

Confronting Standard Model Limitations Through Non-SUSY SO(10)

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- 1 **SUSY SO(10) unifies couplings, gives m_ν by type-I seesaw, predicts Higgsino, Neutralino, wino as DM; predicts R-Parity as Gauged Discrete Sym. for the DM stability, Baryogenesis via Leptogenesis, and explains the origin of P-violation in weak interaction. It has natural solutions to gauge hierarchy problem.**
- 2 **Supersymmetry appears to be indispensable for fundamental understanding of particle physics and cosmology.**
- 3 **BUT “We should also think outside the SUPERBOX”:
John Ellis (CANDARK2017).**

- 1 IN THE ABSENCE OF ANY EXPERIMENTAL CONFIRMATION OF SUSY SO FAR AND WITH FINE TUNING APPROACH TO GAUGE HIERARCHY WE (MKP, B.P. Nayak, R.SATPATHY, R. L. AWASTHI, JHEP 04 (2017)075) ATTEMPT TO SEE IF SIMILAR SOLUTIONS EXIST IN NON-SUSY SO(10) FOR m_ν , DM, BAU, and UNIFICATION WITHIN SM PARADIGM.
- 2 FOR FINE TUNING:
S. Weinberg, “Living in the Universe, in Universe or Multiverse”, B.J. Carr ed.(2007) .

A MINIMAL MODEL

- ① **CONSIDER $SO(10) \rightarrow SU(5) \times U(1)_X \rightarrow SM$:
AFTER SECOND STEP BR.:**

$$X = 4T_{3R} + 3(B - L)$$

POSSIBLE WITH (i) $45_H \oplus 16_H$ OR (ii) $45_H \oplus 126_H$

- ② **IN BOTH CASES: $Z_X = (-1)^{4T_{3R} + 3(B-L)} = (-1)^{3(B-L)} \equiv Z_{MP}$, CALLED MATTER PARITY.**
- ③ **$Z_{MP} = -1$ in (i) SINCE $|B - L| = \text{ODD}$; i.e MATTER PARITY HAS BROKEN SPONTANEOUSLY.**
- ④ **CHOICE (ii) RH. TRIPLET WITH $T_{3R} = 1, B - L = -2$ GETS VEV. $\nu_R \sim M_U$: GIVES Type-I \oplus Type-II for m_ν . HERE $Z_{MP} = +1$, SM LAGRANGIAN CONSERVES MP AS GAUGED DISCRETE SYM.**

- ① SM FERMIONS HAVE -1 AND SM HIGGS HAS $+1$ MATTER PARITY.
- ② MATTER PARITY OF SO(10) REPRESENTATIONS:

$$\begin{aligned}
 Z_{MP} &= \text{Even : } 10, 45, 54, 120, 126, 210; \dots \\
 Z_{MP} &= \text{Odd : } 16, 144, \dots \dots \dots
 \end{aligned}
 \tag{1}$$

- ③ SM FERMIONS $\supset 16_F$ HAVE $Z_{MP} = -1$; SM HIGGS $\supset 10_H$ HAS $Z_{MP} = +1$; THE YUK. INTS. $16_F.16_F.10_H$ AND $16_F.16_F.126_H^\dagger$ CONSERVE MP.
- ④ DM FERMIONS MUST BE IN NON-STANDARD MP-EVEN REPS.
- ⑤ DM SCALARS MUST BE IN NON-STANDARD MP-ODD REPS.

- 1 Ref for DM in SO(10): Kadastik et.al,PRD(2009);Frigerio & Hambye,PRD(2010); MKP et. al, PRD(2010); MKP PLB(2011); Mambrini et.al PRD(2015).

