

The damped wave equation with unbounded and singular damping

Petr Siegl

School of Mathematics and Physics, Queen's University Belfast, UK
<http://mathsites.unibe.ch/siegl/>

Based on

- [1] P. Freitas, P. Siegl, and C. Tretter (2018). “Damped wave equation with unbounded damping”. In: *J. Differential Equations* 264, pp. 7023–7054
- [2] A. Arifoski and P. Siegl (2018). *Pseudospectra of damped wave equation with unbounded damping*.

1. Basic concepts
2. Examples (singular and unbounded damping)
3. Results (existence & uniqueness, general spectral features, diverging damping, essential spectrum, pseudospectrum)

Damped wave equation in $\Omega \subset \mathbb{R}^d$

$$u_{tt}(t, x) + 2a(x) u_t(t, x) = (\Delta - q(x))u(t, x)$$

- $t \in \mathbb{R}_+$ time, $x \in \Omega$ spatial coordinate
- $a : \Omega \rightarrow \mathbb{R}_+$ damping, $q : \Omega \rightarrow \mathbb{R}_+$ potential
- initial conditions $u(0, \cdot) = u_0$, $u_t(0, \cdot) = u_1$
- Dirichlet boundary conditions at $\partial\Omega$ (if $\partial\Omega \neq \emptyset$)

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- almost exclusively $a \in L^\infty(\Omega)$ or “small” w.r.t to Δ

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Towards spectral analysis and semigroups

$$\partial_t \begin{pmatrix} u \\ v \end{pmatrix} = \underbrace{\begin{pmatrix} 0 & I \\ \Delta - q & -2a \end{pmatrix}}_G \begin{pmatrix} u \\ v \end{pmatrix}$$

- G acts in a suitable Hilbert space $\mathcal{H} = \mathcal{W}(\Omega) \times L^2(\Omega)$
- $a = 0 \implies iG$ is symmetric w.r.t. $\langle \nabla \cdot, \nabla \cdot \rangle_1 + \langle q^{\frac{1}{2}} \cdot, q^{\frac{1}{2}} \cdot \rangle_1 + \langle \cdot, \cdot \rangle_2$

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Associated quadratic pencil

- formally for eigenvalues

$$\begin{pmatrix} 0 & I \\ \Delta - q & -2a \end{pmatrix} \begin{pmatrix} u \\ v \end{pmatrix} = \lambda \begin{pmatrix} u \\ v \end{pmatrix}$$

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$$T(\lambda) := -\Delta + q + 2\lambda a + \lambda^2$$

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- generalization of spectra of operators: $\lambda \in \sigma(A) \iff 0 \in \sigma(A - \lambda)$

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- time-dependent solution: $e^{\lambda t} u(x)$

$$T(\lambda) = -\Delta + q + 2\lambda a_0 + \lambda^2, \quad a_0 \in \mathbb{R}_+$$

Finite interval $\Omega = (-1, 1)$ and $q = 0$

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- $\lambda_n = -a_0 \pm i\sqrt{\mu_n - a_0^2}$

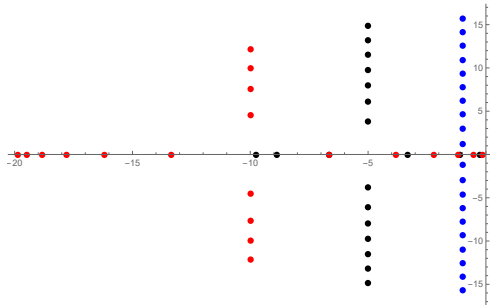
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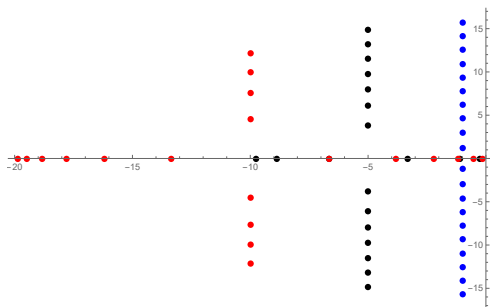
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Finite interval $\Omega = (-1, 1)$, $a \in L^\infty(\Omega)$

- asymptotic behavior of eigenvalues²

$$\lambda_n = \pm i \frac{n\pi}{2} - \frac{\int a(x) dx}{2} + \mathcal{O}(n^{-1}), \quad n \rightarrow \infty$$

- eigenfunctions form a suitable (Riesz) basis³

$\Rightarrow$ decay rate of a general solution determined by $\min\{\operatorname{Re} \lambda_n : n \in \mathbb{N}\}$

²P. Freitas and E. Zuazua (1996). *J. Differential Equations* 132, pp. 338–352.

