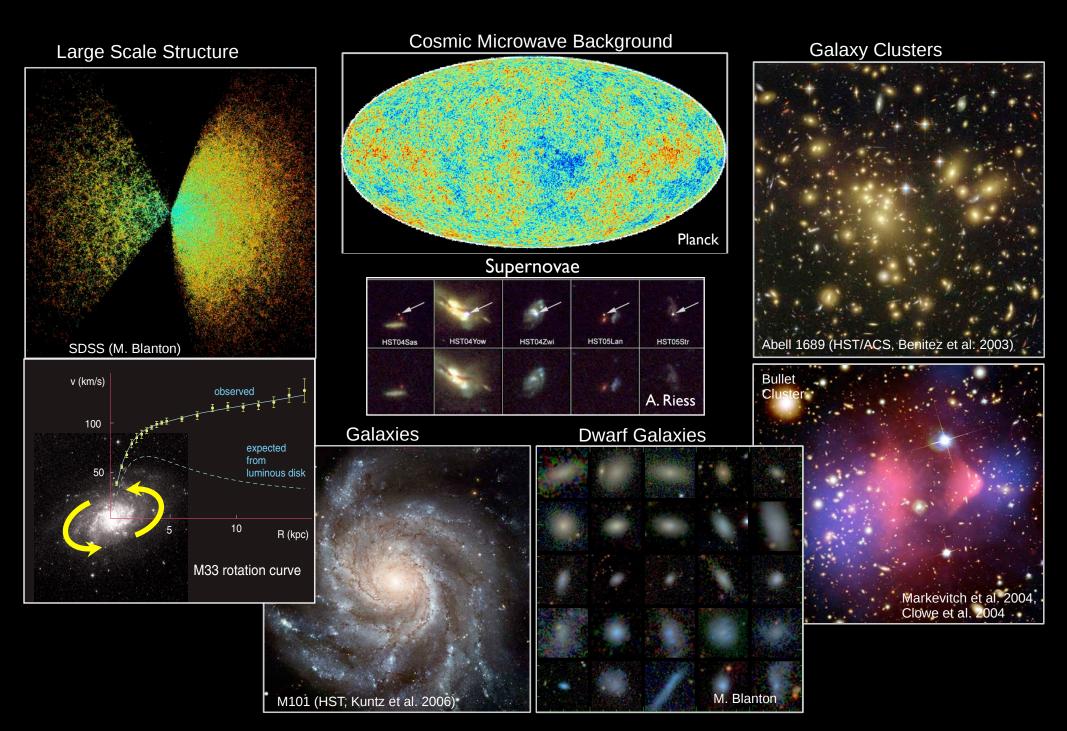
Overview of dark matter

Paolo Gondolo University of Utah



Dwarf galaxies

Dwarf galaxies are dominated by dark matter.

