

N-body equations of motion:

Worked example: 2 bodies and the perihelion shift

Components of the disturbing force

$$\mathcal{R} = \frac{Gm}{c^2 r^2} \left[-(1 + 3\eta)v^2 + \frac{1}{2}(8 - \eta)\dot{r}^2 + 2(2 + \eta)\frac{Gm}{r} \right],$$

$$\mathcal{S} = \frac{Gm}{c^2 r^2} \left[2(2 - \eta)\dot{r}(r\dot{\phi}) \right],$$

$$\mathcal{W} = 0$$

Integrate the Lagrange planetary equations:

$$\Delta e = \Delta a = 0$$

$$\Delta \Omega = \Delta \iota = 0$$

$$\Delta \omega = \frac{6\pi G(M_1 + M_2)}{a(1 - e^2)c^2}$$

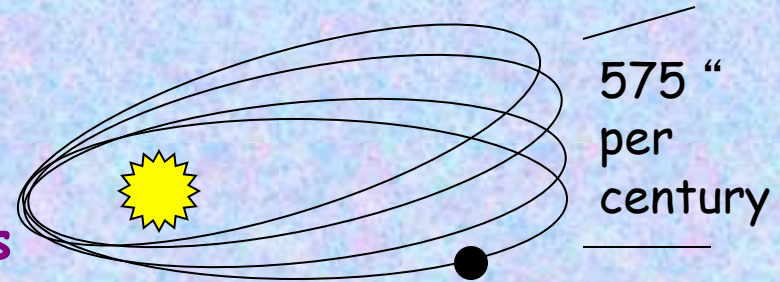
42.98 "/c for
Mercury

4.226598 °/yr
for PSR 1913+16



Mercury's Perihelion: Trouble to Triumph

- 1687 Newtonian triumph
- 1859 Leverrier's conundrum
- 1900 A turn-of-the century crisis



Planet	Advance
Venus	277.8
Earth	90.0
Mars	2.5
Jupiter	153.6
Saturn	7.3
Total	531.2
Discrepancy	42.9
Modern measured value	42.98 ± 0.001
General relativity prediction	42.98



Deflection of light

$$\mathbf{v} = c \left(1 - \frac{2}{c^2} U \right) \mathbf{n} \quad \frac{dn^j}{dt} = \frac{2}{c} (\delta^{jk} - n^j n^k) \partial_k U$$

$$\mathbf{r}(t) = \mathbf{r}_e + c\mathbf{k}(t - t_e) + O(c^{-2})$$

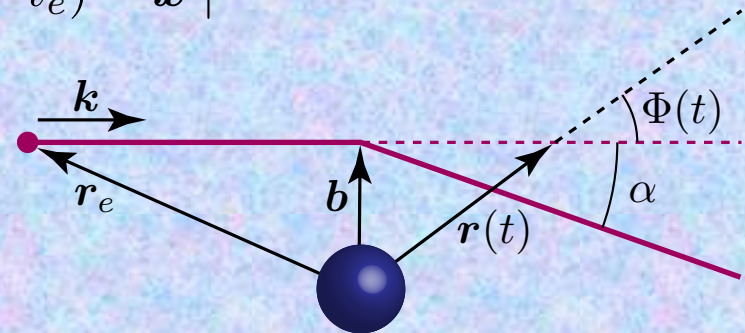
$$\mathbf{n} = \mathbf{k} + \boldsymbol{\alpha} + O(c^{-4})$$

$$\mathbf{s}_e = \mathbf{r}_e - \mathbf{x}'$$

$$\mathbf{s} = \mathbf{r}(t) - \mathbf{x}'$$

$$\mathbf{b} = \mathbf{s}_e - \mathbf{k}(\mathbf{s}_e \cdot \mathbf{k})$$

$$\begin{aligned} \frac{d\alpha^j}{dt} &= -\frac{2G}{c} (\delta^{jk} - k^j k^k) \int \rho' \frac{r_e^k + ck^k(t - t_e) - x'^k}{|\mathbf{r}_e + c\mathbf{k}(t - t_e) - \mathbf{x}'|^3} d^3x' \\ &= -\frac{2G}{c} \int \rho' \frac{\mathbf{b}}{s^3} d^3x' \\ &= -\frac{2G}{c^2} \int \rho' \frac{\mathbf{b}}{b^2} \frac{d}{dt} \left(\frac{\mathbf{s} \cdot \mathbf{k}}{s} \right) d^3x' \end{aligned}$$



$$\boldsymbol{\alpha}(t) = -\frac{2G}{c^2} \int \rho(t, \mathbf{x}') \frac{\mathbf{b}}{b^2} \left(\frac{\mathbf{s} \cdot \mathbf{k}}{s} - \frac{\mathbf{s}_e \cdot \mathbf{k}}{s_e} \right) d^3x'$$



