

# Jets@ATLAS

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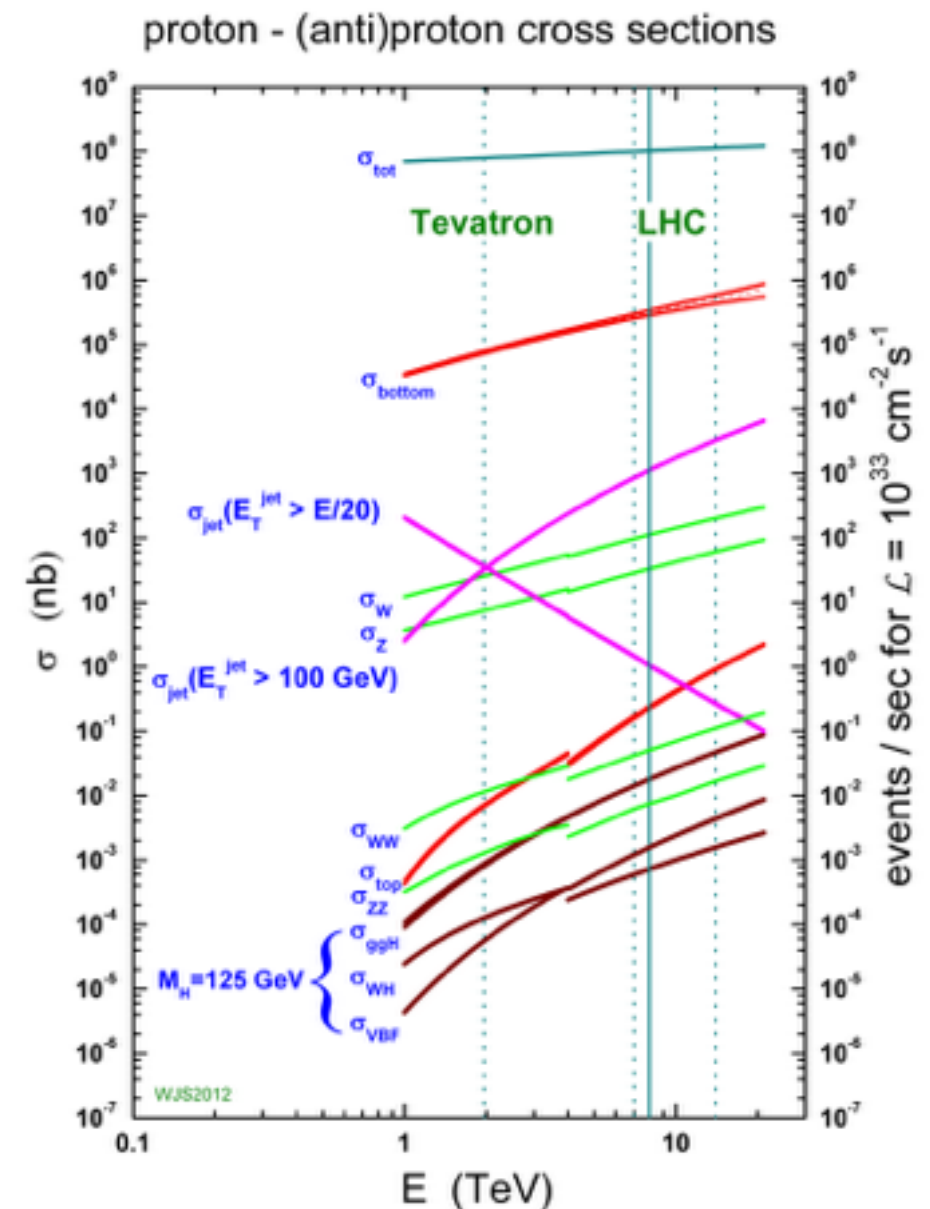


LHC is a jet factory!

As a signal, test of  
QCD predictions

Background for most analyses

Large-radius jets offer a  
powerful way to identify boosted  
hadronically decaying particles

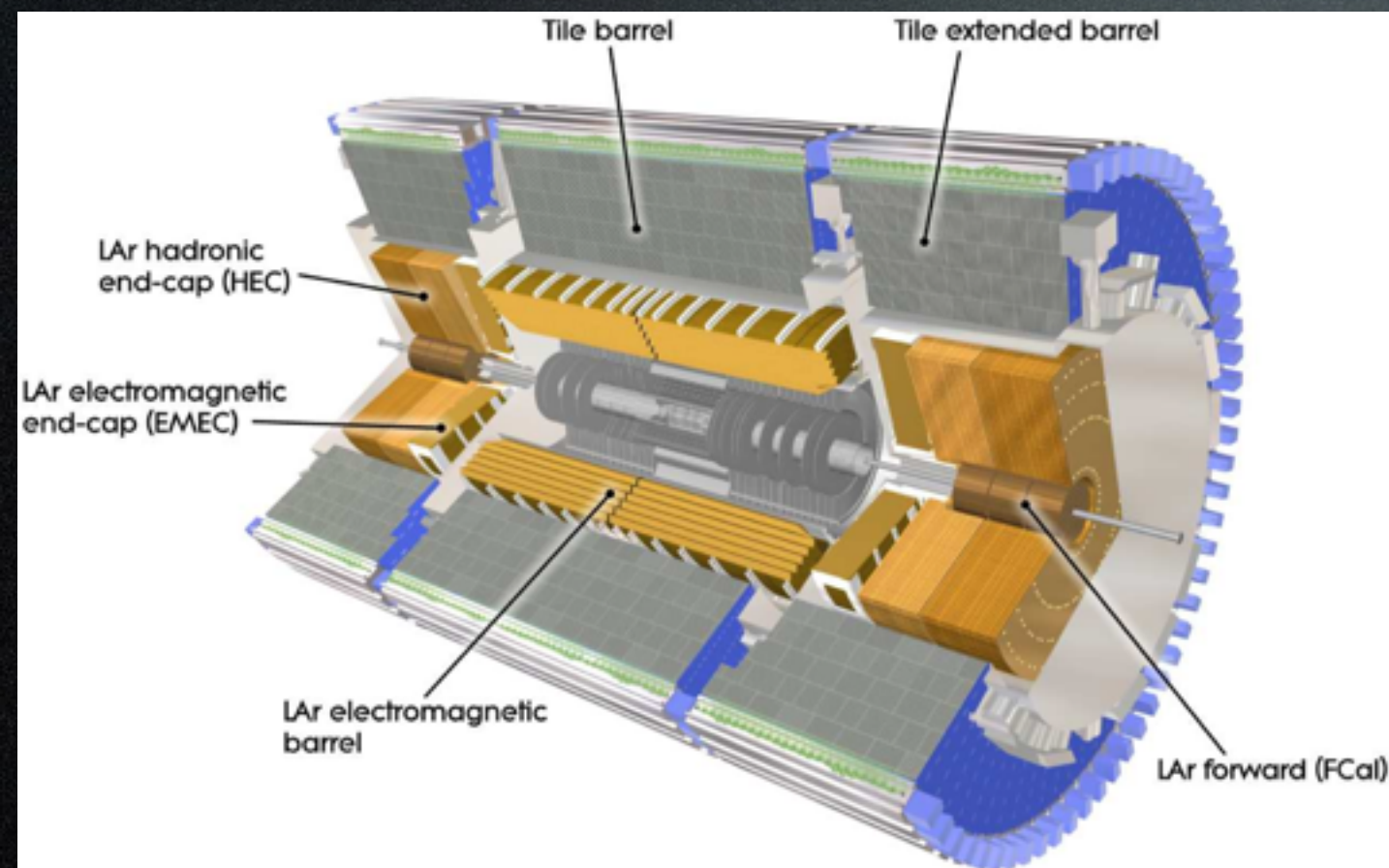




# Part 1: Jet making and Performance



# ATLAS Calorimeters



- High precision calorimetry
  - Highly granular electromagnetic calorimeter up to  $|\eta| < 3.2$
  - Hadronic tile calorimeter barrel and endcaps up to  $|\eta| < 3.2$
  - Forward calorimeters for  $3.2 < |\eta| < 4.9$ , granularity of  $\Delta\eta \times \Delta\phi \approx 0.2 \times 0.2$

Better jet energy resolution than CMS (worse momentum resolution in tracking)

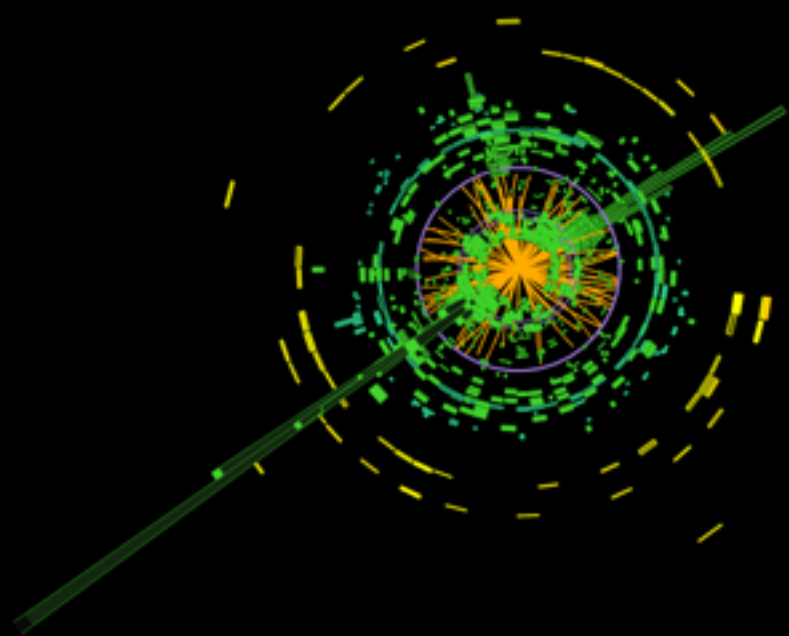




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Event: 690925592

2015-09-18 02:47:06 CEST



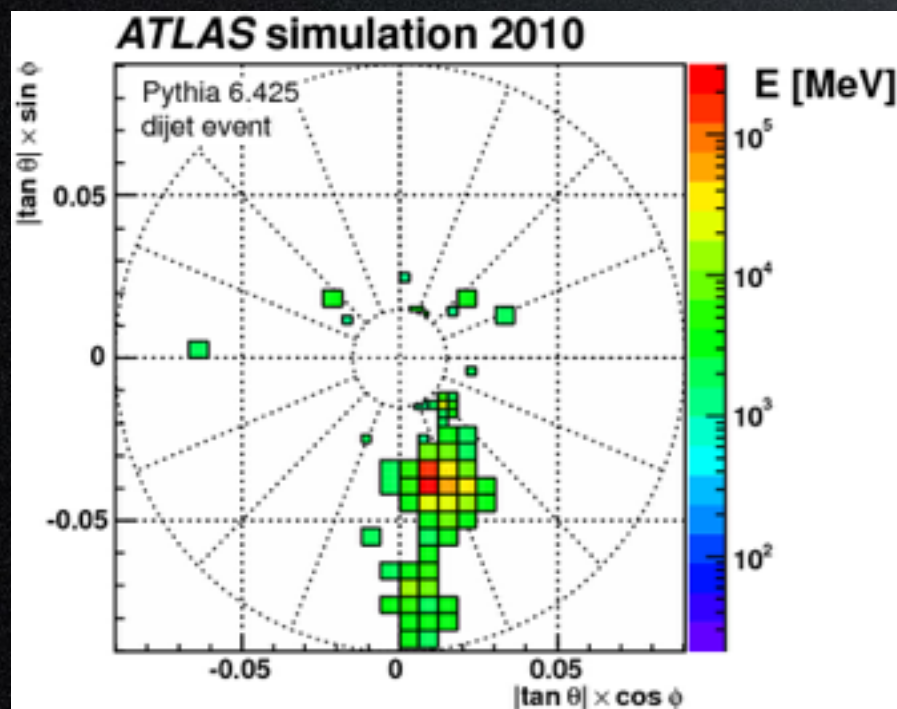
A high mass dijet system



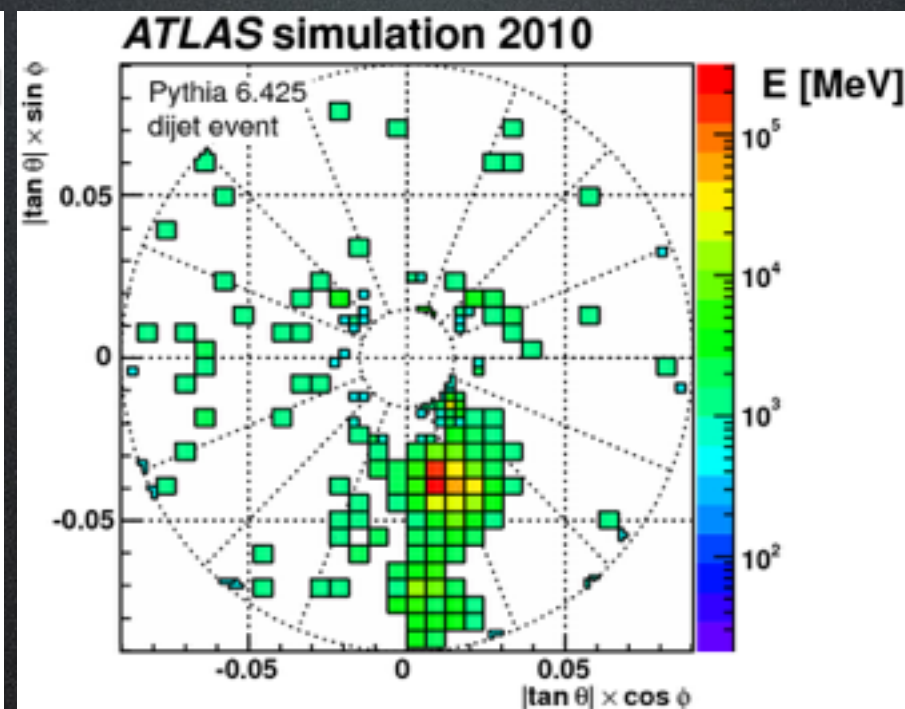
# Inputs

- 3 dimensional clusters of calorimeter cells, or topoclusters (treated massless)
- Attempt to reconstruct hadronic showers while suppressing noise (electronic+pileup)

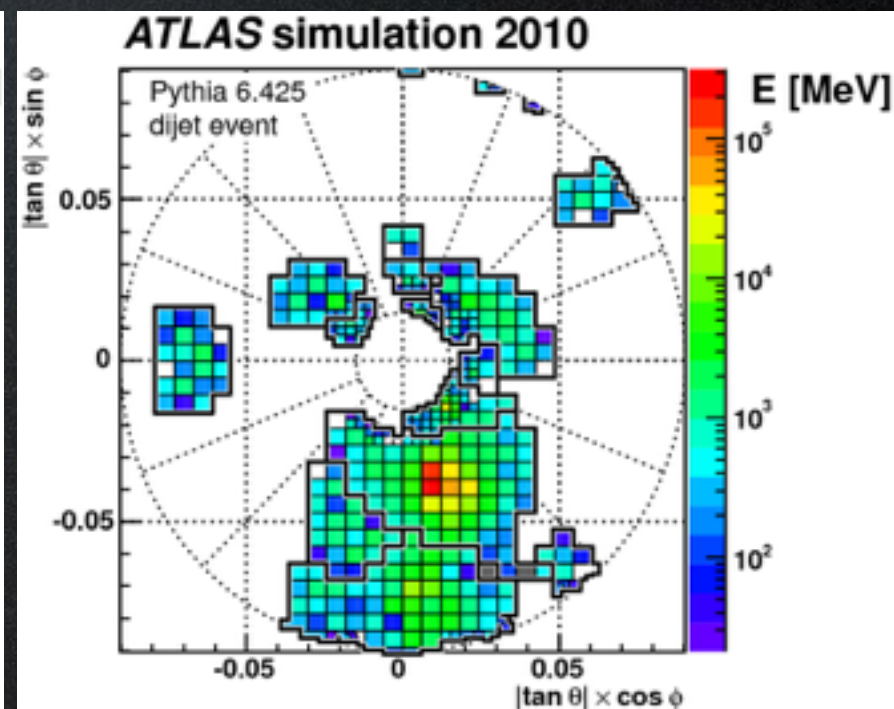
Seeded by: Built from all surrounding+ single layer of all adjacent cell (in 3D)



$E > 4\sigma_{\text{noise}}$



$E > 2\sigma_{\text{noise}}$

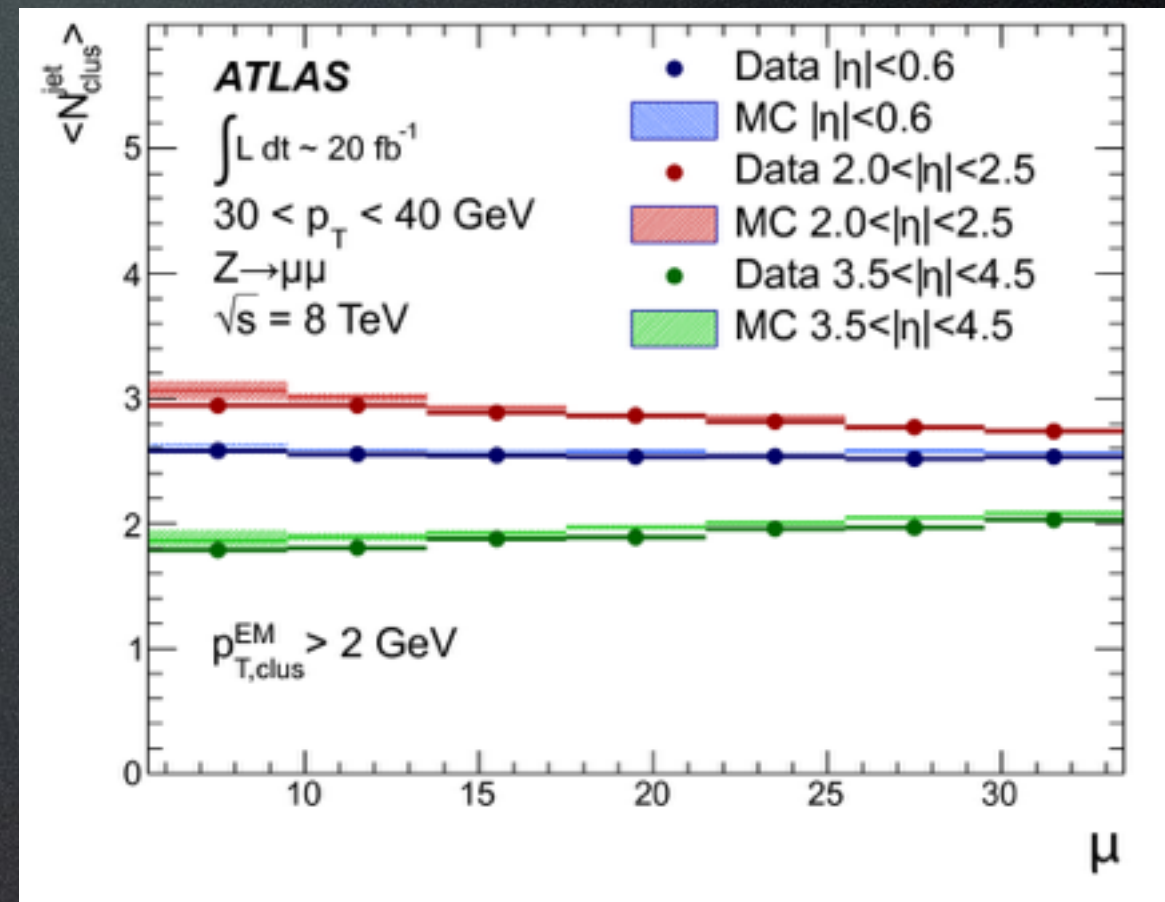


All



# Topoclusters

- Robust against pileup
- Topoclusters initially reconstructed at EM scale.

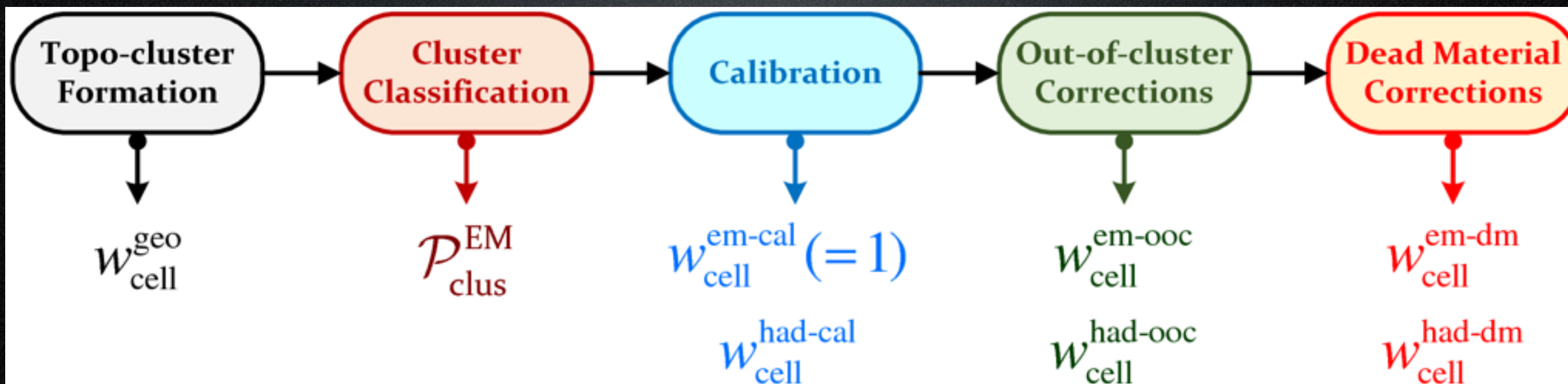




# Local Hadronic Cell Weighting

Then calibrate them to account for hadron response.

From charged hadron and neutral pion MC.





# Tracks

- Inner detector tracks can be ghost associated to jets, for pileup suppression and improving jet energy resolution (for ghost matching, tracks are treated as 4-vectors with infinitesimal magnitudes)
- Particle flow



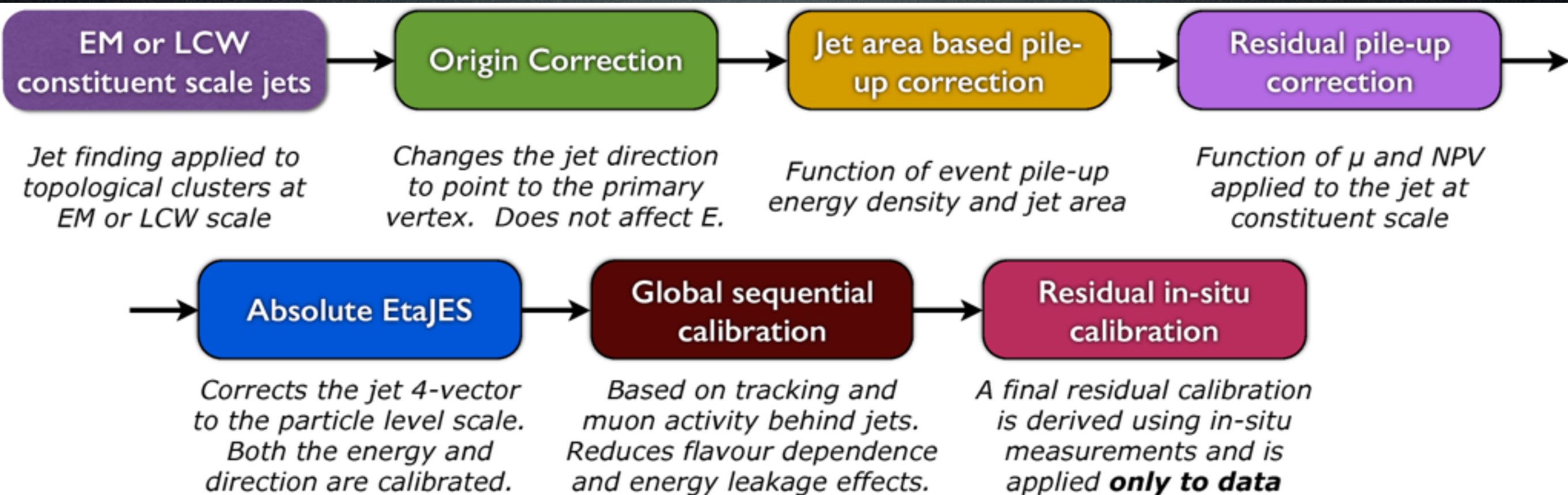


# Jets Need Calibration

1. **Calorimeter non-compensation:** correction for the different scales of the energy measured from hadronic and electromagnetic showers.
2. **Dead material:** energy lost in inactive areas of the detector.
3. **Leakage:** showers reaching the outer edge of the calorimeters.
4. **Out of calorimeter jet:** energy of particles which are included in the truth jet but which are not included in the reconstructed jet.
5. **Energy deposits below noise thresholds:** clusters are only formed by energy deposits which are well above the background noise. Therefore the correction is required to correct for particles that do not form clusters. Additionally some part of a shower may fall outside of the topological clusters such that this also needs to be corrected for.
6. **Pile-up:** energy deposition in jets is affected by the presence of multiple  $pp$  collisions in the same bunch crossing as well as residual signals from other bunch crossings.



# Jet Calibration



Reference is (stable) particle jets, without muon or neutrino



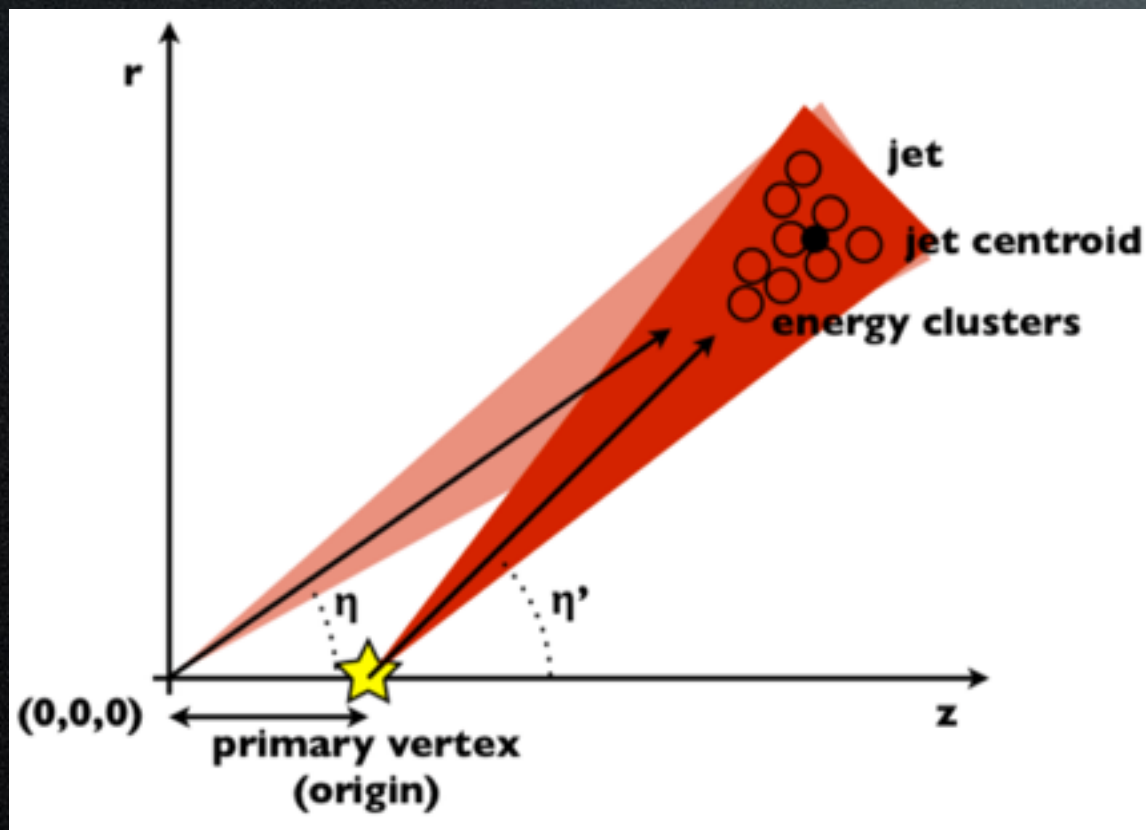
# Definitions

- Calculate  $p_T^{\text{jet}}/p_T^{\text{matched-truth-jet}}$
- Fit a Gaussian
- Mean of the Gaussian:  $R$  or jet response.
- Standard deviation of the Gaussian:  $\sigma/R$  or jet resolution.



# Origin Correction

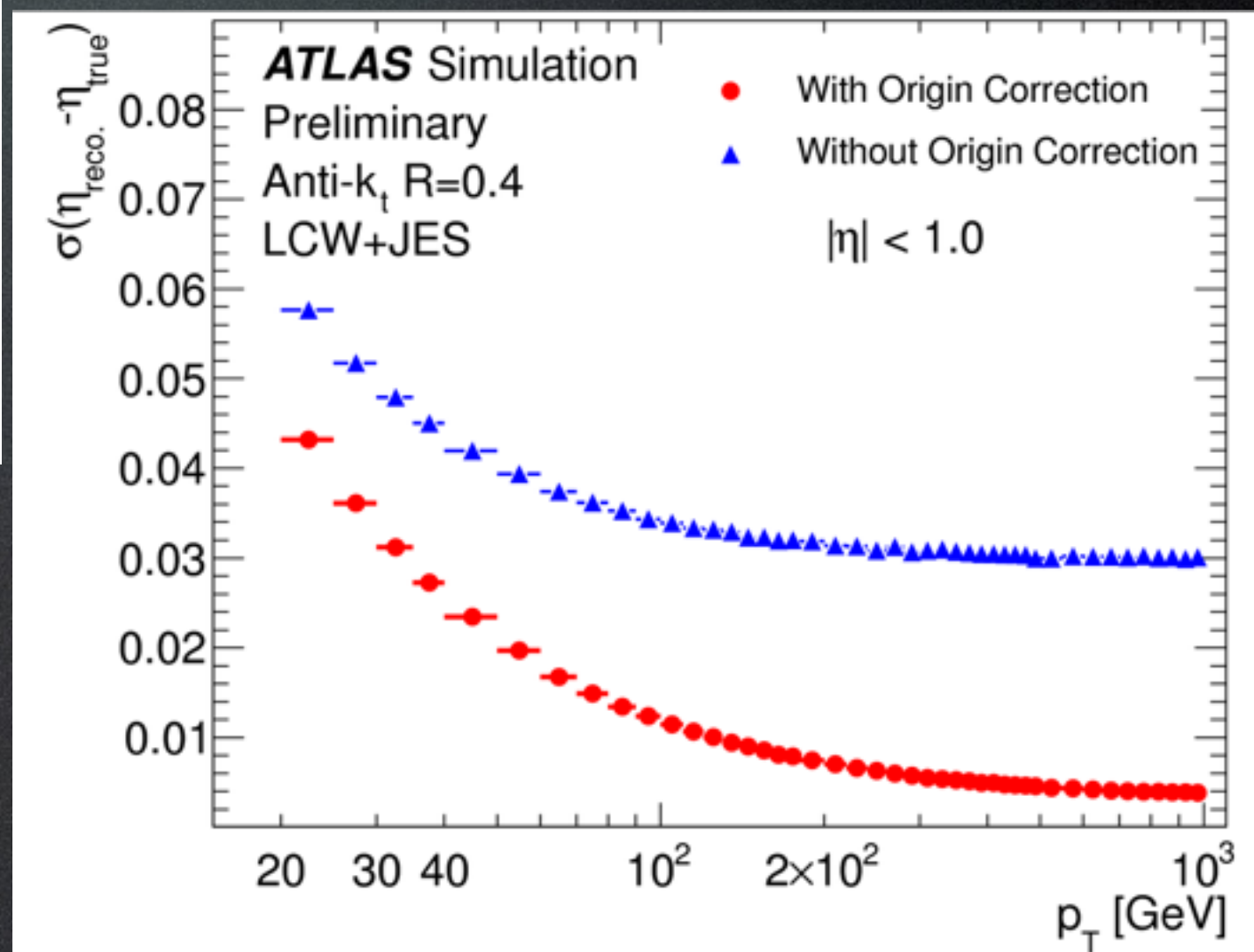
(to point back to hard-scatter PV)



Changes the jet direction  
and improves  $\eta$  resolution

No change in energy

Accounts for the hard  
scattering primary vertex.