40 20

-20

-40

-60

40

20

-20

-40

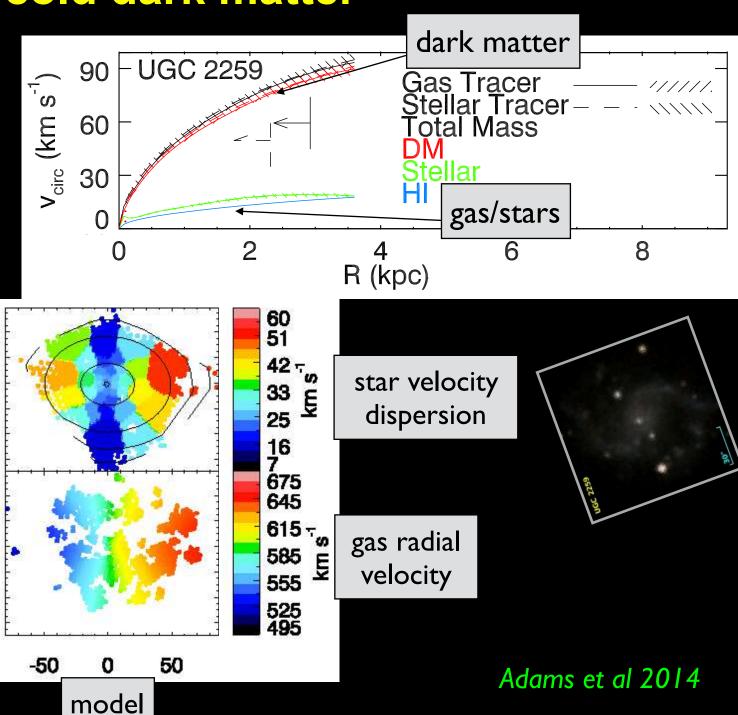
-60

-50

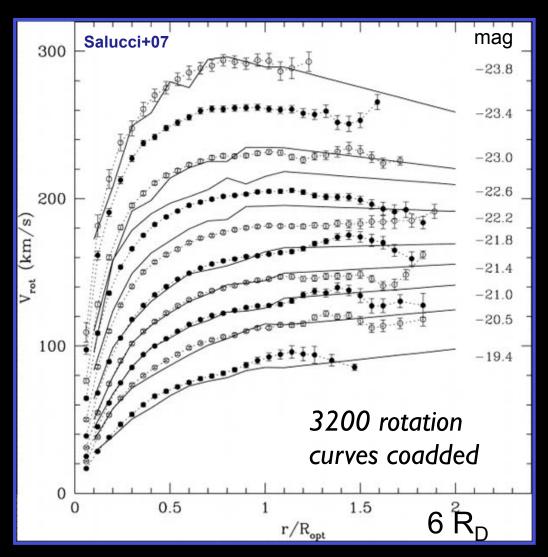
0

data

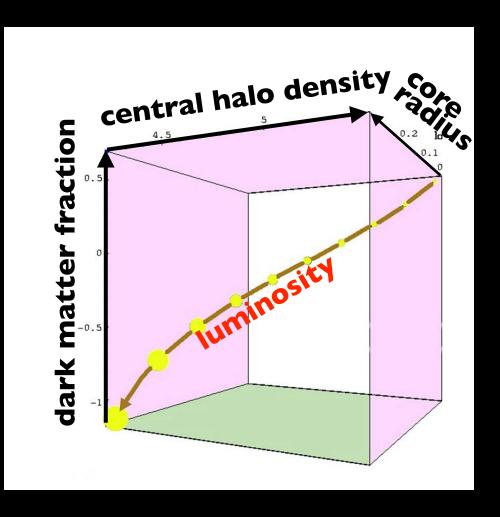
50



Spiral galaxies



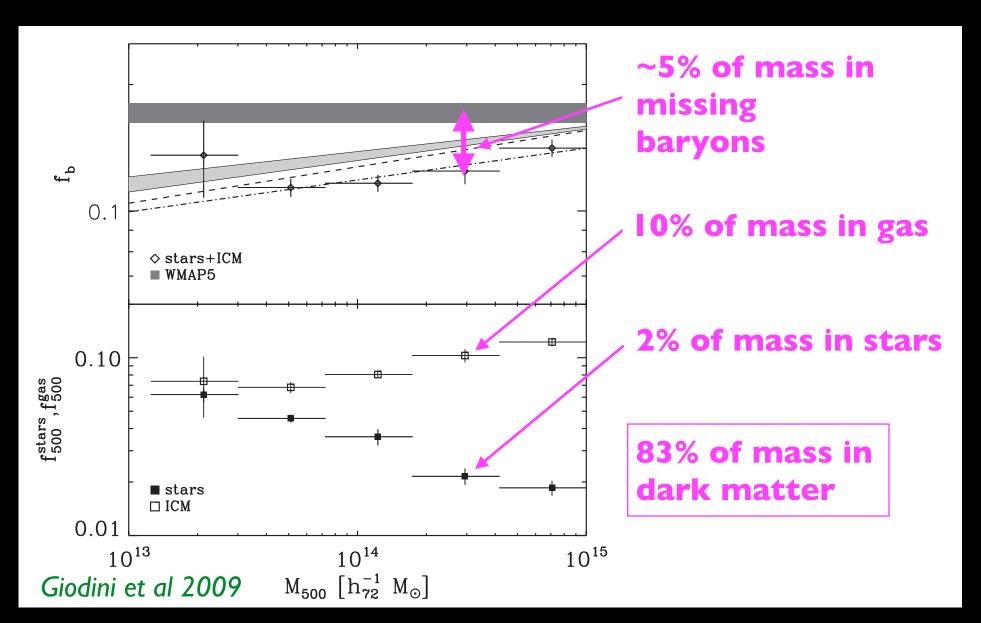
Empirical correlations found from thousands of spiral galaxy rotation curves



Salucci et al 2007

Galaxy clusters

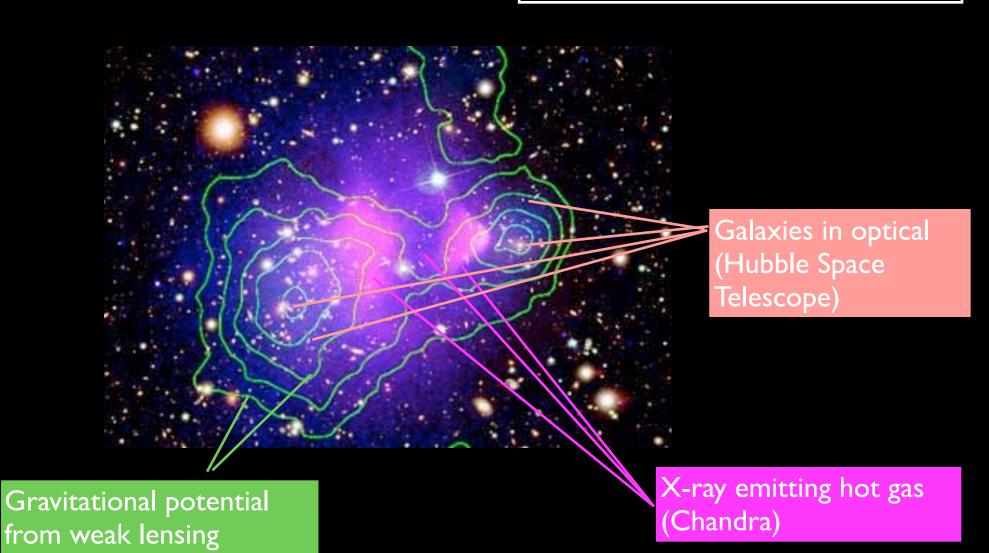
Galaxy clusters are mostly dark matter with some gas and a sprinkle of galaxies



Cold dark matter, not modified gravity

The Bullet Cluster

Symmetry argument: gas is at center, but potential has two wells.

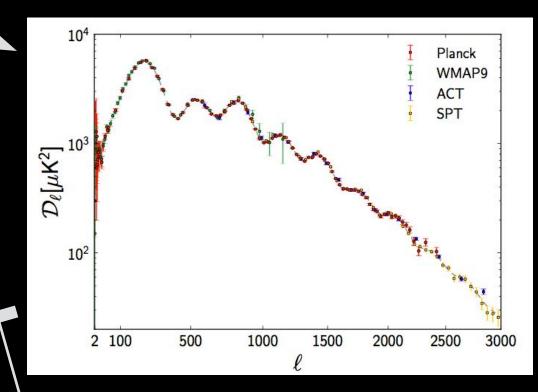


Planck (2013)

Evidence for cold dark matter

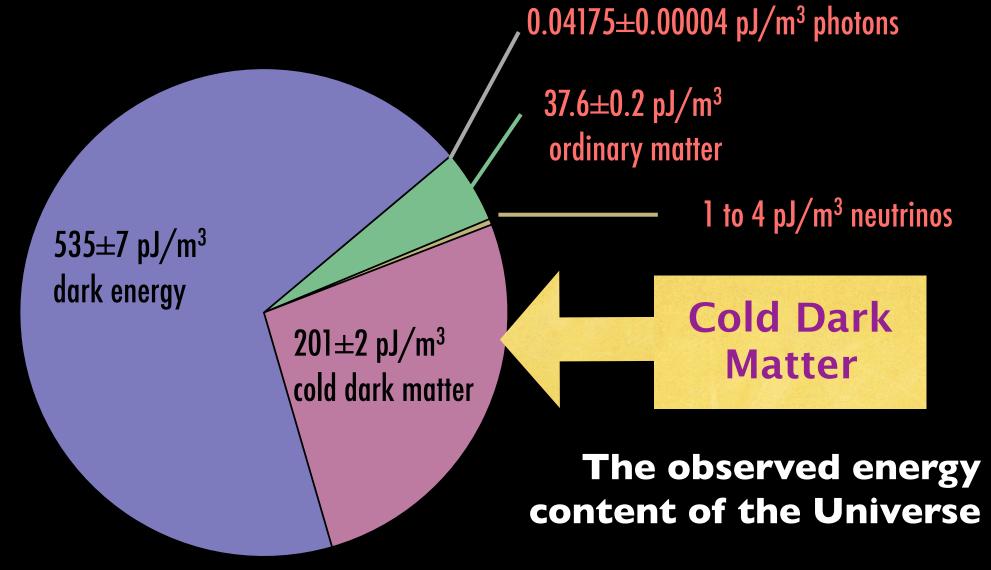
Cosmic Microwave Background fluctuations

