Phase transitions and symmetry breaking in current distributions of diffusive systems

Yariv Kafri

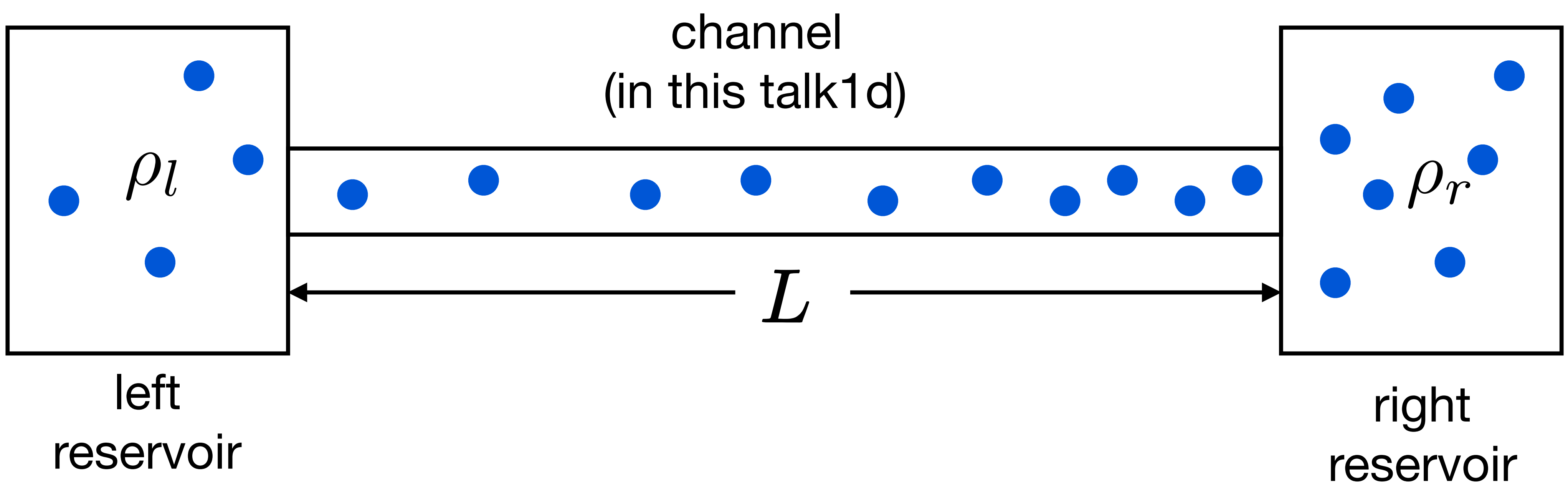
Department of Physics, Technion

Yongjoo Baek (Technion -> Cambridge), Vivien Lecomte (Grenoble)

PRL **118**, 030604 (2017)

Settings: **boundary** driven diffusive systems

- Diffusive *interacting+conserving* channel (`**disordered**' phase - think gas)
- Channel connected to two reservoirs at given densities



Question: Consider current probability distribution

$$P(J) \sim \exp[-TL\underline{\Phi(J)}] \text{ for large } T \text{ and } L$$

Large deviation function (LDF)

Here J is the time-averaged current

T the window of time over which we average

Are there cases where $\Phi(J)$ is singular?

- Know to occur for driven-diffusive-systems with periodic boundary conditions

  WASEP 1D - Bodineau, Derrida, PRE **72**, 066110 (2005) Espigares et al., PRE **87**, 032115 (2013)
  WASEP 2D - Tizón-Escamilla *et al.*, arXiv:1606.07507
  KMP 1D - Bertini *et al.*, JSP **123**, 237 (2006), Hurtado, Garrido, PRL **107**, 180601 (2011)

- Suggested to be possible in boundary driven in Bertini *et al.*, PRL **94**, 030601 (2005)

  no microscopic model, scenario actually different

Answer:

- Two types of possible phase transitions:
  1. symmetry breaking (continuous)
  2. first-order

- Mechanism different from periodic boundary conditions

- Give general conditions for which models exhibit phase transitions

- Identify microscopic models

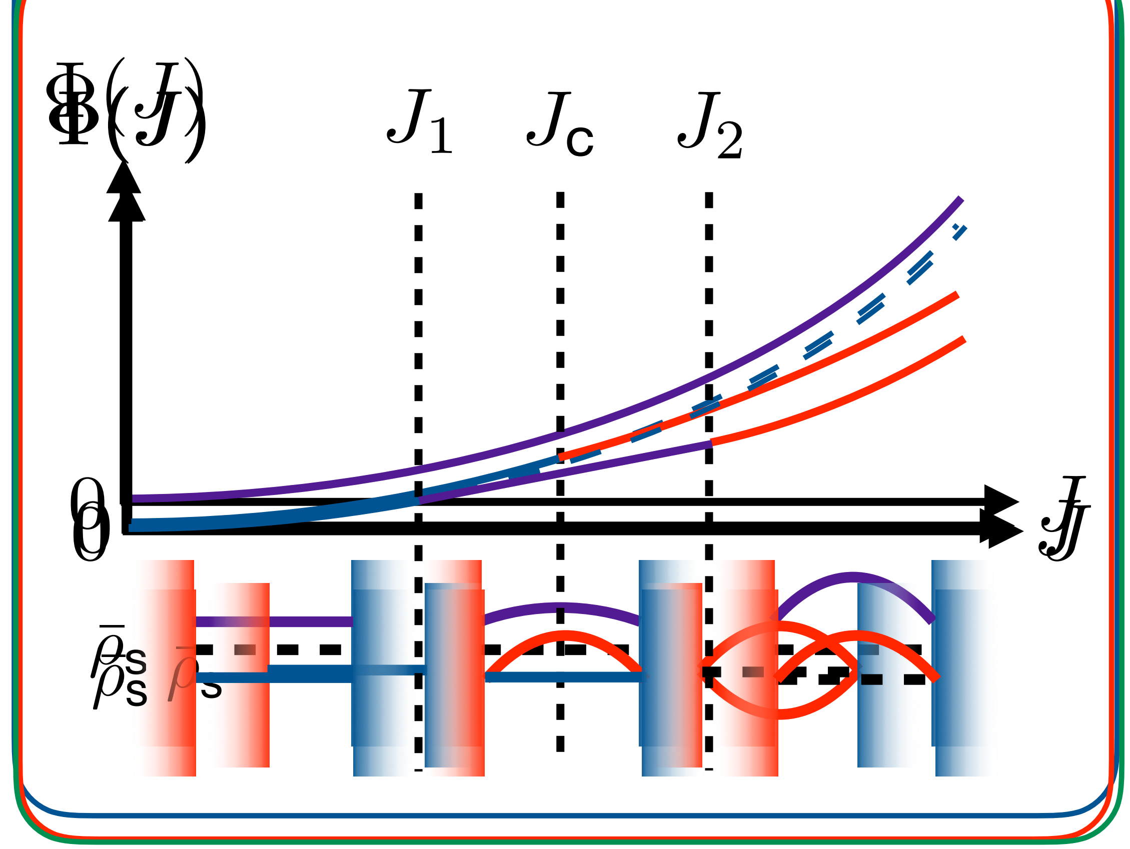
- Transitions occur even when system is in equilibrium
  (equal reservoir density, no bulk field - reversible dynamics)

comment:

another mechanism identified in Shpielberg, Don, Akkermans, PRE **95**, 032137 (2017)

Cartoon of transition scenarios

First-order DPTs



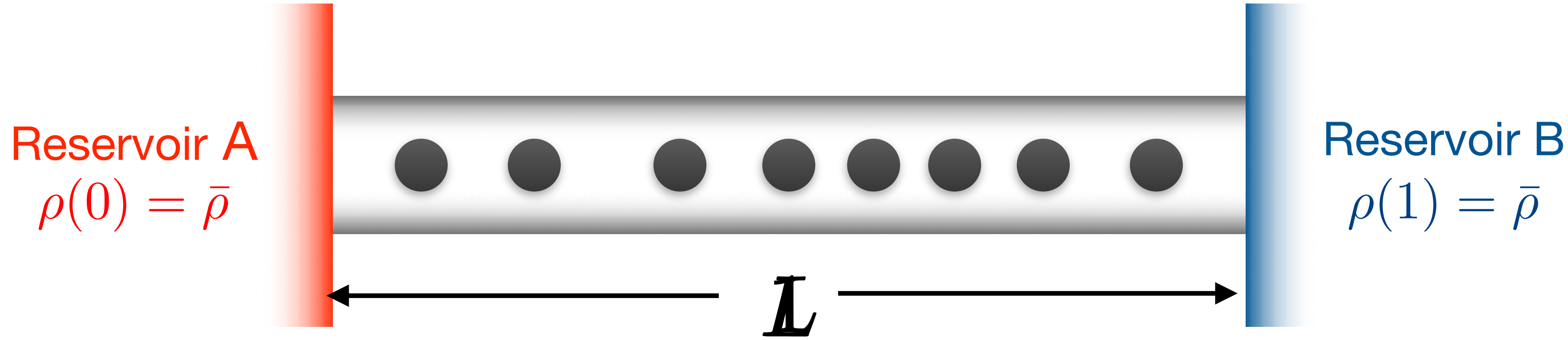
Outline

- Quick recap - formalism, some models, macroscopic fluctuation theory, ensembles, additivity principle.

