Negative mobility in interacting particle systems



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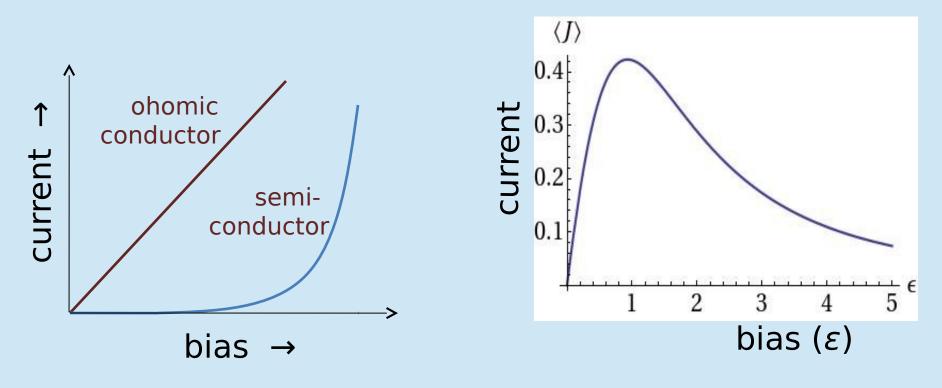
NEGATIVE RESPONSE

normally

current always increases with bias

negative response

current decreases with *increasing bias*



What I intend to discuss.....

Mechamism of **NEGATIVE RESPONSE** in **NON-EQUILIBRIUM SYSTEMS**

In equilibrium $\langle J \rangle = 0$ but current fluctuates (*microscopic*) $\langle J^2 \rangle = finite$

Like in a paramagnet $\langle M \rangle = 0$ but $\langle M^2 \rangle = finite$

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Description:
$$Z(z) = \sum_{M=0}^{\infty} z^{M} Q_{M}$$
 $z = e^{-h}$ $\langle M \rangle = z \frac{d}{dz} \ln Z$
(of para-magnet) $\frac{d}{dh} \langle M \rangle = \langle M^{2} \rangle - \langle M \rangle^{2} \equiv \chi$

With a small field: $\langle M \rangle_h = \langle M \rangle_0 + h \frac{d}{dh} \langle M \rangle_0 = h \langle M^2 \rangle_0$

Response would be propertional to the fluctuations at zero-filed.

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Same analogy

$$J(\epsilon) = J(0) + \epsilon (\langle J^2 \rangle - \langle J \rangle^2 \rangle_{\epsilon=0}$$

= $\epsilon \langle J^2 \rangle_0$

Direction of current = direction of fileld

$$\frac{dJ}{d\epsilon} = \langle J^2 \rangle - \langle J \rangle^2 \ge 0$$

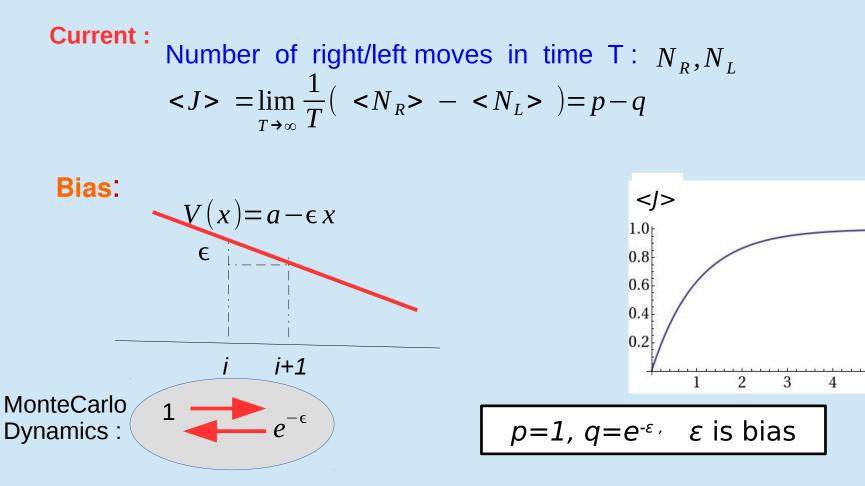
Response is positive

Calculating Response:

