GR and QG: The Next 100 Years

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IUCAA, Pune

T.P., arXiv:1603.08658

T.P., arXiv:1611.03505

20th Century Physics

Top Two Discoveries

Matter is discrete and has microscopic degrees of freedom.

The Universe has dynamics and is expanding.

21st Century Physics

Top Two Discoveries

Spacetime is discrete and has microscopic degrees of freedom.

The Universe has its own dynamics (and is not to be treated as a special solution in GR).

SINGULARITIES: e.g., BLACK HOLES

COSMOLOGICAL CONSTANT

THE THERMODYNAMIC CONNECTION

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Classical GR is incomplete; we need QG!

HOW DO WE PUT TOGETHER THE PRINCIPLES OF GENERAL RELATIVITY AND QUANTUM THEORY?

GR: THE NEXT 100 YEARS

Needs another paradigm shift

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The equations governing classical gravity has the same conceptual status as those describing elasticity/hydrodynamics.

Why did/will the usual approaches to QG fail?

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If you quantize the equations of elasticity you will only get phonons — not the physics of the atoms!

THE PARADIGM

Key ingredient - 1

Spacetime dynamics should/can be described in a thermodynamic language; not in a geometrical language.

See e.g. TP, [gr-qc/0308070]; [arXiv:0911.5004]; [arXiv:0912.3165]; [arXiv:1003.5665]; [arXiv:1312.3253]; [arXiv:1405.5535]; Chakraborty, T. P, [arXiv:1408.4679]

The gravity-thermodynamics connection transcends GR even when the entropy is not proportional to area.

See e.g., Paranjape et al, [hep-th/0607240]; Kothawala, TP [arXiv:0904.0215]; Kolekar et al. [arXiv:1111.0973], [arXiv:1201.2947], Chakraborty, T. P, [arXiv:1408.4679]; [arXiv:1408.4791]

You should get more than what you put in!

Physics is all about numbers and making falsifiable predictions!

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Physics is all about numbers and making falsifiable predictions!

Provides new insights about: (i) classical gravity (ii) the microscopic structure of spacetime and (iii) cosmological constant.

More IS Different

The key new variable which distinguishes thermodynamics from point mechanics

$$Heat\ Density = \mathcal{H} = rac{Q}{V} = rac{TS}{V} = rac{1}{V}(E-F)$$
 $rac{TS}{V} = Ts = p +
ho = T_{ab}\ell^a\ell^b$

Normal matter has a heat density

Spacetime also has a heat density!

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One can associate a T and s with every event in spacetime just as you could with a glass of water!

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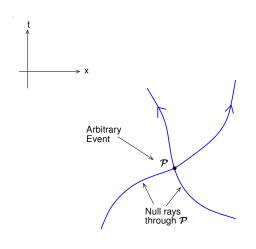
This fact transcends black hole physics and Einstein gravity.

Spacetime also has a heat density!

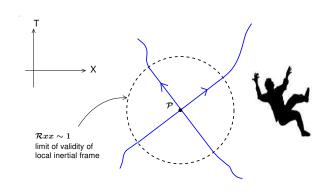
One can associate a T and s with every event in spacetime just as you could with a glass of water!

The T is independent of the theory of gravity; s depends on/determines the theory.

Spacetime in Arbitrary Coordinates

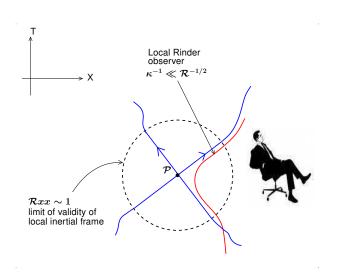


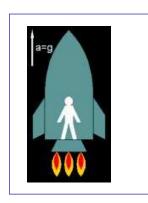
Local Inertial Observers



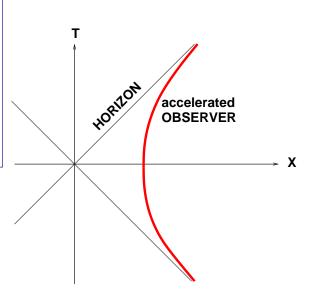
Validity of laws of SR \Rightarrow How gravity affects matter

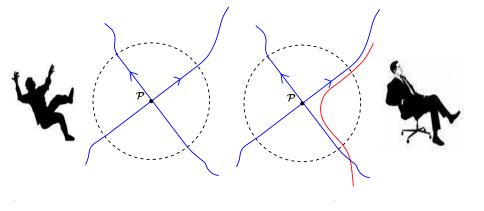
Local Rindler Observers





FLAT SPACETIME

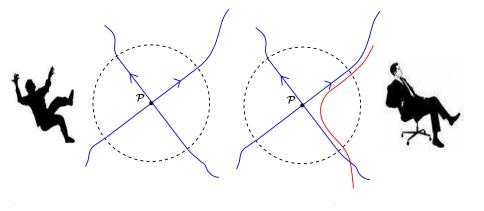




Vacuum fluctuations



Thermal fluctuations



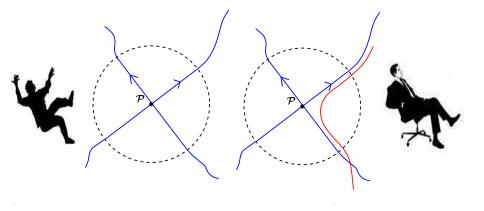
Vacuum fluctuations



Thermal fluctuations

$$k_BT=rac{\hbar}{c}\left(rac{g}{2\pi}
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A VERY NON-TRIVIAL EQUIVALENCE!



Vacuum fluctuations



Thermal fluctuations

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A VERY NON-TRIVIAL EQUIVALENCE!

QFT in FFF introduces \hbar ; we now have (\hbar/c) in the temperature

Spacetimes, Like Matter, can be Hot

The most beautiful result in the interface of quantum theory and gravity

OBSERVERS WHO PERCEIVE A HORIZON ATTRIBUTE A TEMPERATURE TO SPACETIME

$$k_BT=rac{\hbar}{c}\left(rac{g}{2\pi}
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[Davies (1975), Unruh (1976)]

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$$k_BT=rac{\hbar}{c}\left(rac{g}{2\pi}
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[Davies (1975), Unruh (1976)]

This allows you to associate a heat density $\mathcal{H} = Ts$ with every event of spacetime!

Regions of spacetime can be inaccessible to certain class of observers in any spacetime!

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Inaccessibility of Information ⇔ Entropy

Information Content of spacetime plays a central role in the new paradigm!

The Importance Of Being Hot

You could have figured out that water is made of discrete atoms without ever probing it at Angstrom scales!