- Perturbative description of transitions - develop a **Landau theory**

Results general for any model

The formalism



On large length scales can characterize the system by two linear-response quantities

Diffusivity $D(\rho)$ mobility $\sigma(\rho)$

which obey $\frac{2D(\rho)}{\sigma(\rho)} = f''(\rho)$ $f(\rho)$ - free-energy density

- After diffusive rescaling $i \to Lx, \quad t \to L^2 t$ the density field $\rho(x)$ obeys

$$\partial_t \rho = -\partial_x \left[\underbrace{-D(\rho)\, \partial_x \rho}_{\text{Diffusion}} + \underbrace{\sqrt{\sigma(\rho)}\, \eta}_{\text{Noise}} \right]$$

- The noise is weak in the thermodynamic limit $L \to \infty$

$$\langle \eta(x,t)\, \eta(x',t') \rangle = \frac{1}{L}\, \delta(x - x')\, \delta(t - t')$$

The generating function

- Instead of calculating $P(J) \sim \exp[-TL\Phi(J)]$ calculate the generating function

$$\langle e^{TL\lambda J} \rangle \sim \exp[TL\Psi(\lambda)]$$

where as usual $\Psi(\lambda) = \sup_J \left[\lambda J - \Phi(J)\right]$

- Using Martin-Siggia-Rose

$$\langle e^{TL\lambda J} \rangle \sim \int \mathcal{D}\rho \, \mathcal{D}\hat{\rho}_\lambda \, \exp\left\{ -L \int_0^T \mathrm{d}t \int_0^1 \mathrm{d}x \left[\hat{\rho}_\lambda \dot{\rho} - H(\rho, \hat{\rho}_\lambda) \right] \right\}$$

with

$$\rho(0,t) = \rho(1,t) = \bar{\rho}$$

$$\hat{\rho}_\lambda(0,t) = 0, \quad \hat{\rho}_\lambda(1,t) = \lambda$$

and the Hamiltonian $H(\rho, \hat{\rho}_\lambda) = -D(\rho)(\partial_x \rho)(\partial_x \hat{\rho}_\lambda) + \dfrac{\sigma(\rho)}{2}(\partial_x \hat{\rho}_\lambda)^2$

Large L so calculate saddle point.

$$\Psi(\lambda) = -\lim_{T\to\infty} \frac{1}{T} \inf_{\rho(t),\hat{\rho}_\lambda(t)} \int_0^T \mathrm{d}t \int_0^1 \mathrm{d}x \left[\hat{\rho}_\lambda \dot{\rho} - H(\rho, \hat{\rho}_\lambda)\right]$$

or solve (with boundary conditions) - note momentum related to noise

$$\partial_t \rho = \frac{\delta}{\delta \hat{\rho}_\lambda} \int_0^1 \mathrm{d}x\, H(\rho, \hat{\rho}_\lambda) = \partial_x \left[D(\rho)\partial_x \rho - \sigma(\rho)\partial_x \hat{\rho}_\lambda\right]$$

$$\partial_t \hat{\rho}_\lambda = -\frac{\delta}{\delta \rho} \int_0^1 \mathrm{d}x\, H(\rho, \hat{\rho}_\lambda) = -\partial_x \left[D(\rho)\partial_x \hat{\rho}_\lambda\right] - \frac{\sigma'(\rho)}{2} \left(\partial_x \hat{\rho}_\lambda\right)^2$$

Simplification - the solutions which minimize action are *time-independent*

= **additivity principle**

Bodineau, Derrida, PRL **92**, 180601 (2004)

$$\Psi(\lambda) = -\lim_{T\to\infty} \frac{1}{T} \inf_{\rho(t),\hat{\rho}_\lambda(t)} \int_0^T \mathrm{d}t \int_0^1 \mathrm{d}x \left[\hat{\rho}_\lambda \dot{\rho} - H(\rho, \hat{\rho}_\lambda)\right]$$

$$= \sup_{\rho,\hat{\rho}_\lambda} \int_0^1 \mathrm{d}x\, H(\rho, \hat{\rho}_\lambda)$$

Maximize energy

In sum -

To calculate the generating function

$$\langle e^{TL\lambda J}\rangle \sim \exp[TL\Psi(\lambda)]$$

Look for *time-independent solutions (with bc) of*

$$\partial_t \rho = \frac{\delta}{\delta \hat{\rho}_\lambda}\int_0^1 \mathrm{d}x\, H(\rho,\hat{\rho}_\lambda) = \partial_x\left[D(\rho)\partial_x\rho - \sigma(\rho)\partial_x\hat{\rho}_\lambda\right] \qquad = 0$$

$$\partial_t \hat{\rho}_\lambda = -\frac{\delta}{\delta \rho}\int_0^1 \mathrm{d}x\, H(\rho,\hat{\rho}_\lambda) = -\partial_x\left[D(\rho)\partial_x\hat{\rho}_\lambda\right] - \frac{\sigma'(\rho)}{2}\left(\partial_x\hat{\rho}_\lambda\right)^2 \qquad = 0$$

Result -

Typical density and noise profile which realize the fluctuations

We are focused on *looking for singularities* (when? where?)

Comments:

1. Prior to this work phase transitions in current large deviations were constrained to cases where the *additivity principle was broken*.
   (non-stationary optimal profile)

2. For our transitions can prove that the additivity principle holds

3. Condition for applicability of additivity principle to hold
   Shpielberg & Akkermans, PRL **116**, 240603 (2016)

Next - *Show that transitions can occur*

derive Landau theory for transitions

to make discussion easier break into different types

- Symmetry breaking transitions (continuous)
- First order phase transitions
- For each case identify microscopic model
- DERIVE IN EQUILIBRIUM AND THEN DISCUSS WHAT WE KNOW OUT OF EQUILIBRIUM

Note, transitions occur even in equilibrium

where, say, density large-deviation is smooth

Symmetry breaking phase transitions

To observe symmetry breaking transition need an underlying symmetry

Particle-Hole symmetry (about, say, $\rho = \bar{\rho} = 1/2$)

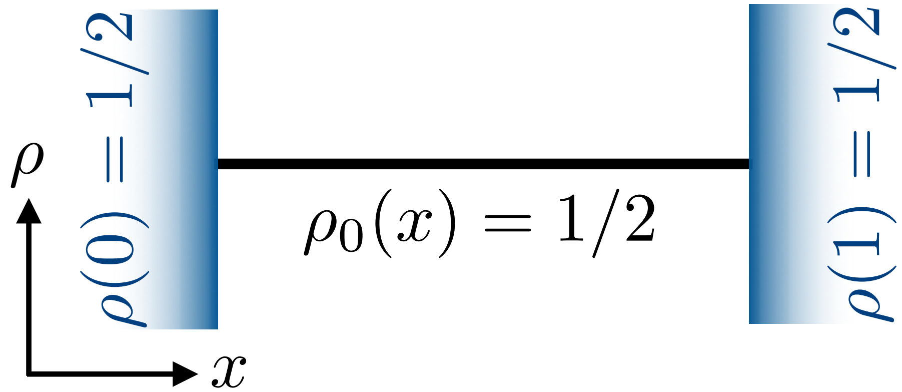
$$D(1/2 - \delta\rho) = D(1/2 + \delta\rho)$$

$$\sigma(1/2 - \delta\rho) = \sigma(1/2 + \delta\rho)$$

recall: consider boundary conditions
at equilibrium point

Consider possible solutions

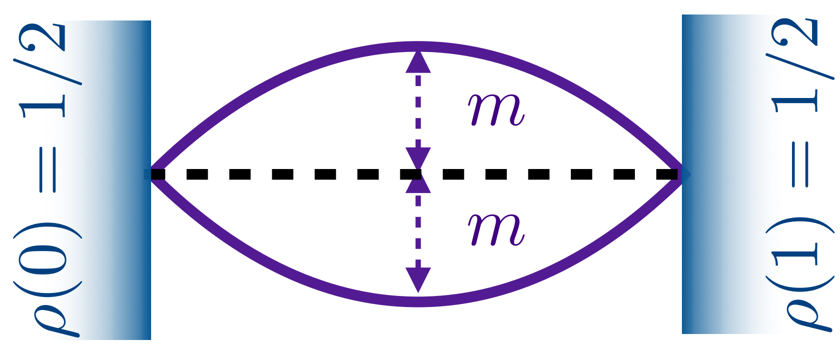
One solution - symmetric profile (bc obey symmetry)



denote this solution

$$\rho_0(x), \hat{\rho}_{\lambda,0}(x)$$

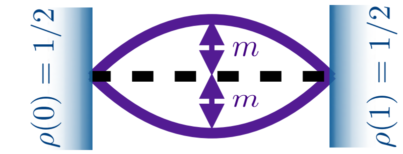
Near transition (if one occurs)
can imagine a deviation whose longest wave length component is



$$\rho_m(x) = 1/2 + m\sin(\pi x)$$

If they occur must be in pairs - *symmetry-breaking profiles*

denote this solution $\rho_m(x), \hat{\rho}_{\lambda,m}(x)$



$\rho_m(x) = 1/2 + m\sin(\pi x)$

With this in mind calculate

Landau theory (expansion in m, skipping details)

$$\mathcal{L}_\lambda(m) = \int_0^1 \mathrm{d}x\,[H(\rho_0,\hat\rho_{\lambda,0}) - H(\rho_m,\hat\rho_{\lambda,m})]$$