# (Jet-area based) Pileup Correction

A uniform population of infinitesimally soft ghost particles are spread ...

Pileup  $p_T$  density:

$$\rho = \text{median} \left\{ \frac{p_{T,i}^{\text{jet}}}{A_i^{\text{jet}}} \right\}$$

First step:

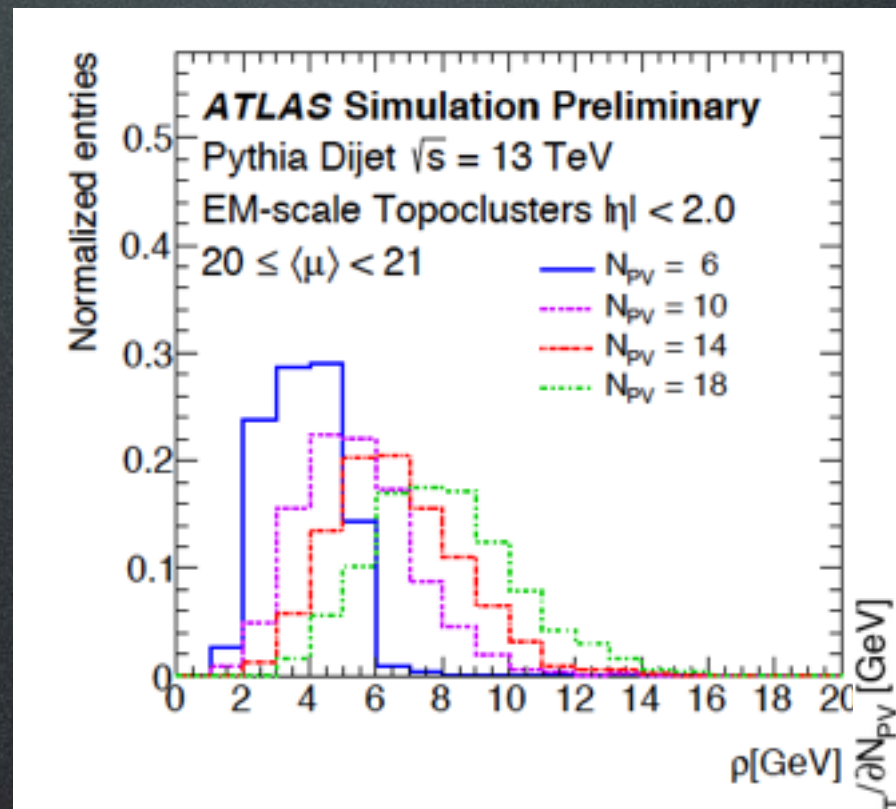
$$p_T^{\text{jet,corr}} = p_T^{\text{jet}} - \rho \cdot A$$

Residual correction:

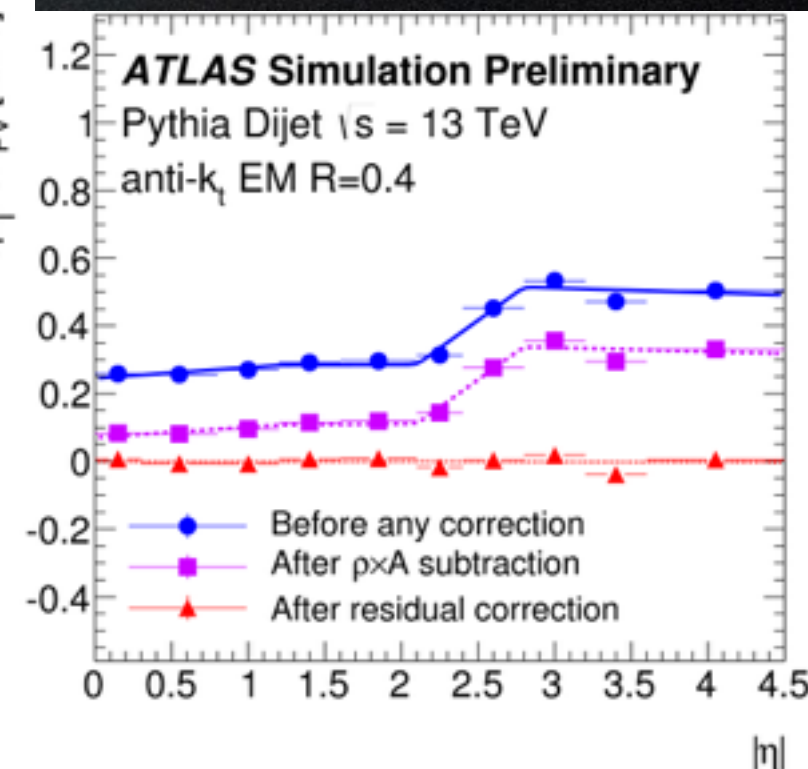
$$\alpha \times (N_{PV} - 1) - \beta \times \langle \mu \rangle$$

In time pu      Out of time pu

Simulation based



Takes into account event-by-event fluctuations





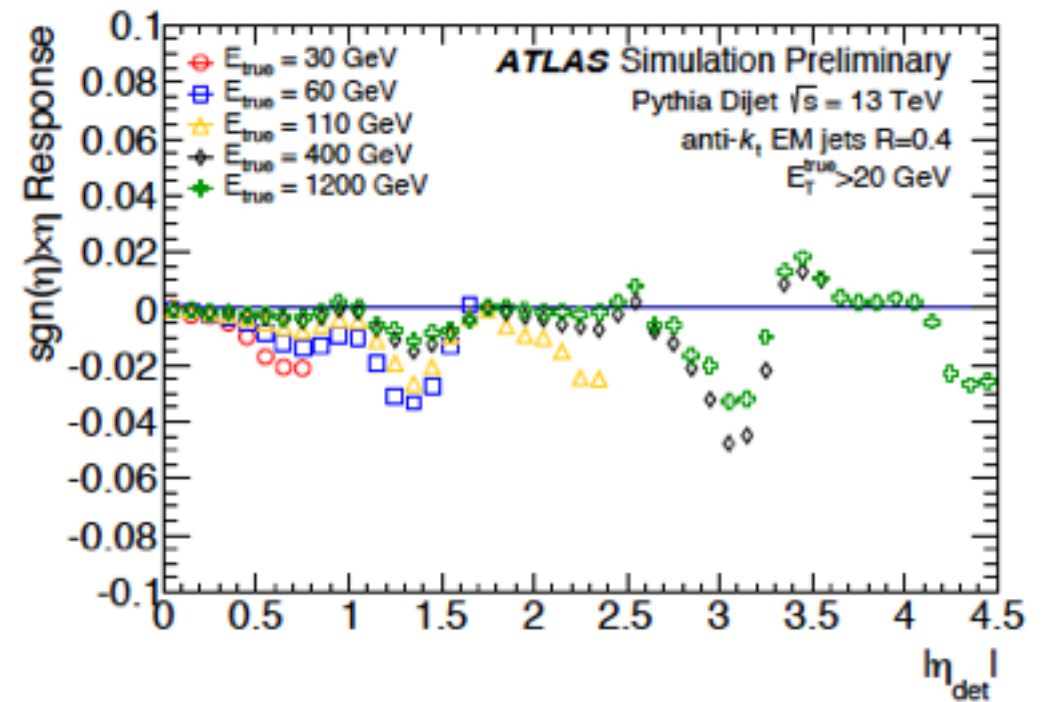
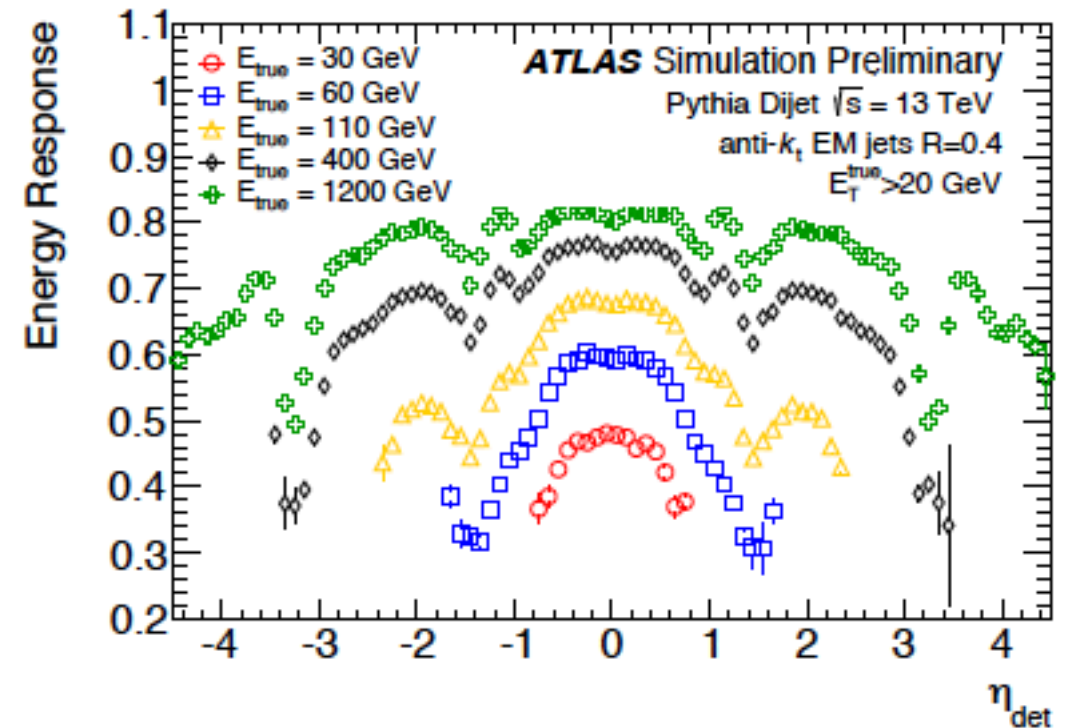
# (MC-based) Jet Energy Scale

Jet response:  
 $p_T$  and  $\eta$  dependant

$$\mathcal{R} = \langle p_T^{\text{jet}} / p_T^{\text{truth}} \rangle$$

In particular regions of the detector there is a bias in the  $\eta$  distribution (different calorimeter type) with respect to the truth jets.

Therefore, an additional correction in purely the angle of the jet is applied to resolve this bias.



Using dijet MC



# Global Sequential Calibration

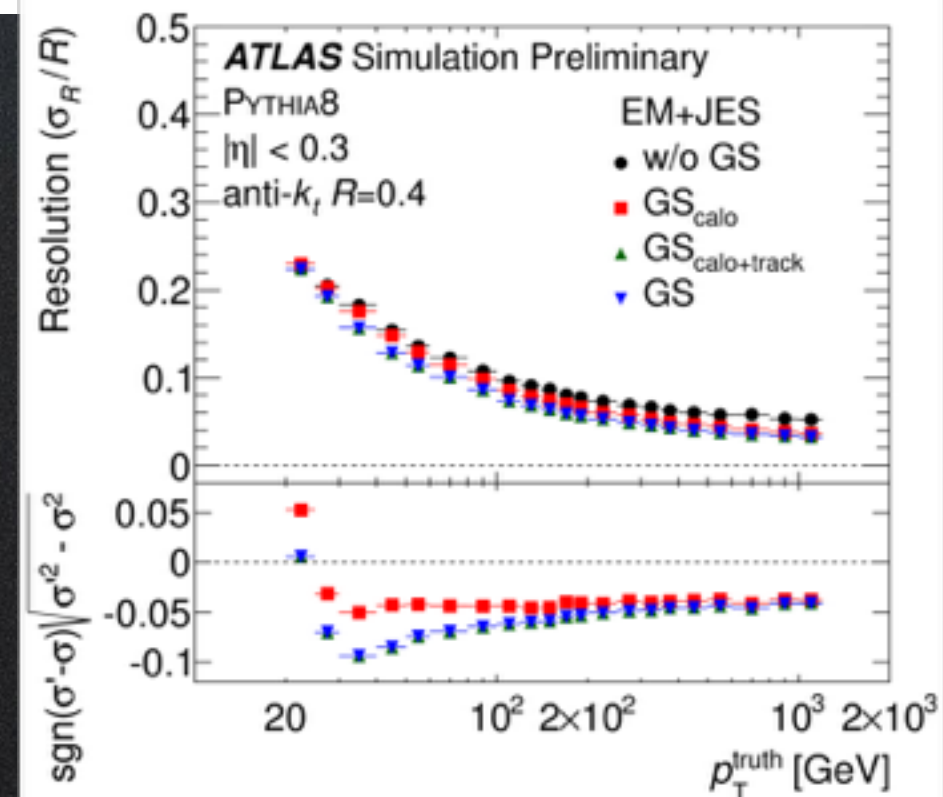
Observed difference between q/g initiated jets (closure, response), leads to large uncertainty on JES.

1. The fraction of energy deposited in the first layer of the tile calorimeter.
2. The fraction of energy deposited in the third layer of the electromagnetic calorimeter.
3. The number of tracks with  $p_T > 1$  GeV associated to the jet.
4. The  $p_T$ -weighted transverse width of the jet measured using tracks with  $p_T > 1$  GeV associated to the jet.
5. The amount of activity behind the jet as measured in the muon spectrometer.

**for LCW last two steps only**

Finally, correct for jets that are not fully contained in the calorimeter.

Mean response not affected

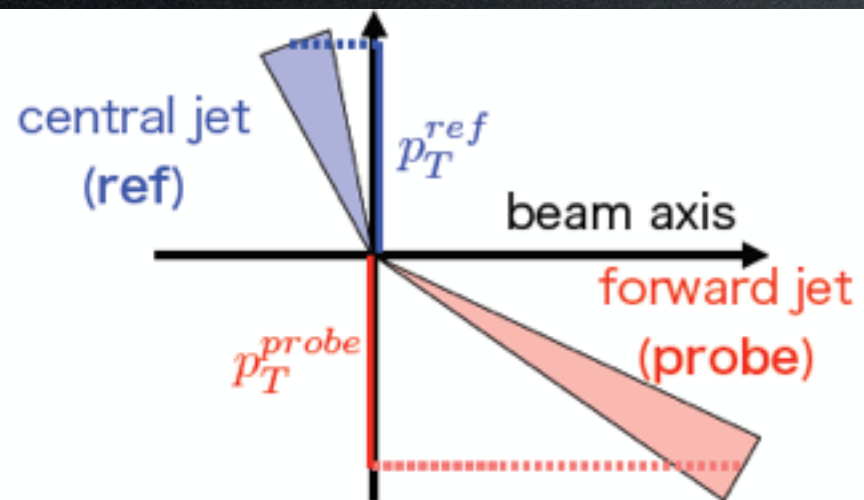




# Residual In situ Corrections

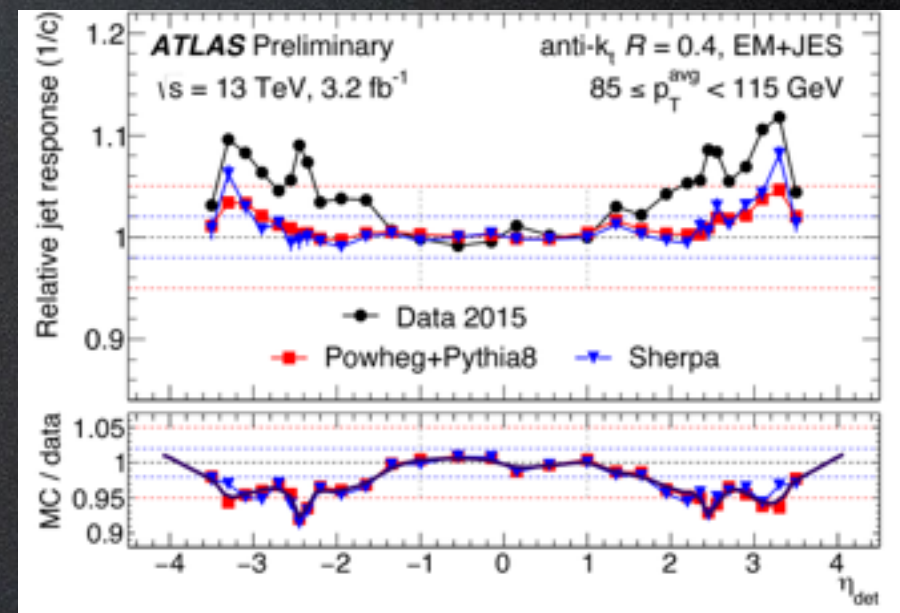
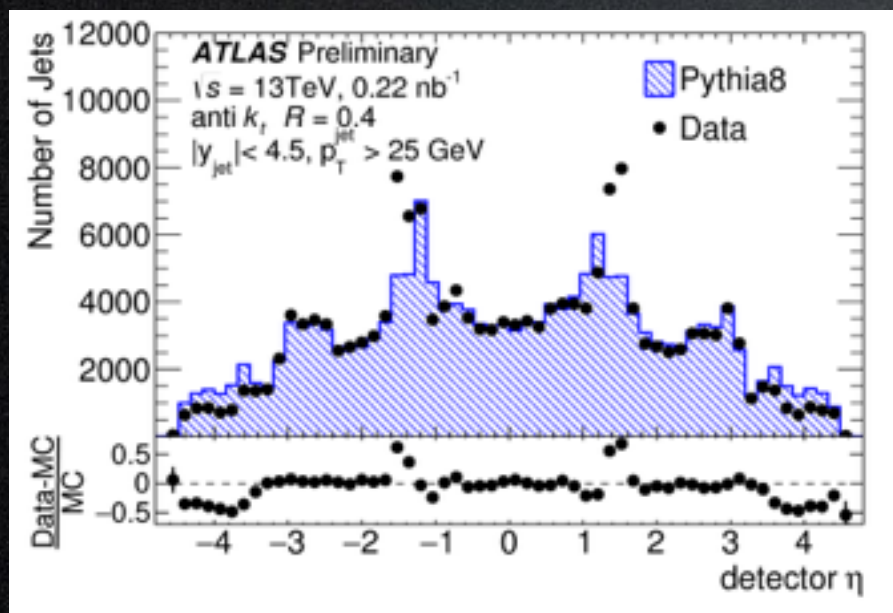
To account for difference in data and simulation  
(applied to data only)

1. Relative (inter)-calibration in eta, derived from dijet events
2. Absolute scale correction from  $\gamma/Z$ +jet and multijet events



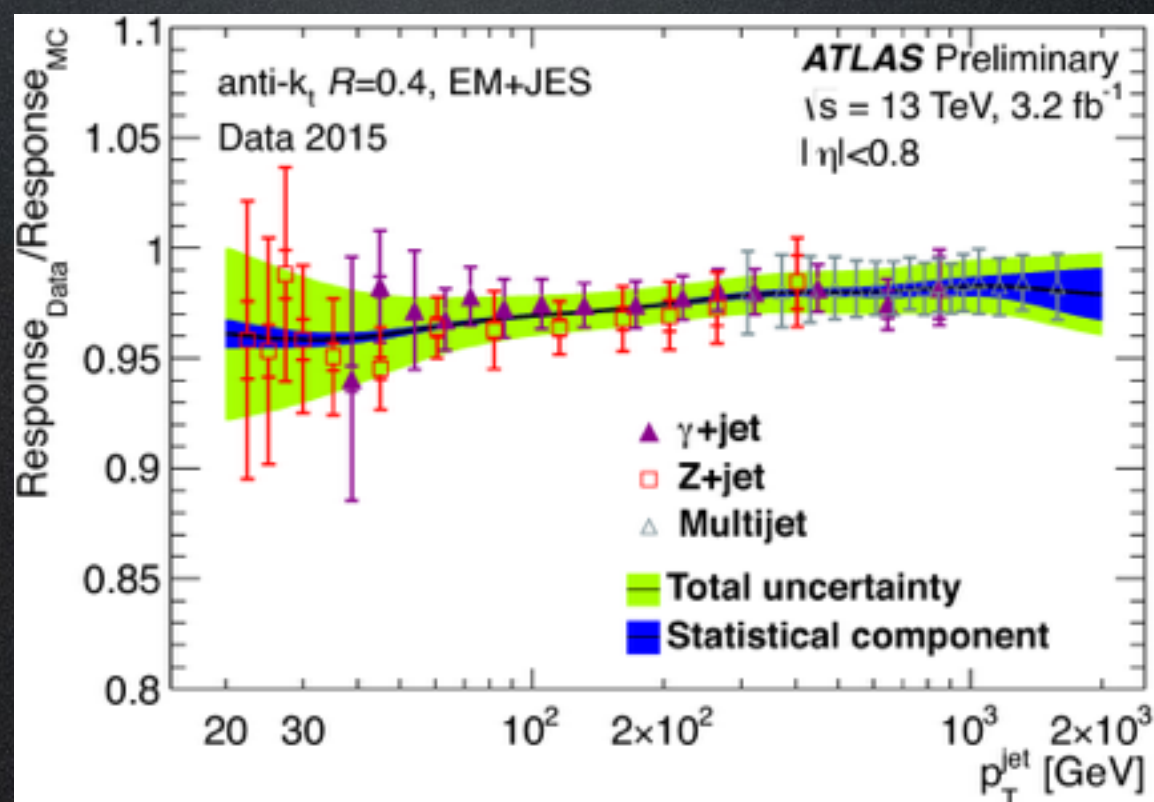
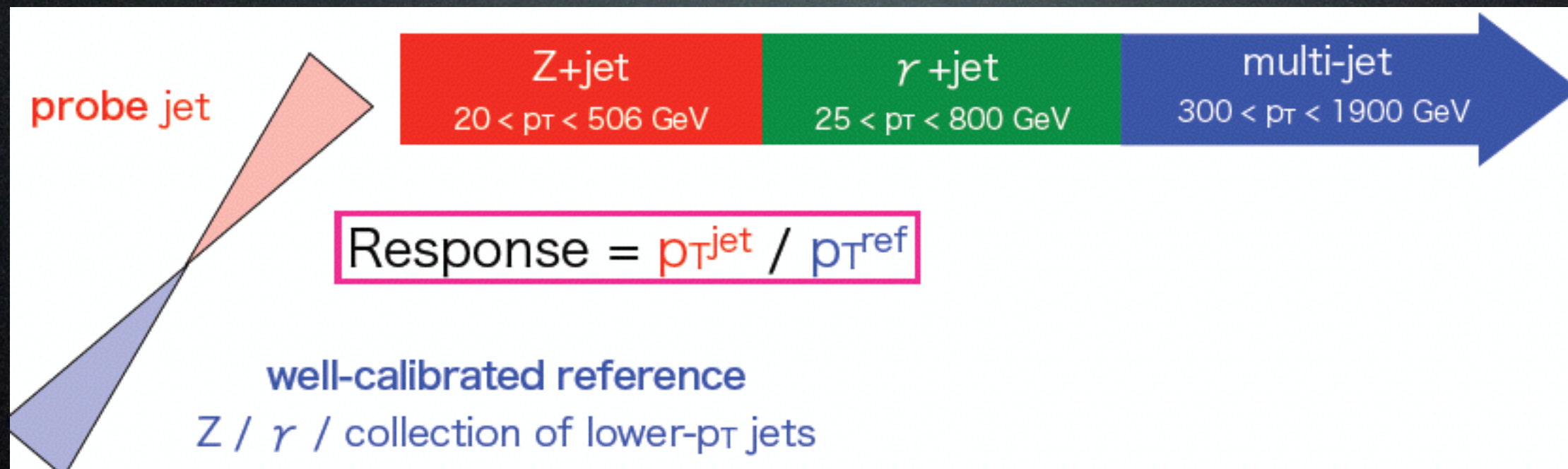
$$\frac{1}{c} = \frac{p_T^{probe}}{p_T^{ref}} = \frac{2 + \langle \mathcal{A} \rangle}{2 - \langle \mathcal{A} \rangle}$$

$$\mathcal{A} = \frac{p_T^{probe} - p_T^{ref}}{p_T^{avg}}, \quad p_T^{avg} = \frac{p_T^{probe} + p_T^{ref}}{2}$$



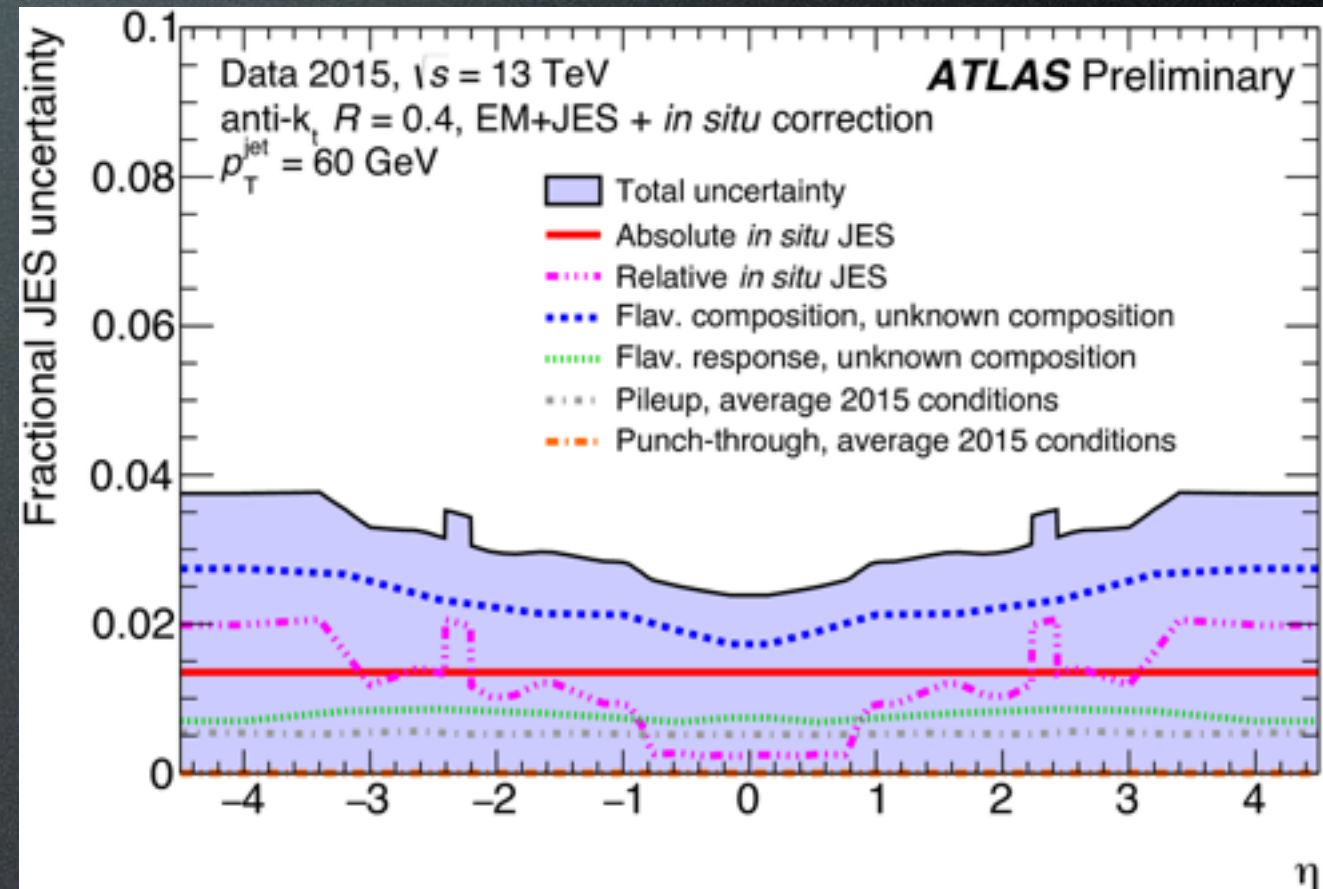
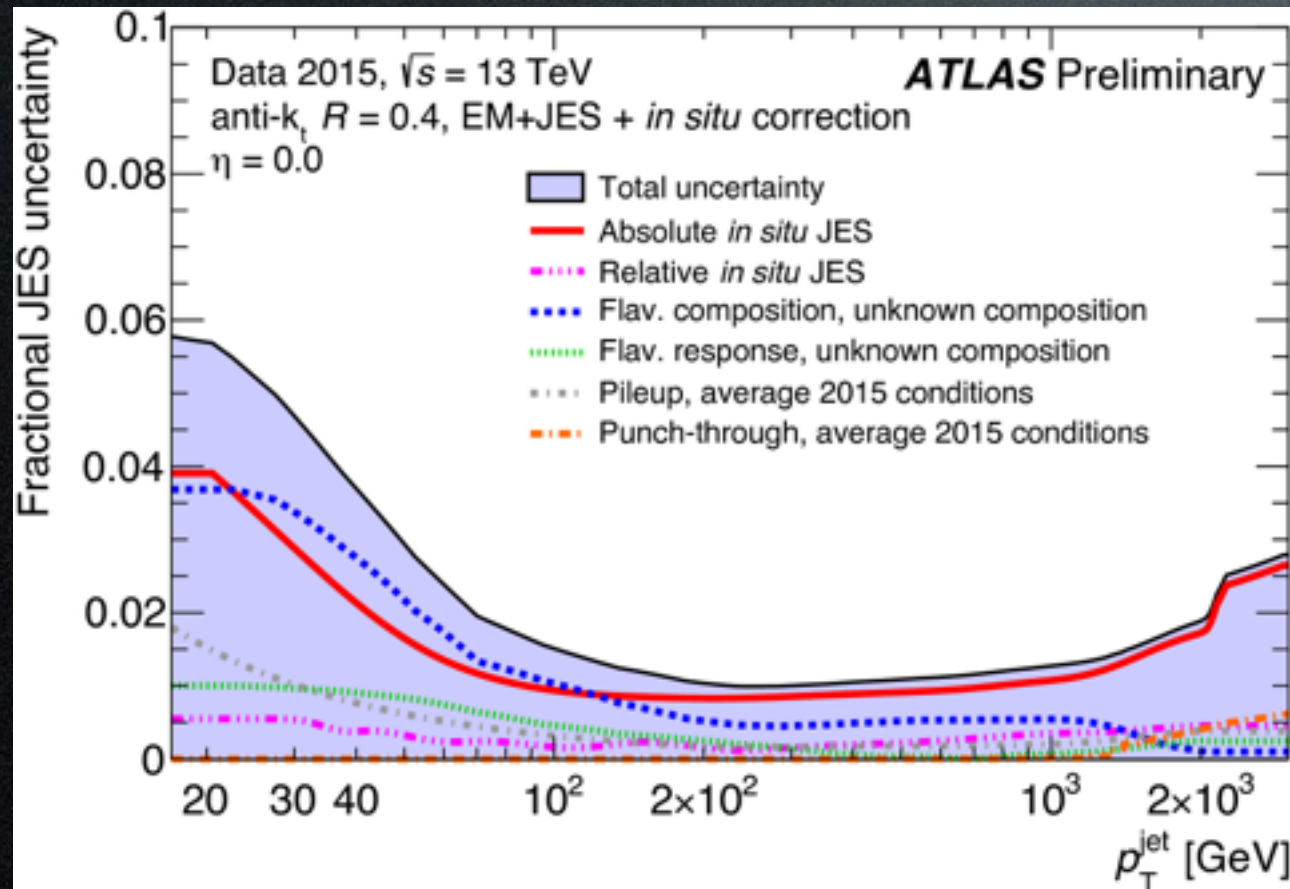