Atoms of Matter

Boltzmann: If you can heat it, it must have micro-structure!

To store energy ΔE at temperature T, you need

$$\Delta n = rac{\Delta E}{(1/2)k_BT}$$

degrees of freedom. Microphysics leaves its signature at the macro-scales

Atoms of Spacetime

Boltzmann: If you can heat it, it must have micro-structure!

You can heat up spacetime!

Do we have an equipartition law for the microscopic spacetime degrees of freedom?

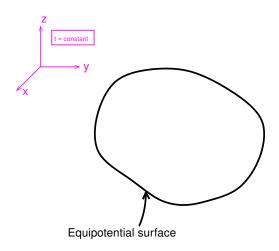
Can you count the atoms of space?

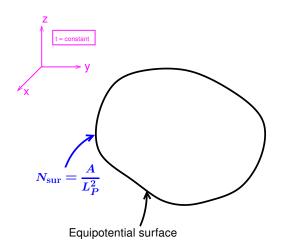
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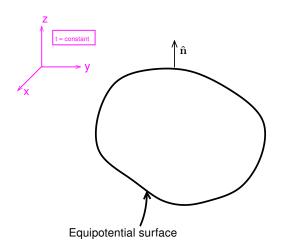
Equipartition with a surface-bulk correspondence

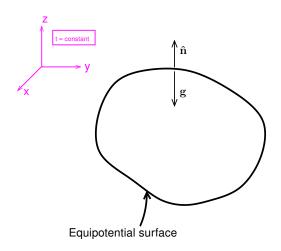
$$E_{
m bulk} = \int_{\partial \mathcal{V}} rac{dA}{L_P^2} \left(rac{1}{2} k_B T_{loc}
ight) \equiv rac{1}{2} k_B \int_{\partial \mathcal{V}} dn \, T_{
m loc}$$

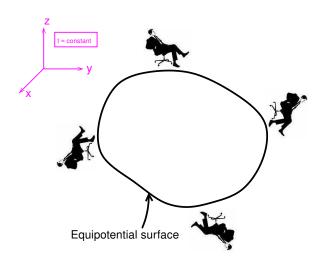
An area dA is endowed with $dn=dA/L_P^2$ microscopic degrees of freedom ('atoms')

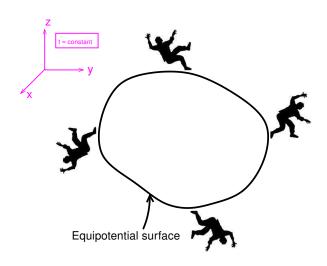


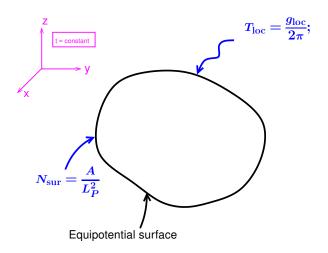


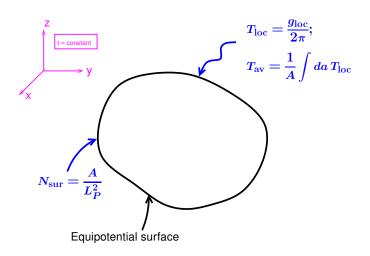


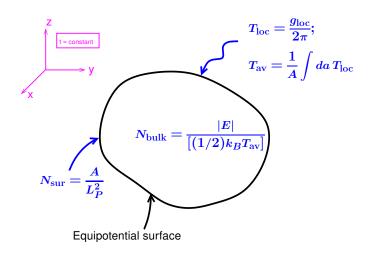


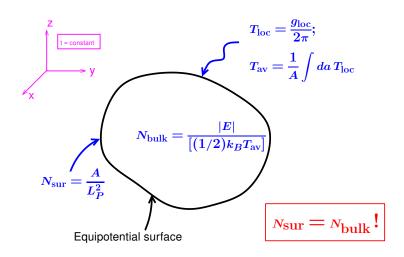












Evolution of spacetime

Review: TP, [arXiv:1410.6285]

All static geometries have

$$N_{
m sur}=N_{
m bulk}$$

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Evolution arises from departure from 'holographic equipartition':

$$Time\ evolution\ \propto\ Heating/Cooling \ \propto\ (N_{
m sur}-N_{
m bulk})$$

Geometry ⇔ **Thermodynamics**

K. Parattu, B.R. Majhi, T.P. [arXiv:1303.1535]

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$$q^{ab} \equiv \sqrt{-g} g^{ab}$$

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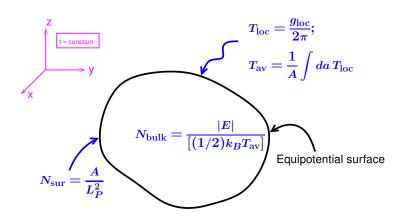
$$p_{bc}^a \equiv -\Gamma_{bc}^a + rac{1}{2}(\Gamma_{bd}^d\delta_c^a + \Gamma_{cd}^d\delta_b^a) \, .$$

These variables have a thermodynamic interpretation

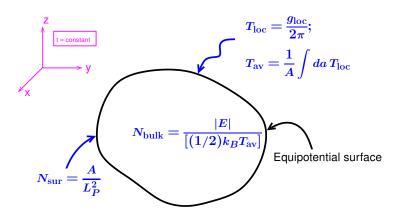
$$(q\delta p, p\delta q) \Leftrightarrow (s\delta T, T\delta s)$$

TP, Gen.Rel.Grav (2014) [arXiv:1312.3253]

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$$\int \frac{d\sigma_a}{8\pi L_P^2} q^{\ell m} \mathcal{L}_{\xi} p_{\ell m}^a = -\frac{1}{2} k_B T_{\text{av}} \left(N_{\text{sur}} - N_{\text{bulk}} \right)$$

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time evolution of spacetime

= heating of spacetime

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time evolution of spacetime

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deviation from holographic equipartition

Evolution of spacetime is described in thermodynamic language; not in geometric language!

Newton's Law of Gravitation

T.P. [hep-th/0205278]

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Three constants: \hbar, c, L_P^2

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$$m{F} = \left(rac{m{c}^3m{L}_P^2}{\hbar}
ight)\left(rac{m{m}_1m{m}_2}{m{r}^2}
ight)$$

T.P. [hep-th/0205278]

Three constants: \hbar, c, L_P^2 Temperature $\Rightarrow (\hbar/c)$; Entropy $\Rightarrow L_P^2$

$$F=\left(rac{c^3L_P^2}{\hbar}
ight)\left(rac{m_1m_2}{r^2}
ight)$$

Gravity, like matter, is intrinsically quantum and cannot exist in the limit of $\hbar o 0$!