	Planck+WP+highL+BAO	
Parameter	Best fit	68% limits
$\Omega_{ m b} h^2 \ldots \ldots$	0.022161	0.02214 ± 0.00024
$\Omega_{ m c} h^2$	0.11889	0.1187 ± 0.0017
$100\theta_{\mathrm{MC}}$	1.04148	1.04147 ± 0.00056
au	0.0952	0.092 ± 0.013
$n_{\rm S}$	0.9611	0.9608 ± 0.0054
$ln(10^{10}A_s)$	3.0973	3.091 ± 0.025
$\overline{\Omega_{\Lambda} \ldots \ldots \ldots }$	0.6914	0.692 ± 0.010
σ_8	0.8288	0.826 ± 0.012
z_{re}	11.52	11.3 ± 1.1
H_0	67.77	67.80 ± 0.77
Age/Gyr	13.7965	13.798 ± 0.037
$100\theta_*$	1.04163	1.04162 ± 0.00056
$r_{\rm drag}$	147.611	147.68 ± 0.45



linear perturbation theory

general relativity and statistical mechanics at 10^4 K ~ 1 eV/k



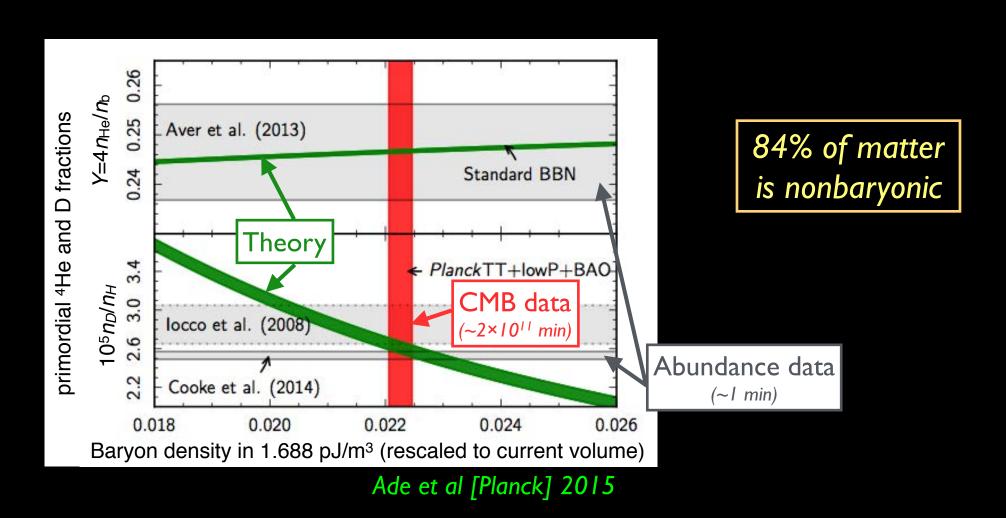
matter $p \ll \rho$ radiation $p=\rho/3$ vacuum $p=-\rho$

Planck (2015)
TT,TE,EE+lowP+lensing+ext

 $1 \text{ pJ} = 10^{-12} \text{ J}$ $\rho_{\text{crit}} = 1688.29 \ h^2 \text{ pJ/m}^3$

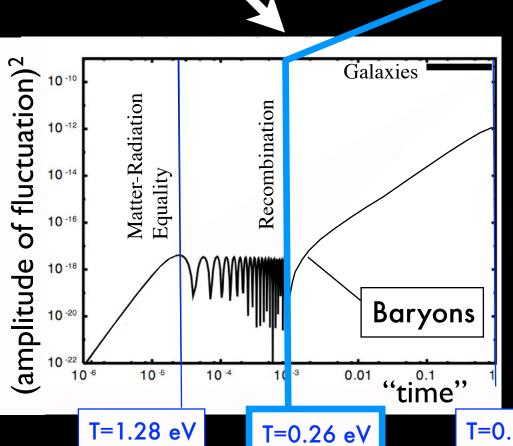
BIG BANG NUCLEOSYNTHESIS

The baryon-to-photon ratio has been the same since ~ 1 minute after the Big Bang. Baryons are $\lesssim 15.7\%$ of the mass in matter.

