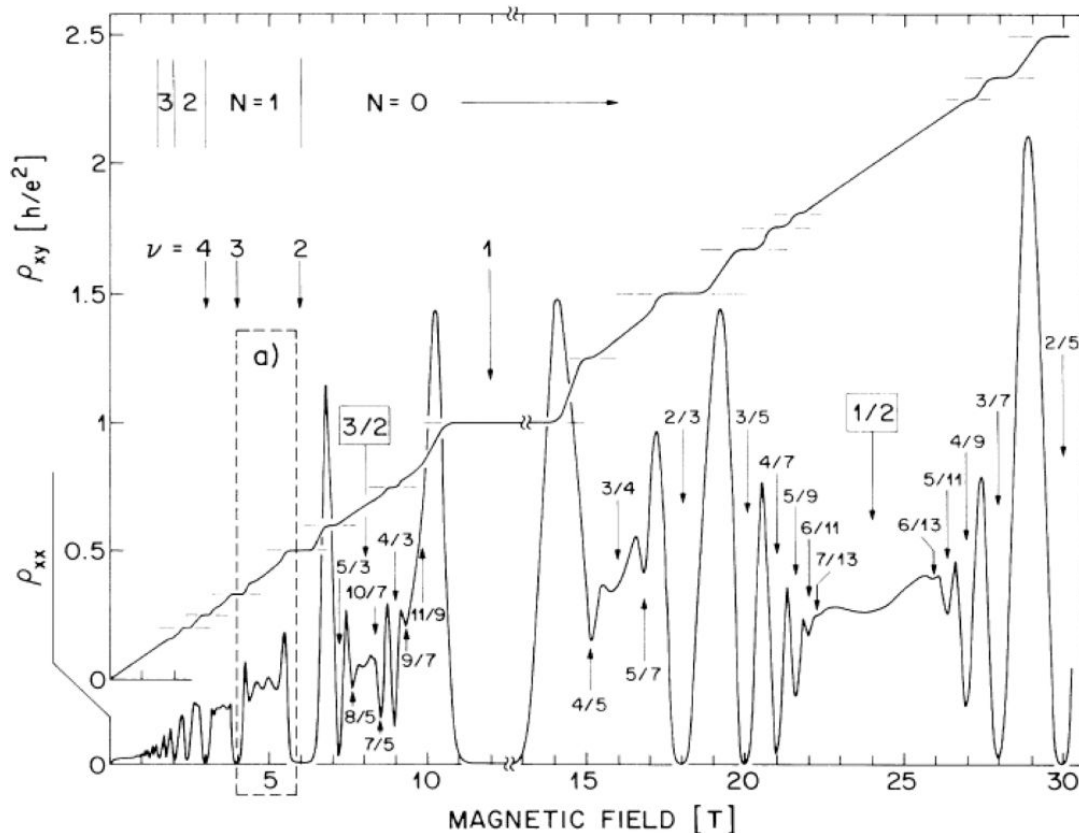


Mean field approximations to n body Hamiltonians in FQHE

Sreejith G J
IISER Pune

G J Sreejith, Y Zhang, J K Jain, Phys. Rev. B 96, 125149 (2017)
B Kusmierz, G J Sreejith (in preparation)

FQH States



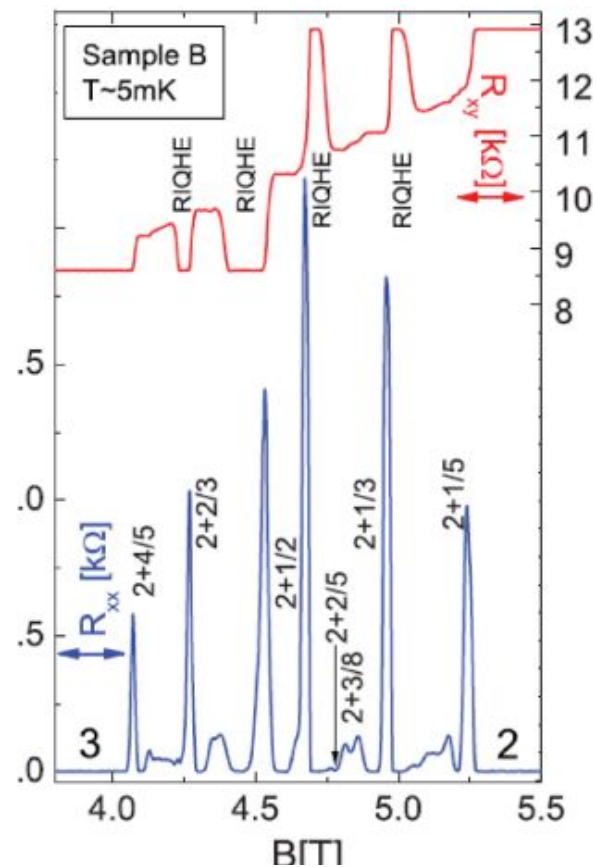
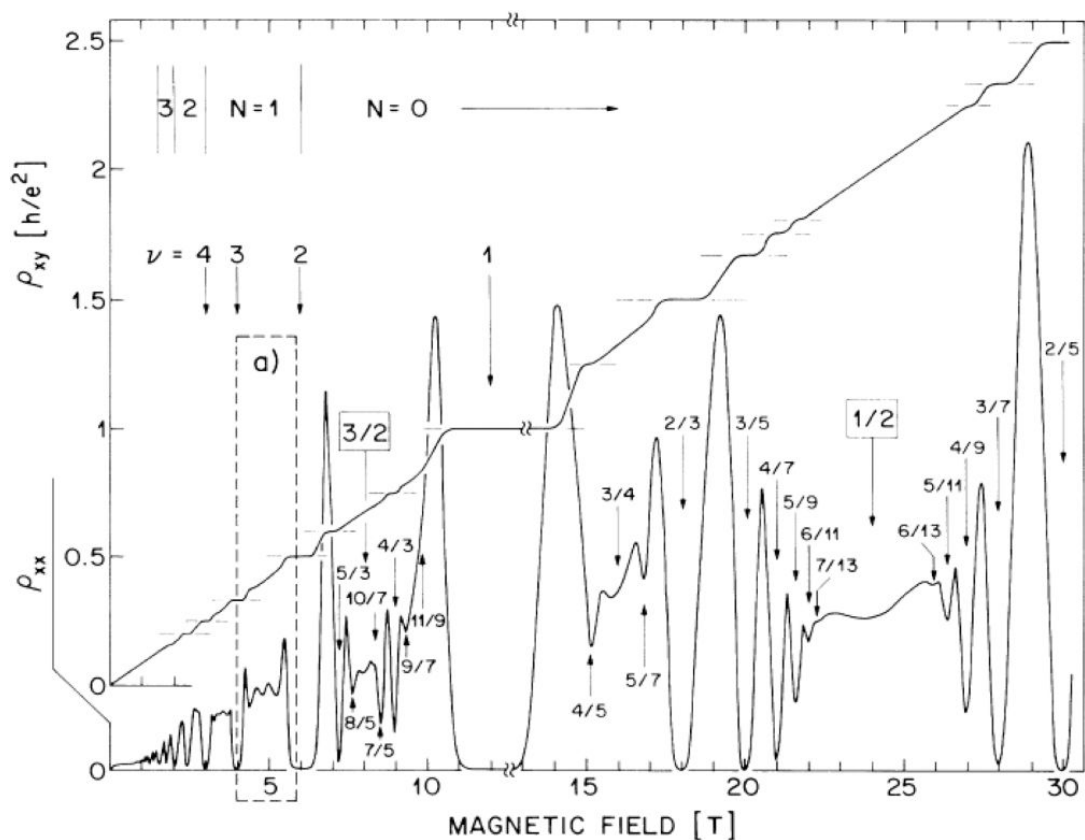
FQHE in the lowest Landau level:

