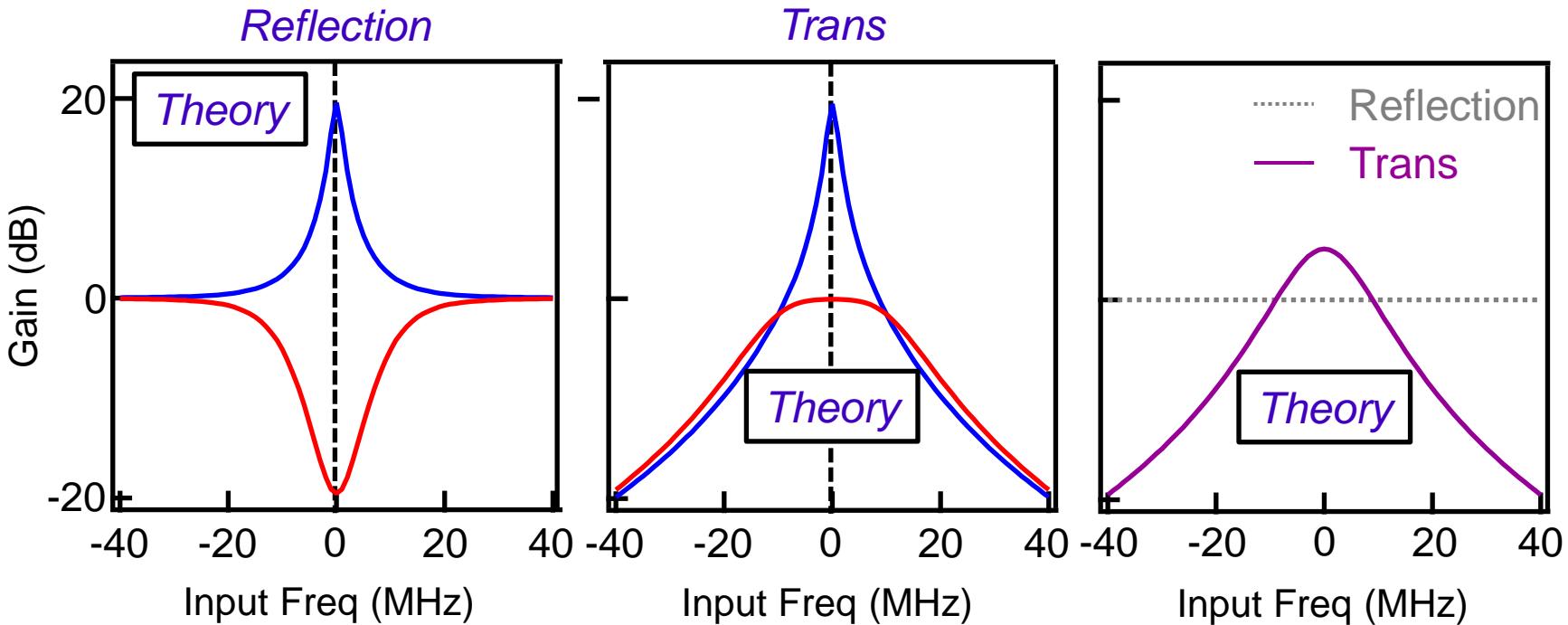


$$\frac{H_{\text{Coupling}}}{\hbar} = g_G (a^\dagger b^\dagger + a b) + g_C (a^\dagger b + a b^\dagger) - K_{ac} a^\dagger a \langle c^\dagger c \rangle$$

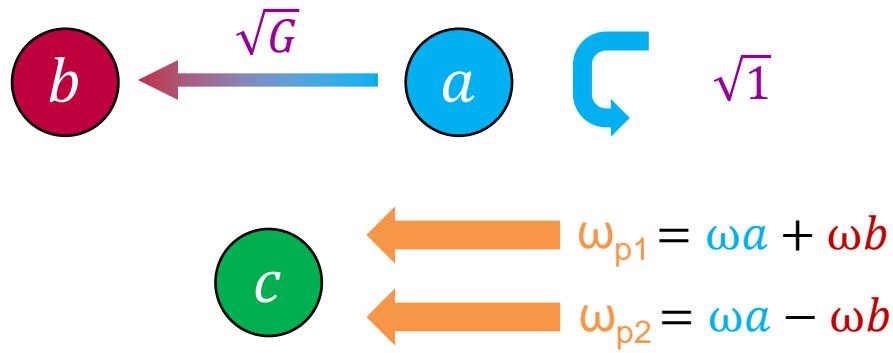
# G-C amplification w/o 4<sup>th</sup> order - Theory