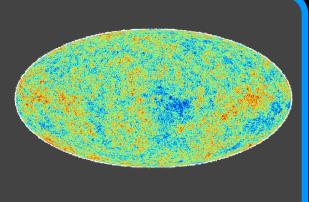


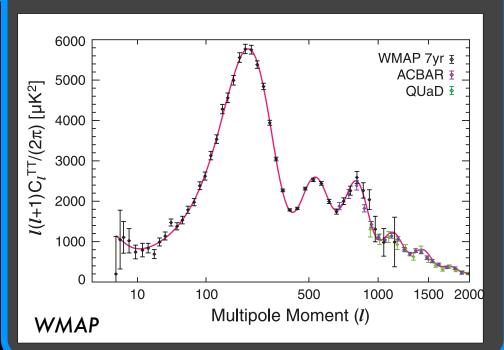


13 billion years ago



Cosmic Microwave Background fluctuations

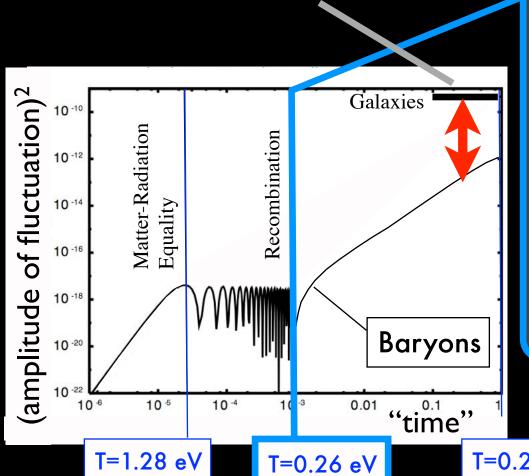




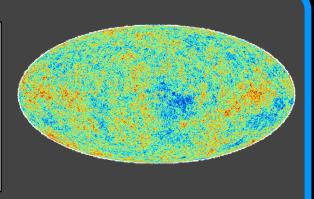
T=1.28 eV

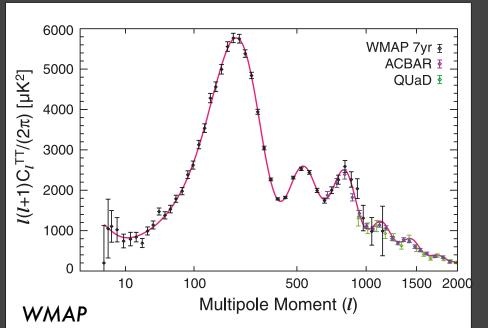
T=0.2348 meV

Without dark matter, fluctuations are too small to gravitationally grow into galaxies in the given 13 billion years.



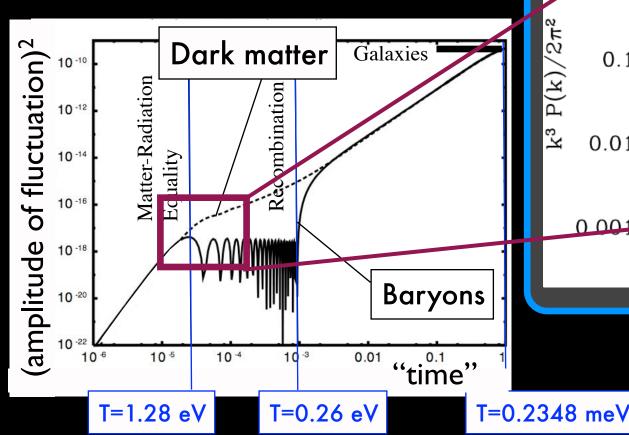






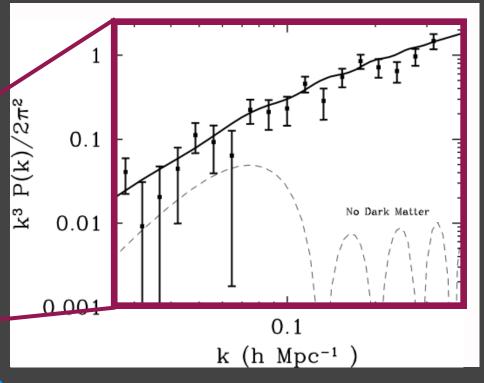
T=0.2348 meV

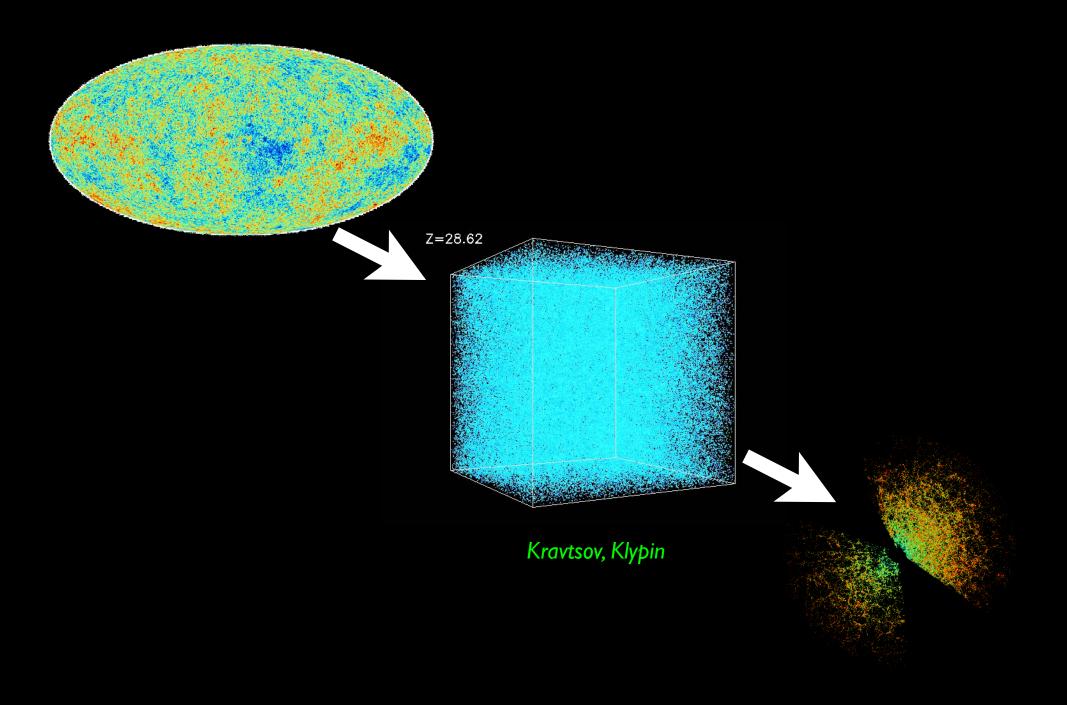
Dark matter fluctuations, uncoupled to the plasma, start growing early and have enough time to grow into galaxies

