

Active Brownian Motion in 2d

Urna Basu
Raman Research Institute

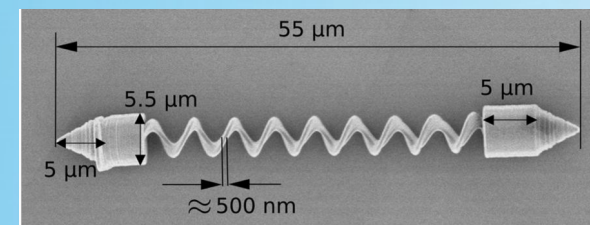
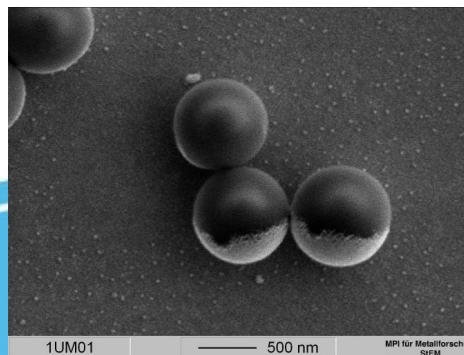
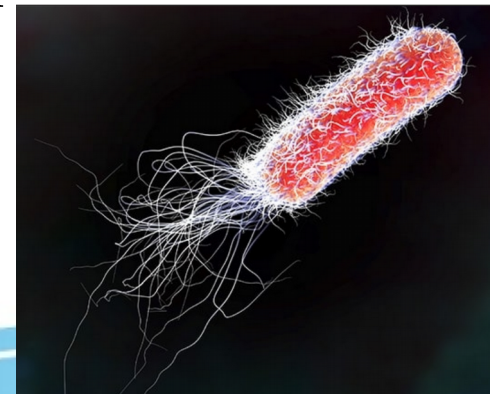


Joint work with Satya N. Majumdar, Alberto Rosso, Gregory Schehr
LPTMS, CNRS, Univ. Paris Sud, France

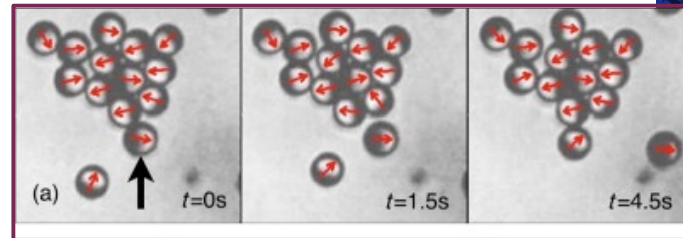
ISPCM, ICTS Bengaluru
February 2019

Active Matter

- Collection of self-propelled ‘active’ agents
- Generate directed motion consuming energy from environment
- Intrinsically driven *out of equilibrium*
- Relevant in a wide range of disciplines
- Many examples : nature and artificial



- Novel collective behaviour
 - Swarming/flocking
 - Motility induced phase separation
 - Absence of well defined pressure



- Even at the single particle level: Rich behaviour
 - Non-Boltzmann stationary state
 - Anomalous dynamics

Study a simple model ...

Active Brownian Particle (ABP)

- Overdamped particle moving with a constant speed v_0
- Two dimensions: position (x,y) and orientation ϕ
- Orientation undergoes rotational diffusion

Langevin equations

$$\dot{x} = v_0 \cos \phi(t)$$

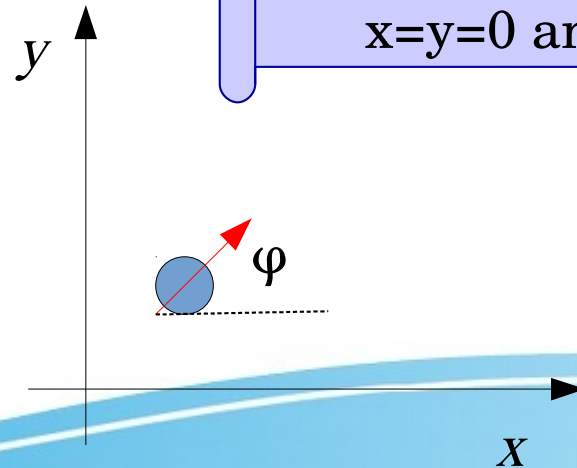
$$\dot{y} = v_0 \sin \phi(t)$$

$$\dot{\phi} = \sqrt{2D_R} \eta_\phi(t)$$

Rotational diffusion
constant

White noise
(delta-correlated)

Starting at
origin, oriented along x:
 $x=y=0$ and $\phi=0$



[Bechinger et. al. Rev Mod Phys 2016]

[Cates & Tailleur, EPL 2013]

and many more...

