Permanent spin currents in cavity-qubit systems

Manas Kulkarni

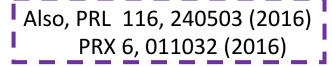
International Centre for Theoretical Sciences

Tata Institute of Fundamental Research, Bangalore

PRB 97, 064506 (2018)

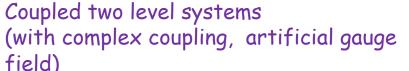
Collaboration:

- C. Aron (ENS Paris / Belgium)
- S. Hein (Technische Univ, Berlin)
- E. Kapit (Tulane, USA)

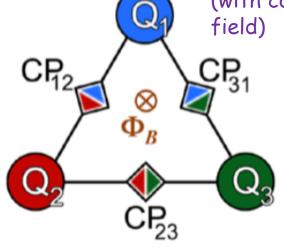


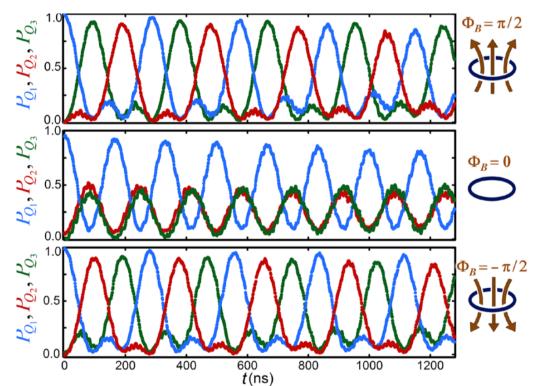


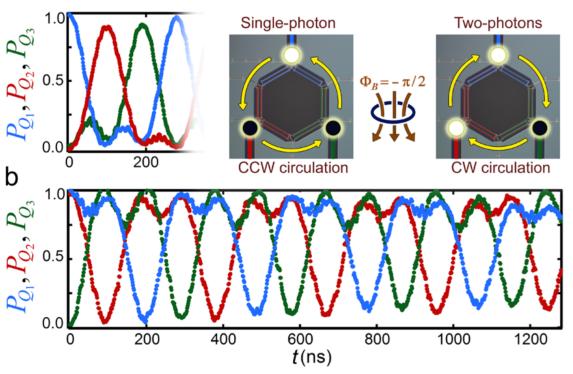




Google Inc / Santa Barbara Experiment







$$H(t) = \hbar \sum_{j=1}^{3} \omega_j(\hat{n}+1/2) + \hbar \sum_{j,k} g_{jk}(t) (a_j^{\dagger} a_k + a_j a_k^{\dagger}) + H_{\text{int}}$$

J. Martinis (Google Inc / Santa Barbara)

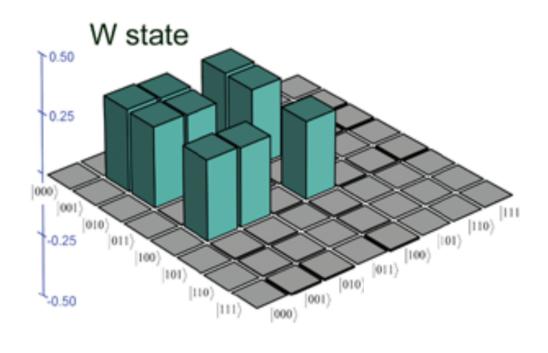
P. Roushan et al,

Nature Physics 13, 146–151 (2017)

Chiral currents

b Chiral current $<\hat{I}_{chiral}>$ -3 Q_1 • two-photon -2π -π 2π $\Phi_{_{\mathrm{B}}}(\mathsf{rad})$

Creating W-states



P. Roushan et al Nature Physics 13, 146–151 (2017)

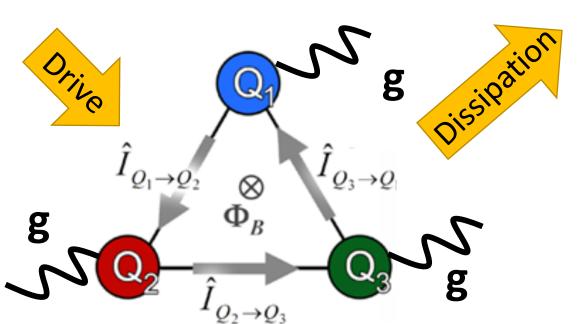
Quantum State Tomography

However, currents decay and are not persistent.

