BETHE ANSATZ SOLUTION FOR BLOCH ELECTRONS IN MAGNETIC FIELD

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Magnetic field is good for electrons!

BLOCH ELECTRONS IN MAGNETIC FIELD

Lattice electrons in magnetic field

$$H=\sum_{\mathrm{nm}}t_{\mathrm{nm}}c_{\mathrm{n}}^{\dagger}c_{\mathrm{m}}, \quad |t_{\mathrm{nm}}|=1$$

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Magnetic translation

$$H = T_x + T_x^{-1} + T_y + T_y^{-1}$$

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► In Landau gauge:

$$T_x|\mathbf{n}\rangle = |\mathbf{n} + \mathbf{1}_x\rangle, \ T_y = q^2, \quad \psi_{\mathbf{n}} = e^{\mathrm{i}kn_y}\psi_{n_x}(k)$$

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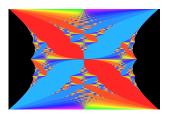
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$$\psi_{n+1} + \psi_{n-1} + 2\cos(k + 2\pi\Phi)\psi_n = E\psi_n$$

▶ One of the most celebrated problem of spectral theory, with applications to localization theory, chaos, quantum Hall effect, etc.

$$\psi_{n+1} + \psi_{n-1} + 2\cos\left(k + 2\pi\Phi\right) \psi_n = E\psi_n$$

- ▶ If the flux is a rational number $\Phi = P/Q$, the spectrum consists of Q bands metal
- If Φ = irrational, the spectrum is a peculiar Cantor set *singular continuous* (a set without isolated points but with zero measure). The total bandwidth is zero



▶ incomplete list of works before 1990:

Zak 1964, Azbel 1964, Hofstadter 1976, Wannier 1978, Aubry-Andre 1980, Zak-Avron 1985, Bellissard-Simon 1980-1990, Thouless-Kohmoto 1982, Wilkinson 1987

HALL CONDUCTANCE=FIRST CHERN CLASS

Hall conductance or the First Chern class is the topological characteristic of the spectrum.

The Hall conductance σ_k of the kth gap is the solution of the Diophantine equation

$$P\sigma_k = k \pmod{Q}$$

Example:
$$\frac{P}{Q} = \frac{4}{15}$$
,

$$\sigma_k = 4, -7, -3, 1, 5, -6, -2, 2, 6, -5, -1, 3, 7, 4$$

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- in 1997 together with Alexander Abanov and J. Talstra we determined the hierarchical structure of the spectrum

HOFSTADTER PROBLEM IS INTEGRABLE! SPECTRUM

$$\psi_{n+1} + \psi_{n-1} + 2\cos\left(k + 2\pi\Phi\right)\psi_n = E\psi_n$$

Equations for the middle of the bands k = 0,

$$q = e^{i\pi\Phi}, \qquad \Phi = \frac{P}{Q}$$

$$E = 2 (-1)^{p} \sin(\pi \Phi) \sum_{l=1}^{Q-1} z_{l}$$

Roots $z_1, \dots z_{Q-1}$ obeys the Bethe Ansatz equations

$$\frac{z_l^2 + q}{q z_l^2 + 1} = (-1)^p \prod_{m \neq l}^{Q-1} \frac{q z_l - z_m}{z_l - q z_m}$$

HOFSTADTER PROBLEM IS INTEGRABLE! WAVE FUNCTIONS

Polynomials

$$\Psi(z) = \prod_{l=1}^{Q-1} (z-z_l)$$

play a special role. Some of them have names: $\Psi_{E=0}(z)$ - is q-Legendre polynomial.

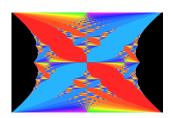
The wave function

$$\psi_n = \sum_{m=1}^{Q-1} c_{nm} \Psi(z) \Big|_{z=q^m}$$

The coefficients are quantum di-logarithm

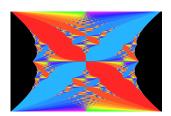
$$c_{nm} = q^{2nm + \frac{m}{2}} \prod_{j=0}^{m-1} \frac{1 + q^{-j - \frac{1}{2}}}{1 - q^{j + \frac{1}{2}}}$$

$$\frac{z_l^2 + q}{qz_l^2 + 1} = (-1)^p \prod_{m \neq l}^{Q-1} \frac{qz_l - z_m}{z_l - qz_m}$$



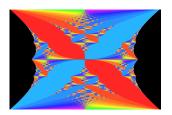
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√ How to obtain these equations?



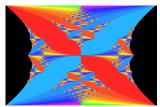
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- ▶ How to solve them? We are interested in $Q \to \infty$, $P \to \infty$, $\Phi = P/Q$ fixed



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- ✓ How to obtain these equations?
- ▶ How to solve them? We are interested in $Q \to \infty$, $P \to \infty$, $\Phi = P/Q$ fixed
- ▶ What does the limit mean in such complicated spectrum?