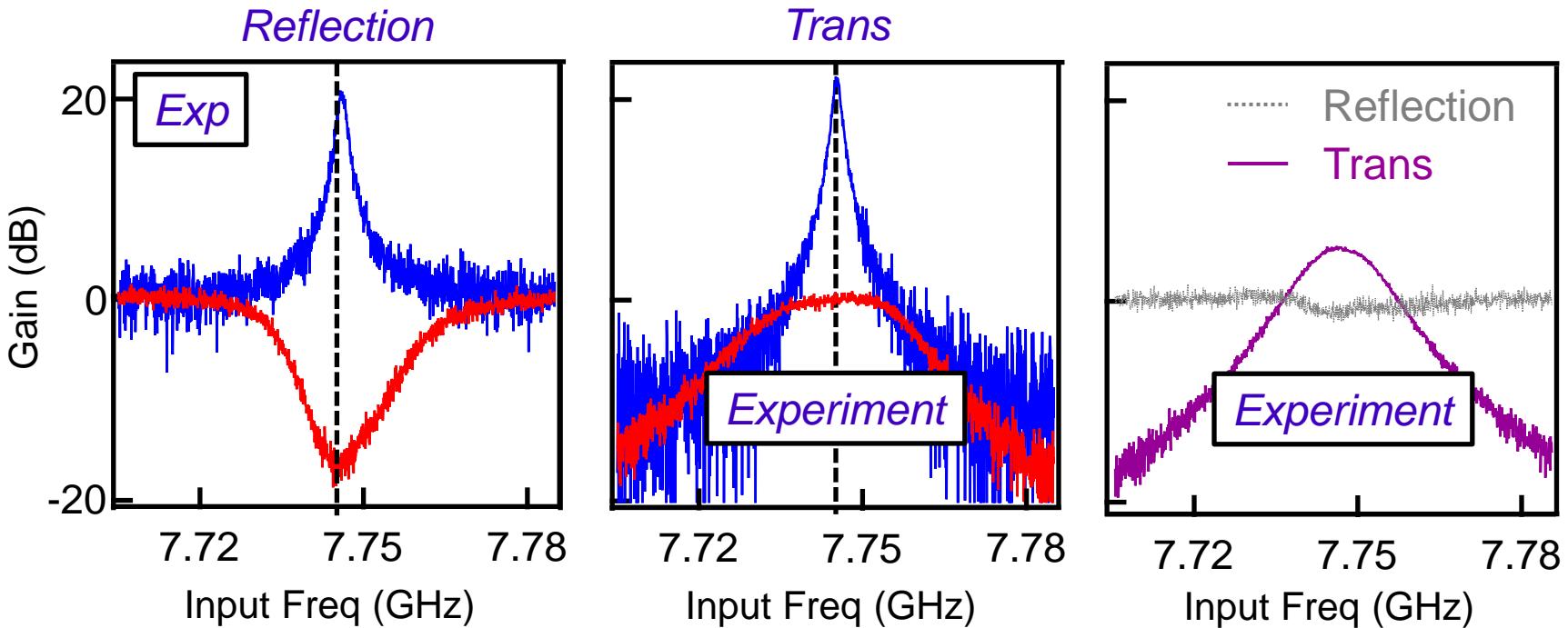
$$\frac{H_{Coupling}}{\hbar} = g_G(a^\dagger b^\dagger + a b) + g_C(a^\dagger b + a b^\dagger)$$



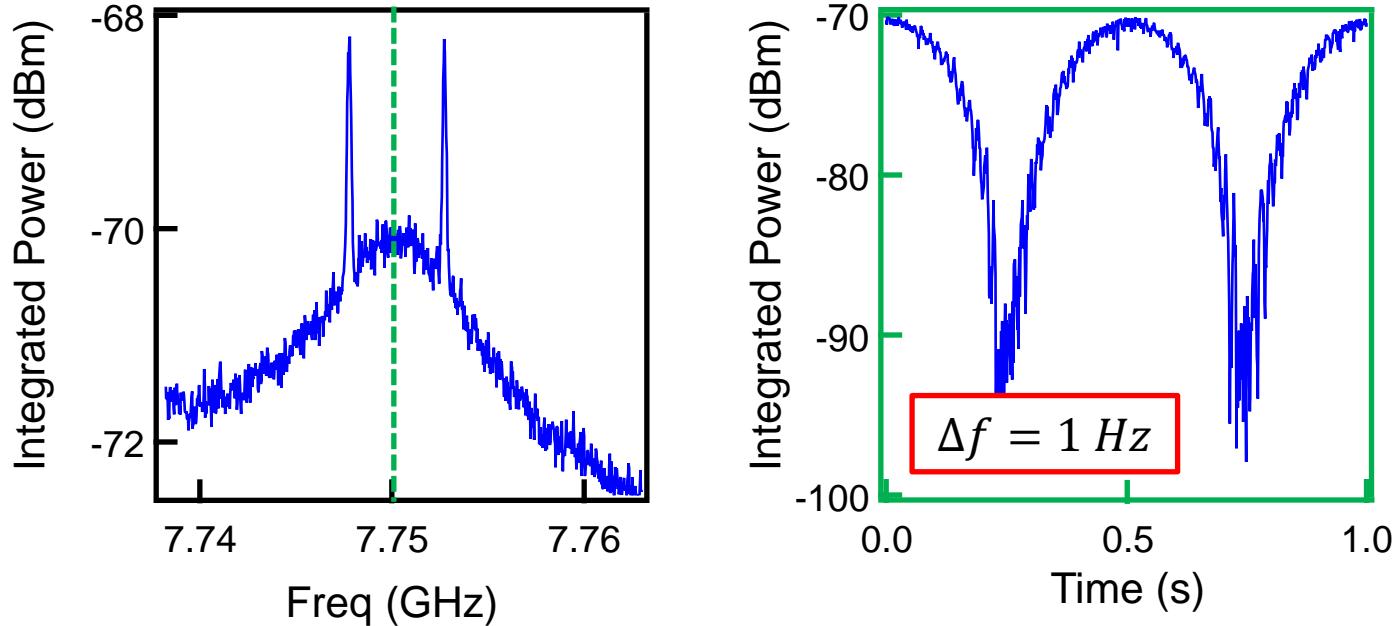
# G-C amplification w/o 4<sup>th</sup> order - Experiment