# Jet Energy Calibration





# JES uncertainty



70 correlated components

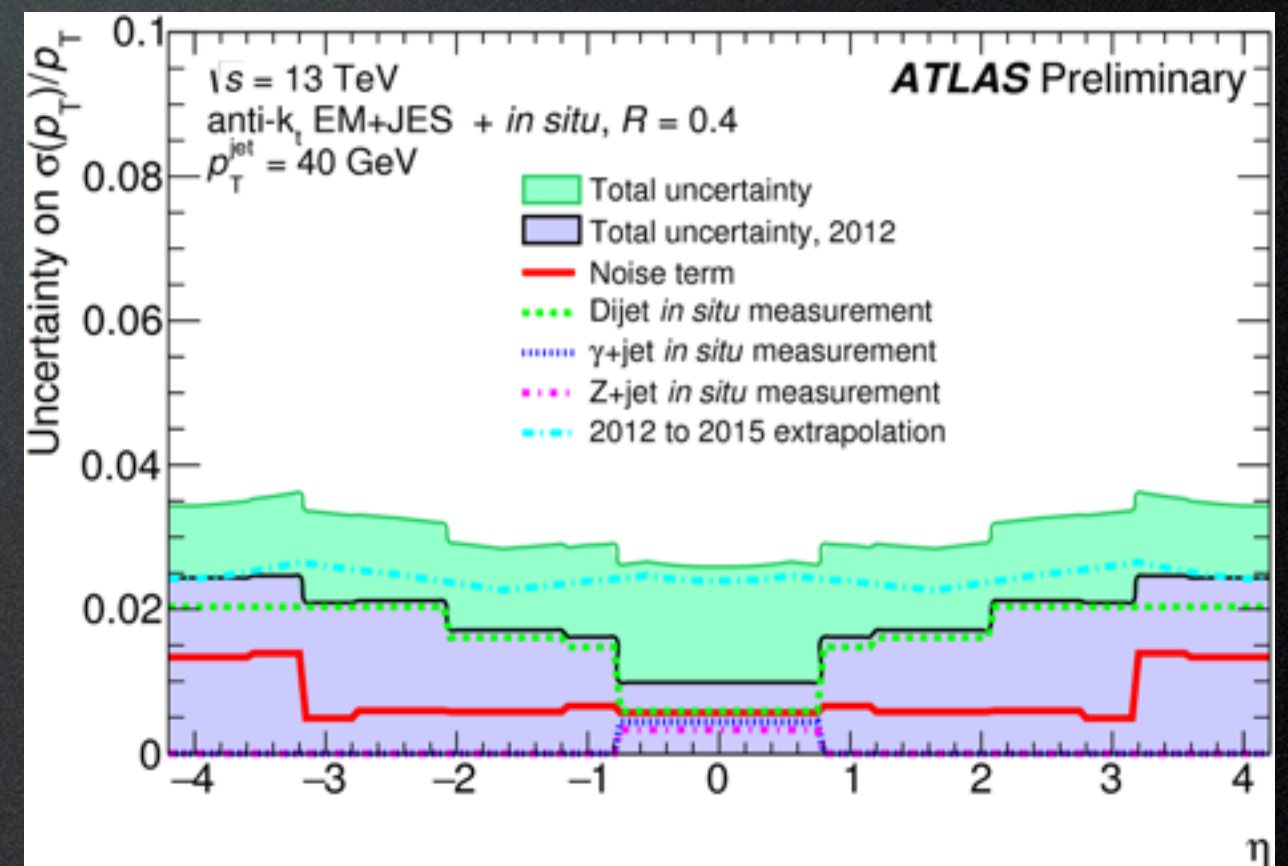
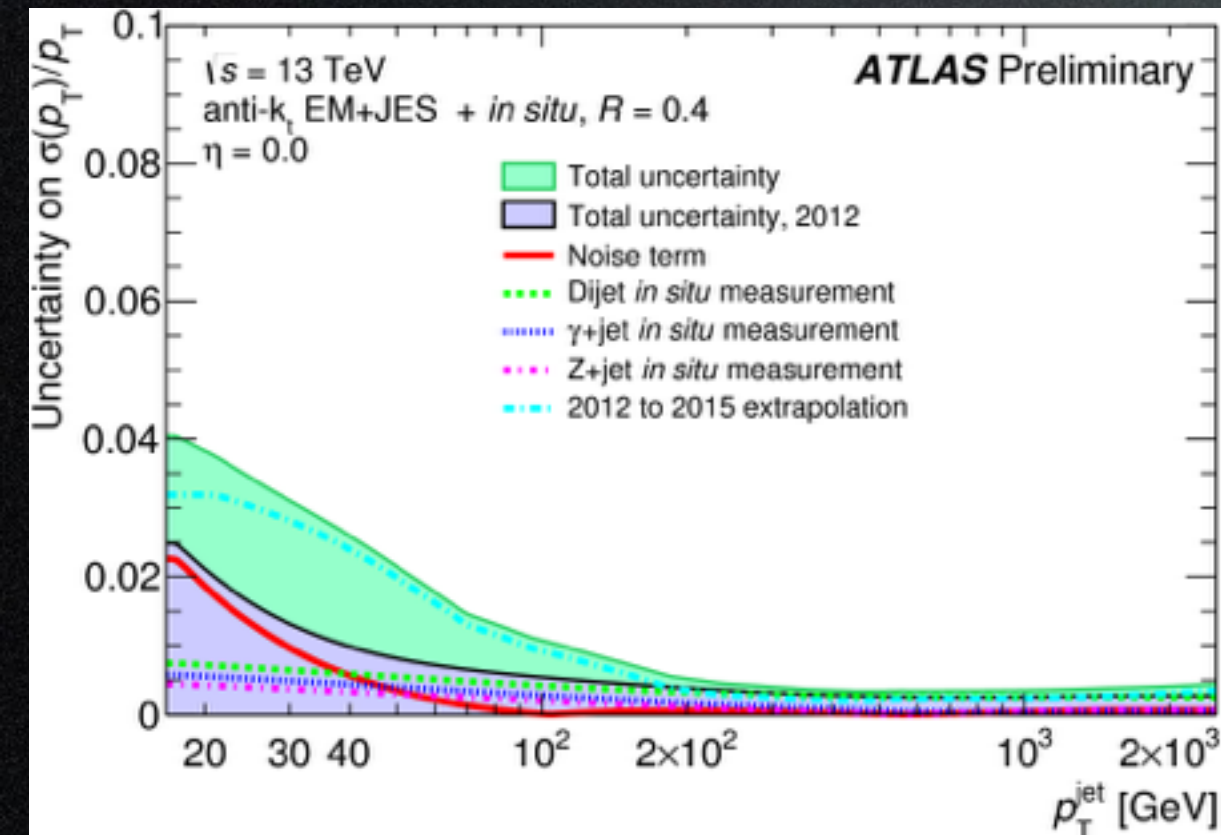
Reduced to smaller number in analyses



# Jet Energy Resolution

$$\sigma(\mathcal{A}) = \frac{\sqrt{\sigma(p_T^{ref})^2 + \sigma(p_T^{probe})^2}}{p_T^{avg}}$$

All uncertainties are combined into one nuisance parameter in a correlated way





# Large-R Jet

- No explicit pileup removal or origin correction.
- After MC based JES, data-MC double ratio for residual correction:

$$\frac{\mathcal{R}_{\text{data}}}{\mathcal{R}_{\text{MC}}} = \frac{\langle p_{\text{T}}^{\text{jet}} / p_{\text{T}}^{\text{ref}} \rangle_{\text{data}}}{\langle p_{\text{T}}^{\text{jet}} / p_{\text{T}}^{\text{ref}} \rangle_{\text{MC}}}$$

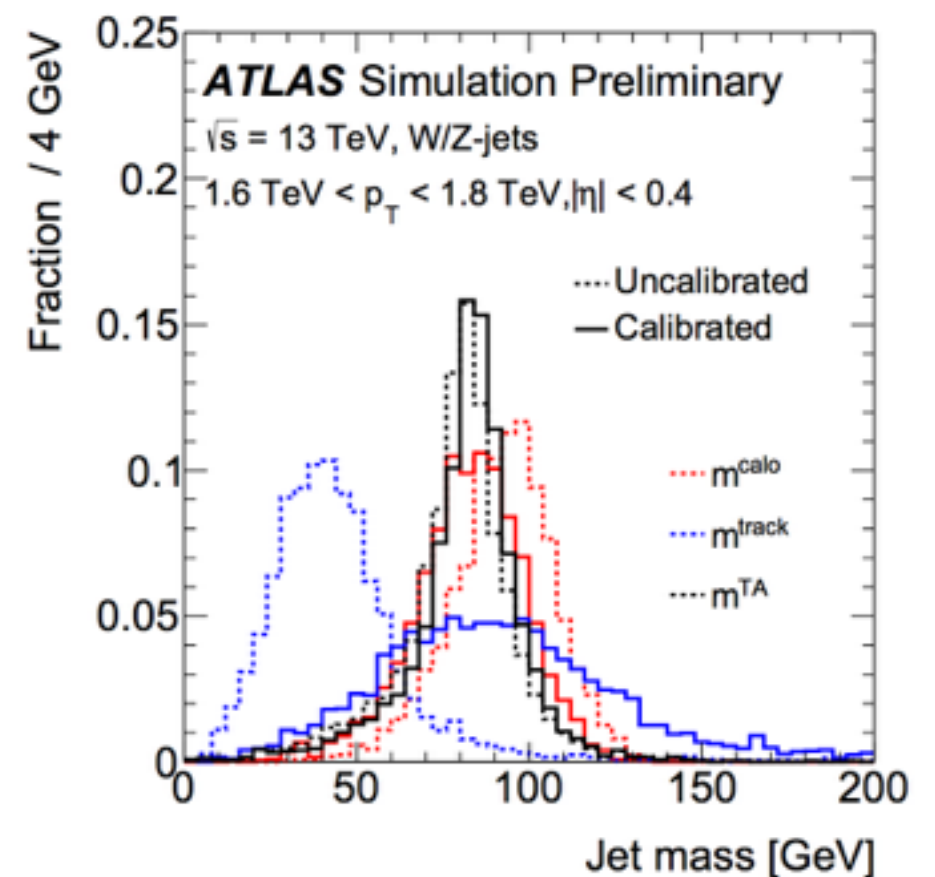
- More important: mass calibration.



# (Large-R) Jet Mass

- $m^{\text{calo}}$ : from constituent four vectors
- $m^{\text{TA}}$ : track-assisted-mass, ratio corrects for charge to neutral fluctuations. Better control of uncertainties.

$$m^{\text{TA}} = \frac{p_T^{\text{calo}}}{p_T^{\text{track}}} \times m^{\text{track}}$$



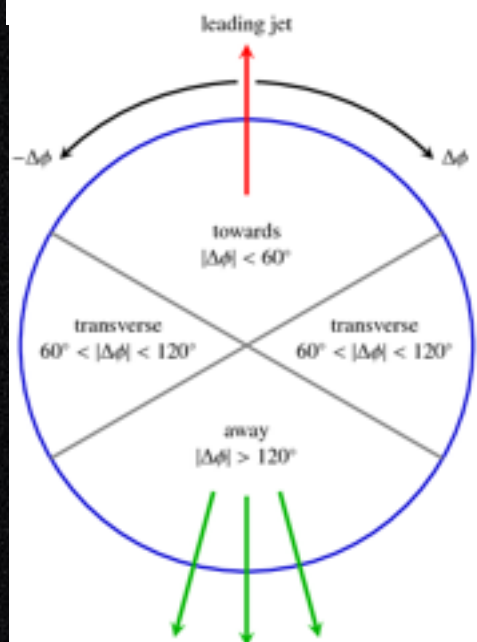
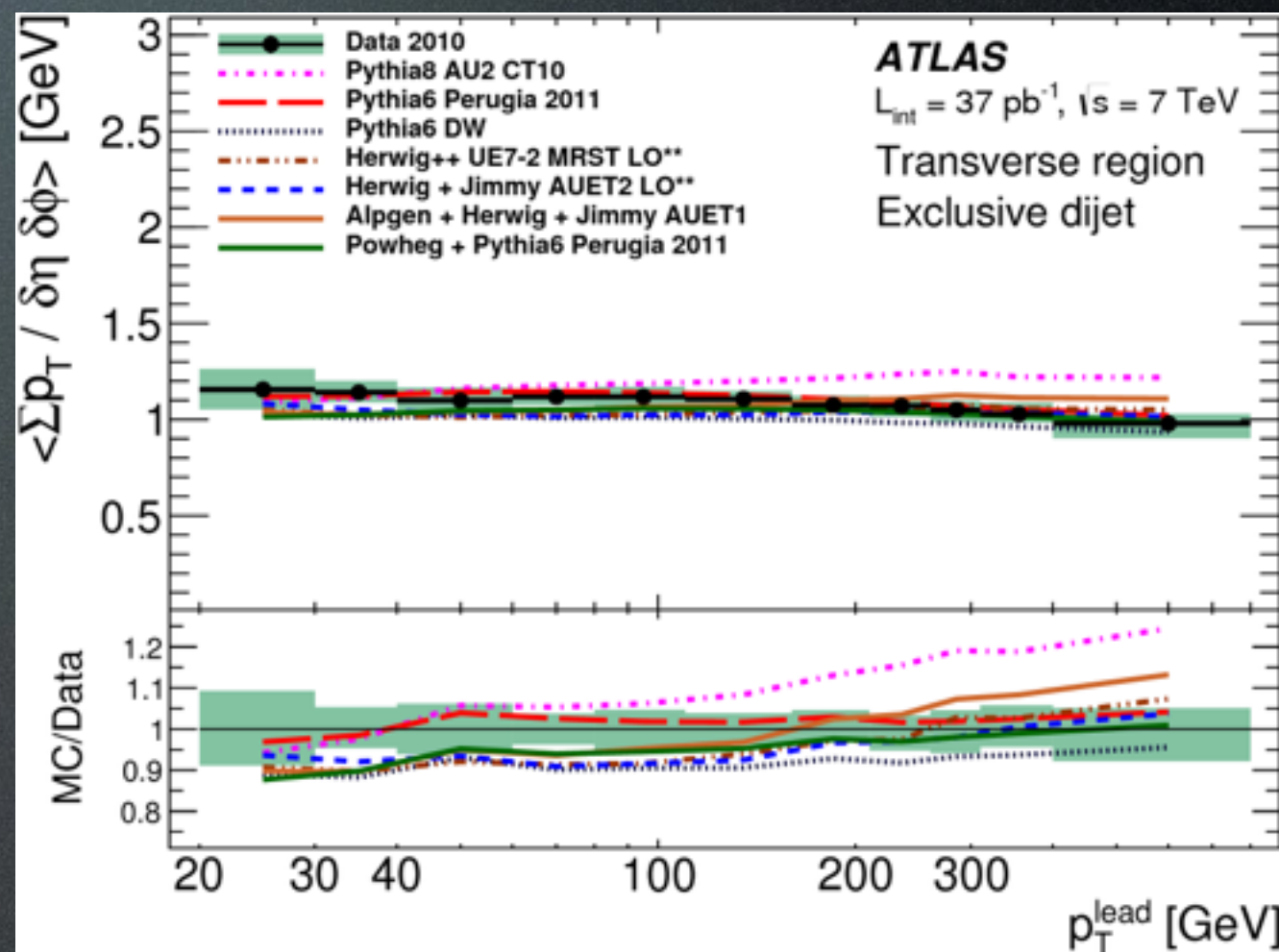
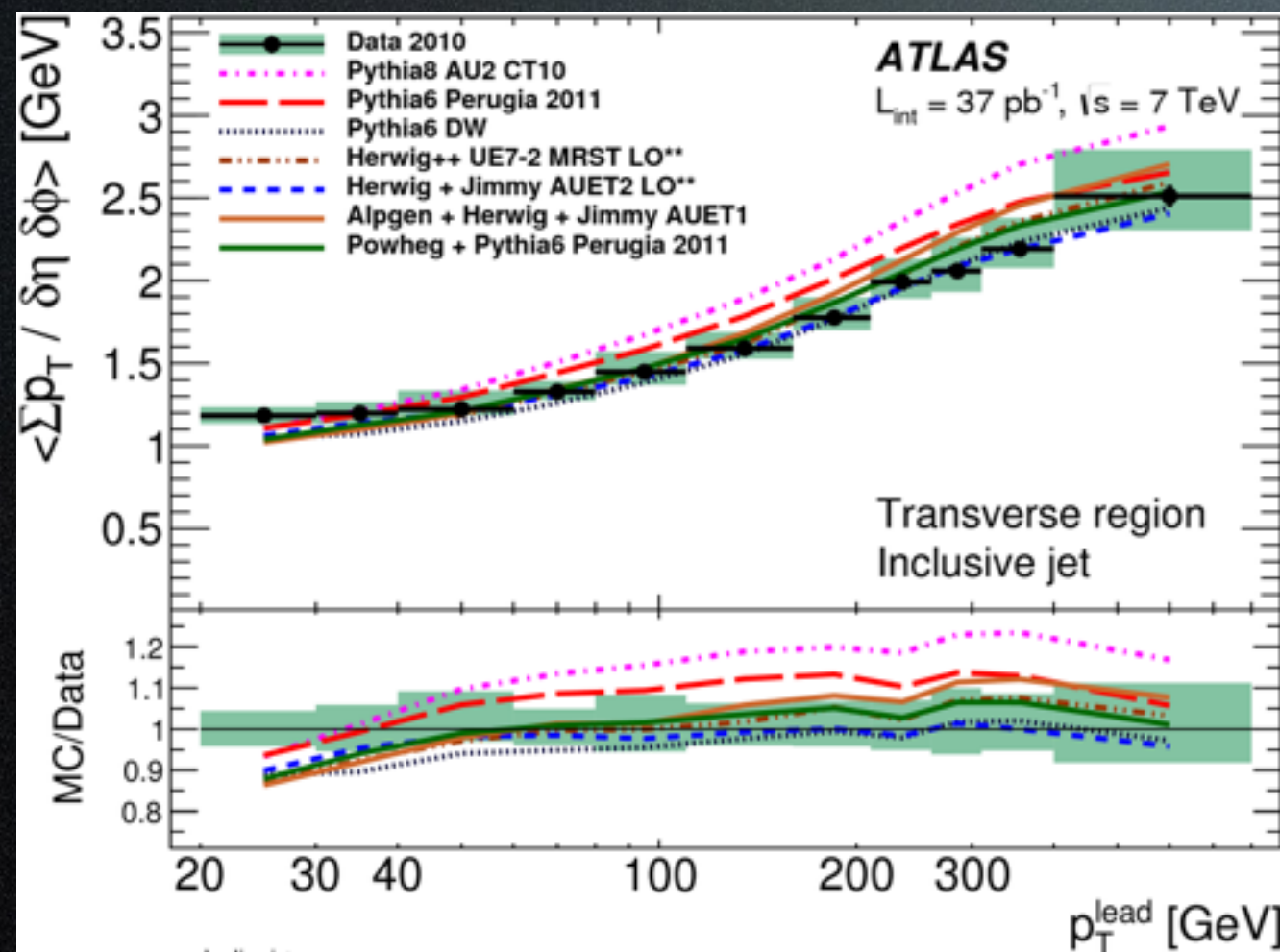


# Jet Measurements

- Many interesting physics results
- Only a very small number covered in this talk



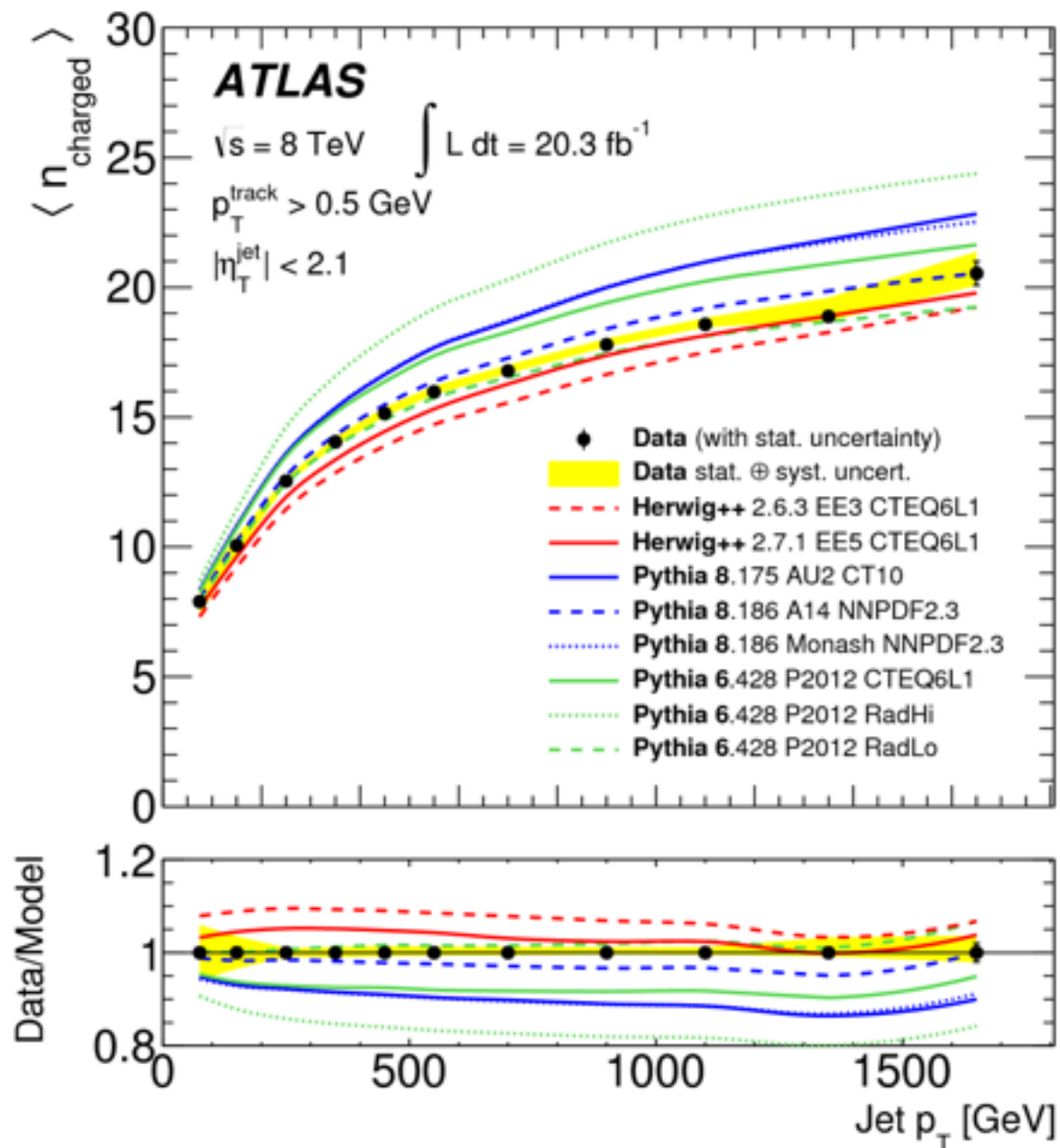
# UE in Jet Events



Exclusive topology more sensitive to MPI, modelled better



# Charged Multiplicity inside Jets

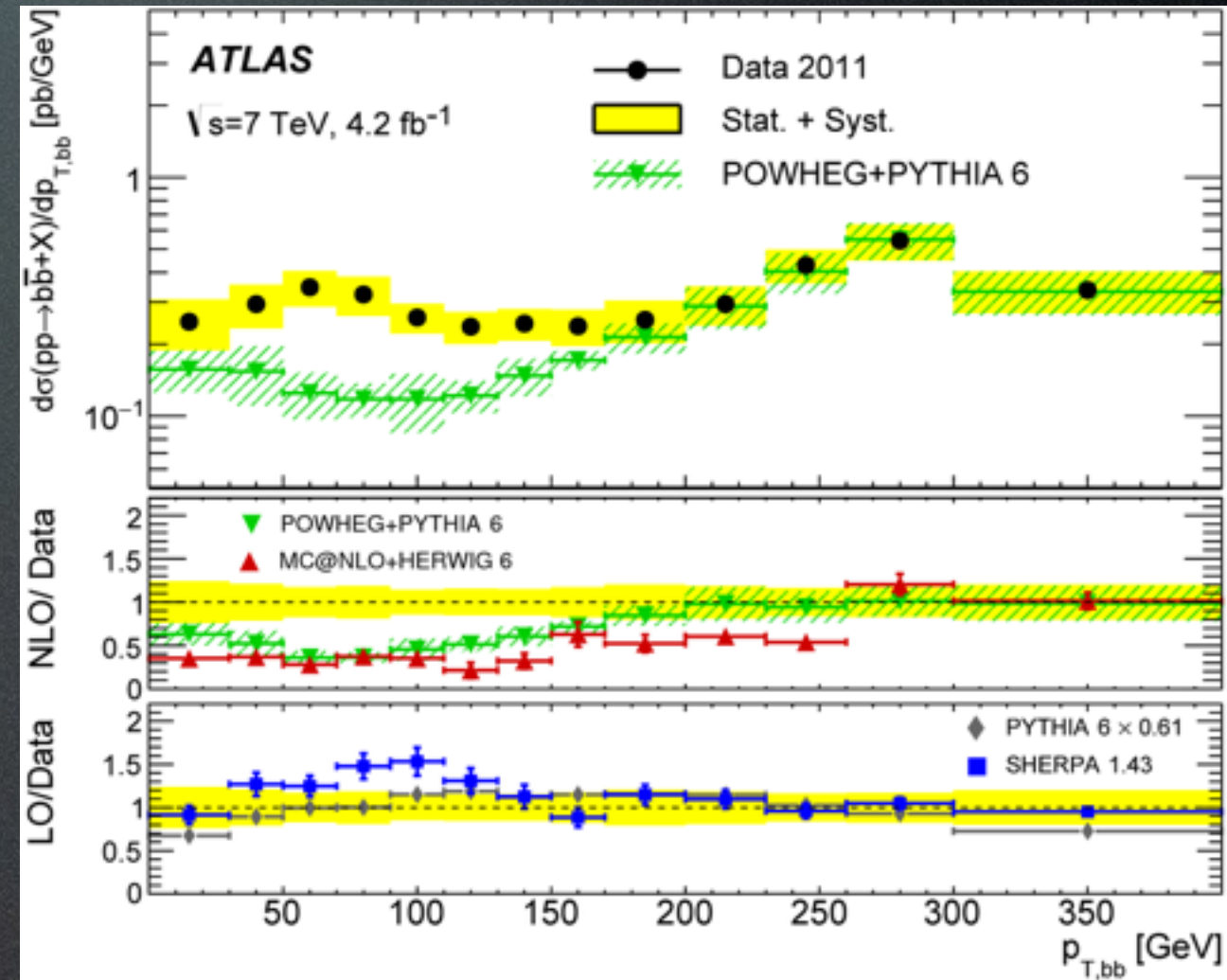
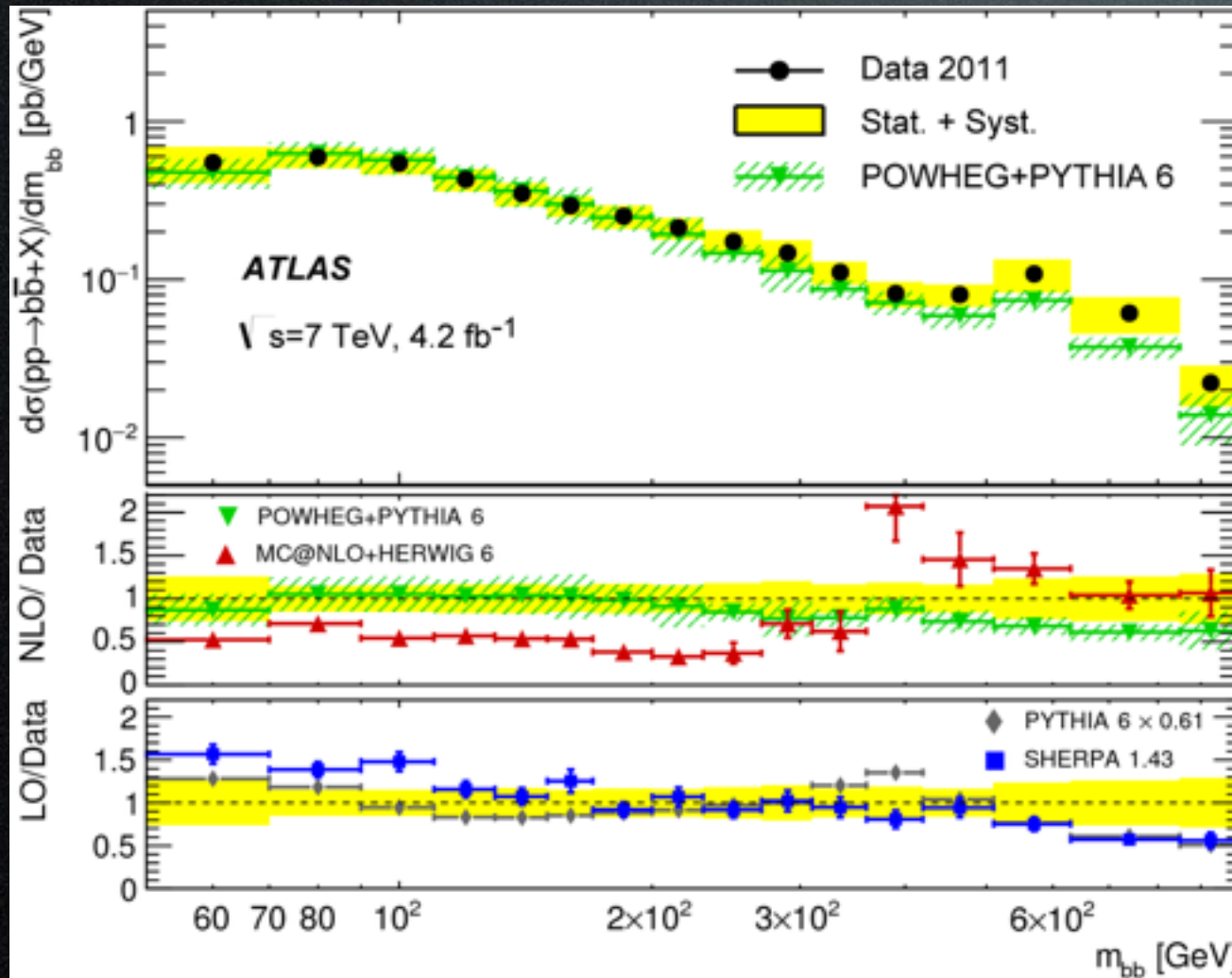


Sensitive to q/g tagging

Significant differences  
between  
generator predictions.



