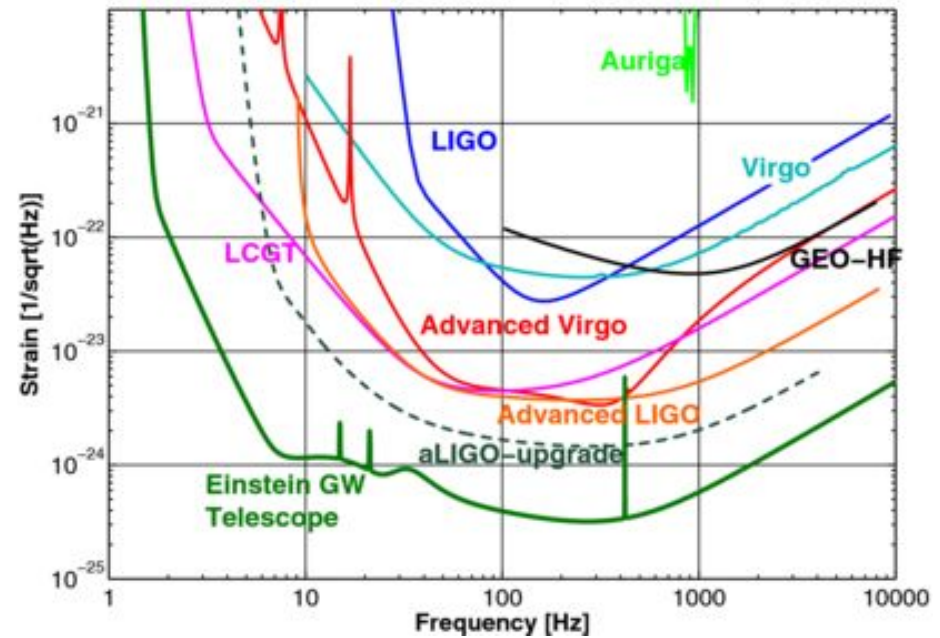
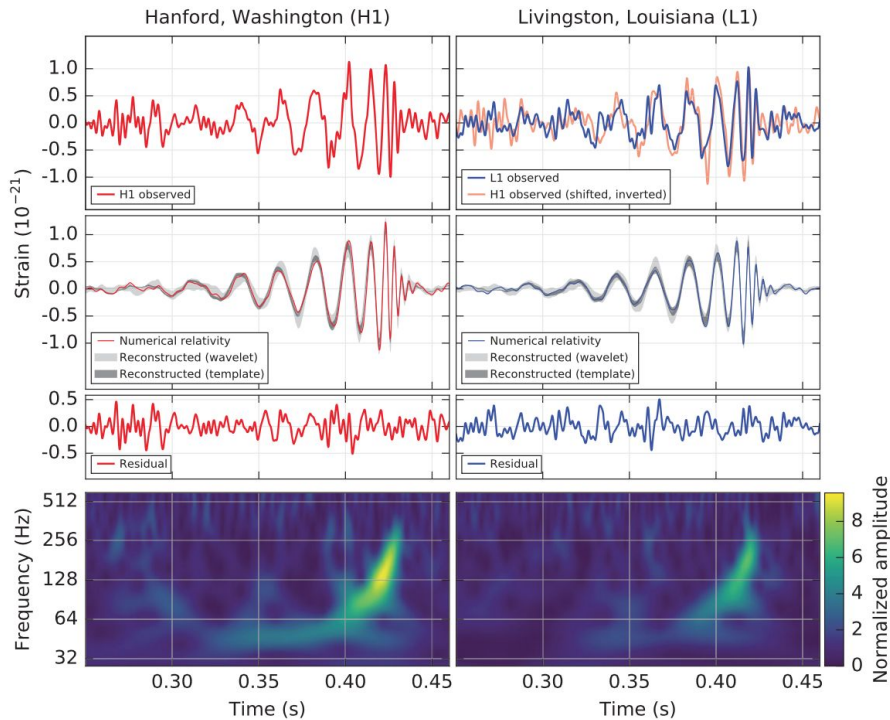


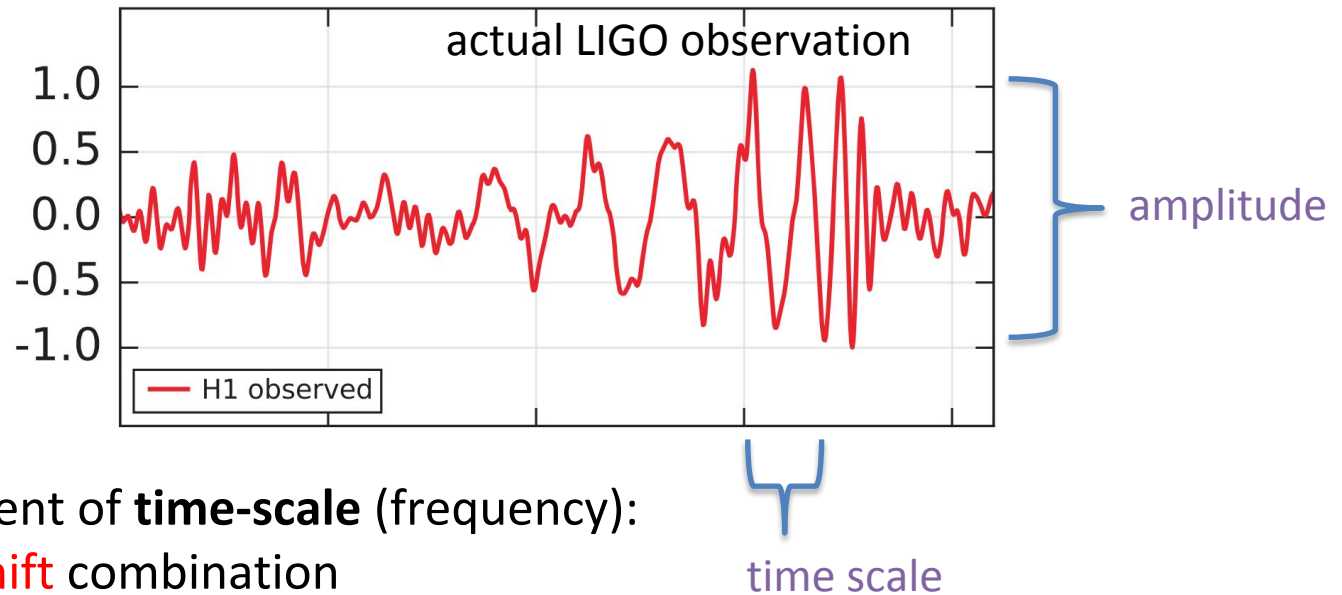
Gravitational waves from binary black holes