Can be described using a model of single Landau level of electrons interacting through a two-body Coulomb interaction.

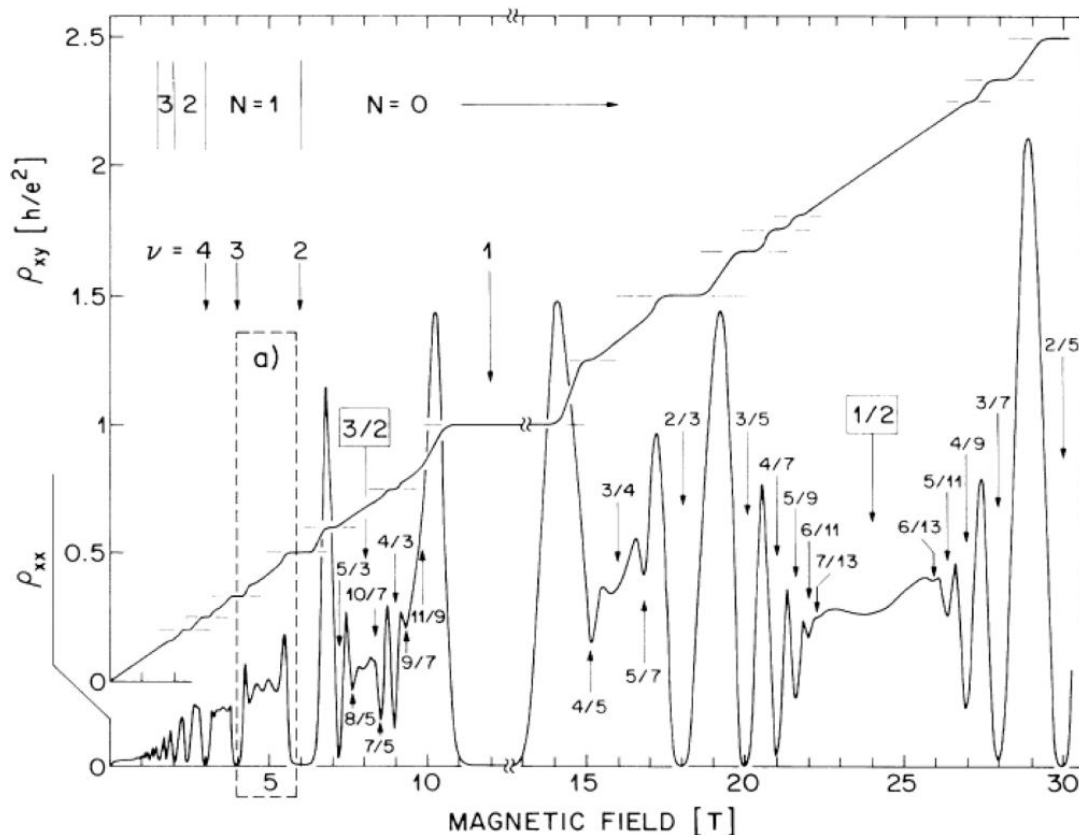
FQH occurs at odd-denominator filling fractions -- well-understood in terms of emergent particles called composite fermions

No FQHE at filling fraction $1/2$

FQH States



Even denominator FQHE



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Second Landau level:

Strong FQH state at filling fraction $1/2$.