T.P., A. Paranjape [gr-qc/0701003]; T.P. [arXiv:0705.2533]

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Field equations arise from maximizing entropy/heat density of gravity plus matter on all null surfaces.

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$$Q=\int d{\cal V}\,\left({\cal H}_g+{\cal H}_m
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Field equations arise from maximizing entropy/heat density of gravity plus matter on all null surfaces.

$$Q=\int d{\cal V}\,\left({\cal H}_g+{\cal H}_m
ight)$$

Works for a wide class gravitational theories; entropy decides the theory.

Macroscopic Nature Of Gravity

Gravity responds to heat density $(Ts = p + \rho)$ — not energy density!

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Cosmological constant arises as an integration constant

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Gravity responds to heat density $(Ts = p + \rho)$ — not energy density!

Cosmological constant arises as an integration constant

Its value is determined by a new conserved quantity for the universe!

Can we understand these results at a deeper level?

Can we get something new?

The Challenge

$$Q=\int d{\cal V}\,\left({\cal H}_g+{\cal H}_m
ight)$$

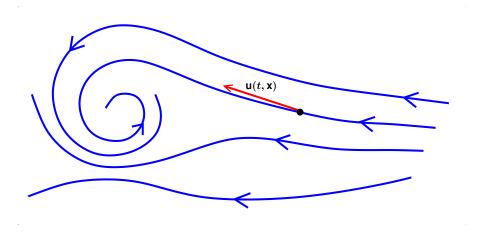
We need to count the microscopic degrees of freedom without knowing the full QG!

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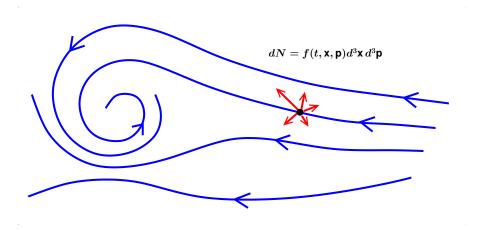
We need to recognize discreteness and yet use continuum mathematics!

Continuum description



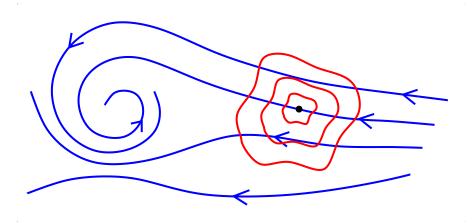
Continuum fluid mechanics: $\rho(x^i)$, $U(x^i)$, ignores discreteness and *velocity dispersion*.

Kinetic theory: Distribution function



Kinetic theory: Recognizes discreteness, yet uses continuum maths! Several atoms with different p_i can exist at same x^i

Kinetic theory: Distribution function



 $f(x^i,p_i)=$ number of atoms "at" $x^i,$ "with" momentum p_i ; to be determined by a limiting process.

Atoms Of Space

The distribution function for 'atoms of space' provides the microscopic origin for the variational principle The distribution function for 'atoms of space' provides the microscopic origin for the variational principle

You can determine it in a spacetime with zero-point length using equi-geodesic surfaces

Geodesic Interval

D. Kothawala, T.P. [arXiv:1405.4967]; [arXiv:1408.3963]

The geodesic interval $\sigma^2(x,x')$ and g_{ab} have the same information about geometry:

$$\sigma(x,x')=\int_{x}^{x'}\sqrt{g_{ab}n^{a}n^{b}}d\lambda$$

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$$egin{aligned} \sigma(x,x') &= \int_x^{x'} \sqrt{g_{ab} n^a n^b} d\lambda \ &\lim_{x' o x} rac{1}{2}
abla_a
abla_b \sigma^2 &= g_{ab} \end{aligned}$$

Zero-Point Length

T.P. Ann.Phy. (1985), 165, 38; PRL (1997), 78, 1854

Zero-Point Length

T.P. Ann. Phy. (1985), 165, 38; PRL (1997), 78, 1854

Discreteness arises through a quantum of area, which is a QG effect

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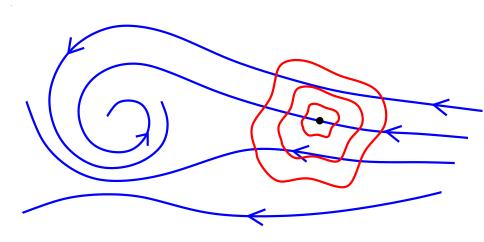
Quantum spacetime has a zero-point length:

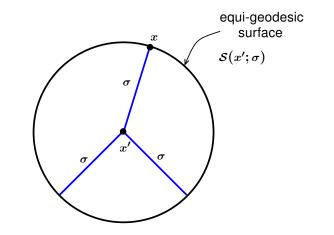
 $D.\ Kothawala, T.P.\ [arXiv:1405.4967]; [arXiv:1408.3963]$

$$egin{array}{lll} \sigma^2(x,x') & o & S(\sigma^2) = \sigma^2(x,x') + L_0^2 \ g_{ab}(x) & o & q_{ab}(x,x';L_0^2) \end{array}$$

In conventional units,
$$L_0^2=(3/4\pi)L_P^2$$
.

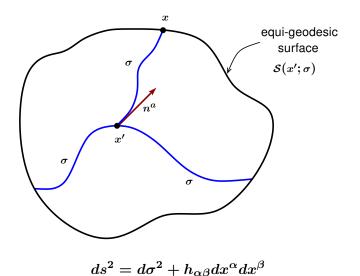
Kinetic theory: Distribution function



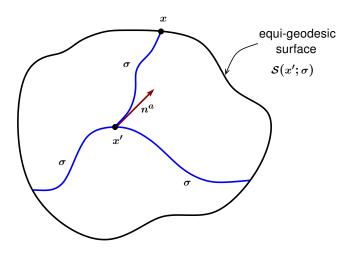


$$\sqrt{g} \propto \sigma^3 ~~ \sqrt{h} \propto \sigma^3$$

 $ds^2 = d\sigma^2 + \sigma^2 d\Omega_{(S3)}^2$



The $\sqrt{g} = \sqrt{h}$ will pick up curvature corrections



$$ds^2 = d\sigma^2 + h_{\alpha\beta} dx^{\alpha} dx^{\beta}$$

$$\sqrt{h}(x,x') = \sqrt{g}(x,x') = \sigma^3 \left(1 - rac{\sigma^2}{6} \mathcal{E}
ight) \sqrt{h_\Omega}; \;\;\; \mathcal{E} \equiv R_{ab} n^a n^b$$