Single random walker(A): 1-D periodic lattice, L sites

$$A0 \mathop{\rightleftharpoons}\limits_{q}^{p} 0A$$

5



Many walker & exclusion :

$$A0 \stackrel{p}{\underset{q}{\rightleftharpoons}} 0A$$
$$p=1, q=e^{-\varepsilon}$$

$$\langle J \rangle = p - q$$

$$\stackrel{d}{\Longrightarrow} \frac{d J}{d \epsilon} \ge 0$$

Many walker & exclusion :

<J> = p-q

 $\rightarrow \frac{dJ}{d\epsilon} \ge 0$

$$A0 \stackrel{p}{\rightleftharpoons} 0A$$
$$q = 0A$$
$$p=1, q=e^{-\varepsilon}$$

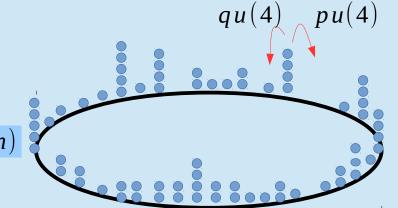
n

Many walker, no exclusion :

- Right/left hop rates: $u_R(n) = pu(n)$ $u_L(n) = qu(n)$
- Zero Range Process (ZRP), exactly solvable.
- Steadystate weights: $P(\{n_i\}) \sim \prod_{i=1}^{L} f(n_i)$, where $f(m) = \prod_{k=1}^{m} \frac{1}{u(k)}$

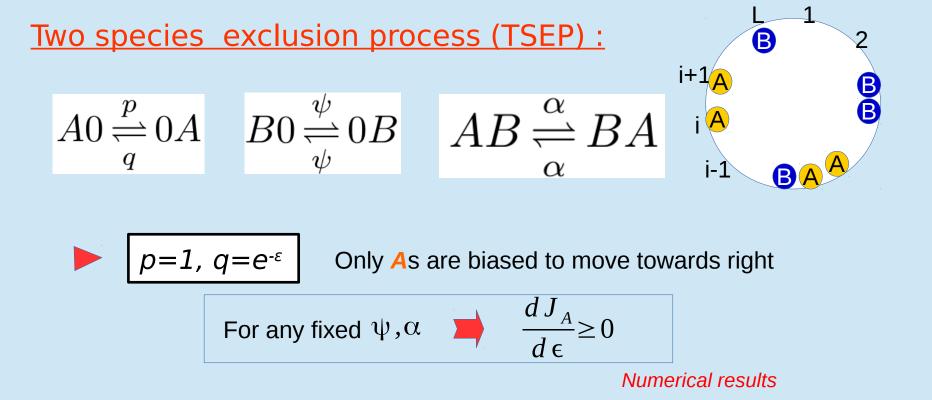
• Current :
$$J = \langle u_R(n) \rangle - \langle u_L(n) \rangle = (p-q)z$$