General structure of the FQH states not very well understood


Pfaffian, Anti-pfaffian and Particle hole symmetry

$$\psi \sim \text{Pf} \left[\frac{1}{z_i - z_j} \right] \prod_{i < j=1}^N (z_i - z_j)^2$$


Antisymmetrize $\left[\frac{1}{z_1 - z_2} \frac{1}{z_3 - z_4} \frac{1}{z_5 - z_6} \dots \right]$

P-wave superconductor
of emergent particles
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Pfaffian which occurs at a filling fraction of 1/2 is not particle hole symmetric

‘Anti Pfaffian’ defined as the PH conjugate of the Pfaffian is a **different** topologically ordered state at the same filling fraction $\frac{1}{2}$ (Levin, Halperin, Rosenow PRL 2007, Lee et al PRL 2007)

Two-body interaction within a Landau level is PH symmetric \Rightarrow System spontaneously breaks PH symmetry. Inter Landau level mixing could provide a breaking field.

Aside: Actual resolution of this might be slightly more exotic (Banerjee et al arXiv:1710.00492, Mross et al arXiv 1711.06278; Wang et al arXiv 1711.11557; Simon arXiv:1801.09687)

LL Mixing - ineffective particle hole symmetry breaking

$$H_{\text{eff}} = H_0 + \kappa \sum_{p \in \text{intermediate states}} \frac{V |p\rangle \langle p| V}{E_0 - E_p}$$

Leads to correction to the two-body interaction Hamiltonian (which do not break particle hole symmetry)

+

PH symmetry breaking three-body interactions between the electrons.

(Bishara, Nayak PRB 2009, Simon, Rezayi PRL 2011, Sodeman, MacDonald 2013)

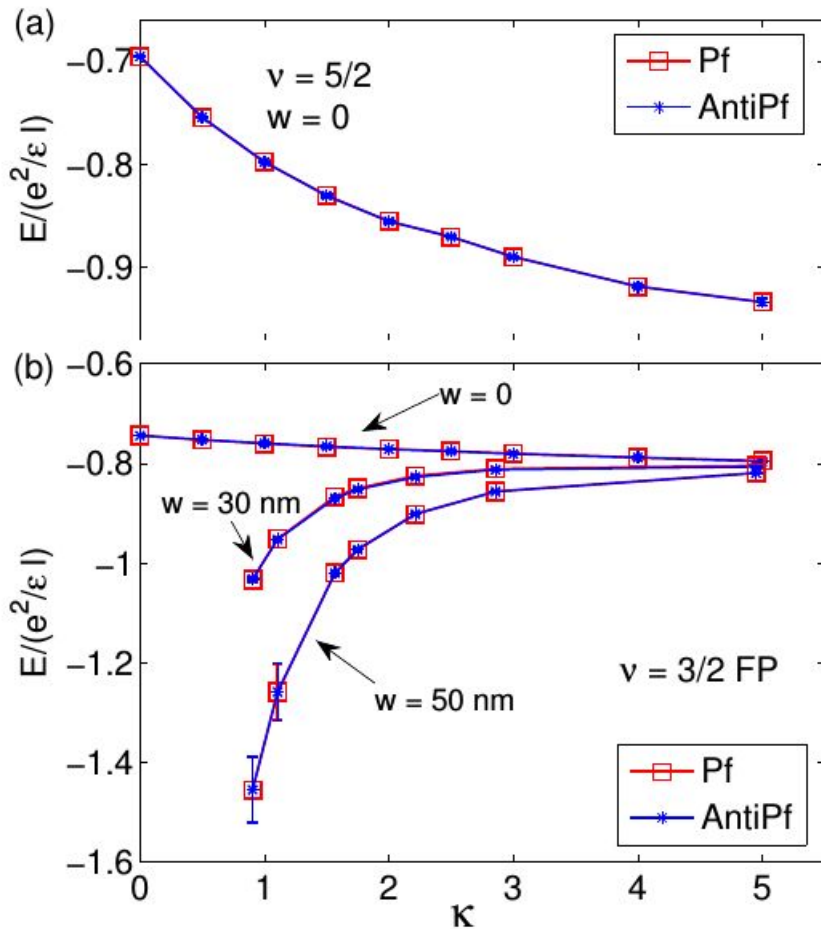
Some calculations point to a Pf ground state (Wojs et al PRL 2010, Pakrouski et al PRX 2015)

Others suggest a APf ground state (Rezayi, Simon PRL 2011, Zalatel et al 2015, Rezayi PRL 2017)

They all point to the fact that the effect of particle hole symmetry breaking is weak.