1

$$-\mathcal{L}_{\text{Yuk}} = \mathbf{Y} \mathbf{16}_F \cdot \mathbf{16}_F \cdot \mathbf{10}_H + \mathbf{f} \mathbf{16}_F \cdot \mathbf{16}_F \cdot \mathbf{126}_H^\dagger + \dots \quad (2)$$

2

$$\mathbf{m}_\nu = \mathbf{f} \mathbf{v}_L - \mathbf{M}_D \frac{1}{\mathbf{f} \mathbf{v}_R} \mathbf{M}_D^\mathbf{T}, \quad (3)$$

$$\mathbf{M}_D = M_u \text{ (UP QUARK MASS MATRIX),}$$

$$\mathbf{v}_L = \lambda \mathbf{v}_R \mathbf{v}_{\text{ew}}^2 / M_\Delta^2$$

$$M_N = \mathbf{f} \mathbf{v}_R \sim f M_{\text{GUT}},$$

(4)

- 3 REF: E. Akhmedov, M. Frigerio, JHEP 01(2007)043:
Eq. quadratic in f , $2^3 = 8$ solutions; We find 2 distinct solutions for each chosen basis

1

$$\mathbf{M}_D^{(d)}(\text{GeV}) =$$

$$(5) \quad \begin{pmatrix} 0.01832 + 0.00441i & 0.08458 + 0.01114i & 0.6588 + 0.27319i \\ 0.08458 + 0.01114i & 0.38538 + 1.56 \times 10^{-5}i & 3.3278 + 0.00019i \\ 0.65882 + 0.27319i & 3.32785 + 0.00019i & 81.8543 - 1.64 \times 10^{-5} \end{pmatrix}$$

1

$$\mathbf{M}_D^{(u)}(\text{GeV}) =$$

$$(6) \begin{pmatrix} 0.00054 & (1.5027 + 0.0038i)10^{-9} & (7.51 + 3.19i)10^{-6} \\ (1.5027 + 0.0038i)10^{-9} & 0.26302 & 9.63 \times 10^{-5} \\ (7.51 + 3.19i)10^{-6} & 9.63 \times 10^{-5} & 81.9963 \end{pmatrix}$$

1

$$\mathbf{f} = 10^{-6} \times$$

$$\begin{pmatrix} 0.385 + 0.1291i & 0.4617 - 0.4922i & 3.509 + 1.080i \\ 0.4617 - 0.4922i & 4.626 + 0.1567i & 22.80 + 0.3317i \\ 3.509 + 1.080i & 22.80 + 0.3317i & 511.6 + 0.47i \end{pmatrix}. \quad (7)$$

①

$$\mathbf{f} = 10^{-6} \times$$

$$\begin{pmatrix} 0.3175 + 0.0904i & 0.1232 - 0.6089i & -0.4869 - 0.6918i \\ 0.1232 - 0.6089i & 3.610 - 0.0724i & 1.587 + 0.2599i \\ -0.4869 - 0.6918i & 1.587 + 0.2599i & 511.8 + 0.6524i \end{pmatrix} \quad (8)$$

1

$$\mathbf{f} = 10^{-6} \times$$

$$(9) \quad \begin{pmatrix} -0.0690 + 0.0147i & -0.341 + 0.0164i & -4.0194 + 1.5783i \\ -0.341 + 0.0164i & -1.5745 - 0.2133i & -20.2464 - 0.3306i \\ -4.0194 + 1.5783i & -20.2464 - 0.3306i & -507.895 - 0.4034i \end{pmatrix}$$

①

$$\mathbf{f} = 10^{-6} \times$$

$$\begin{pmatrix} -0.000025 + 0.000008i & -0.00019 - 0.00215i & -0.00538 - 0.00177i \\ -0.00019 - 0.00215i & -0.56091 + 0.0092i & 0.95702 - 0.27084i \\ -0.00538 - 0.00177i & 0.95702 - 0.27084i & -508.16 - 0.60957i \end{pmatrix}$$

(10)

Fig1: Hierarchical Spectrum of $RH\nu$

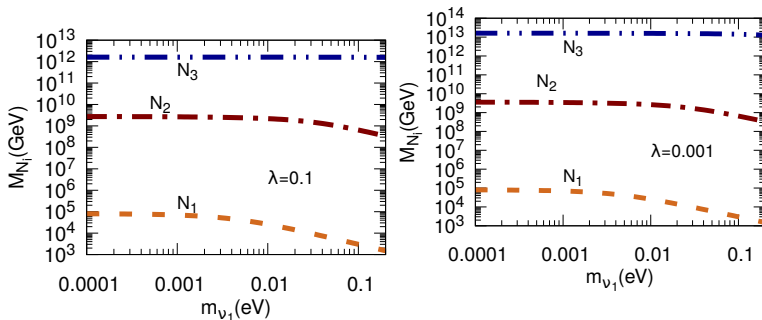


Figure: Prediction of Hierarchical Spectrum of heavy RH neutrino masses as a function of the lightest neutrino mass and when the three neutrino masses are normally ordered. The values of $M_{\Delta_L} = 10^{12}$ GeV and $\nu_R = 10^{15.5}$ GeV have been kept fixed.

