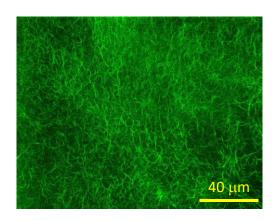
Memory retention in disordered bio-polymer networks

Sayantan Majumdar

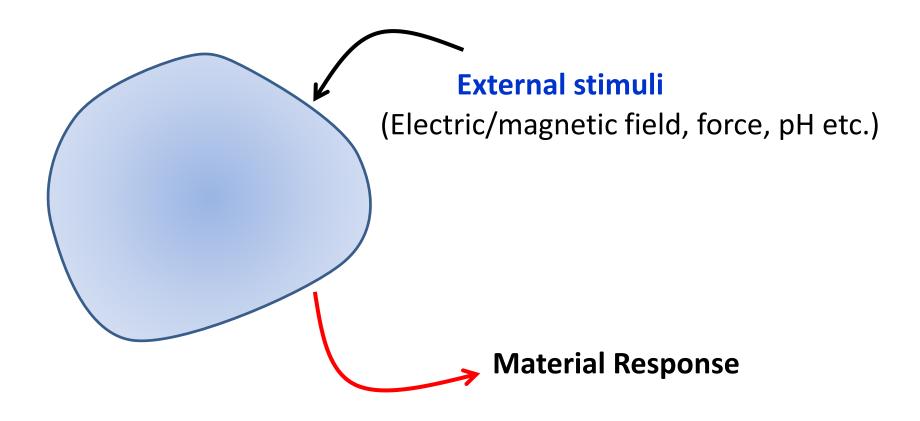
Raman Research Institute, Bangalore





EIOSM, ICTS 2018

Adaptive materials

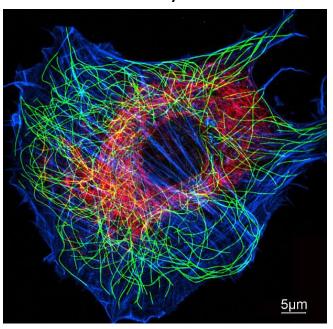


Rapid, controlled reversible changes: adapting to changing environment

Designing such materials is extremely challenging

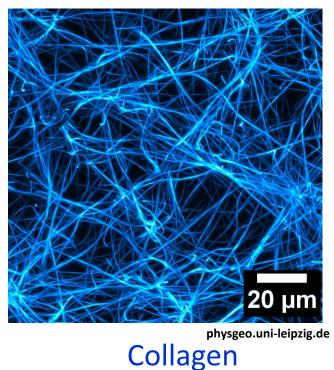
Biological polymers: new design motif

Inside Eukaryotic cell



Actin, Microtubule, IF

Extracellular matrix

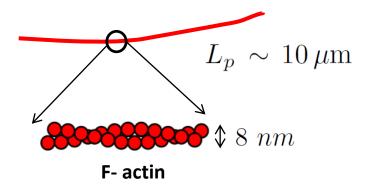


Unique materials: no synthetic analogue

Cross-linked actin networks: design motif for programmable materials

Persistence Length quantifies bending stiffness

