

Light-cone spreading of perturbations in the Heisenberg spin chain

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- Chaos in classical systems.
- Propagation of chaos in many-body systems.
- The out-of-time-ordered-correlator (OTOC) as a probe of quantum chaos and of ballistic propagation of information.
- Numerical results for classical Heisenberg spin-chain.

Chaos in classical systems

A butterfly fluttering its wings in Beijing causes a tornado in Kansas!!



Classical Chaos - Sensitivity to initial conditions

Consider a dynamical system with deterministic time-evolution

$$\frac{dz_x}{dt} = f_x(\mathbf{z}), \quad x = 1, 2, \dots, N.$$

Consider two different initial conditions

$$z^A(0) \text{ and } z^B(0) = z^A(0) + \delta z(0).$$

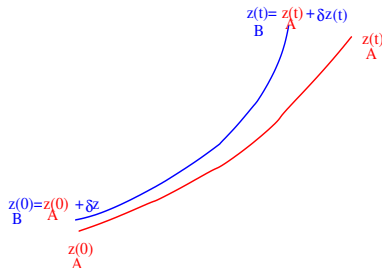
After time t let

$$\delta z(t) = z^B(t) - z^A(t)$$

Definition of chaos — at large t

$$\lim_{\delta z(0) \rightarrow 0} \frac{|\delta z(t)|}{|\delta z(0)|} \sim e^{\lambda(\mathbf{z})t}$$

with $\lambda > 0$ — LYAPUNOV EXPONENT



Propagation of chaos

Suppose that the initial perturbation is localized in space

$$\delta z_x(0) = \epsilon \delta_{x,0} .$$

We expect that the differences in phase-space variables at any point in space should eventually diverge with time. Thus

$$\lim_{\delta z_0(0) \rightarrow 0} \frac{|\delta z_x(t)|}{|\delta z_0(0)|} \sim \phi(x, t) e^{\lambda(z)t} .$$

Some questions:

- How long does it take before the perturbation is felt at point x ?
- How does the spatio-temporal evolution of the perturbation take place?

Spread of perturbations in a Hamiltonian system

Consider a chain of coupled oscillators (e.g coupled pendula).



$z_x = (q_x, p_x)$ — (position, momentum) of x^{th} particle. $[x = -N/2, \dots, N/2]$.

- Disturb system locally.
- The spread of the **perturbation** can be expressed as a **Poisson bracket**:

$$\frac{\partial q_x(t)}{\partial q_0(0)} = \{q_x(t), p_0(0)\} \quad \{A, B\} = \sum_x \frac{\partial A}{\partial q_x} \frac{\partial B}{\partial p_x} - \frac{\partial A}{\partial p_x} \frac{\partial B}{\partial q_x}$$

- Linear response $\langle \partial q_x(t) / \partial q_0(0) \rangle \sim \langle q_x(t) q_0(0) \rangle$ — correlation functions
- Chaos and non-linear response — $\left\langle \left(\frac{\partial q_x(t)}{\partial q_0(0)} \right)^2 \right\rangle$

Chaos in quantum systems

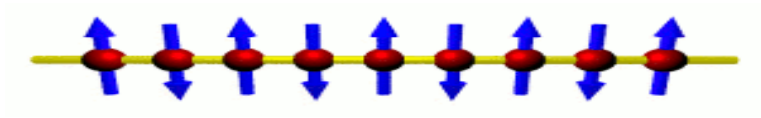
- $z \rightarrow \psi$ — No chaos since linear dynamics.
- Usual characterization of quantum chaos: level-spacing statistics
- Maldacena, Shenker, Stanford — replace Poisson bracket by commutator. Hence look at

$$-\langle [q_x(t), p_0(0)]^2 \rangle / \hbar^2 ,$$

$\langle \dots \rangle$ represents expectation over pure or thermal state.

- No long time divergence but perhaps the short-time growth shows a exponential growth regime and one can extract a Lyapunov exponent.
- Ehrenfest time and scrambling time.
- Ballistic spread characterized by butterfly velocity, light-cone velocity, Lieb-Robinson velocity.