³S. Cox and E. Zuazua (1994). *Comm. Partial Differential Equations* 19, pp. 213–243; P. Djakov and B. Mityagin (2010). *Math. Nachr.* 283, pp. 443–462.

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- the spectrum of

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is discrete, but **empty**

- every solution of

$$\begin{cases} u_{tt} + \frac{2}{x}u_t = u_{xx} \\ u(\cdot, 0) = u(\cdot, 1) = 0 \\ (u(0, \cdot), u_t(0, \cdot)) = (u_0, u_1) \in W^{1,2}(0, 1) \times L^2(0, 1) \end{cases}$$

vanishes in a finite time⁴

$$u(t, \cdot) = 0, \quad t > 2$$

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1D examples on $\Omega = \mathbb{R}$

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- eigenvalues $\lambda_k = 2^{\frac{1}{2n+1}} e^{\pm i\pi \frac{n+1}{2n+1}} \mu_{k;n}^{\frac{n+1}{2n+1}} - \frac{2(n+1)}{2n+1} a_0 + o_k(1), \quad k \rightarrow \infty$
- $\{\mu_{k;n}\}_k$ are eigenvalues of the self-adjoint $-\frac{d^2}{dx^2} + x^{2n}$ in $L^2(\mathbb{R})$
- for $n = 1$: $\mu_{k;1} = 2k + 1, \quad k = 0, 1, 2, \dots$

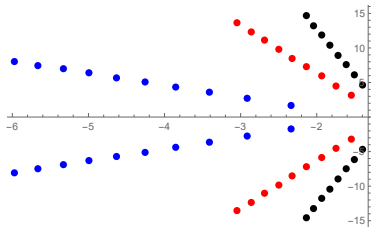
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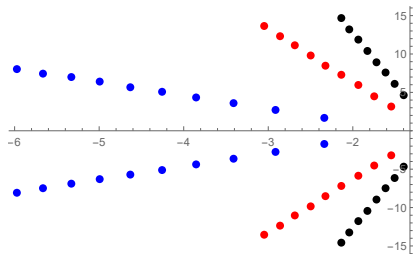
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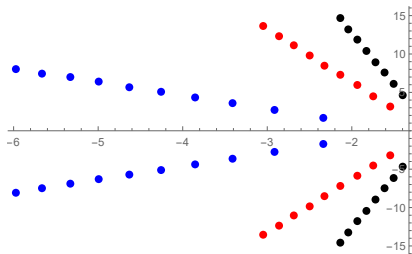
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- it seems that EV's converge to those of $a_\infty(x) = a_0$ on $(-1, 1)$
- similar effect as for Schrödinger operators (convergence to the “square-well”)

$$-\frac{d^2}{dx^2} + x^{2n} \quad \text{in } L^2(\mathbb{R}) \xrightarrow{n \rightarrow \infty} -\frac{d^2}{dx^2} \quad \text{with Dirichlet BC in } L^2(-1, 1)$$

Infinite 2D strip

- $\Omega = \mathbb{R} \times (-1, 1)$ and $a(x, y) = x^2 + a_0$
- separation of variables: algebraic equation for eigenvalues λ

$$2\lambda\mu_k = (\lambda^2 + \sigma_j^2 + 2\lambda a_0)^2$$

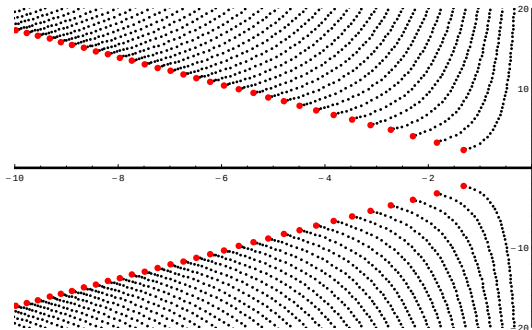
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Theorem (semigroup)

Under suitable regularity assumptions:

- $-G$ is m-accretive

$\Rightarrow G$ generates a contraction semigroup in $\mathcal{H} = \mathcal{W}(\Omega) \times L^2(\Omega)$

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Theorem (relation of spectra)

Under suitable regularity assumptions: $\forall \lambda \in \mathbb{C} \setminus (-\infty, 0]$

$$\begin{aligned} \lambda \in \sigma(G) &\iff 0 \in \sigma(T(\lambda)) \\ \lambda \in \sigma_p(G) &\iff 0 \in \sigma_p(T(\lambda)) \\ \lambda \in \sigma_{\text{ess}}(G) &\iff 0 \in \sigma_{\text{ess}}(T(\lambda)) \end{aligned}$$

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If a is unbounded at infinity:

- $\sigma(G) \cap \mathbb{C} \setminus (-\infty, 0] =$ isolated EV's of finite multiplicity
- possible accumulations only to $\mathbb{R}_-$