Then the scaled CGF

$$\Psi(\lambda) = \sup_{\rho,\hat\rho_\lambda}\int_0^1 \mathrm{d}x\,H(\rho,\hat\rho_\lambda) = \int_0^1 \mathrm{d}x\,H(\rho_0,\hat\rho_{\lambda,0}) - \inf_m \mathcal{L}_\lambda(m)$$

FIND TO LEADING ORDER

$$\mathcal{L}_\lambda(m) = -\frac{\lambda_c\bar\sigma''}{4}\,\delta\lambda\, m^2 + \frac{\pi^2\bar D(4\bar D''\bar\sigma'' - \bar D\bar\sigma^{(4)})}{64\bar\sigma\bar\sigma''}m^4$$

$\delta\lambda = \lambda - \lambda_c$ and $\lambda_c = \pm\sqrt{\dfrac{2\pi^2\bar D}{\bar\sigma\bar\sigma''}}$

To have transition

Condition 1 $\bar\sigma'' > 0$

Condition 2 $4\bar D''\bar\sigma'' > \bar D\bar\sigma^{(4)}$

TO HAVE A TRANSITION NEED A MODEL WITH A LOCAL MINIMA IN σ

Recap -

Landau theory shows a symmetry-breaking transitions when

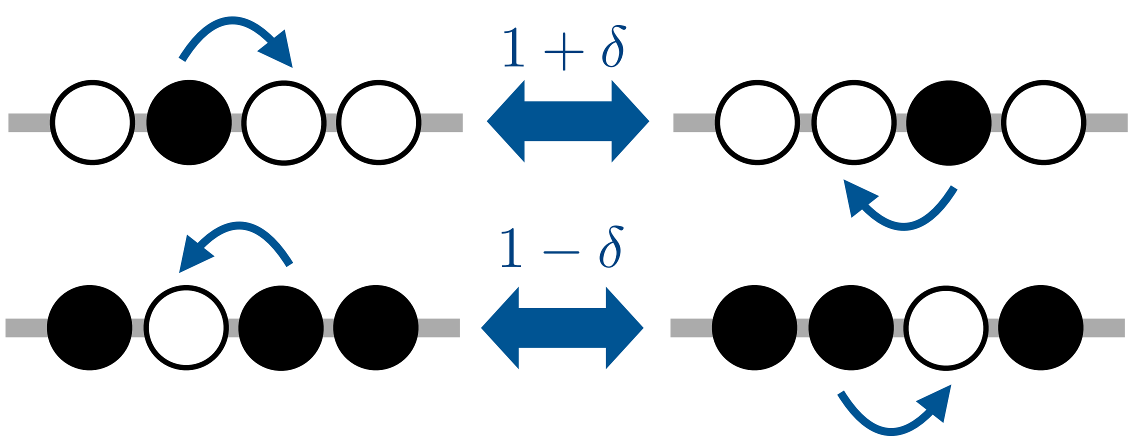
1. **Particle-hold symmetry (in b.c. and model)**
2. **mobility σ at this point has a local minima**

Microscopic model

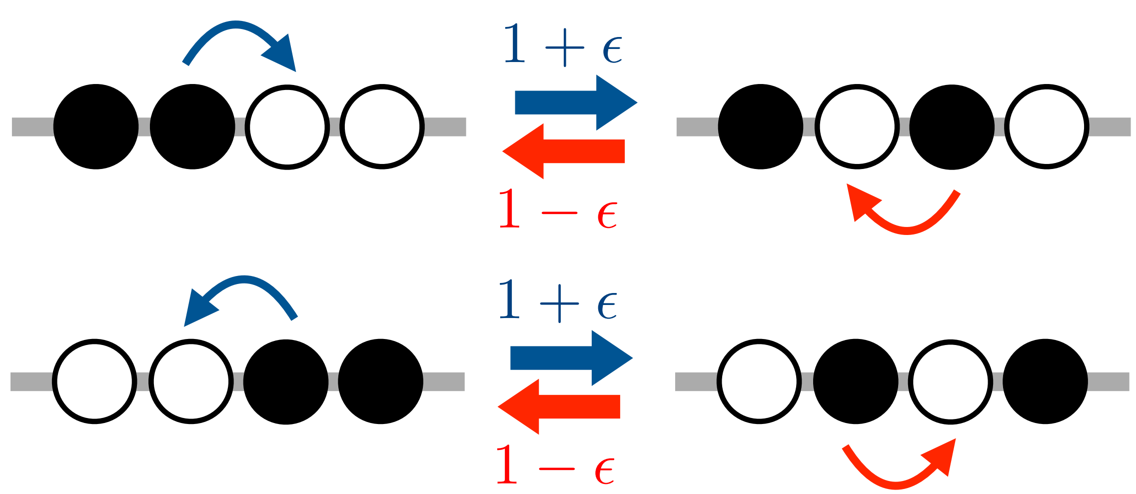
KLS model

Katz, Lebowitz, Spohn, JSP (1984)