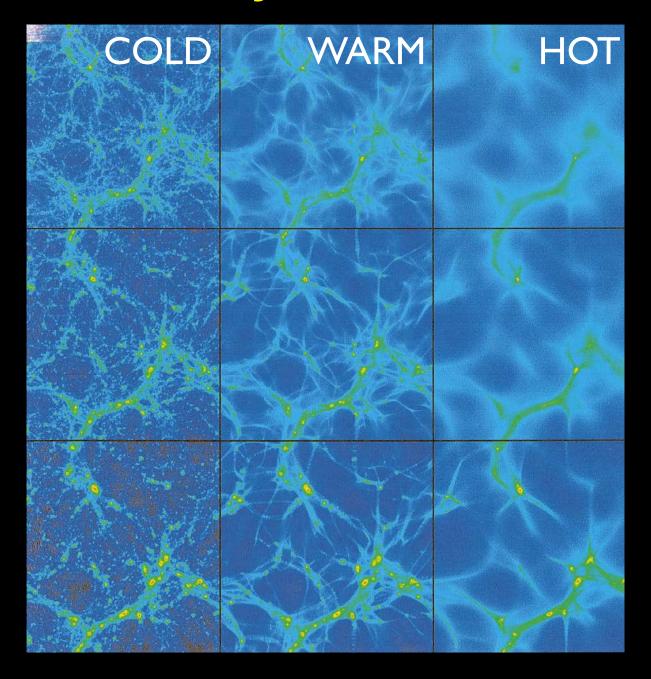


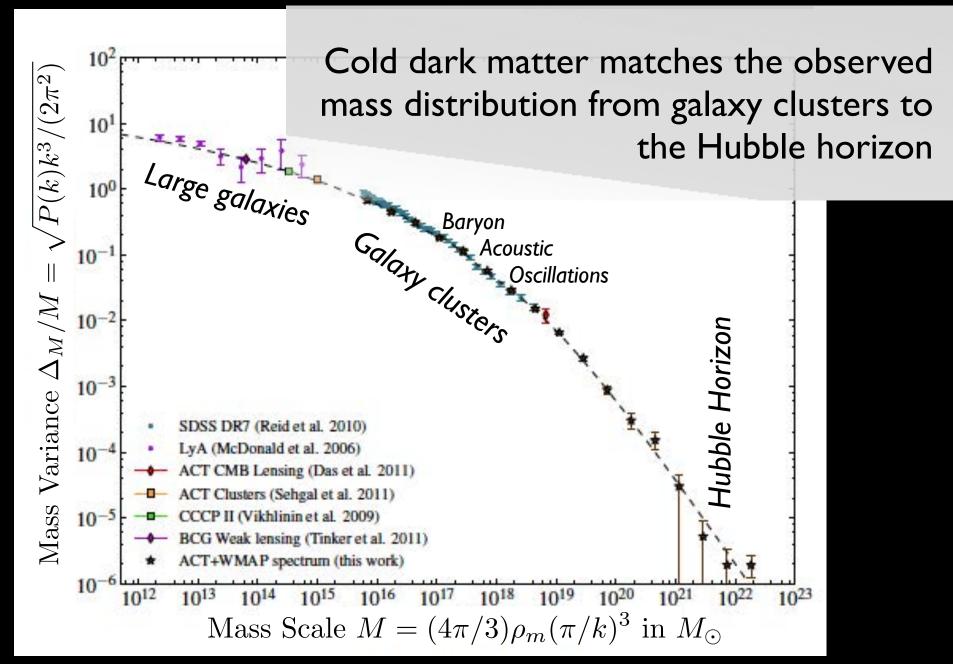
More than 80% of all matter does not couple to the primordial plasma!

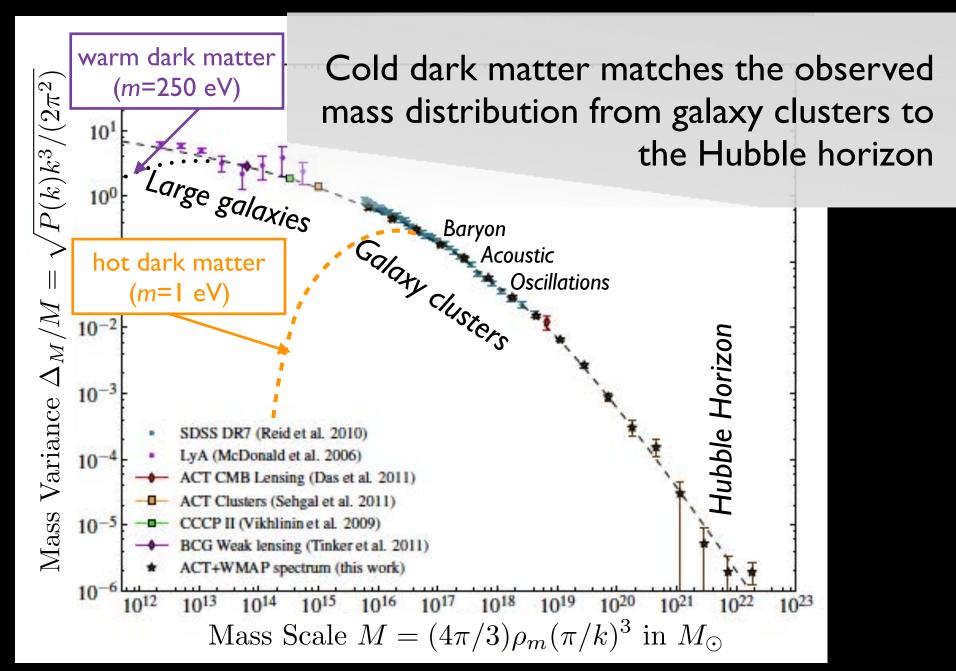
SDSS





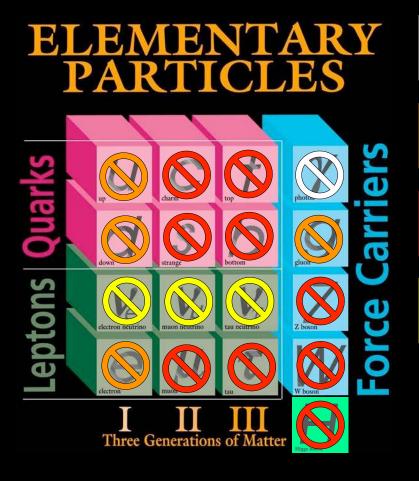








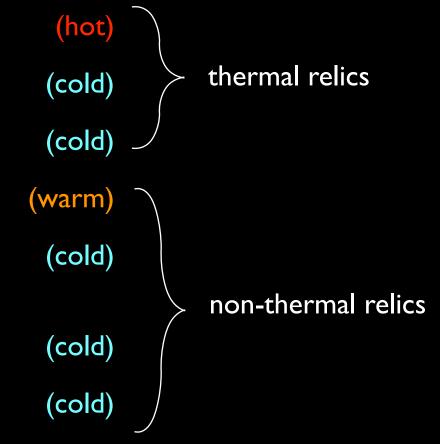
Is dark matter an elementary particle?