Assumption: limiting damping

- sequence of unbounded dampings $\{a_n\}$: $\lim_{|x| \rightarrow \infty} a_0(x) = \infty$ and
$$\forall n \in \mathbb{N}, \quad a_n \geq a_0$$
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Under the assumption above:

- every $\lambda \in \mathbb{C} \setminus (-\infty, 0] \cap \sigma_p(T_\infty)$ is approximated:

$$\forall \lambda \in \mathbb{C} \setminus (-\infty, 0] \cap \sigma_p(T_\infty), \quad \exists \{\lambda_n\}, \quad \lambda_n \in \sigma_p(T_n), \quad \lambda_n \rightarrow \lambda$$

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- no pollution: if

$$\{\lambda_n\}_{n \in \mathbb{N}} \subset \mathbb{C} \setminus (-\infty, 0], \quad \lambda_n \in \sigma_p(T_n)$$

has an accumulation point $\lambda \in \mathbb{C} \setminus (-\infty, 0]$, then $\lambda \in \sigma_p(T_\infty)$

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- the corresponding eigenfunctions converge in L^2

The issue

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$$\sigma\left(-\frac{d^2}{dx^2} - x^2\right) = \mathbb{R} \quad \text{vs.} \quad \sigma\left(-\frac{d^2}{dx^2} - x^4\right) = \sigma_{\text{disc}}\left(-\frac{d^2}{dx^2} - x^4\right)$$

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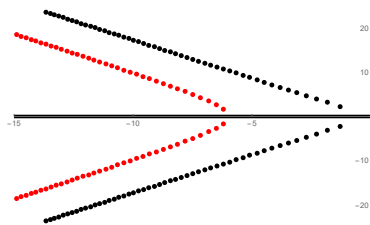
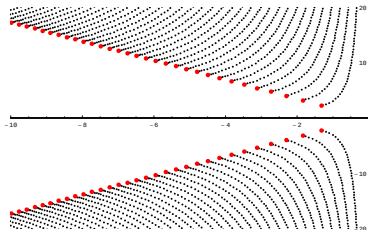
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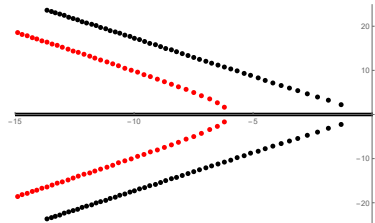
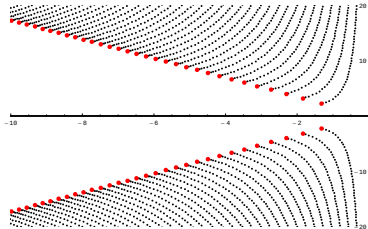
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Examples in $d = 1$: $a_n(x) = x^{2n}$ and $q(x) = q_0$

- $\sigma_{\text{ess}}(G) = (-\infty, 0]$



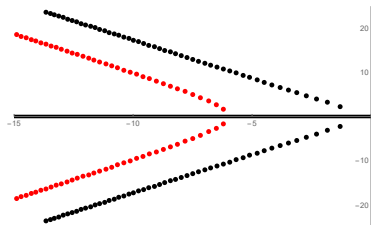
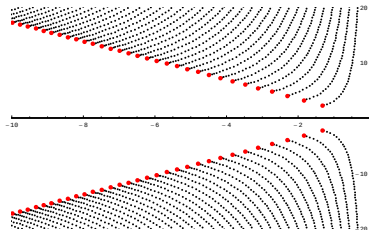
⁵R. Ikehata and H. Takeda (2018). *Funkcialaj Ekvacioj*.



Overdamping due to $\sigma_{\text{ess}}(G)$

- no exponential decay of the semigroup possible

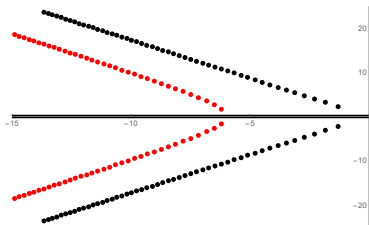
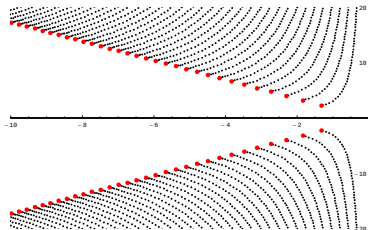
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Overdamping due to $\sigma_{\text{ess}}(G)$

- no exponential decay of the semigroup possible
- shifting σ_{ess} of 0: adding “larger” q than a : e.g. $q(x) = 10x^2$, $a(x) = x^2$