from ET conceptual design paper

Potential to detect a lot of massive binary (BH) mergers from far away

Gravitational wave observables



- ◆ Measurement of **time-scale** (frequency):
mass-redshift combination

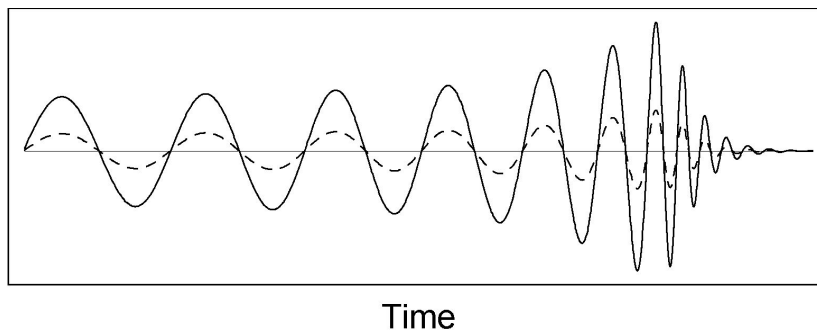
$$M(1 + z)$$

- ◆ Measurement of **strain amplitude**: (multiple detectors for orientation)
Luminosity distance-magnification combination

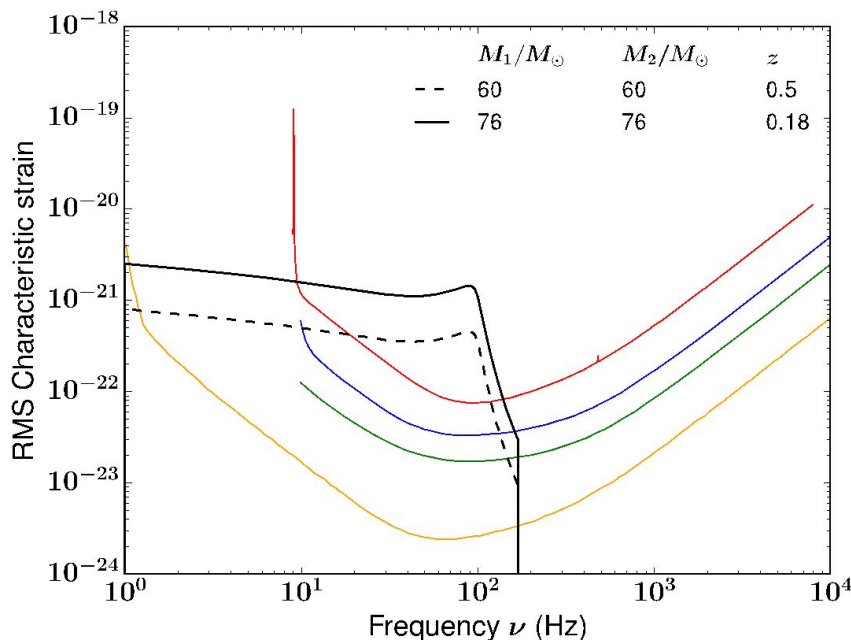
$$\frac{\sqrt{\mu}}{d_L(z)} \left(\frac{dE}{df_s} \right)^{1/2}_{M, f_s = f_o(1+z)}$$

- ◆ **Detailed shape**: mass ratio, spin, etc.

Gravitational lensing and GW parameter estimation



Consider an event with physical mass scale M , source redshift z , lensing magnification μ



Assume ignorance about magnification: **Inferred** mass scale \tilde{M} , **inferred** source redshift \tilde{z}

