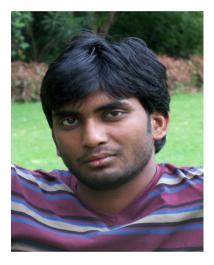
Aggregation Processes in Systems of Fractal Objects

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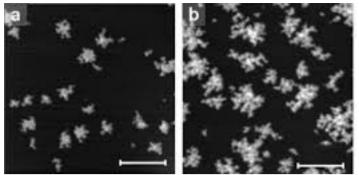
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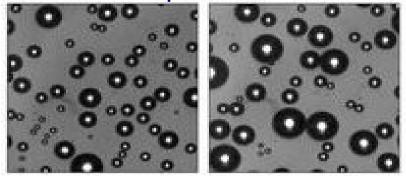
Plan: Introduction

Ballistic Aggregation in Hard Sphere model in Vapor-Solid Phase Transition Aggregation in a model Active Matter system Conclusion. lce



Heidom et al. (CPL)

Liquid water



t = 20 min t = 45 min

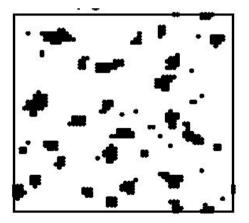
Mouterde et al. (Nature Materials)

Possible ways of growth in disconnected morphology:

Droplets consist of small particles.

Particles from smaller droplets get detached to be deposited on a bigger droplet.

Droplets <u>themselves move</u> and undergo ``sticky Collisions" with each other to form bigger droplets.



Motion of droplets/clusters can be diffusive, ballistic,

Ballistic aggregation: Materials Science, Granular Matter, Astrophysics,

Carnevale et al., Phys. Rev. Lett. (1990). Hansen and Trizac, J. Stat. Phys. (1996) Trizac and Krapivsky, Phys. Rev. Lett. (2003). J. Blum, Astronomy and Astrophyics (2010).

2 types:

Ballistic deposition – seed is fixed (vapor deposition on a substrate, etc.).

Ballistic aggregation: all clusters move (translation and rotation) and collide.

Source: ERCIM Source: ERCIM Market States Substates Substates

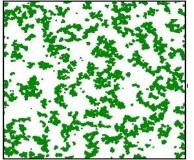
Heidom et al.

Fractality is usually a natural outcome.

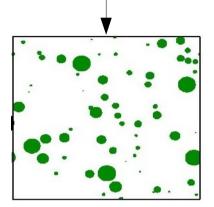
Typically simulations are done with ``sticky'' hard spheres

- fractal structures are expected.

In event driven simulations, it's necessary to keep track of exact locations as well as orientations of fractal clusters and identify exact points of contact during collision – technically challenging – so spherical cluster approximation becomes necessary.



S. Paul and SKD, Phys. Rev. E (2017)



Theory: M. Smoluchowski (1916)
G.F. Carnevale, Y. Pomeau and W.R. Young, PRL (1990)
Hansen and Trizac, J. Stat. Phys. (1996)

$$\frac{dn}{dt} = -Collision - cross - section \times v_{rms} \times n^2$$

Collision - cross - section $\sim M^{\frac{d-1}{d}} n \propto 1/M \qquad v_{rms} \sim M^{-1/2}$
 $\frac{dM}{dt} \sim M^{\frac{d-2}{2d}} \qquad M \sim t^{\beta}; \ \beta = 2d/(d+2)$

In event driven simulations ... spherical cluster approximation is necessary.

Hansen and Trizac, J. Stat. Phys. (1996). Trizac and Krapivsky, Phys. Rev. Lett. (2003). S.N. Pathak, Z. Jabeen, D. Das and R. Rajesh, Phys. Rev. Lett. (2014). S. Paul and SKD, Phys. Rev. E (2017).

SO the role of fractality in growth is NOT properly known because of spherical structural approximation – for same average mass, collision cross-section should be higher for (rotating) fractal clusters.