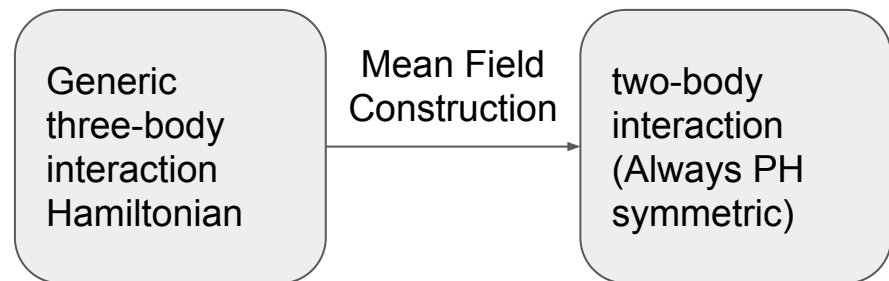
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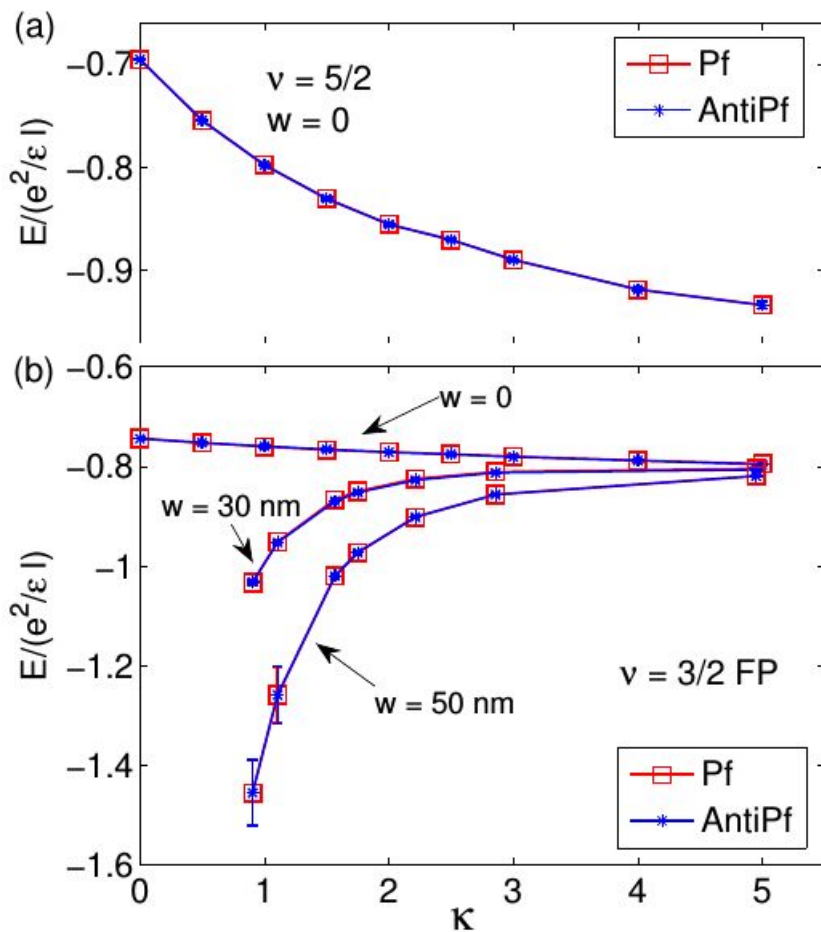
Effect of LL mixing on particle hole symmetry breaking is weak

⇒ The generic three-body interactions induced by LL mixing should effectively act as a PH symmetric two-body interaction.



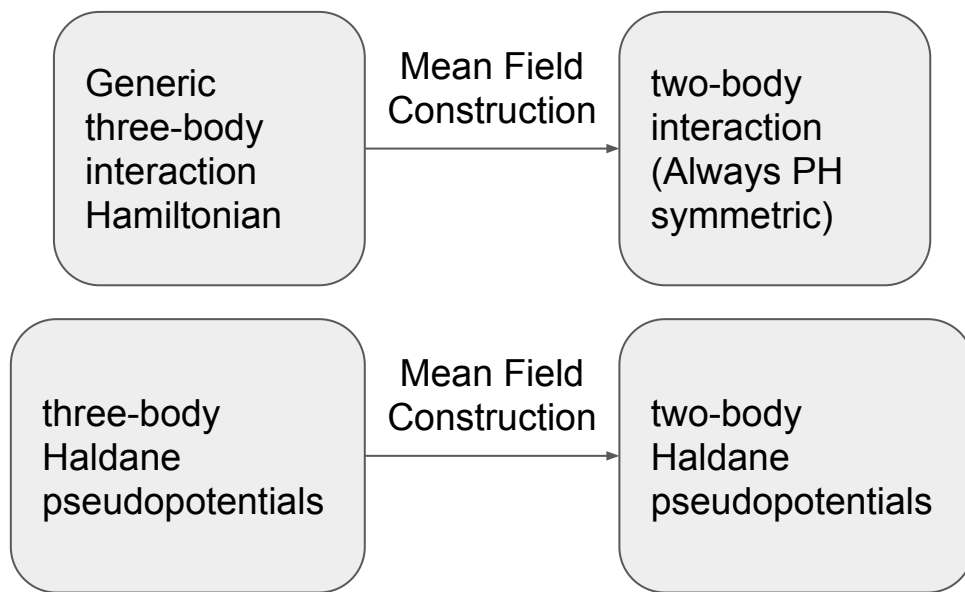
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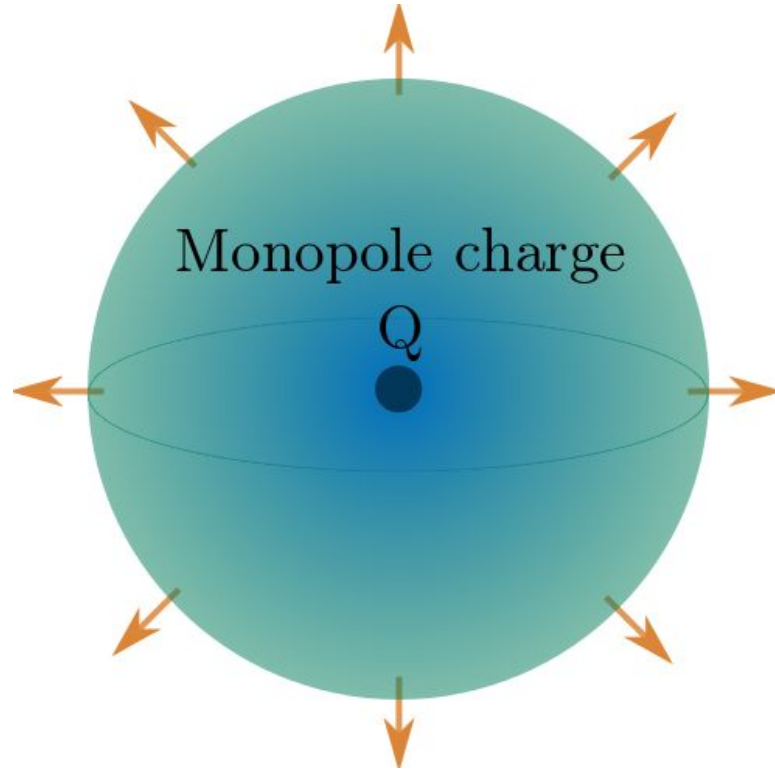
Spherical geometry and Haldane pseudopotentials

Bulk physics of the system can be effectively studied in a spherical geometry

Translational, rotational symmetry of the bulk of the system manifests as complete rotational symmetry on the sphere.