Microscopic dynamics

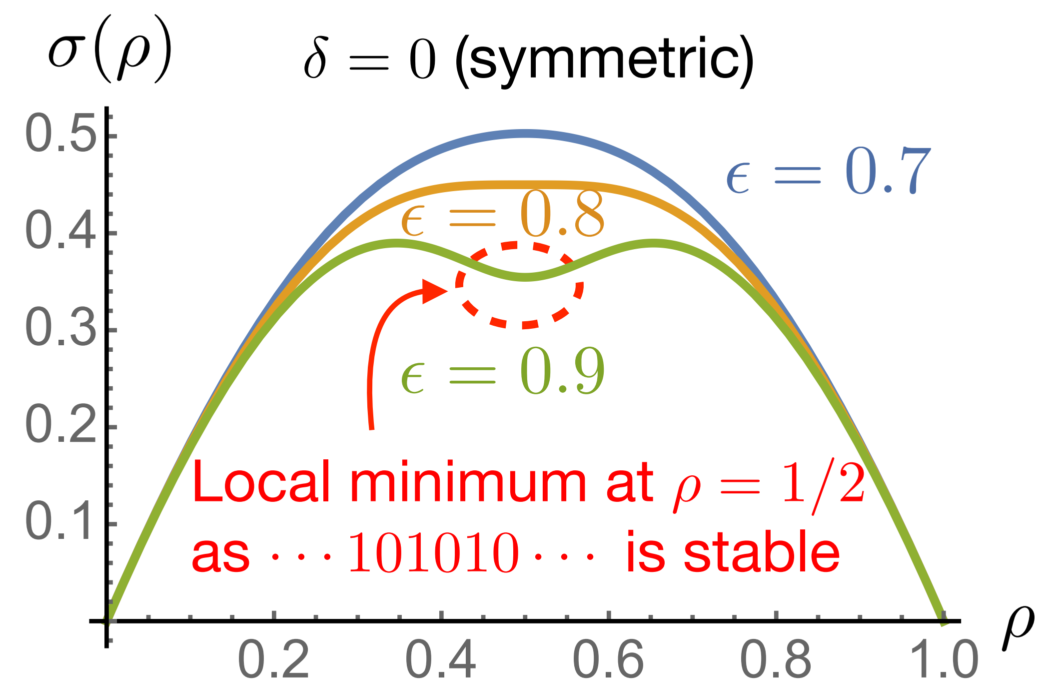


$\delta > 0:$ Particles faster than holes



$\epsilon > 0:$ Short-range repulsion

Transport coefficients

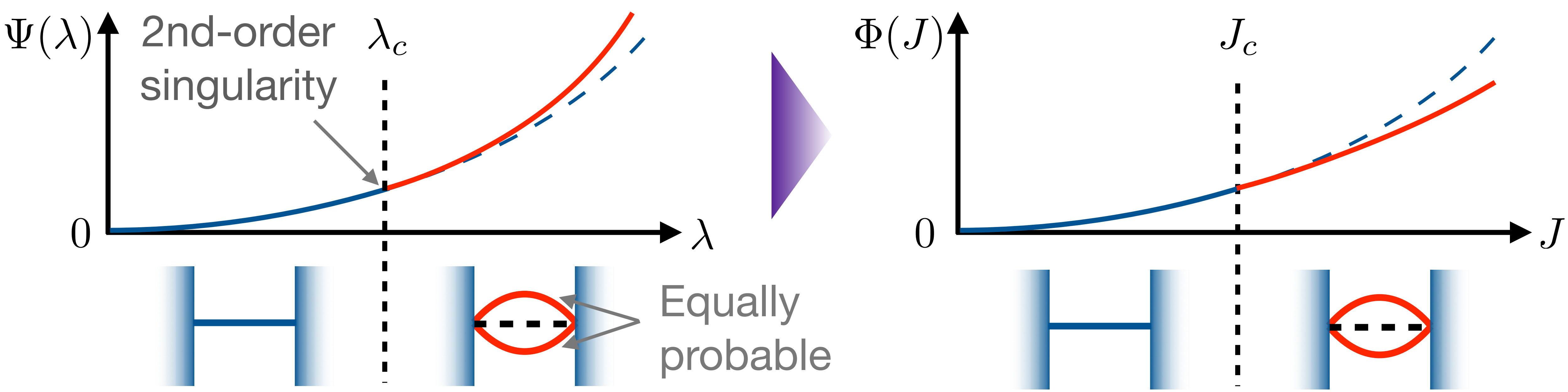


Hager et al., PRE **63**, 056110 (2001); Krapivsky (unpublished)

SUMMARY OF SYMMETRY BREAKING TRANSITION

- System with local minima of σ at `symmetric' point $\bar{\rho}$
- Up to now in equilibrium
- Results unchanged to leading order for boundary conditions

$$\rho(0) = 1/2 + \delta\rho \qquad \rho(1) = 1/2 - \delta\rho$$



First order phase transitions

Now - models with **no particle-hole symmetry**
again **in equilibrium at minima of** σ

Landau theory (exactly along the lines outlined before)

$$\mathcal{L}_\lambda(m) = -\frac{\lambda_c \bar{\sigma}''}{4}\,\delta\lambda\, m^2 - \frac{2\pi \bar{D}(\bar{D}\bar{\sigma}^{(3)} - 3\bar{D}'\bar{\sigma}'')}{9\bar{\sigma}\bar{\sigma}''} m^3 + \frac{\pi^2 \bar{D}(4\bar{D}''\bar{\sigma}'' - \bar{D}\bar{\sigma}^{(4)})}{64\bar{\sigma}\bar{\sigma}''} m^4$$

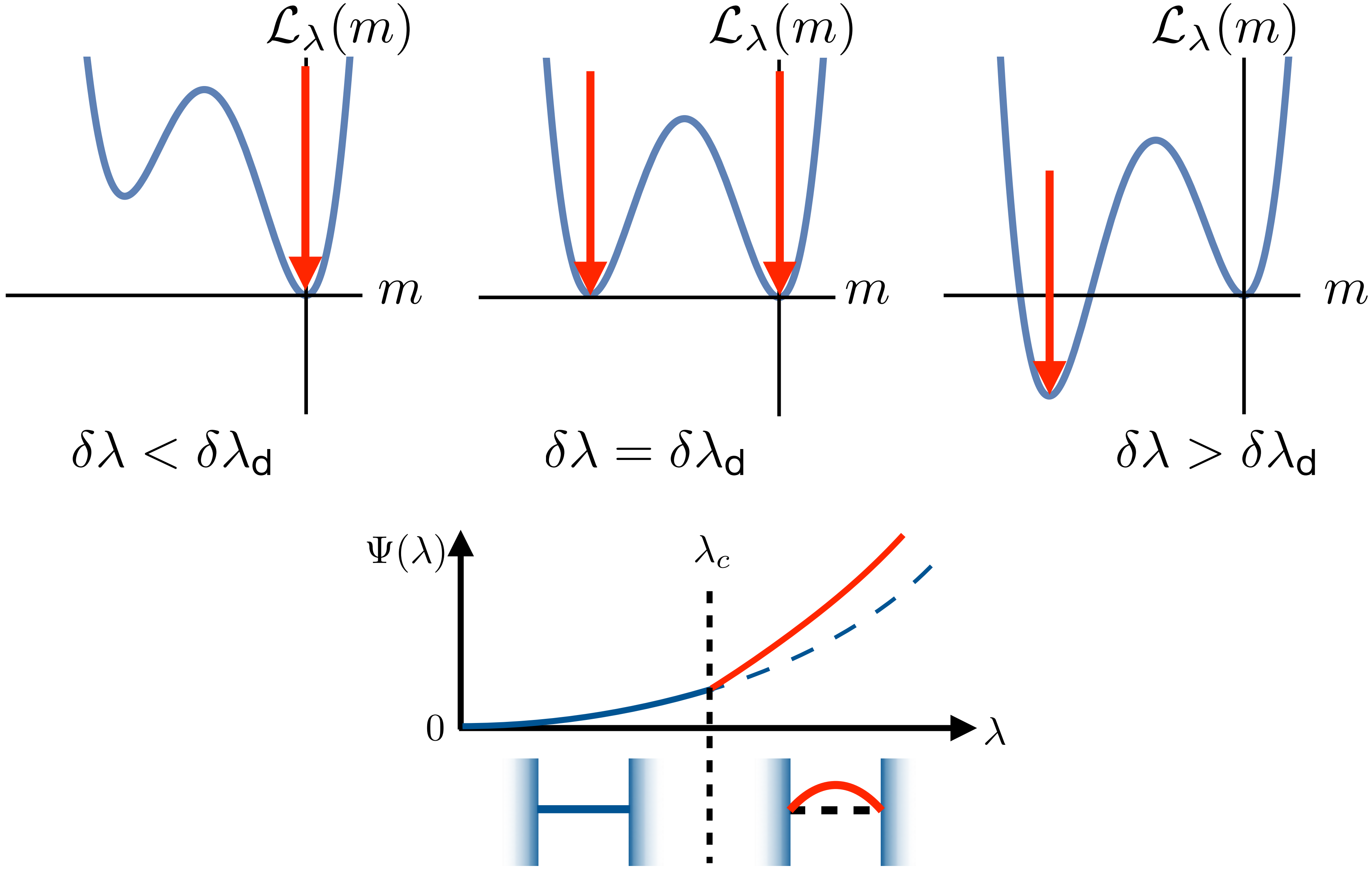
To have transition

Condition 1 $\bar{\sigma}'' > 0$

Condition 2 $\bar{D}\bar{\sigma}^{(3)} \neq 3\bar{D}'\bar{\sigma}''$

Condition 3 $4\bar{D}''\bar{\sigma}'' > \bar{D}\bar{\sigma}^{(4)}$

$$\mathcal{L}_\lambda(m) = -\frac{\lambda_c\bar{\sigma}''}{4}\,\delta\lambda\, m^2 - \frac{2\pi\bar{D}(\bar{D}\bar{\sigma}^{(3)} - 3\bar{D}'\bar{\sigma}'')}{9\bar{\sigma}\bar{\sigma}''}m^3 + \frac{\pi^2\bar{D}(4\bar{D}''\bar{\sigma}'' - \bar{D}\bar{\sigma}^{(4)})}{64\bar{\sigma}\bar{\sigma}''}m^4$$