Area Of A Point

T.P. [arXiv:1508.06286]

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$$\sqrt{q} = \sigma \left(\sigma^2 + L_0^2
ight) \left[1 - rac{1}{6} \mathcal{E} \left(\sigma^2 + L_0^2
ight)
ight] \sqrt{h_\Omega}$$

$$\sqrt{h} = L_0^3 \left[1 - rac{1}{6} \mathcal{E} L_0^2
ight] \sqrt{h_\Omega}$$



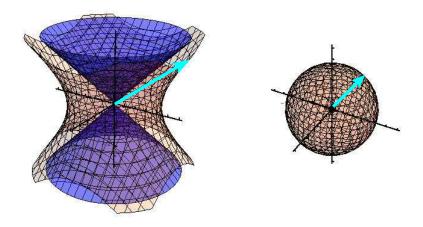
Fig. The number of atoms of space at x^i scales as area measure of the equigeodesic surface when $x' \to x$

$$f(x^i,\ell_a) \propto \lim_{\sigma o 0} \sqrt{h(x,\sigma)} pprox 1 - rac{1}{8\pi} L_P^2 R_{ab} \ell^a \ell^b$$

ullet The number of atoms of space at x^i scales as area measure of the equigeodesic surface when x' o x

$$f(x^i,\ell_a) \propto \lim_{\sigma o 0} \sqrt{h(x,\sigma)} pprox 1 - rac{1}{8\pi} L_P^2 R_{ab} \ell^a \ell^b$$

Fuclidean origin maps to local Rindler horizon. The $\sigma^2 \to 0$ limit picks the null vectors!



Null horizon ⇔ Euclidean origin

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$$f(x^i,\ell_a) \propto \lim_{\sigma o 0} \sqrt{h(x,\sigma)} pprox 1 - rac{1}{8\pi} L_P^2 R_{ab} \ell^a \ell^b$$

- Euclidean origin maps to local Rindler horizon. The $\sigma^2 \to 0$ limit picks the null vectors!
- The variational principle leads to

$$G_{ab} = \kappa T_{ab} + \Lambda g_{ab}$$

Quantum of Information

T.P, Chakraborty, Kothawala, [arXiv:1507.05669], T.P. [arXiv:1508.06286]

Spacetime becomes two-dimensional close to Planck scales!

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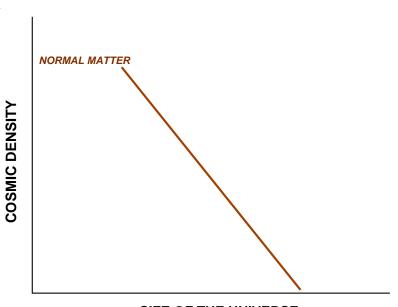
See e.g., Carlip et al., [arXiv:1103.5993]; [arXiv:1009.1136]; G. Calcagni et al, [arXiv:1208.0354]; [arXiv:1311.3340]; J. Ambjorn, et al. [arXiv:hep-th/0505113]; L. Modesto, [arXiv:0812.2214]; V. Husain et al., [arXiv:1305.2814] etc.

- Spacetime becomes two-dimensional close to Planck scales!
- Basic quantum of information to count spacetime degrees of freedom is

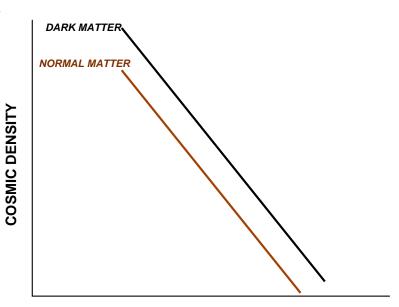
$$I_{QG}=rac{4\pi L_P^2}{L_P^2}=4\pi$$

You should get more than what you put in!

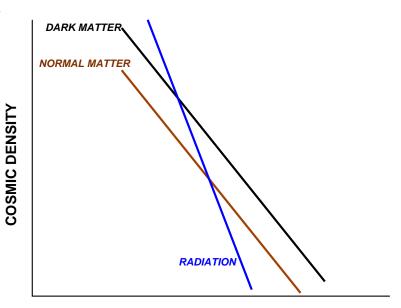
Physics is all about numbers and making falsifiable predictions!