$$L_p = \frac{\kappa_{bend}}{k_B T}$$

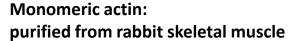


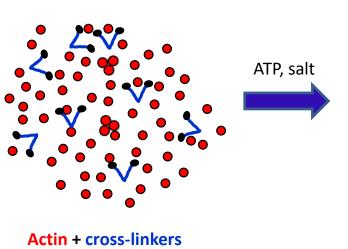




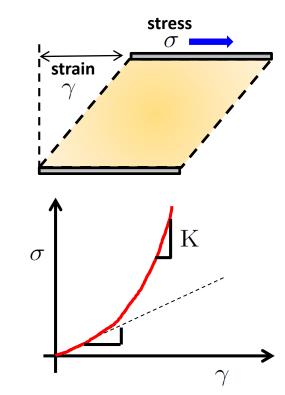
Poly-ethylene







Cross-linked actin network



Stress history significantly alters the material moduli



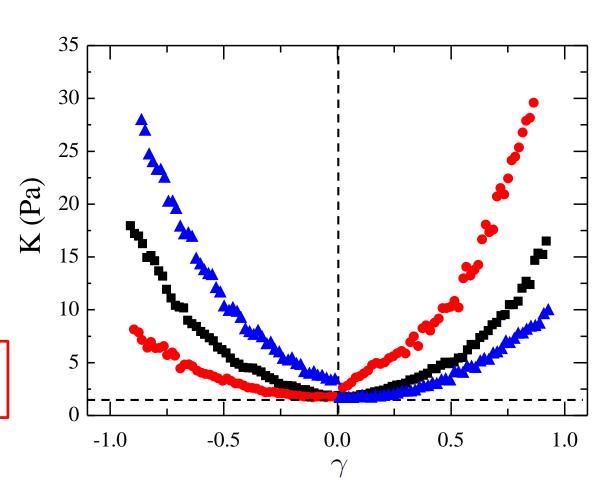


Deformation History

- 1. No Training
- 2. + 4 Pa Training
- 3. 4 Pa Training

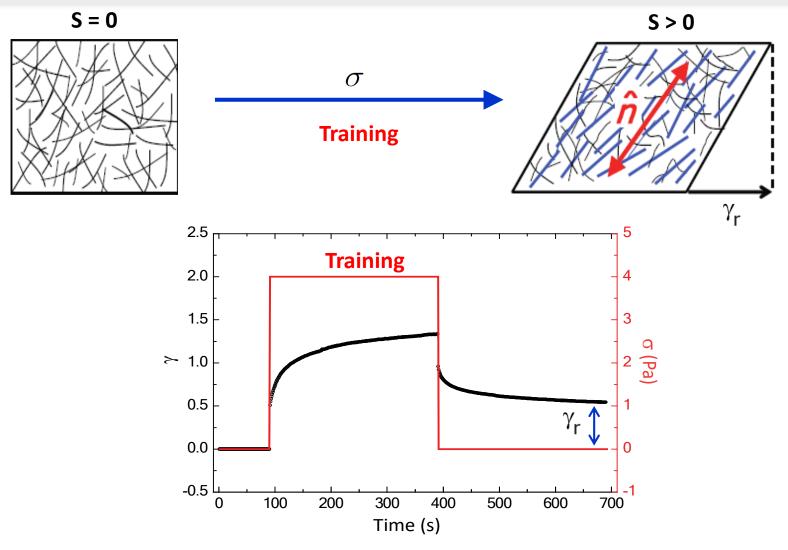
Stress history can be reversibly encoded.

Mechanical memory



Mechano-memory Frozen nematic order 20 **▼** | $\gamma > 0$ **Trained** γ < 0 **Trained** (anisotropic), 15 Stretch Bend 10 dominated dominated 5 **Untrained (isotropic)** 0 -0.2 0.2 0.4 0.6 8.0 -0.4 -0.8 -0.6 0 Majumdar et al., Soft Matter (2018)