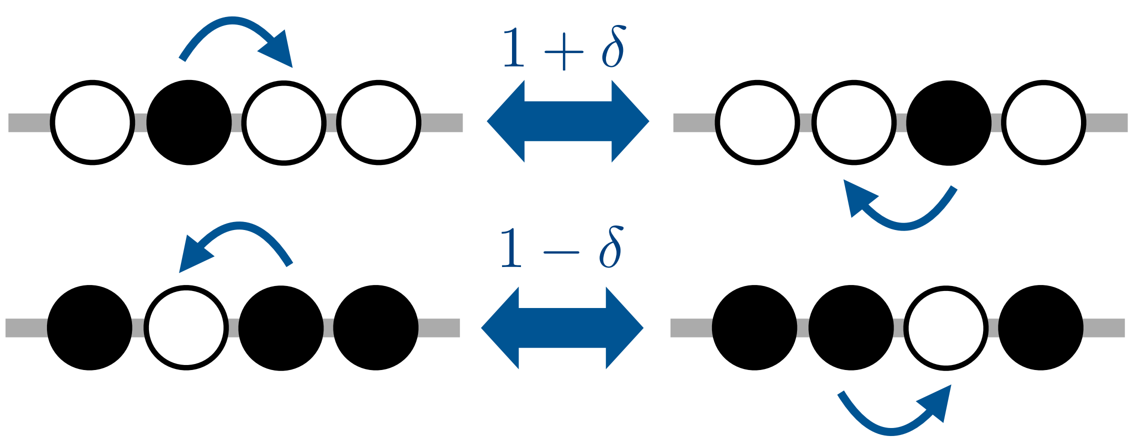


Microscopic model

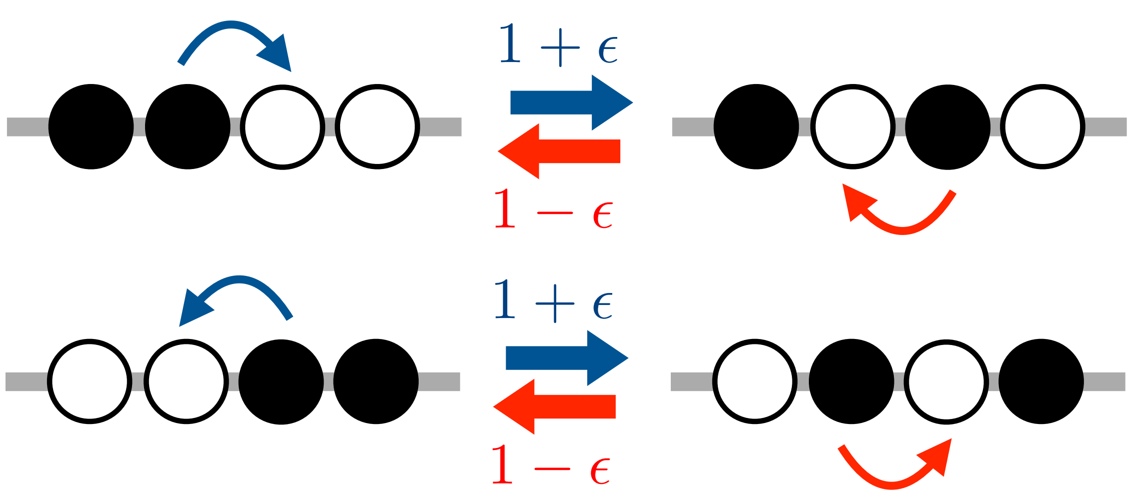
KLS model

Katz, Lebowitz, Spohn, JSP (1984)

Microscopic dynamics

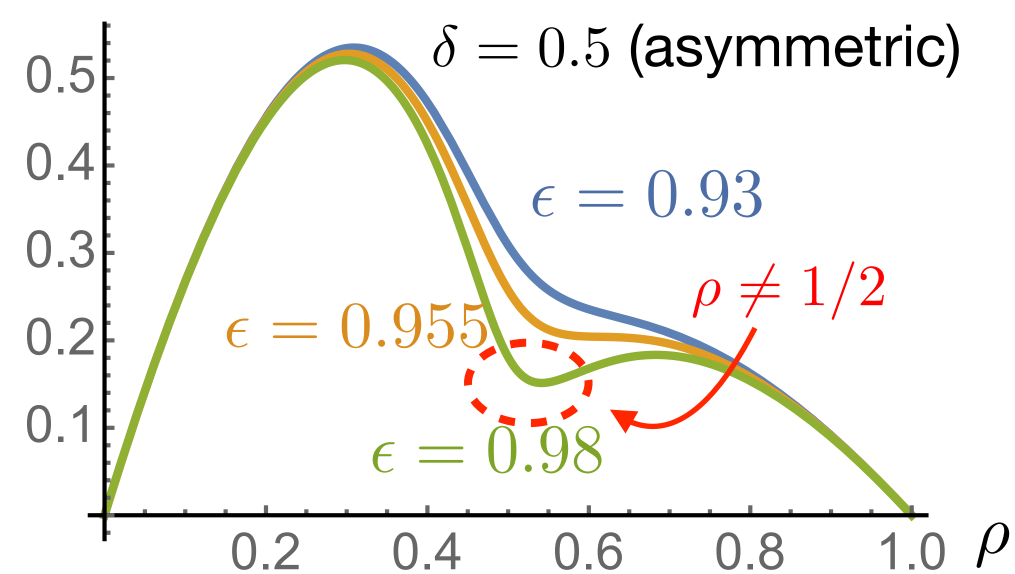


$\delta > 0:$ Particles faster than holes



$\epsilon > 0:$ Short-range repulsion

Transport coefficients

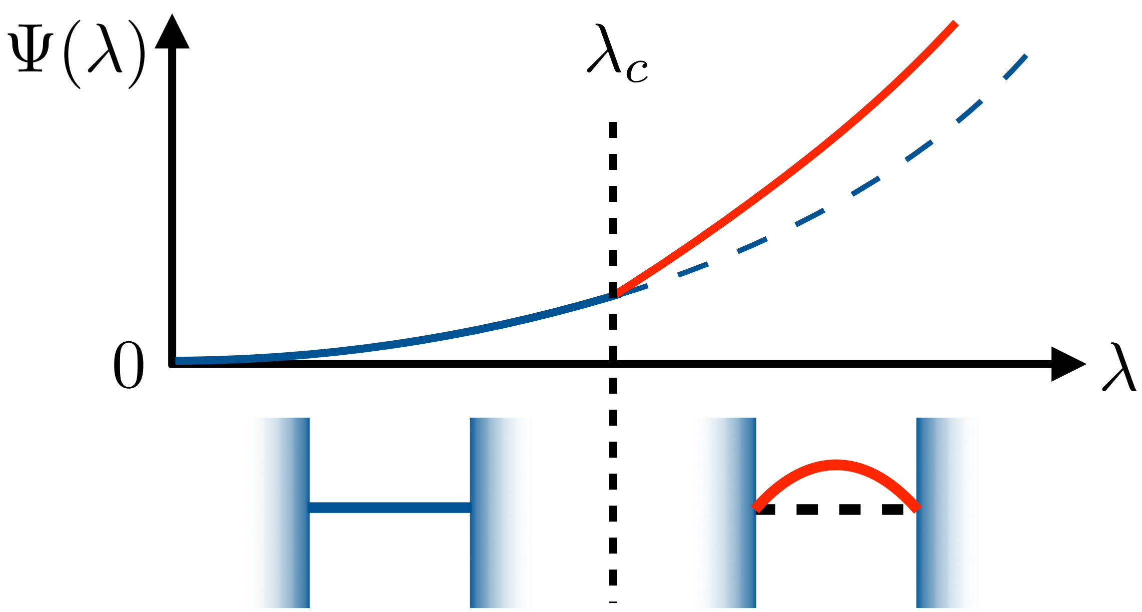


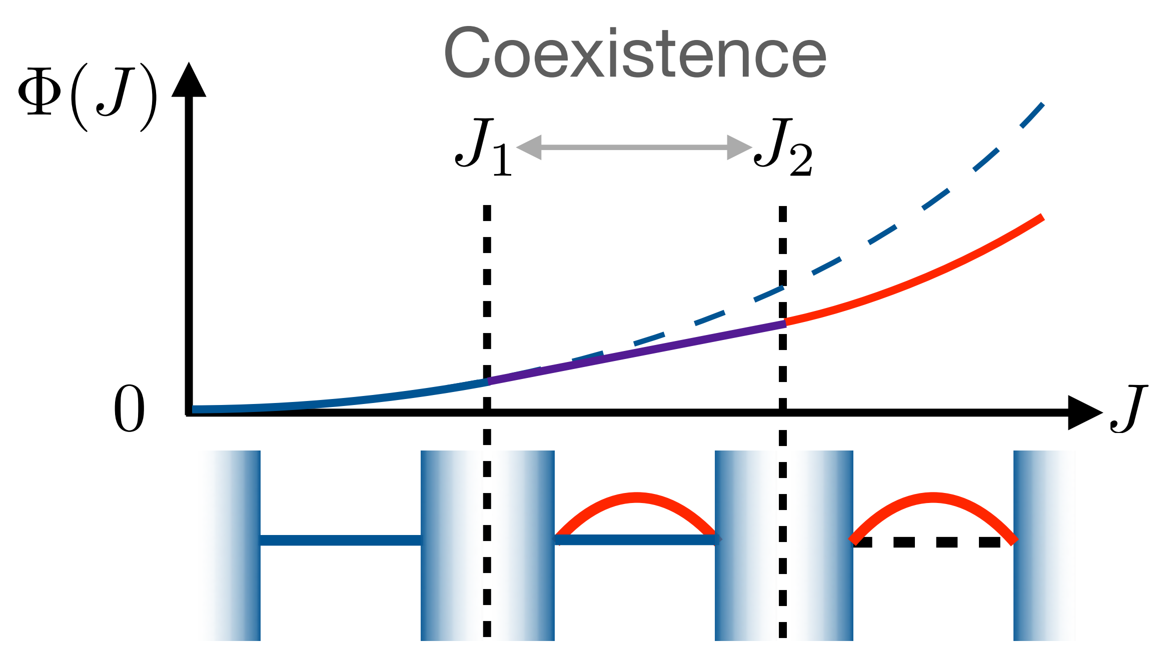
Hager et al., PRE **63**, 056110 (2001); Krapivsky (unpublished)