- Recast ABP dynamics

$$\dot{x} = \xi_x(t) = v_0 \cos \phi(t)$$

$$\dot{y} = \xi_y(t) = v_0 \sin \phi(t)$$

$$\dot{\phi} = \sqrt{2D_R} \eta_\phi(t)$$

- Effective noise is different than white-noise

Bounded $|\xi| \leq v_0$

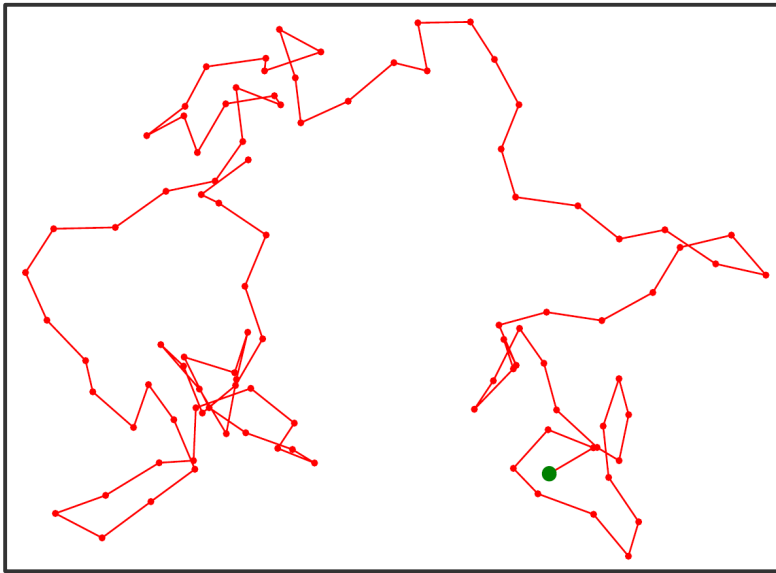
Correlated $\langle \xi_x(t) \xi_x(t') \rangle \simeq v_0^2 e^{-D_R |t-t'|}$

Similar form for the y-y and x-y noise correlations



Time-scale: $1/D_R$

Trajectories

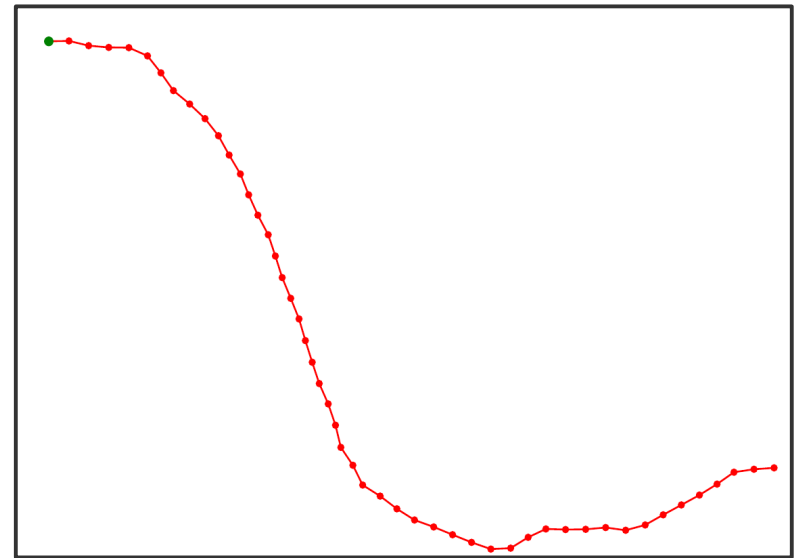
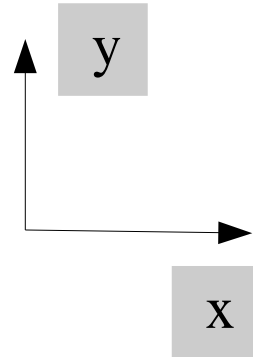


Long time ($t \gg 1/D_R$)

Diffusive, similar to
passive Brownian motion

Effective diffusion constant

$$D_{\text{eff}} = v_0^2 / 2D_R$$



Short time ($t \ll 1/D_R$)

Very different:
signature of activity

● Explore and characterize ...