Systems are subject to inevitable environmental effects

Designing a protocol to create persistent currents

- striking a delicate balance between drive and dissipation to activate, specific entangled states with high fidelity, which are capable of carrying current
- power of quantum bath engineering approaches to realize highly non-trivial nonequilibrium steady states in Open Quantum Systems.



- Sources of dissipation are qubit dissipation/dephasing and cavity decay
- Qubits are coupled with complex hoping
- Each Qubit is coupled to a respective photon mode (g)
- Photons are driven by a classical microwave drive

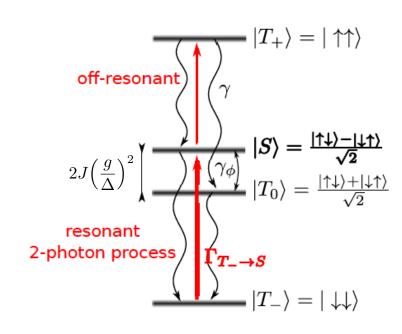
Hamiltonians with complex coupling realized

P. Roushan et al Nature Physics (2017)

Many fascinating variations possible

E. Kapit, PRA (2015)





(Experiment realized @Berkeley, PRL 2016)

Protocols to strike a delicate balance between drive and dissipation (especially in context of entanglement)

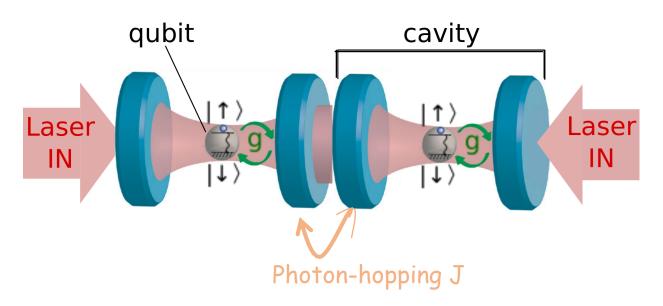
Aron, **Kulkarni**, Tureci (PRX 2016) Schwartz, Martin, Flurin, Aron, **Kulkarni**, Tureci, Siddiqi (PRL 2016) Aron, **Kulkarni**, Tureci (PRA 2014)

We try to unite these two ideas

Engineering a quantum device and the environment

Quantum Nanoelectronics Laboratory, Berkeley

Schwartz, Martin, Flurin, Aron, **Kulkarni**, Tureci, Siddiqi (PRL 2016)



Driven unitary evolution

- ullet two-level systems $\omega_{
 m q} rac{\sigma_i^{
 m z}}{2}$, $\omega_{
 m q} \sim 10$ GHz
- ullet far-detuned cavities $\omega_{
 m c}
 eq \omega_{
 m q}$
- ullet light-matter coupling $g\sim$ 0.1 GHz
- ullet inter-cavity coupling $J\sim 0.1~{
 m GHz}$
- microwave drives: $\epsilon_{\rm d} \cos(\omega_{\rm d} t)(a_i + a_i^{\dagger})$

Dissipative processes

- ullet cavity decay $\kappa\sim 10^{-3}~{
 m GHz}$
- qubit decay $\gamma \sim 10^{-4}~{\rm GHz}$
- ullet qubit dephasing $\gamma_{\phi}\sim 10^{-5}~\mathrm{GHz}$