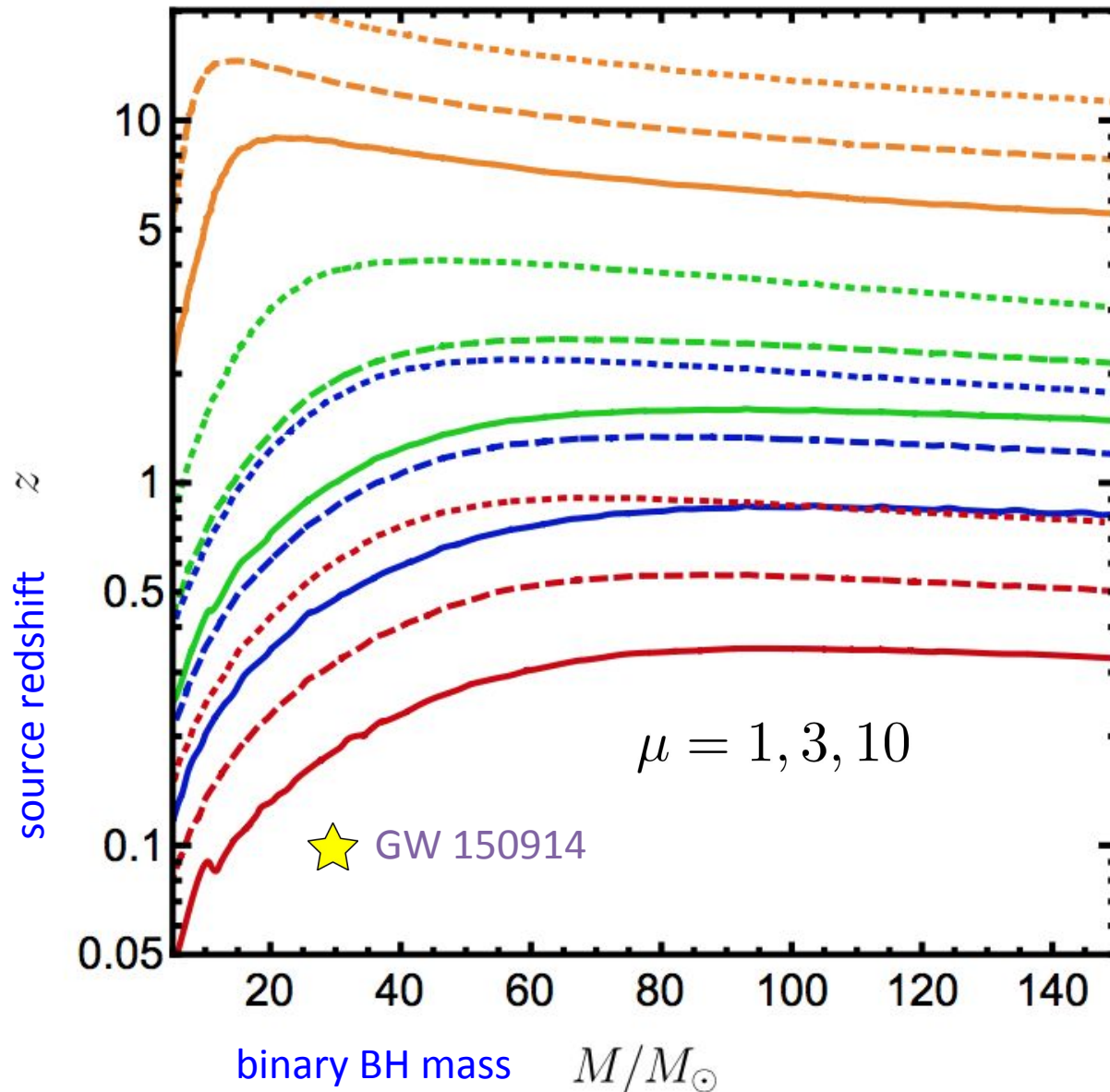
$$M(1+z) = \tilde{M}(1+\tilde{z})$$

$$d_L(\tilde{z}) = d_L(z)/\sqrt{\mu}$$

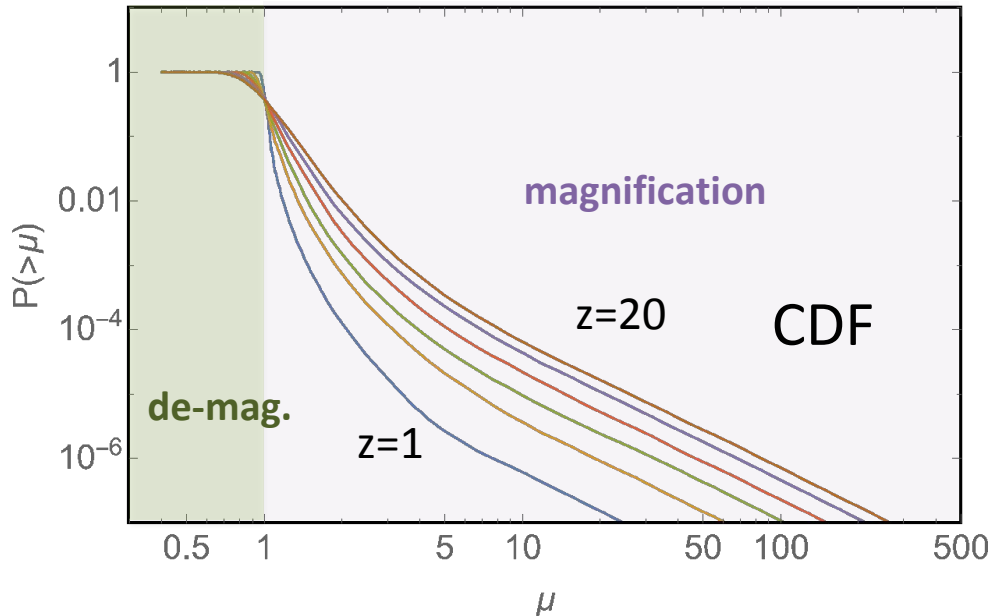
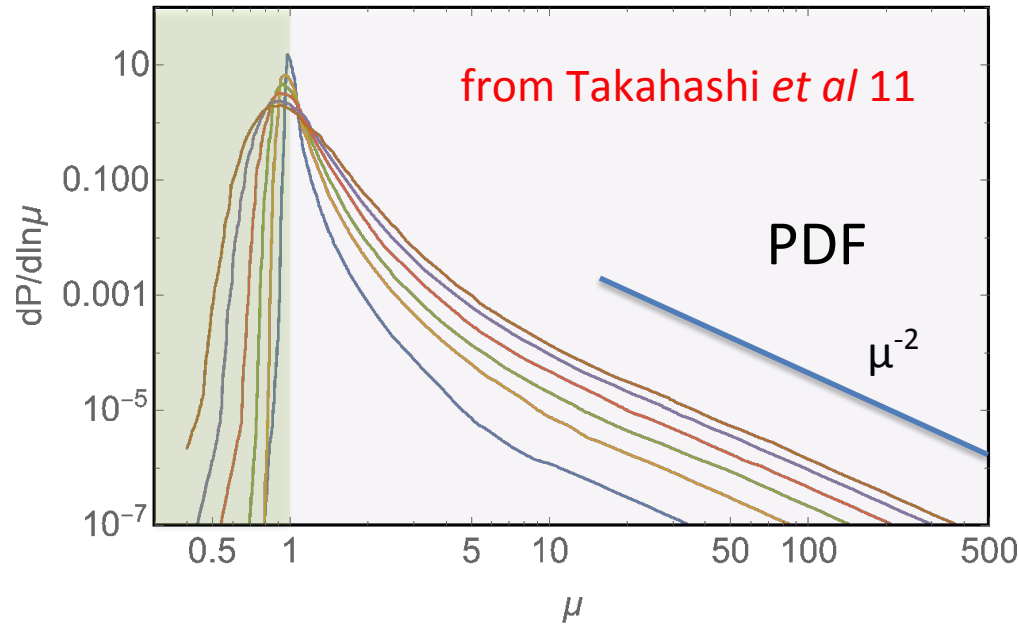
Availability of redshift information vs. lensing