Short-time Regime ($t \ll 1/D_R$)

U.B., S. N. Majumdar, A. Rosso, G. Schehr, Phys. Rev. E 2018

Short-time: Anisotropy

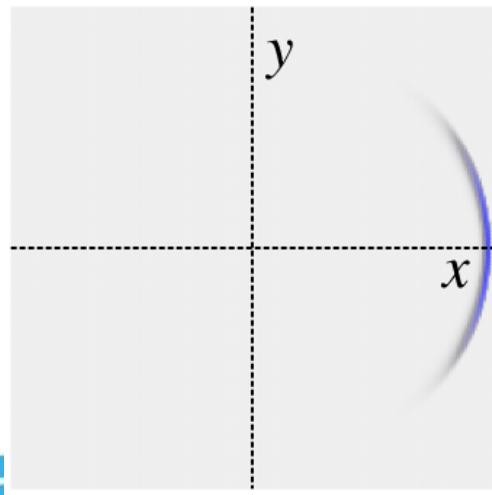
- Effective equation for short-time $t \ll 1/D_R$

$$\begin{aligned}\dot{x} &= v_0 \cos \phi(t) \\ \dot{y} &= v_0 \sin \phi(t) \\ \dot{\phi} &= \sqrt{2D_R} \eta_\phi(t)\end{aligned}$$

Small t ,
small ϕ

$$\begin{aligned}\dot{x} &\approx v_0 \left(1 - \frac{\phi^2}{2}\right) \\ \dot{y} &\approx v_0 \phi\end{aligned}$$

Position
distribution
 $P(x,y;t)$



$D_R=0.1$
 $t=1$

Analyze marginal distributions
 $P(x,t)$ and $P(y,t)$ separately

- Use Path Integral techniques ...

Short-time: $P(x,t)$

$$D_R = 0.01$$

$$v_0 = 1$$

Exact probability distribution

$$P(x, t) = \frac{1}{v_0 D_R t^2} f\left(\frac{v_0 t - x}{v_0 D_R t^2}\right)$$

$$f(u) = \frac{1}{2\sqrt{\pi}u^3} \sum_{m=0}^{\infty} (-1)^m \frac{(4m+1)}{2^{2m}} e^{-\frac{(4m+1)^2}{8u}}$$

