

ICTS TALK:

D2-Branes from Periodic M2-brane

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INTRO - Usual reduction of M2 \rightarrow D2
by "Naive Higgs Mechanism"



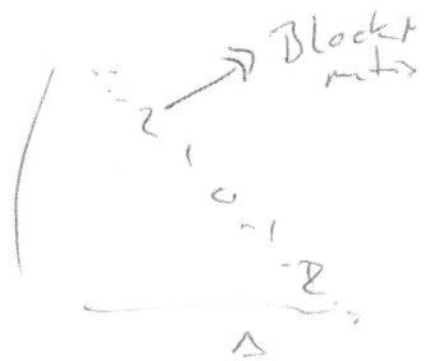
- But the most naive way would be to compactify the transverse space $\mathbb{R}^8 \rightarrow \mathbb{R}^7$

Easy to arrange for a periodic array of ADJM

$$\langle Z^{A'} \rangle = 0$$

$$Z^A = 2\pi i R$$

$$A = 1, 2, 3, 4, \quad A' = 1, 2, 3$$



But the problem is to impose the discrete orbifold or fluctuations.

$$Z^A = 2\pi i S_4^A R \Delta + \phi^A$$

wait ~~ϕ~~ $\phi_{m1} = \phi_{m-1}$

Translation is not a symmetry of ABJM for

But on the other hand, at least for large N ,
ABJM is weakly coupled.

• something should work

→ So here we proceed with "Grand finite"

~~ABJM~~ Motivations: 1) Resolve a puzzle

- cf. W. Taylor + D. Gross.
periodic way.

(T-duality) $D(p+1)$ -brane as S^1

- In fact KK tower.

→ Don't expect such states here

2) Explore T-duality in M-theory,

in particular $M2 \cong \Pi^3$ way,

$\cong M5$ on Π^3 .

Close to published but a few points

Remarks:

PLAN:

1) INTRODUCTION ✓

2) SET-UP

3) Evaluate the Lagrangian

4) Compare to SDMSYM

5) Further Π^2 arg to ~~SDSYM~~ ^{D4-5d}

6) Conclusions.

2) So we impose.

$$Z_{mn}^A = 2\pi i R S_4^A Z_{m-1, n-1}^A$$

$$A_{\mu\nu}^{LR} = A_{\mu\nu-m}^{LR} \text{ etc.}$$

→ leads to divergences.

e.g. $\text{tr}(D_r Z D^r \bar{Z})$

$$= \sum_{m,n} \text{tr}(\partial Z_{mn} \partial Z_{mn}^+)$$

$$= \sum_{m,n} \text{tr}(\partial Z_{m-n,0} \partial Z_{m-n,0}^+)$$

$$= \sum_p \sum_q \text{tr}(\partial Z_{p,0} \partial Z_{+p,0}^+)$$

$$= (\sum_p) \sum_q \text{tr}(\partial Z_{p,0}^A \partial Z_{+p,0}^+)$$

$$\sum_p \text{tr} = \infty.$$

This arises in the D-brane case but is harmless as $|\hat{1}|$ just multiplies the whole Lagrangian.

We will find other divergences e.g.

$$\sum_i q^2 = \cancel{\sum_i q^2} \dots$$

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~~and $\sum_i q^2 = \dots$~~

To handle these we consider a large but finite cross $n = -N, \dots, 0, \dots, +N$
 - normalize fluctuations. eg. $\sum_i q^2 = \frac{2}{3} N^3 \dots$

$$Z_m^A = 2\pi i R \cdot S_4^A \times S_{mm} + \frac{1}{\sqrt{2N}} \phi_{n-m}^A$$

$$A_{\mu m}^{LIR} = a_{\mu n-m}^{LIR}$$

$$\Rightarrow \text{tr}(\partial \partial \partial \partial) = \sum_P \text{tr}(\partial \phi_P^A \partial \phi_{A1}^A)$$

But crucially also let $k \propto R$
 scale $\propto N$.

We ^{still} find. divergences

1st eg. for $\text{Tr}(DZ D\bar{Z})$

$$\sim \sum p^L a_{\mu p}^+ a_{-p}^{\bar{R}}$$

$$a_{\mu}^+ = \frac{1}{2}(a_{\mu}^L + a_{\mu}^R)$$

$$\Rightarrow A_{\mu}^L = A_{\mu}^R$$

$U(M) \times U(N) \rightarrow U(M)$ adjoint.
 Bifund.
 \downarrow

2nd

for $\text{Tr}(DZ D\bar{Z})$

$$\sim \sum_{\substack{p^L \\ \bar{L}N}} \left(\phi_{\mu p}^{A1} \phi_{A1 p} \right) \sim \left(\frac{R^4 N^2}{L^2} \right)$$

$$\Rightarrow \phi_{\mu p}^{A1} = 0 \quad \forall p$$

- resolves the no puzzle

~~So we set~~

$$\text{So we set } \begin{cases} \phi_{\mu p}^{A1} = 0 & p \neq 0 \\ a_{\mu p}^L - a_{\mu p}^R = 0 \end{cases}$$

\Rightarrow Constraints.