	w/o lensing (de-)magnification	w/ lensing (de-)magnification
w/ EM counterpart or host identification	Independently measure M , z , d_L ; standard sirens for measuring $d_L(z)$ (Holz & Hughes 05)	assuming fiducial cosmology $d_L(z)$; measure M , z and magnification μ
w/o EM counterpart or host identification (probably BH-BH merger)	assuming fiducial cosmology $d_L(z)$; still able to fix M , z (e.g. GW 150914)	Cannot uniquely determine M , z , μ DEGENERACY

High magnification tail



Lensing probability in Λ CDM



normalization

$$\int d \ln \mu \frac{dP(\mu, z)}{d \ln \mu} = 1$$

conservation of solid angle

$$\langle \mu \rangle = \int d \ln \mu \frac{dP(\mu, z)}{d \ln \mu} \mu = 1$$

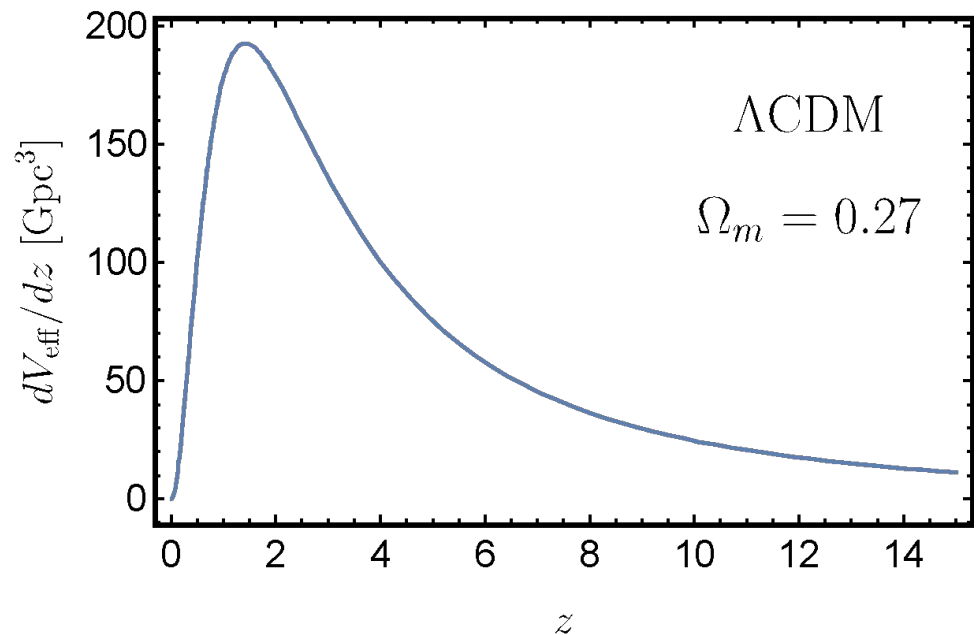
- Strong lensing dominated by galaxies; cluster lensing subdominant. Li & Ostriker 03
- Stellar lens population: diffraction regime. Takahashi 03