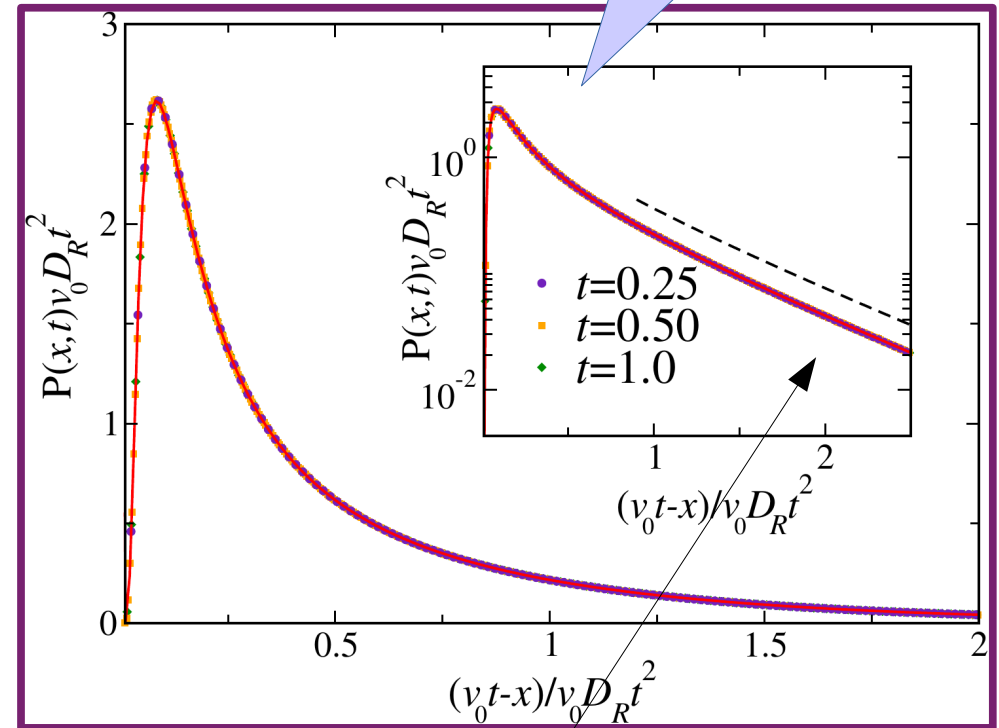
Asymptotes

$$\lim_{u \rightarrow 0} f(u) \simeq \frac{1}{2\sqrt{\pi}u^3} e^{-\frac{1}{8u}}$$

$$\lim_{u \rightarrow \infty} f(u) \simeq \frac{1}{\sqrt{2u}} e^{-\frac{\pi^2 u}{8}}$$

Variance:

$$\sigma_x^2 \approx \frac{1}{3} v_0^2 D_R^2 t^4$$

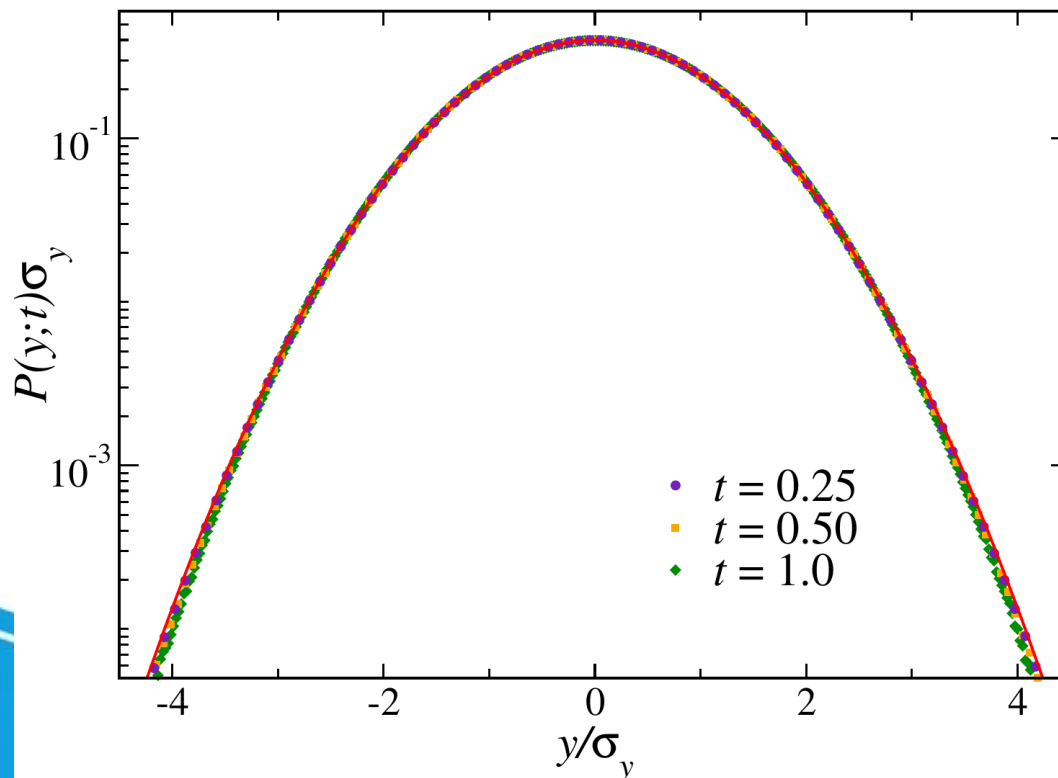


Short-time: $P(y,t)$

$$\dot{y} \approx v_0 \phi$$

- Mapping to Random Acceleration Process
- Gaussian distribution

$$P(y, t) = \sqrt{\frac{3}{4\pi v_0^2 D_R t^3}} \exp \left[-\frac{3y^2}{4v_0^2 D_R t^3} \right]$$