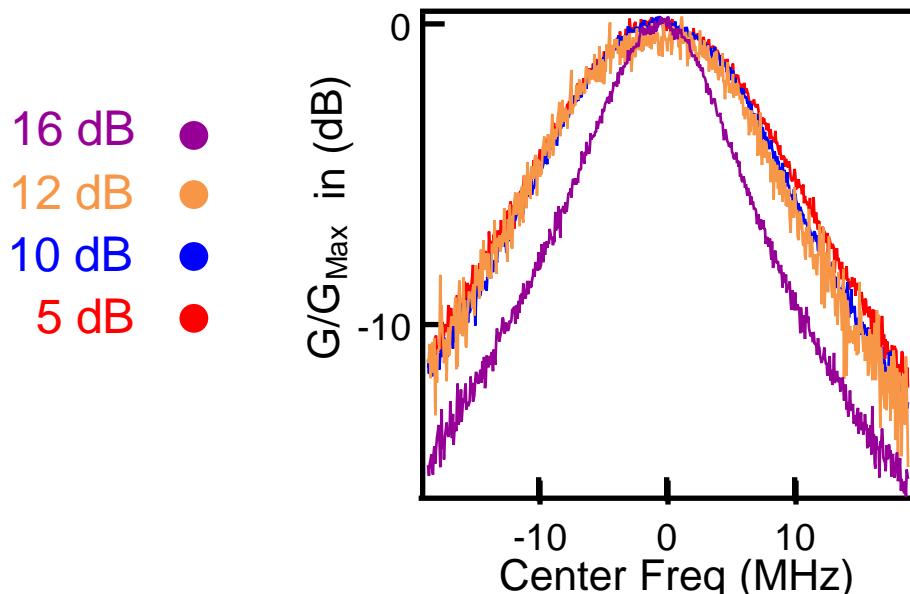
$$\frac{H_{Coupling}}{\hbar} = g_G(a^\dagger b^\dagger + a b) + g_C(a^\dagger b + a b^\dagger)$$



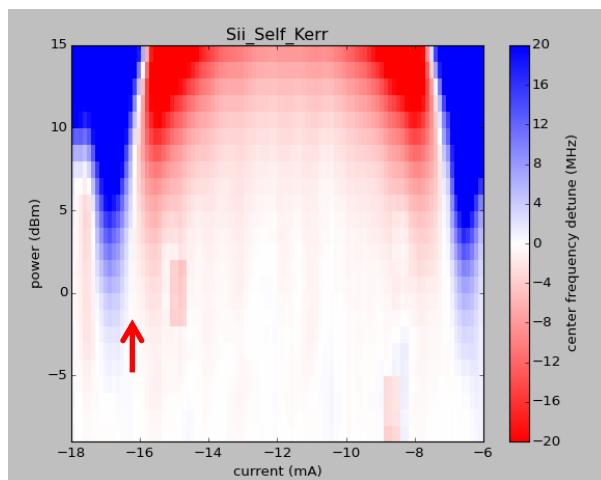
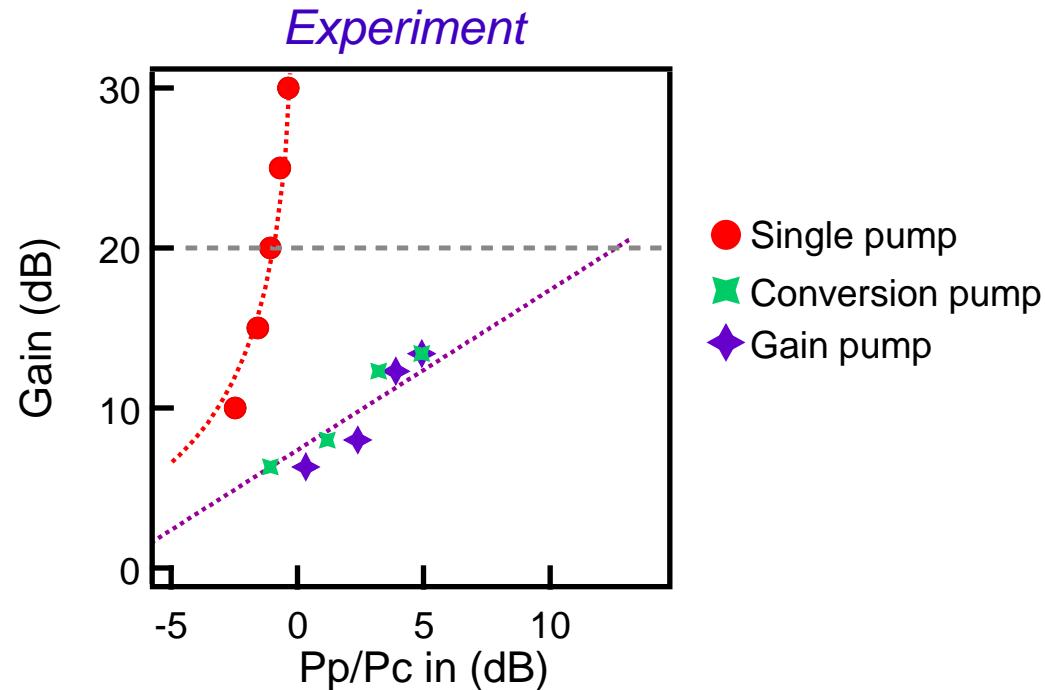
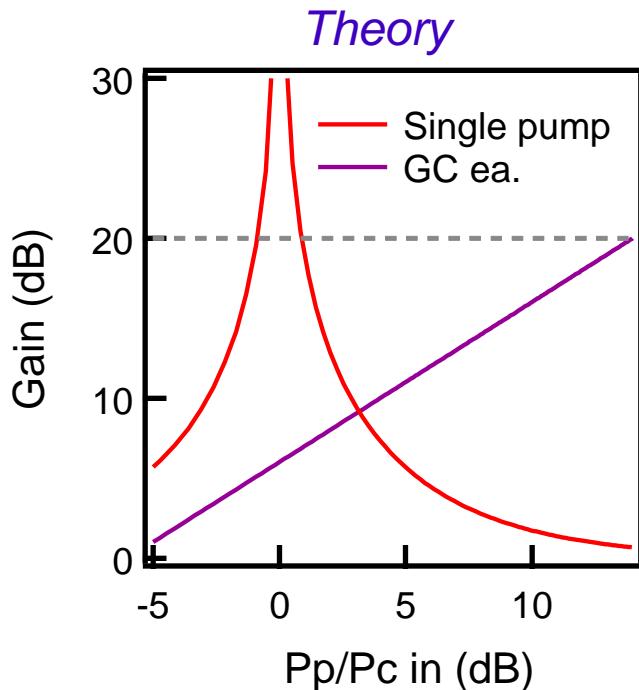
# Phase sensitive amplifier