Observed merger distribution

Only **inferred mass and redshift** are observable

$$\underbrace{\frac{d^3 N(\tilde{M}, \tilde{z})}{d\tilde{M} d\tilde{z} dt}}_{\text{Observed rate in terms of } \tilde{M}' \text{ and } \tilde{z}'} = \int d \ln \mu \underbrace{\frac{dP(\mu, z)}{d \ln \mu}}_{\text{lensing probability}} \underbrace{\frac{d^3 N(M, z)}{dM dz dt}}_{\text{unlensed rate in terms of } M \text{ and } z} \underbrace{\left[\frac{\partial(M, z)}{\partial(\tilde{M}, \tilde{z})} \right]_{\mu}}_{\text{mapping Jacobian}}$$

$$\frac{d^3 N(M, z)}{dM dz dt} = \underbrace{\frac{d^2 n(M, z)}{dM dt_s}}_{\text{intrinsic rate astrophysics models}} \underbrace{\frac{4\pi c \chi^2(z)}{(1+z) H_0 E(z)}}_{\text{effective volume}}$$

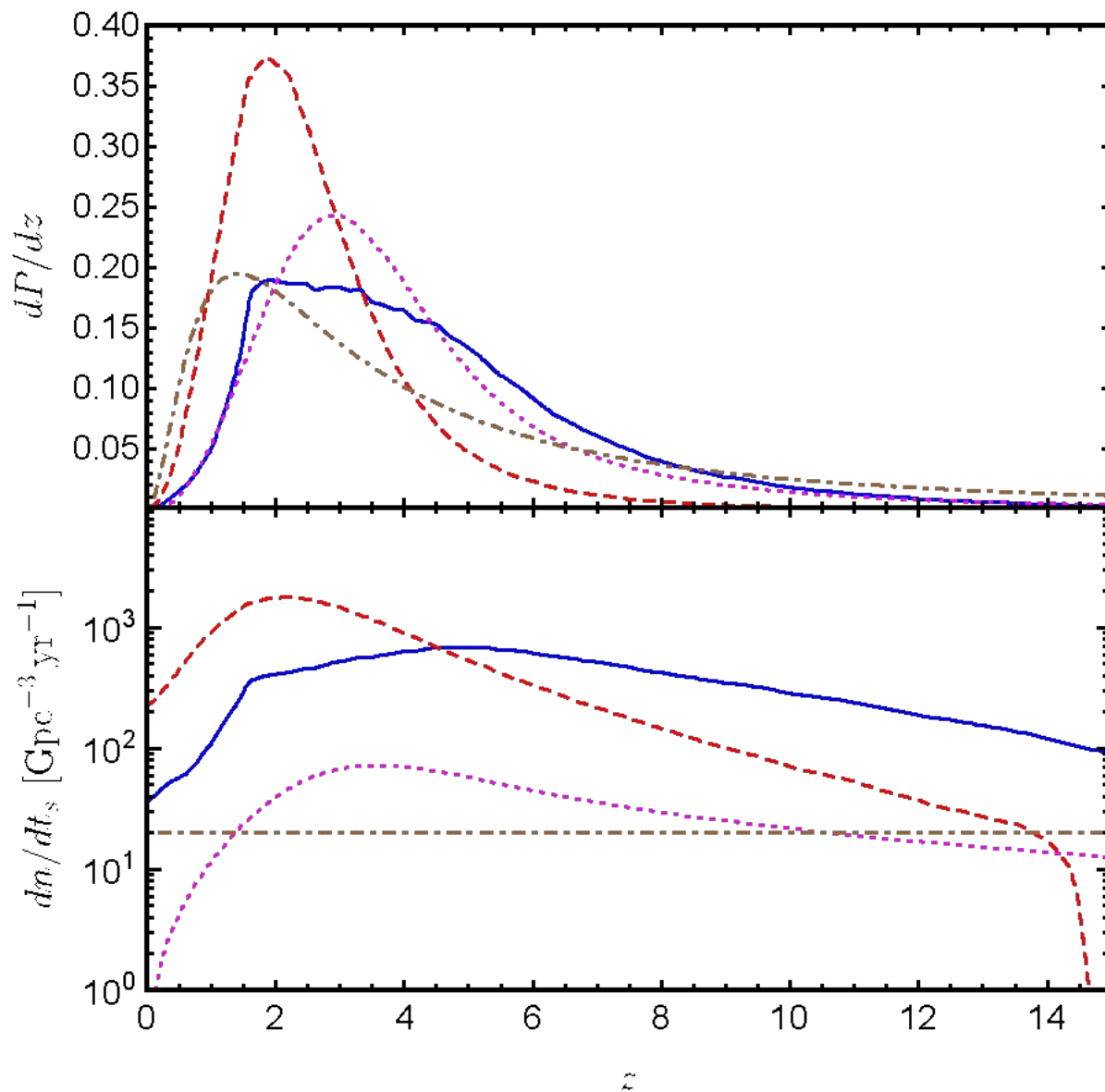


Amplification bias

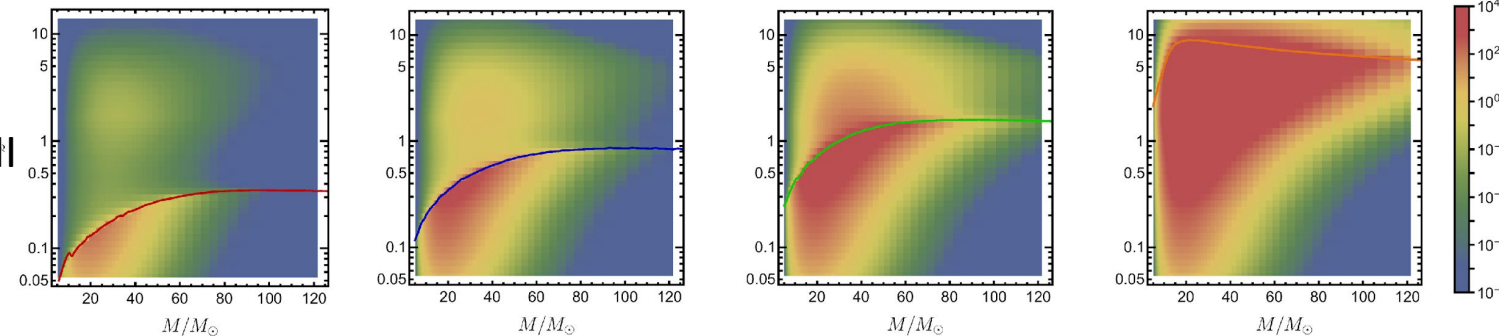
$$\begin{aligned} \frac{dN_{\mathcal{D}}}{dt} \left(\tilde{M} > M_{\min} \right) &= \int_{M_{\min}}^{+\infty} d\tilde{M} \int d\tilde{z} \Theta \left(\overbrace{\mathcal{S}_{\mathcal{D}}(\tilde{M}, \tilde{z}) - \mathcal{S}_0}^{\text{Sensitivity cut}} \right) \\ &\times \int d\ln \mu \frac{dP(\mu; z)}{d\ln \mu} \frac{d^3 N(M, z)}{dM dz dt} \left| \frac{\partial(M, z)}{\partial(\tilde{M}, \tilde{z})} \right|_{\mu}. \end{aligned}$$

- Detector-dependent
- Increases the strongly lensed fraction for given parameters

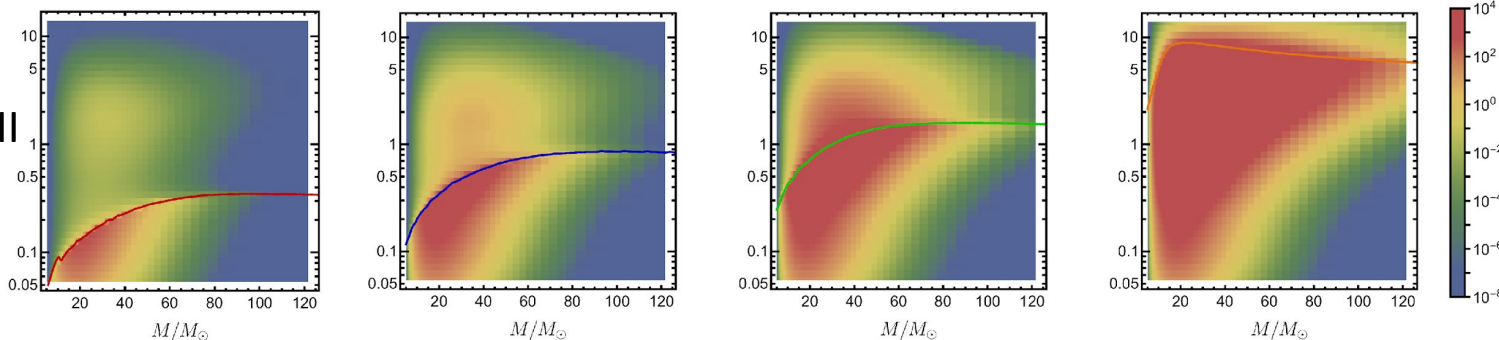
Origin of massive BH binaries?