CHIRAL GAUGE

► Chiral gauge $t_{\mathbf{n},\mathbf{n}+\mathbf{1}_{\mathbf{x}}} = e^{-i\frac{\Phi}{2}n_{+}}, \quad t_{\mathbf{n},\mathbf{n}+\mathbf{1}_{\mathbf{y}}} = e^{+i\frac{\Phi}{2}(n_{+}+1)}, \quad n_{+} = n_{x} + n_{y}, \quad \Psi_{\mathbf{n}} = e^{i\mathbf{p}\mathbf{n}}\Psi_{n_{+}}$ $iq^{-1/2}(1+q^{2n+1})\Psi_{n+1} - iq^{1/2}(1+q^{-2n+1})\Psi_{n-1} = E\Psi_{n}, \quad \mathbf{p} = (\frac{\pi}{2}, \frac{\pi}{2})$

▶ What is the advantage this gauge?

Consider a difference equation, such that $\Psi_n = \Psi(z)\Big|_{z=q^n}$

$$i(z^{-1} + qz)\Psi(qz) - i(z^{-1} + q^{-1}z)\Psi(q^{-1}z) = E\Psi(z)$$

▶ All solutions of the difference equation are polynomials

$$\Psi(z) = \prod_{l=1}^{Q-1} (z - z_l)$$

► Equation for the roots

$$\frac{z_l^2 + q}{qz_l^2 + 1} = (-1)^p \prod_{m \neq l}^{Q-1} \frac{qz_l - z_m}{z_l - qz_m}$$



SPECTRAL ALGEBRAIZATION AND REPRESENTATION THEORY

▶ Q: When all solutions of the second order ODE are polynomials?

$$H\Psi = \left[a(z)\frac{d^2}{dz^2} + b(z)\frac{d}{dz} + c(z)\right]\Psi(z) = E\Psi(z)$$

A: If the operator is equivalent to the Euler top

$$H = \sum_{i,j=1,2,3} \alpha_{ij} S_i S_j + \sum_{i=1,2,3} \beta_i S_i$$

where

$$S_3 = z \frac{d}{dz} - j, \ S_+ = z \left(2j - z \frac{d}{dz} \right), \ S_- = \frac{d}{dz}$$

are finite dimension representation of SU(2)

HOFSTADER PROBLEM AND REPRESENTATION THEORY

▶ When all solutions of the difference equation

$$a(z)\Psi(q^2z) + d(z)\Psi(q^{-2}z) + v(z)\Psi(z) = E\Psi(z)$$

are polynomials $\Psi(z) = \prod_{i} (z - z_i)$?

Setting $z = q^n$ we obtain solvable discrete equation

$$a_n \psi_{n+1} + d_n \psi_{n-1} + \mathbf{v}_n \psi_n = E \psi_n$$

► Lie group \rightarrow quantum group $SL(2) \rightarrow SL_a(2)$

$$SL(2) \rightarrow SL_q(2)$$

$$\{{\bf 1},\,S_x,\,S_y,\,S_z\} \to \{A,\,B,\,C,\,D\}$$

$$[S_a, S_b] = i\epsilon^{abc}S_c$$

$$AB = qBA, BD = qDB,$$

$$DC = qCD, CA = qAC,$$

$$AD = 1, [B, C] = \frac{A^2 - D^2}{a - a^{-1}}$$

$SL_q(2)$ QUANTUM VERSION OF SL(2)

▶ Universal *R*-matrix, obeying Yang-Baxter equation

$$R(u) = \begin{bmatrix} \frac{uA - u^{-1}D}{q - q^{-1}} & C\\ B & \frac{uD - u^{-1}A}{q - q^{-1}} \end{bmatrix}$$

Magnetic translations embedded into $SL_q(2)$

Hamiltonian happens to be equal

$$H = \sum_{\mathbf{nm}} t_{\mathbf{nm}} c_{\mathbf{n}}^{\dagger} c_{\mathbf{m}} = \text{Tr}[\sigma_1 R] = B + C$$

- ▶ How did we find this?
- Embedding

$$T_{n}T_{m} = q^{-n \times m}T_{n+m}$$

$$T_{-x} + T_{-y} = B, \qquad T_{x} + T_{y} = C,$$

$$T_{-y}T_{x} = q^{-1}A^{2}, \qquad T_{-x}T_{y} = qD^{2}$$
 (1)

STRINGS

▶ How to solve the Bethe Ansatz equations?

$$\frac{z_l^2 + q}{qz_l^2 + 1} = (-1)^p \prod_{m \neq l}^{Q-1} \frac{qz_l - z_m}{z_l - qz_m}$$