Classical Heisenberg spin chains



Heisenberg spins $\mathbf{S} = (S^x, S^y, S^z)$ — unit three-dimensional vectors.
Hamiltonian given by

$$H = - \sum_{\ell=1}^N \mathbf{S}_{\ell} \cdot \mathbf{S}_{\ell+1} .$$

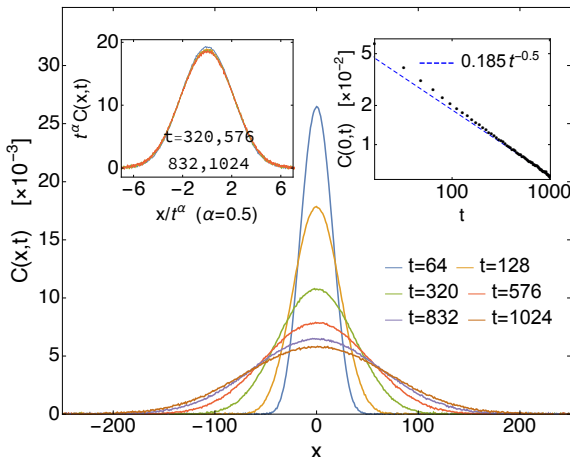
Equations of motion [Structure of usual Hamiltonian dynamics]

$$\dot{\mathbf{S}}_{\ell} = \mathbf{S}_{\ell} \times \mathbf{B}_{\ell}^{\text{eff}} = \{\mathbf{S}_{\ell}, H\} .$$

Linear response at high temperatures

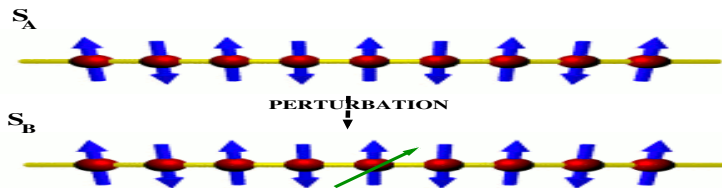
Expect diffusive hydrodynamic equations.

Spreading of a spin perturbation $C_{ss}(x, t) = \langle \mathbf{S}_x(t) \cdot \mathbf{S}_0(0) \rangle_{eq}$.



Energy correlations are similar.

Chaos and ballistic propagation in the XXX Heisenberg chain at high temperature



- Consider finite perturbation $\delta \mathbf{S}_0(0) = \epsilon \hat{\mathbf{n}}$ and look at

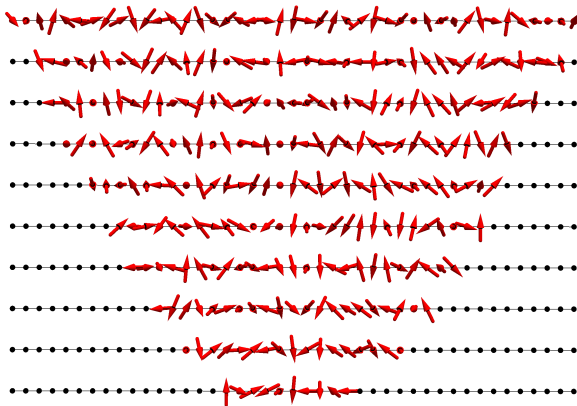
$$D(x, t) = \langle [\delta \mathbf{S}_x(t)]^2 \rangle / 2 = 1 - \langle \mathbf{S}_x^A(t) \cdot \mathbf{S}_x^B(t) \rangle$$

- Clearly $D(x, t) \rightarrow 1$ at large times
- However in the limit $\epsilon \rightarrow 0$

$$\left\langle \left(\frac{\partial \mathbf{S}_x(t)}{\partial \mathbf{S}_0(0)} \right)^2 \right\rangle = \frac{D(x, t)}{\epsilon^2} \sim e^{2\lambda t}$$

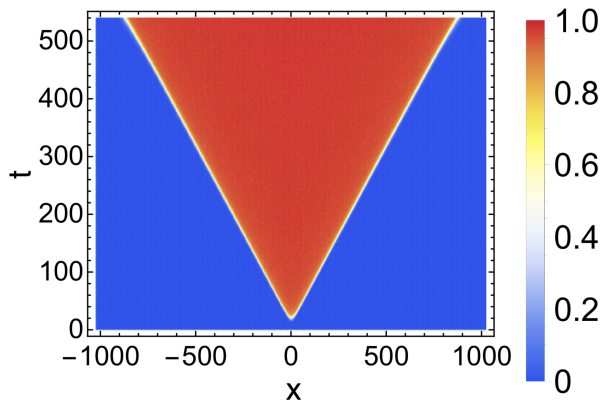
Spreading of perturbation for single realization

Space-time evolution of localized perturbation.



Light-cone evolution of $D(x, t)$

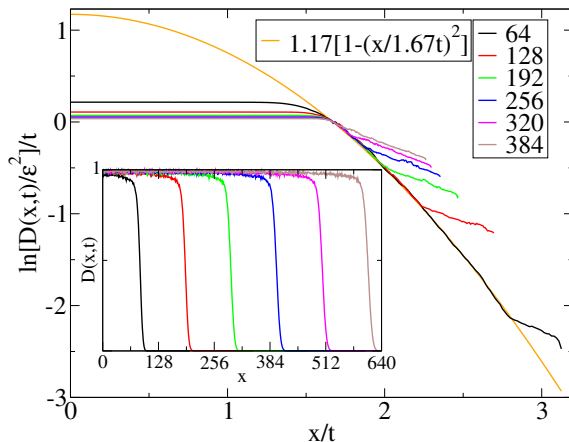
System of $N = 2048$ spins, average over random initial conditions.