Deflection of light

For a point mass at the origin

$$\alpha(t) = -\frac{2GM}{c^2} \frac{b}{b^2} \left[\frac{\mathbf{r}(t) \cdot \mathbf{k}}{r(t)} - \frac{\mathbf{r}_e \cdot \mathbf{k}}{r_e} \right]$$

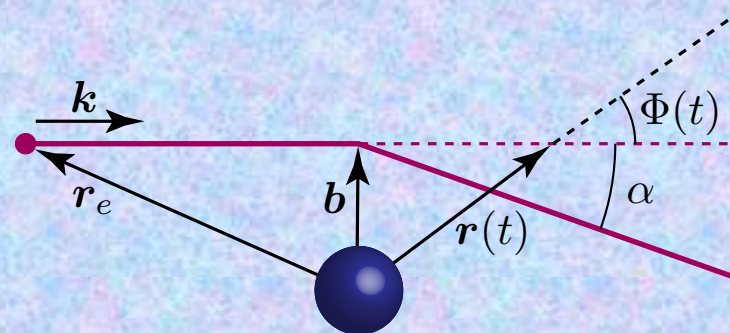
$$\approx -\frac{4GM}{c^2 b} \hat{\mathbf{b}} \left[\frac{\cos \Phi(t) + 1}{2} \right]$$

For $\Phi \approx 0$

$$\Delta\theta \approx |\alpha|$$

$$= \frac{4GM}{c^2 b}$$

$$= 1''.7505 \left(\frac{M/M_\odot}{b/R_\odot} \right)$$



Deflection of light: The 1919 Eclipse



A. S. Eddington

LIGHTS ALL ASKEW IN THE HEAVENS

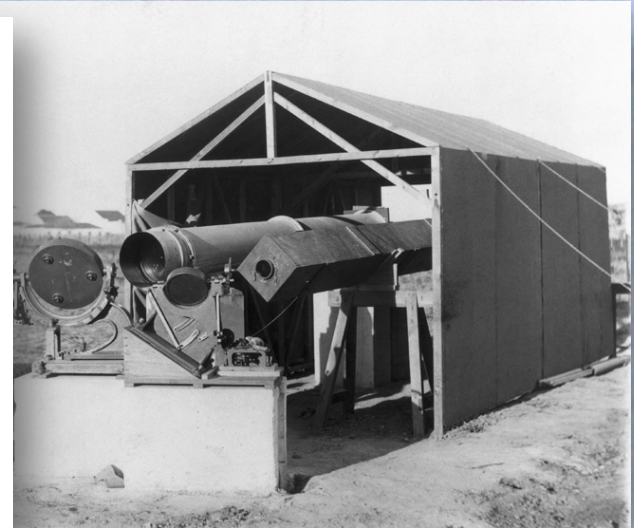
Men of Science More or Less
Agog Over Results of Eclipse
Observations.

EINSTEIN THEORY TRIUMPHS

Stars Not Where They Seemed
or Were Calculated to be,
but Nobody Need Worry.