# bbbar-dijet Production

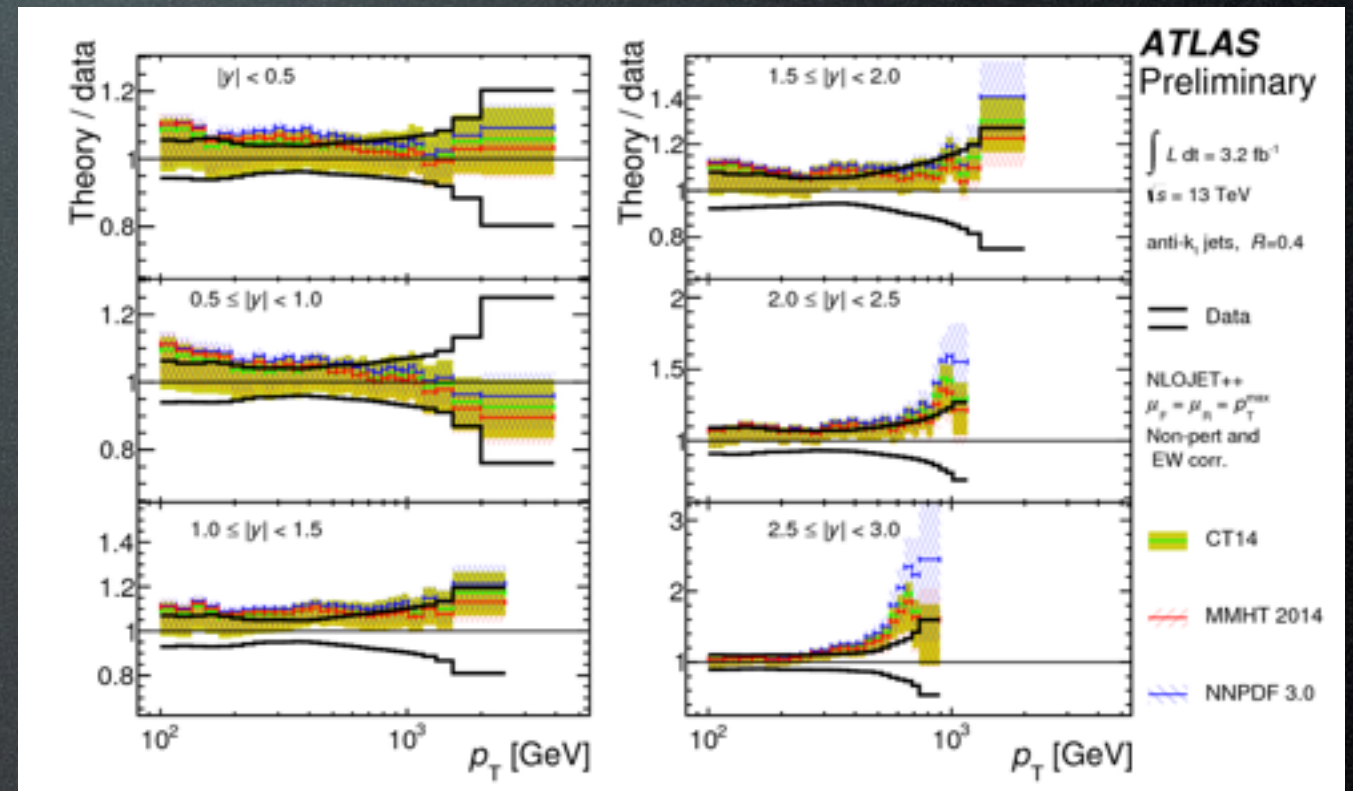
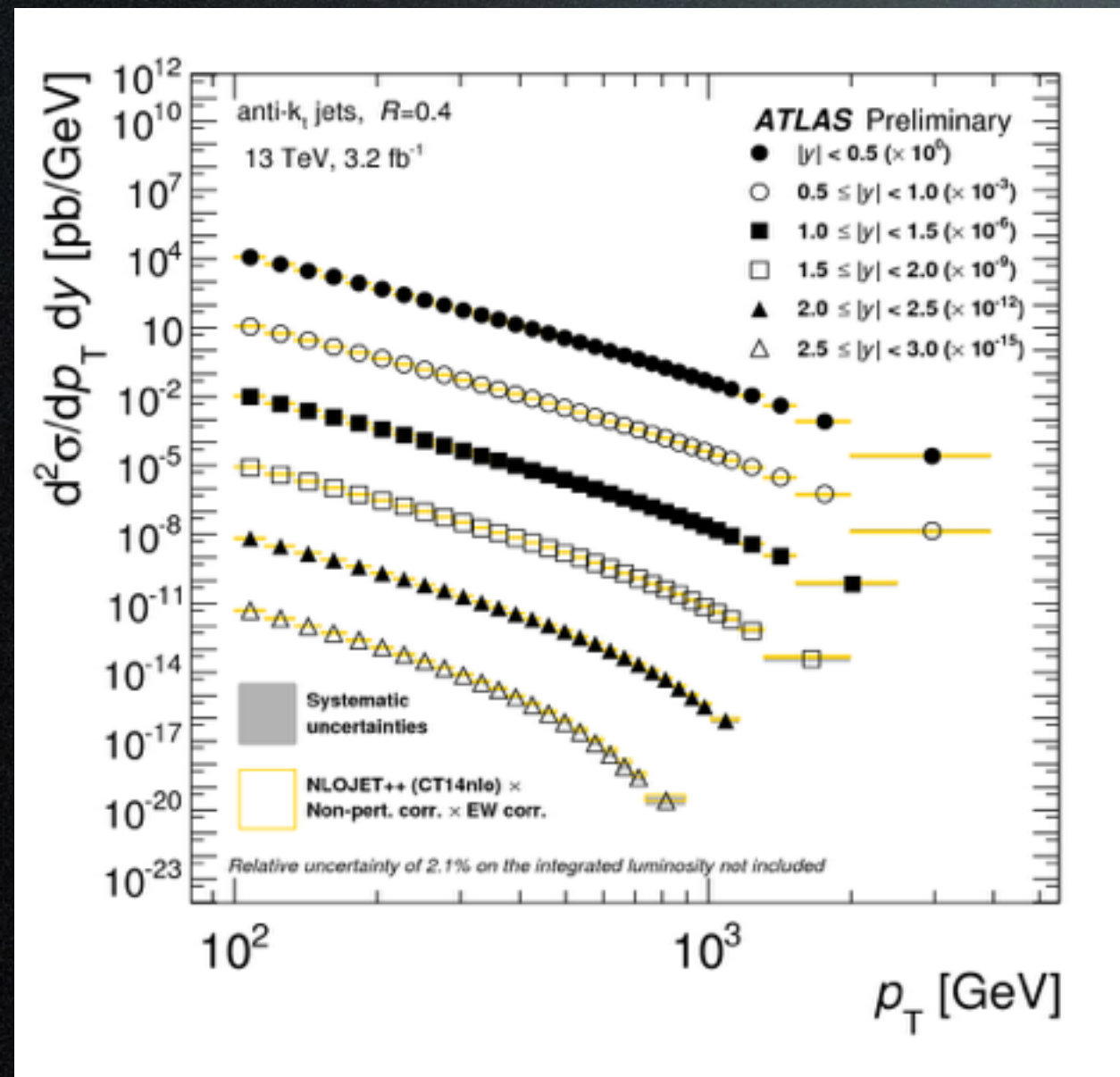


b-tagged by vertex and track information in a neural network, 70% efficiency

Difference in data and predictions



# Inclusive Jet Cross-section at 13 TeV

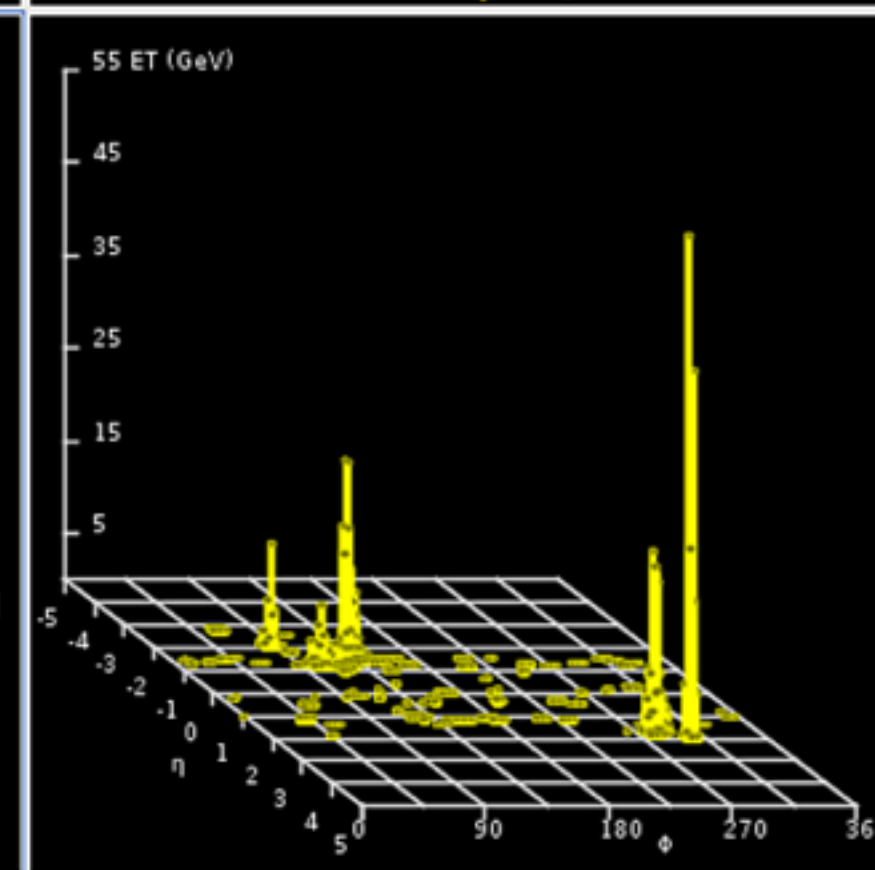
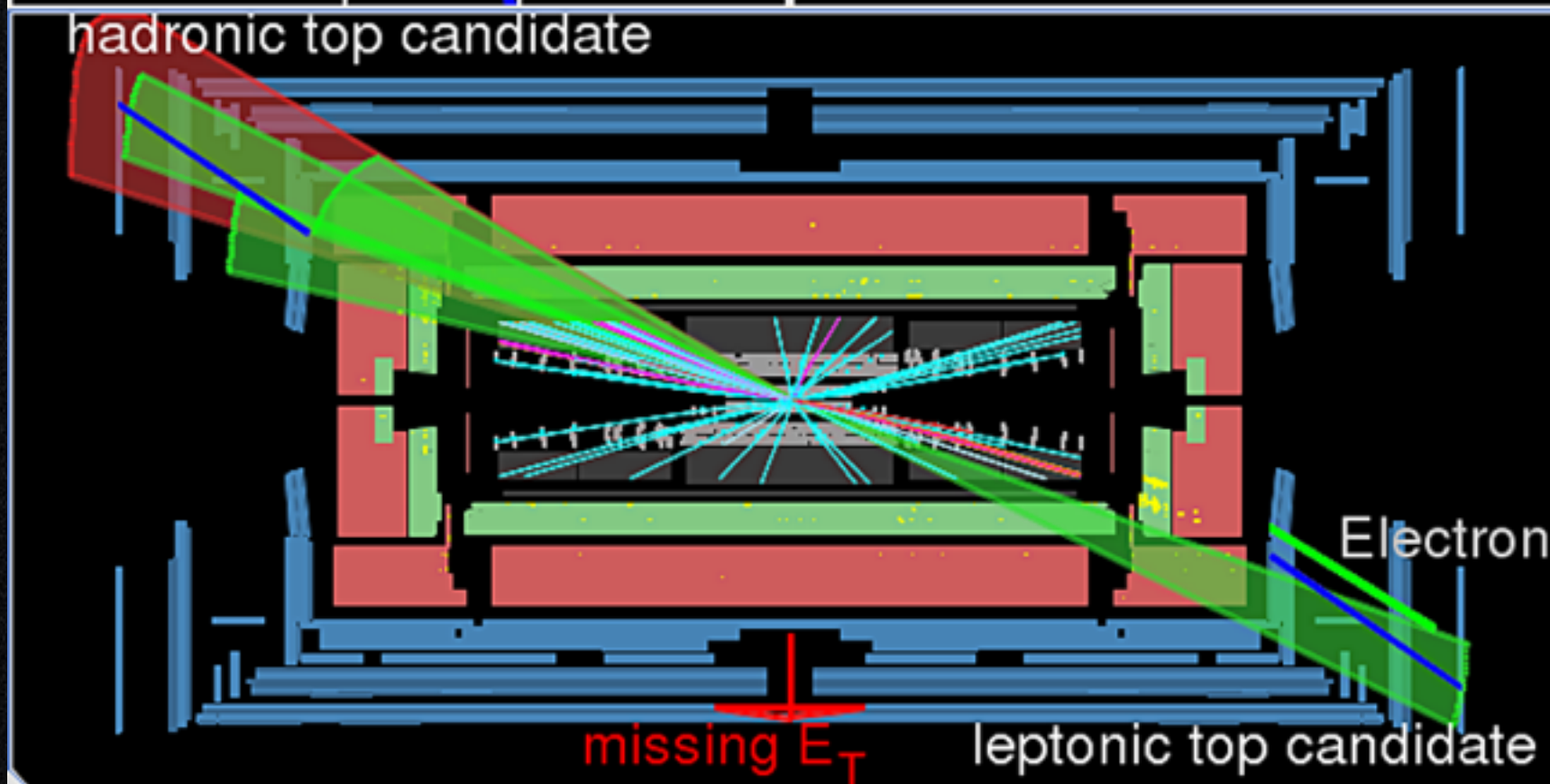
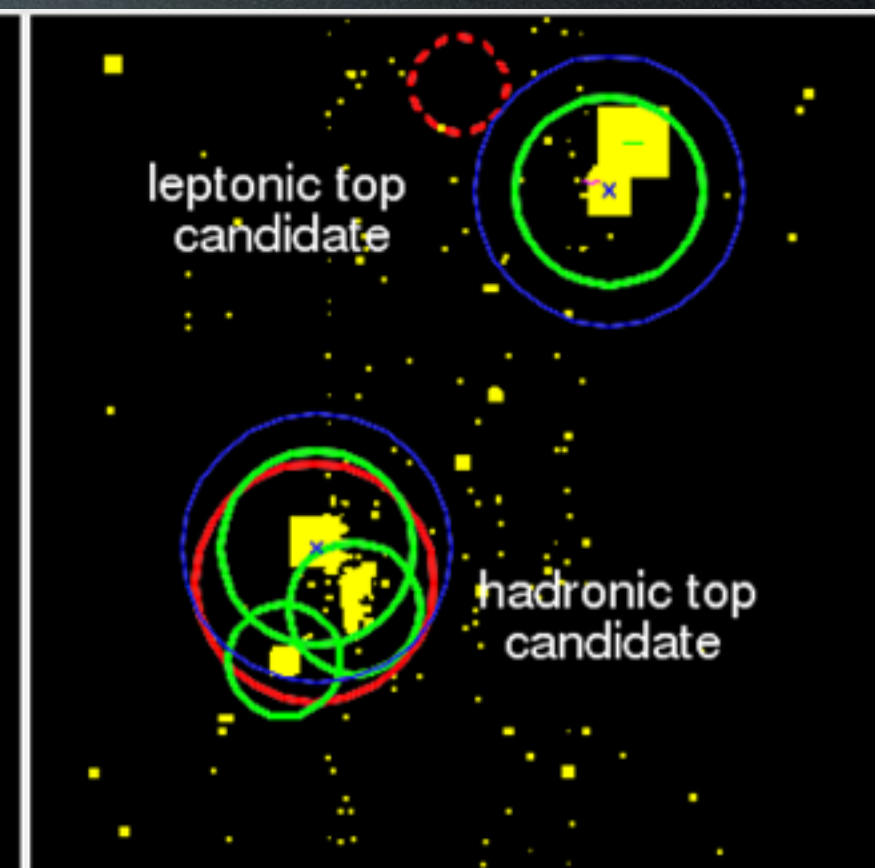
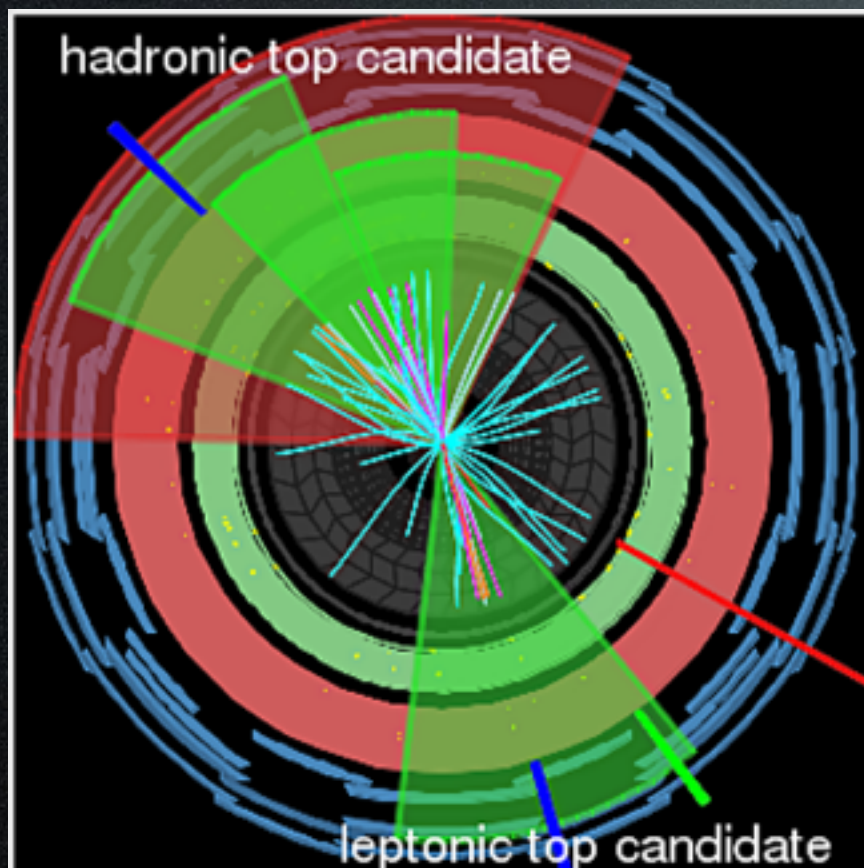


100 GeV to 3.2 TeV!



# Part 2: Jet Substructure, techniques and uses









hadronic top candidate

This visualization shows a hadronic top candidate as a series of concentric, semi-transparent colored rings (red, green, blue) in a detector-like coordinate system. A blue line segment is drawn across the rings.



ATLAS

The ATLAS logo is displayed in a stylized, metallic font.



leptonic top candidate

This visualization shows a leptonic top candidate with a central yellow square, a red dashed circle, and a blue solid circle. A blue 'x' marks a point within the circles. Yellow dots are scattered in the background.

Although reduced cross-section, better acceptance, smaller background

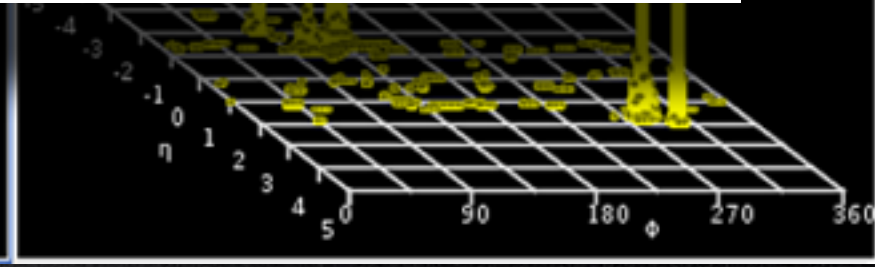
Boosted top ( $qq'b$ ), Higgs ( $bb$ ).  $W/Z$  ( $qq'$ ) are key signatures for many search and measurements



missing  $E_T$

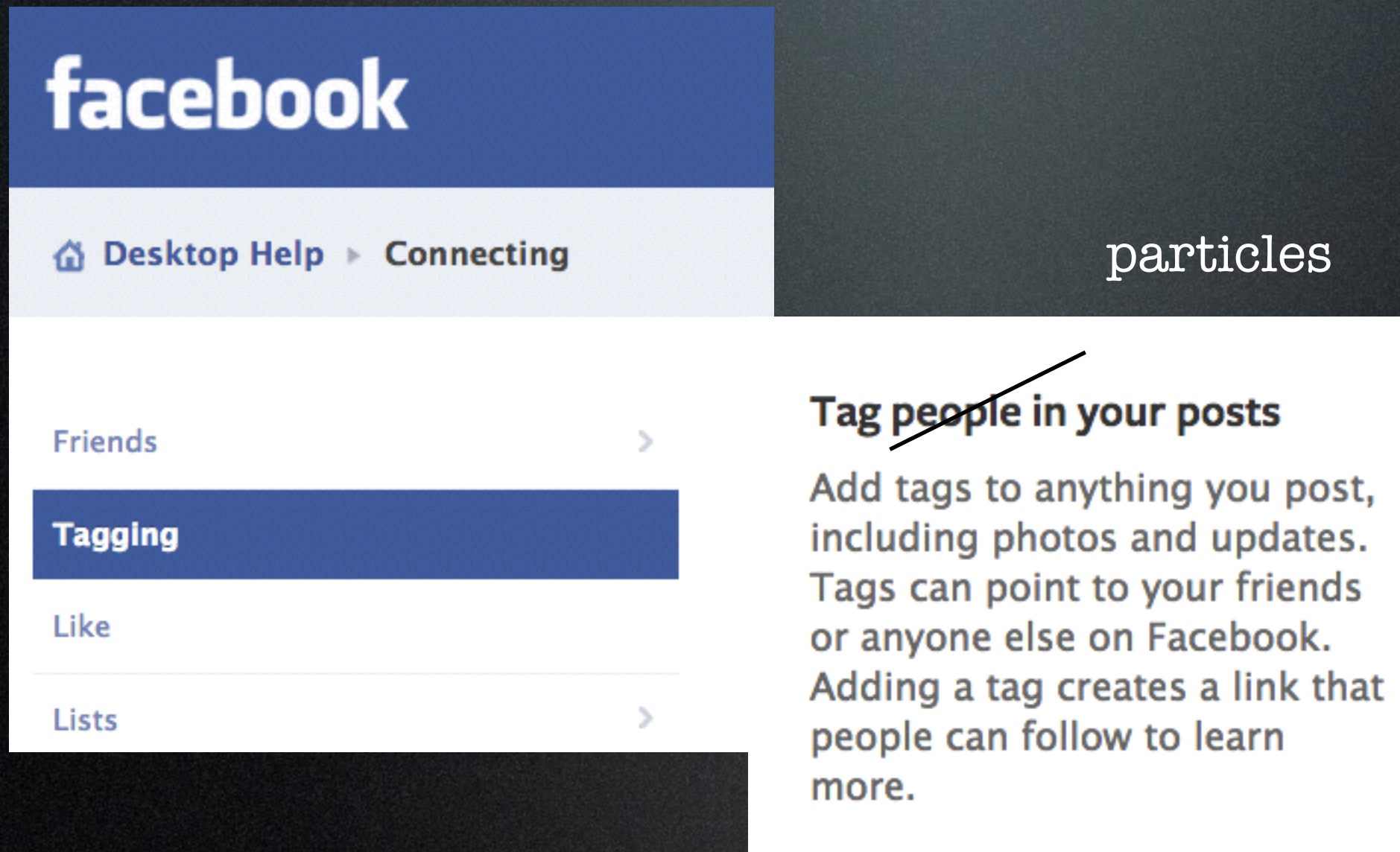
This visualization shows a series of horizontal blue bars of varying heights. A red vertical line with a downward-pointing arrow is positioned over one of the bars.

leptonic top candidate





# Tagging Boosted Objects: observables and taggers



particles

Target is to identify jets resulting from the decay of top quark or Higgs against jets coming from light quark/gluons.



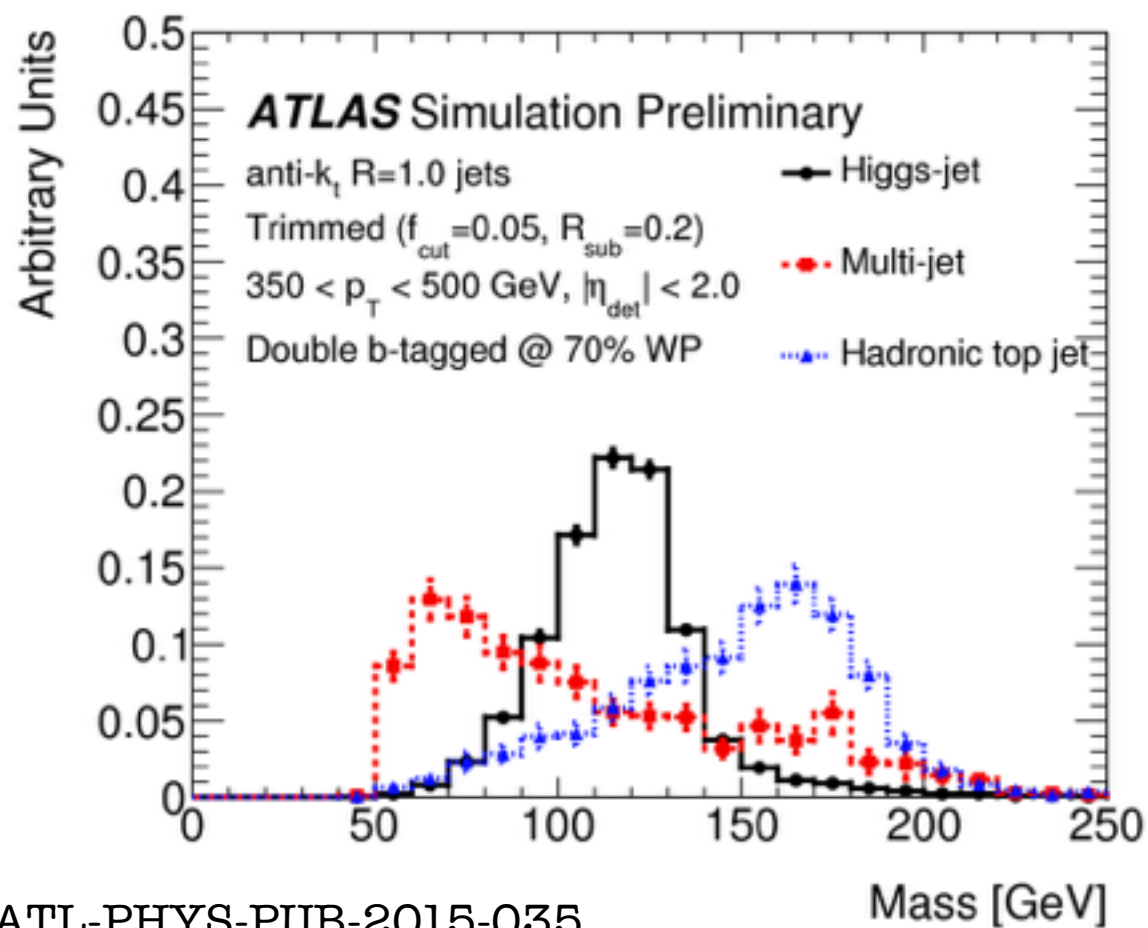
# What we use in ATLAS

Typical starting point: Trimmed ( $f_{\text{cut}} = 5\%$ ,  $R_{\text{sub}} = 0.2$ )  
antikt<sub>t</sub> R=1.0 jet with  $p_T > 200 \text{ GeV}$ ,  $|\eta| < 2.0$