Jaynes-Cummings Array to XY Spin chain coupled to photons

$$\begin{split} H_{\sigma} &= \sum_{i} \omega_{\mathrm{q}} \frac{\sigma_{i}^{z}}{2}, \qquad H_{\sigma a} = g \sum_{i} [a_{i}^{\dagger} \sigma_{i}^{-} + \mathrm{H.c.}] \\ H_{a}(t) &= \sum_{i} [\omega_{\mathrm{c}} a_{i}^{\dagger} a_{i} - J(a_{i}^{\dagger} a_{i+1} + \mathrm{H.c.}) \\ &+ 2\epsilon_{i}^{\mathrm{d}} \cos(\omega_{\mathrm{d}} t + \Phi_{i})(a_{i} + a_{i}^{\dagger})]. \end{split}$$

$$H_a(t) = \sum_{i} [\omega_{c} a_i^{\dagger} a_i - J(a_i^{\dagger} a_{i+1} + \text{H.c.})]$$

$$+2\epsilon_i^{\mathrm{d}}\cos(\omega_{\mathrm{d}}t+\Phi_i)(a_i+a_i^{\dagger})].$$

RWA+ Schrieffer-Wolff Transformation

$$H_{\sigma} = \sum_{i} h_{i} \cdot \frac{\sigma_{i}}{2} - \frac{J}{2} \left(\frac{g}{\Delta} \right)^{2} [\sigma_{i}^{x} \sigma_{i+1}^{x} + \sigma_{i}^{y} \sigma_{i+1}^{y}],$$

$$H_{\sigma a} = \sum_{i} \left(\frac{g}{\Delta}\right)^{2} \sigma_{i}^{z} (\Delta a_{i}^{\dagger} a_{i} + \epsilon_{i}^{\mathrm{d}} a_{i}^{\dagger} + \epsilon_{i}^{\mathrm{d}*} a_{i}).$$

Baths/Noise/Imperfections

$$H_{total} = H_{system} + H_{bath}$$

$$H_{\sigma,b} = \eta \sum_{i\,n} \sigma_i^x (b_{in} + b_{in}^\dagger), \quad H_b = \sum_{i\,n} \omega_n b_{in}^\dagger b_{in}$$
 Qubits coupled to Bath

$$H_{a, ilde{b}}= ilde{\kappa}\sum_{i,n}(a_i^\dagger ilde{b}_{in}+ ilde{b}_{in}^\dagger a_i), H_{ ilde{b}}=\sum_{i,n}\omega_n ilde{b}_{in}^\dagger ilde{b}_{in}^{} ilde{b}_{in}^{}$$
 to bath

$$H_{\sigma} = \sum_{i} h_{i} \cdot \frac{\sigma_{i}}{2} - \frac{J}{2} \left(\frac{g}{\Delta} \right)^{2} [\sigma_{i}^{x} \sigma_{i+1}^{x} + \sigma_{i}^{y} \sigma_{i+1}^{y}],$$

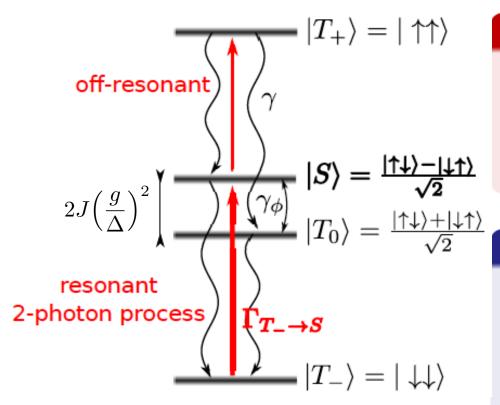
$$H_{\sigma a} = \sum_{i} \left(\frac{g}{\Delta}\right)^{2} \sigma_{i}^{z} (\Delta a_{i}^{\dagger} a_{i} + \epsilon_{i}^{\mathrm{d}} a_{i}^{\dagger} + \epsilon_{i}^{\mathrm{d}*} a_{i}).$$

System Hamiltonian

Aron, **Kulkarni**, Tureci (PRA, 2014) Aron, **Kulkarni**, Tureci (PRX 2016)