<u>Time step driven MD (TDMD)</u> can take care of such difficulty – events occur naturally.

But TDMD applies usually to non-hard sphere kind of systems, like Lennard-Jones inter-particle/atomic potential.

How will such particles stick and keep fractality?

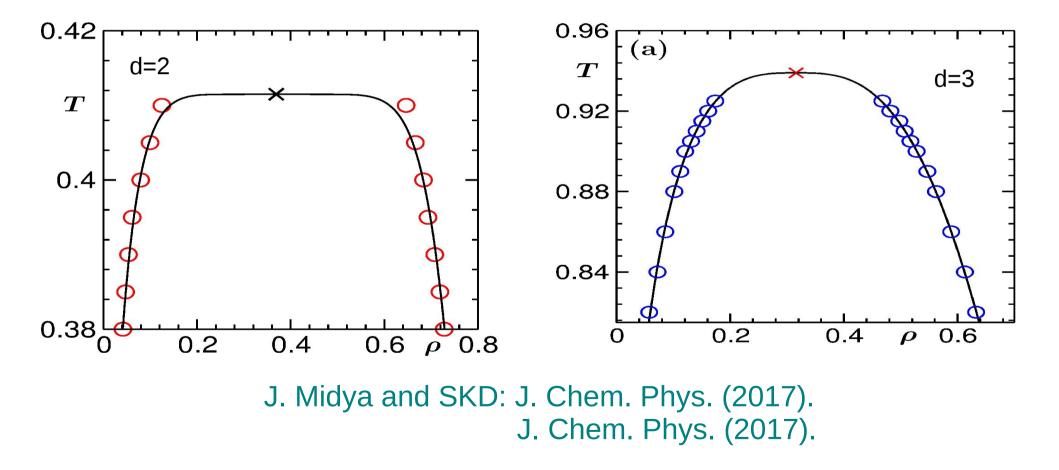
Even if clustering occurs they will try to gain spherical structure to minimize (interfacial) free energy. **??**

Typically one uses Model H, Lattice Boltzmann, Molecular Dynamics.

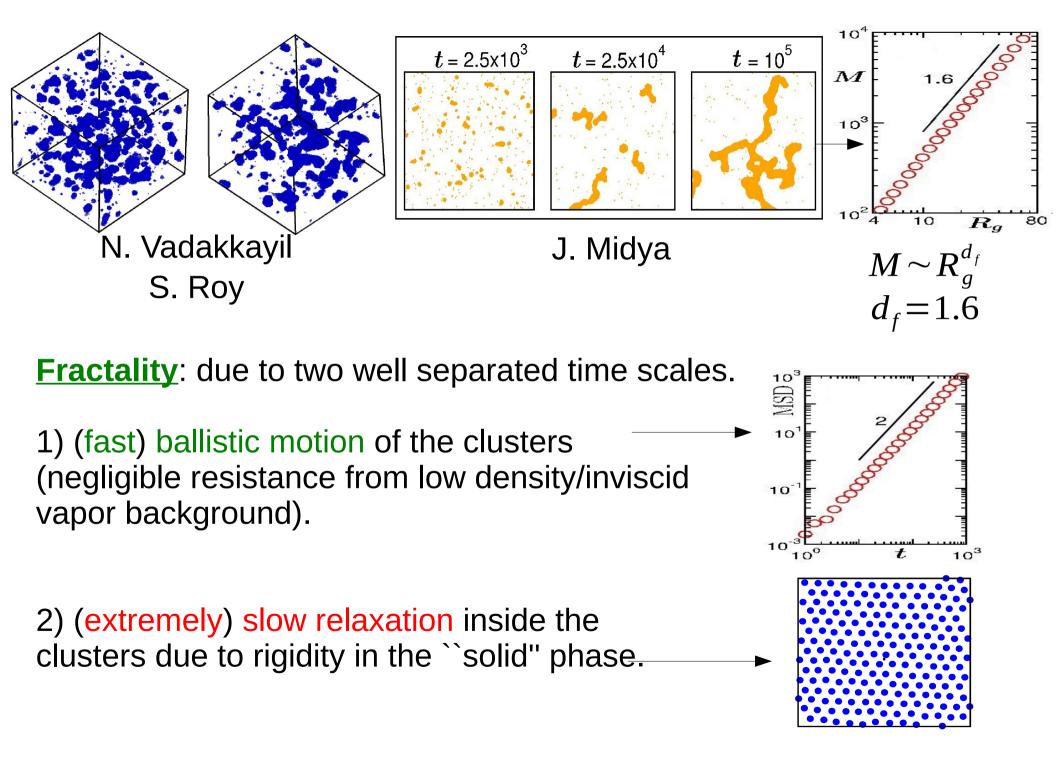
Our Model:
$$\varphi(r) = 4 \epsilon \left[\left(\frac{\sigma}{r} \right)^{12} - \left(\frac{\sigma}{r} \right)^{6} \right]$$

