



Subroto Mukerjee Department of Physics, Indian Institute of Science

Current frontiers in condensed matter ICTS Jun 29 2016





#### Sayonee Ray

Vijay Shenoy

Funding: ISF-UGC Indo-Israeli grant arxiv:1603.09478

How is superconductivity destroyed by a boost?

How is superconductivity destroyed by a boost?

Superconducting gap decreases with boost and goes to zero at  $v_c$ 

How is superconductivity destroyed by a boost?

Superconducting gap decreases with boost and goes to zero at  $v_c$ 

Landau critical velocity

$$v_c = \operatorname{Min.}\left(v_s, \left(\left[\sqrt{\Delta^2 + \mu^2} - \mu\right]/m\right)^{1/2}\right)$$

Sound velocity of bosonic collective mode

Fermionic (single-particle) excitations

Baym and Pethick, Phys. Rev. A 86 023602 (2012)

What happens if the gap is zero to begin with?

What happens if the gap is zero to begin with?

1D superconductor at T=0

No long-range order (Mermin-Wagner-Coleman theorem)

Algebraic long-range order

What happens if the gap is zero to begin with?

1D superconductor at T=0

No long-range order (Mermin-Wagner-Coleman theorem)

Algebraic long-range order

$$\langle O_{SU} \rangle = 0 \qquad \langle O_{SU}(x) O_{SU}(x') \rangle \sim 1/|x - x'|^{\eta}$$

What happens if the gap is zero to begin with?

1D superconductor at T=0

No long-range order (Mermin-Wagner-Coleman theorem)

Algebraic long-range order

$$\langle O_{SU} \rangle = 0 \qquad \langle O_{SU}(x) O_{SU}(x') \rangle \sim 1/|x - x'|^{\eta}$$

What does a boost do to this kind of superconductor?

## What is known?

Mean-field

Boost destroys superconductivity via Clogston-Chandrasekhar type mechanism. Critical velocity < Landau critical velocity

T.-C. Wei and P. M. Goldbart, Phys. Rev. B 80, 134507 (2009).

Fluctuations

Interaction with walls can induce phase slips "dynamically" destroying superconductivity, finite  $\omega$  and T

T. Eggel, M. A. Cazalilla and M. Oshikawa, Phys. Rev. Lett.107, 275302 (2011).

Bosonization treatment taking into account quantum fluctuations

Bosonization treatment taking into account quantum fluctuations

Superconductivity can be strengthened upon applying a boost

Bosonization treatment taking into account quantum fluctuations

Superconductivity can be strengthened upon applying a boost

A boost can open or close gaps depending on whether the system has spinless or spin 1/2 fermions

Low energy theory for interacting 1D spinless Fermi system

$$\mathcal{H} = \frac{v}{2} \int dx \left[ K \Pi(x)^2 + \frac{1}{K} \left( \nabla \phi(x) \right)^2 \right]$$

Low energy theory for interacting 1D spinless Fermi system

$$\mathcal{H} = \frac{v}{2} \int dx \left[ K \Pi(x)^2 + \frac{1}{K} \left( \nabla \phi(x) \right)^2 \right]$$
$$\mathcal{L} = \frac{1}{2K} \int dx dt \left[ \frac{1}{v} \left( \partial_t \phi \right)^2 - v \left( \partial_x \phi \right)^2 \right]$$

Low energy theory for interacting 1D spinless Fermi system

$$\mathcal{H} = \frac{v}{2} \int dx \left[ K \Pi(x)^2 + \frac{1}{K} \left( \nabla \phi(x) \right)^2 \right]$$
$$\mathcal{L} = \frac{1}{2K} \int dx dt \left[ \frac{1}{v} \left( \partial_t \phi \right)^2 - v \left( \partial_x \phi \right)^2 \right]$$

*K*-Luttinger parameter  $\langle O_{SU}(x)O_{SU}(x')\rangle \sim 1/|x-x'|^{1/K}$ 

