# Relative Canonical Bundles for families of Calabi-Yau manifolds, twisted Hodge Bundles, and Positivity 

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Marburg

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Hodge metric on $f_{*} \Omega_{\mathcal{X} / S}^{n}$
$L^{2}$-inner product/norm of canonical forms $\eta$ :

$$
(-1)^{n} \int_{\mathcal{X}_{s}} \eta \wedge \bar{\eta}
$$

Hodge bundles $R^{n-p} f_{*} \Omega_{\mathcal{X} / S}^{p}$

## Griffiths: Curvature of Hodge bundle

Kodaira-Spencer classes: obstructions against splitting of

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\rho: T_{S, s_{0}} & \longrightarrow & H^{1}\left(X, \mathcal{T}_{X}\right) \\
\frac{\partial}{\partial s} & \mapsto & {\left[A_{s}\right]=\left[A_{\bar{\beta}}^{\alpha}(z) \frac{\partial}{\partial z^{\alpha}} \overline{d z^{\beta}}\right]}
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The cup product

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A \otimes \phi & \mapsto A \cup \phi \\
H^{1}\left(\mathcal{X}_{s}, \mathcal{T}_{\mathcal{X}_{s}}\right) \otimes H^{0}\left(\mathcal{X}_{s}, \Omega_{\mathcal{X}_{s}}^{n}\right) & \rightarrow H^{1}\left(\mathcal{X}_{s}, \Omega_{\mathcal{X}_{s}}^{n-1}\right)
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induces

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\sigma_{0}: f_{*} \Omega_{\mathcal{X} / S}^{n} \rightarrow R^{1} f_{*} \Omega_{\mathcal{X} / S}^{n-1} \otimes \mathcal{T}_{\mathcal{S}}^{\vee}
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## Curvature $\Theta$ of $f_{*} \Omega_{\mathcal{X} / S}^{n}$

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\left(\Theta e, e^{\prime}\right)=\left(\sigma_{0} e, \sigma_{0} e^{\prime}\right)
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R\left(\partial / \partial s, \partial / \partial s^{\prime}, e, e^{\prime}\right)=\left(A \cup e, A^{\prime} \cup e^{\prime}\right)
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Notation

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\begin{aligned}
\omega_{X}^{n} & =g d V \text { Ricci-flat volume form } \\
0 & =\operatorname{Ric}\left(\omega_{X}\right)=-\sqrt{-1} \partial \bar{\partial} \log \left(\omega_{X}^{n}\right)
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Let $X$ Kähler.
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\lambda_{\mathcal{X} / \mathcal{S}} \in R^{1} f_{*}\left(\Omega_{\mathcal{X} / \mathcal{S}}^{1}\right)(S)
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Yau's theorem states the unique existence of a unique Ricci-flat Kähler form $\omega_{X}$ in any Kähler class $\lambda_{X}$.

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Known:

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\omega_{W P} \text { is Kähler. }
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G. Sch. '85 for hol. symplectic manifolds Tian '86 for general CY manifolds

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\left(A_{s}, A_{s}\right)_{W P}=\frac{\int_{\mathcal{X}_{s}}\left(A_{s} \cdot A_{\bar{s}}\right) \phi \wedge \bar{\phi}}{\int_{\mathcal{X}_{s}} \phi \wedge \bar{\phi}}=\frac{\int_{\mathcal{X}_{s}}\left(A_{s} \cup \phi\right) \wedge\left(A_{\bar{s}} \cup \bar{\phi}\right)}{\int_{\mathcal{X}_{s}} \phi \wedge \bar{\phi}}
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i.e. the Weil-Petersson form is the curvature of the Hodge metric on $f_{*}\left(\mathcal{K}_{\mathcal{X} / S}\right)$

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Now turn to $\mathcal{K}_{\mathcal{X} / \mathcal{S}}$ rather than $f_{*}\left(\mathcal{K}_{\mathcal{X} / S}\right)$.

## Curvature of $\mathcal{K}_{\mathcal{X} / \mathcal{S}}$

Let $f: \mathcal{X} \rightarrow S$ be a (polarized) family of Calabi-Yau manifolds with family of Ricci-flat relative volume forms

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\Theta_{\mathcal{X}}=\frac{1}{\operatorname{vol}\left(\mathcal{X}_{s}\right)} f^{*} \omega_{W P}
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( $f: \mathcal{X} \rightarrow S, \lambda_{\mathcal{X} / S}$ ) a polarized family of Calabi-Yau manifolds, and $\omega_{\mathcal{X}_{s}} \in \lambda_{\mathcal{X}_{s}}$ the Kähler-Einstein forms. Then locally with respect to $S$ there exists a $d$-closed ( 1,1 )-form $\omega_{\mathcal{X}}$ on the total space $\mathcal{X}$ such that

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\omega_{\mathcal{X}} \mid \mathcal{X}_{s}=\omega_{\mathcal{X}_{s}} .
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If the polarization can be represented by a closed, real ( 1,1 )-form, then $\omega_{\mathcal{X}}$ can be chosen globally as a (1,1)-form.
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\int_{\mathcal{X} / S} \omega_{\mathcal{X}}^{n+1}=0 .
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## Validity of the assumptions

## Cheeger '70, Cheeger - Yau '80

The Green's function is bounded, if the diameter (of the fibers $\mathcal{X}_{s}$ ) is bounded from above.