# Fixed-bandwidth



# Why not 30 dB ??



# Future work

- More modes, more drives, better devices!
- New theory tools (is expansion of modes inadequate?)
- Better Hamiltonian engineering via improved fabrication
- Medium term: broadband, high-bandwidth directional amplifier

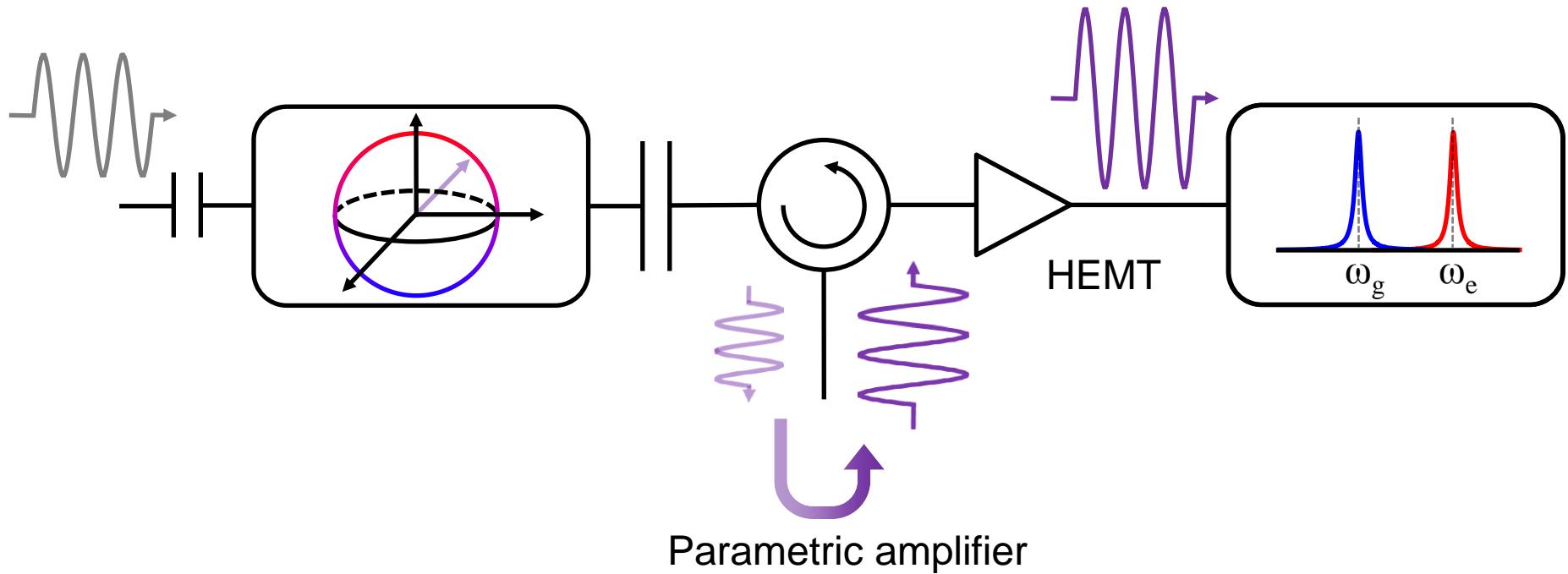
D. Pekker



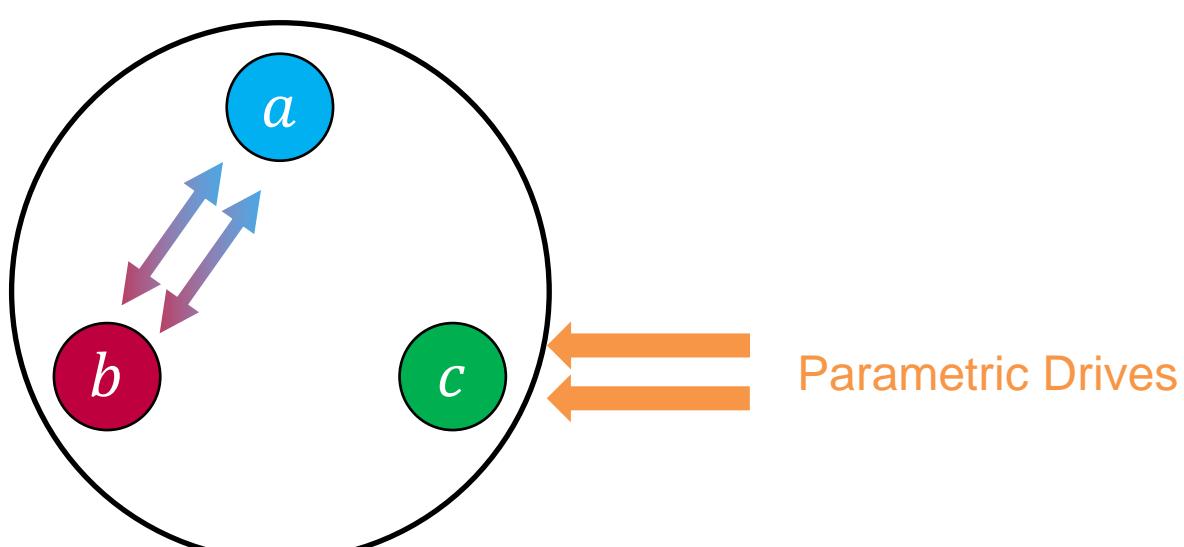
<http://hatlab.pitt.edu/>

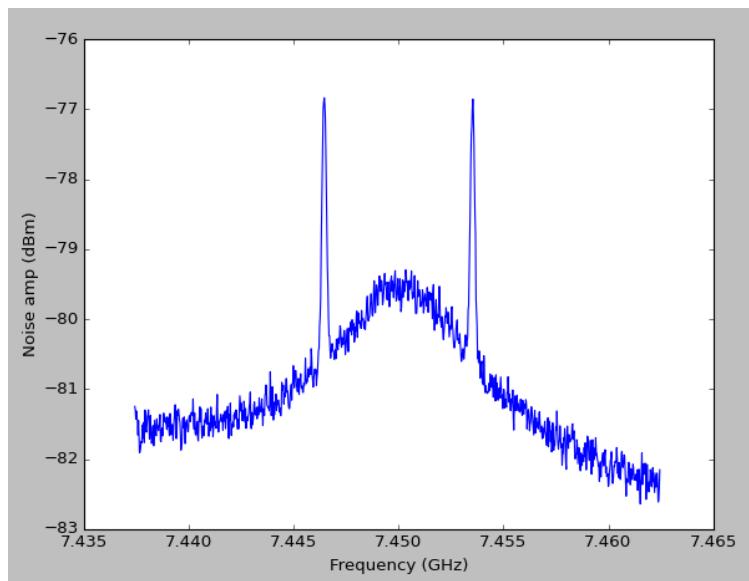
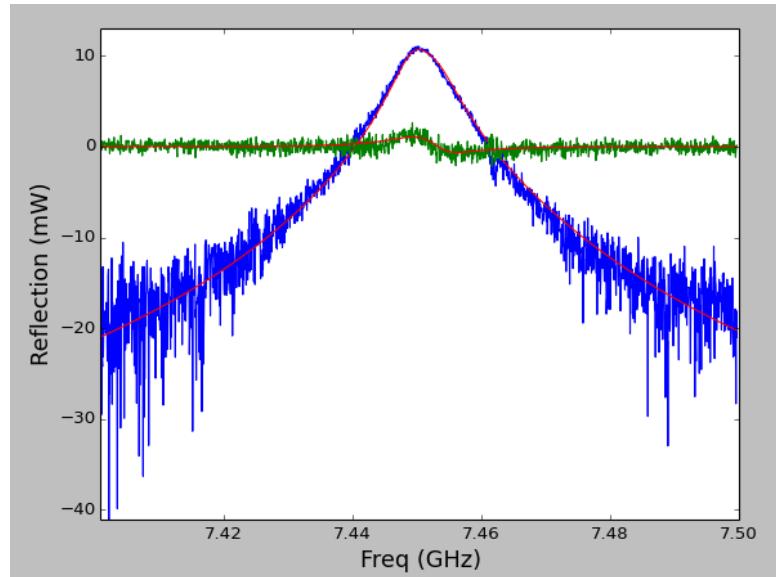
# **Back up**