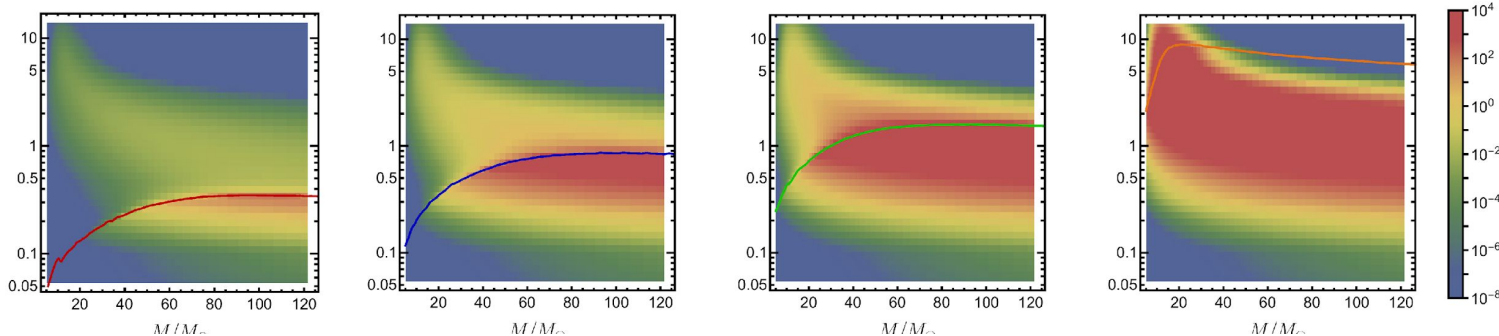
Pop I and II



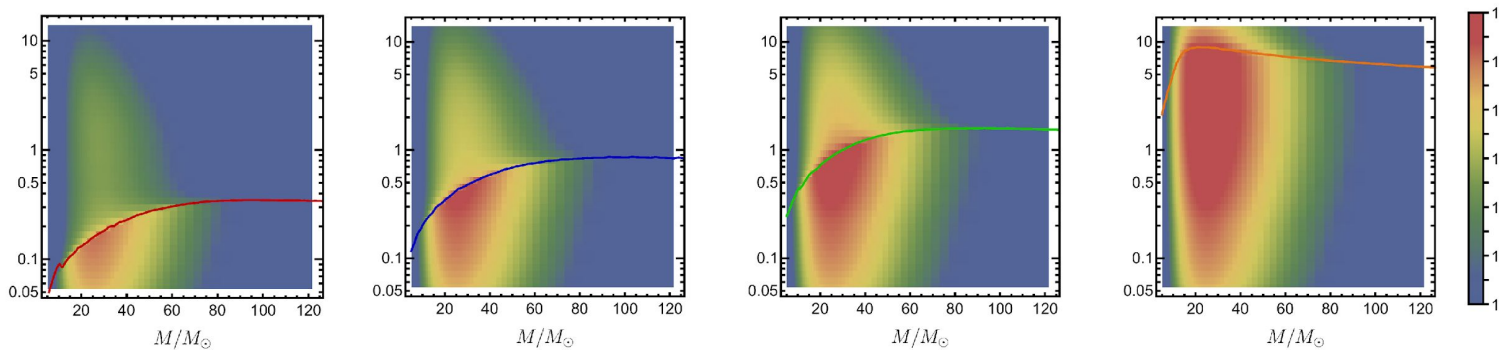
Pop I and II
(new)

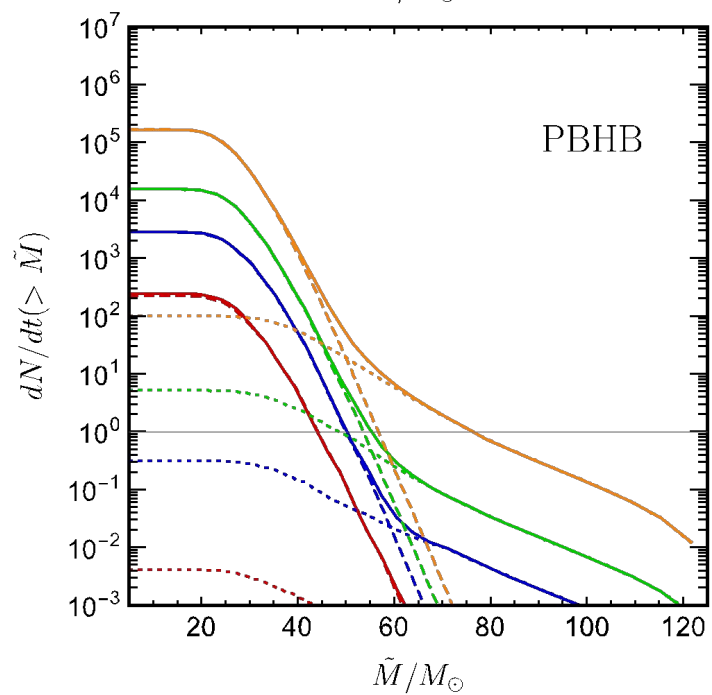