 $\langle O_{CDW}(x)O_{CDW}(x')\rangle \sim 1/|x-x'|^K$ 

Low energy theory for interacting 1D spinless Fermi system

$$\mathcal{H} = \frac{v}{2} \int dx \left[ K \Pi(x)^2 + \frac{1}{K} \left( \nabla \phi(x) \right)^2 \right]$$
$$\mathcal{L} = \frac{1}{2K} \int dx dt \left[ \frac{1}{v} \left( \partial_t \phi \right)^2 - v \left( \partial_x \phi \right)^2 \right]$$
$$K \text{- Luttinger parameter}$$

 $\langle O_{SU}(x) O_{SU}(x') \rangle \sim 1/|x - x'|^{1/K}$ 

 $\langle O_{CDW}(x)O_{CDW}(x')\rangle \sim 1/|x-x'|^K$ 

K > 1 attractive interactions (dominant SU order) K < 1 repulsive interactions (dominant CDW order)



$$v = \sqrt{\left(v_F + \frac{g_4}{2\pi}\right)^2 - \left(\frac{g_2}{2\pi}\right)^2}$$

$$K = \sqrt{\frac{v^F - \left(\frac{g_2}{2\pi} - \frac{g_4}{2\pi}\right)}{v^F + \left(\frac{g_2}{2\pi} + \frac{g_4}{2\pi}\right)}}$$

Ref: Quantum physics in one dimension, Giamarchi

Momentum:  $k \rightarrow k + u$ 

Momentum:  $k \rightarrow k + u$ 

Different Fermi speeds for left and right movers

$$v_L^F(u) = v^F(u) - w(u); v_R^F(u) = v^F(u) + w(u)$$

Momentum:  $k \rightarrow k + u$ 

Different Fermi speeds for left and right movers

$$v_L^F(u) = v^F(u) - w(u); v_R^F(u) = v^F(u) + w(u)$$



$$\mathcal{H} = \int dx \left\{ \frac{v(u)}{2} \left[ K(u) \Pi^2 + \frac{1}{K(u)} \left( \partial_x \phi \right)^2 \right] - w(u) \left[ \left( \Pi \partial_x \phi \right) + \left( \partial_x \phi \Pi \right) \right] \right\}$$

$$\mathcal{H} = \int dx \left\{ \frac{v(u)}{2} \left[ K(u) \Pi^2 + \frac{1}{K(u)} \left( \partial_x \phi \right)^2 \right] - w(u) \left[ \left( \Pi \partial_x \phi \right) + \left( \partial_x \phi \Pi \right) \right] \right\}$$

$$v(u) = \sqrt{\left(v^F(u) + \frac{g_4}{2\pi}\right)^2 - \left(\frac{g_2}{2\pi}\right)^2},$$
  

$$K(u) = \sqrt{\frac{v^F(u) - \left(\frac{g_2}{2\pi} - \frac{g_4}{2\pi}\right)}{v^F(u) + \left(\frac{g_2}{2\pi} + \frac{g_4}{2\pi}\right)}}$$

$$\mathcal{H} = \int dx \left\{ \frac{v(u)}{2} \left[ K(u) \Pi^2 + \frac{1}{K(u)} \left( \partial_x \phi \right)^2 \right] - w(u) \left[ \left( \Pi \partial_x \phi \right) + \left( \partial_x \phi \Pi \right) \right] \right\}$$

$$v(u) = \sqrt{\left(v^F(u) + \frac{g_4}{2\pi}\right)^2 - \left(\frac{g_2}{2\pi}\right)^2},$$
  

$$K(u) = \sqrt{\frac{v^F(u) - \left(\frac{g_2}{2\pi} - \frac{g_4}{2\pi}\right)}{v^F(u) + \left(\frac{g_2}{2\pi} + \frac{g_4}{2\pi}\right)}}$$

$$v_L^F(u) = v^F(u) - w(u); v_R^F(u) = v^F(u) + w(u)$$

$$\mathcal{H} = \int dx \left\{ \frac{v(u)}{2} \left[ K(u) \Pi^2 + \frac{1}{K(u)} \left( \partial_x \phi \right)^2 \right] - w(u) \left[ \left( \Pi \partial_x \phi \right) + \left( \partial_x \phi \Pi \right) \right] \right\}$$

$$v(u) = \sqrt{\left(v^F(u) + \frac{g_4}{2\pi}\right)^2 - \left(\frac{g_2}{2\pi}\right)^2},$$
  