Energy depends only on the L^2 quantum number.

A rotational symmetric Hamiltonian can be parametrized using a few 'Haldane pseudopotentials' instead of all matrix elements.



$$\sum_{i < j=1}^N \sum_L V_L P_L^{ij}$$

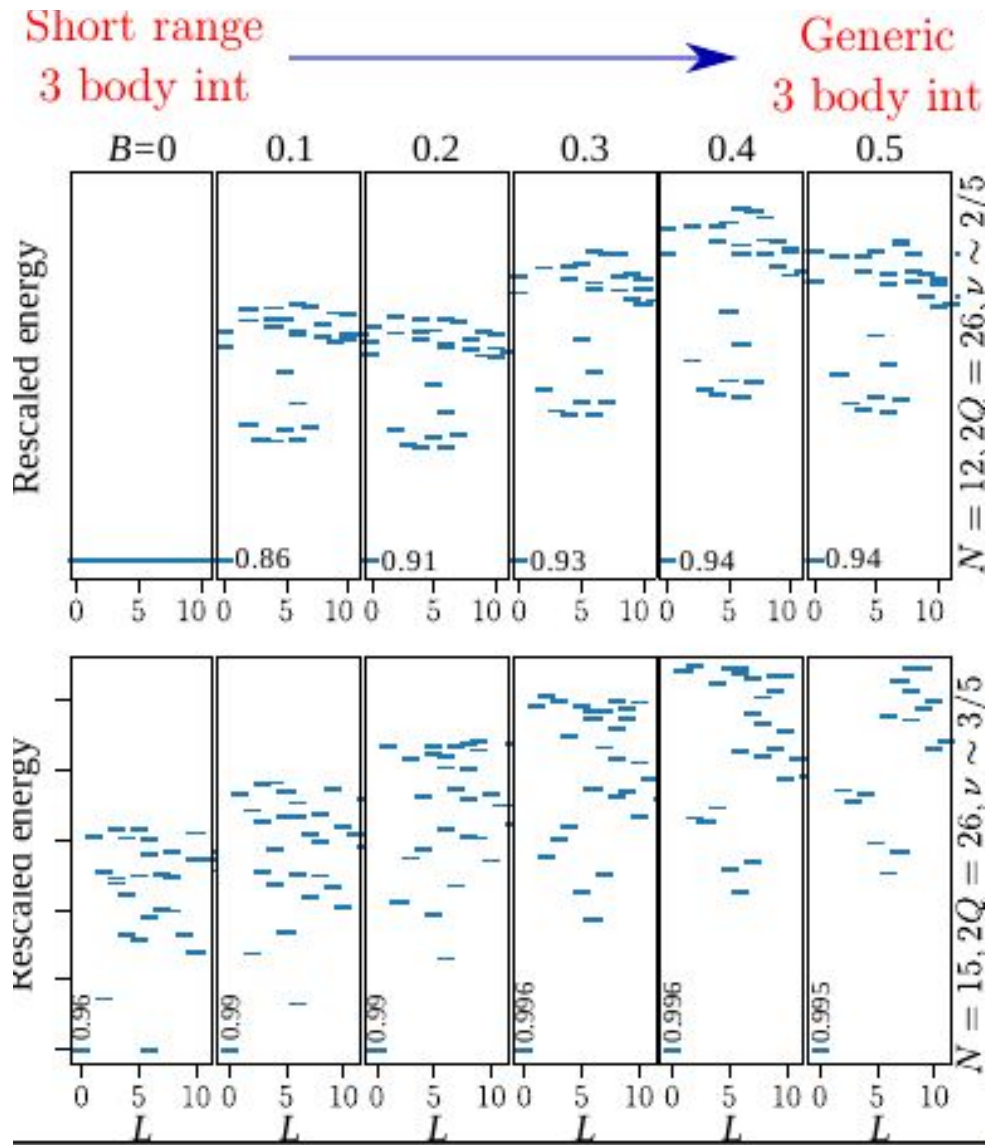
three-body interaction beyond the shortest range

Short range three-body interaction

$$H^{3BI-short} = \sum_{i < j < k=1}^N P_{L=3}^{ijk}$$

Generic three-body interaction

$$H^{3BI} = \sum_{i < j < k=1}^N \sum_L V_L P_L^{ijk}$$



Mean field approximation to the 3-body interaction

Can the generic three-body interaction be approximated with a two-body interaction ?

$$\begin{aligned} H^{3BI} &= \sum_{i < j < k = 1}^N \sum_L V_L P_L^{ijk} \\ &= \sum_{\mathbf{p}, \mathbf{q}} V_{p_1, p_2, p_3; q_1, q_2, q_3} c_{p_1}^\dagger c_{p_2}^\dagger c_{p_3}^\dagger c_{q_3} c_{q_2} c_{q_1} \end{aligned}$$