$\{H, C\} = C' \quad \Rightarrow$ more constraints

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To make a lagrangian short we find only zero-modes:

$$\phi_0^A, a_{\mu 0}^+$$

$$a_{rP}^+ + \ln \phi_P^A \propto \mathbb{1}$$

and can be integrated out

$$a_{rP}^+ = \frac{1}{2\pi i} \frac{1}{R} \frac{1}{5N} \partial_r \ln \phi_P^A$$

So in the end we find:

$$\mathcal{L} = -(\partial_\mu \phi_0^A \partial^\mu \phi_{A0}) - V$$

$$V = \frac{1}{2} g^2 \left(\left| [\phi_0^A, \phi_{B0}^+] \right|^2 + \left| \phi_0^{A'}, \phi_{B0}^+ \right|^2 + 2 \left| [\phi_0^{A'}, \ln \phi_0^A] \right|^2 \right)$$

$$g^2 \propto \frac{R^2 N}{k^2}$$

$$M_{\text{pl}}^2 = \frac{R^4 N^2}{k^2}$$

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Consider the Fermions

how we see stars like

$$\sum_i g_i = ?$$

let us assume $\sum_i g_i = \Omega N^2$

might say $\Omega = 0$

on the other hand tower of Fermions

$$\frac{\Omega^2 \Omega N^2}{L^3 N} \sum_P P \psi_P^A \psi_{AP}$$

$$\Rightarrow M_S \propto \Omega N \frac{\Omega^2}{L^2} = M_S \text{ if } \Omega = \frac{L^2}{N}$$

then

$$Z = Z_S + -i \text{tr}(\psi_0 \gamma^{\mu} \psi_{\mu} \psi_0)$$

- Z_{Yukawa}

Furthermore more it turns out that Z_{Yukawa}

and U are exactly the same as

3D N=4 SYM.

$$\Rightarrow \mathcal{L}_{\text{array}} = -\frac{1}{2} \nabla_r X^I \nabla X^I - \frac{1}{2} \nabla_r Y \nabla Y + \frac{1}{4} g^2 ([X^I, X^J])^2 + \text{Fermions.}$$

$$Y = \text{Re } \psi^4, \quad I = 3, \dots, 9$$

\Rightarrow $SO(7)$ invariant.

but no kinetic term for $A_r = \epsilon^+_{r0}$

But all is not lost: A_r eq. of M.

$$i [X^I, \nabla_r X^I] + i [Y, \nabla_r Y] = 0.$$

$$\nabla^2 Y = 0.$$

with $\nabla_r Y = -\frac{1}{2} g^2 \epsilon_{r \rightarrow} F^{IJ}$

$$\Rightarrow [\nabla_r, \nabla_r Y] = \frac{1}{g^2} \nabla^{\mu\nu} F_{\mu\nu}.$$

\Rightarrow recover 3D SYM.

i.e. all solutions to 3D SYM solve our system.

but in general there are extra "modes".

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e.g. $A_r = ig' \partial_r g$

$$Y = g \chi_0 g^{-1} \quad \chi_0 = \text{"modul."}$$

\Rightarrow presence of M2's in 11 FL Dim.

$Y \sim$ Dual gauge.

Consider a further T-duality $\rightarrow \mathbb{T}^2$

\rightarrow follow usual procedure of W. Taylor

$$S = \frac{-1}{g^2} \int d^5x \left[\frac{1}{2} (\nabla^{\mu'} X^{I'}) (\nabla_{\mu'} X^{I'}) - \frac{1}{4} \sum (X^A)^2 \right. \\ \left. + \frac{1}{2} (\nabla_r Y) (\nabla^r Y) + \frac{1}{2} \bar{F}_{\mu\nu} \bar{F}^{\mu\nu} + \frac{1}{4} \bar{F}_r^2 \right]$$

\rightarrow not Lorentz invariant.

but again, $\nabla_r Y = -\frac{1}{2} \epsilon_{r\mu\nu} F^{\mu\nu}$

gives. eqn of motion of SDMSYM.

This here might say

$$\nabla_r Y = \frac{1}{2} \epsilon^{\mu\nu\lambda} F_{\mu\nu} \partial_\lambda$$

$$Y \approx \int_{-\infty}^{\infty} B \quad ?$$