$$K(u) = \sqrt{\frac{v^F(u) - \left(\frac{g_2}{2\pi} - \frac{g_4}{2\pi}\right)}{v^F(u) + \left(\frac{g_2}{2\pi} + \frac{g_4}{2\pi}\right)}}$$

$$v_L^F(u) = v^F(u) - w(u); v_R^F(u) = v^F(u) + w(u)$$

g's unchanged by boost because they only depend on momentum differences

$$\mathcal{H} = \int dx \left\{ \frac{v(u)}{2} \left[ K(u) \Pi^2 + \frac{1}{K(u)} \left( \partial_x \phi \right)^2 \right] - w(u) \left[ \left( \Pi \partial_x \phi \right) + \left( \partial_x \phi \Pi \right) \right] \right\}$$

$$\mathcal{H} = \int dx \left\{ \frac{v(u)}{2} \left[ \frac{K(u)\Pi^2 + \frac{1}{K(u)} \left(\partial_x \phi\right)^2 \right] - w(u) \left[ \left(\Pi \partial_x \phi\right) + \left(\partial_x \phi \Pi\right) \right] \right\}$$

$$\mathcal{L} = \frac{1}{2K(u)} \int dx dt \left[ \frac{1}{v(u)} \left( \partial_t \tilde{\phi} \right)^2 - v(u) \left( \partial_x \tilde{\phi} \right)^2 \right]$$

$$\mathcal{H} = \int dx \left\{ \frac{v(u)}{2} \left[ K(u) \Pi^2 + \frac{1}{K(u)} \left( \partial_x \phi \right)^2 \right] - w(u) \left[ \left( \Pi \partial_x \phi \right) + \left( \partial_x \phi \Pi \right) \right] \right\}$$

$$\mathcal{L} = \frac{1}{2K(u)} \int dx dt \left[ \frac{1}{v(u)} \left( \partial_t \tilde{\phi} \right)^2 - v(u) \left( \partial_x \tilde{\phi} \right)^2 \right]$$

$$\tilde{\phi}(x,t) = \phi(x+w(u)t,t)$$

Galilean transformation with speed w(u)

$$\mathcal{H} = \int dx \left\{ \frac{v(u)}{2} \left[ K(u) \Pi^2 + \frac{1}{K(u)} \left( \partial_x \phi \right)^2 \right] - w(u) \left[ \left( \Pi \partial_x \phi \right) + \left( \partial_x \phi \Pi \right) \right] \right\}$$

$$\mathcal{L} = \frac{1}{2K(u)} \int dx dt \left[ \frac{1}{v(u)} \left( \partial_t \tilde{\phi} \right)^2 - v(u) \left( \partial_x \tilde{\phi} \right)^2 \right]$$

$$\tilde{\phi}(x,t) = \phi(x+w(u)t,t)$$

Galilean transformation with speed w(u)

$$\langle O_{SU}(x)O_{SU}(x')\rangle \sim 1/|x-x'|^{1/K(u)}$$

$$\mathcal{H} = \int dx \left\{ \frac{v(u)}{2} \left[ K(u) \Pi^2 + \frac{1}{K(u)} \left( \partial_x \phi \right)^2 \right] - w(u) \left[ \left( \Pi \partial_x \phi \right) + \left( \partial_x \phi \Pi \right) \right] \right\}$$

$$\mathcal{L} = \frac{1}{2K(u)} \int dx dt \left[ \frac{1}{v(u)} \left( \partial_t \tilde{\phi} \right)^2 - v(u) \left( \partial_x \tilde{\phi} \right)^2 \right]$$

$$\tilde{\phi}(x,t) = \phi(x+w(u)t,t)$$

Galilean transformation with speed w(u)

$$\langle O_{SU}(x)O_{SU}(x')\rangle \sim 1/|x-x'|^{1/K(u)}$$

 $\langle O_{CDW}(x)O_{CDW}(x')\rangle \sim 1/|x-x'|^{K(u)}$
$$\langle O_{\rm SU}(x,t)O_{\rm SU}^{\dagger}(x',t')\rangle \sim e^{i2u(x-x')} \left(\frac{1}{\sqrt{[x-x'+w(u)(t-t')]^2+[v^F(u)]^2(t-t')^2}}\right)^{1/K(u)}$$