A BOOK FOR 12 WISE MEN

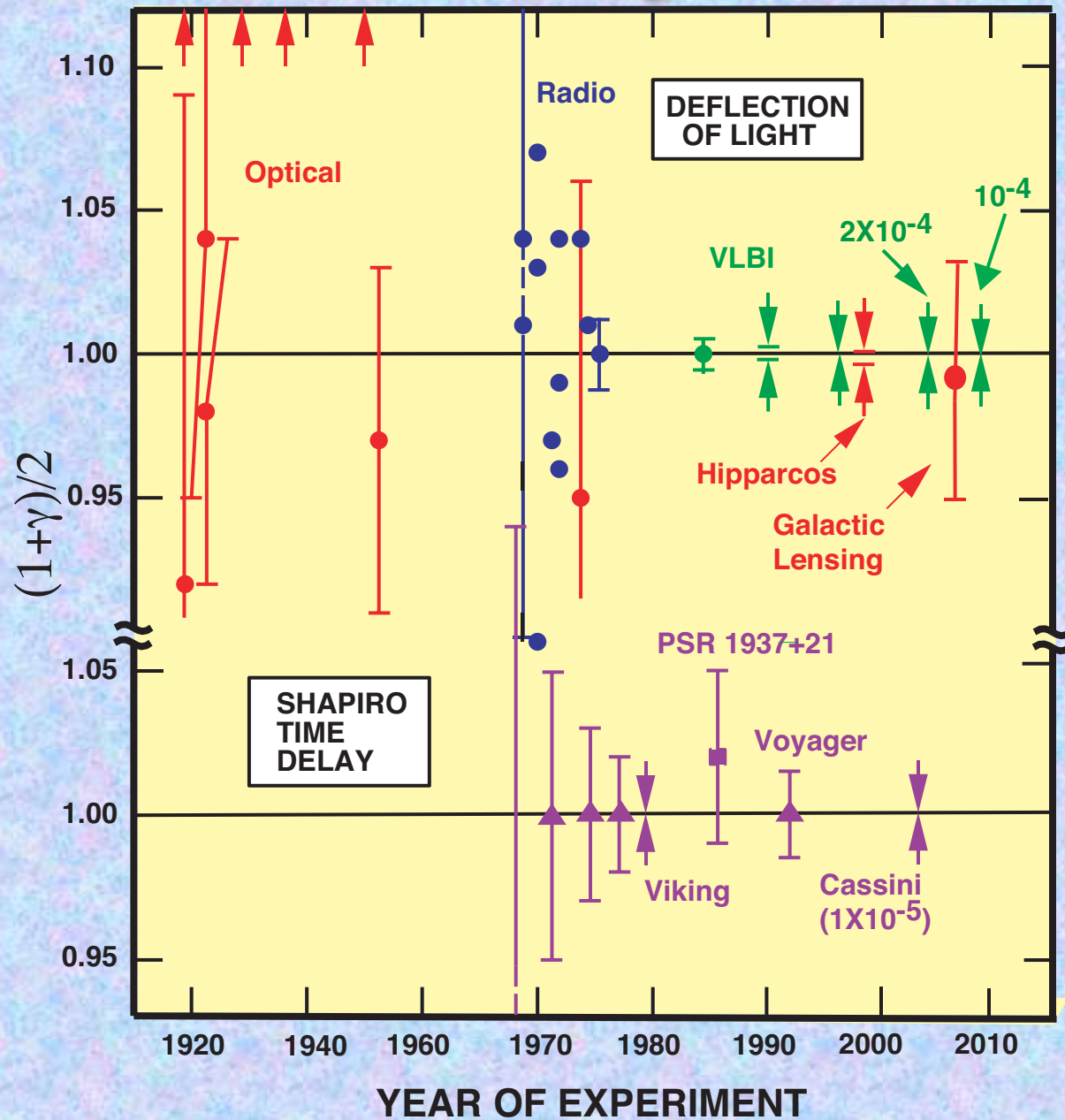
No More in All the World Could
Comprehend It, Said Einstein When
His Daring Publishers Accepted It.



Sobral site

Photo from Principe

Deflection of light: Results



Gravitational lenses

$$\begin{aligned}\alpha(t) &= -\frac{2G}{c^2} \int \rho(t, \mathbf{x}') \frac{b}{b^2} \left(\frac{\mathbf{s} \cdot \mathbf{k}}{s} - \frac{\mathbf{s}_e \cdot \mathbf{k}}{s_e} \right) d^3 x' \\ &= -\frac{4G}{c^2} \int \rho(t, \mathbf{x}') \frac{b}{b^2} d^3 x' \\ &= -\frac{4G}{c^2} \int \Sigma(t, \boldsymbol{\xi}') \frac{\boldsymbol{\xi} - \boldsymbol{\xi}'}{|\boldsymbol{\xi} - \boldsymbol{\xi}'|^2} d^2 \xi'\end{aligned}$$