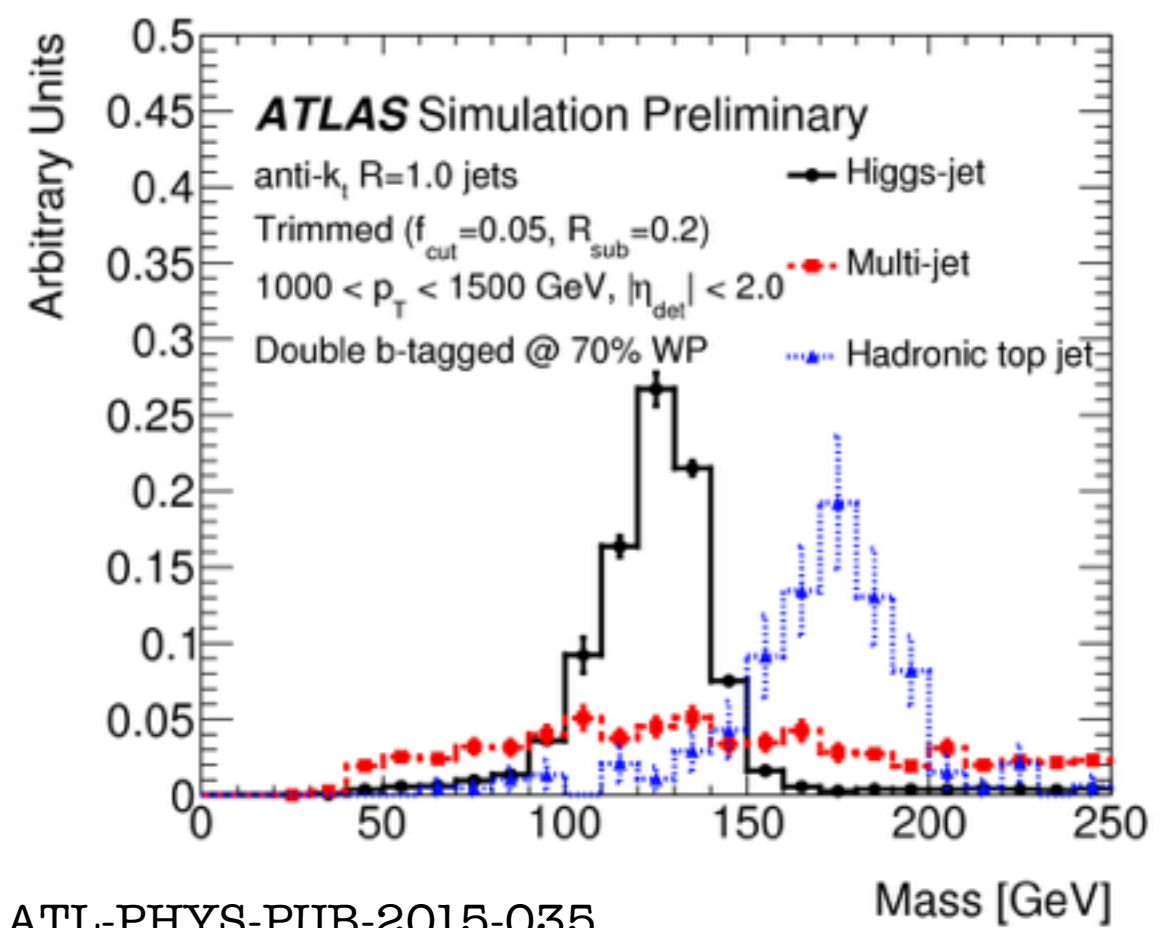
- Jet mass
- Splitting Scale
- Nsubjettiness (discussed yesterday)
- Energy Correlation
- HEPTopTagger (discussed yesterday)
- Shower Deconstruction



# Boosted (tt)Higgs to (tt)bbbar



ATL-PHYS-PUB-2015-035

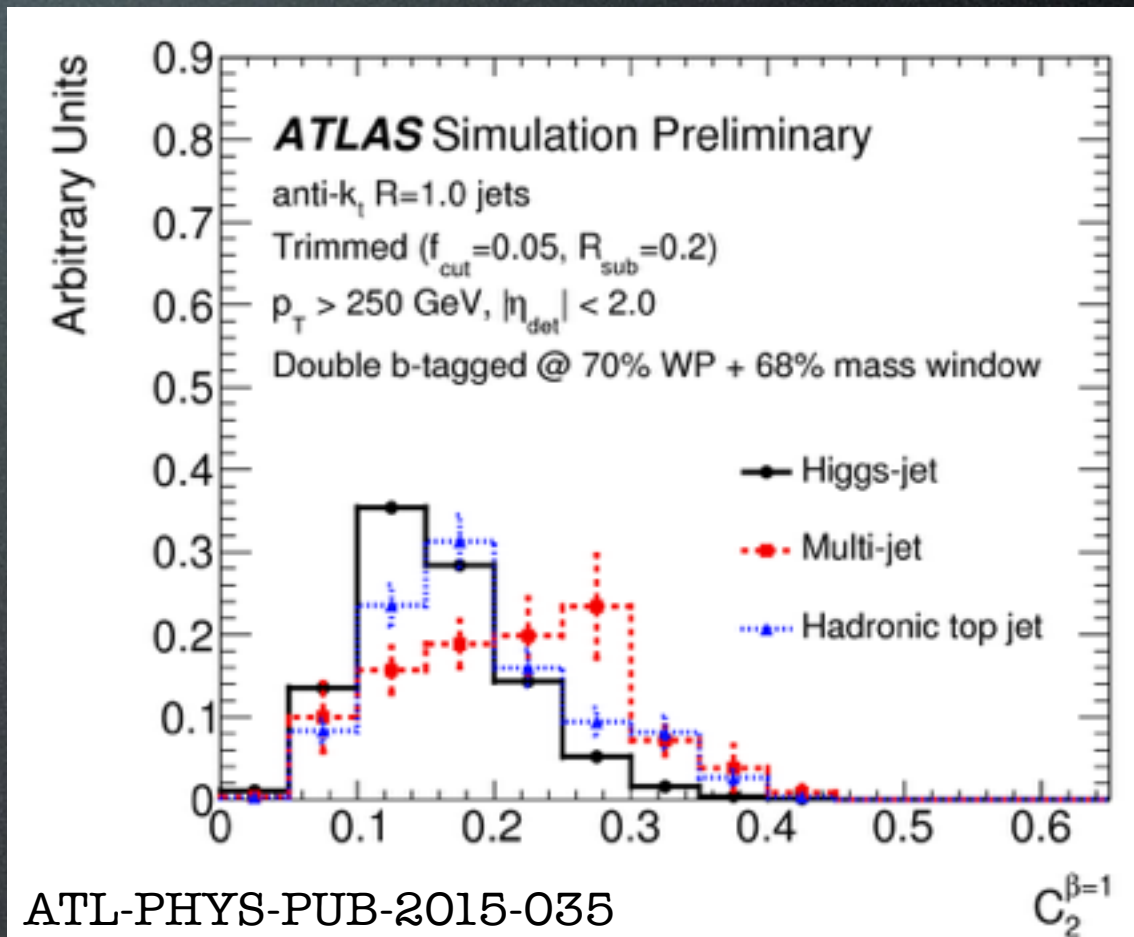
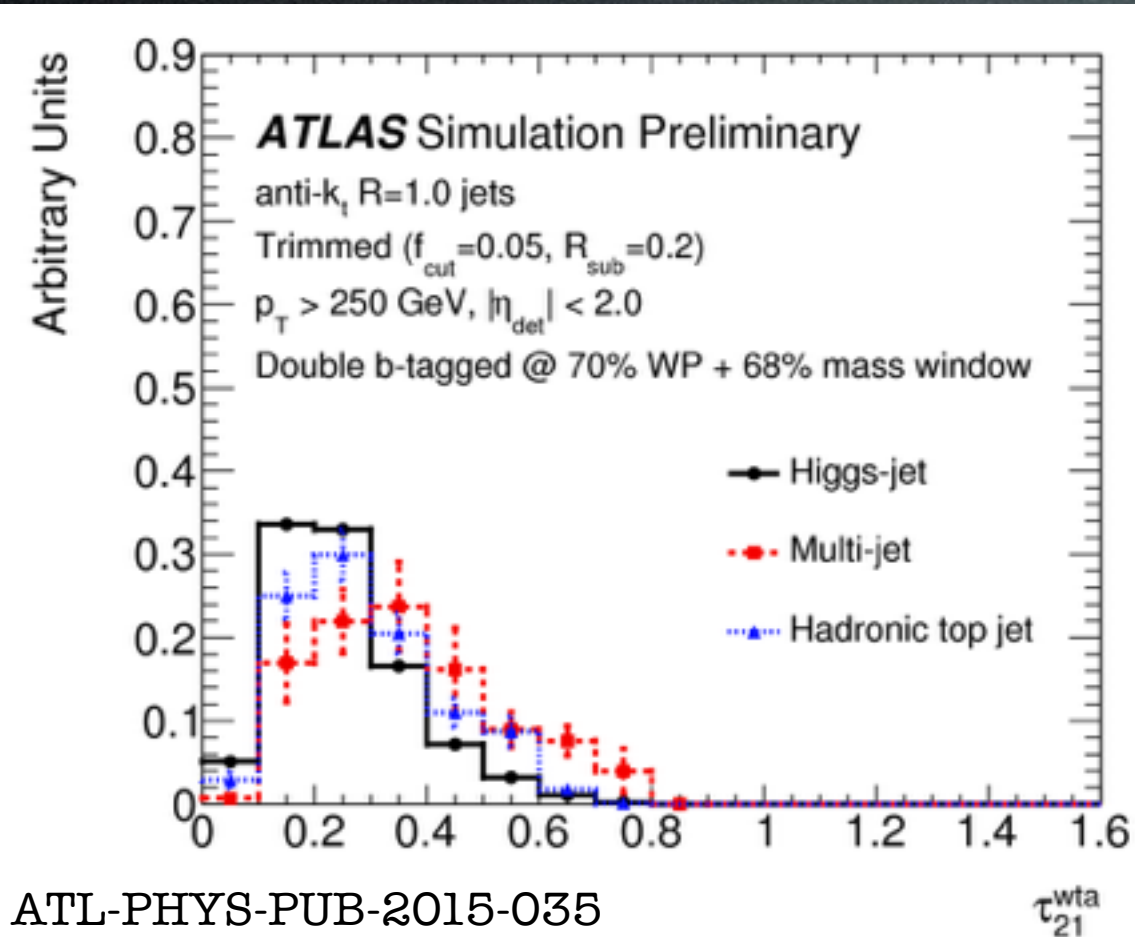


ATL-PHYS-PUB-2015-035

Fatjet mass can be a useful discriminating variable



# Boosted (tt)Higgs to (tt)bbbar



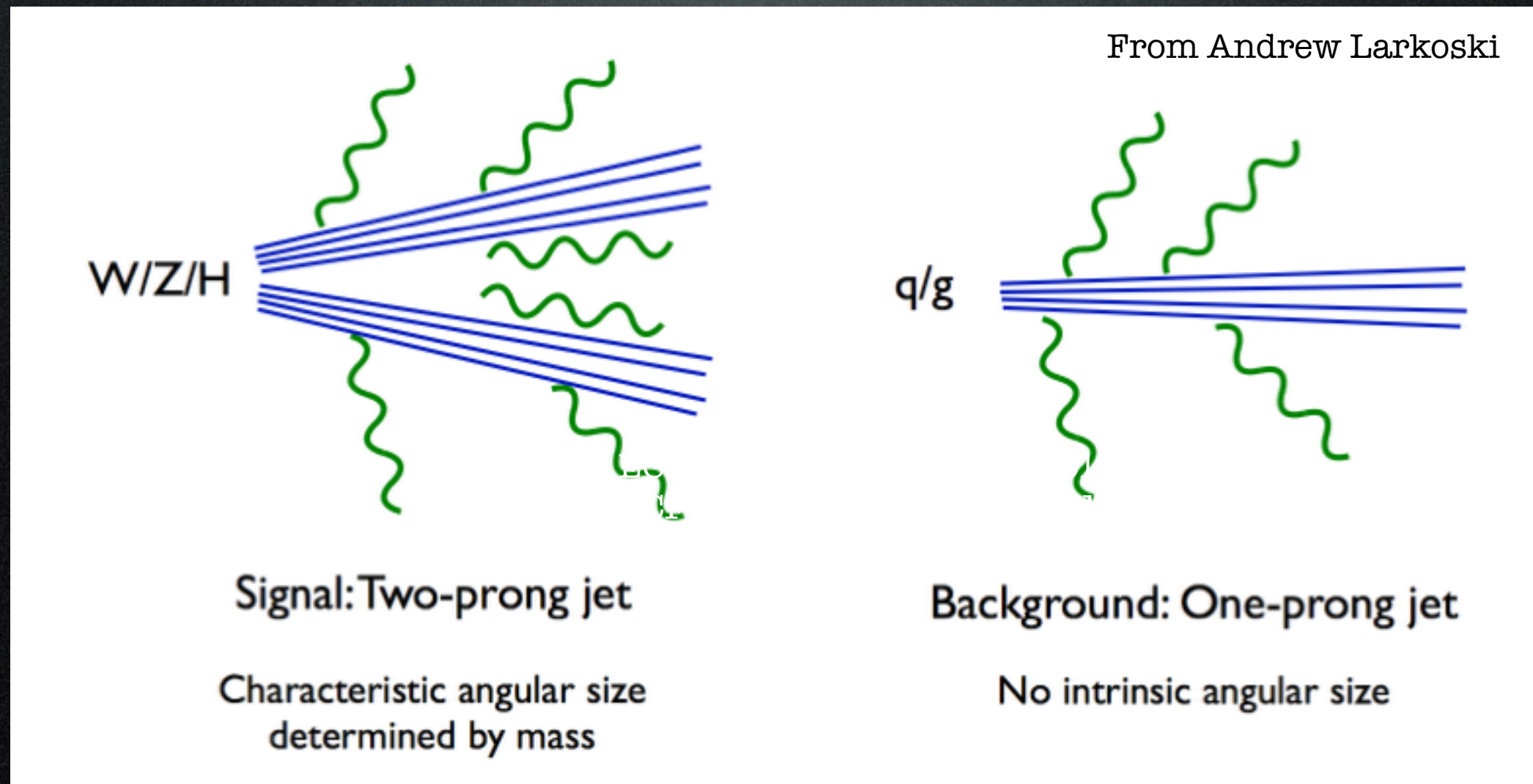
Mass can be combined with other substructure variables

Depends on specific analysis what efficiency/rejection  
working point is needed



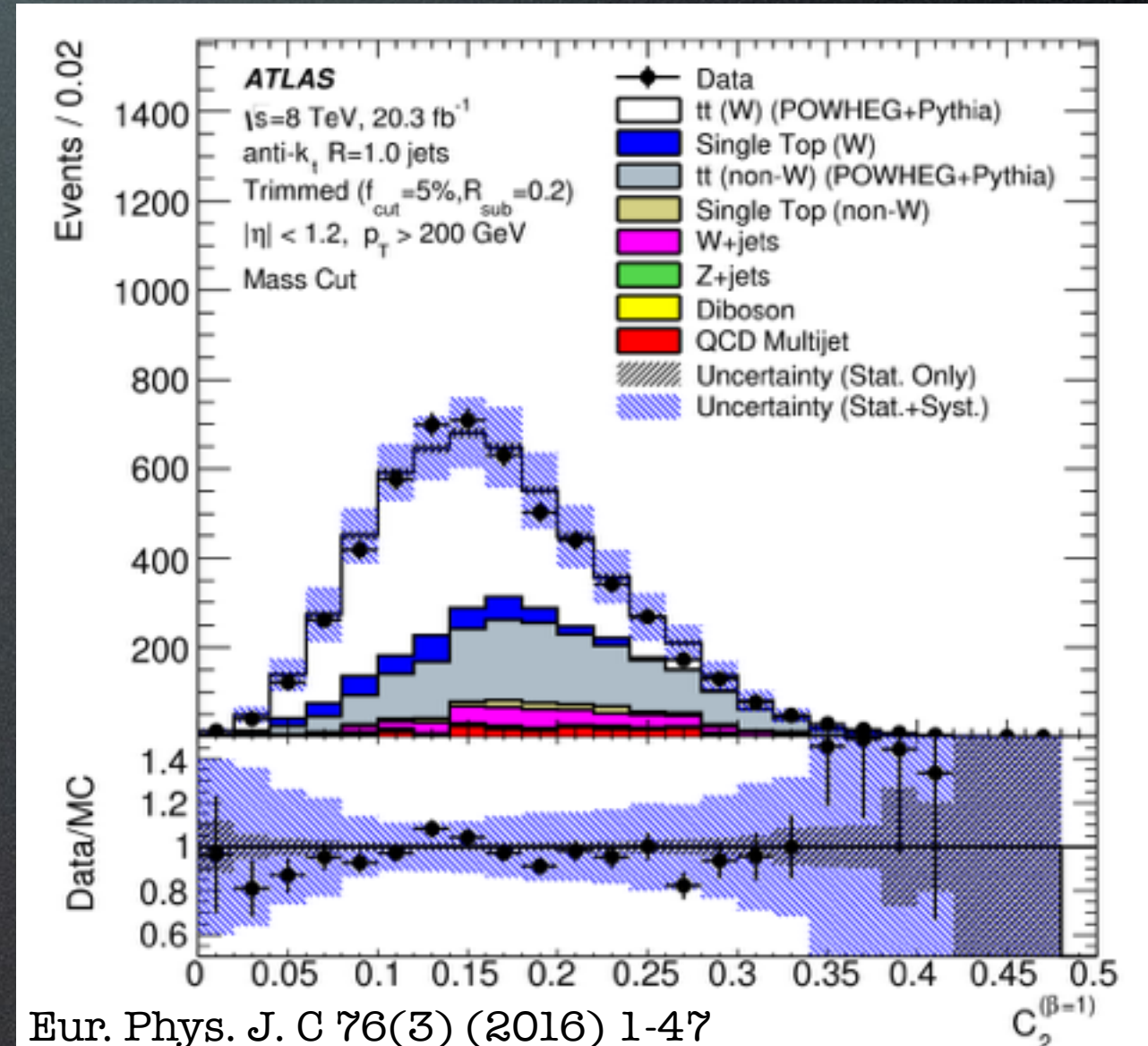
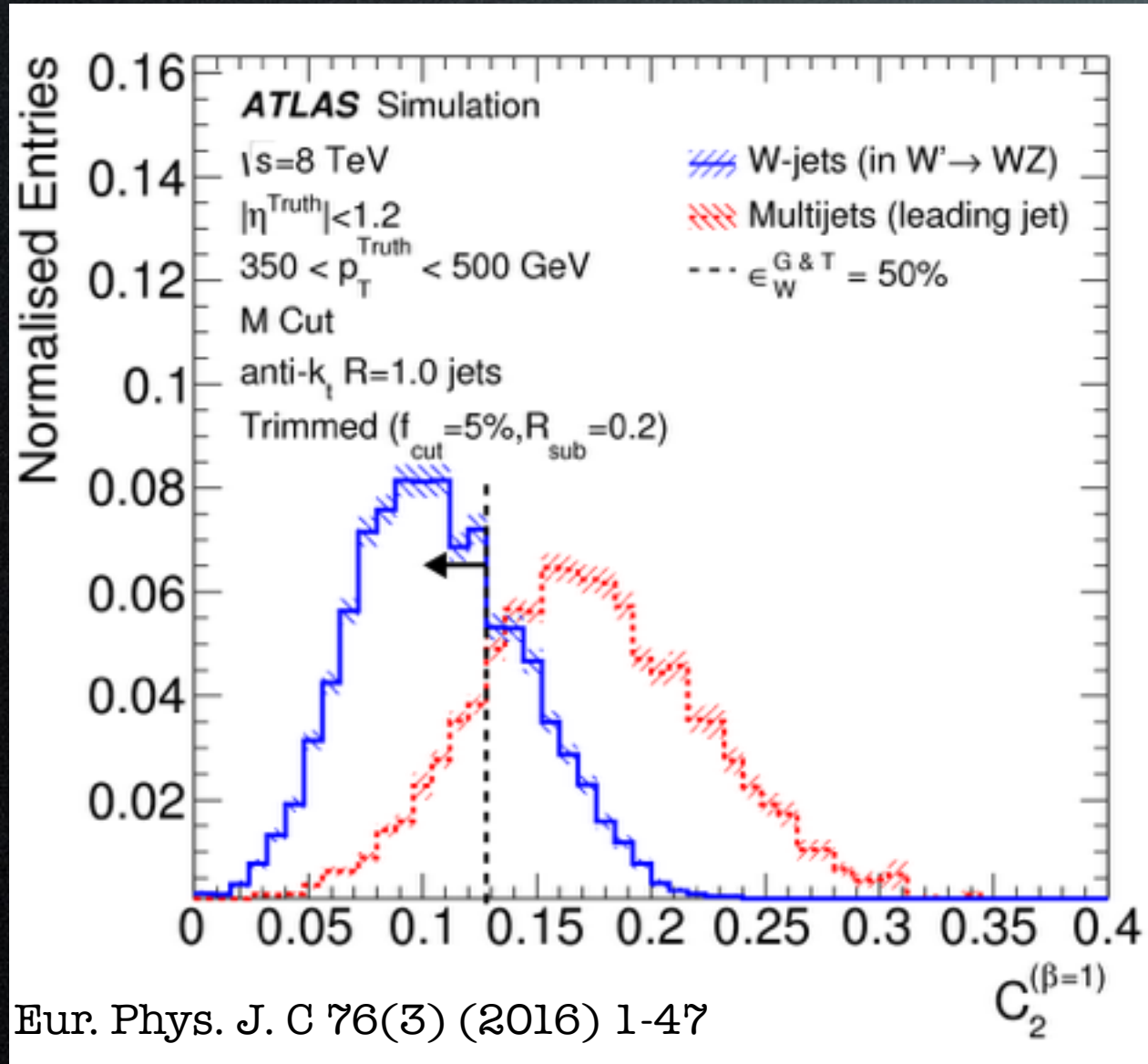
# Energy Correlation Function (or jet substructure without trees)

Discriminate between:



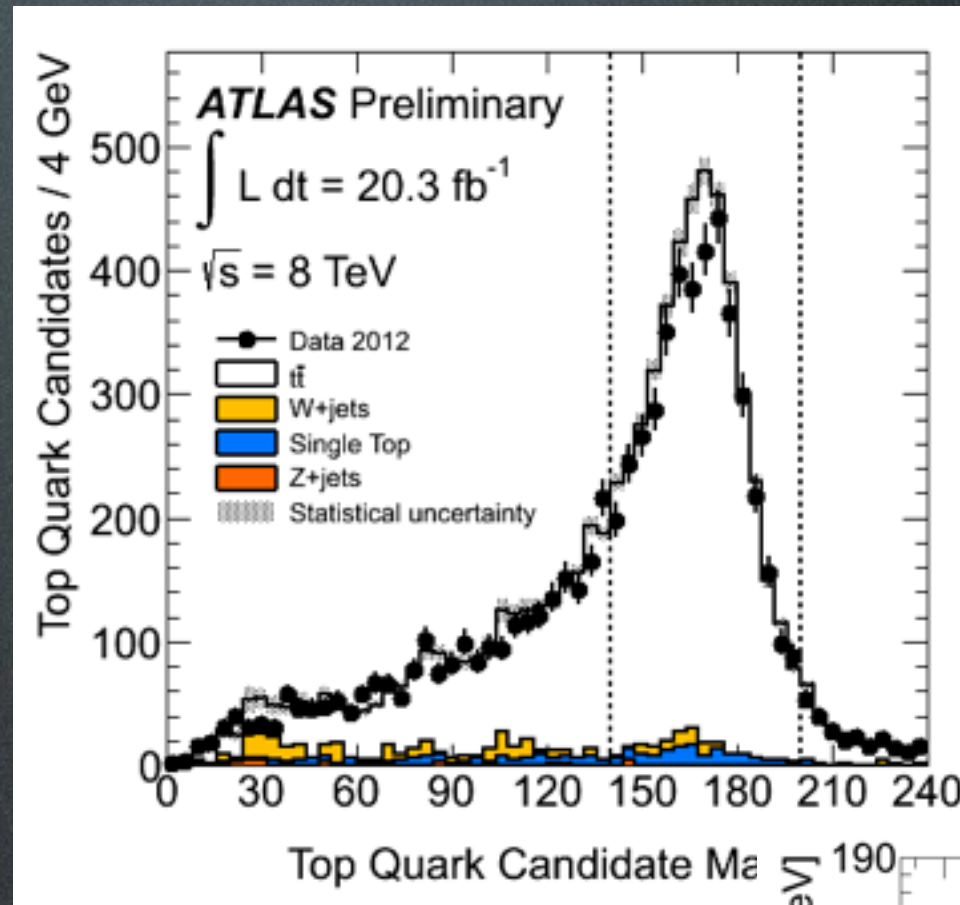
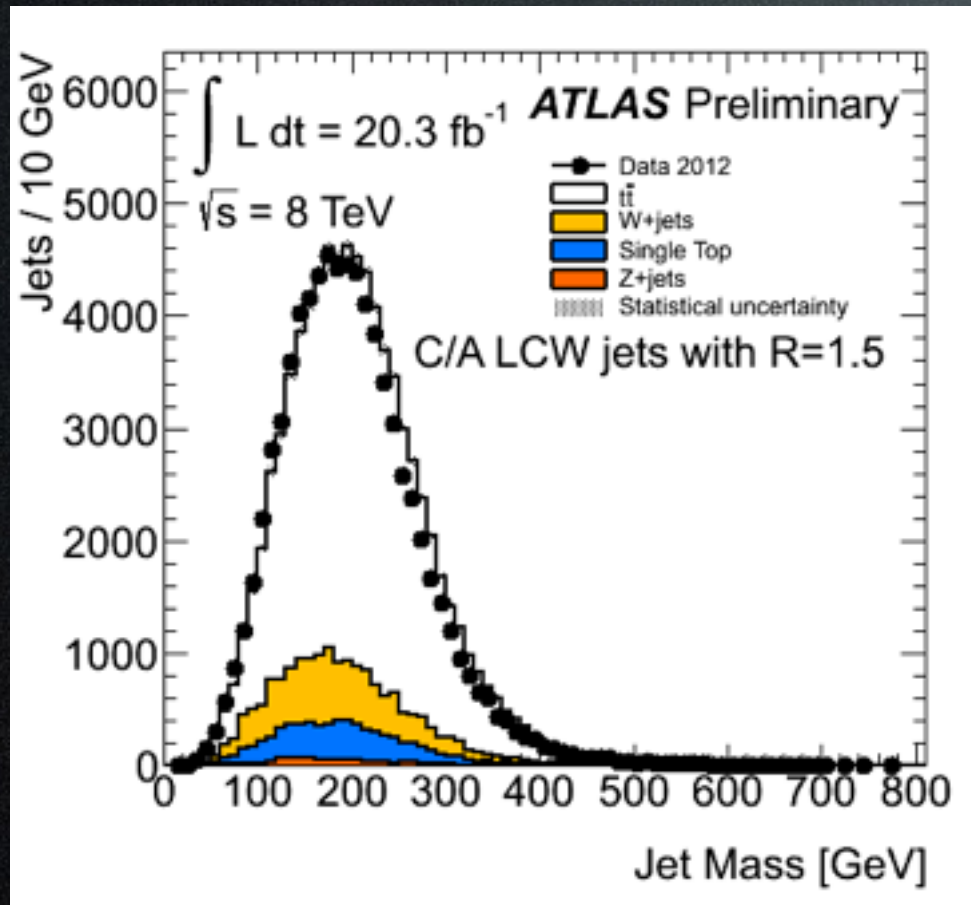