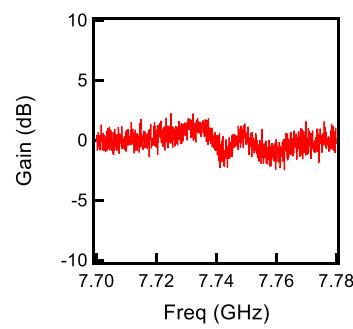
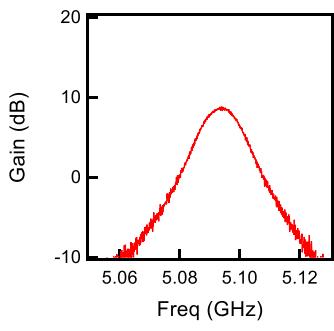
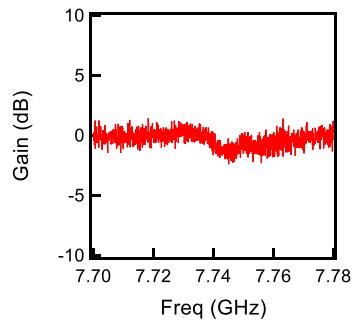
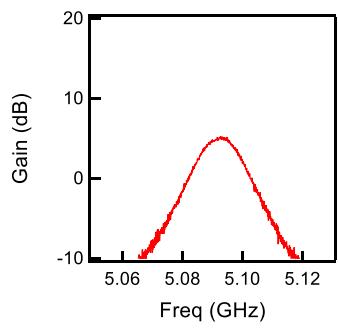
# The ubiquitous parametric amplifier

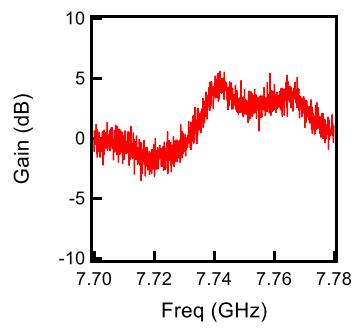
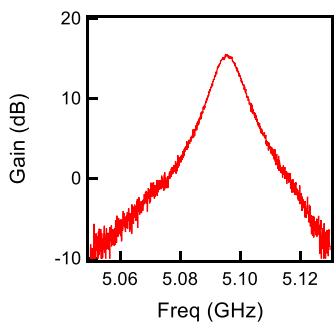
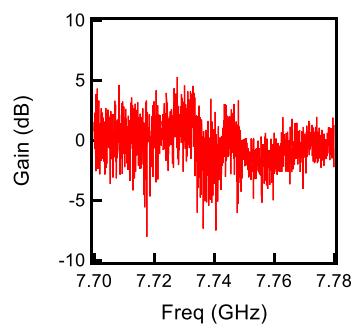
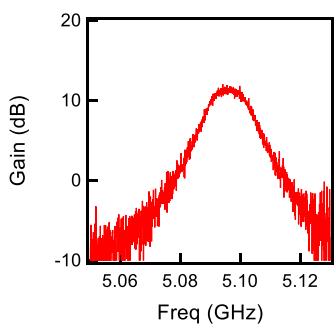


- Phase sensitive amplifier
- Fixed bandwidth
- Gain in transmission only
- Unity reflection gain

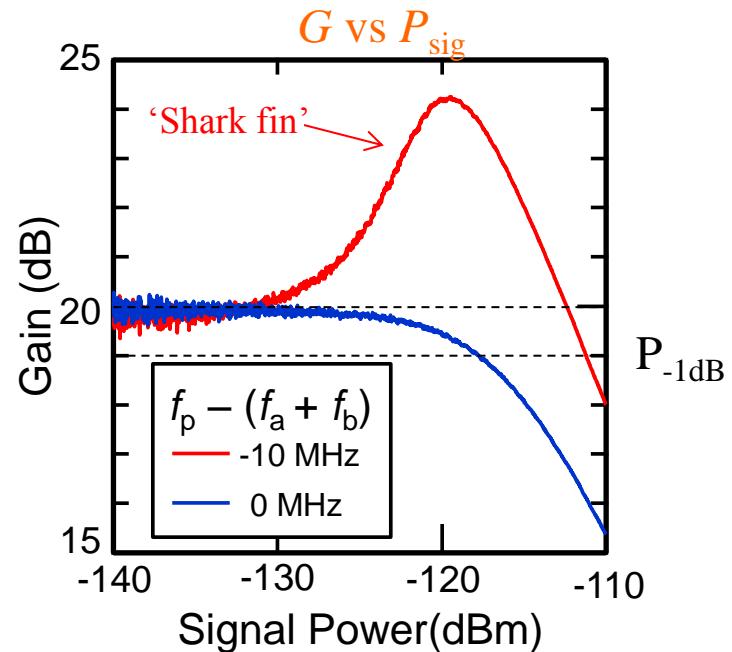
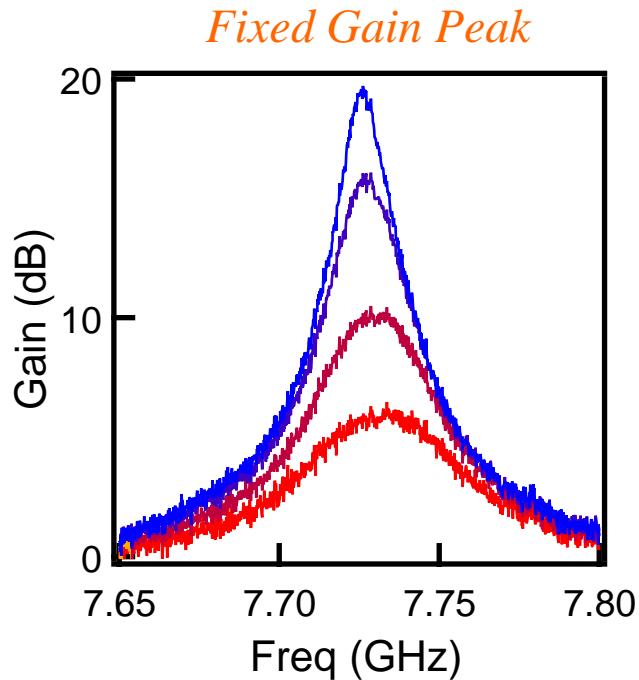