Variance:

$$\sigma_y^2 \approx \frac{2}{3} v_0^2 D_R t^3$$

$$D_R = 0.01$$

$$v_0 = 1$$

Short-time: Radial distribution

- Radial distance r :

$$P(r^2, t) = \frac{1}{2D_R v_0^2 t^3} f_r \left(\frac{v_0^2 t^2 - r^2}{2D_R v_0^2 t^3} \right)$$

$$f_r(z) = \frac{1}{4\pi z^{3/2}} \sum_{k=0}^{\infty} \frac{(2k)!}{(k!)^2 2^{2k}} \sqrt{4k+1} e^{-\frac{(4k+1)^2}{16z}} \left(\left(\frac{(4k+1)^2}{4z} - 1 \right) K_{1/4} \left(\frac{(4k+1)^2}{16z} \right) + \frac{(4k+1)^2}{4z} K_{3/4} \left(\frac{(4k+1)^2}{16z} \right) \right)$$

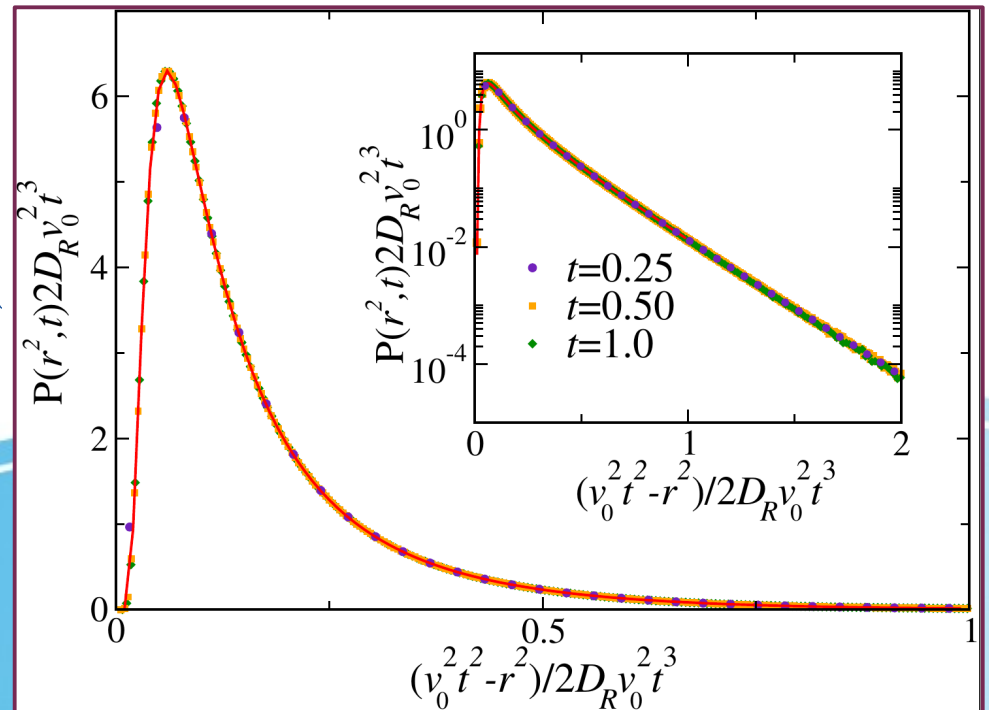
$$D_R = 0.01$$

$$v_0 = 1$$

Asymptotes

$$f_r(z) \simeq \frac{1}{2\sqrt{2\pi} z^2} e^{-\frac{1}{8z}}, \quad \text{as } z \rightarrow 0$$

$$f_r(z) \simeq \sqrt{\frac{\pi}{z}} e^{-\frac{\pi^2}{2} z} \quad \text{as } z \rightarrow \infty$$



Summary so far...

- Strong signatures of activity at short-times
 - Anisotropy
 - Non-trivial position distribution
- Crossover to passive Brownian behaviour at times larger than $1/D_R$

How to see signatures of activity
in the long-time regime?