Fig2: Compact Spectrum of $RH\nu$

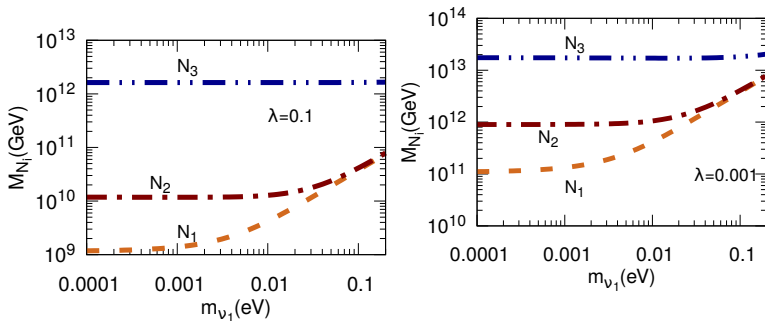


Figure: Prediction of Compact Spectrum of heavy RH neutrino masses as a function of the lightest neutrino mass when the three neutrino masses are normally ordered. The values of $M_{\Delta_L} = 10^{12}$ GeV and $v_R = 10^{15.5}$ GeV have been kept fixed.

1

$$\varepsilon_{i\alpha} = \frac{\Gamma(\mathbf{N}_i \rightarrow \mathbf{l}_\alpha + \mathbf{H}^*) - \Gamma(\mathbf{N}_i \rightarrow \bar{\mathbf{l}}_\alpha + \mathbf{H})}{\sum_\beta [\Gamma(\mathbf{N}_i \rightarrow \mathbf{l}_\beta + \mathbf{H}^*) + \Gamma(\mathbf{N}_i \rightarrow \bar{\mathbf{l}}_\beta + \mathbf{H})]}. \quad (11)$$

One loop decay contributions of N_i are mediated by either $N_{k \neq i}$ or Δ_L [?], as shown in Fig. 3. The total asymmetry is sum of the two contributions

$$\varepsilon_{i\alpha} = \varepsilon_{i\alpha}^{\mathbf{N}} + \varepsilon_{i\alpha}^{\mathbf{\Delta}}. \quad (12)$$

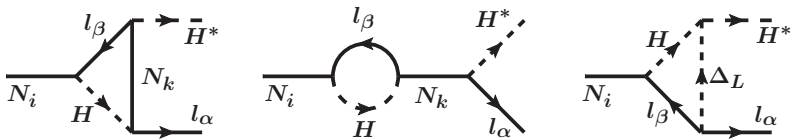


Figure: One-loop Feynman diagrams for the decay of RH neutrino N_i . The first and the third diagrams represent vertex corrections and the second diagram represents self-energy correction.

Compact Spectrum: u -diagonal basis:BAU:fig6

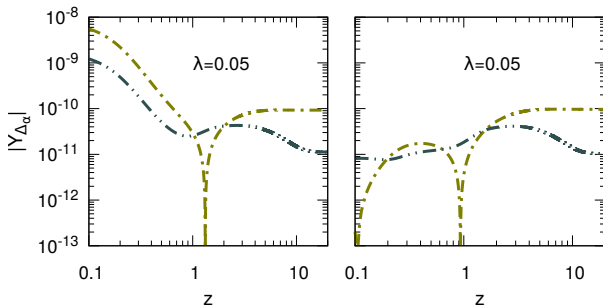


Figure: The baryon asymmetry in $e + \mu$ flavours (double-dot-dashed blue curve) and τ flavor (dot-dashed curve) for the u -quark diagonal basis and compact spectrum $RH\nu$ mass scenario. Left (right) panel correspond to non-zero (zero) initial thermal abundance.