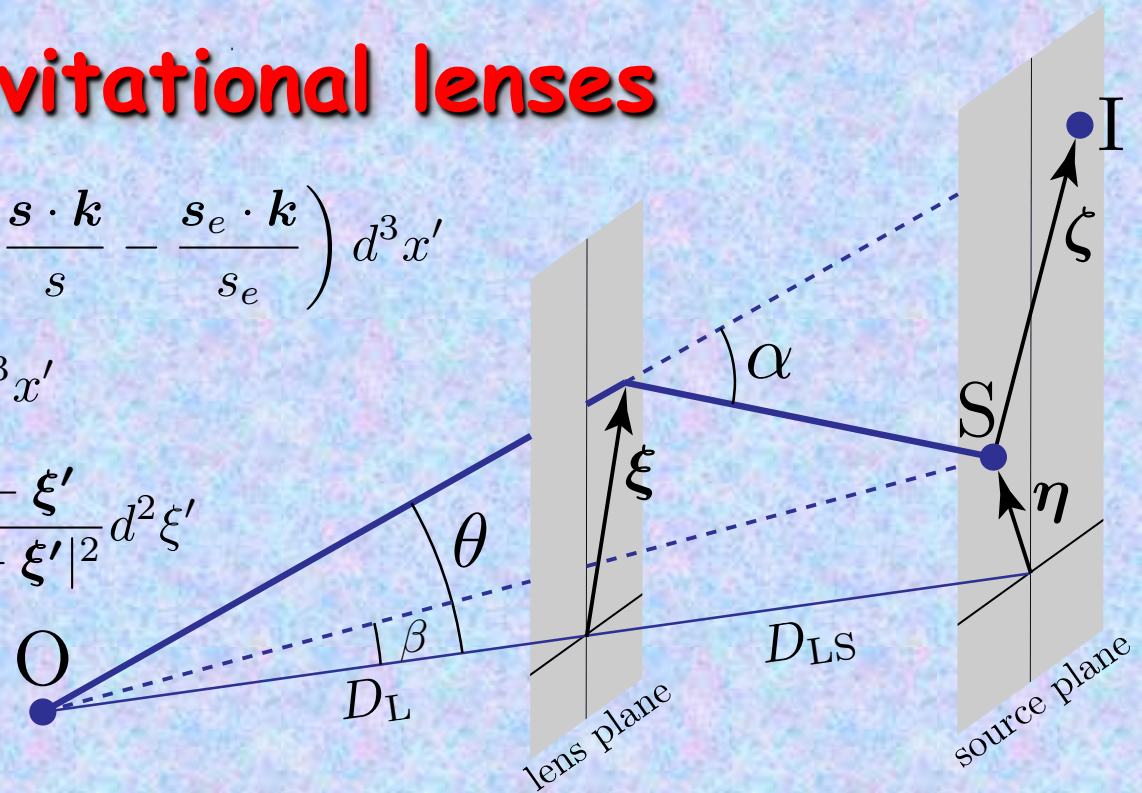
Lens equation

$$\theta + \frac{D_{LS}}{D_S} \alpha = \beta$$

$$\boldsymbol{\xi} = D_L \boldsymbol{\theta}$$

Schwarzschild lens

$$\alpha(\boldsymbol{\xi}) = -\frac{4GM}{c^2} \frac{\boldsymbol{\xi}}{\xi^2}$$



$$D_S \theta = \eta + \zeta = D_S \beta - D_{LS} \alpha$$

$$\theta - \frac{\theta_E^2}{\theta} = \beta$$

$$\theta_E^2 := \frac{4GM}{c^2} \frac{D_{LS}}{D_S D_L}$$



Gravitational lenses

$$\theta_{\pm} = \frac{1}{2} \left(\beta \pm \sqrt{\beta^2 + 4\theta_E^2} \right)$$

$$\begin{aligned} \theta_E &\sim 2 \text{ as} & M &\sim 10^{12} M_{\odot} \\ &\sim 0.5 \text{ mas} & M &\sim M_{\odot} \end{aligned}$$

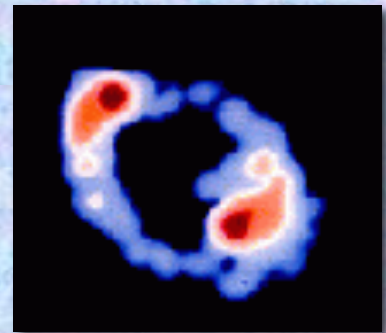
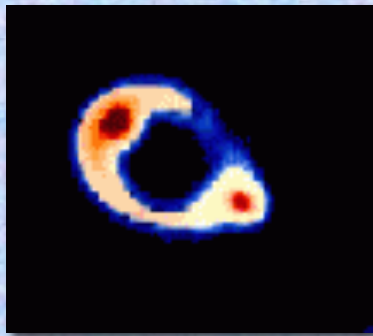
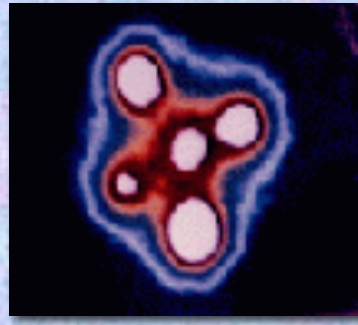
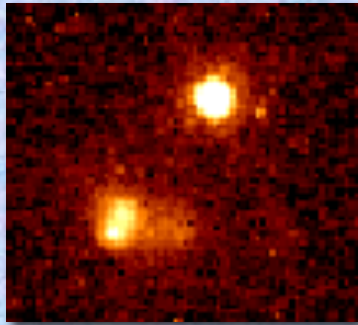
$$\beta \ll \theta_E, \quad \theta_{\pm} = \pm\theta_E + \beta/2$$