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$$MF(H^{3BI}) = H^{(2)} = \sum_{\mathbf{p}, \mathbf{q}} V_{p_1, p_2, p_3; q_1, q_2, q_3} c_{p_1}^\dagger c_{p_2}^\dagger \langle c_{p_3}^\dagger c_{q_3} \rangle c_{q_2} c_{q_1}$$

For any gapped
homogeneous
FQH state

$$= \sum_{\mathbf{p}, \mathbf{q}} V_{p_1, p_2, p_3; q_1, q_2, q_3} c_{p_1}^\dagger c_{p_2}^\dagger \delta_{p_3 q_3} c_{q_2} c_{q_1}$$


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For any
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$$= \sum_{\mathbf{p}, \mathbf{q}} V_{p_1, p_2, p_3; q_1, q_2, q_3} c_{p_1}^\dagger c_{p_2}^\dagger \delta_{p_3 q_3} c_{q_2} c_{q_1}$$

$$MF(H^{3BI}) = H^{(2)} = \sum_{\mathbf{p}, \mathbf{q}} \bar{V}_{p_1, p_2; q_1, q_2} c_{p_1}^\dagger c_{p_2}^\dagger c_{q_2} c_{q_1}$$

$$\bar{V}_{p_1, p_2; q_1, q_2} = \sum_x V_{p_1, p_2, x; q_1, q_2, x}$$

Mean field mapping: Properties

Linear
$$MF(H_a) + MF(H_b) = MF(H_a + H_b)$$

Rotational symmetry:

If H is rotationally symmetric, the interaction $MF(H)$ is rotationally symmetric.

$$MF(H) = \sum_{i < j=1}^N V_L^{2\text{body}} P_L^{ij}$$

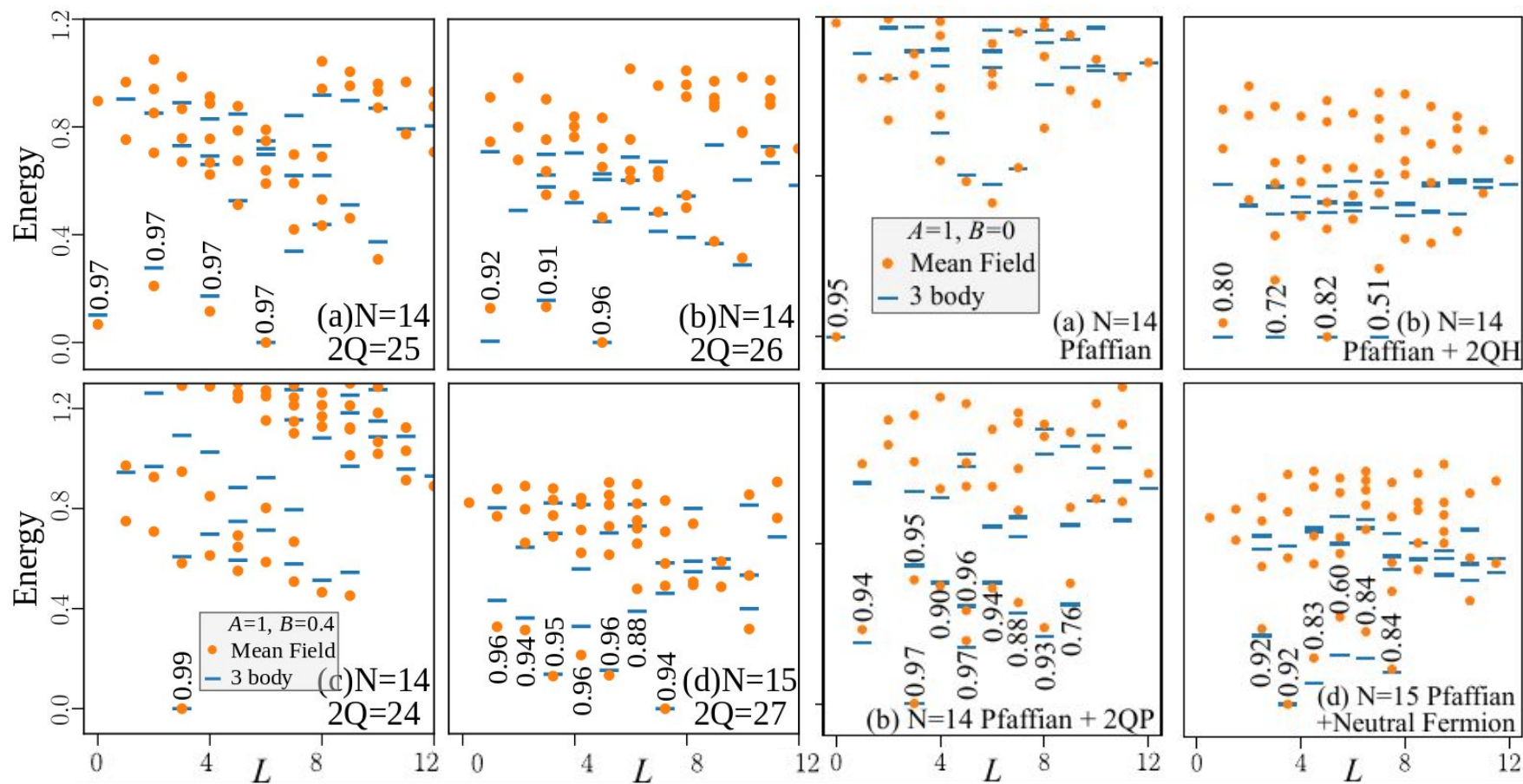
Thus MF is a well-defined linear mapping from three-body Haldane pseudopotentials to two-body Haldane pseudopotentials

PH symmetry: The mapping indeed gives a particle hole symmetric Hamiltonian by definition (it is a two-body interaction).