# ECF Results



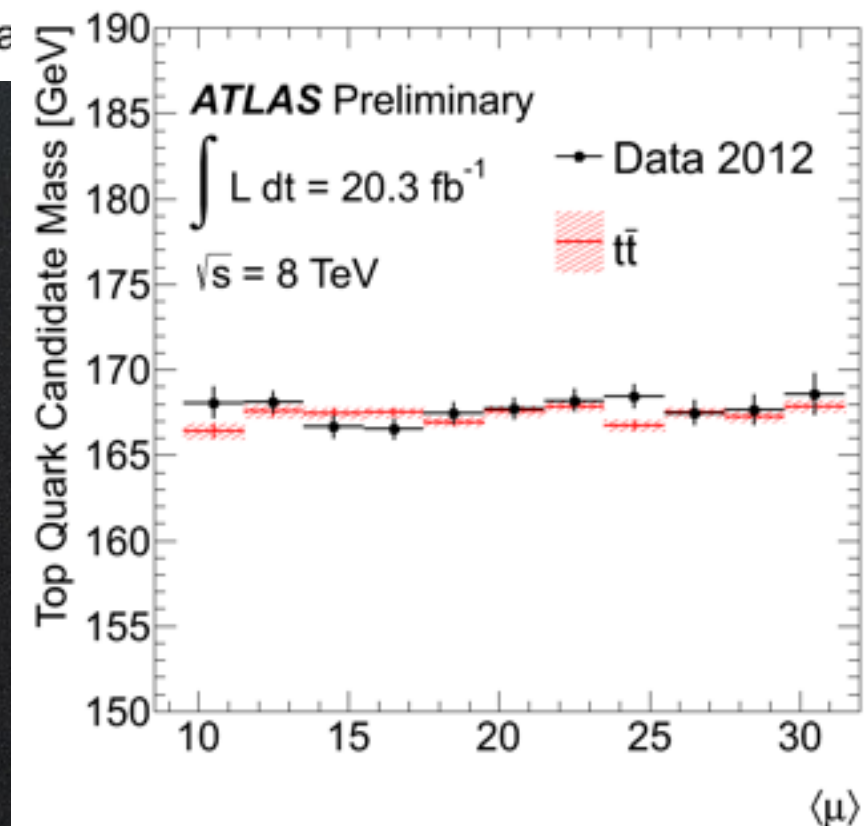


# HEPTopTagger Performance



Before and after tagging by  
HepTopTagger

Pileup resilience

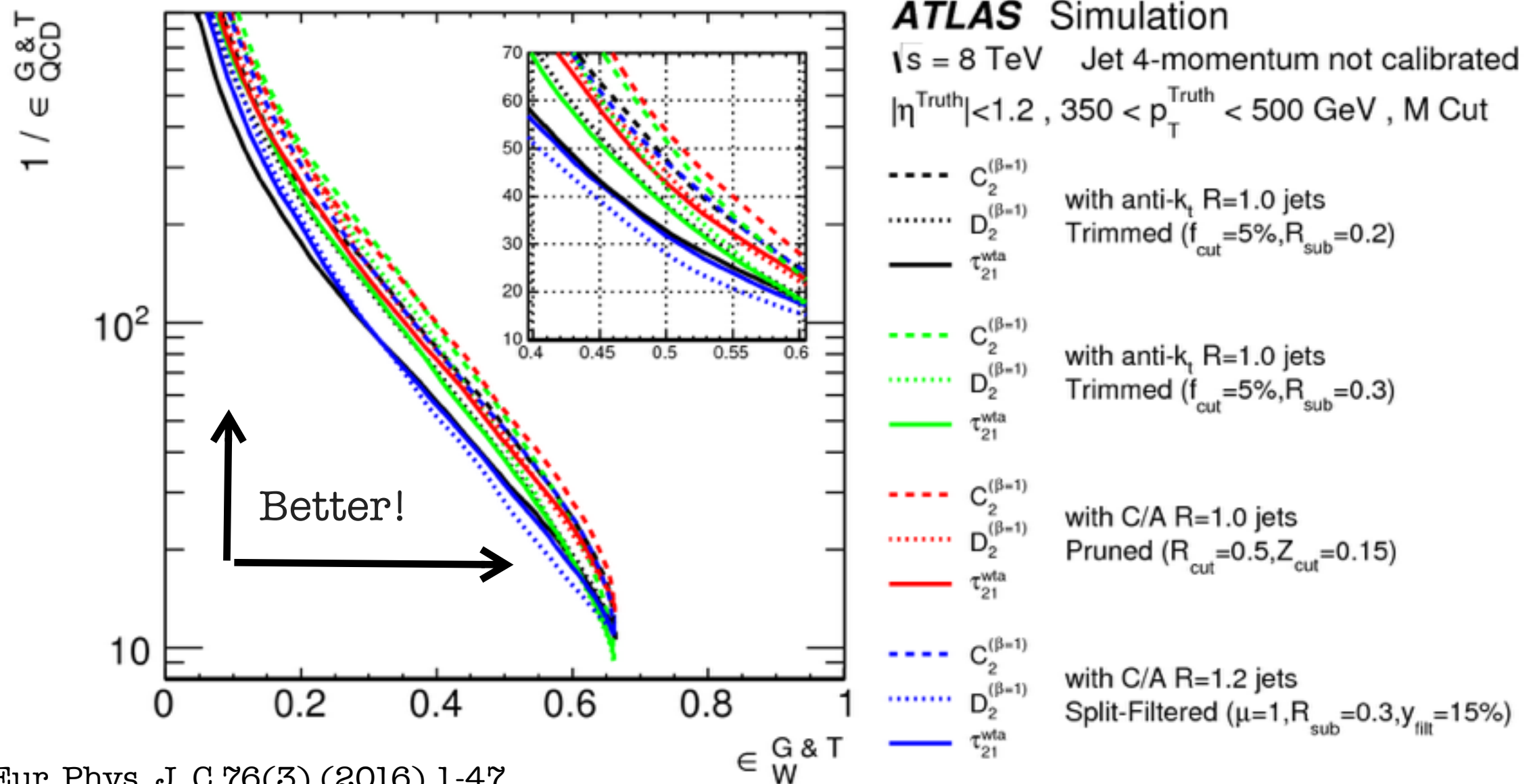




# Tagger Comparisons

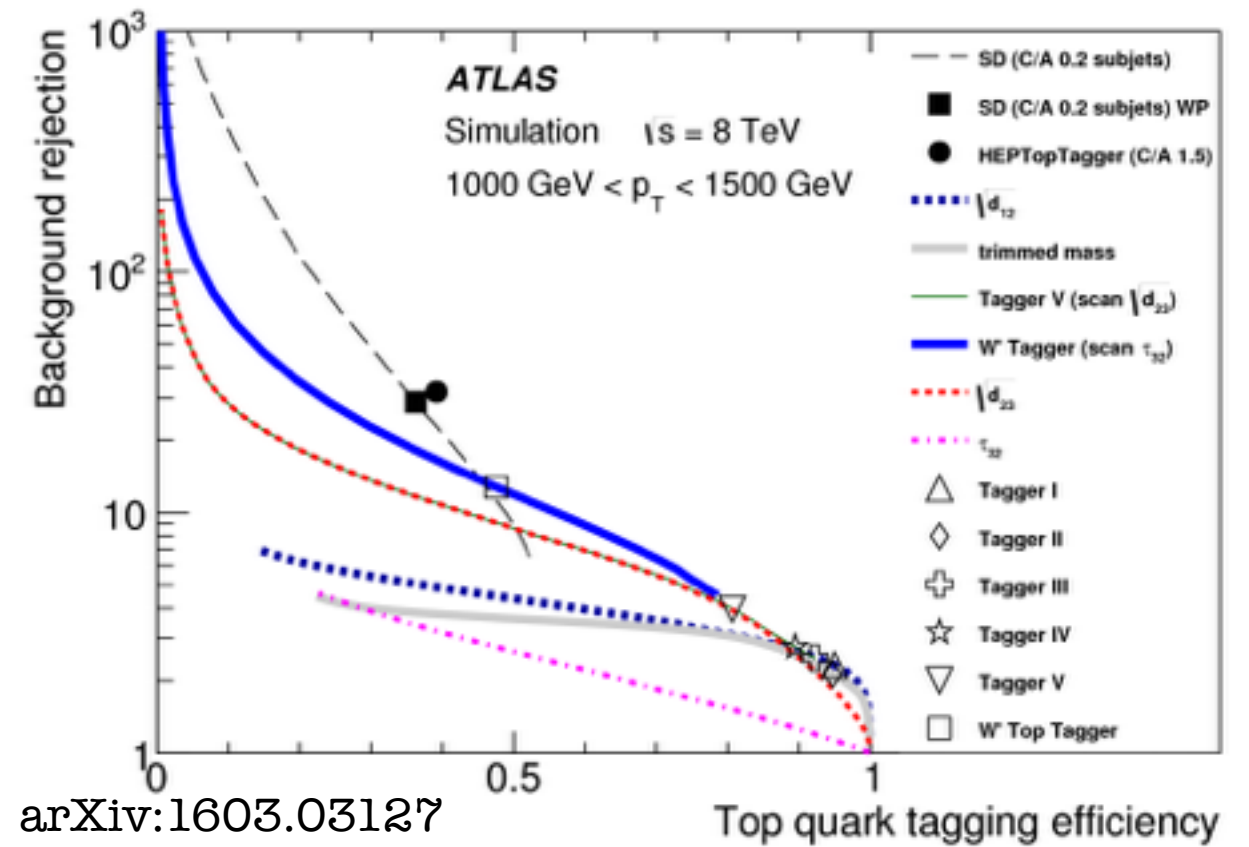
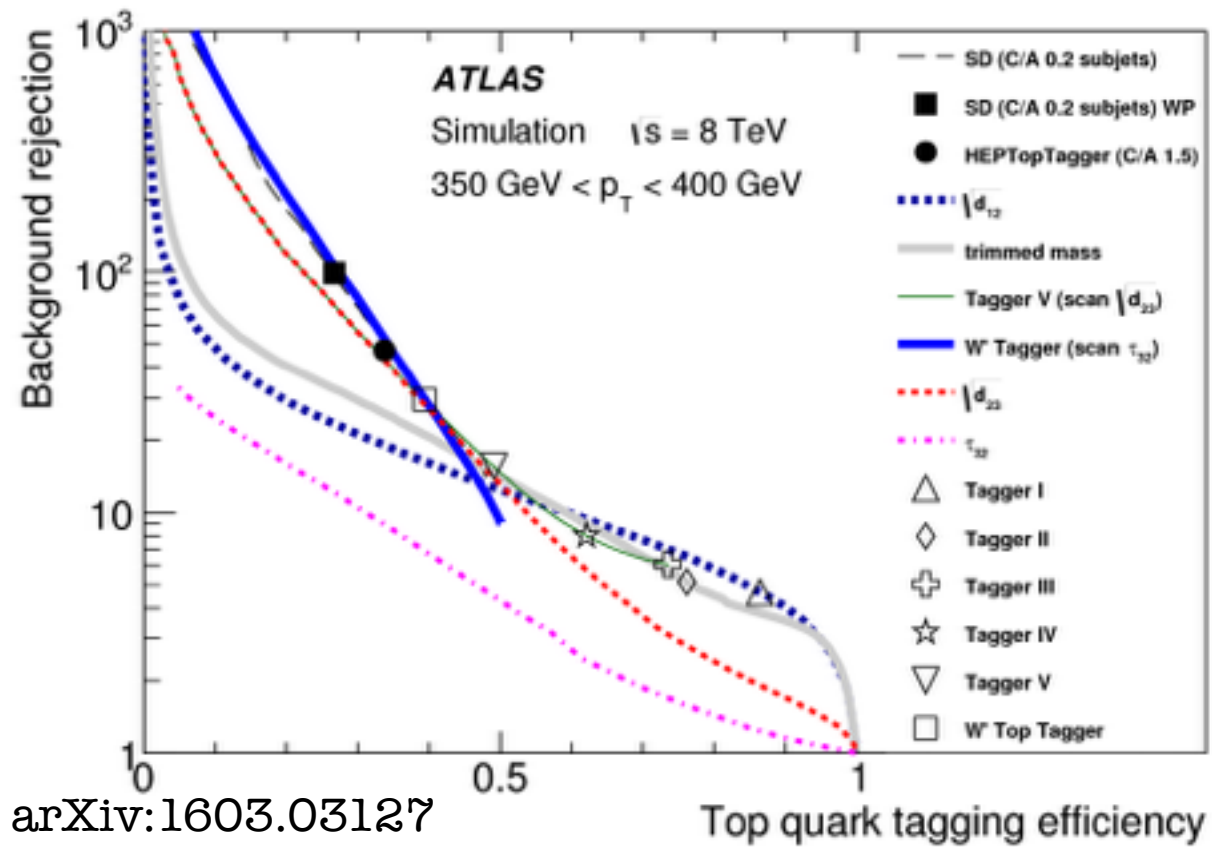


# W tagging Summary





# Top tagging Summary



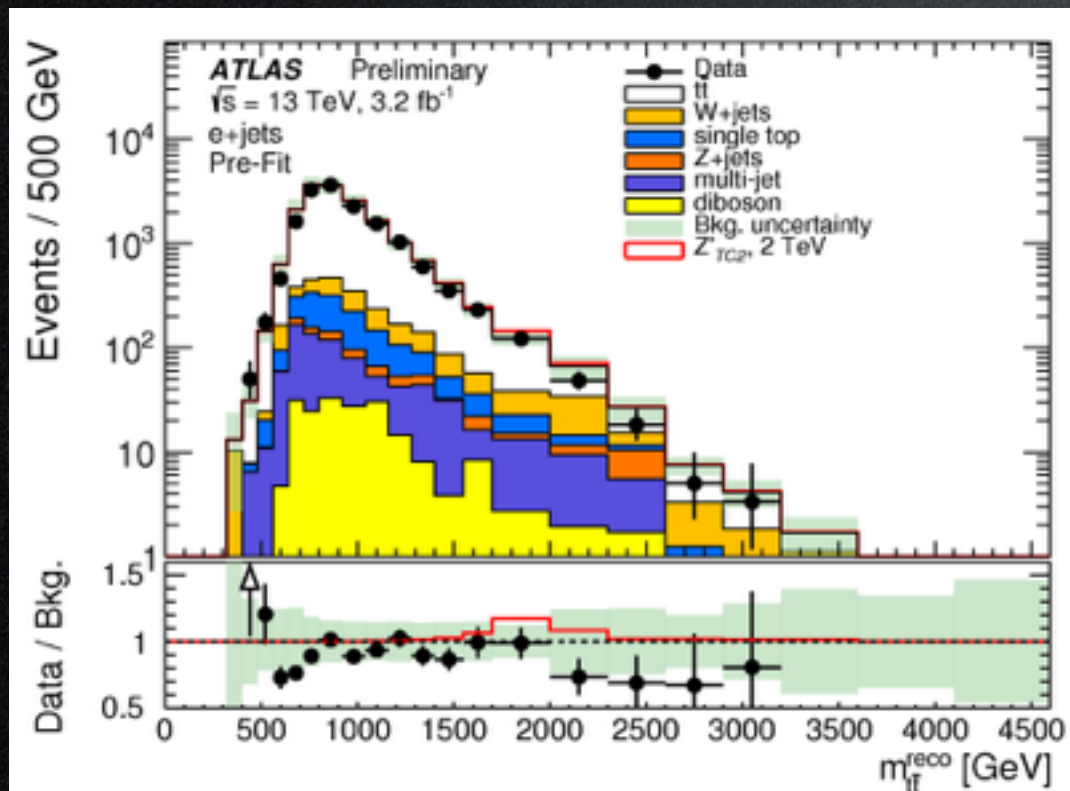
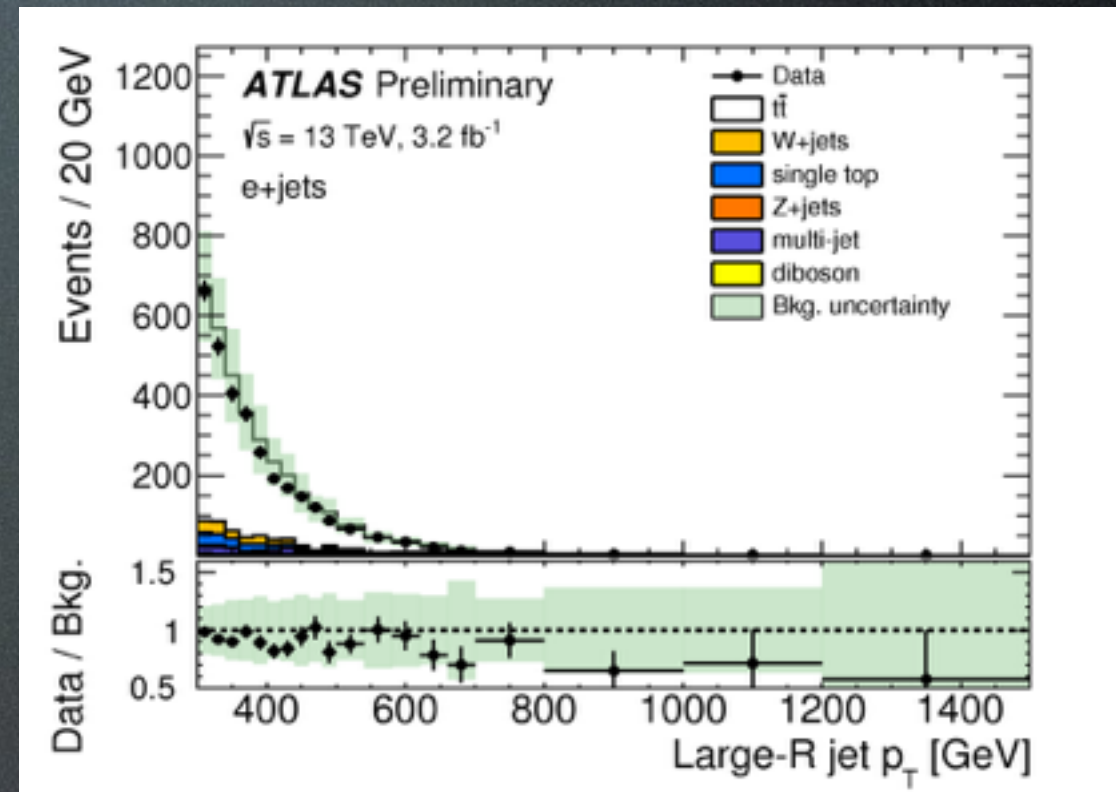
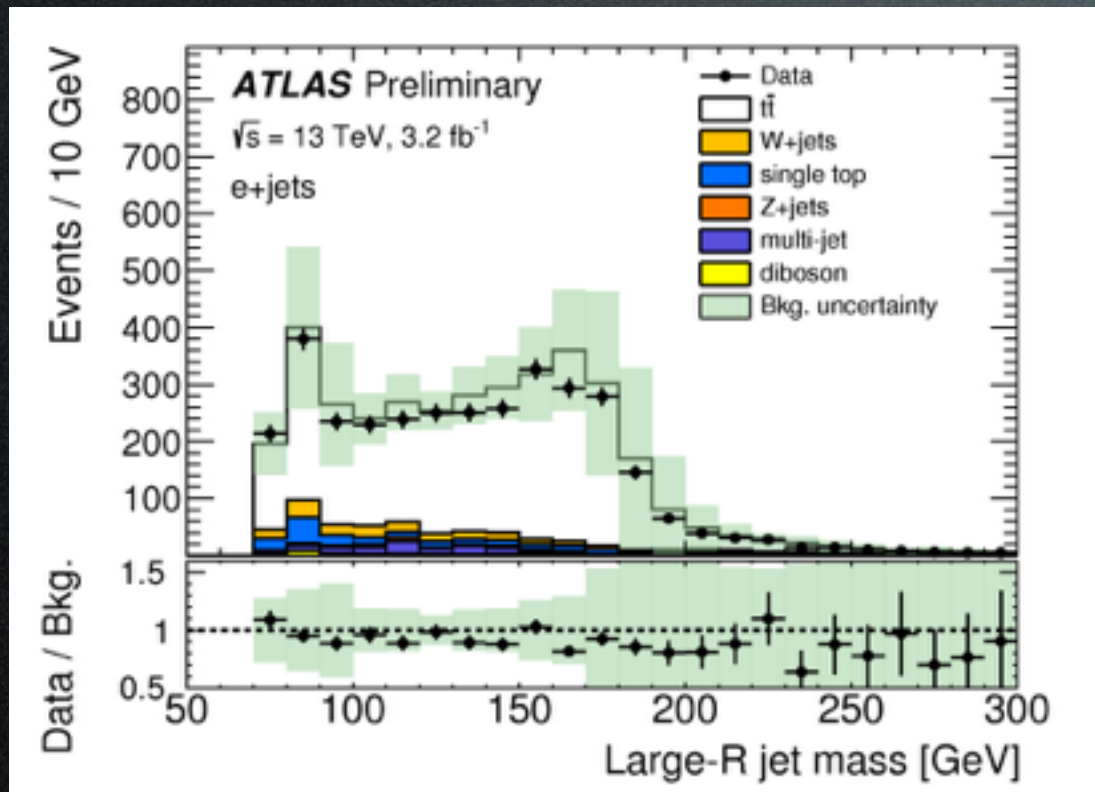
Better top quark finding efficiency with SD at the same rejection of multijets when compared to other taggers



# Example Test Cases



# ttbar Resonance



Top tagged by using  $\text{mass} + \tau_{32}$

80% efficiency working point

Used  $R=0.2$  track jets for b-tagging



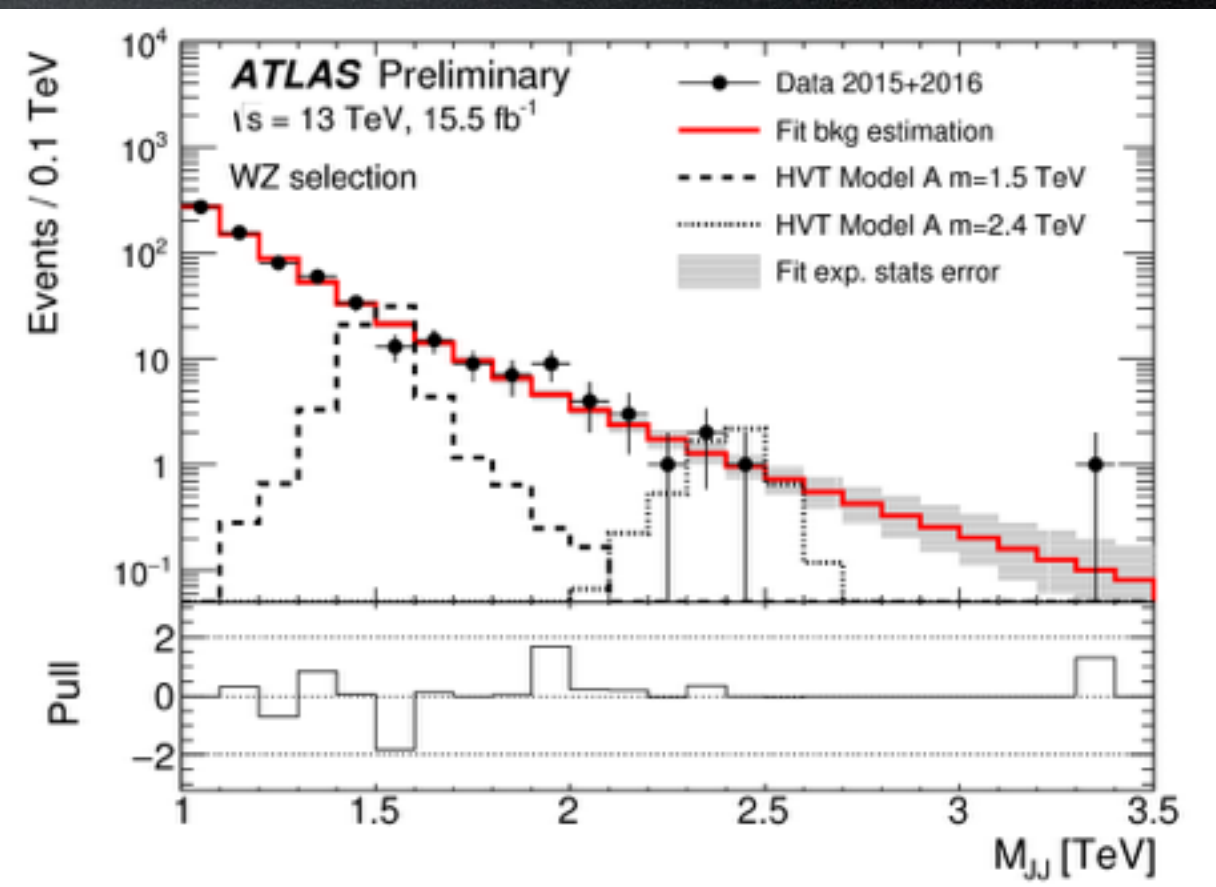
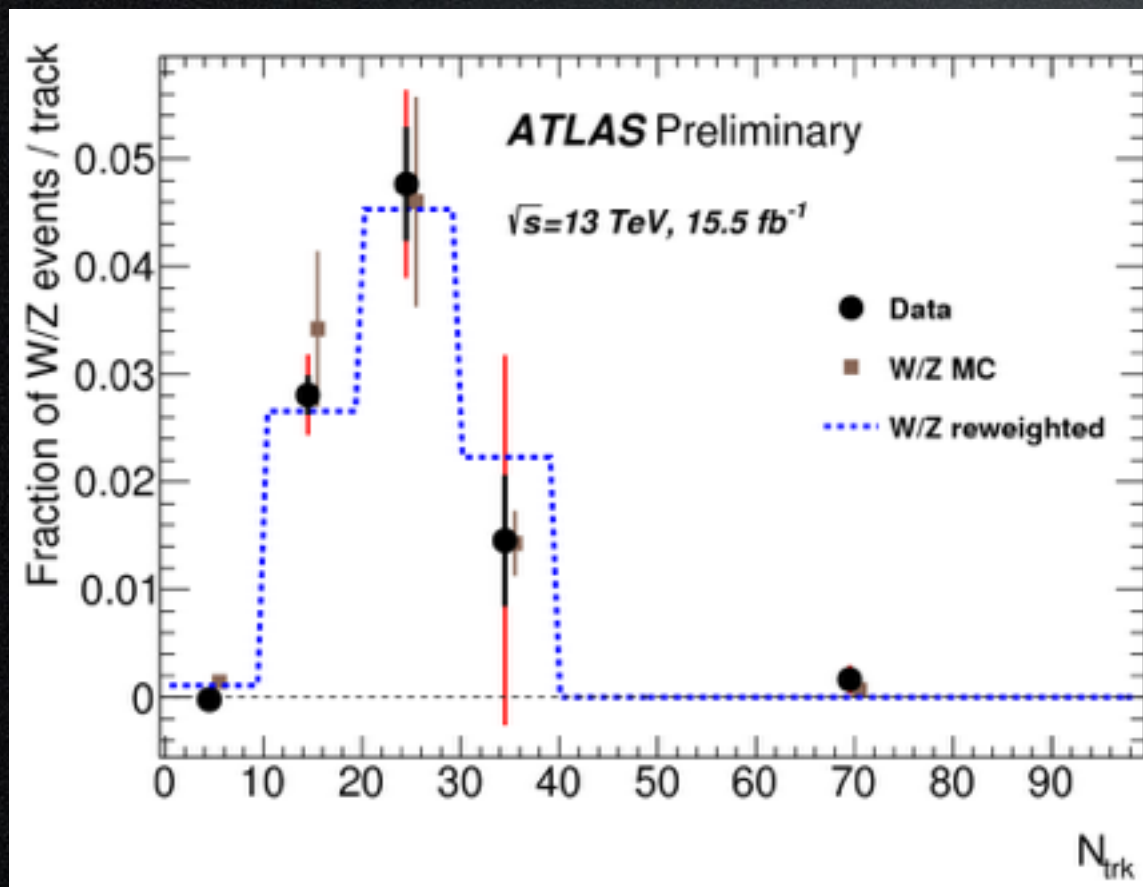
# Vector Boson Pair Production

Event Selection: pair of fatjets, leading  $p_T > 450$  GeV,  $m_{JJ} > 1$  TeV,  $\Delta y_{12} < 1.2$   
(t-channel bg, s-channel signal)

Boson tagged by using mass ( $\pm 15$  GeV around W/Z peak) + D2 ( $\beta=1$ )

50% efficiency working point ( $p_T$  dependent)

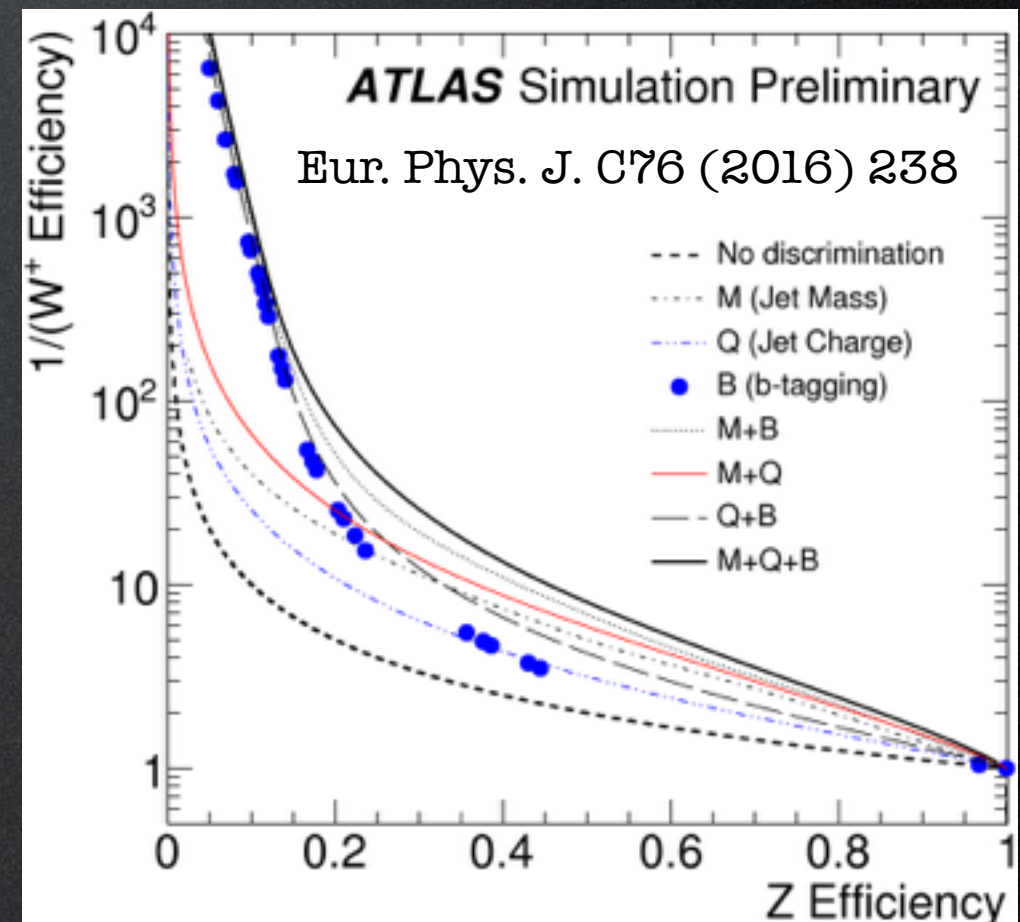
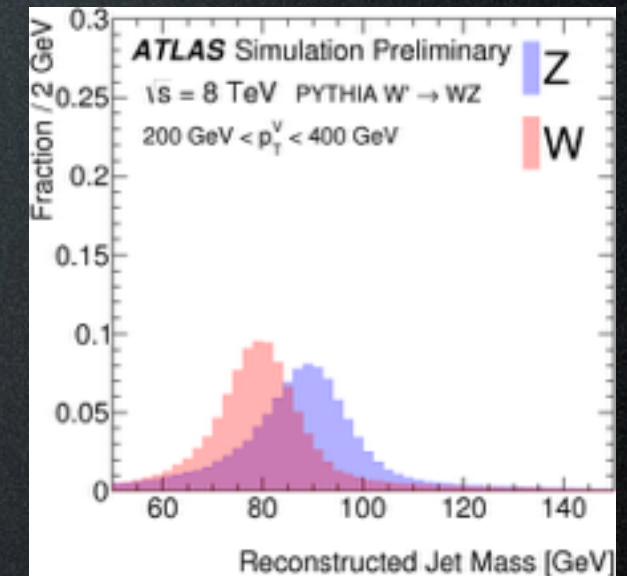
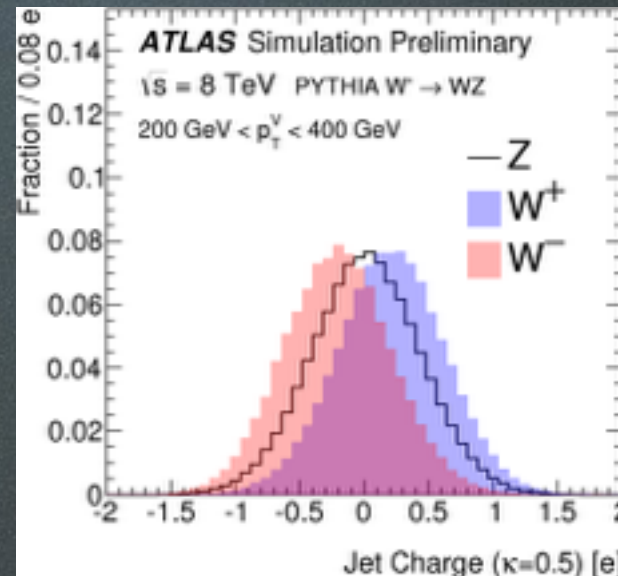
Additional requirement:  $N_{\text{trk}} < 30$  (higher fraction of gluon jets in background, different fragmentation)





# W/Z Discrimination

- Jet mass, charge and b-tagging used to construct a discriminant
- Only in simulation, although data models the observables well.





# Summary

- Many signatures of new physics involve hadronically decaying boosted top quark or Higgs/ $W/Z$  boson(s), either by itself or adding sensitivity to the resolved channel.
- Higher kinematic reach at LHC Run 2 (and beyond) for such boosted objects.
- Pileup remains a significant challenge.
- Many theoretical tools, but commissioning them in experiment takes a lot of effort.