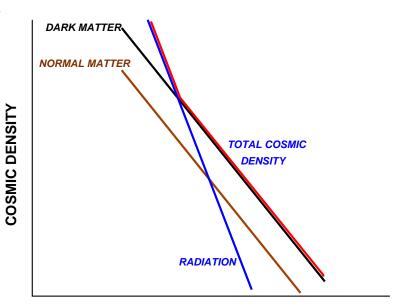
SIZE OF THE UNIVERSE



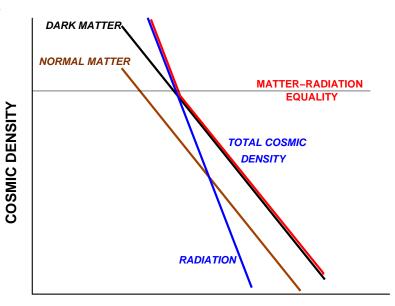
SIZE OF THE UNIVERSE

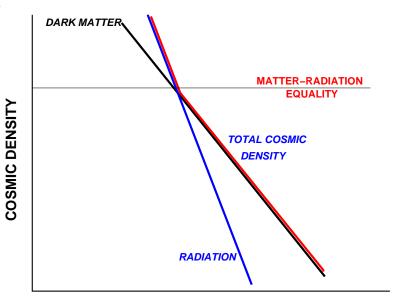


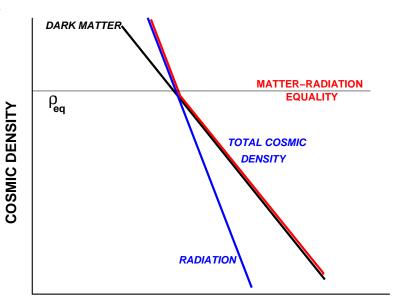
SIZE OF THE UNIVERSE



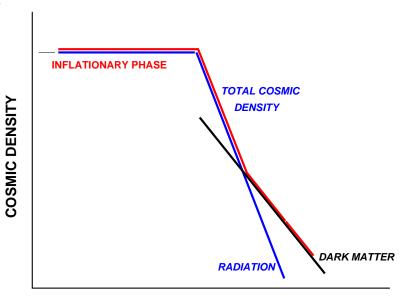
SIZE OF THE UNIVERSE

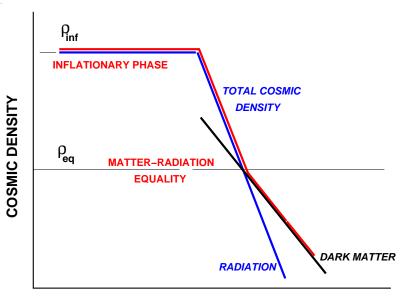


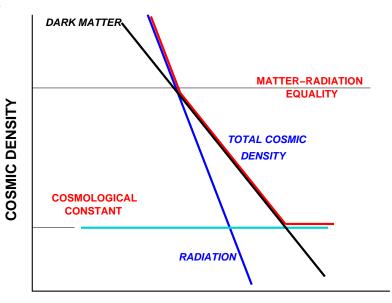




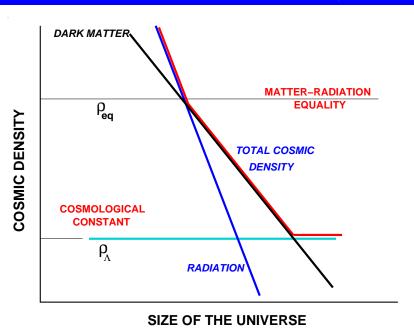
SIZE OF THE UNIVERSE

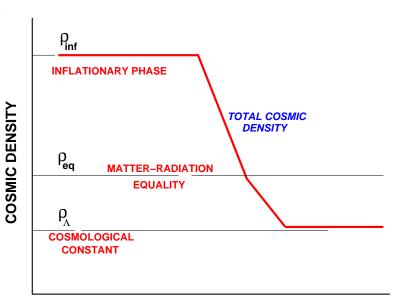






SIZE OF THE UNIVERSE





$$ho_{
m inf} < (1.94 imes 10^{16} \, {
m GeV})^4$$

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m inf} < (1.94 imes 10^{16}~{
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$$ho_{
m eq} = rac{
ho_m^4}{
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m eV}]^4$$

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$$ho_{\Lambda} = [(2.26 \pm 0.05) imes 10^{-3} ext{eV}]^4$$



Hope: High energy physics will (eventually!) fix $ho_{inf}^{1/4}pprox 10^{15} { m GeV}$ and

$$\left|
ho_{eq}^{1/4} \propto \left|rac{n_{DM}}{n_{\gamma}}m_{DM} + rac{n_{B}}{n_{\gamma}}m_{B}
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ight| pprox 0.86 ext{ eV}$$

But we have no clue why

$$ho_{\Lambda} L_{P}^{4} pprox 1.4 imes 10^{-123} pprox 1.1 imes e^{-283}$$
.

► Compute the combination:

$$I=rac{1}{9\pi}\,\ln\left(rac{4}{27}rac{
ho_{
m inf}^{3/2}}{
ho_{\Lambda}\,
ho_{
m eq}^{1/2}}
ight)$$

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You will find that

$$Ipprox 4\pi\left(1\pm\mathcal{O}\left(10^{-3}
ight)
ight)=I_{QG}$$

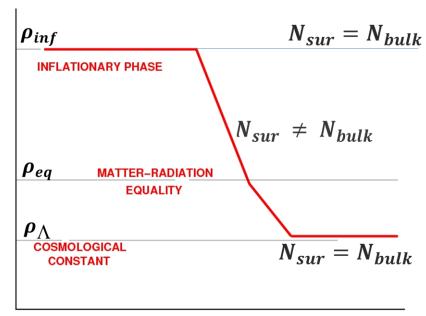
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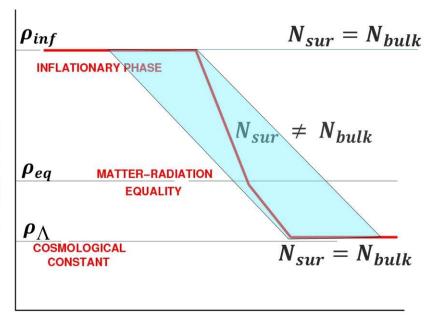
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$$x(a_2,a_1)=\int_{t_1}^{t_2}rac{dt}{a(t)}=\int_{a_1}^{a_2}rac{da}{a^2H(a)}$$

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Boundary of accessible cosmic information for eternal observer

$$x(\infty,a)\equiv x_\infty(a)=\int_a^\infty rac{dar a}{ar a^2 H(ar a)}$$

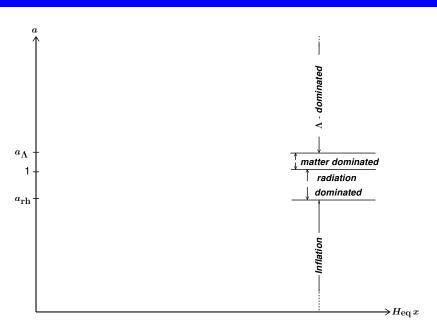
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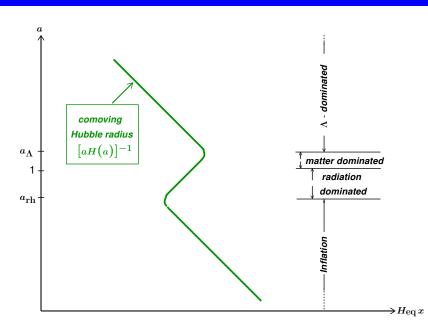
$$x(\infty,a)\equiv x_\infty(a)=\int_a^\infty rac{dar a}{ar a^2 H(ar a)}$$

This is infinite if $\Lambda=0$; finite if $\Lambda \neq 0$

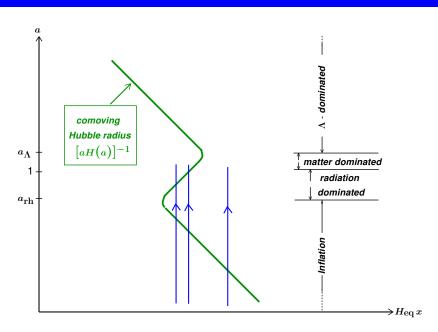
Cosmic Information: CosmIn

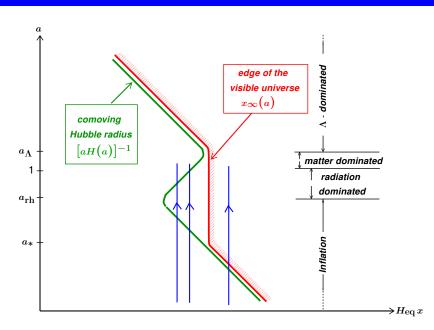


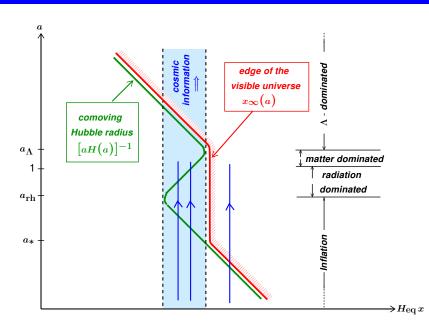
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T.P, Hamsa Padmanabhan [arXiv:1404.2284]