Nonequilibrium rate from $|T_{-}\rangle$ to $|S\rangle$

$$\Gamma_{T_- \to S} \propto 2\pi (g^3 \epsilon_{\rm d}^2)^2 \rho_- (\omega_{\rm d} + E_{T_-} - E_S)$$



Transverse-field isotropic XY-model

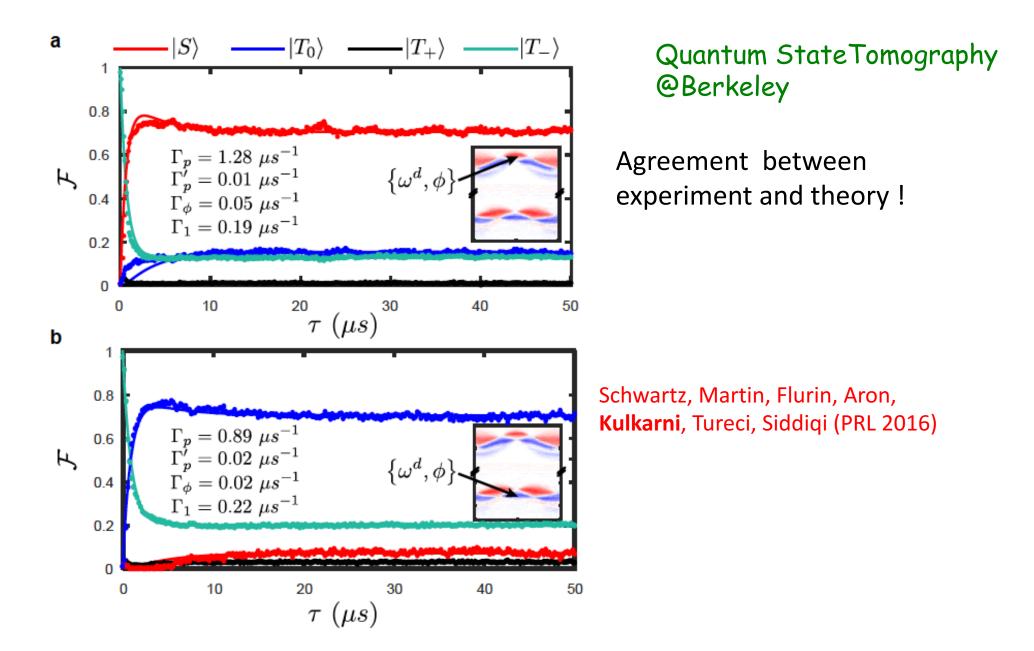
$$H_{\sigma} = \sum_{i=1}^{2} \mathbf{h} \cdot \frac{\sigma_{i}}{2} - \frac{J}{2} \left(\frac{g}{\Delta}\right)^{2} \left[\sigma_{1}^{x} \sigma_{2}^{x} + \sigma_{1}^{y} \sigma_{2}^{y}\right]$$
$$h^{x} \equiv \mathbf{g} \epsilon_{d}, h^{y} \equiv 0, h^{z} \equiv \omega_{d} - \omega_{d}$$

Weakly-coupled photon-fluctuation bath

$$H_{\sigma d} = \#(\frac{g}{\Delta})^{2} \epsilon_{\mathbf{d}} \sum_{i=1}^{2} \sigma_{i}^{z} (d_{i} + d_{i}^{\dagger})$$

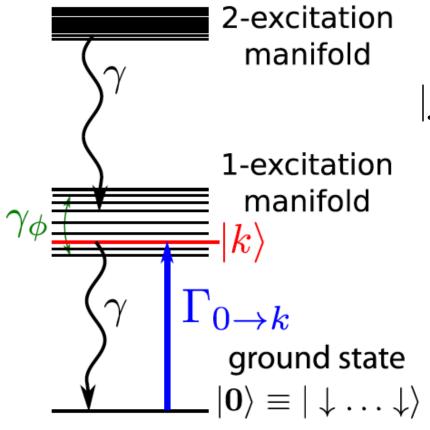
$$H_d = \sum_{i=1}^{2} (\omega_c - \omega_d) d_i^{\dagger} d_i - J(d_1^{\dagger} d_2 + \text{h.c.})$$