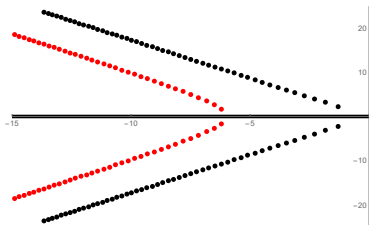
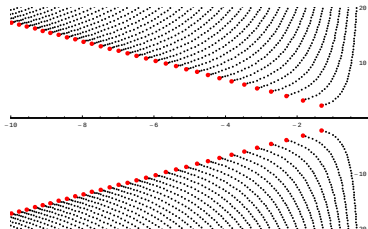
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- crucial spectral difference:
 $a_n(x) = x^{2n} + a_0$ on $\mathbb{R}$ vs. $a_\infty(x) = a_0$ on $(-1, 1)$
- recall: eigenvalues and eigenmodes do converge vs. “jump” in σ_{ess}

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- recall: eigenvalues and eigenmodes do converge vs. “jump” in σ_{ess}
- polynomial decay⁵ for $\Omega = \mathbb{R}^d$, $d \geq 3$ and $a(x) \geq a_0 > 0$

$$\|e^{itG}\| = \mathcal{O}(t^{-1}), \quad t \rightarrow +\infty$$

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- ε -pseudospectrum

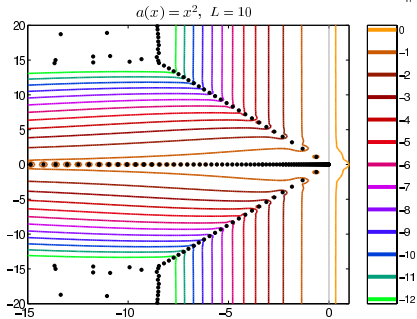
$$\sigma_\varepsilon(G) := \sigma(G) \cup \left\{ z \in \mathbb{C} : \|(G - z)^{-1}\| > \frac{1}{\varepsilon} \right\}$$

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Eigenvalues inside the $\varepsilon = 10^{-12}$ -pseudospectrum (top and bottom left) are not reliable.

Theorem [Arifoski & PS] (Pseudospectrum)

Let

- $d = 1$, $\Omega = \mathbb{R}$
- $q(x) = 0$, $a(x) = x^{2n}$ with some (fixed $n \in \mathbb{N}$)

Then for every $\omega \in (\pi/2, \pi) \cup (\pi, 3\pi/2)$

$$\lim_{t \rightarrow +\infty} \|(G - e^{i\omega}t)^{-1}\|_{\mathcal{H}} = \infty.$$

Generation of the semigroup

- accretivity of $-G$ is easy (numerical range)

⁷W. N. Everitt and M. Gierztz (1978). *Proc. Roy. Soc. Edinburgh Sect. A* 79, pp. 257–265.

Generation of the semigroup

- accretivity of $-G$ is easy (numerical range)
- crucial step for m-accretivity:

$$\|(-\Delta + q + \gamma a)f\| \approx \|(-\Delta + q)f\| + |\gamma|\|af\|$$

if $\gamma \in \mathbb{C} \setminus (-\infty, 0]$ and the regularity assumption **I** holds

$$|\nabla a| \leq \varepsilon a^{\frac{3}{2}} + M_{\nabla}(q^{\frac{1}{2}} + 1)$$

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- defined through the form

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Spectral convergence

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- norm convergence from “collective compactness”¹⁰:
 $\{\|T_n(\lambda)\phi_n\| + \|\phi_n\|\}_{n \in \mathbb{N}}$ is bounded $\implies$
 $\{\|T_n(\lambda)\phi_n\| + \|\phi_n\|\}_{n \in \mathbb{N}}$ has a convergent subsequence in $L^2(\Omega)$

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- based on WKB approximation

$$\phi_n(x) = \varphi_n(x) \exp\left(i \int_0^x \sqrt{-(q(t) + 2\lambda a(t) + \lambda^2)} dt\right)$$

- φ_n : dispersing cut-off with $\text{supp } \varphi_n \rightarrow \infty$

Pseudospectrum

$$a(x) = x^2 \quad \implies \quad \|(G-\lambda)^{-1}\|_{\mathcal{H}} \geq C_N |\operatorname{Re} \lambda|^N, \quad \operatorname{Re} \lambda \rightarrow -\infty \text{ \& \, } |\operatorname{Im} \lambda| \geq \delta > 0.$$

- pseudomodes $\{\Psi_\lambda\}_\lambda = \{(f_\lambda, \lambda f_\lambda)^t\}_\lambda \subset C_0^\infty(\mathbb{R})^2$ such that

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$$f_\lambda := \xi_\lambda g_\lambda, \quad g_\lambda(x) = e^{-i \int_0^x \sqrt{-\lambda^2 - 2\lambda a(t) - q(t)} \, dt}.$$

- $\xi_\lambda \in C_0^\infty(\mathbb{R})$: λ -dependent cut-off