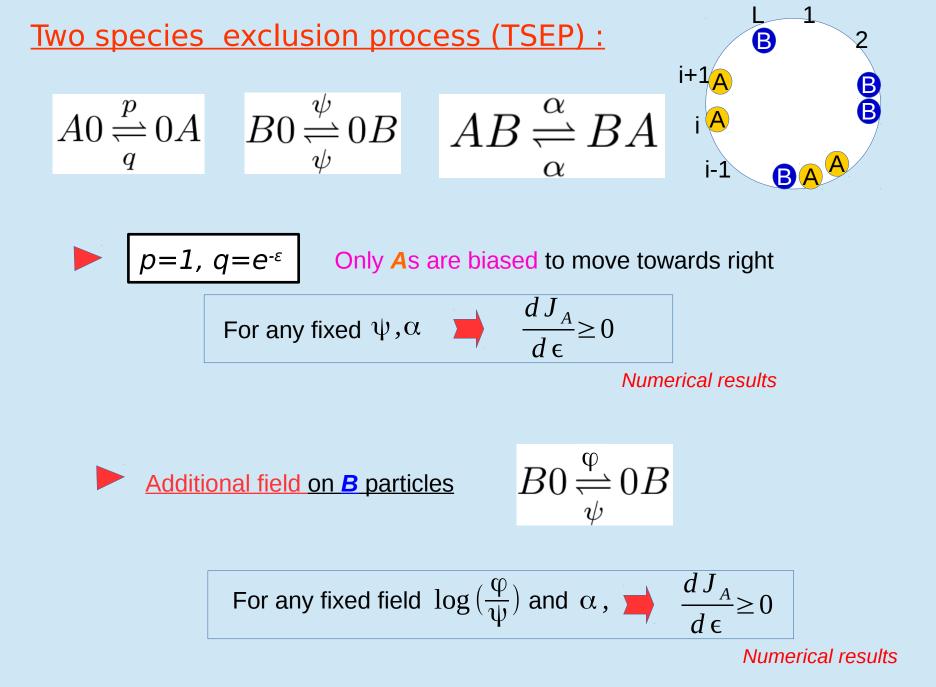
$$\stackrel{d}{\longrightarrow} \frac{dJ}{d\epsilon} \ge 0$$



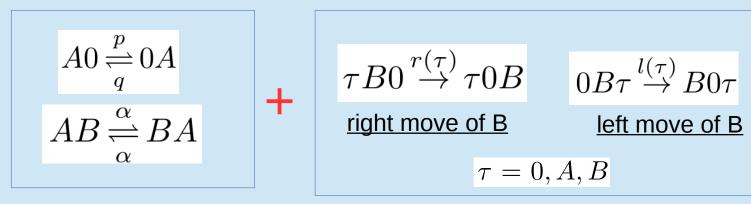
ρ

max





<u>Additional interaction</u> among A,B particles :

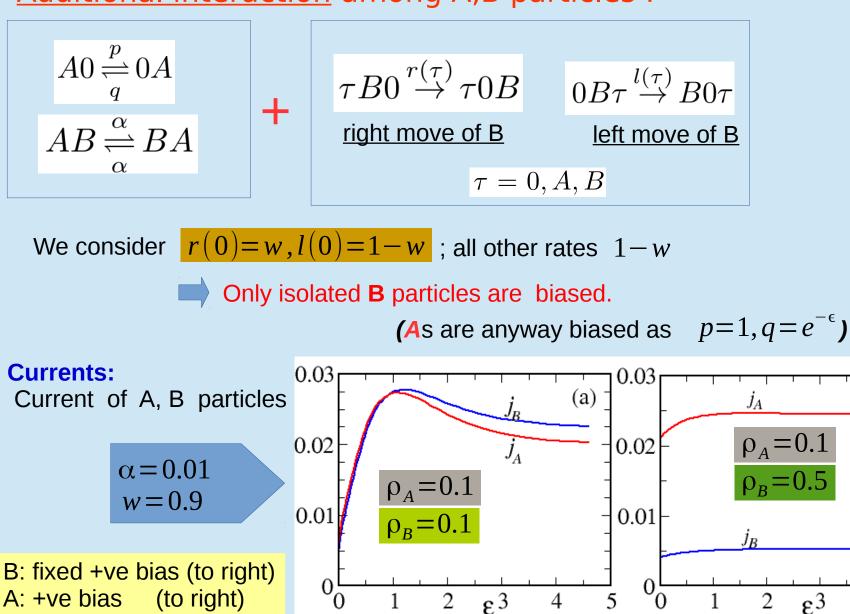


+ We consider r(0) = w, l(0) = 1 - w; all other rates 1 - w

Only isolated **B** particles are biased.