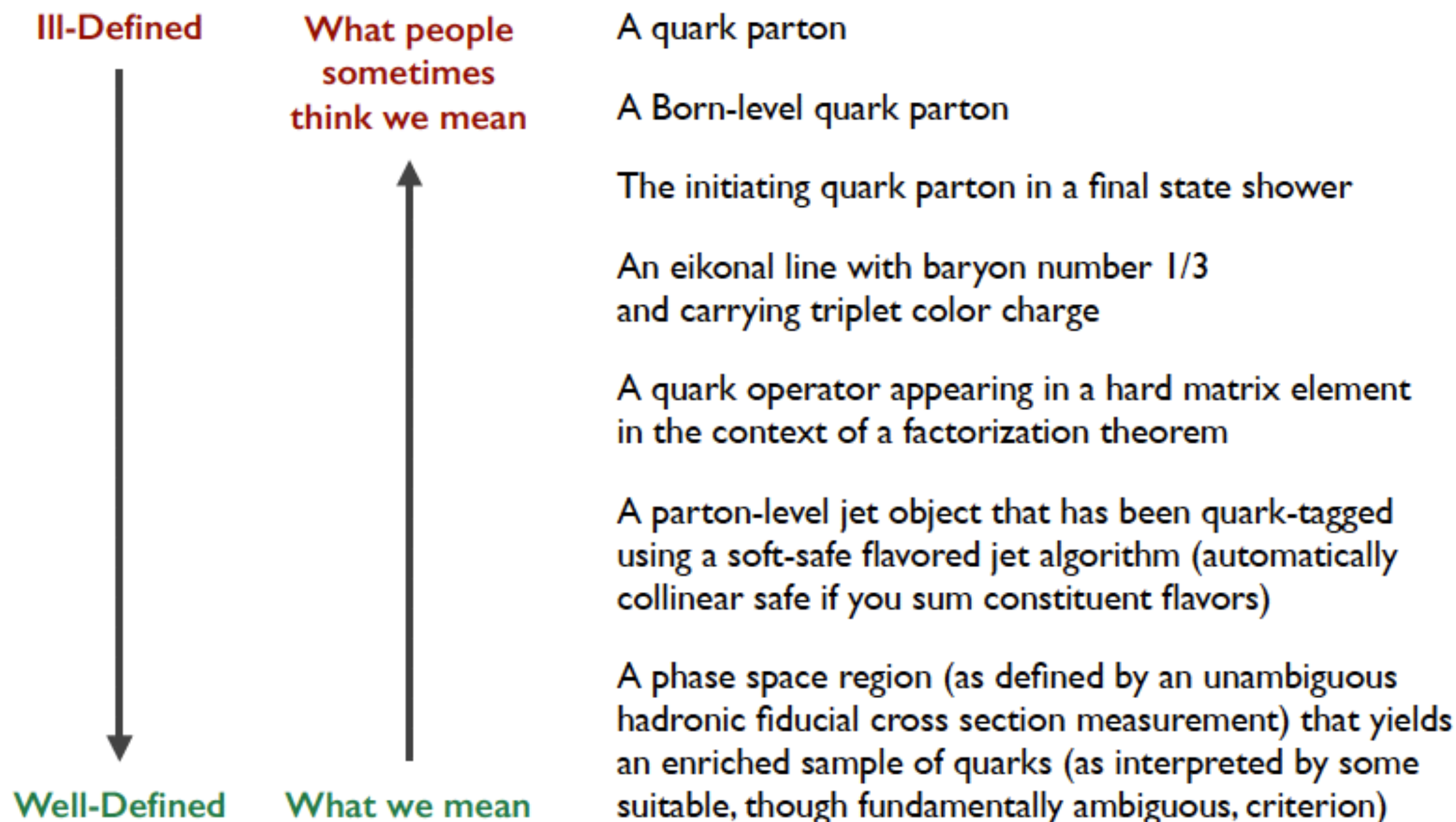
# Supporting Material





# What is a Quark Jet?

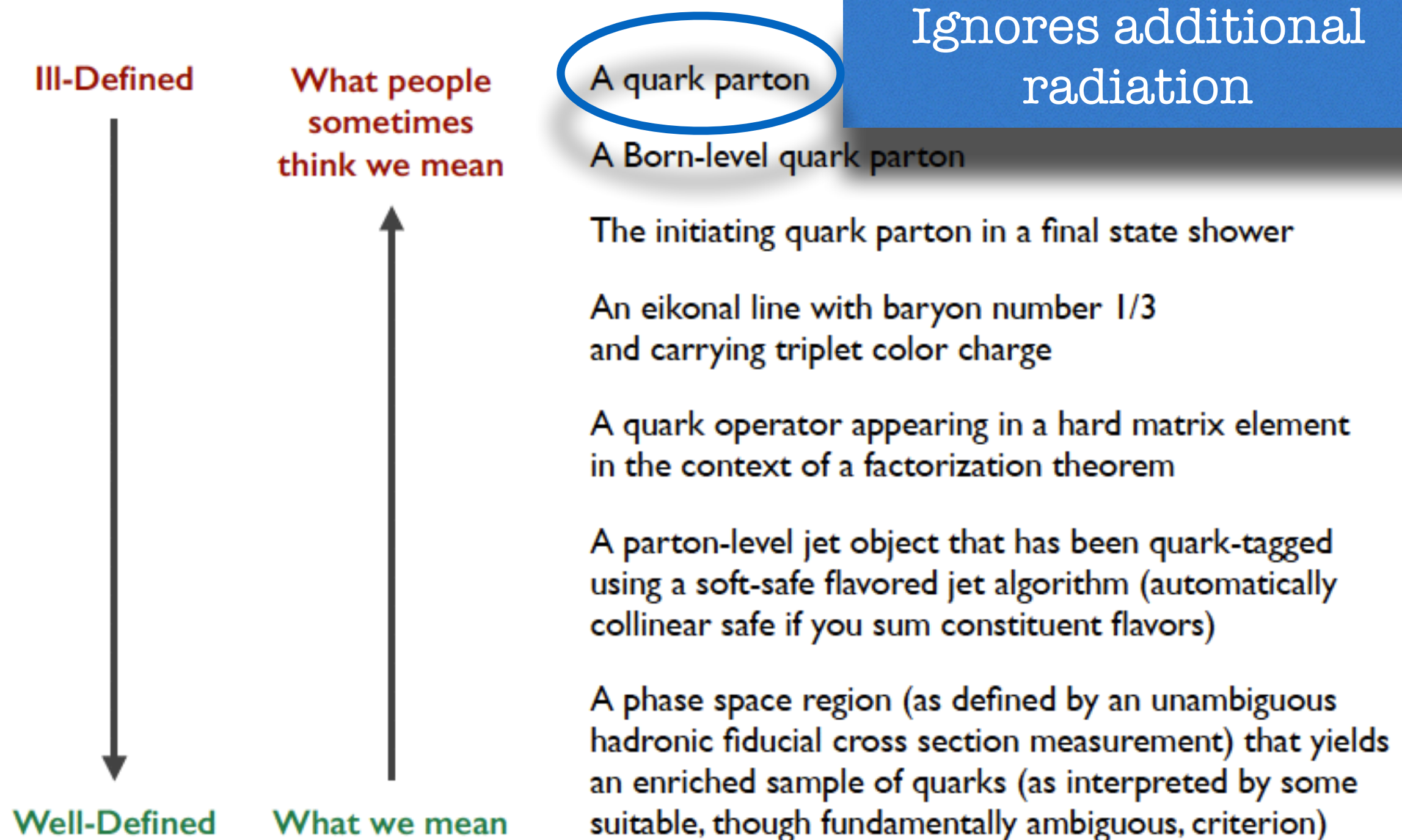
*From lunch/dinner discussions*





# What is a Quark Jet?

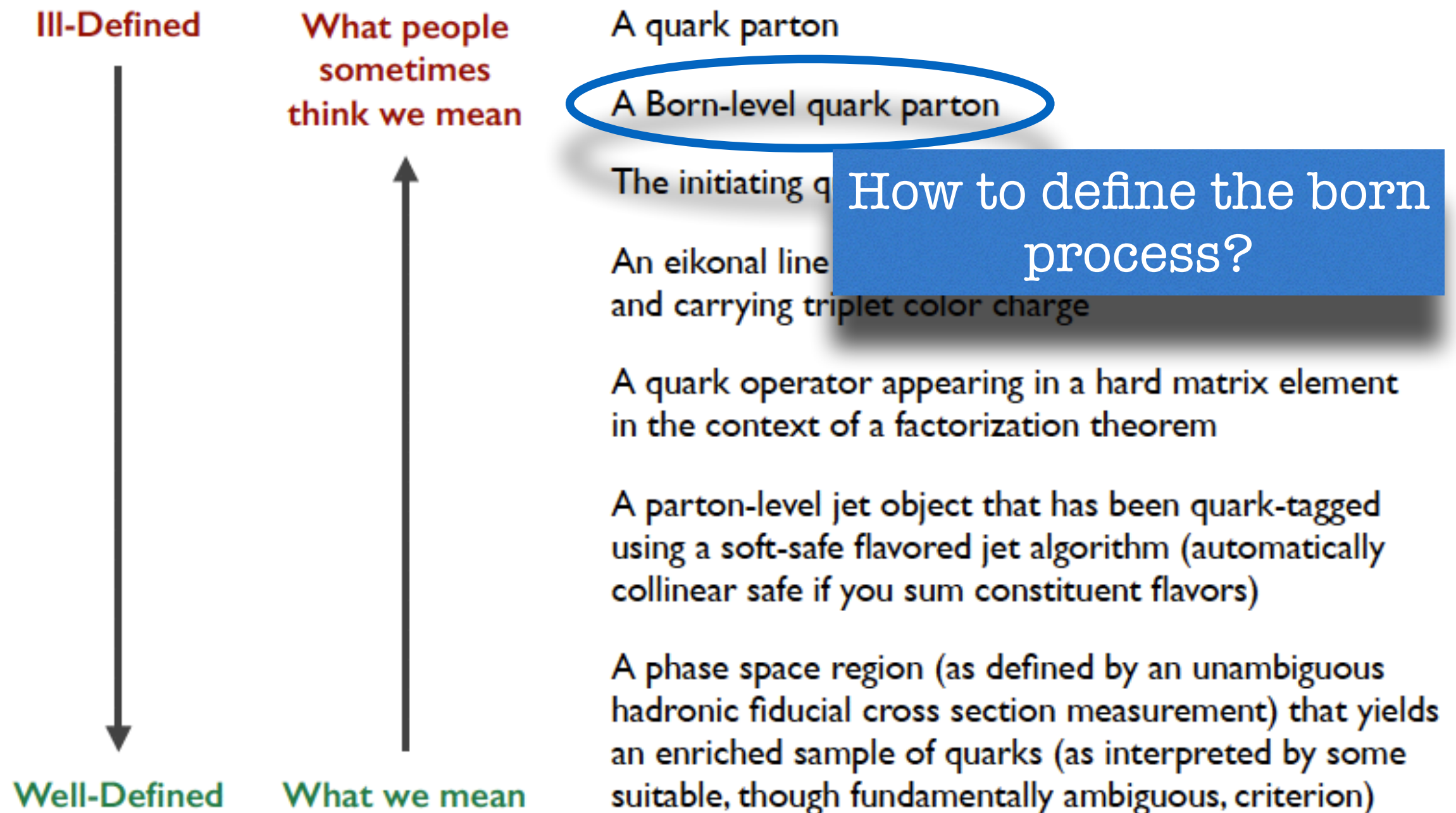
*From lunch/dinner discussions*





# What is a Quark Jet?

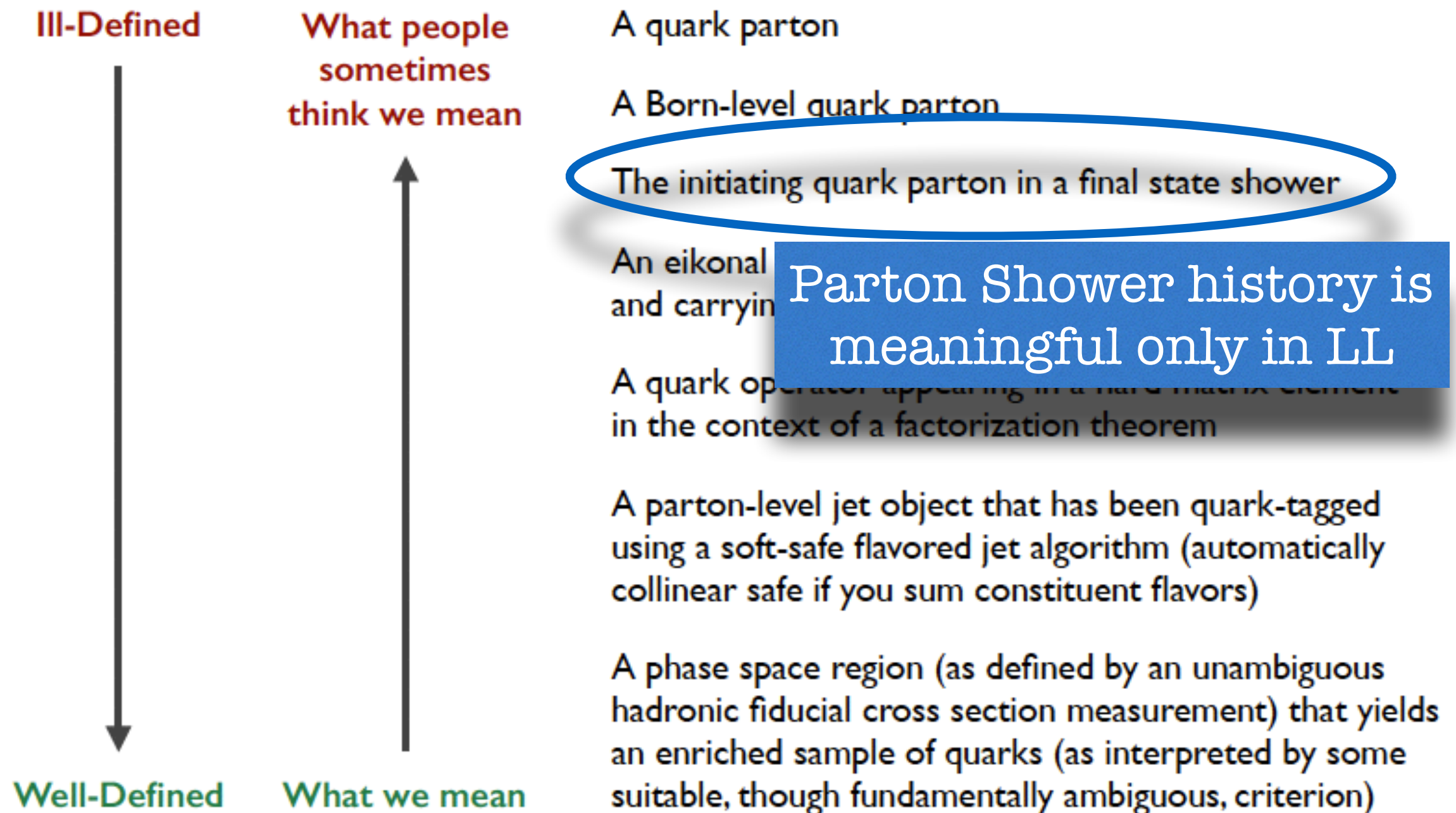
*From lunch/dinner discussions*





# What is a Quark Jet?

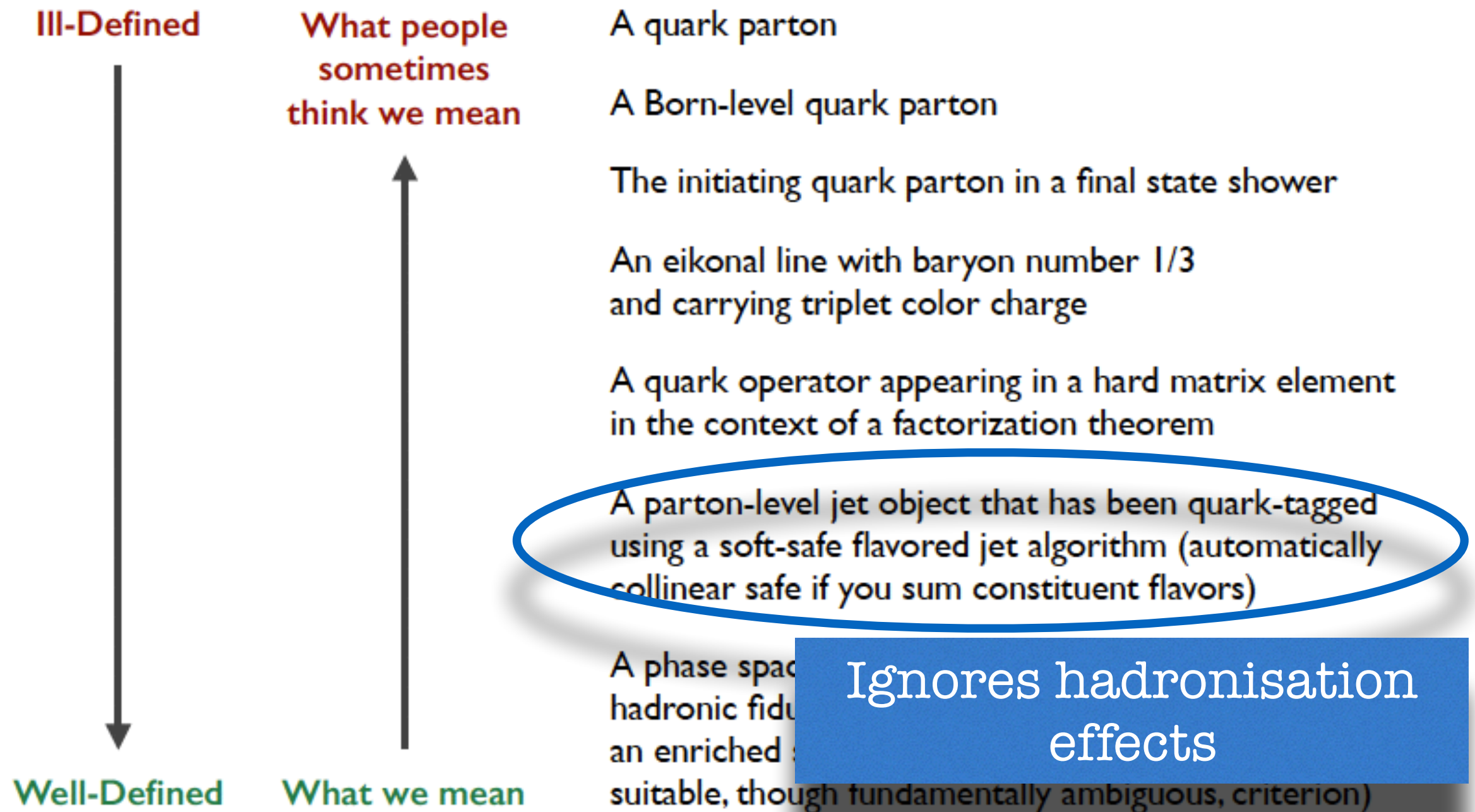
*From lunch/dinner discussions*





# What is a Quark Jet?

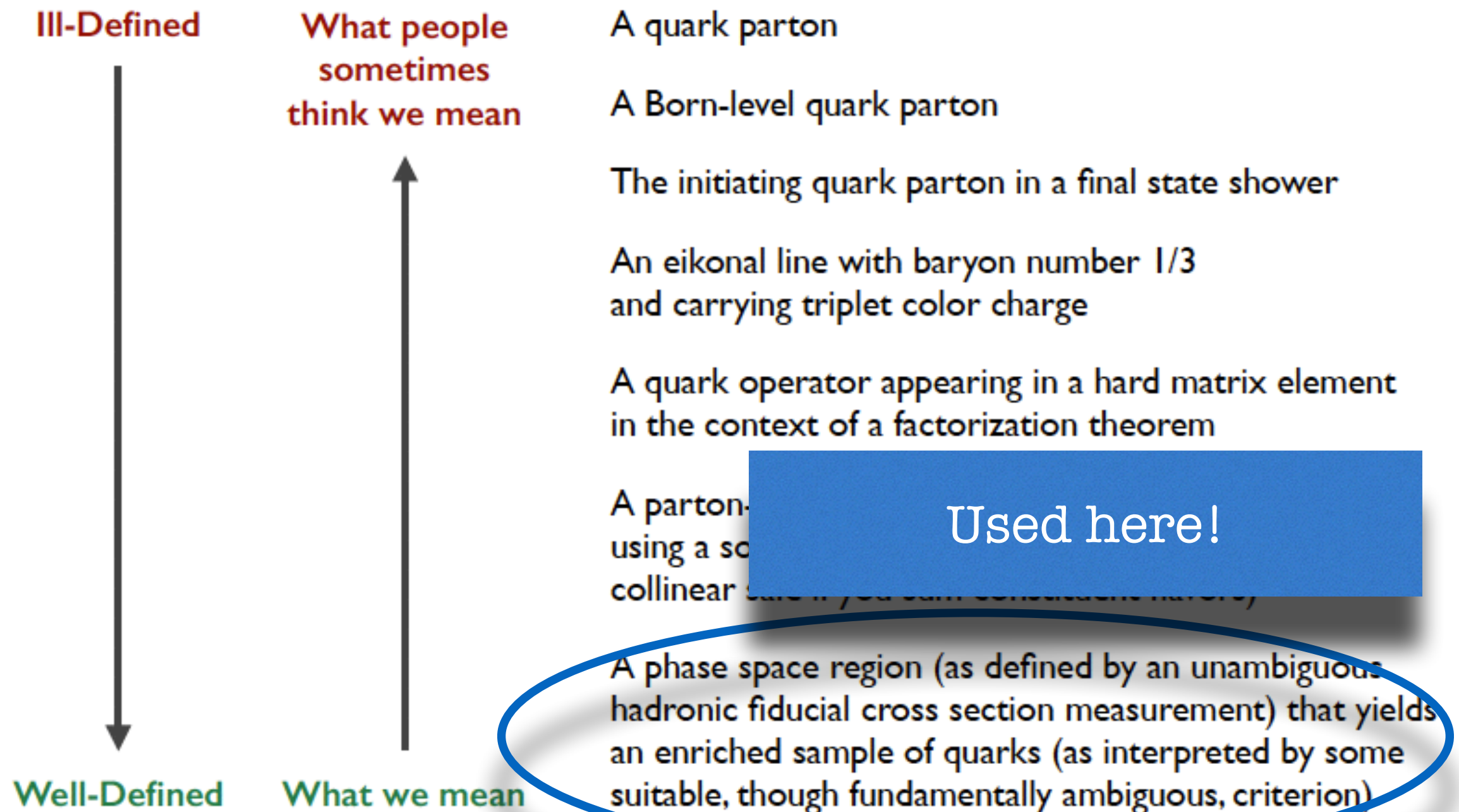
*From lunch/dinner discussions*





# What is a Quark Jet?

*From lunch/dinner discussions*





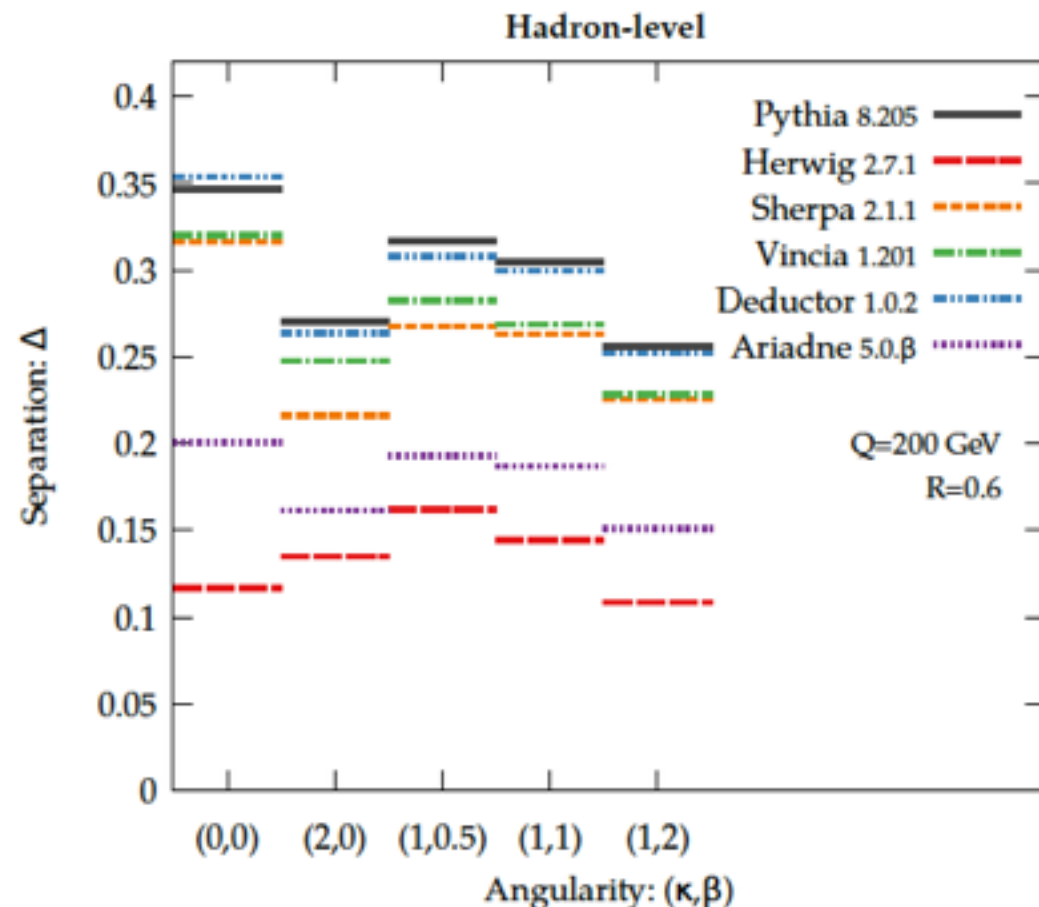
# Les Houches Study

Generalised angularities:

$$\lambda_{\beta}^{\kappa} = \sum_{i \in \text{jet}} z_i^{\kappa} \theta_i^{\beta},$$

Discriminator:

$$\Delta = \frac{1}{2} \int d\lambda \frac{(p_q(\lambda) - p_g(\lambda))^2}{p_q(\lambda) + p_g(\lambda)},$$



Precise radiation pattern  
is very model dependent

Need more measurements



# ECF

Over all constituents (beta: angular exponent):

$$\text{ECF}(1, \beta) = \sum_i p_{Ti}$$

$$\text{ECF}(2, \beta) = \sum_{i < j} p_{Ti} p_{Tj} (R_{ij})^\beta \leftarrow \begin{array}{l} \text{[see Banfi, Salam, Zanderighi;} \\ \text{Jankowiak, Larkoski]} \end{array}$$

$$\text{ECF}(3, \beta) = \sum_{i < j < k} p_{Ti} p_{Tj} p_{Tk} (R_{ij} R_{jk} R_{ki})^\beta$$

$$\text{ECF}(N, \beta) = \sum_{\text{sets of } N} (N \text{ energies}) \times \left( \binom{N}{2} \text{ angles} \right)^\beta$$

$$\text{ECF}(N+1) \ll \text{ECF}(N) \\ \text{for } N \text{ subjects}$$

Define (double) ratio =  $[\text{ECF}(N+1)/\text{ECF}(N)]/[\text{ECF}(N)/\text{ECF}(N-1)]$

$$C_N^{(\beta)} = \frac{\text{ECF}(N+1, \beta) \text{ECF}(N-1, \beta)}{\text{ECF}(N, \beta)^2}$$

Analogous to Nsubjettiness ratio

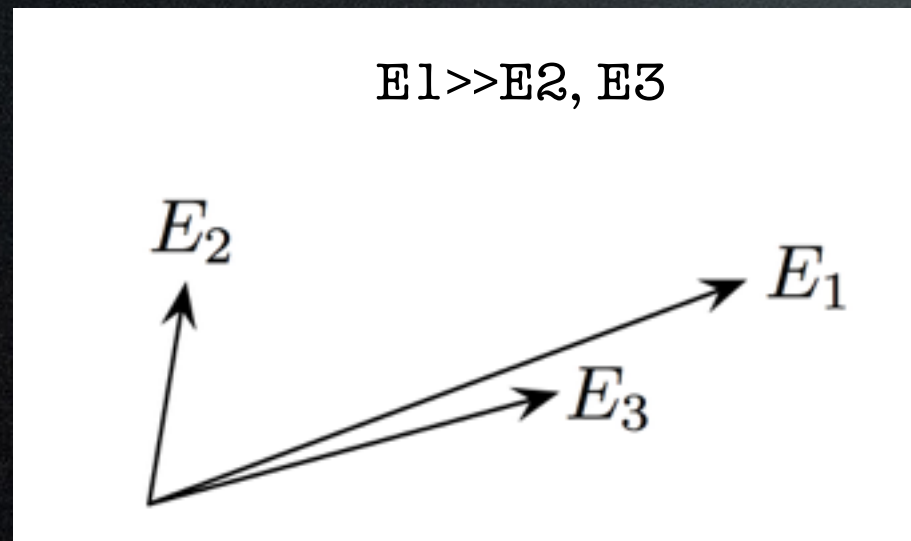
Large  $C_N$ : more than  $N$  subjects, extra radiation is not correlated with leading order  $N$  subjects.