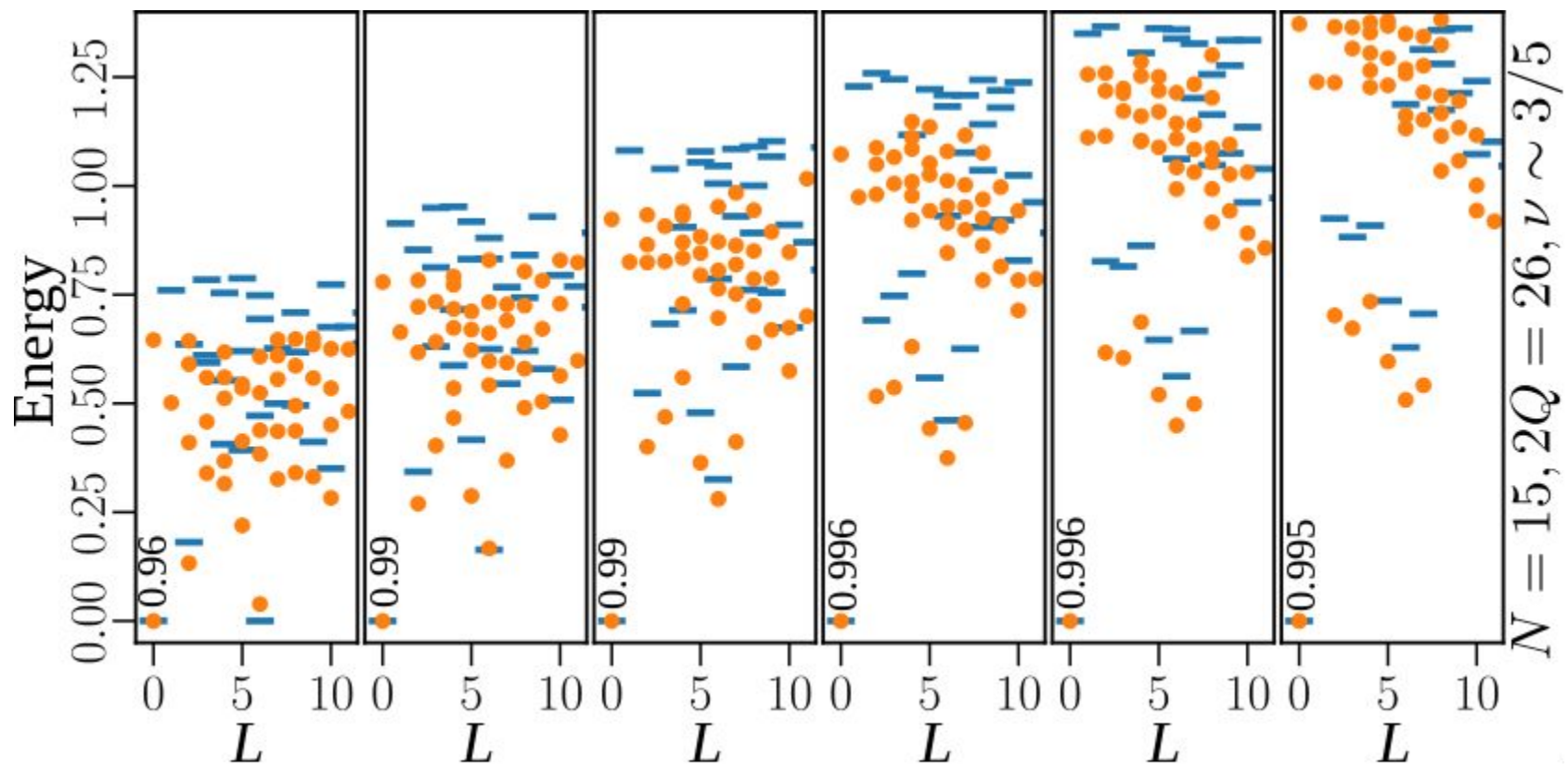
It indeed gives the particle hole symmetrization of the three-body interaction.

$$MF(H) = H + PHC(H)$$

Testing the mean field mapping



Testing the mean field mapping



Summary

Various short-range n -body terms in the Hamiltonian specify constraints on the many-particle correlations in the low-energy wavefunctions.

Generic n -body constraints on the many-particle wavefunctions imply and are implied by constraints on a smaller number of particles.

Fewer-body terms should reproduce the low-energy spectrum --> mean field approximation

Can be generalized to the $n > 3$ body case. (Kusmierz, GJS (in preparation))