$$\langle O_{\rm CDW}(x,t)O_{\rm CDW}(x',t')\rangle \sim \cos\left[2k_F(x-x')\right] \left(\frac{1}{\sqrt{[x-x'+w(u)(t-t')]^2+[v^F(u)]^2(t-t')^2}}\right)^{K(u)}$$

$$\langle O_{\rm SU}(x,t)O_{\rm SU}^{\dagger}(x',t')\rangle \sim e^{i2u(x-x')} \left(\frac{1}{\sqrt{[x-x'+w(u)(t-t')]^2+[v^F(u)]^2(t-t')^2}}\right)^{1/K(u)}$$

$$\langle O_{\rm CDW}(x,t)O_{\rm CDW}(x',t')\rangle \sim \cos\left[2k_F(x-x')\right] \left(\frac{1}{\sqrt{[x-x'+w(u)(t-t')]^2+[v^F(u)]^2(t-t')^2}}\right)^{K(u)}$$

#### No conformal invariance

$$\langle O_{\rm SU}(x,t)O_{\rm SU}^{\dagger}(x',t')\rangle \sim e^{i2u(x-x')} \left(\frac{1}{\sqrt{[x-x'+w(u)(t-t')]^2+[v^F(u)]^2(t-t')^2}}\right)^{1/K(u)}$$

 $\langle O_{\rm CDW}(x,t)O_{\rm CDW}(x',t')\rangle \sim \cos\left[2k_F(x-x')\right] \left(\frac{1}{\sqrt{[x-x'+w(u)(t-t')]^2+[v^F(u)]^2(t-t')^2}}\right)^{K(u)}$ 

No conformal invariance

$$\tilde{\phi}(x,t) = \phi(x+w(u)t,t)$$

Galilean transformation with speed w(u)

$$\langle O_{\rm SU}(x,t)O_{\rm SU}^{\dagger}(x',t')\rangle \sim e^{i2u(x-x')} \left(\frac{1}{\sqrt{[x-x'+w(u)(t-t')]^2+[v^F(u)]^2(t-t')^2}}\right)^{1/K(u)}$$

 $\langle O_{\rm CDW}(x,t)O_{\rm CDW}(x',t')\rangle \sim \cos\left[2k_F(x-x')\right] \left(\frac{1}{\sqrt{[x-x'+w(u)(t-t')]^2+[v^F(u)]^2(t-t')^2}}\right)^{K(u)}$ 

#### No conformal invariance

$$\tilde{\phi}(x,t) = \phi(x+w(u)t,t)$$

Galilean transformation with speed w(u)

Conformal invariance restored in terms of  $\tilde{\phi}$ 

 $K(u) = K(v^F(u))$   $v^F(u) \text{ is an even function of } u$ 

 $K(u) = K(v^F(u))$  $v^F(u) \text{ is an even function of } u$ 

$$K(u) \approx K(0) + \frac{u^2}{2} \frac{dK}{dv^F(u)} \frac{d^2 v^F(u)}{du^2} \Big|_{u=0} + \dots$$
  
for  $u \ll k_F$ 

 $K(u) = K(v^F(u))$  $v^F(u) \text{ is an even function of } u$ 

$$K(u) \approx K(0) + \frac{u^2}{2} \frac{dK}{dv^F(u)} \frac{d^2 v^F(u)}{du^2} \Big|_{u=0} + \dots$$
  
for  $u \ll k_F$ 



 $K(u) = K(v^F(u))$  $v^F(u) \text{ is an even function of } u$ 

$$K(u) \approx K(0) + \frac{u^2}{2} \frac{dK}{dv^F(u)} \frac{d^2 v^F(u)}{du^2} \Big|_{u=0} + \dots$$
  
for  $u \ll k_F$ 



 $K(u) = K(v^F(u))$  $v^F(u) \text{ is an even function of } u$ 

$$\begin{split} K(u) &\approx K(0) + \frac{u^2}{2} \frac{dK}{dv^F(u)} \frac{d^2 v^F(u)}{du^2} \Big|_{u=0} + \dots \\ & \text{for } u \ll k_F \end{split}$$

$$\frac{dK}{dv^F} < 0 \text{ when } K > 1 \qquad \qquad \frac{dK}{dv^F} > 0 \text{ when } K < 1$$
$$\frac{d^2v^F}{du^2} = 0 \Rightarrow K(u) = K(0)$$