Experimental Results: Time Dynamics (time evolution of entanglement)



Creating a target generalized-W state

Target state



$$|k\rangle \equiv \sum_{j} e^{ikj} |j\rangle$$

$$|j\rangle = |\downarrow_0 \dots \downarrow_{j-1} \uparrow_j \downarrow_{j+1} \dots \downarrow_{N-1} \rangle$$

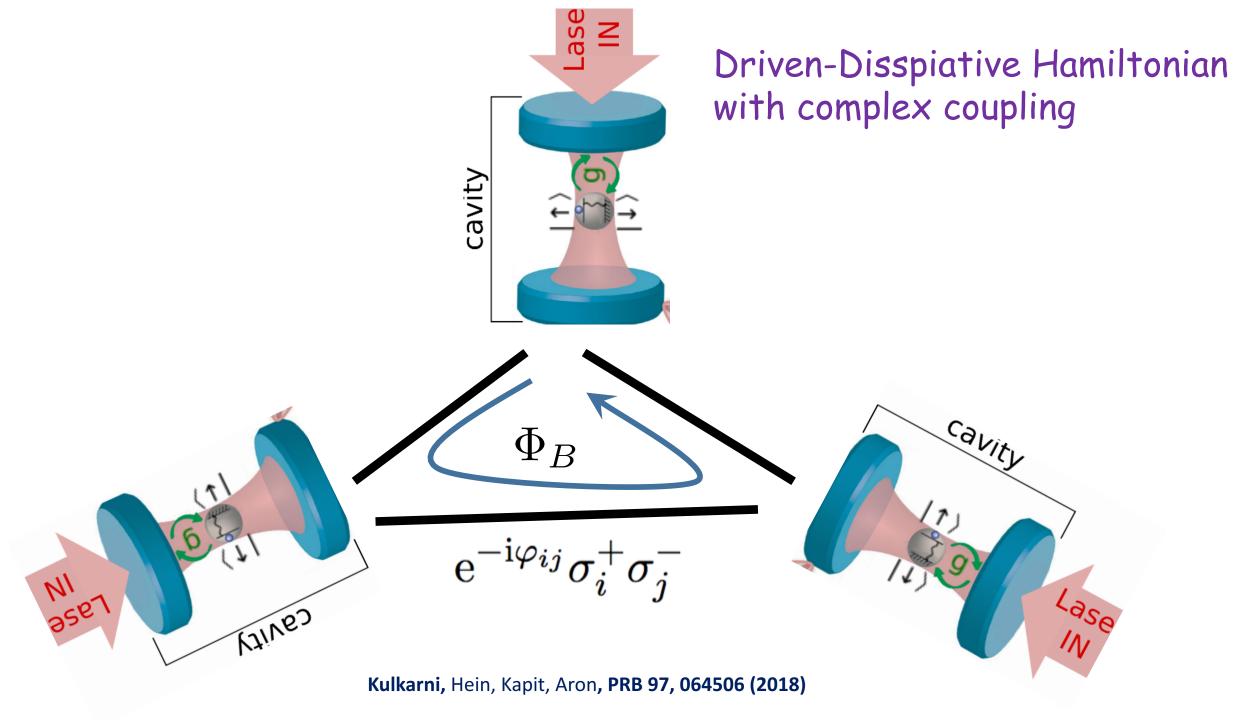
Single magnon excitations

For e.g., N=2 and k=0, target state is

$$|T_0>=\frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle+|\downarrow\uparrow\rangle)$$