Presence of external potential

- ABP diffusing in an external potential
 - Numerical studies
 - Experiments
- Strong signatures of activity at long-times: Non-Boltzmann stationary state
- Crossover to Boltzmann-like state

[Takatori et. al., Nat. Comm. 2015]

[Pototsky & Stark, EPL 2012]

[Solon, Cates & Tailleur, EPJ Spl. Topics 2015]

Qualitative understanding, few exact results

[Malakar et. al., arXiv: 1902.04171]

ABP in a Harmonic trap

- 2d harmonic potential
- Langevin equation

$$U(x, y) = \frac{\mu}{2}(x^2 + y^2)$$

$$\begin{aligned}\dot{x} &= -\mu x + v_0 \cos \phi(t) \\ \dot{y} &= -\mu y + v_0 \sin \phi(t) \\ \dot{\phi} &= \sqrt{2D_R} \eta_\phi(t).\end{aligned}$$

- Two competing time-scales: trap ($1/\mu$) and rotational diffusion ($1/D_R$)



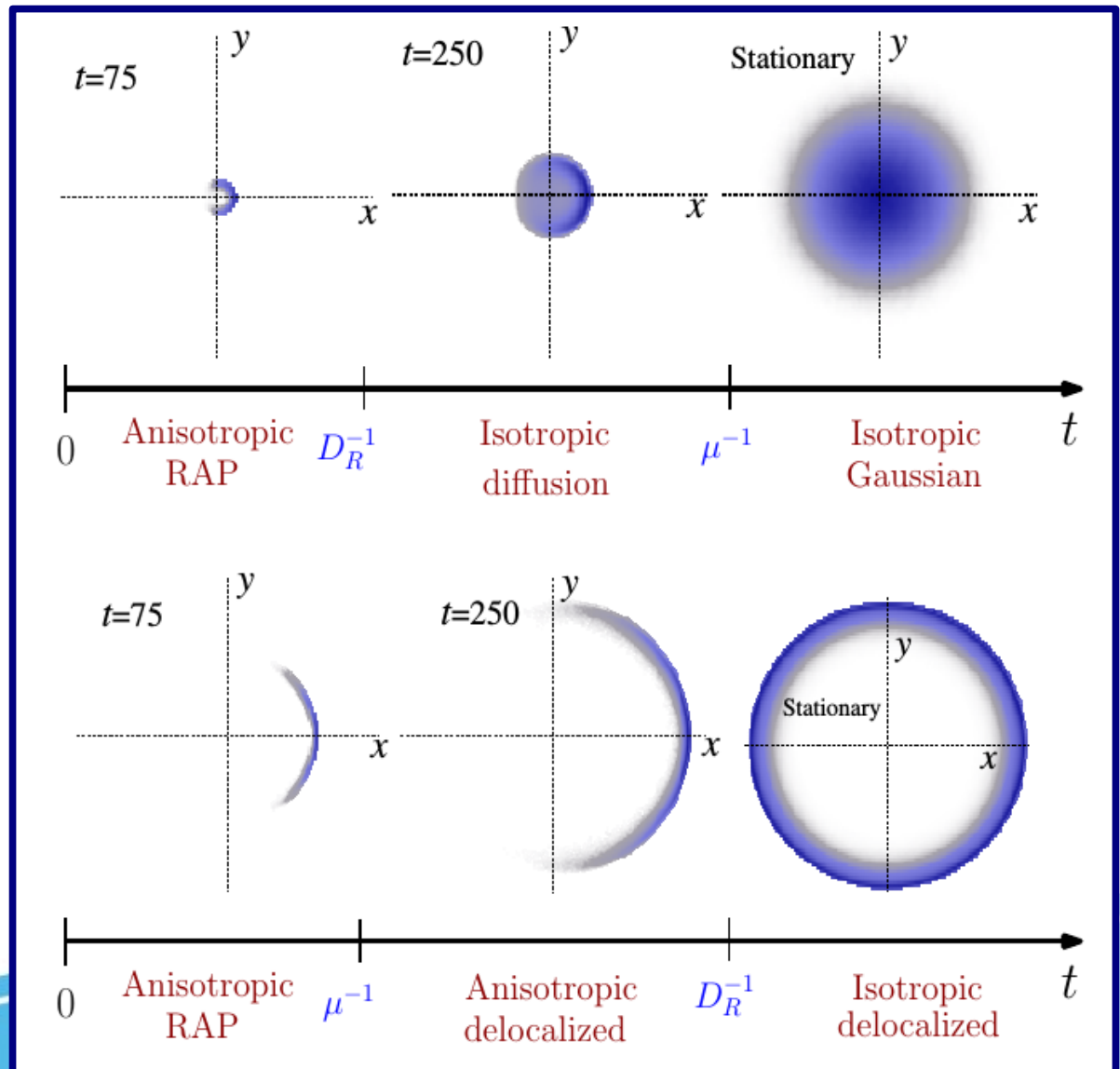
Passive Brownian particle $\rightarrow$ Ornstein-Uhlenbeck process