Boost has no effect when the dispersion is parabolic

Simplest dispersion on a lattice:  $\epsilon_k = -2t \cos k$ 

$$\frac{d^2 v^F}{du^2} < 0$$

Simplest dispersion on a lattice:  $\epsilon_k = -2t \cos k$ 



 $\Rightarrow K(u) > K(0)$  when K > 1

dominant superconductivity is strengthened upon boosting

Simplest dispersion on a lattice:  $\epsilon_k = -2t \cos k$ 



 $\Rightarrow K(u) > K(0)$  when K > 1

dominant superconductivity is strengthened upon boosting

 $\Rightarrow K(u) < K(0)$  when K < 1

dominant CDW order is strengthened upon boosting

Simplest dispersion on a lattice:  $\epsilon_k = -2t \cos k$ 



 $\Rightarrow K(u) > K(0)$  when K > 1

dominant superconductivity is strengthened upon boosting

 $\Rightarrow K(u) < K(0)$  when K < 1

dominant CDW order is strengthened upon boosting

The same effect for spin 1/2 fermions

Superconducting pairing susceptibility

$$\chi_{\text{pair}} = \lim_{\omega \to 0} \frac{1}{\Omega} \sum_{k} \frac{f(\xi_k) - f(-\xi_{-k})}{\omega - \xi(k) - \xi(-k) + i\delta}$$

Superconducting pairing susceptibility

$$\chi_{\text{pair}} = \lim_{\omega \to 0} \frac{1}{\Omega} \sum_{k} \frac{f(\xi_k) - f(-\xi_{-k})}{\omega - \xi(k) - \xi(-k) + i\delta}$$

 $\chi_{\text{pair}} \sim \log T \text{ as } T \to 0$ 

Superconducting pairing susceptibility

$$\chi_{\text{pair}} = \lim_{\omega \to 0} \frac{1}{\Omega} \sum_{k} \frac{f(\xi_k) - f(-\xi_{-k})}{\omega - \xi(k) - \xi(-k) + i\delta}$$

 $\chi_{\text{pair}} \sim \log T \text{ as } T \to 0$ 



#### 

# Left mover $\longleftarrow$ Right mover $LL \longleftrightarrow RR$ $q = 2k_F$

operative only at half-filling on a lattice

Left mover  $\longleftarrow$  Right mover  $LL \longleftrightarrow RR$   $q = 2k_F$ 

operative only at half-filling on a lattice

 $L\uparrow, R\downarrow\longleftrightarrow R\uparrow, L\uparrow$ 

for spin 1/2 fermions

At any filling so exists even in the continuum

Left mover  $\longleftarrow$  Right mover  $LL \longleftrightarrow RR$   $q = 2k_F$ 

operative only at half-filling on a lattice

 $L\uparrow, R\downarrow\longleftrightarrow R\uparrow, L\uparrow$ 

for spin 1/2 fermions

At any filling so exists even in the continuum

Sine-Gordon Hamiltonian

# $Umklapp \mathcal{H} = \frac{v_{\nu}}{2} \int dx \left[ K_{\nu} \Pi_{\nu}^{2} + \frac{1}{K_{\nu}} \left( \nabla \phi_{\nu} \right)^{2} + \frac{2g_{\nu}}{2\pi a^{2}} \cos \left( \alpha_{\nu} \phi_{\nu} \right) \right] \\ \nu = \text{charge, spin}$

$$\begin{aligned} & \mathsf{Umklapp} \\ \mathcal{H} = \frac{v_{\nu}}{2} \int dx \left[ K_{\nu} \Pi_{\nu}^{2} + \frac{1}{K_{\nu}} \left( \nabla \phi_{\nu} \right)^{2} + \frac{2g_{\nu}}{2\pi a^{2}} \cos \left( \alpha_{\nu} \phi_{\nu} \right) \right] \\ & \nu = \text{charge, spin} \end{aligned}$$