- is the particle of light
- Occuples to the plasma
- disappears too quickly
- is hot dark matter

No known particle can be nonbaryonic cold dark matter!

- SM neutrinos
- lightest supersymmetric particle
- lightest Kaluza-Klein particle
- sterile neutrinos, gravitinos
- Bose-Einstein condensates, axions, ultralight scalars, axion clusters
- solitons (Q-balls, B-balls, ...)
- supermassive wimpzillas



Mass range

 $10^{-22} \, {\rm eV} \, (10^{-59} {\rm kg}) \, {\rm B.E.C.s}$ $10^{-8} \, M_{\odot} \, (10^{+22} {\rm kg}) \, {\rm axion \, clusters}$

Interaction strength range

Only gravitational: wimpzillas Strongly interacting: B-balls

Hot dark matter

- relativistic at kinetic decoupling (last scattering, start of free streaming)
- big structures form first, then fragment

light neutrinos

Cold dark matter

- non-relativistic at kinetic decoupling
- small structures form first, then merge

neutralinos, axions, WIMPZILLAs, solitons

Warm dark matter

- semi-relativistic at kinetic decoupling
- smallest structures are erased

sterile neutrinos, gravitinos

Thermal relics

- in thermal equilibrium with the plasma in the early universe
- produced in collision of plasma particles
- insensitive to initial conditions

```
neutralinos, other WIMPs, ....
```

Non-thermal relics

- not in thermal equilibrium with the plasma in the early universe
- produced in decays of heavier particles or extended structures
- have a memory of initial conditions

```
axions, WIMPZILLAs, solitons, ....
```

DM production	in plasma reactionsfrom decays of decoupled speciesemitted from extended objects	collider searches cosmic density
$\begin{array}{c} DM\text{-}\overline{DM} \text{ annihilation} \\ \chi^+\overline{\chi} \longrightarrow \text{anything} \end{array}$	self-conjugate DMasymmetric DM	indirect detection cosmic density
DM—SM scattering $\chi + SM \rightarrow \chi' + SM$	elastic/inelastic scatteringshort-/long-range interactions	hot/cold/warm halo (sub)structure direct detection
DM—DM scattering $\chi^+\chi \rightarrow \chi^+\chi$	- collisionless - self-interacting	dark halo structure
DM decay χ —anything	- stable - long-lived - ensemble of short-lived particles	indirect detection

Some factors affecting the particle dark matter cosmic density

— Production mechanism:

- produced in reactions of plasma (thermal) particles
 - reaching reaction equilibrium
 - not reaching reaction equilibrium
 - coannihilating with similar mass particles
- produced in decays of non-thermal particles
- emitted from extended objects

axions, ...

WIMP freeze-out, ...

FIMP freeze-in, ...

neutralinos, .

gravitinos, ...

— Dark matter-antimatter asymmetry:

• self-conjugate

Majorana fermions, neutralinos, axions, gravitinos, ...

not self-conjugate

Dirac fermions, asymmetric dark matter, ...

— Hubble expansion rate before nucleosynthesis:

• standard vs nonstandard cosmology

low temperature reheating, kination, ...

QCD axions

QCD axions as dark matter

Hot

Produced thermally in early universe Important for $m_a > 0.1 eV$ ($f_a < 10^8$), mostly excluded by astrophysics

Cold

Produced by coherent field oscillations around mimimum of $V(\theta)$ (Vacuum realignment)

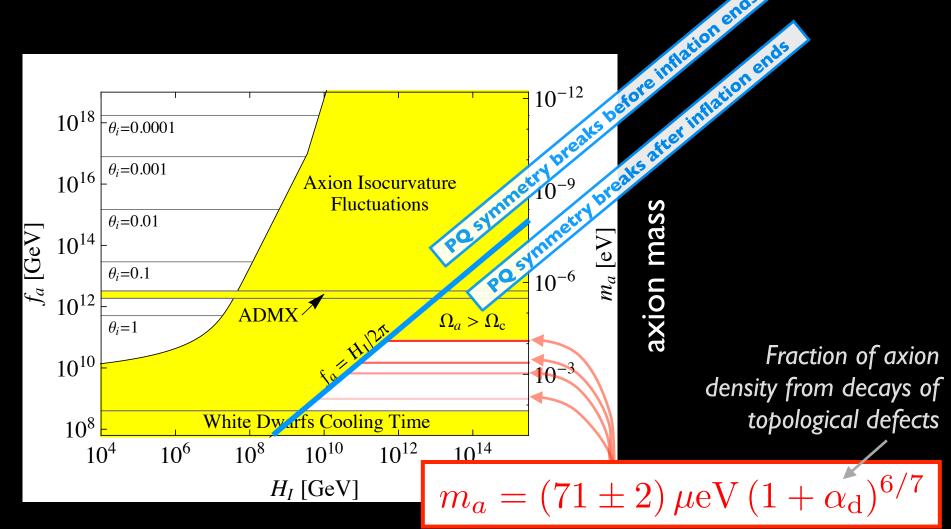
Produced by decay of topological defects

(Axionic string decays)

Still a very complicated and calculation!