The signal decays exponentially with distance away from the light cone $x = v_b t$, with $v \approx 1.67$.

Evolution of $D(x, t)$

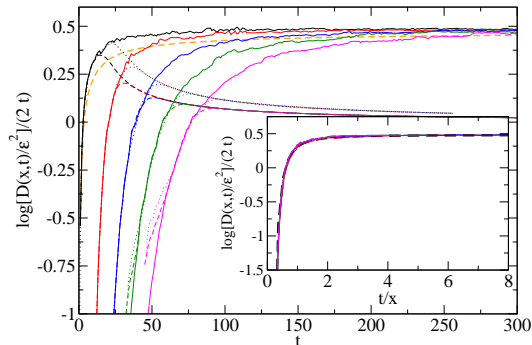
Plot of $D(x, t)$ profile at different times



From numerics, the front behaves as $D(x, t) \sim \epsilon^2 e^{2\lambda t[1-(x/v_b t)^2]}$

Behaviour of $D(x, t)$

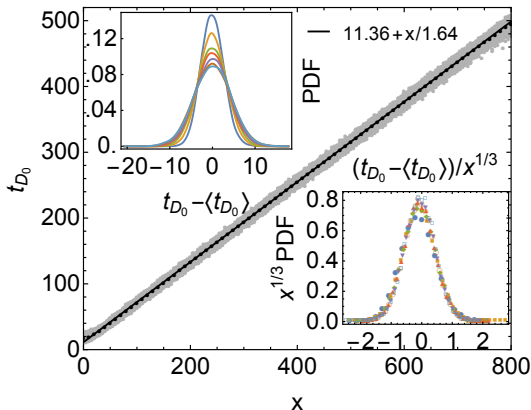
Growth of $D(x, t)$ profile at different positions – from linear and non-linear dynamics.



From numerics, the signal grows as $D(x, t) \sim \epsilon^2 e^{2\lambda t[1-(x/v_b t)^2]}$, with $\lambda \approx 0.5$, $v_b \approx 1.67$.

Behaviour of $D(x, t)$

Estimating the butterfly speed - At what time does the signal reach some threshold value D_0 at a given point x .



$x^{1/3}$ broadening of the front.

Connection to KPZ surface growth

- Exponential growth at early times (Butterfly effect)
- Ballistic front (Butterfly velocity)
- Saturation at long times
- Broadening of growth front - possibly related to KPZ

Define

$$h(x, t) = \lim_{\epsilon \rightarrow 0} \log[\delta \mathbf{S}^2(x, t)/2\epsilon^2]/2$$

Our observations:

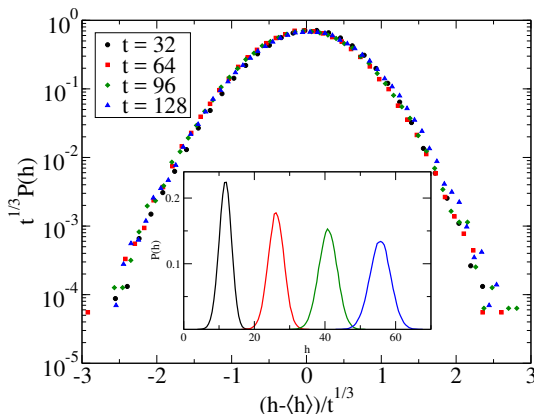
$$h(x, t) = at + bx^2/t + t^{1/3}\eta$$

Numerically $a \approx \lambda_L$ and $b \approx -\lambda_L/v_b^2$.

This suggests that $h(x, t)$ is like the KPZ height field.

Distribution of η — Tracy-Widom ?

Connections to KPZ



Clear $t^{1/3}$ -scaling but no signs yet of Tracy-Widom.

Conclusions

- Investigated the classical version of “Out-of-Time-Ordered-Correlator” in the Heisenberg spin chain.

—Ballistic front propagation in a diffusive system, exponential growth, broadening of front and possible connections to KPZ universality.

- $D(x, t) \sim e^{2\mu(x/t)t}$, $\mu(v) = \lambda[1 - (v/v_b)^2]$ — velocity-dependent Lyapunov exponent (R. J. Deissler and K. Kaneko, 1987)
Integrable systems: $\mu(v) \sim -(1 - v/v_b)^{3/2}$.

- What physical observable ? Information propagation, Loschmidt echo,...

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