Dynamical Evolution

$$\mu^{-1} > D_R^{-1}$$

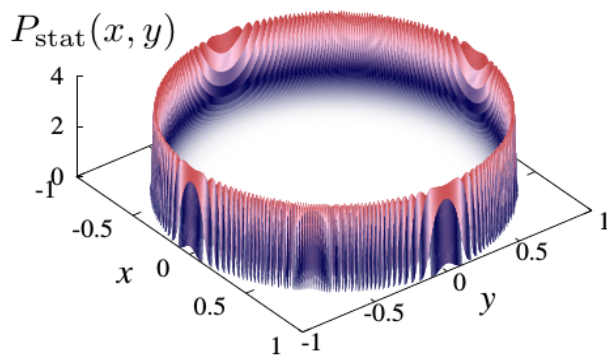
Two regimes

$$\mu^{-1} < D_R^{-1}$$

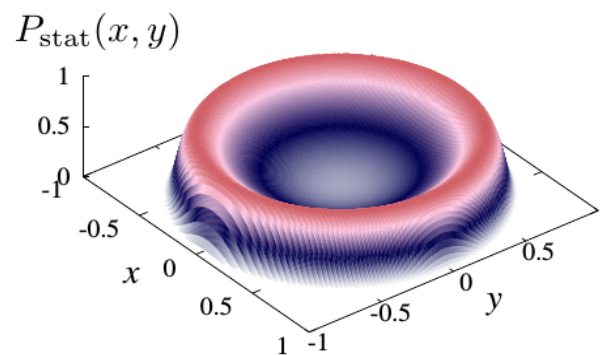


Stationary state

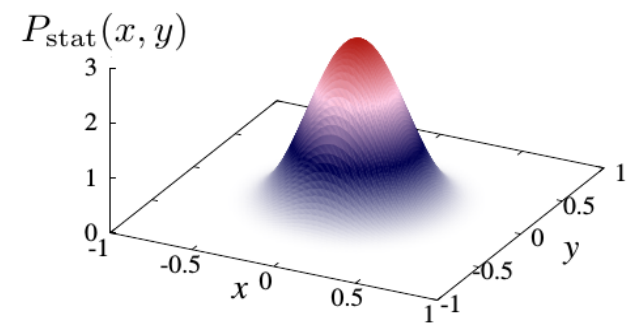
- Crossover from single peak to delocalized ring away from trap center



$D_R = 0.1$



$D_R = 1$



$D_R = 10$

$\mu=1, v_0=1$

Analytical understanding ?

Position distribution

- Moment evolution equation

$$\dot{M}_{k,l} = -D_R(k-l)^2 M_{k,l} + v_0 e^{-\mu t} [k M_{k-1,l} + l M_{k,l-1}]$$

- Difficult to solve in general, even in the stationary state
- Limiting cases:
 - Strongly passive : $D_R \rightarrow \infty$
 - Strongly active : $D_R \rightarrow 0$