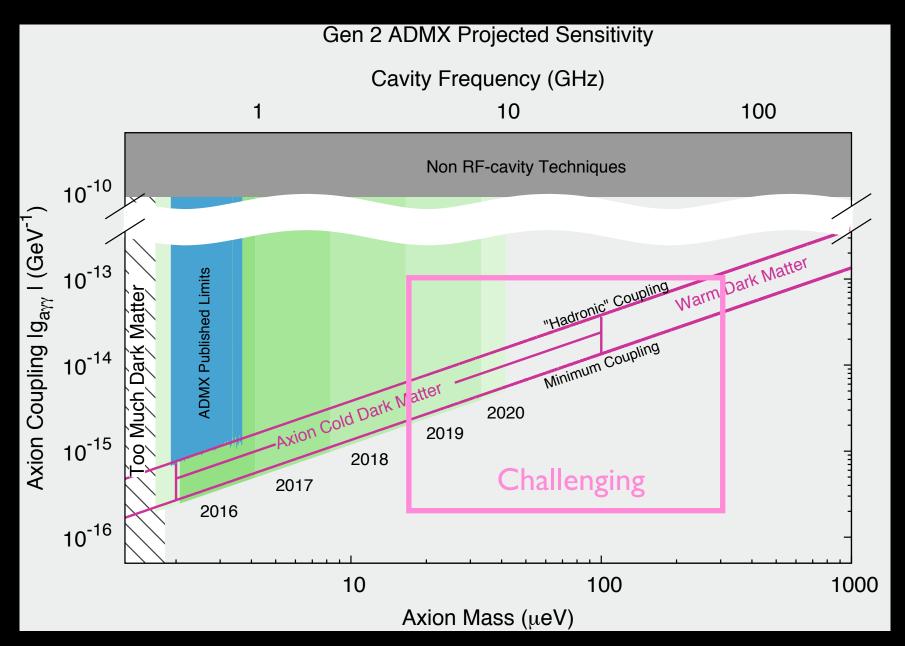
uncertain calculation!

e.g. Hiramatsu et al 2012



Expansion rate at end of inflation

QCD axions as cold dark matter: searches



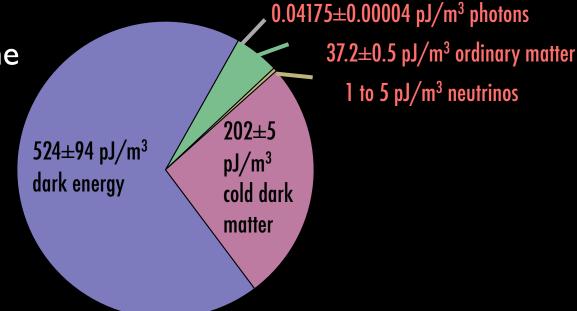
Weakly Interacting Massive Particles (WIMPs)

The magnificent WIMP

(Weakly Interacting Massive Particle)

 One naturally obtains the right cosmic density of WIMPs

Thermal production in hot primordial plasma.



One can experimentally test the WIMP hypothesis

The same physical processes that produce the right density of WIMPs make their detection possible

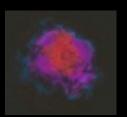
Cosmic density of thermal WIMPs

• At early times, WIMPs are produced in e^+e^- , $\mu^+\mu^-$, etc collisions in the hot primordial soup [thermal production].

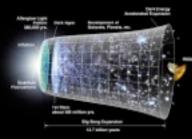
- WIMP production ceases when the production rate becomes smaller than the Hubble expansion rate [freeze-out].
- After freeze-out, there is a constant number of WIMPs in a volume expanding with the universe.

Indirect detection

the WIMP



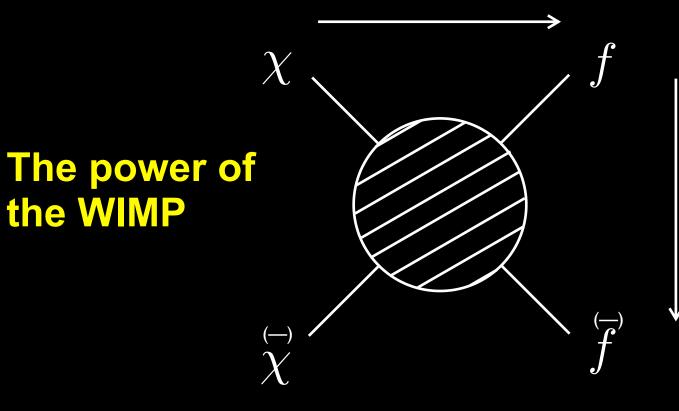




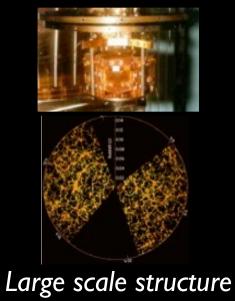
Cosmic density

Scattering

Annihilation



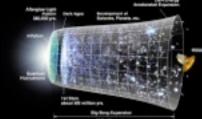
Direct detection



Production

Colliders

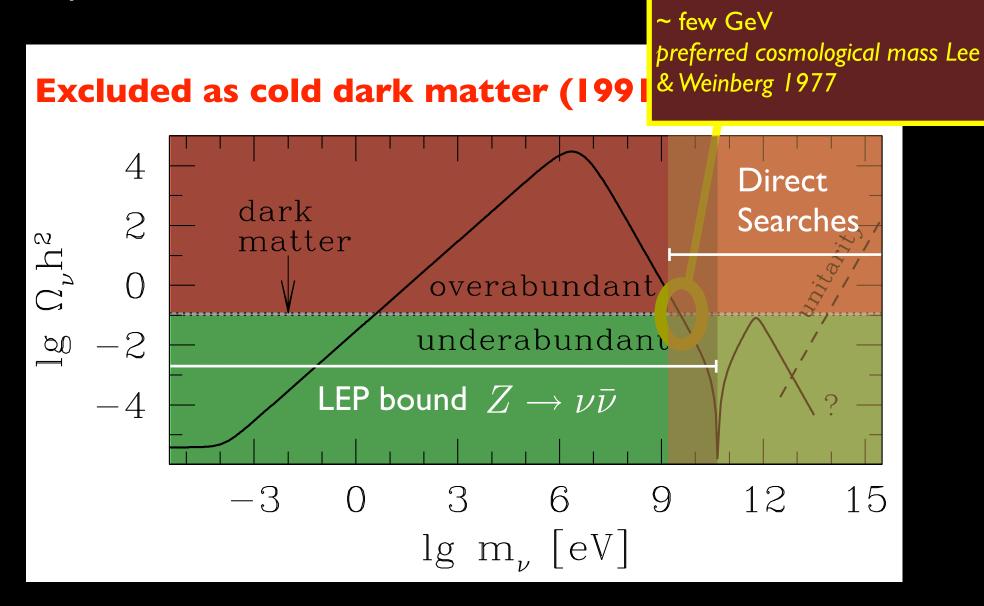




Cosmic density

Massive neutrinos as dark matter

Heavy active neutrino



Sterile neutrino dark matter

Standard model + right-handed neutrinos

Active and sterile neutrinos oscillate into each other.