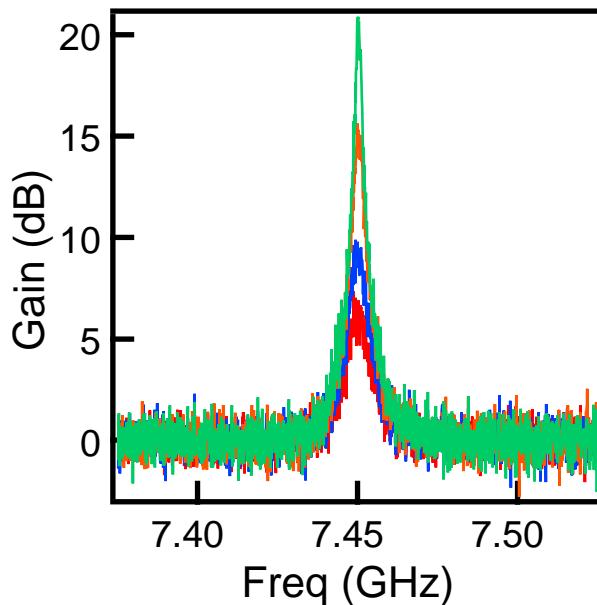
# Higher order



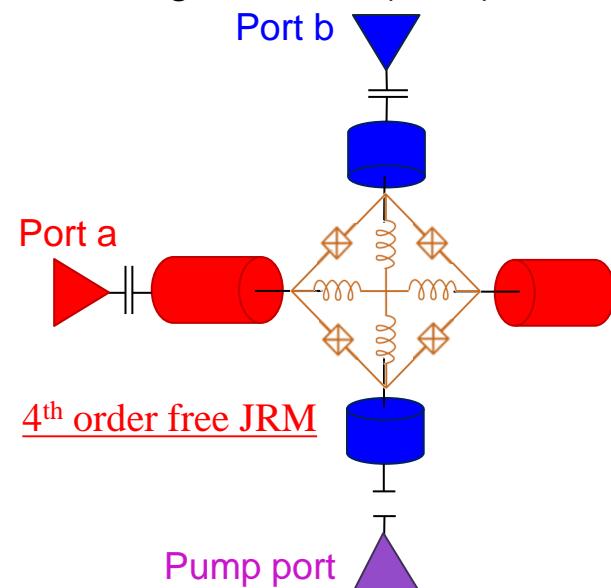
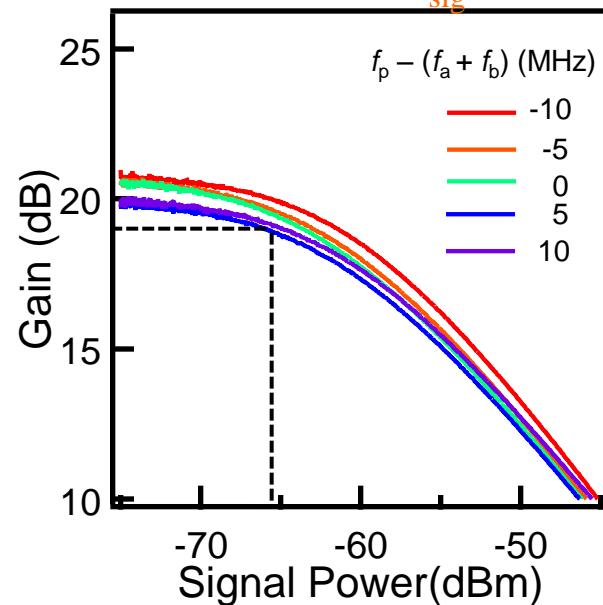
- 4th order terms shift mode freqs
- Shark-fin dynamic range