(As are anyway biased as $p=1, q=e^{-\epsilon}$)

<u>Additional interaction</u> among A,B particles :



(a)Negative diffrential mobility/respose (NDM) for j_A , j_B , $j = j_A + j_B$

b)

4

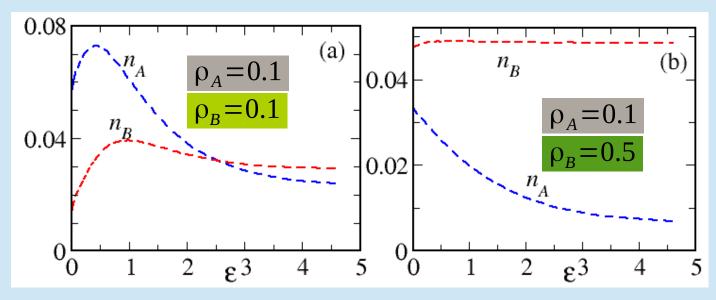
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Why does it exhibit NDM ?



We measure other observables n_A

along with $j_{A,B} = \langle n_r^{A,B} + n_l^{A,B} \rangle$



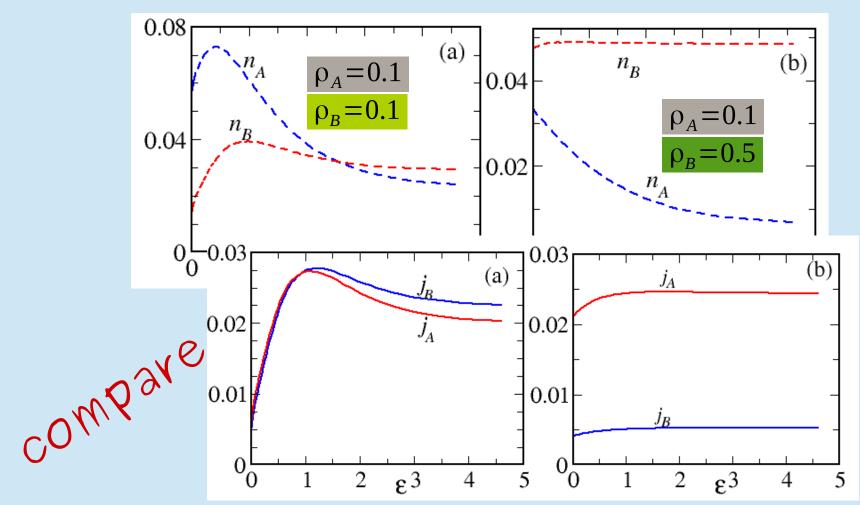
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We measure oher observables

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 $n_{A,B} = \langle n_r^{A,B} + n_l^{A,B} \rangle$ $j_{A,B} = \langle n_r^{A,B} - n_l^{A,B} \rangle$

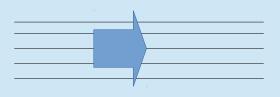


Current of A particles (driven-sector) follows the traffic of B (not driven by ε)

Conjecture :

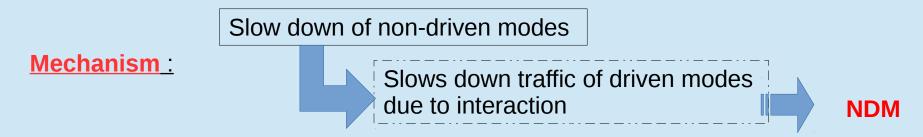
Possible mechanism for NDM in interacting systems:

particle current in a driven many particle system might show a nonmonotonic behaviour if some modes, which are not driven by the external field, slow down with increased driving.



Interacting particle system has many current-carrying modes (here A, B)

What is necessary for NDM: traffic of driven degrees must slow down



What, if we explicitly slow down Bs in TSEP?