A measure of cosmic information accessible to eternal observer ('Cosmln')

 I_c = Number of modes (geodesics) which cross the Hubble radius during the radiation + matter dominated phase .

$$I_c = rac{2}{3\pi} \ln \left(rac{a_{
m rh}}{a_*}
ight)$$

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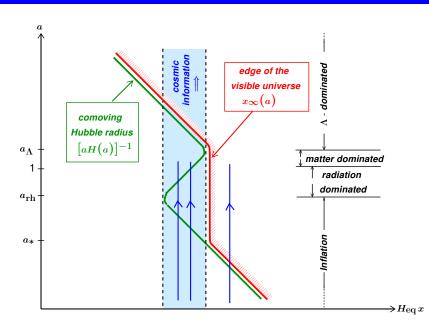
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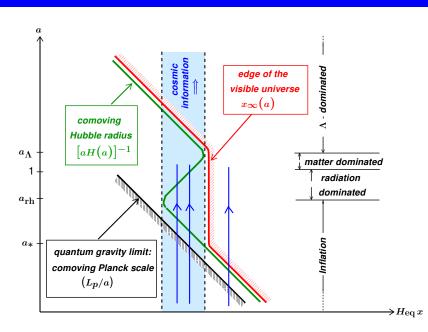
$$I_c = rac{1}{9\pi}\,\ln\left(rac{4}{27}rac{
ho_{
m inf}^{3/2}}{
ho_\Lambda\,
ho_{
m eq}^{1/2}}
ight)$$

Cosmln and the Λ

T.P, Hamsa Padmanabhan [arXiv:1404.2284]

$$ho_{\Lambda} = rac{4}{27} rac{
ho_{
m inf}^{3/2}}{
ho_{
m eq}^{1/2}} \, \exp\left(-9\pi I_c
ight)$$





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m inf}^{3/2}}{
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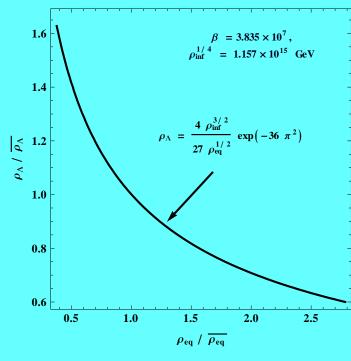
$$ho_{\Lambda} = rac{4}{27} rac{
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m inf}^{3/2}}{
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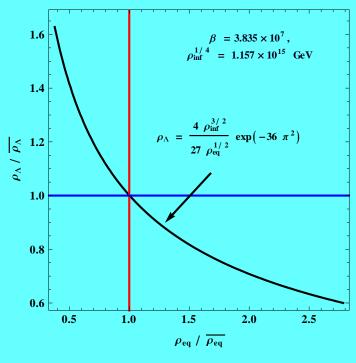
Using $I_c=I_{QG}=4\pi$ gives the numerical value of ho_{Λ}

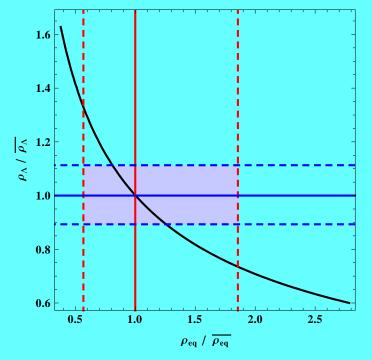
The Magical Relation!

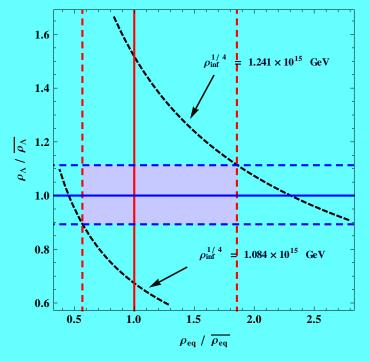
$$ho_{\Lambda} = rac{4}{27} rac{
ho_{inf}^{3/2}}{
ho_{eq}^{1/2}} \exp(-36\pi^2)$$

Hamsa Padmanabhan, **T.P**, *CosMIn: Solution to the Cosmological constant problem* [arXiv:1302.3226]









Cosmic Conundrums

- Large scale universe defines a Lorentz frame
- The universe made a spontaneous quantum to classical transition
- Cosmic evolution is towards holographic equipartition

We should not describe the cosmos as a particular solution to gravitational field equations.

The FRW universe is described by the two equations (with $T=H/2\pi$)

$$rac{dV_H}{dt} = L_P^2(N_{sur}-N_{bulk})$$

$$\rho V_H = TS$$

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► Gravitational dynamics arises from a thermodynamic variational principle principle. Its form can be related to the kinetic theory of atoms of space.

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- ► The evolution equation has a purely thermodynamic interpretation related to the information content of the spacetime.
- ► The cosmological constant is related to the amount of information accessible to an eternal observer.
- ▶ At Planck scales spacetime is 2-dimensional with 4π units of information; this allows the determination of the value of the cosmological constant.

References

T.P, General Relativity from a Thermodynamic Perspective, Gen. Rel. Grav., **46**, 1673 (2014) [arXiv:1312.3253].

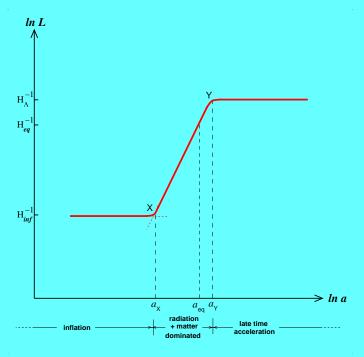
Review: T.P, The Atoms Of Space, Gravity and the Cosmological Constant, IJMPD, 25, 1630020 (2016) [arXiv:1603.08658].

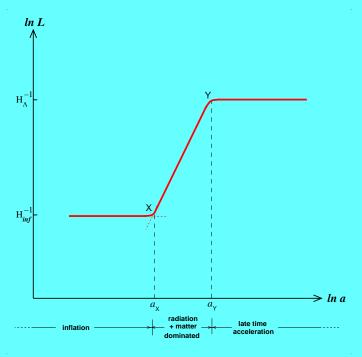
Review: T.P, Do we really understand the cosmos?, (2016) [arXiv:1611.03505].