For e.g., N=3 and k=0, target state is

$$W = \frac{1}{\sqrt{3}} \left(|\downarrow\downarrow\uparrow\rangle + |\downarrow\uparrow\downarrow\rangle + |\uparrow\downarrow\downarrow\rangle \right)$$



$$H_{\sigma} = \sum_{i} (\omega_{\mathbf{q}} - \omega_{\mathbf{d}}) \frac{\sigma_{i}^{z}}{2} - \sum_{\langle ij \rangle} J_{0} \left[e^{-i\varphi_{ij}} \sigma_{i}^{+} \sigma_{j}^{-} + \text{h.c.} \right]$$

$$H_{\sigma a} = \sum_{i} g \left[a_i^{\dagger} \sigma_i^{-} + \text{h.c.} \right] ,$$

$$H_a = \sum_{k} (\omega_{\rm c} - \omega_{\rm d}) a_k^{\dagger} a_k + (\epsilon_k^{\rm d} a_k^{\dagger} + \text{h.c.}).$$



SW transformation

$$H_{\sigma} = \sum_{i} \boldsymbol{h_i} \cdot \frac{\boldsymbol{\sigma_i}}{2} - \sum_{i} J_0 \left[e^{i\phi} \sigma_i^+ \sigma_{i+1}^- + \text{h.c.} \right],$$

$$H_{\sigma a} = \sum_{i} \left(\frac{g}{\Delta}\right)^{2} \sigma_{i}^{z} \left(\Delta a_{i}^{\dagger} a_{i} + \frac{1}{2} \epsilon_{i}^{d} a_{i}^{\dagger} + \frac{1}{2} \epsilon_{i}^{d*} a_{i}\right)$$

$$H_a = \sum_{k} (\omega_k - \omega_d) a_k^{\dagger} a_k + (\epsilon_k^d a_k^{\dagger} + \text{h.c.})$$

$$H_d = \sum_{k} (\omega_{\mathrm{c}} - \omega_{\mathrm{d}}) d_k^{\dagger} d_k$$
 $H_{\sigma} = \sum_{k} E_k |k\rangle\langle k| + \left(\frac{g}{\Delta}\right) \left(\epsilon_k^{\mathrm{d}} |k\rangle\langle \mathbf{0}| + \mathrm{h.c.}\right)$

$$H_{\sigma d} = \left(\frac{g}{\Delta}\right)^2 \sum_{i} \sigma_i^z \left[\left(\Delta \bar{a}_i + \frac{1}{2} \epsilon_i^{\mathrm{d}}\right) d_i^{\dagger} + \text{h.c.} \right]$$

$$\partial_t \rho_{\sigma}^{\text{NESS}} = 0 = -\mathrm{i} \left[H_{\sigma}, \rho_{\sigma}^{\text{NESS}} \right] + \sum_k \Gamma_{\mathbf{0} \to k} \mathcal{D}[|\widetilde{k}\rangle\langle \widetilde{\mathbf{0}}|] \rho_{\sigma}^{\text{NESS}}$$

$$+ \gamma \sum_{k} \mathcal{D}[|\widetilde{\mathbf{0}}\rangle\langle \widetilde{k}|] \rho_{\sigma}^{\text{NESS}} + \frac{2\gamma_{\phi}}{N} \sum_{k,q} \mathcal{D}[|\widetilde{q}\rangle\langle \widetilde{k}|] \rho_{\sigma}^{\text{NESS}}.$$