B

B

i+1

i-1

2

Β

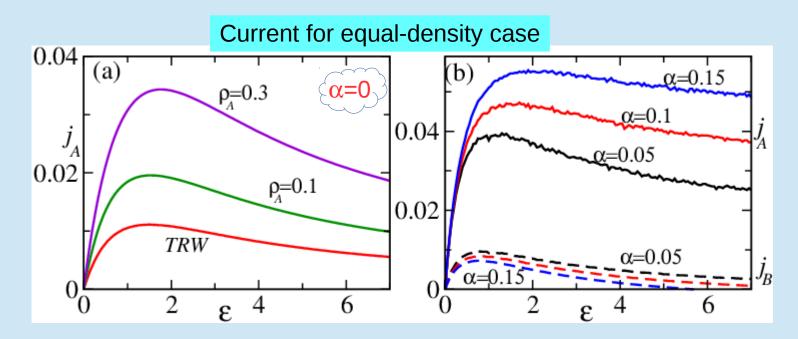
$$A0 \stackrel{p}{\underset{q}{\rightleftharpoons}} 0A \qquad B0 \stackrel{\psi}{\underset{\psi}{\rightrightarrows}} 0B \qquad AB \stackrel{\alpha}{\underset{\alpha}{\rightrightarrows}} BA$$

Now $p=1, q=e^{-\epsilon}$, but ψ is a decreasing function of ϵ , say $\psi=1/(1+\epsilon)$

What, if we explicitly slow down Bs in TSEP?

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Now $p=1, q=e^{-\epsilon}$, but Ψ is a decreasing function of ϵ , say $\psi=1/(1+\epsilon)$



Indeed TSEP now exhibits NDR

TSEP is exactly solvable for α =0

B

B/

B

i+1

i-1

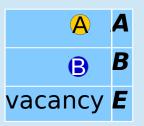
Just two walkers : (no exchange) $A0 \stackrel{p}{\rightleftharpoons} 0A = B0 \stackrel{\psi}{\underset{\psi}{\to}} 0B$

$$s_i = A, B, 0$$

Exact solution using Matrix Product Ansatz

$$P(C) = P(\lbrace s_i \rbrace) = Tr(\prod_{i=1}^{L} X_i)$$
$$X_i = A\delta_{s_i,A} + B\delta_{s_i,B} + E\delta_{s_i,0}$$

 $P(000 \triangle 0... \triangle 0...) = Tr(EEEAE...BE...)$



B

i+1

i-1

• Just two walkers:
(no exchange)
$$A0 \stackrel{p}{\rightleftharpoons} 0A B0 \stackrel{\psi}{\hookrightarrow} 0B$$

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Matrix Algebra :

$$p A E - q E A = x_0 A$$

$$\psi (B E - E B) = x_0 B$$

$$A^2 = 0 , B^2 = 0$$

Matrix representation :

$$A = \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} \qquad B = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \qquad E = \begin{pmatrix} r & 0 \\ 0 & 1 \end{pmatrix}$$

Here, $r = \frac{p+\psi}{q+\psi}$

• Average particle current :

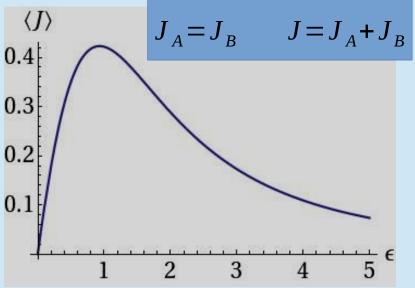
$$=(p-q) + \psi(-) = \frac{2(p-q)\psi}{p+\psi}$$

$$\langle J \rangle(\varepsilon) = \frac{2(1-e^{-\varepsilon})}{(2+\varepsilon^2)}$$

$$p = 1, q = e^{-\varepsilon}, \psi = \frac{1}{(1 + \varepsilon^2)}$$

NDM for any \psi, as long as \psi < 0

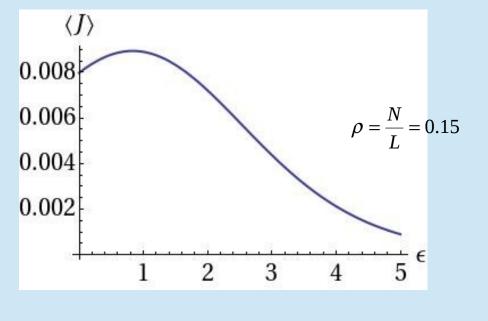




More examples: Asymmetric misanthrope process1-D periodic latticei = 1, ..., L sites $u_L(3,1)$ $u_R(3,2)$ $n_i \ge 0$ $u_R(3,2)$ $u_R(3,2)$ $\sum_{i=1}^L n_i = N$