Training stress reorganizes the network



$$S = \frac{\gamma_{\rm r}}{\sqrt{\gamma_{\rm r}^2 + 4}}$$

Residual strain

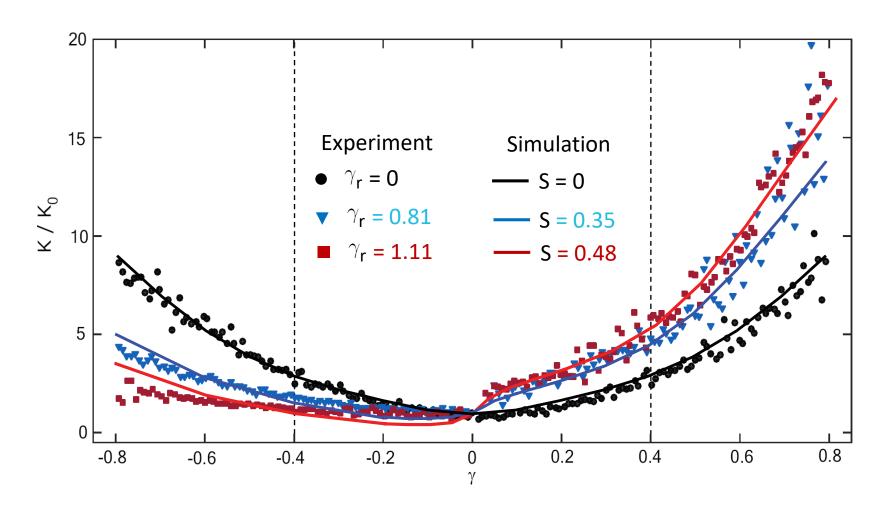


Nematic order

Simple model captures the mechanical response with no fitting parameter

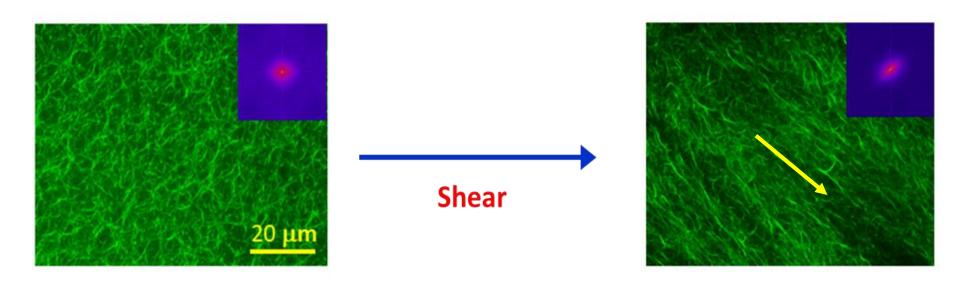
$$S = \frac{\gamma_{\rm r}}{\sqrt{\gamma_{\rm r}^2 + 4}}$$

Nematic order estimated from experimentally measured residual strain

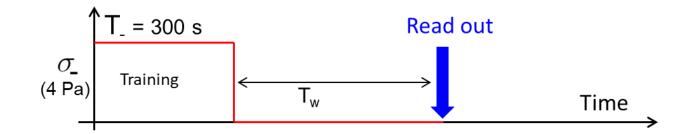


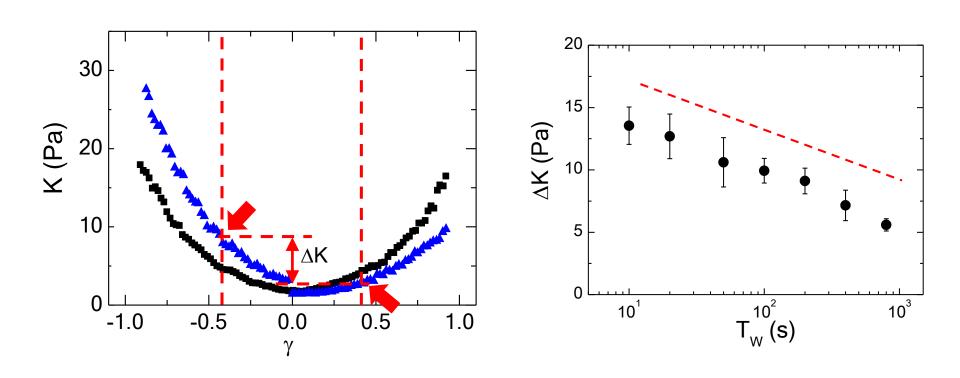
Experimentally observed frozen nematic order

Unidirectional stress (training stress) Frozen nematic order



Mechano-memory: Dynamics

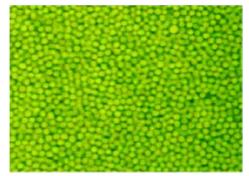




Slow non-exponential relaxation

Wide range of disordered condensed matter systems show slow non-exponential relaxations:

Colloidal glass



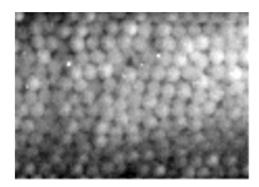
Ghosh et al. (2010)

Crumpled sheet



Lahini et al. (2017)