$$\beta = 0, \quad \theta_{\pm} = \pm\theta_E \quad \text{Einstein ring}$$

$$\beta \gg \theta_E, \quad \theta_{\pm} = \begin{cases} \beta + \theta_E^2/\beta \\ -\theta_E^2/\beta \end{cases}$$



Gravitational lenses



Gravitational lenses: Magnification

$$\theta_{\pm} = \frac{1}{2} \left(\beta \pm \sqrt{\beta^2 + 4\theta_E^2} \right)$$

Brightness proportional to $\int \beta d\beta d\varphi$ at the source, $\int \theta d\theta d\varphi$ at the observer

$$\mu_{\pm} = \frac{\theta_{\pm} d\theta_{\pm}}{\beta d\beta} = \pm \frac{1}{4} \left(\frac{\beta}{\sqrt{\beta^2 + 4\theta_E^2}} + \frac{\sqrt{\beta^2 + 4\theta_E^2}}{\beta} \pm 2 \right).$$

Microlensing

$$|\mu_+| + |\mu_-| = \frac{1}{2} \left(\frac{\beta}{\sqrt{\beta^2 + 4\theta_E^2}} + \frac{\sqrt{\beta^2 + 4\theta_E^2}}{\beta} \right) > 1$$



Post-Newtonian time

Clock synchronization on the Earth

$$d\tau = \left[1 - \frac{1}{c^2} \left(\frac{1}{2} \bar{v}^2 + U \right) + O(c^{-4}) \right] d\bar{t}$$

$$= \left[1 - \frac{1}{c^2} \left(\frac{1}{2} v^2 + \mathbf{v} \cdot (\boldsymbol{\omega} \times \mathbf{r}) + \Phi \right) + O(c^{-4}) \right] d\bar{t}$$

$$\Phi = U + \frac{1}{2} \omega^2 (x^2 + y^2) \quad \text{Geoid potential}$$

Two clocks at rest

$$\tau_A = \left(1 - \frac{\Phi_{\text{geoid}}}{c^2} \right) \bar{t}$$

$$\tau_B = \left(1 - \frac{\Phi_{\text{geoid}}}{c^2} \right) \bar{t} + \tau_B^0$$

For travelling clock with $v \ll c$

$$\tau_{\text{trav}} = \tau_A(\bar{t}_1) + \left(1 - \frac{\Phi_{\text{geoid}}}{c^2} \right) (\bar{t}_2 - \bar{t}_1) - \frac{1}{c^2} \int_A^B (\boldsymbol{\omega} \times \mathbf{r}) \cdot d\mathbf{r}$$

$$\tau_B^0 = -\frac{1}{c^2} \int_A^B (\boldsymbol{\omega} \times \mathbf{r}) \cdot d\mathbf{r}$$

Sagnac effect:
207 ns for closed path
at the equator



Post-Newtonian time: GPS

$$d\tau_{\text{ground}} = \left(1 - \frac{\Phi_{\text{geoid}}}{c^2} \right) d\bar{t}$$

$$\begin{aligned} d\tau_{\text{sat}} &= \left[1 - \frac{1}{c^2} \left(\frac{1}{2} v^2 + U \right) \right] d\bar{t} \\ &= \left(1 - \frac{3}{2} \frac{Gm}{c^2 a} \right) d\bar{t} \end{aligned}$$

$$d\tau_{\text{sat}} \left(1 + \frac{\Phi_{\text{geoid}}}{c^2} - \frac{3}{2} \frac{Gm}{c^2 a} \right) d\tau_{\text{ground}}$$