No hard core exclusion **Dynamics** - right hop: $u_R(n_i, n_{i+1})$ left hop: $u_L(n_i, n_{i+1})$

 $u_{R}(m,n) \quad p0 , \quad n=0$ 1, otherwise $u_{L}(m,n) \quad p0 , \quad m=1$ $e^{-\varepsilon} , \\m>1, \quad n=0$ $\frac{1/2}{2} , \\m>1, \quad n>0$



But, what went wrong with the response formula ?

 \boldsymbol{N}_{R} , \boldsymbol{N}_{L} fluctuates independently

Correspondingly, $J = N_R - N_L$ $S = N_R + N_L$ fluctuates, but J and S *are correlated*

Correct response formula :

$$\frac{d}{d \epsilon} < J > = -const. < J; S >$$

In equilibrium: < J > = 0,

Baiesi et. al., J. Stat. Phys 2009

$$\frac{d}{d \epsilon} < J > = =$$

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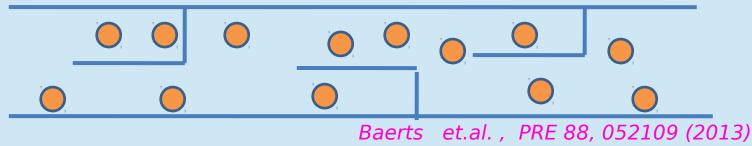
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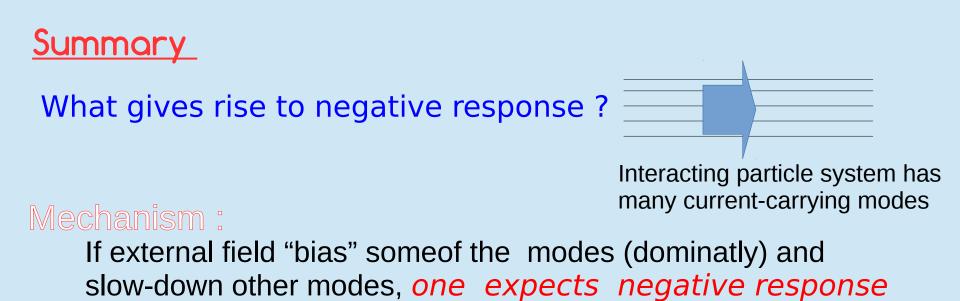
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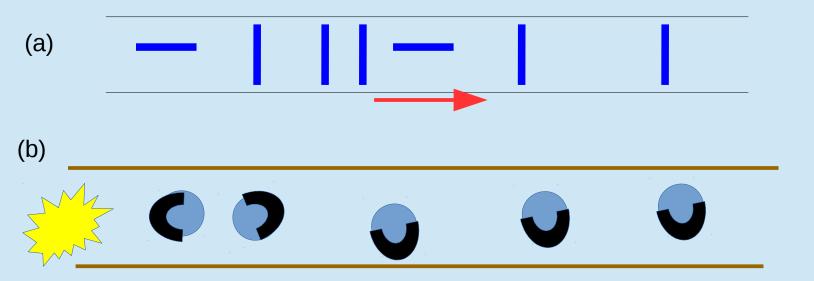
Baiesi et. al., J. Stat. Phys 2009

Non- equilibrium: driven non-interacting particles





Experimrnts ?