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The diameter is bounded for polarized families of Calabi-Yau manifolds, under mild assumptions for the type of degeneration.

## Twisted Hodge bundles, canonically polarized case

## Definition

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## Notation

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\omega_{X} \text { Kähler form } \\
\omega_{X}^{n}=g d V \text { volume form. } \\
\operatorname{Ric}\left(\omega_{X}\right)=-\sqrt{-1} \partial \bar{\partial} \log \left(\omega^{n}\right)=-\sqrt{-1} \partial \bar{\partial} \log g \\
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\left(\square_{s}+1\right)\left(v_{i}, v_{j}\right)=\left(A_{i}, A_{j}\right) w_{P} \\
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\begin{align*}
R(A, \bar{A}, \psi, \bar{\psi})== & m \int_{\mathcal{X}_{s}^{\prime}}(\square+1)^{-1}(A \cdot \bar{A}) \cdot(\psi \cdot \bar{\psi}) g d V \\
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$$

The only contribution in (1), which may be negative, originates from the harmonic parts in the third term. It equals

$$
-\int_{\mathcal{X}_{s}} H(A \cup \bar{\psi}) \overline{H\left(A_{j} \cup \bar{\psi}\right)} g d V .
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## Previous results

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The moduli stack of canonically polarized manifolds is (Kobayashi-)hyperbolic.

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## Corollary

The locally free sheaf $f_{*} \mathcal{K}_{\mathcal{X} / S}^{\otimes(m+1)}$ is Nakano-positive.

Using the above positivity of the relative canonical bundle, we see that the Corollary follows from

Theorem (Bo Berndtsson '09)
Let $L$ be a positive line bundle on $\mathcal{X}$, then $f_{*}\left(\mathcal{K}_{\mathcal{X} / \mathcal{S}} \otimes L\right)$ is
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Further results by Sh. Takayama - Chr. Mourougane and K. Liu X. Yang.

## Twisted Hodge bundles - general case

Let $f: \mathcal{X} \rightarrow S$ be a holomorphic family of compact complex manifolds, and $(L, h)$ be a relatively positive line bundle on $\mathcal{X}$.

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## Corollary

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Use Lie-derivatives for differential forms with values in hermitian line bundles.

## Methods (Calabi-Yau manifolds)

## Calabi's Theorem '57

Holmorphic 1-forms and holomorphic vector fields on Calabi-Yau manifolds are parallel.

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f: \mathcal{X} \longrightarrow S \\
(z, s) \mapsto s \\
z=\left(z^{1}, \ldots, z^{n}\right) \\
s=\left(s^{1}, \ldots, s^{r}\right) \\
\text { components } z^{\alpha}, s^{i}
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Given $\omega_{\mathcal{X}}$ like in proposition

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## Lemma

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Lemma

$$
\Theta_{i \bar{\beta}} g^{\bar{\beta} \alpha} \partial_{\alpha}=\bar{\partial}^{*} A_{i \bar{\beta}}^{\alpha} \partial_{\alpha} d z^{\bar{\beta}}
$$

Let

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\chi_{i \bar{\jmath}}:=\left\langle v_{i}, v_{j}\right\rangle_{\Theta}=\Theta_{i \bar{\jmath}}-a_{i}^{\alpha} \Theta_{\alpha \bar{\jmath}}-\Theta_{i \bar{\beta}} a_{\bar{\jmath}}^{\bar{\beta}}
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and $\Theta_{i \bar{\jmath}}$ must be fiberwise constant. The value of $\Theta_{i \bar{\jmath}}=\Theta_{i \bar{\jmath}}(s)$ is determined by integration over $\mathcal{X}_{s}$ according to the following Lemma.

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$\Rightarrow$ Fact.

## On the other hand (cf. proposition)

$$
0=\int_{\mathcal{X} / S} \omega_{\mathcal{X}}^{n+1}=\sqrt{-1}\left(\int_{\mathcal{X} / S} \varphi_{i \bar{\jmath}} g d V\right) d s^{i} \wedge d s^{\bar{\jmath}}
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$$
G_{s}\left(A_{i} \cdot A_{\bar{\jmath}}\right) \geq-c \cdot \operatorname{vol}\left(\mathcal{X}_{s}\right) \Theta_{i \bar{\jmath}}
$$

(in the sense of matrices/hermitian forms)

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For a suitable constant $c^{\prime}>0$ the form

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Namely, with $c^{\prime}=c \cdot \operatorname{vol}\left(\mathcal{X}_{s}\right)+1$

$$
\begin{aligned}
\widetilde{\omega}_{\mathcal{X}}^{n+1} & =\left(\omega_{\mathcal{X}}+c^{\prime} f^{*} \omega_{W P}\right)^{n+1}=\sqrt{-1}\left(\varphi_{i \bar{\jmath}}+c^{\prime} \Theta_{i \bar{\jmath}}\right) d s^{i} \wedge d s^{\bar{\jmath}} \\
& \geq \sqrt{-1} \Theta_{i \bar{\jmath}} d s^{i} \wedge d s^{\bar{j}}
\end{aligned}
$$

whereas $\widetilde{\omega}_{\mathcal{X}} \mid \mathcal{X}_{s}=\omega_{\mathcal{X}_{s}}>0$.