Boltzmann-like

Non-Boltzmann

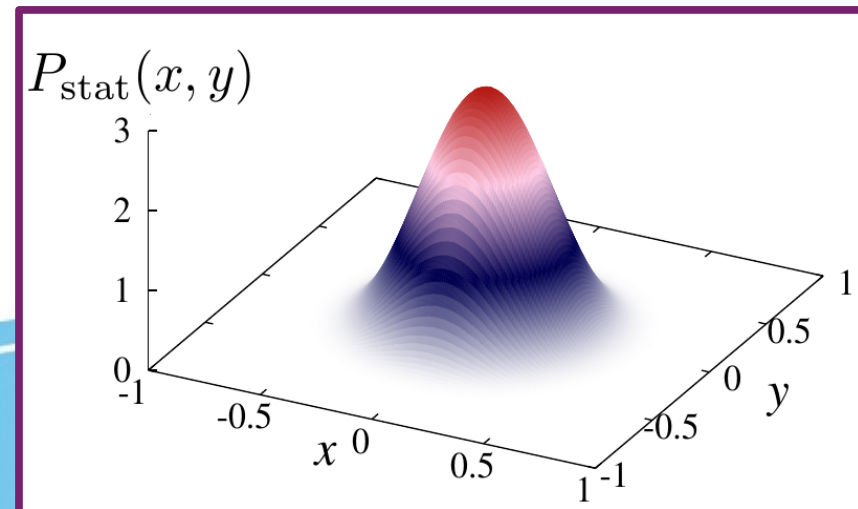
Passive limit ($D_R \rightarrow \infty$)

- Radial distribution: Gaussian *at all times*

$$P_{\text{rad}}(r, t) \simeq \frac{2 \mu D_R}{v_0^2 (1 - e^{-2\mu t})} \exp \left[-\frac{\mu D_R r^2}{v_0^2 (1 - e^{-2\mu t})} \right]$$

- Stationary state given by large t limit
- Rotationally symmetric

$$P_{\text{stat}}(x, y) = \frac{\mu D_R}{\pi v_0^2} \exp \left[-\frac{\mu D_R (x^2 + y^2)}{v_0^2} \right]$$



Strongly active limit ($D_R \rightarrow 0$)

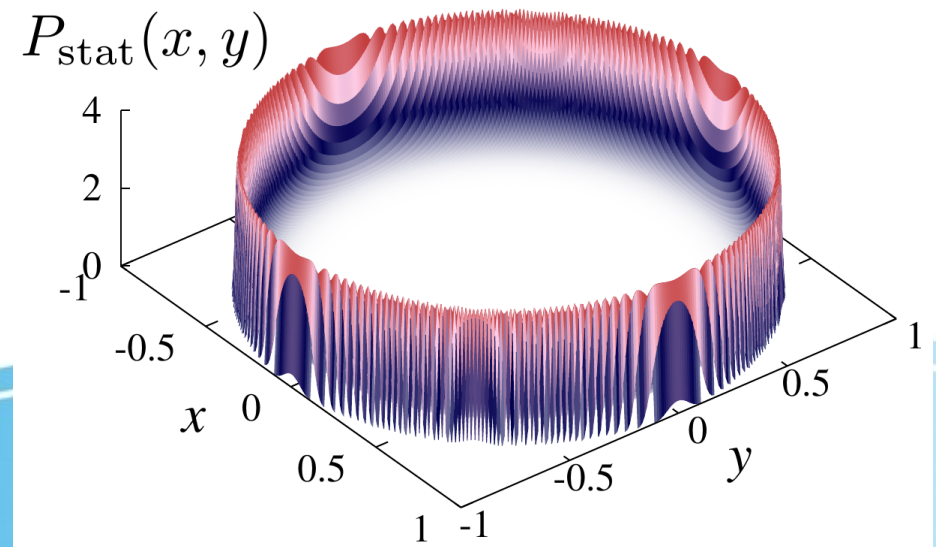
- Time-dependent radial distribution

$$P_{\text{rad}}(r, t) = \frac{\mu}{v_0(1 - e^{-\mu t})} \delta \left[r - \frac{v_0(1 - e^{-\mu t})}{\mu} \right]$$

The particle is likely to be away from the center of the trap, the distance growing with time until ‘boundary’ is reached

- Stationary distribution: rotationally symmetric for $D_R \rightarrow 0+$

$$P_{\text{stat}}(x, y) = \frac{\mu}{2\pi v_0} \delta \left[\sqrt{x^2 + y^2} - \frac{v_0}{\mu} \right]$$



Conclusions

- Active Brownian motion in 2d
 - Diffusive behaviour at long-times
 - Non-trivial anisotropic short-time properties
- Exact distribution at short-times
- ABP in presence of a trap
 - Exact results in two limiting cases

Thank you!