SUMMARY OF FIRST ORDER TRANSITIONS

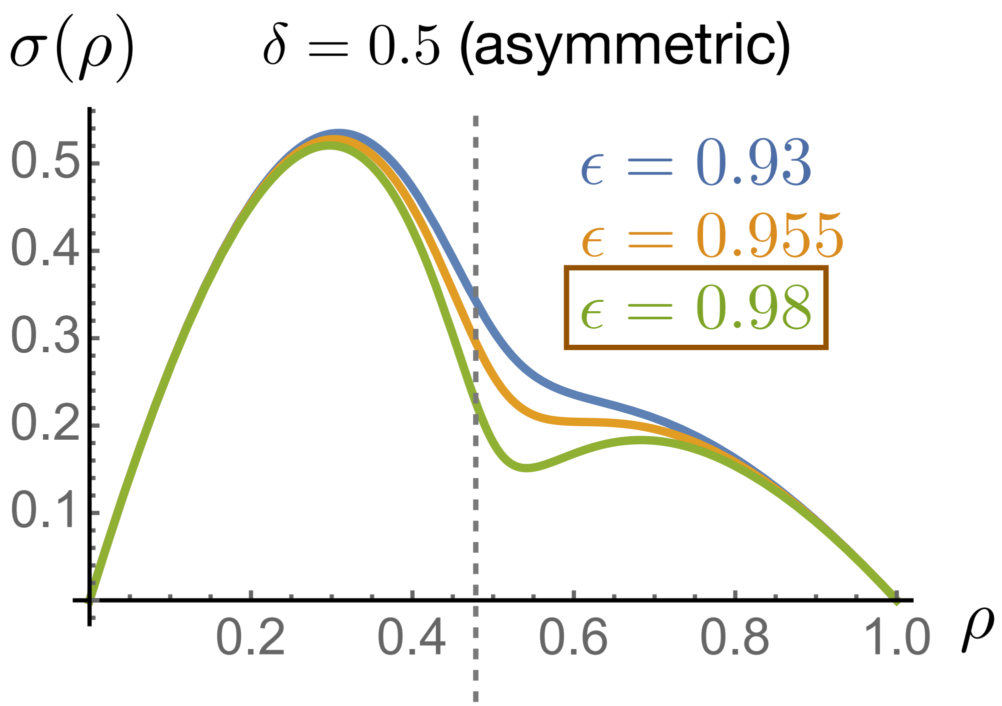
- System with local minima of σ at `symmetric' point $\bar{\rho}$
- Up to now in equilibrium
- Results unchanged to leading order for boundary conditions

$$\rho(0) = 1/2 + \delta\rho \qquad\qquad \rho(1) = 1/2 - \delta\rho$$





What happens when not at minima of σ ?



$$\rho(0) = \rho(1) = \bar{\rho} \neq 1/2$$
$$\bar{\sigma}' \neq 0$$

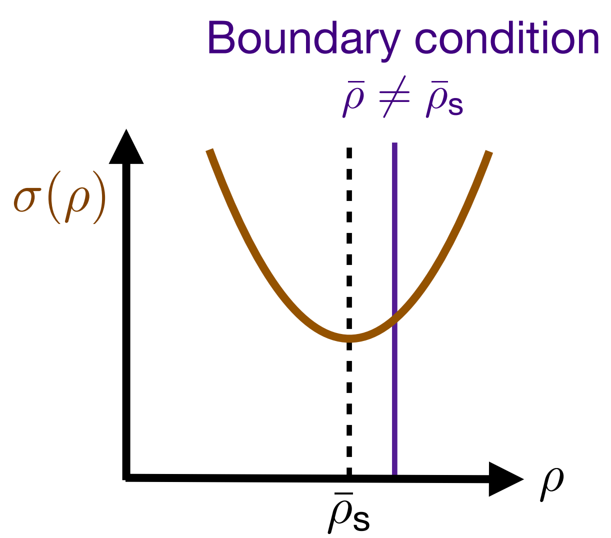
Landau theory

$$\mathcal{L}_\lambda(m) = -\frac{2\pi \bar{D}^2}{\bar{\sigma}\bar{\sigma}''}\bar{\sigma}' m - \frac{\lambda_c \bar{\sigma}''}{4}\,\delta\lambda\, m^2 - \frac{2\pi \bar{D}(\bar{D}\bar{\sigma}^{(3)} - 3\bar{D}'\bar{\sigma}'')}{9\bar{\sigma}\bar{\sigma}''} m^3 + \frac{\pi^2 \bar{D}(4\bar{D}''\bar{\sigma}'' - \bar{D}\bar{\sigma}^{(4)})}{64\bar{\sigma}\bar{\sigma}''} m^4$$