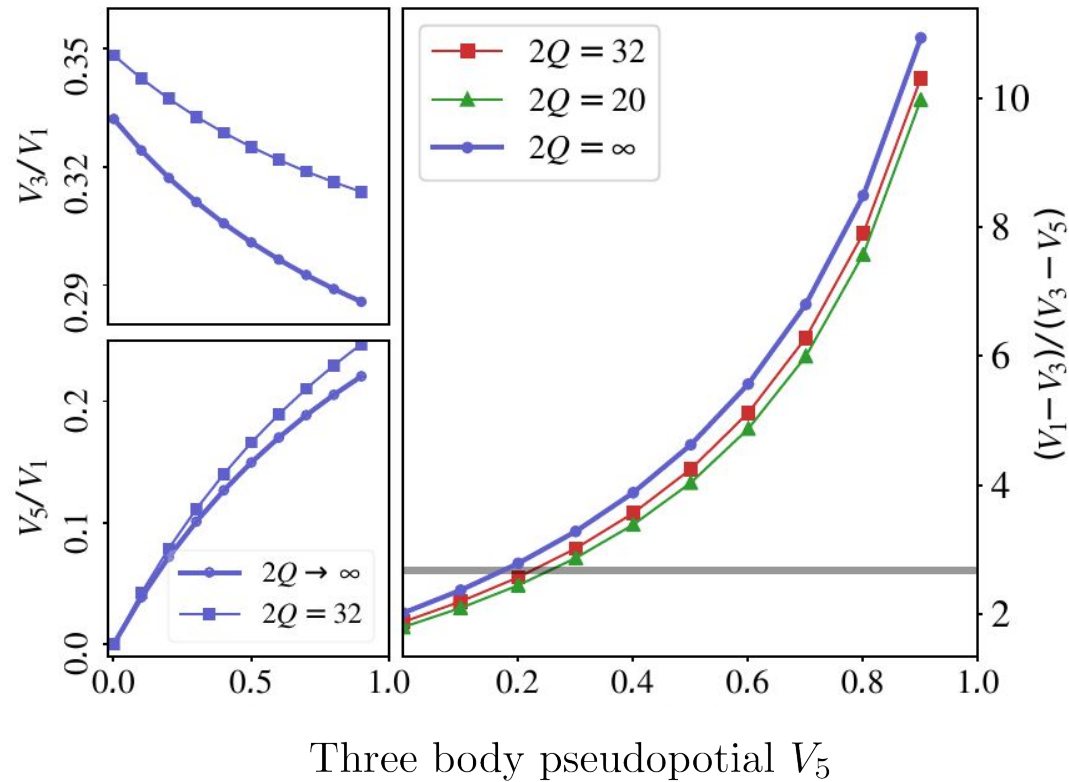
Mean field pseudopotentials

MF approximation to the short range three-body interaction gives the two-body Hamiltonian:

$$H = \sum_{ij} 3P_{L=1}^{ij} + 1P_{L=3}^{ij}$$

ie $V_3=1$ and $V_1=3$

With increasing range, the mean field Hamiltonian contains one more pseudopotential.



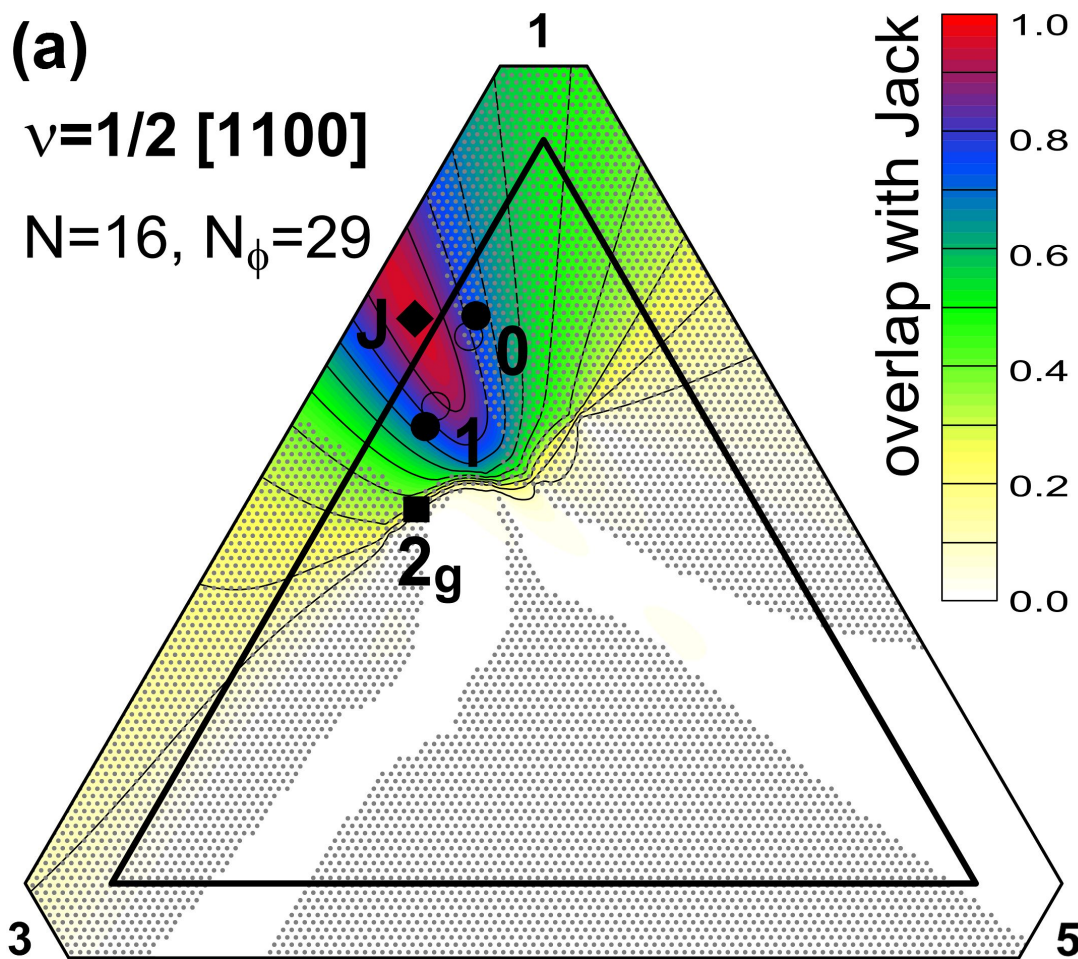
Optimal two-body interaction for the Pfaffian state

Scan for an optimal two-body interaction that produces the Pfaffian ground state

(a)

$\nu=1/2$ [1100]

$N=16, N_\phi=29$



Mean field approximation to three-body interaction gives the optimal parameters for two-body interaction that produces the Pfaffian state

A Wojs, B Kusmierz (unpublished)
B Kusmierz, G J Sreejith (in preparation)