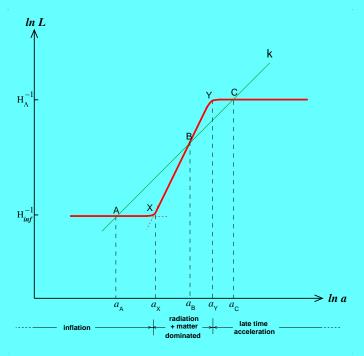
Acknowledgements

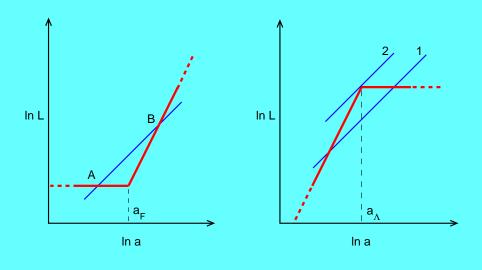
Sunu Engineer Dawood Kothawala Bibhas Majhi Krishna Parattu Sumanta Chakraborty James Bjorken Aseem Paranjape Hamsa Padmanabhan Donald Lynden-Bell

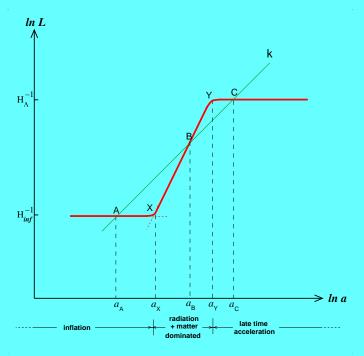
THANK YOU FOR YOUR TIME!

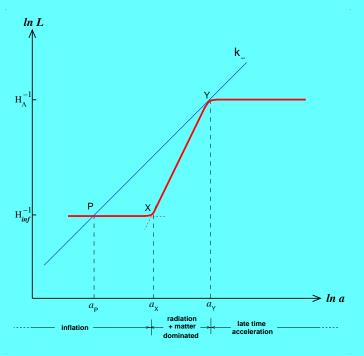


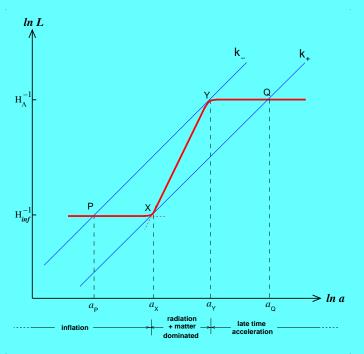




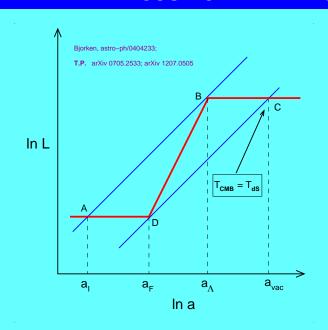




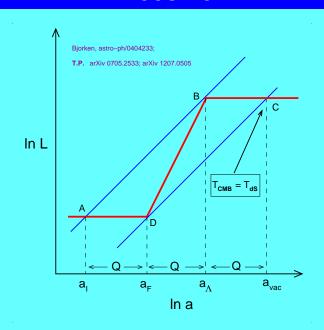




COSMIC PARALLELOGRAM



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Key Open Question Possibility for Matter Sector

Matter and Geometry need to emerge together for proper interpretation of $T^{ab}n_an_b$ at the microscopic scale. How do we do this?

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$$\langle n^a n^b
angle pprox (4\pi/\mu L_P^2) R_{ab}^{-1}$$

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$$\langle n^a n^b
angle pprox (4\pi/\mu L_P^2) R_{ab}^{-1}$$

$$2\mu L_P^4 ~\langle ar{T}_{ab} n^a n^b
angle pprox 2\mu L_P^4 ~\langle ar{T}_{ab}
angle \langle n^a n^b
angle = 1$$

COSMIC EXPANSION: A QUEST FOR HOLOGRAPHIC EQUIPARTITION

T.P., [1207.0505]

$$egin{align} rac{dR_H}{dt} &= (1-rac{\epsilon N_{
m bulk}}{N_{
m sur}}) \hspace{0.5cm} \epsilon = \pm 1 \ N_{
m sur} &= 4\pirac{R_H^2}{L_P^2}; \hspace{0.5cm} N_{
m bulk} = -\epsilonrac{E}{(1/2)k_BT}; \hspace{0.5cm} T = rac{H}{2\pi} \ \end{array}$$

Remarkably enough, this leads to the standard FRW dynamics!

$$N_{
m sur} = rac{A}{L_P^2} = N_{
m bulk} = rac{E}{[(1/2)k_BT_{
m av}]}$$

$$E = rac{1}{2L_P^2} \int da \; k_B T_{
m loc}$$

$$E=rac{1}{2L_{P}^{2}}\int da\;k_{B}T_{\mathrm{loc}}$$

$$\int
ho \; dV = rac{1}{2L_P^2} \int da \; \left(rac{\hbar}{c}
ight) \left(rac{g}{2\pi}
ight)$$

$$E=rac{1}{2L_{P}^{2}}\int da\;k_{B}T_{\mathrm{loc}}$$

$$\int
ho \; dV = rac{1}{2L_{P}^{2}} \int da \; rac{\hbar}{c} \left(rac{- ext{g}\cdot\hat{ ext{n}}}{2\pi}
ight)$$

$$E=rac{1}{2L_{P}^{2}}\int da\;k_{B}T_{\mathrm{loc}}$$

$$\int
ho \; dV = rac{1}{2L_P^2} \int da \; rac{\hbar}{c} \left(rac{-{f g} \cdot \hat{f n}}{2\pi}
ight)
onumber \ = -rac{\hbar}{4\pi c L_P^2} \int dV \; {f
abla} \cdot {f g}$$

$$E=rac{1}{2L_{P}^{2}}\int da\;k_{B}T_{\mathrm{loc}}$$

$$\int
ho \; dV = rac{1}{2L_P^2} \int da \; rac{\hbar}{c} \left(rac{-\mathbf{g}\cdot\hat{\mathbf{n}}}{2\pi}
ight)
onumber \ = -rac{\hbar}{4\pi c L_P^2} \int dV \; \mathbf{
abla} \cdot \mathbf{g}$$

$$abla \cdot \mathrm{g} = -rac{4\pi c L_P^2}{\hbar}
ho = -4\pi G \left(rac{
ho}{c^2}
ight)$$