Renormalization group flows can be calculated for K and g g relevant  $\Rightarrow$  gapped, g irrelevant  $\Rightarrow$  gapless

$$\begin{aligned} & \mathsf{Umklapp} \\ \mathcal{H} = \frac{v_{\nu}}{2} \int dx \left[ K_{\nu} \Pi_{\nu}^{2} + \frac{1}{K_{\nu}} \left( \nabla \phi_{\nu} \right)^{2} + \frac{2g_{\nu}}{2\pi a^{2}} \cos \left( \alpha_{\nu} \phi_{\nu} \right) \right] \\ & \nu = \text{charge, spin} \end{aligned}$$

Renormalization group flows can be calculated for K and g g relevant  $\Rightarrow$  gapped, g irrelevant  $\Rightarrow$  gapless



$$\begin{aligned} & \mathsf{Umklapp} \\ \mathcal{H} = \frac{v_{\nu}}{2} \int dx \left[ K_{\nu} \Pi_{\nu}^{2} + \frac{1}{K_{\nu}} \left( \nabla \phi_{\nu} \right)^{2} + \frac{2g_{\nu}}{2\pi a^{2}} \cos \left( \alpha_{\nu} \phi_{\nu} \right) \right] \\ & \nu = \text{charge, spin} \end{aligned}$$

Renormalization group flows can be calculated for K and g g relevant  $\Rightarrow$  gapped, g irrelevant  $\Rightarrow$  gapless



Ref: Quantum physics in one dimension, Giamarchi

$$\mathcal{H} = \frac{v_{\nu}}{2} \int dx \left[ K_{\nu} \Pi_{\nu}(x)^2 + \frac{1}{K_{\nu}} \left( \nabla \phi_{\nu}(x) \right)^2 + \frac{2g_{\nu}}{2\pi a^2} \cos(\alpha_{\nu} \phi_{\nu}) \right]$$

$$K_{\nu}^{c} = \frac{8\pi}{\alpha_{\nu}^{2}}$$
$$h_{\nu} = 2\left(\frac{K_{\nu}}{K_{\nu}^{c}} - 1\right)$$
$$g_{\nu}^{\perp} = K_{\nu}^{c}g_{\nu}$$

$$\frac{dg_{\nu}^{\perp}}{dl} = -h_{\nu}g_{\nu}^{\perp},$$

$$\frac{dh_{\nu}}{dl} = -\left(g_{\nu}^{\perp}\right)^{2}.$$

# gv(u) is independent of the boost RG flow equations are unaffected by the boost

# gv(u) is independent of the boost RG flow equations are unaffected by the boost

The boost only changes bare values of K and g

gv(u) is independent of the boost RG flow equations are unaffected by the boost

The boost only changes bare values of K and g

$$h_{\nu,i}(u) = 2\left(\frac{K_{\nu,i}(0)}{K_{\nu}^{c}} - 1\right) - \frac{v_{i}^{F}(0)f(u)}{v_{\nu}(0)K_{\nu}^{c}}\left[\left(K_{\nu,i}(0)\right)^{2} - 1\right]$$
$$g_{\nu,i}^{\perp}(u) = g_{\nu,i}^{\perp}(0)\left[1 - \frac{v_{i}^{F}(0)f(u)}{2v_{\nu}(0)}\left(K_{\nu,i}(0) + \frac{1}{K_{\nu,i}(0)}\right)\right].$$

$$f(u) = -\frac{v_i^F(0) - v_i^F(u)}{v_i^F(0)}$$

#### Spinless fermions



#### Spinless fermions

Boost can open a gap

#### Spinless fermions








The boost can open a gap for K < 1 Possibly algebraic CDW to long-ranged CDW



Boost can close gap









Interesting possibility

Unboosted system with charge and spin gap

Boost closes one of the gaps

Close charge gap but not spin gap - Luther-Emery fluid

#### Conclusions

- A boost has a non-trivial effect on algebraic order in 1D only for non-parabolic bands
- For a simple lattice dispersion, a boost can strengthen superconductivity (and CDW order)
- At commensurate filling, a boost can open or close gaps depending on whether the fermions are spinless or spinful
- Possible to obtain a Luther-Emery fluid by boosting starting from fully gapped state