For small  $C_N$ : the additional radiation is soft/collinear



# ECF Discussion

For this multiple soft radiation case,  
with only 1 **real** subjet



$$C_2 > \tau_{21}$$

Nsubjettiness will identify this as  
more 2 subjet-like while ECF will  
identify more as 1 subjet-like

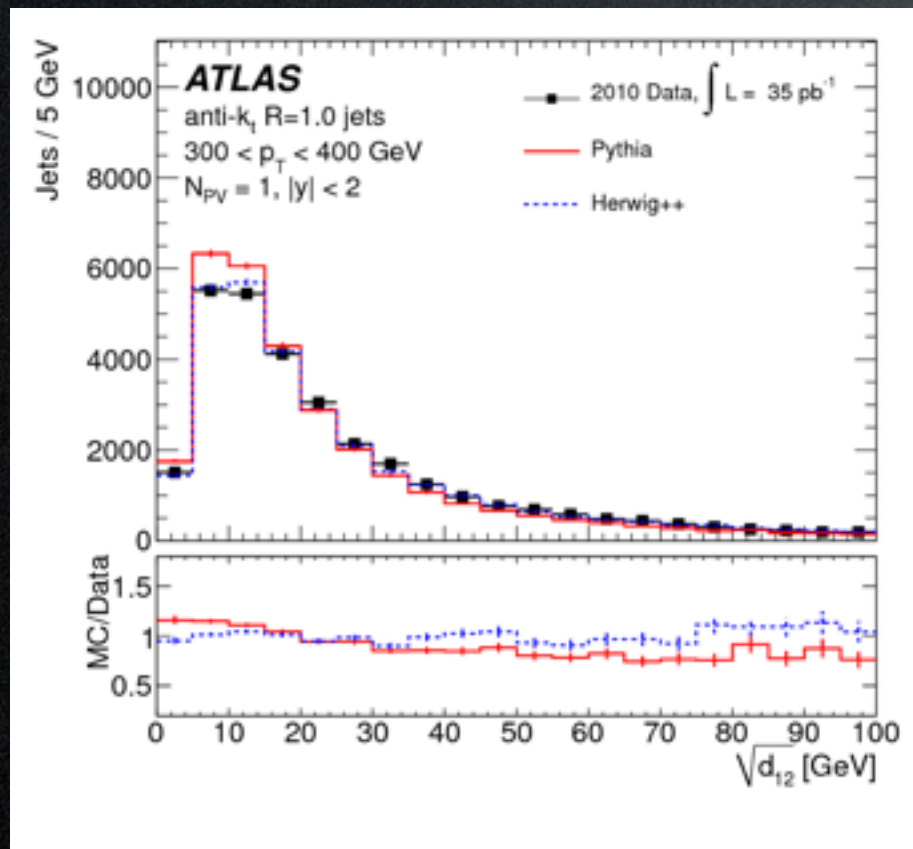
D-observables are further optimised  
by exploiting boost-invariance  
of the difference of one and two prong



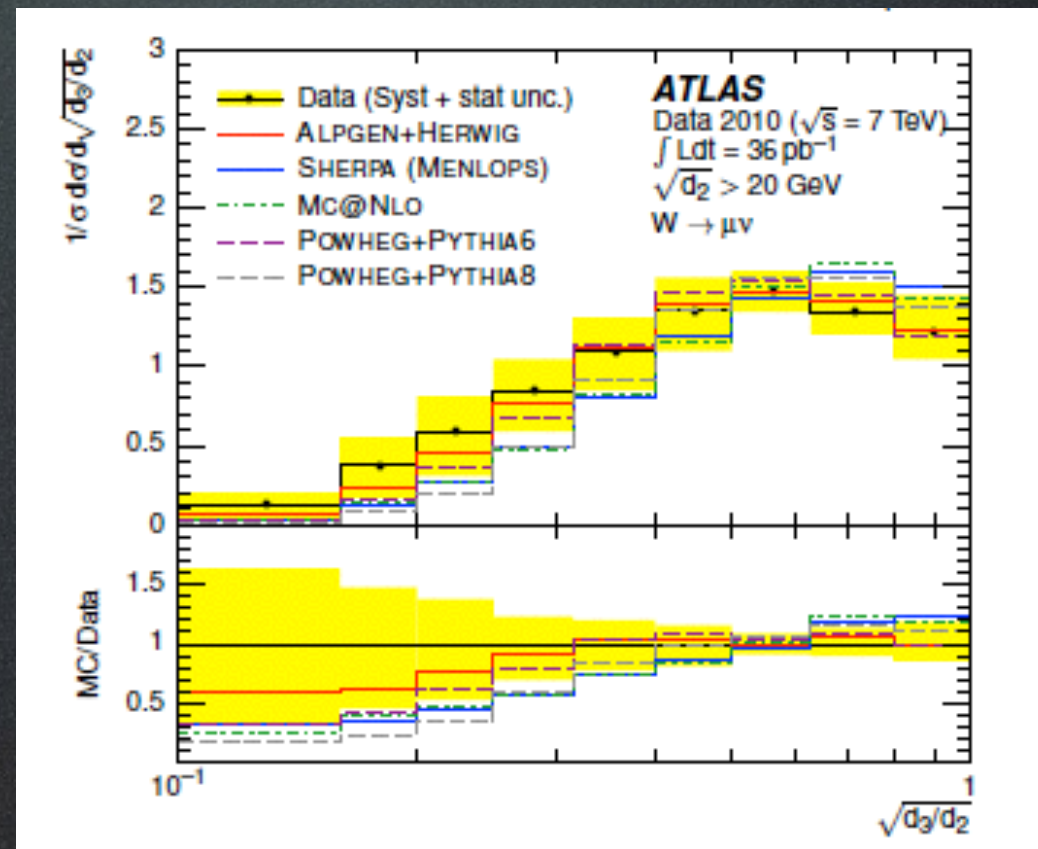
# $k_t$ Splitting Scale

$$\sqrt{d_{ij}} = \min(p_{Ti}, p_{Tj}) \times \Delta R_{ij}$$

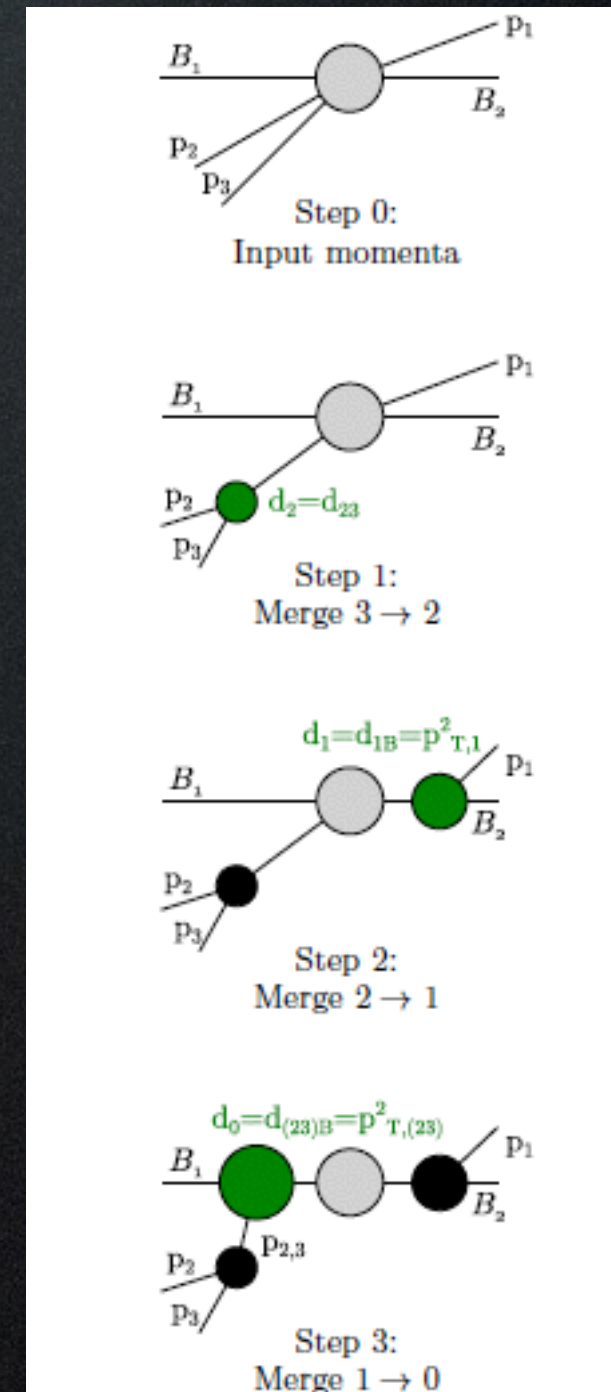
When combining  
two subjets with  $k_t$   
algorithm:



[arXiv:1203.4606](https://arxiv.org/abs/1203.4606)



[arXiv:1302.1415](https://arxiv.org/abs/1302.1415)



[arXiv:1302.1415](https://arxiv.org/abs/1302.1415)

Symmetric for heavy particle two body decay



# N-Subjettiness

Quantify the degree to which jet radiation is aligned along specific subjet axes.

$$\tau_N \equiv \frac{1}{d_0} \sum_{k=1}^M \left( p_{T,k} \times \underbrace{\Delta R_{\min,k}}_{\text{distance to nearest subjet}} \right)$$

Smaller values: N or less energy deposits

Larger values: more than N energy deposits

$$d_0 = R \times \text{sum of } p_T \text{ of all constituents}$$

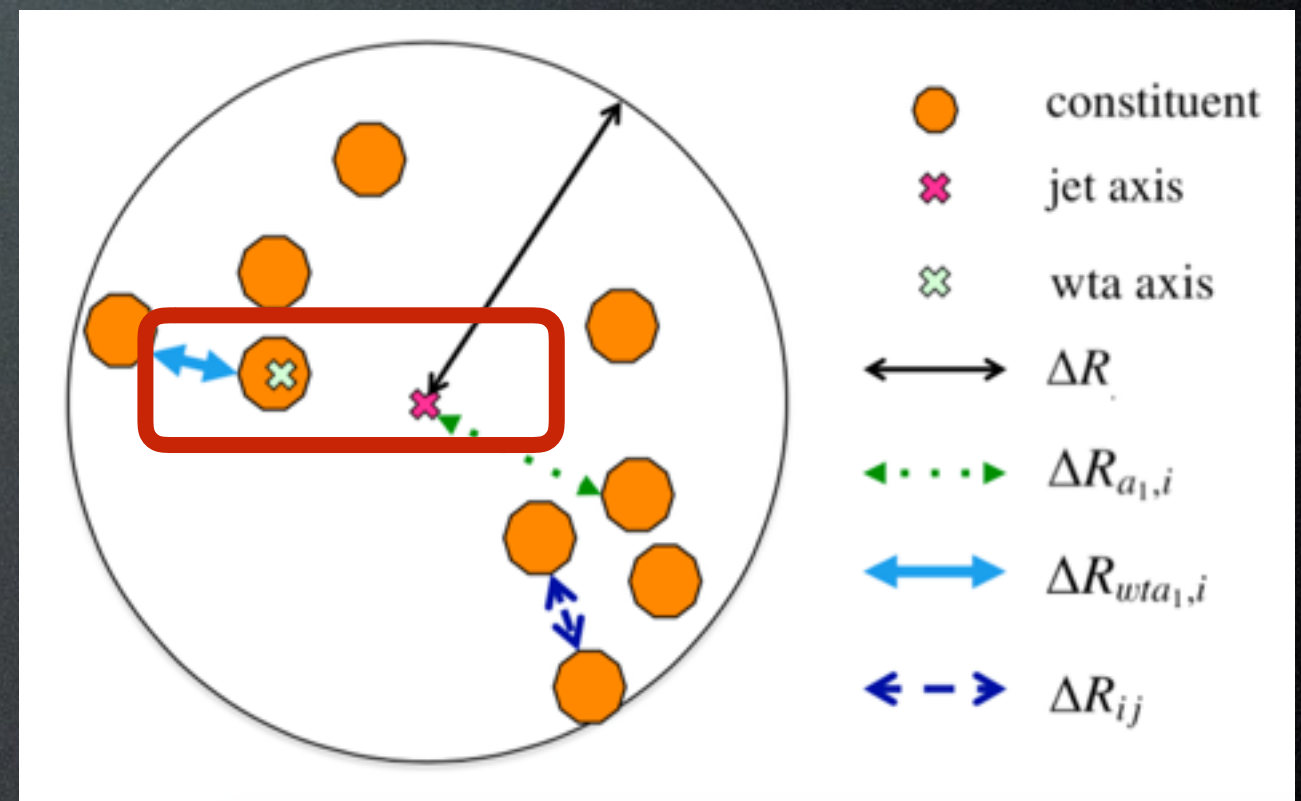
$$\tau_{N-1} > \tau_N \text{ for } N \text{ prong substructure}$$

Calculated by  $k_t$  clustering the constituents, and requiring exactly N subjets



# Winner Takes All

- Choice of subjet directions introduces an inherent ambiguity
- Sum the constituents, but  $p_T$  given by hardest
- Less number of subjets
- Minimize over all possible candidate subjet directions





# N-Subjettiness

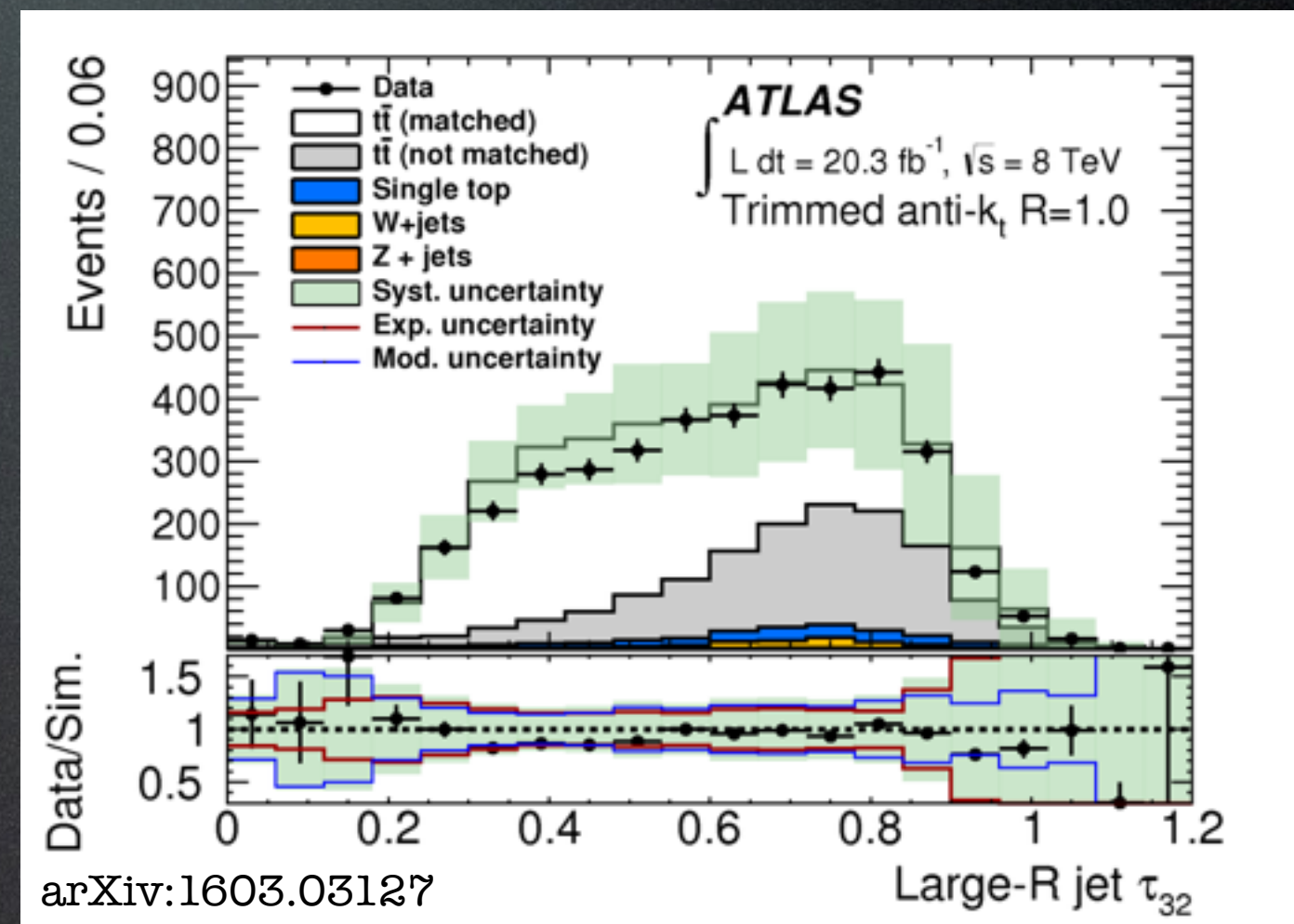
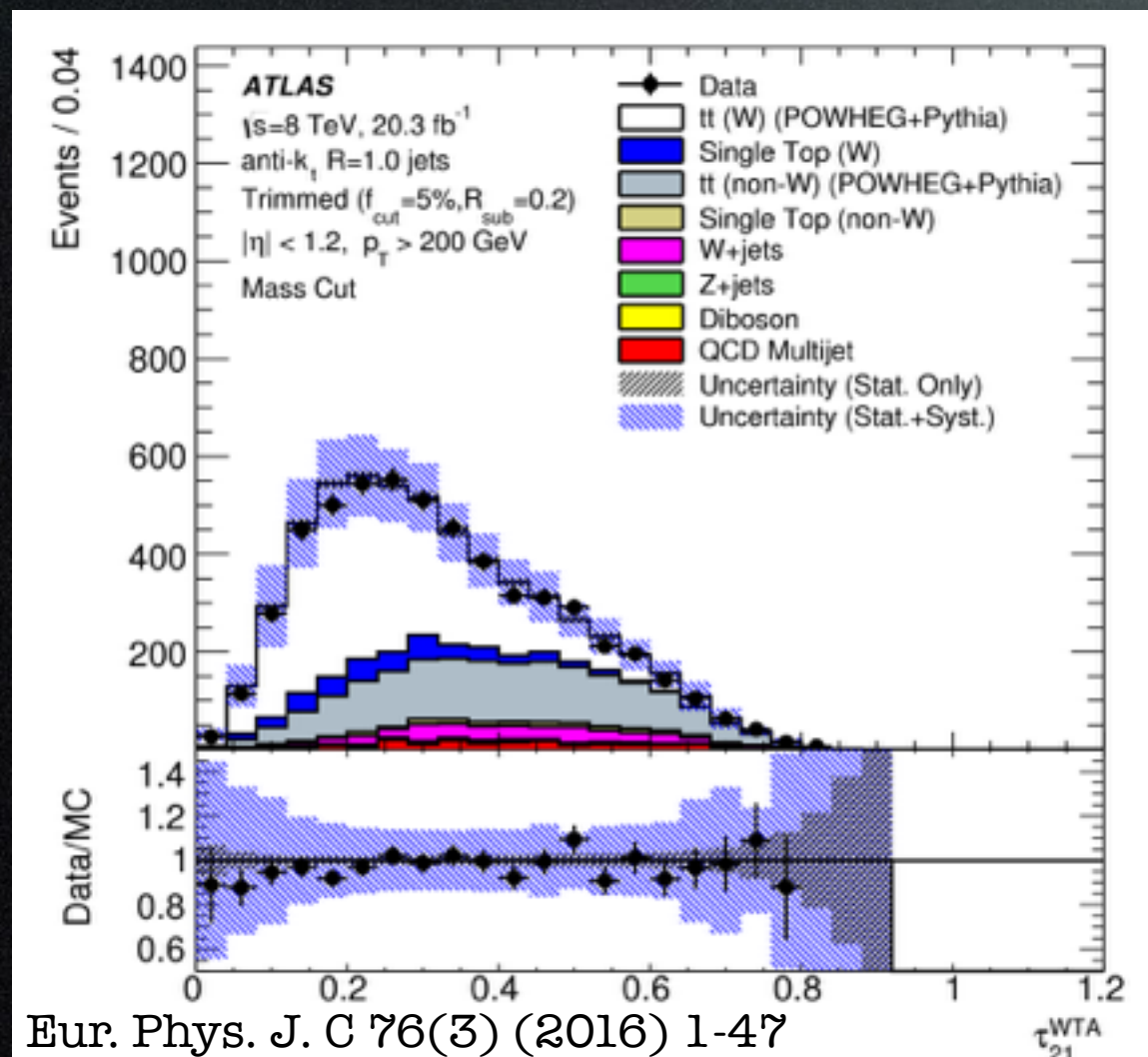
The ratio of  $\tau_N/\tau_{N-1}$  is used as discriminant



More like 2 subjects than 1



More like 3 subjects than 2



W-like



MJ-like

59

Top-like



MJ-like



# N-Subjettiness

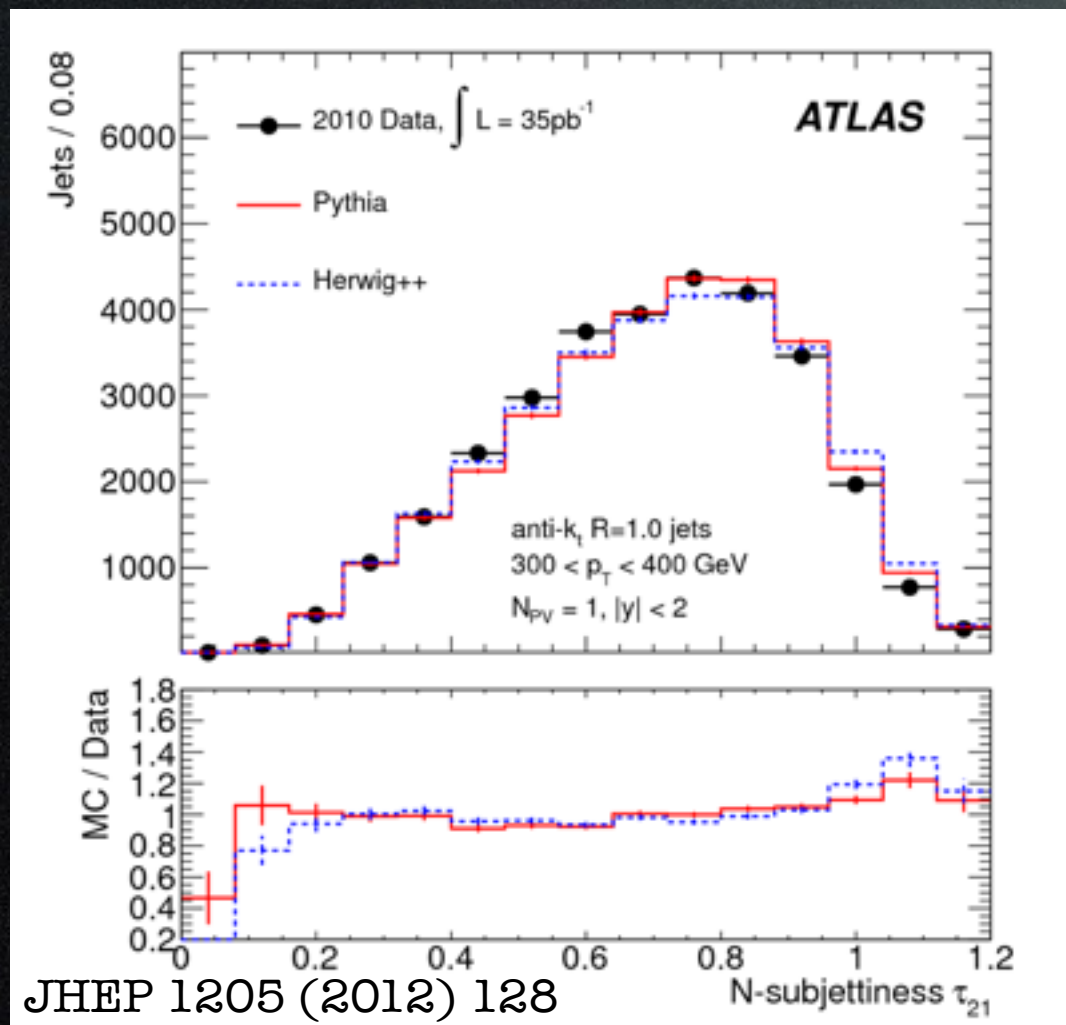
The ratio  $\tau_N/\tau_{N-1}$  is used as discriminant



More like 2 subjects than 1



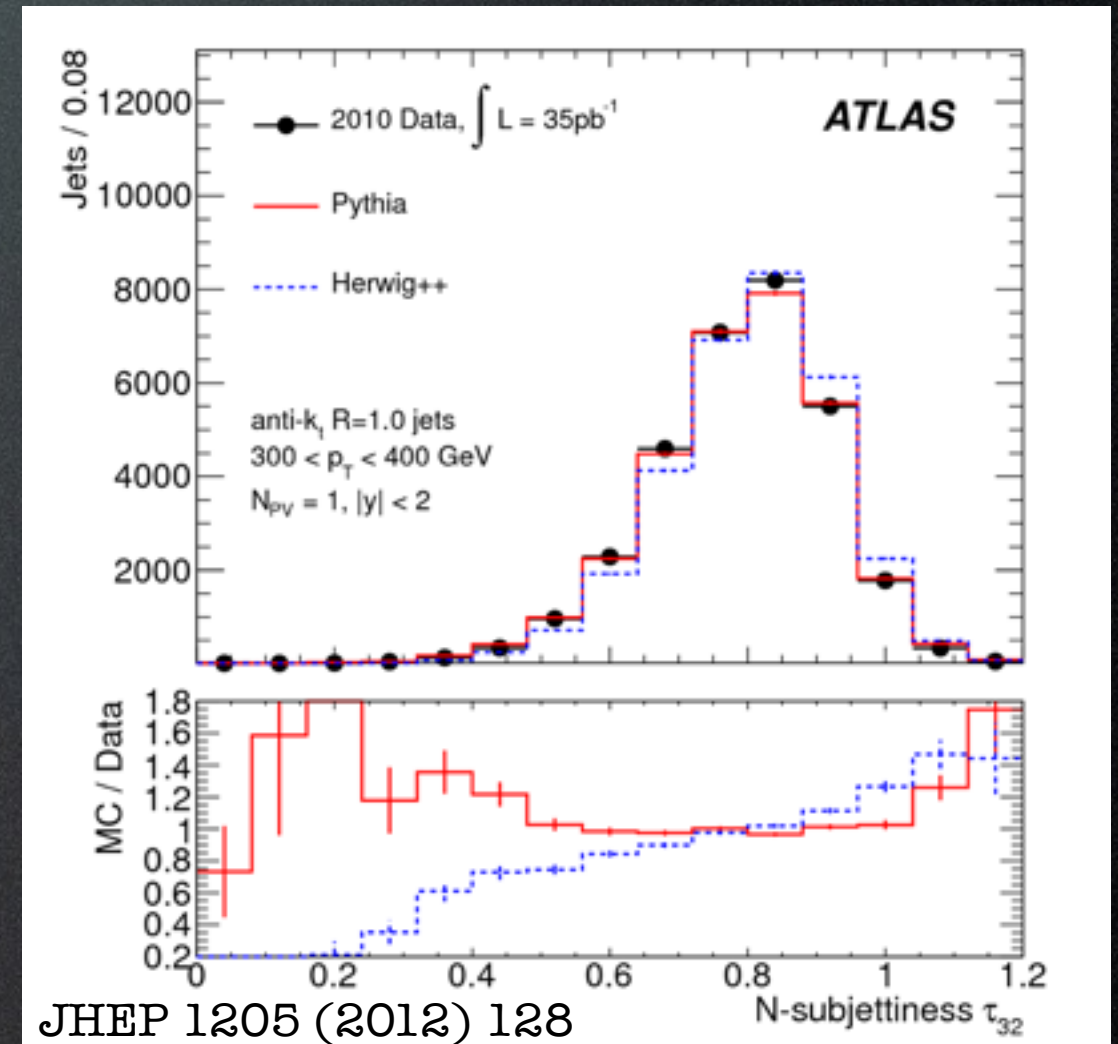
More like 3 subjects than 2



W-like



MJ-like



Top-like



MJ-like



# HEPTopTagger

Browsing through all the branches of jet recombination history

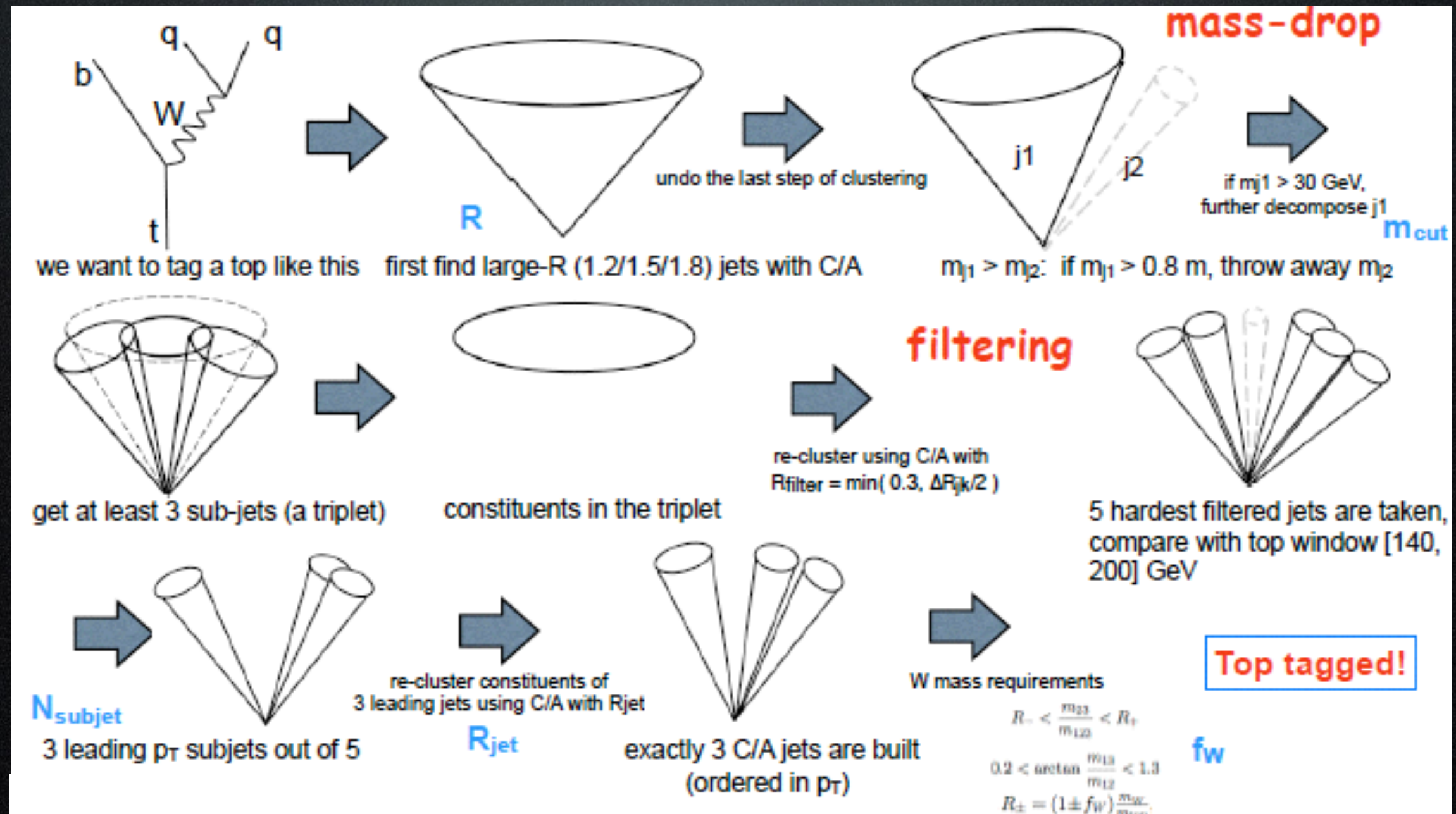


Figure by Xiaoxiao Wang



Ghost association provides a more correct matching of tracks to calorimeter jets. In this technique, tracks are treated as infinitesimally soft, low- $p_T$  particles by scaling their  $p_T$  by a very small number, such as 10–100. These tracks are then added to the list of inputs to the jet algorithm. The low scale means the tracks do not affect the reconstruction of calorimeter jets. However, after the jet algorithm, it is possible to identify which tracks were clustered into which jets. This approach properly accounts for jets with irregular cross-sectional shapes, which would lead to incorrect association in the case of simple  $\Delta R$  matching.