Local Rindler Horizon

► Heat transfered due to matter crossing a null surface:

[T. Jacobson, qr-qc/9504004]

$$Q_m = \int d\mathcal{V} \, (T_{ab}\ell^a\ell^b); \quad \mathcal{H}_m \equiv T_{ab}\ell^a\ell^b$$

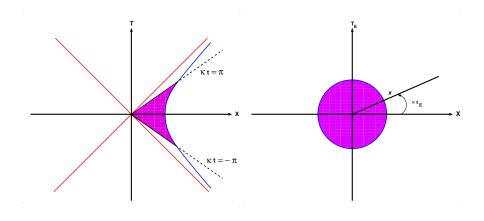
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▶ Note: Null horizon ⇔ Euclidean origin

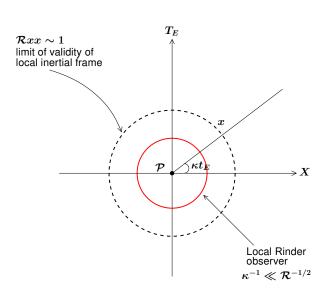
$$X^2 - T^2 = 0 \Leftrightarrow X^2 + T_E^2 = 0$$

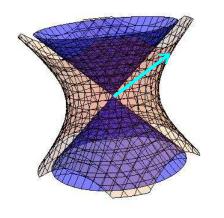


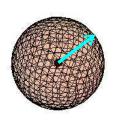
$$T = x \sinh \kappa t, \ X = x \cosh \kappa t$$

 $T = x \sinh \kappa t, \ X = x \cosh \kappa t$ $T_E = x \sin \kappa t_E, \ X = x \cos \kappa t_E$

$$X^2 - T^2 = 0 \Leftrightarrow X^2 + T_E^2 = 0$$







$$X^2-T^2=\sigma^2\Leftrightarrow X^2+T_E^2=\sigma^2$$
 $X^2-T^2=0\Leftrightarrow X^2+T_E^2=0$

$$egin{array}{ll} \Omega_{
m tot} &=& \prod_{\phi_A} \prod_x \,
ho_g(\mathcal{G}_N,\phi_A) \,
ho_m(T_{ab},\phi_A) \ &\equiv& \prod_{n_a} \exp \sum_x \left(\ln
ho_g + \ln
ho_m
ight) \end{array}$$

$$egin{array}{ll} \Omega_{
m tot} &=& \prod_{\phi_A} \prod_x \,
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ight) \end{array}$$

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ight) \end{array}$$

$$egin{array}{ll} \ln
ho_m & \equiv & L_P^4 \mathcal{H}_m = L_P^4 T_{ab} \ell^a \ell^b \ & \ln
ho_g & \equiv & L_P^4 \mathcal{H}_g pprox -rac{L_P^2}{8\pi} R_{ab} \ell^a \ell^b \end{array}$$

$$egin{array}{ll} \Omega_{\mathrm{tot}} &=& \prod_{\phi_A} \prod_x \,
ho_g(\mathcal{G}_N,\phi_A) \,
ho_m(T_{ab},\phi_A) \ &\equiv& \prod_{\ell_a} \exp \int d\mathcal{V} \left(\mathcal{H}_g(\mathcal{G}_N,\ell_a) + \mathcal{H}_m(T_{ab},\ell_a)
ight) \end{array}$$

 $R_b^a \ell_a \ell^b = (8\pi L_P^2) T_b^a \ell_a \ell^b$

$$egin{array}{ll} \Omega_{\mathrm{tot}} &=& \prod_{\phi_A} \prod_x \,
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ho_m(T_{ab},\phi_A) \ &\equiv& \prod_{\ell_a} \exp \int d\mathcal{V} \left(\mathcal{H}_g(\mathcal{G}_N,\ell_a) + \mathcal{H}_m(T_{ab},\ell_a)
ight) \end{array}$$

 $G_h^a = (8\pi L_P^2)T_h^a + (\text{const})\delta_h^a$

T.P., A. Paranjape [gr-qc/0701003]; T.P. [arXiv:0705.2533]

$${\cal H}_m = T^a_b \ell_a \ell^b$$

$$\mathcal{H}_g = -\left(rac{1}{16\pi L_D^2}
ight)(4P_{cd}^{ab}
abla_a\ell^c
abla_b\ell^d)$$

T.P., A. Paranjape [gr-qc/0701003]; T.P. [arXiv:0705.2533]

$${\cal H}_m = T^a_b \ell_a \ell^b$$

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abla_a\ell^c
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$$P_{cd}^{ab} \propto \delta_{cdc_{2}d_{2}...c_{m}d_{m}}^{aba_{2}b_{2}...a_{m}b_{m}} R_{a_{2}b_{2}}^{c_{2}d_{2}}...R_{a_{m}b_{m}}^{c_{m}d_{m}}$$

► The P^{ab}_{cd} is the entropy tensor of the spacetime which determines the theory [[yer and Wald (1994)]

T.P., A. Paranjape [gr-qc/0701003]; T.P. [arXiv:0705.2533]

T.P., A. Paranjape [gr-qc/0701003]; T.P. [arXiv:0705.2533]

▶ The demand $\delta Q/\delta \ell_a=0$ for all null ℓ_a leads to:

$$E^a_b \equiv P^{ai}_{jk} R^{jk}_{bi} - rac{1}{2} \delta^a_b \mathcal{R} = (8\pi L_P^2) T^a_b + \Lambda \delta^a_b,$$

T.P., A. Paranjape [gr-qc/0701003]; T.P. [arXiv:0705.2533]

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$$E^a_b \equiv P^{ai}_{jk} R^{jk}_{bi} - rac{1}{2} \delta^a_b \mathcal{R} = (8\pi L_P^2) T^a_b + \Lambda \delta^a_b,$$

ullet These are Lanczos-Lovelock models of gravity. In d=4, it uniquely leads to GR

$$G_b^a = (8\pi L_P^2)T_b^a + \Lambda \delta_b^a$$

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► On-shell value of Q_{tot}

$$Q_{
m tot}^{
m on-shell} = \int d^2x (T_{
m loc}\,s)igg|_{\lambda_1}^{\lambda_2}$$

Algebraic Aside

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Interestingly enough:

$$2P_{cd}^{ab}
abla_a n^c
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► Alternative, dimensionless, form in GR:

$${\cal K}_g \equiv -rac{1}{8\pi}(L_P^2 R_{ab} n^a n^b)$$