Hierarchical Spectrum; u -Diagonal Basis: BAU: fig9

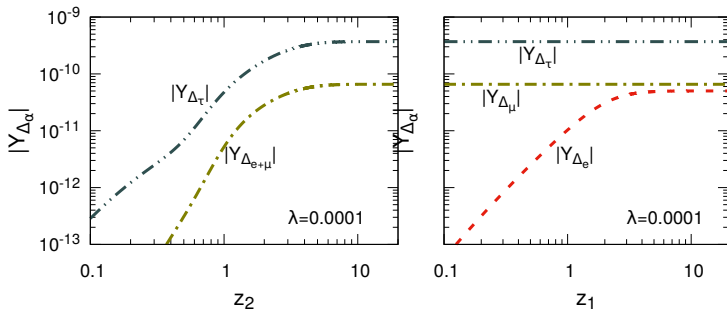


Figure: The baryon asymmetries in $e + \mu$ and τ flavors (left panel) and separately for e, μ and τ flavors (right panel). Hierarchical spectrum for $RH\nu$ has been used.

- ① USING SOLUTIONS FROM BOLTZMANN Eq. WE GET BAU IN AGREEMENT WITH PLANCK2015 DATA.
- ② THIS OCCURS IN THE CASE OF DIRAC NEUTRINO MASSES DERIVED USING UP-QUARK DIAGONAL BASIS
- ③ THE AGREEMENT OCCURS WITH BOTH COMPACT AND HIERRAECHICAL SPECTRA OF RH_ν MASSES.
- ④ WE HAVE USED NORMALLY ORDERED LIGHT m_ν .

- 1 **PROPOSED IN SM**
EXTN:Cirelli,Forengo,Strumia,NPB(2006):
 $\Sigma(3, 0, 1) \subset 45_F$ OF **SO(10)**
MASS DIFFERENCE $m_{\Sigma^\pm} - m_{\Sigma^0} = 166$ MeV
- 2 **Extensive investigation on indirect signals and proposed collider searches: S.Mohanty et. al,IJMP(2012); A. Hryczuk et.al JCAP(2014), M. Cirelli et. al,JHEP(2014), T.Aizawa et.al,PRD(2015) and a large list of authors.**
- 3 **SUGGESTED MASS RANGE 400 – 2700 GeV.**
Substantially affects $SU(2)_L$ gauge coupling evolution.
- 4 **DIRECT SEARCH: LUX, XENON100 > 35 GeV; LHC: CANDARK2017 LECTURE:(1-1000) GeV ?**

- 1 WITH $A_F = 45_F, \Phi = 210_H, E = 54_H$, SO(10) GIVES
- 2 Non-Std. Higgs Maj.Fermion Yuk. Int:

$$-\mathcal{L}_{\text{Yuk}} = \mathbf{A}_F (\mathbf{m}_A + \mathbf{h}_p \Phi + \mathbf{h}_e \mathbf{E}) \mathbf{A}_F, \quad (13)$$

where $m_A \simeq M_U$ and h_i 's are Yukawa couplings.

- 3 THE TRIPLET MAJORANA FERMION DM MASS CAN HAVE ANY VALUE FROM $100 - 10^{15}$ GeV BY TUNING THE PARAMETERS:

$$m_{\Sigma}(3, 0, 1) = m_A + \sqrt{2}h_p \frac{\Phi_1}{3} + \sqrt{\frac{3}{5}}h_e \langle E \rangle. \quad (14)$$

- 1 We use the fields Δ_L , Σ , needed for m_ν , Leptogenesis, and DM. In addition we need a color octet fermion $C_8(1, 0, 8) \subset 45_F$. Its mass is derived in the same way. Aizawa et al.PRD(2015: this can act as a nonthermal source of Σ .