σ' **acts as a \`magnetic-field' killing the transitions**

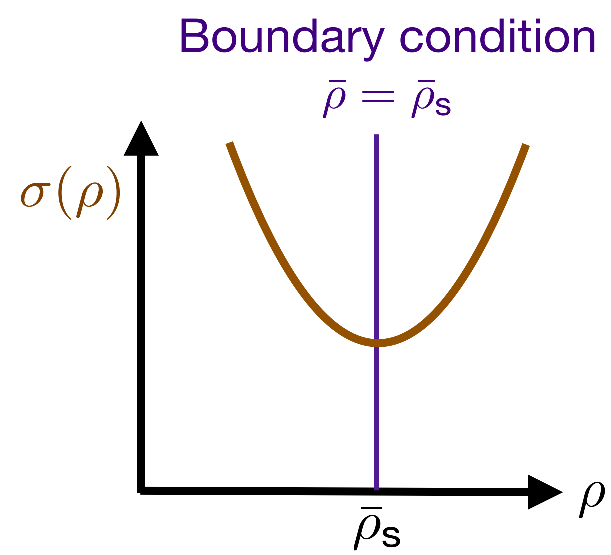
SUMMARY UP TO HERE

No DPTs



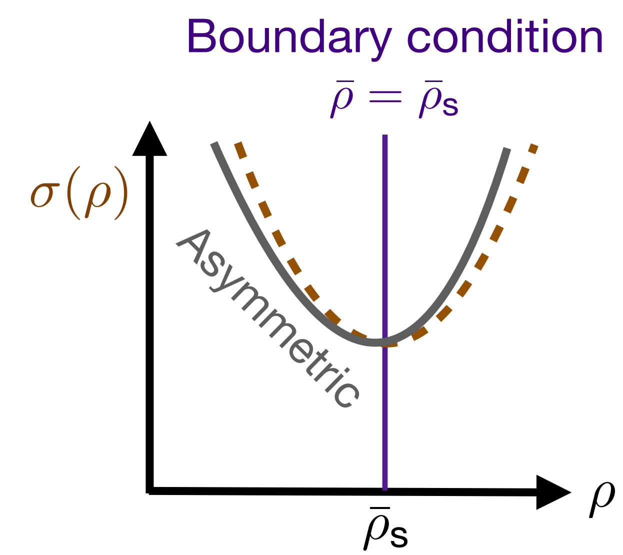
$\bar\sigma' \neq 0$

Symmetry breaking



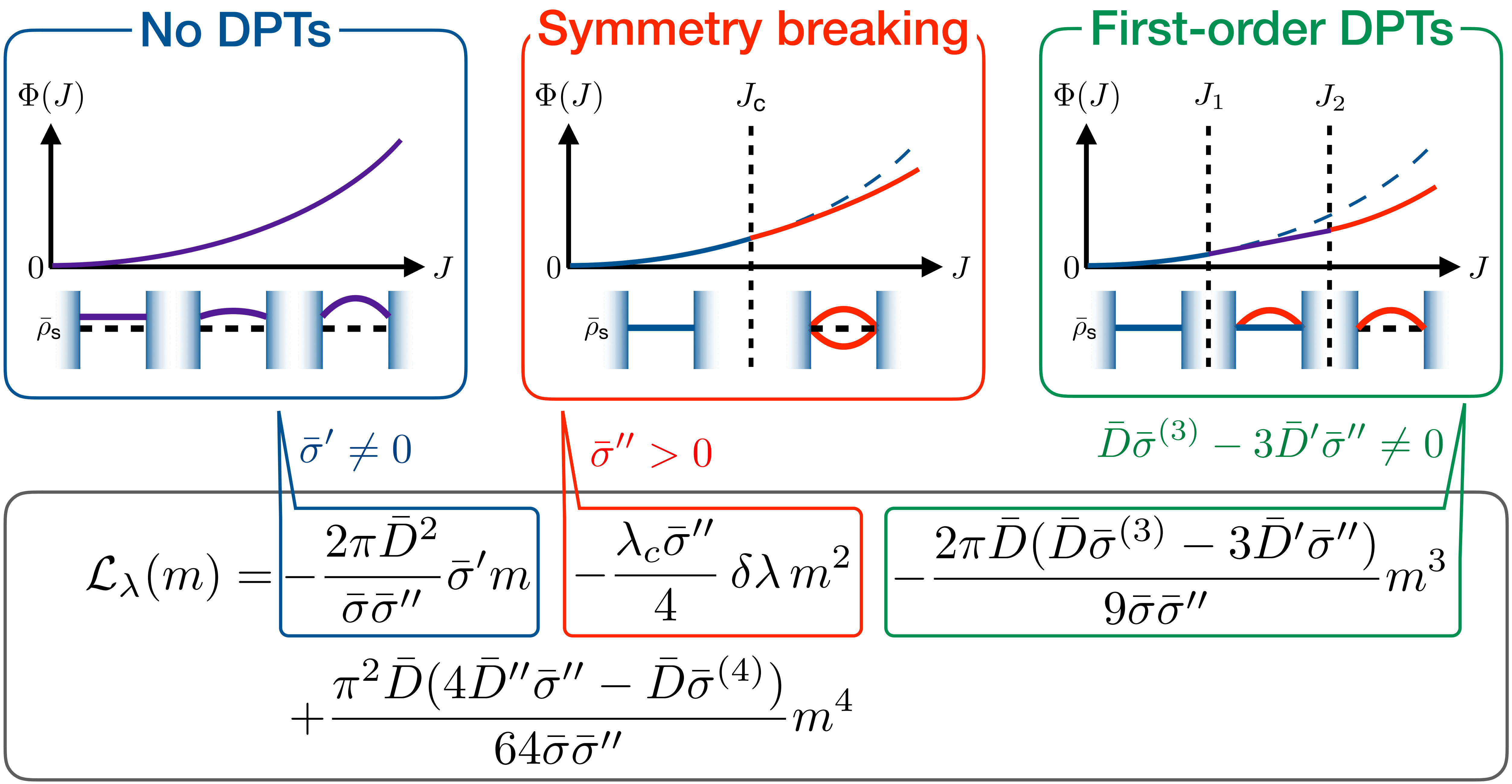
$\bar\sigma'' > 0$

First-order DPTs



$\bar D \bar\sigma^{(3)} - 3\bar D' \bar\sigma'' \neq 0$

$$\mathcal{L}_\lambda(m) = -\frac{2\pi \bar D^2}{\bar\sigma \bar\sigma''}\bar\sigma' m - \frac{\lambda_c \bar\sigma''}{4}\delta\lambda\, m^2 - \frac{2\pi \bar D(\bar D \bar\sigma^{(3)} - 3\bar D' \bar\sigma'')}{9\bar\sigma\bar\sigma''}m^3 + \frac{\pi^2 \bar D(4\bar D''\bar\sigma'' - \bar D \bar\sigma^{(4)})}{64\bar\sigma\bar\sigma''}m^4$$

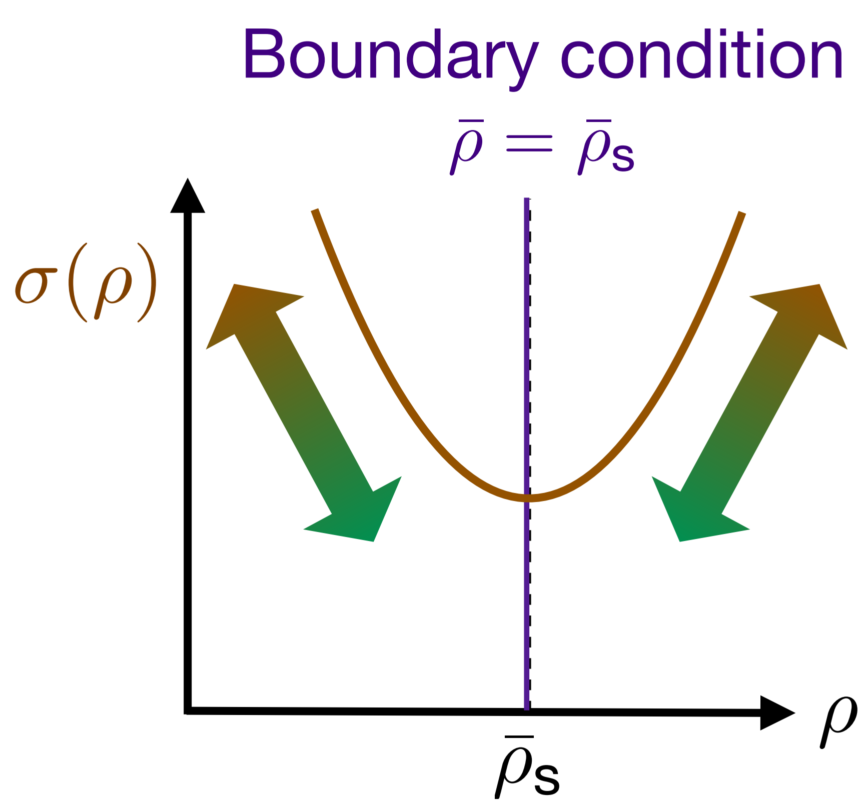

No DPTs
Symmetry breaking
First-order DPTs
$\Phi(J)$
J
0
$\bar{\rho}_{\mathrm{s}}$
J_{c}
J_1
J_2
$\bar{\sigma}' \neq 0$
$\bar{\sigma}'' > 0$
$\bar{D}\bar{\sigma}^{(3)} - 3\bar{D}'\bar{\sigma}'' \neq 0$
$\mathcal{L}_\lambda(m) = -\frac{2\pi \bar{D}^2}{\bar{\sigma}\bar{\sigma}''}\bar{\sigma}' m - \frac{\lambda_c \bar{\sigma}''}{4}\delta\lambda\, m^2 - \frac{2\pi \bar{D}(\bar{D}\bar{\sigma}^{(3)} - 3\bar{D}'\bar{\sigma}'')}{9\bar{\sigma}\bar{\sigma}''} m^3 + \frac{\pi^2 \bar{D}(4\bar{D}''\bar{\sigma}'' - \bar{D}\bar{\sigma}^{(4)})}{64\bar{\sigma}\bar{\sigma}''} m^4$