Thanks to Amit, Urna

Zero range and finite range processes with asymmetric rate functions A Chatterjee, PKM, JSTAT2017

MPA for interacting-particles systems without hardcore-constraints.

A Chatterjee, PKM, Jphys A 2017

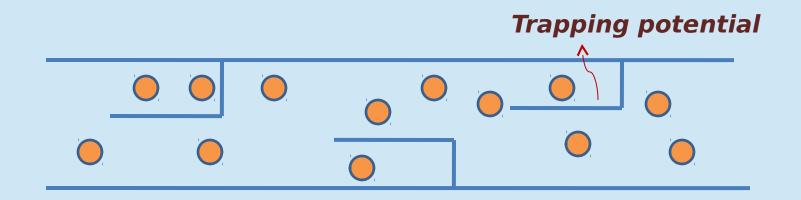
Mechanism of negative response in interacting-particles systems

A Chatterjee, U basu, PKM, (*arXiv:1712.01236*)





• **Negative Response** : for *non-interacting particles*



(a) q=1 J=0

q
$$\xrightarrow{1} 1$$

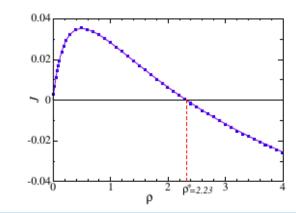
(c) q=0.05 , most of the particles get trapped, negative response

> Baerts et.al., PRE 88, 052109 (2013)

Many other examples

Asymmetric misanthrope process (AMAP):

Hop-rate depends on occupation of dep. and arv. sites. But the rate functions are different for R,L



 \sim Asymmetric finite range process (AFRP):

Hop-rate depends on occupation of dep. Site and its R- nearest neighbours. But the rate functions are different for R,L.

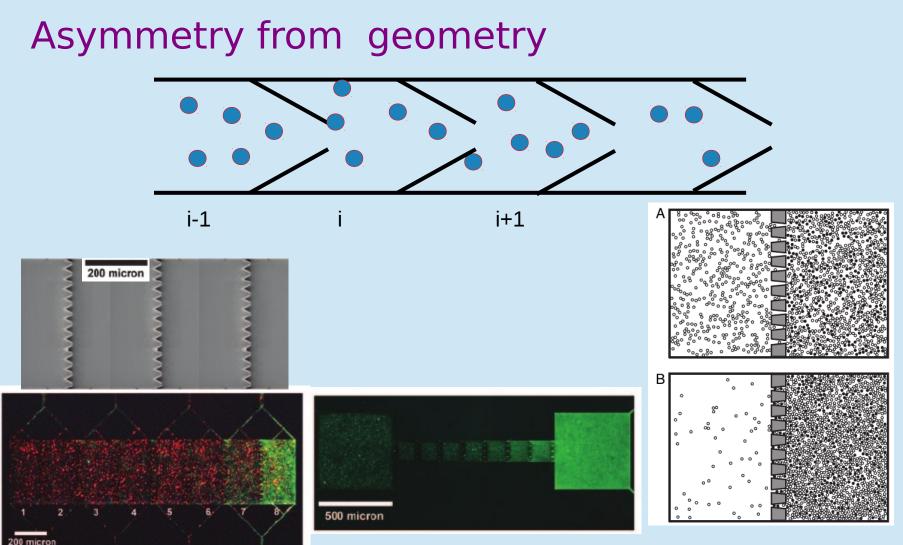
 \therefore Higher dimensions:

AZRP and AMAP shows current reversal in higher dimension.

Requirement

Interacting particle system where some internal degrees experience a different bias then the rest.

Up-down, active-inactive, isolated-crowded



A Wall of Funnels Concentrates Swimming Bacteria, P Galajda, J Keymer, PChaikin, and Robert Austin. JOURNAL OF BACTERIOLOGY, Dec. 2007, p. 8704–8707 Geometry-induced asymmetric Diffusion, RS Shaw, N Packard,

9580–9584 | PNAS | June 5, 2007 | vol. 104 |