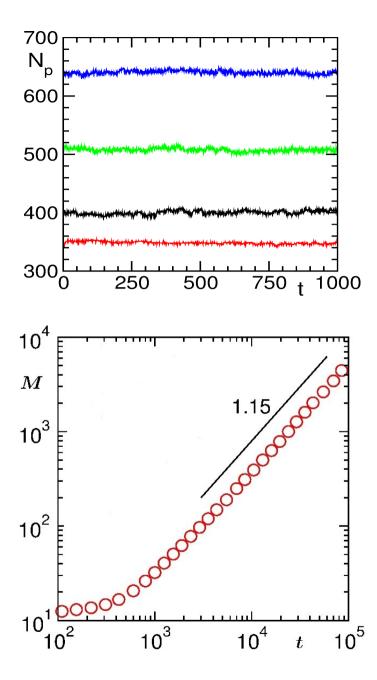
LJ potential

Phase behavior: Monte Carlo simulations.



Kinetics?: Molecular dynamics simulations with Nose-Hoover thermostat.





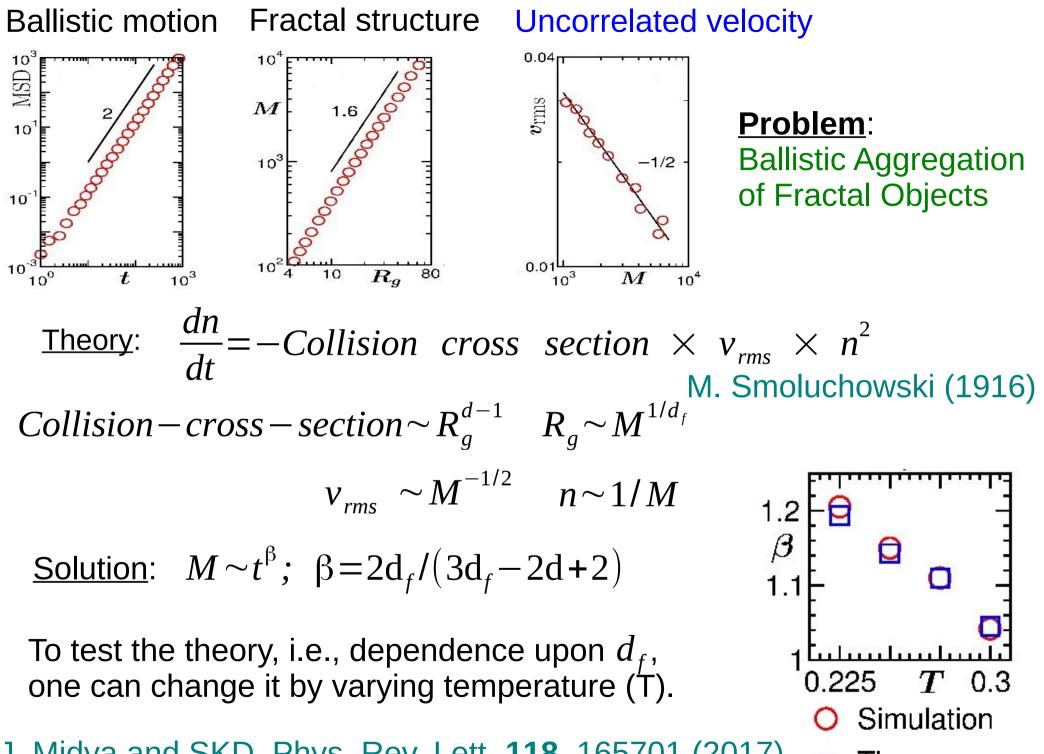
Particle diffusion mechanism (practically) does not contribute to the growth.