$$\sim \frac{Gm}{c^2 a_{\oplus}}$$

$$\sim 6.9 \times 10^{-10}$$

$$\sim \frac{3}{2} \frac{Gm}{c^2 (4.2 a_{\oplus})}$$

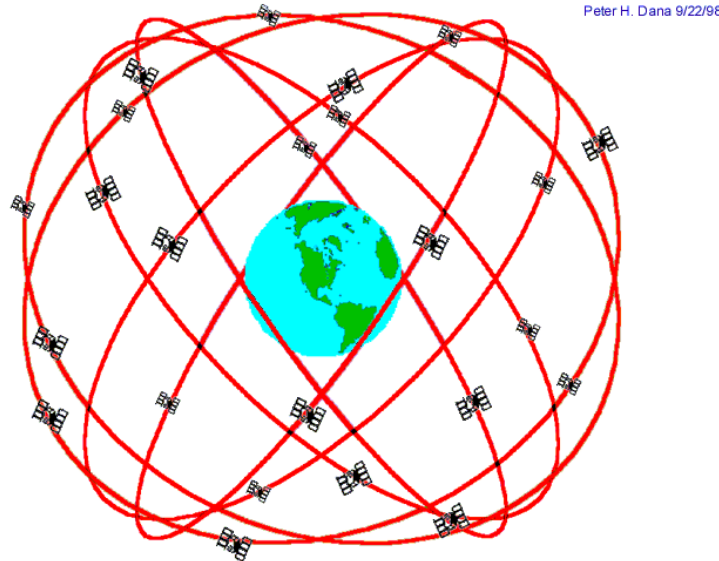
$$\sim 2.5 \times 10^{-10}$$

for GPS satellites

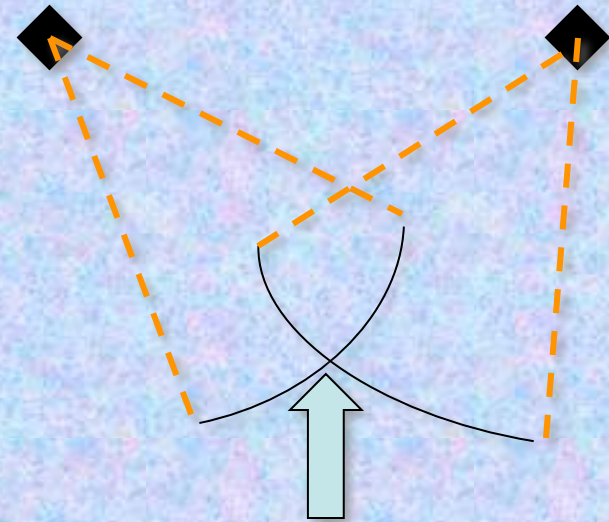


General Relativity and Daily Life

The Global Positioning System (GPS)



GPS Nominal Constellation
24 Satellites in 6 Orbital Planes
4 Satellites in each Plane
20,200 km Altitudes, 55 Degree Inclination



Navigation Requirement: 15 m $\Rightarrow$ 50ns

Relativistic effects: 39,000 ns per day!

GR = 46,000

SR = -7,000

Relativity **must**
be taken into
account, for GPS
to function



PN theory: Far-zone physics

Recall the wave-zone solution to $\square\psi = -4\pi\mu$

$$\frac{\mu(t - |\mathbf{x} - \mathbf{x}'|/c, \mathbf{y})}{|\mathbf{x} - \mathbf{x}'|} = \sum_{\ell=0}^{\infty} \frac{(-1)^\ell}{\ell!} x'^L \partial_L \frac{\mu(t - r/c, \mathbf{y})}{r}.$$

$$\psi_{\mathcal{N}}(t, \mathbf{x}) = \sum_{\ell=0}^{\infty} \frac{(-1)^\ell}{\ell!} \partial_L \left[\frac{1}{r} \int_{\mathcal{M}} \mu(\tau, \mathbf{x}') x'^L d^3x' \right]$$

Must add the integral over the wave zone $\psi_{\mathcal{W}}$

For the field $h^{\alpha\beta}$:

$$h_{\mathcal{N}}^{\alpha\beta}(t, \mathbf{x}) = \frac{4G}{c^4} \sum_{\ell=0}^{\infty} \frac{(-1)^\ell}{\ell!} \partial_L \left[\frac{1}{r} \int_{\mathcal{M}} \tau^{\alpha\beta}(\tau, \mathbf{x}') x'^L d^3x' \right]$$