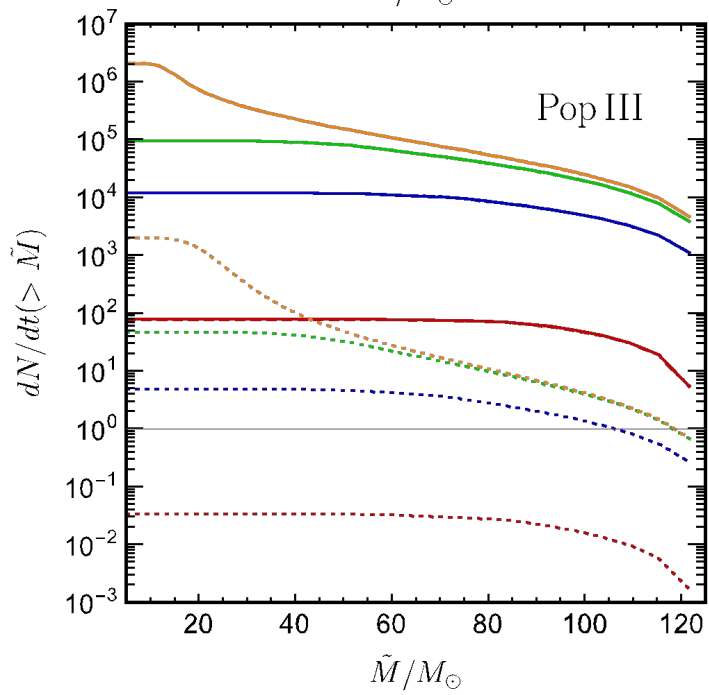
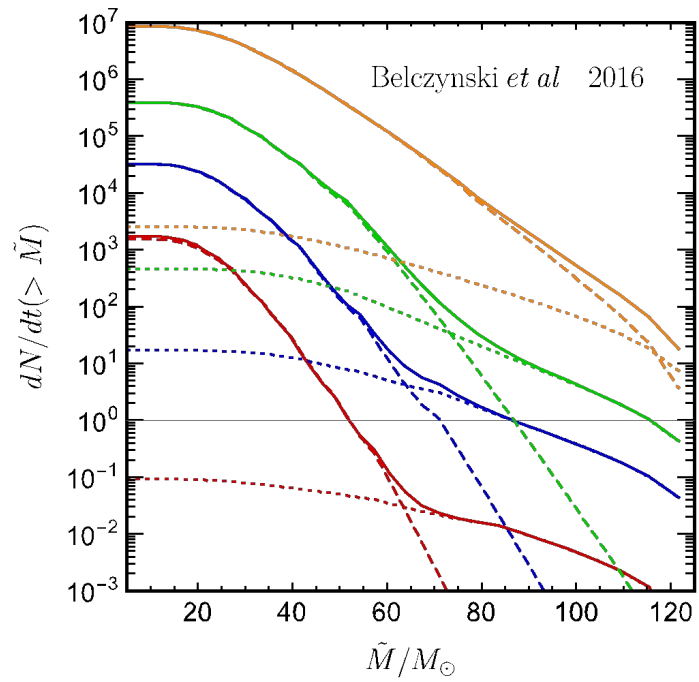
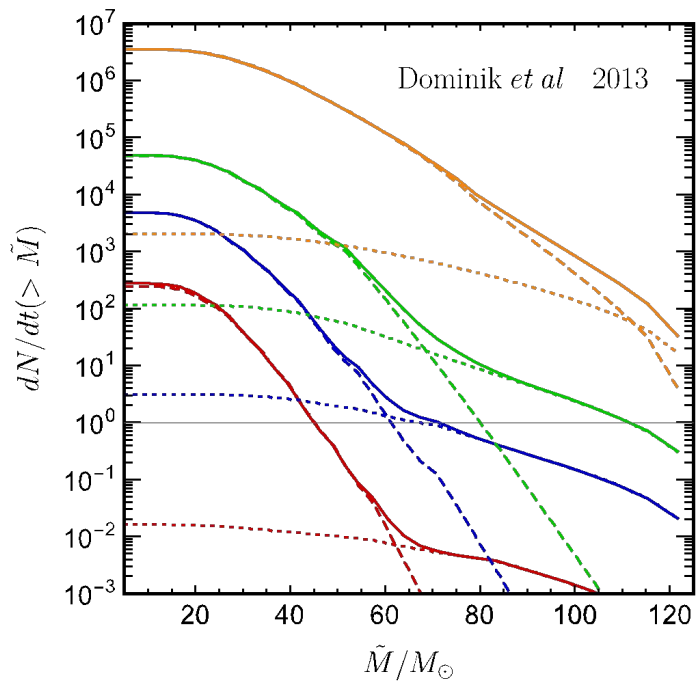