Problem of ballistic aggregation.

But, for ballistic aggregation: $M \sim t^{\beta}$; $\beta = 1.15$ $M \sim t^{\beta}$; $\beta = 2d/(d+2)$ d=2: $\beta=1$

Discrepancy due to fractality?

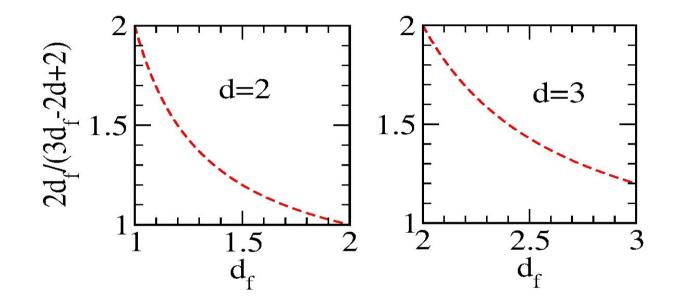
<u>YES.</u>



J. Midya and SKD, Phys. Rev. Lett. 118, 165701 (2017).
Theory

In d dimensions: collision cross section $R_g^{d-1} = M^{\frac{d}{d_f}}$

$$M \sim t^{\beta}; \quad \beta = 2d_f / (3d_f - 2d + 2)$$



2D: Ballistic aggregation exponent is same as Binder-Stauffer exponent for $d_f=2$. 3D: Possibility of exponential growth for extremely filament like structure with $d_f = 4/3$. But not a physical picture.

Phase separation in a Vicsek-like Active matter model

$$m\frac{d^{2}\vec{r}}{dt^{2}} = -\vec{\nabla}U_{i} - \gamma m\frac{d\vec{r}}{dt} + \sqrt{6\gamma k_{B}Tm}\vec{R}(t) + \vec{f}_{i} \qquad U_{i} \rightarrow \text{LJ}$$

$$< R_{i\mu}(t)R_{j\nu}(t') > = \delta_{ij}\delta\mu\nu\delta(t-t') \quad \vec{f}_{i} = f_{A}\vec{D}_{N}; \quad \vec{D}_{N} = \frac{\sum_{j}\vec{v}_{j}}{mag\sum_{j}\vec{v}_{j}}$$

$$\int_{t=5000}^{t=100} \int_{t=9000}^{t=2500} \int_{0}^{10^{4}} \int_{0}^{10} \int_{0}^{10^{4}} \int_{0}^{10^{4}$$

By tuning the temperature and active force all types of combination of fractality and cluster mean-squared-displacements can be obtained.

Role of fractality in processes related to diffusive to ballistic coalescence can be understood.

S. Paul and SKD, to be published.

Summary:

Structure and dynamics related to nucleation and growth in various systems have been discussed.

As opposed to Ostwald ripening mechanism in solid binary mixtures or Binder-Stauffer diffusive coalescence mechanism in liquid-liquid or vapor-liquid transitions, in vapor-solid transition for off-critical density, clusters/droplets grow via ballistic aggregation mechanism.

Similar features have been identified in an active matter model.

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