2

$$\begin{aligned} \frac{1}{\alpha_i(M_Z)} &= \frac{1}{\alpha_i(M_U)} + \frac{a_i}{2\pi} \ln \left(\frac{M_\Sigma}{M_Z} \right) \\ &+ \frac{a'_i}{2\pi} \ln \left(\frac{M_{C_8}}{M_\Sigma} \right) + \frac{a''_i}{2\pi} \ln \left(\frac{M_\Delta}{M_{C_8}} \right) \\ &+ \frac{a'''_i}{2\pi} \ln \left(\frac{M_U}{M_\Delta} \right) + \Theta'_i + \Theta''_i + \Theta'''_i - \frac{\lambda_i}{12\pi}, \quad (15) \end{aligned}$$

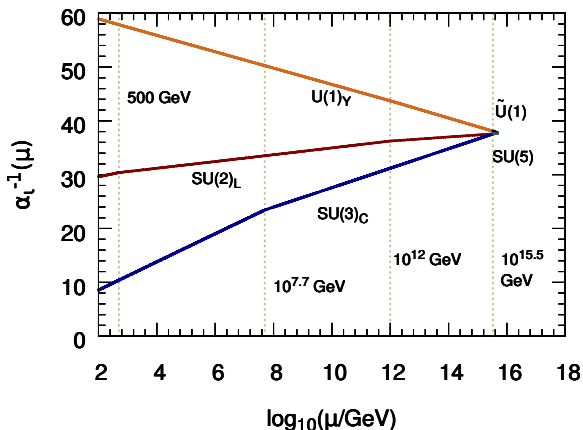


Figure: Unification of gauge couplings of SM with $M_U = 10^{15.56}$ GeV, $g_G = 0.573$.

THRESHOLD EFFECTS

1 Refs:

(1).S.Weinberg PLB(1980), L. Hall, NPB(1981),
B.Ovrut, H.Snitzer, PLB(1981), MKP,C.Hazra,
PRD(1989)

(2).SUPERHEAVY COMPONENTS OF A GIVEN
SO(10) REP. HAVE DEGENERATE MASS: R. N.
Mohapatra, MKP PRD(1993).

2

$$\begin{aligned}\lambda_1 &= 17/5 + 4\eta_{(10)} + (0)\eta_{(45)} + 136\eta_{(126)}, \\ \lambda_2 &= 6 + 4\eta_{(10)} + 2\eta_{(45)} + 140\eta_{(126)}, \\ \lambda_3 &= 8 + 4\eta_{(10)} + 3\eta_{(45)} + 140\eta_{126}, \\ \eta_X &= \ln(M_X/M_U)\end{aligned}\tag{16}$$

PROTON LIFETIME

1

$$\tau_p^{\text{expt.}} \geq 1.4 \times 10^{34} \text{ yrs.} \quad (17)$$

2

$$\tau_p^{\text{SO}(10)} \simeq 1.8 \times 10^{34 \pm 3.712\eta_S \pm 1.012\eta_F} \text{ yrs.} \quad (18)$$

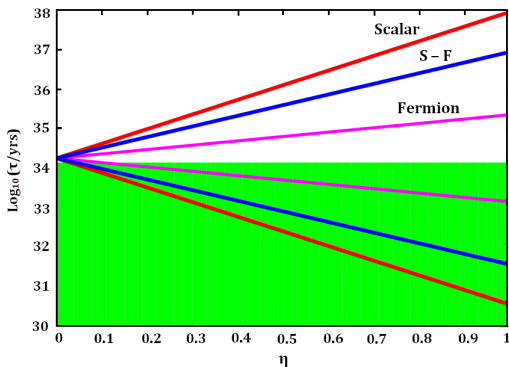


Figure: Proton lifetime prediction for the decay mode $p \rightarrow e^+\pi^0$ shown by slanting solid lines as a function of $\eta = \eta_S(\eta_F) = |\log_{10}(M_{SH}/M_U)|(|\log_{10}(M_F/M_U)|)$ for super-heavy scalar(fermion) components.

- 1 Non-SUSY SO(10) directly breaking to SM can successfully answer the questions of m_ν , BAU, DM, DM stability, Unification, and observed proton stability. It gives asymptotic parity restoration.
- 2 Embedding DM in SO(10) enhances proton lifetime prediction substantially.
- 3 Resolution of gauge hierarchy by fine-tuning is not as natural as in SUSY SO(10).
- 4 In non-SUSY case there is no gravitino problem, or problem associated with Higgsino mediated proton decay.
- 5 Doublet-triplet splitting occurs naturally under extended survival hypothesis.

Thank You