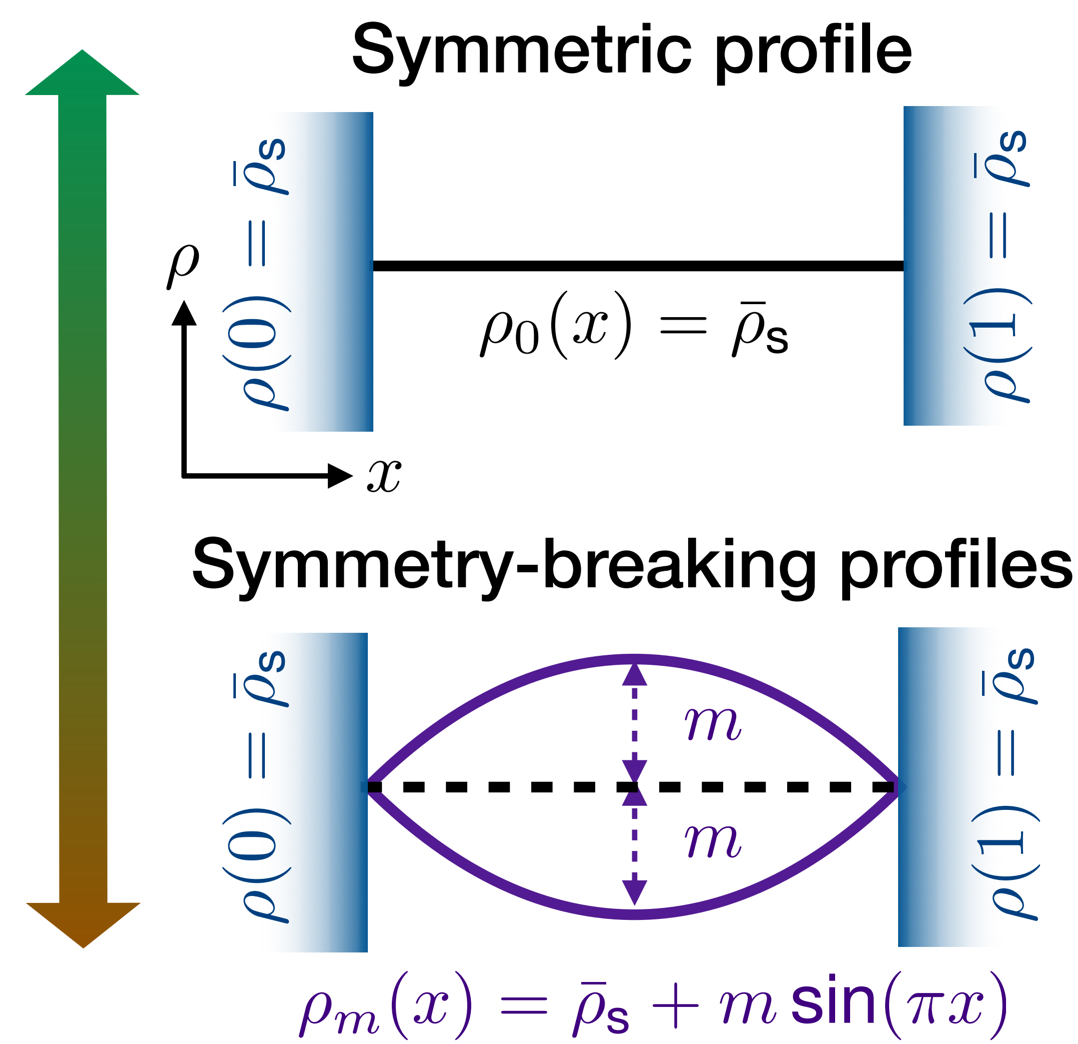
Physical Intuition

$$\partial_t \rho = -\partial_x \left[-D(\rho)\, \partial_x \rho + \sqrt{\sigma(\rho)}\, \eta \right]$$

Diffusion favors flat profile

High J is easier if $\sigma(\rho)$ is high





Lagrangian picture

$$\Phi(J) = \inf_{\rho} \int_0^1 \mathrm{d}x \, \frac{\left[J + D(\rho)\partial_x \rho - \sigma(\rho)E\right]^2}{2\sigma(\rho)} \, .$$

Expand in m

$$\Phi(J) \simeq \frac{\delta J^2}{2\bar{\sigma}} + \inf_m \left[\left(\frac{\bar{D}^2}{2} - \frac{\bar{\sigma}'' \, \delta J^2}{4\bar{\sigma}} \right) m^2 + O(m^4) \right]$$

numerator **denominator**

With $\delta J \equiv J - \bar{\sigma} E$

DENOMINATOR WINS FOR LARGE ENOUGH δJ

Effect of Bulk Field

So far the possibility of a bulk field (with diffusive scaling) was ignored.

Including a bulk field gives the following dynamical equation for the density

$$\partial_t \rho = -\partial_x \left[-D(\rho)\, \partial_x \rho + \sqrt{\sigma(\rho)}\, \eta + \sigma(\rho) E \right]$$

REPEAT SAME ANALYSIS AS BEFORE

Landau theory

$$\mathcal{L}(m) \simeq -\frac{2\pi\bar{D}^2}{\bar{\sigma}\bar{\sigma}''}\,\bar{\sigma}'\,m - \frac{(\lambda_c + E)\bar{\sigma}''}{4}\,\delta\lambda\,m^2 - \frac{2\pi\bar{D}(\bar{D}\bar{\sigma}^{(3)} - 3\bar{D}'\,\bar{\sigma}'')}{9\bar{\sigma}\bar{\sigma}''}\,m^3$$
$$+\left[\frac{\pi^2\,\bar{D}\left(4\bar{D}''\,\bar{\sigma}'' - \bar{D}\bar{\sigma}^{(4)}\right)}{64\bar{\sigma}\bar{\sigma}''} + \frac{\bar{\sigma}''^2\,E^2}{64\bar{\sigma}}\right] m^4\,.$$

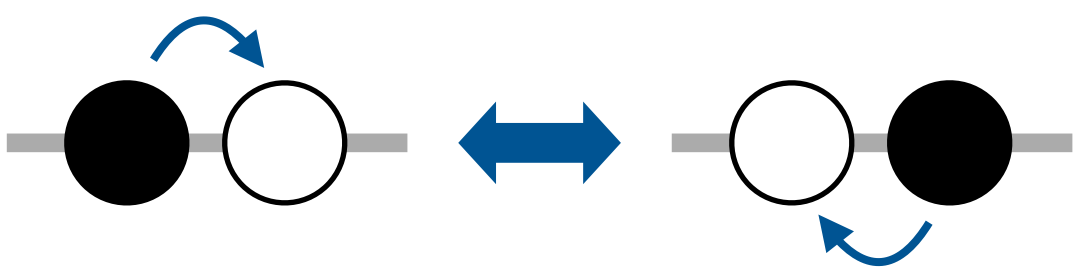
As long as $\bar{\sigma}' = 0$ even if not sitting at minima of σ

for large enough field E ***have a transition***

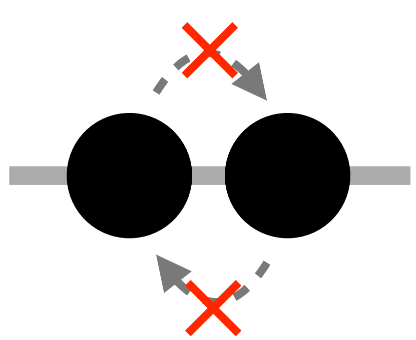
WASEP

Microscopic dynamics

Symmetric random walk



Exclusion

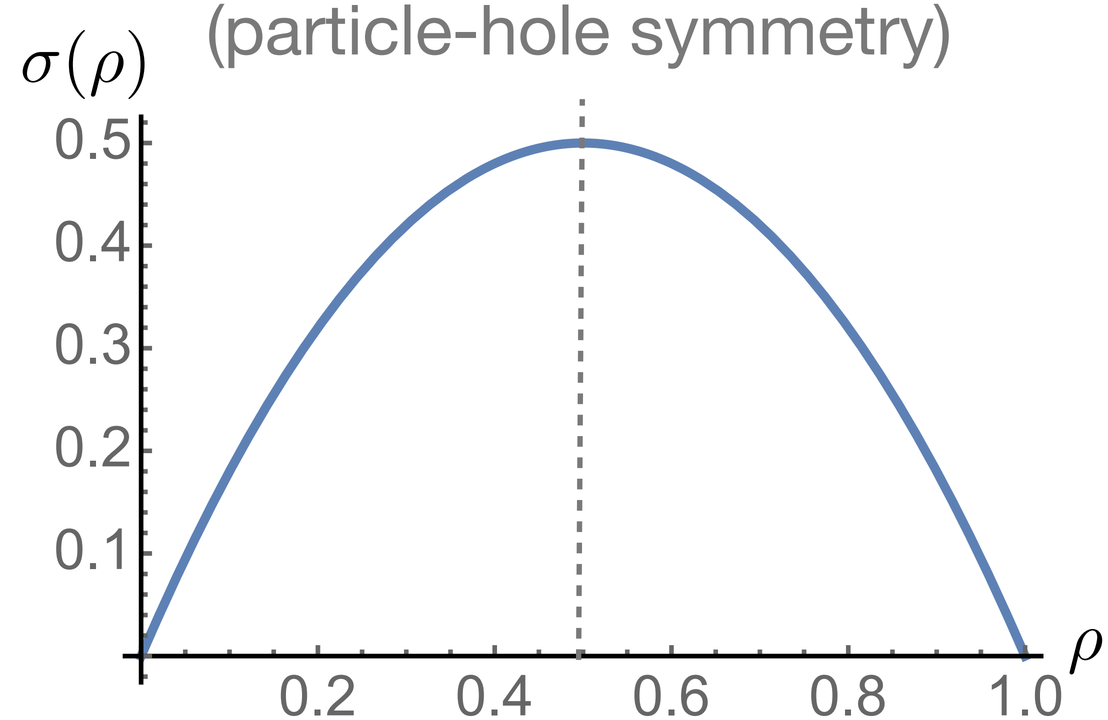


Transport coefficients

$$D(\rho) = 1$$

$$\sigma(\rho) = 2\rho(1-\rho)$$

Symmetric w.r.t.
$\rho - 1/2 \to 1/2 - \rho$
(particle-hole symmetry)



SUMMARY

- Current LDF of boundary driven diffusive systems can have singularities
- Identified models for difference scenarios - 1st, 2nd order
- Transitions *not* associated with breaking of additivity principle

Question - transition in boundary driven which breaks additivity?