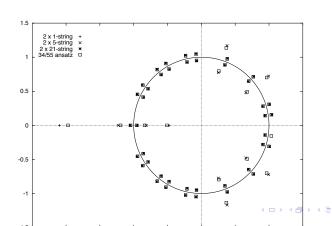
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▶ At large *Q* solutions consist of collections of *strings*

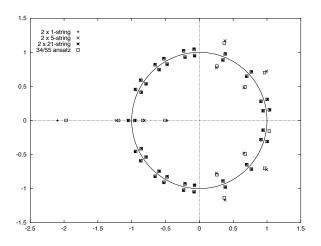
A string of spin l centered at x_l is a set roots of unity $z^{(l)} = x_l \times \{e^{i\pi k/l}\}, \quad k = 1, ..., l$



STRING AND HALL CONDUCTANCE

▶ The length of the longest string of a given band is the **Hall conductance** of the band:

$$(2l+1)_{\text{max}} = |\text{Hall conductance}|$$



- ▶ Each solution is labeled by a content of strings $\{l_i, l_{i-1}, ...\}$
- ▶ The length of the longest string of a given band is the **Hall conductance** of the band:

$$(2l+1)_{max} = |Hall conductance|$$

▶ The remaining roots of the state is a solution of the Bethe equation for the parent state

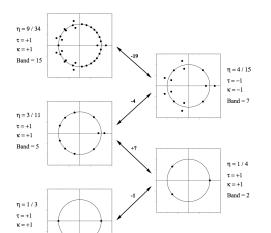
$$\Psi^{\text{daugther}}(z) \approx \prod_{m=-l}^{l} (z - x_l q_l^m) \Psi^{\text{parent}}(z)$$

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Fibonacci Tree: Example

- ▶ Golden mean $\Phi = \frac{\sqrt{5}-1}{2}$
- The sequence of rational approximants is given by ratios of subsequent Fibonacci numbers

$$\Phi_i = \frac{F_{i-1}}{F_i}$$
: $F_i = F_{i-2} + F_{i-1} = 1, 1, 2, 3, 5, 8, 13, 21, 34, 55, 89, \dots$

▶ The set of Hall conductances (lengths of strings) are again Fibonacci numbers: F_{k-1} . The wave function of this state is

$$\Psi\bigg(z|\Phi_k = \frac{F_{3k-1}}{F_{3k}}\bigg) \approx \prod_{n=0}^{k-1} \prod_{i=-\frac{1}{2}(F_{2n}-1)}^{\frac{1}{2}(F_{3n}-1)} \bigg(z - e^{i\pi\frac{F_{3n-1}}{F_{3n}}j}\bigg)^2$$

SCALING HYPOTHESIS (ABANOV, TALSTRA, P.W.)

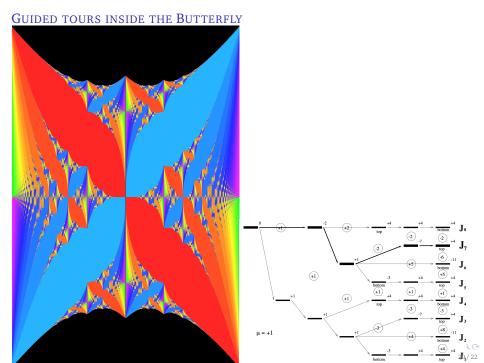
- ✓ A gap width scales along generations as Q^{-2} (proven)
- ▶ Mid band energy scales along generations

$$|E_{\text{daughter band}} - E_{\text{parent band}}| \sim Q^{-2+\Delta_{\text{band}}}$$

 Δ_{band} – scaling dimension - function of flux and chemical potential

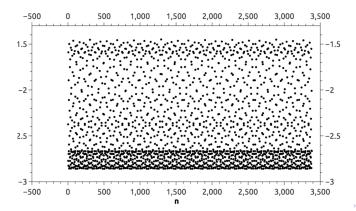
$$\eta = \frac{4}{15} : \begin{cases} \sigma(k) = 4, -11, 4, 4, 4, -11, 4, 4, 4, -11, 4, 4, 4, -11, 4 \\ \sigma_k = 4, -7, -3, 1, 5, -6, -2, 2, 6, -5, -1, 3, 7, 4 \end{cases}$$

$$\eta_6 = \frac{3}{11} : \begin{cases} \sigma(k) = 4, -7, 4,$$



SCALING DIMENSIONS

- ► Once we know the dimensions, we know physics of singular continuum state of matter e.g., a.c. conductance, specific heat, etc. of
- ➤ The main problem: what are the anomalous dimensions remains unsolved!



V. NABOKOV "GIFT" CHAPTER 4

Truth bends her head to fingers curved cupwise; And with a smile and care Examines something she is holding there Concealed by her from our eyes.