Granular materials

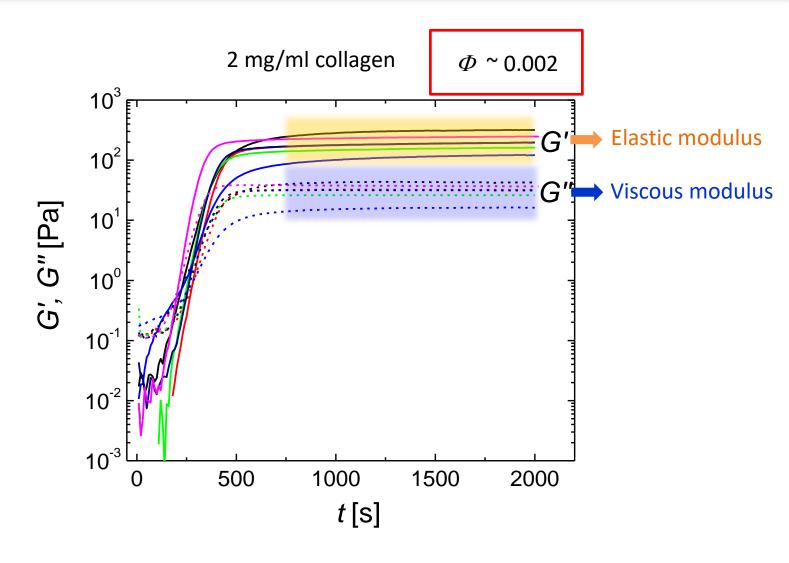


Current / voltage relaxation in CDW conductors

Thermal expansion of glassy polymers

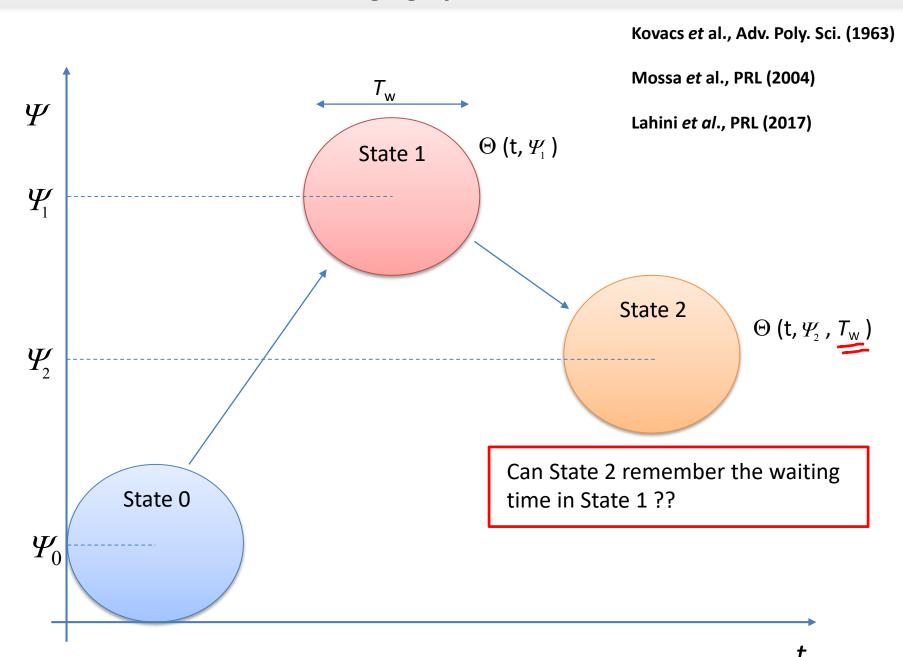
Flux creep in superconductors

Collagen networks: polymerization dynamics

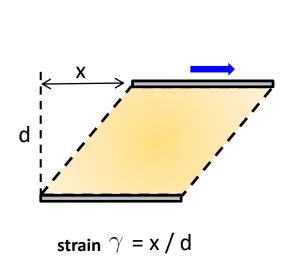


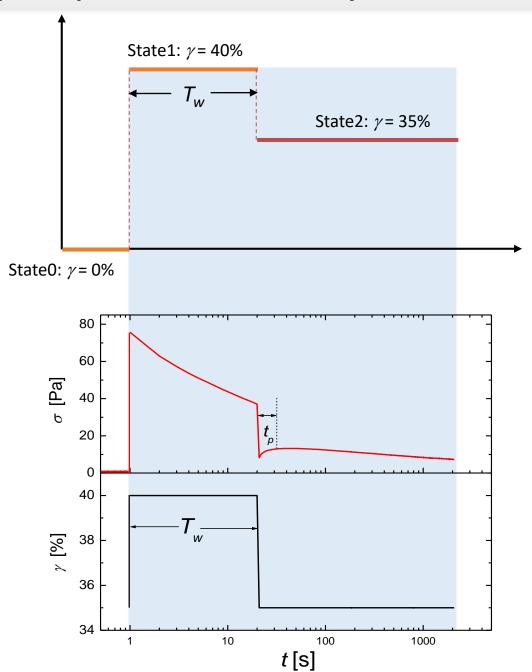
Does it show time dependent properties of glassy systems?