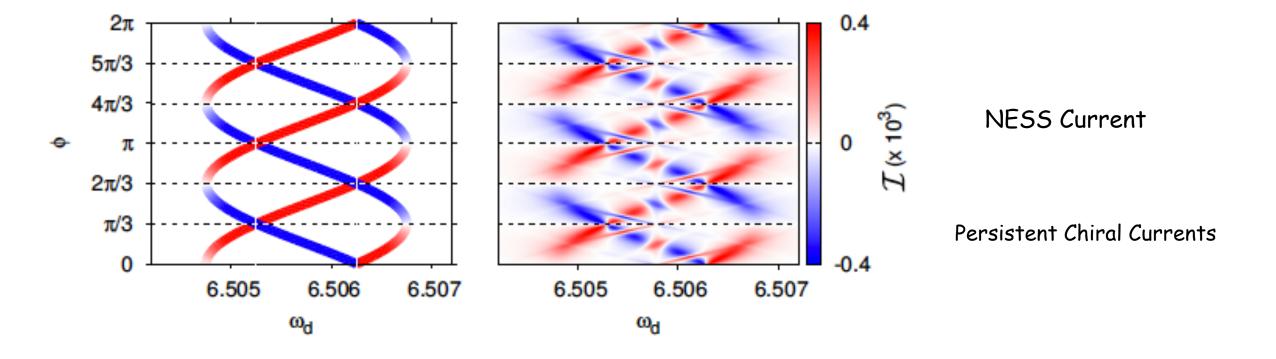
$$\Gamma_{\mathbf{0}\to k} = 2\pi \sum_{\mathbf{q}} \Lambda_{kq}^2 \, \rho(\omega_{\mathrm{d}} + \widetilde{E}_{\mathbf{0}} - \widetilde{E}_{k})$$

$$\omega_{\rm d} = \frac{\omega_{\rm c} + \omega_{\rm q} + \delta\omega_{\rm q}}{2} - J_0 \cos(k + \phi) + \frac{1}{2} \left(\frac{g}{\Delta}\right)^2 \sum_{q \neq k} \frac{|\epsilon_q^{\rm d}|^2}{\Delta_{\rm q}}$$

Once we compute density matrix we can calculate current:

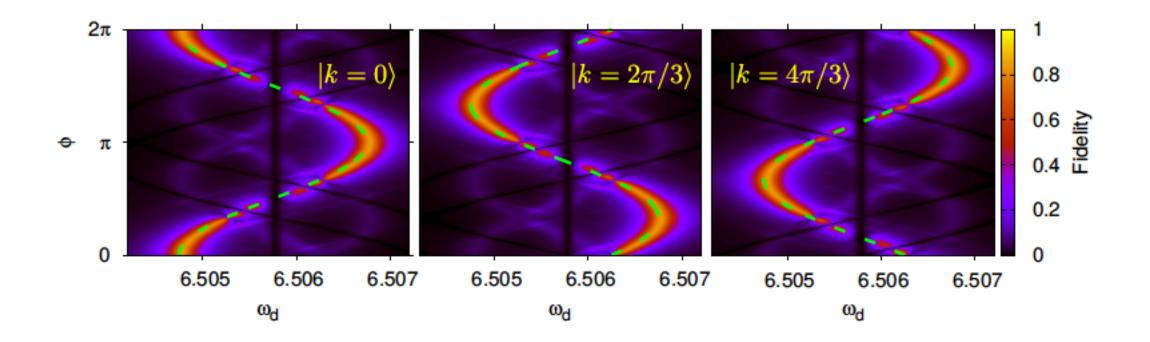
$$\mathcal{J}_i = -iJ_0 \left[e^{i\phi} \sigma_i^+ \sigma_{i+1}^- - \text{H.c.} \right]$$

We also compute the fidelity to go to target eigenstates



$$\mathcal{I}(\omega_{\mathrm{d}}, \phi) = \frac{2}{N} J_0 \sum_{k} \sin(k + \phi) n_k(\omega_{\mathrm{d}}, \phi)$$

Kulkarni, Hein, Kapit, Aron, PRB 97, 064506 (2018)



Steady State Population (Fidelities)

$$\omega_{
m d}^{
m opt}(\phi) = ar{\omega}_{
m d} - J_0 \cos(k + \phi)$$
 Optimal Drive

Conclusions

Kulkarni, Hein, Kapit, Aron, PRB 97, 064506 (2018)

- · Persistent (infinitely long-lived) current
- Preparation of High fidelity states
- Entanglement Dynamics

Outlook

- Extending to ring geometries (many qubits & cavities)
- Performing an exact brute force numerical computation (including spins and photos)
- Treat higher excitation manifolds