Pop III



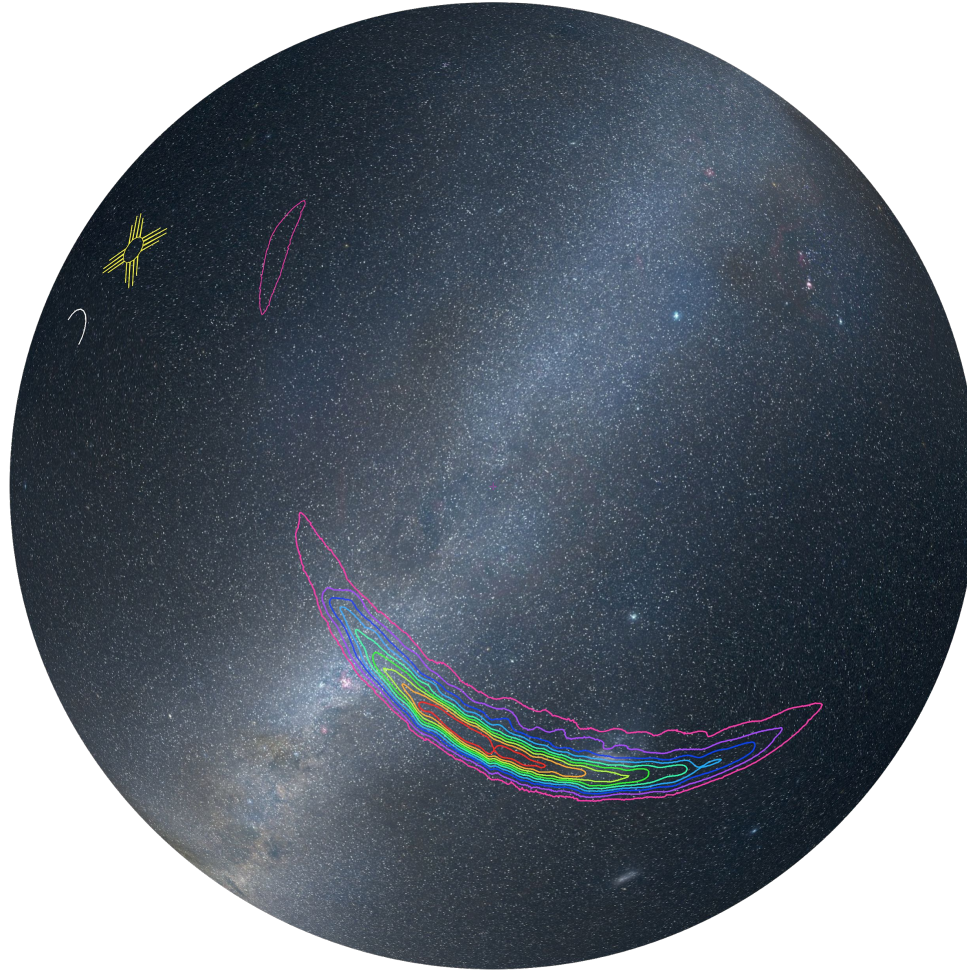
PBH





What can we learn from multiple images?

With two detectors, localization is a function of orientation.
Multiple orientations -> improved location.



GW150914 Localization

See also Seto (14)