# Higher order

*Fixed Gain Peak*

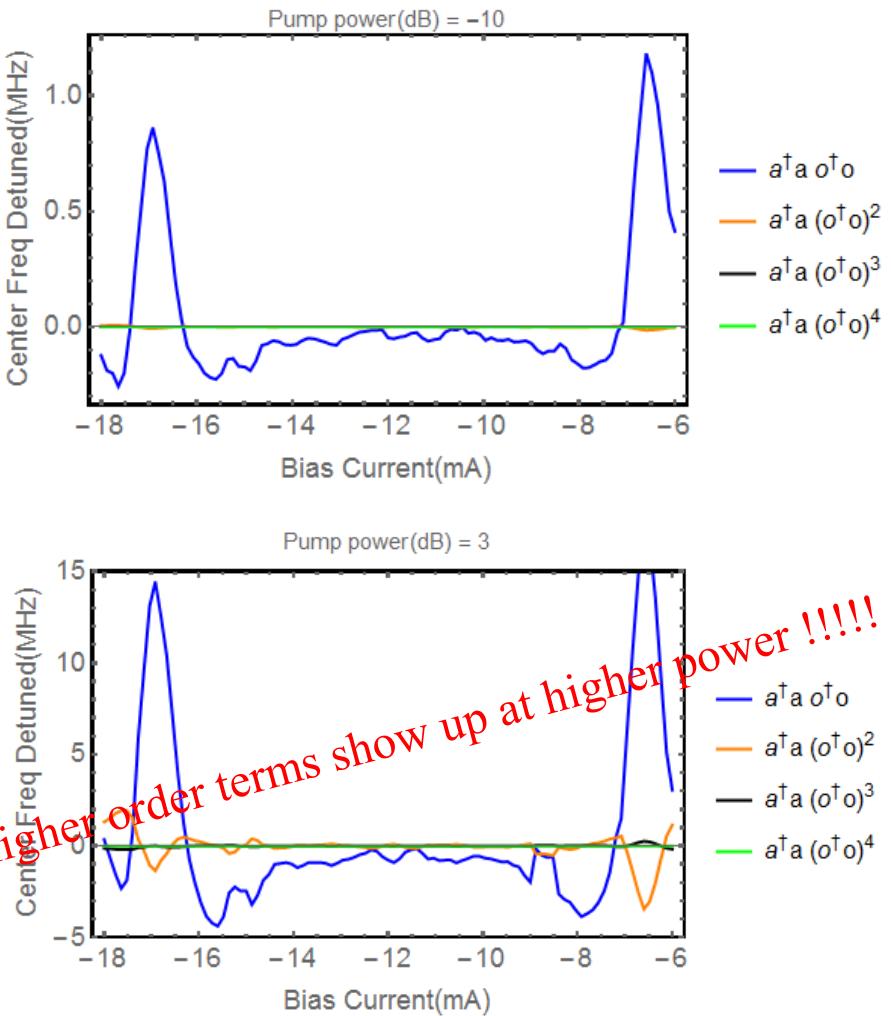
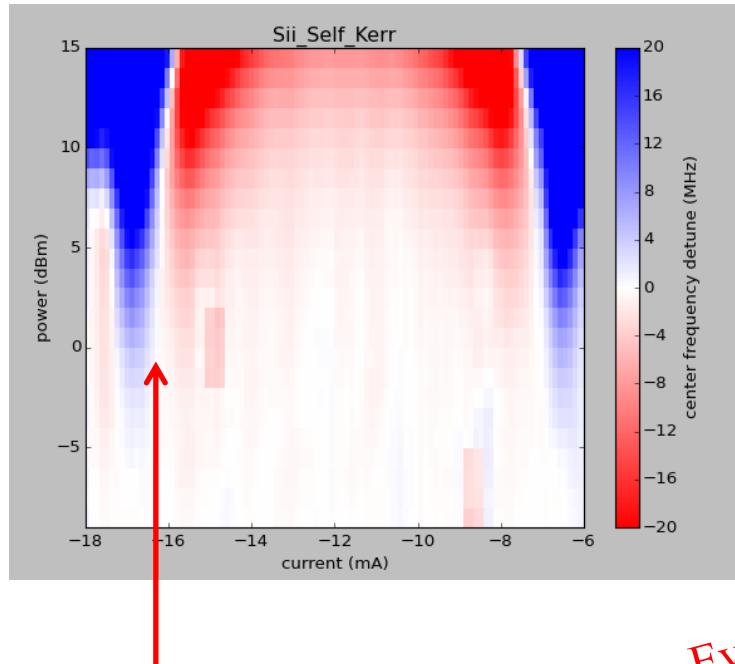


*G vs  $P_{\text{sig}}$*



- 4th order terms shift mode freqs
- Shark-fin dynamic range
- **4<sup>th</sup> order free device**  
( advertise Xi's talk)

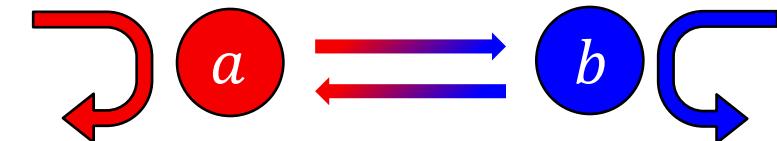
# Why not 30 dB Trans-Gain ?



# Bi-direction phase sensitive gain (XX)

$$\omega_{p1,p2} = \omega_a \pm \omega_b \neq \omega_c$$

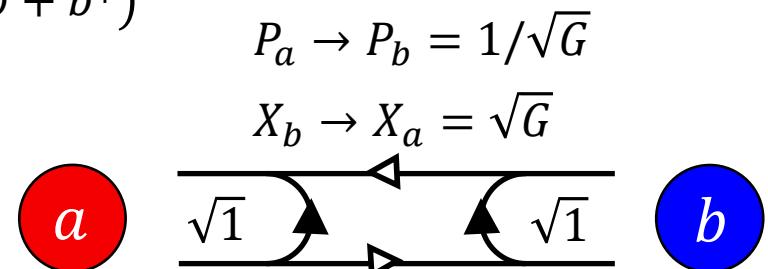
$$H_G = \hbar g(a + a^\dagger)(b + b^\dagger)$$



$$\omega_p \simeq \omega_a + \omega_b$$



$$\omega_p \simeq \omega_a - \omega_b$$



$$P_a \rightarrow P_b = 1/\sqrt{G}$$

$$X_b \rightarrow X_a = \sqrt{G}$$

$$G = \frac{P_P}{P_C}$$

- Enhanced bandwidth  $BW \simeq \frac{1}{2}\kappa/2\pi$  and saturation power!
- Have to match mode bandwidths
- Have to pump HARD!