Non-monotonic aging dynamics: Kocavs Effect

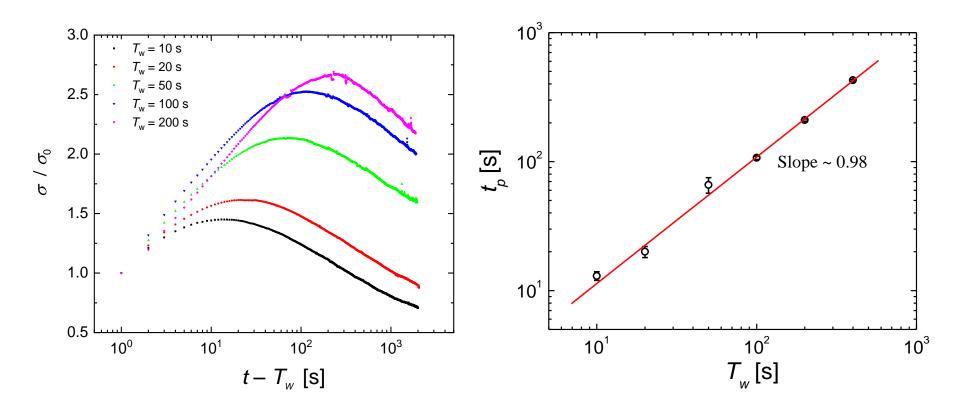


Strain protocol to probe pulse duration memory



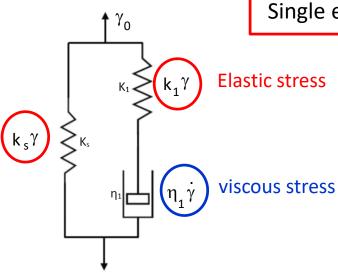


Waiting time can be predicted from peak position



Modelling non-monotonic relaxation

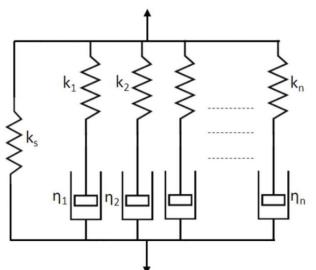




$$\sigma(t) = \sigma_s + \sigma_0 e^{-\frac{t}{\tau}}$$

$$\tau = 1/\lambda = \eta_1/k_1$$

Multi-elements Maxwell's model



$$\sigma(t) = \sigma_{\rm S} + \Sigma \sigma_{i} e^{-t/\tau_{i}}$$

Two elements can give non-monotonic relaxation

With M.K. Firoz (VSP student, RRI)

Competing logarithmic relaxations

Modelling non-monotonic relaxation

$$X(t) = X_0 \sum_{i} e^{-\lambda_i t}$$

$$X(t) = X_0 \int_{\lambda_{\min}}^{\lambda_{\max}} d\lambda P(\lambda) e^{-\lambda t}$$

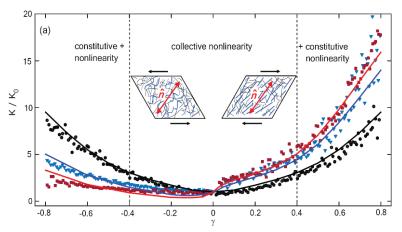
$$P(\lambda) \sim 1/\lambda$$
 Logarithmic relaxation

Slow modes dominate the distribution

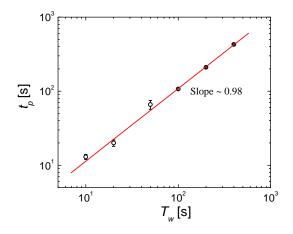
How do we understand it physically?

Summary and outlook

- Disordered bio-polymer networks show reversible mechano-memory.
- Such memory arises from induced nematic order.

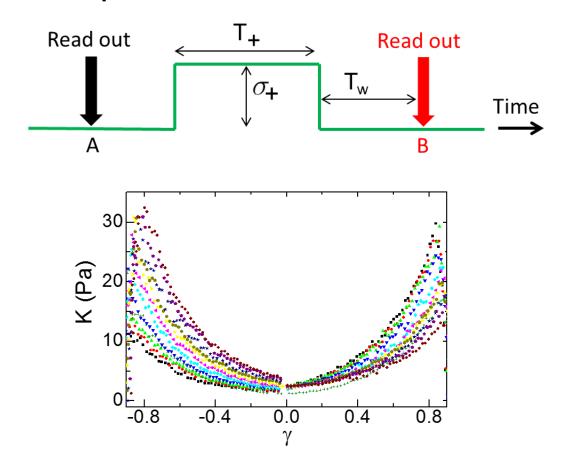


Slow, logarithmic relaxation dynamics give rise to pulse duration memory.



Summary and outlook

Force induced adaptations in bulk soft materials



- Origin, spatial distribution of fast and slow modes in our system ??
- Implications in cell mechanics.

Acknowledgements

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Thank you for your attention

Continuous tuning of mechanical response by mechano-memory