10

 10^{-6}

Sterile neutrinos can be warm dark matter (mass > 0.3 keV)

Dodelson, Widrow 1994; Shi, Fuller 1999; Laine, Shaposhnikov 2008

MW

M31

MW

SPI

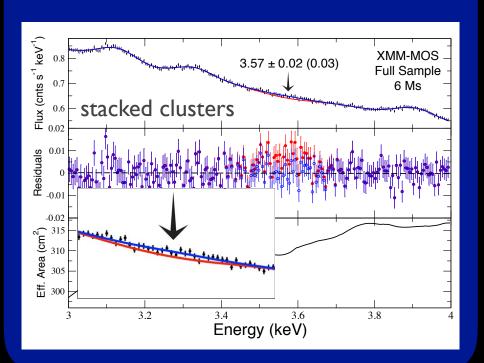
vMSM

Laine, Shaposhnikov 2008

Sterile neutrino dark matter

An unidentified 3.5-keV X-ray line has been reported in galaxy clusters and the Andromeda galaxy.

Bulbul et al 2014; Boyarski et al 2014; lakubovskyi et al 2015

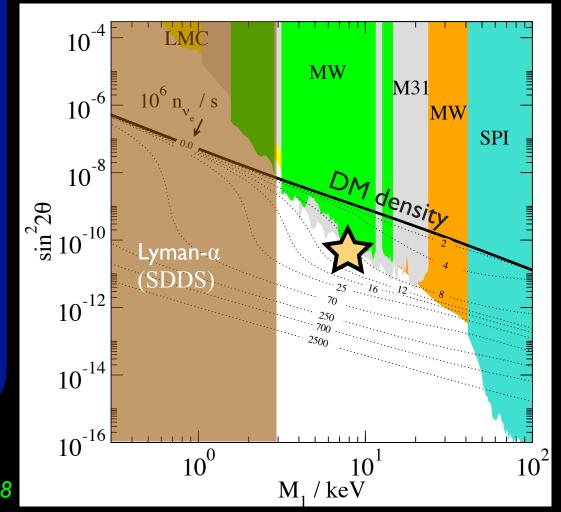


vMSM

Laine, Shaposhnikov 2008

Radiative decay of sterile neutrinos

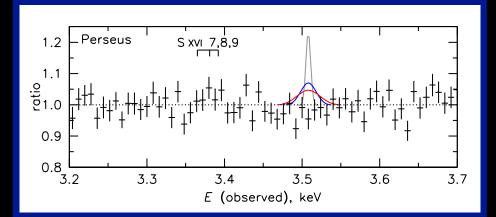
$$u_s o \gamma \nu_a \qquad E_\gamma = m_s/2$$
 $m_v = 7.1 \text{ keV} \qquad \sin^2(2\theta) = 7 \times 10^{-11}$



Sterile neutrino dark matter

The HITOMI data on the Perseus galaxy cluster do not show an X-ray line with the expected strength

Aharonian et al (HITOMI Collab.) 2016

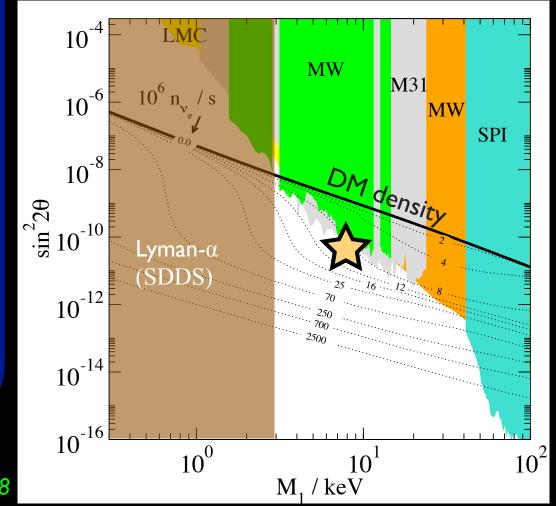


vMSM

Laine, Shaposhnikov 2008

Radiative decay of sterile neutrinos

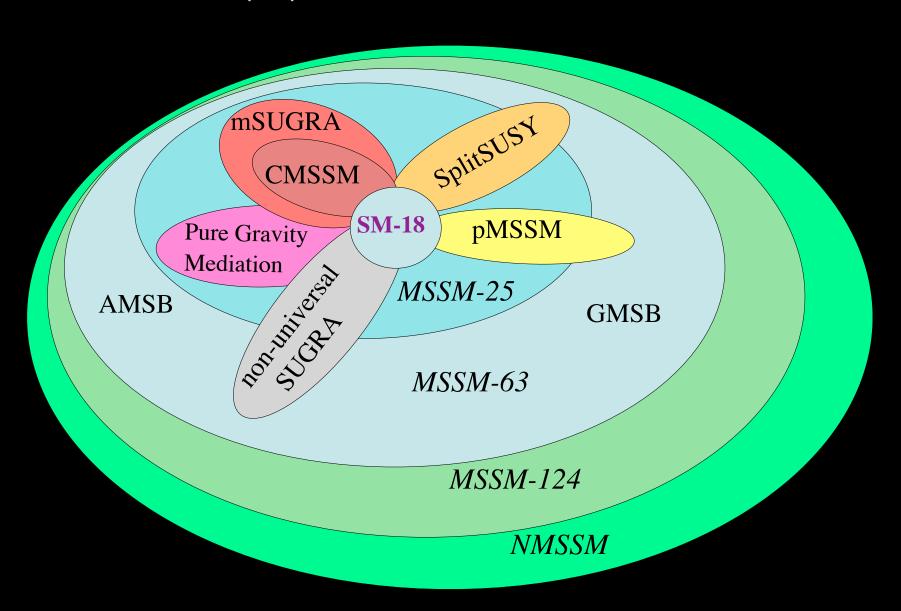
$$u_s o \gamma
u_a \qquad E_\gamma = m_s/2$$
 $m_v = 7.1 \text{ keV} \qquad \sin^2(2\theta) = 7 \times 10^{-11}$



Supersymmetric dark matter

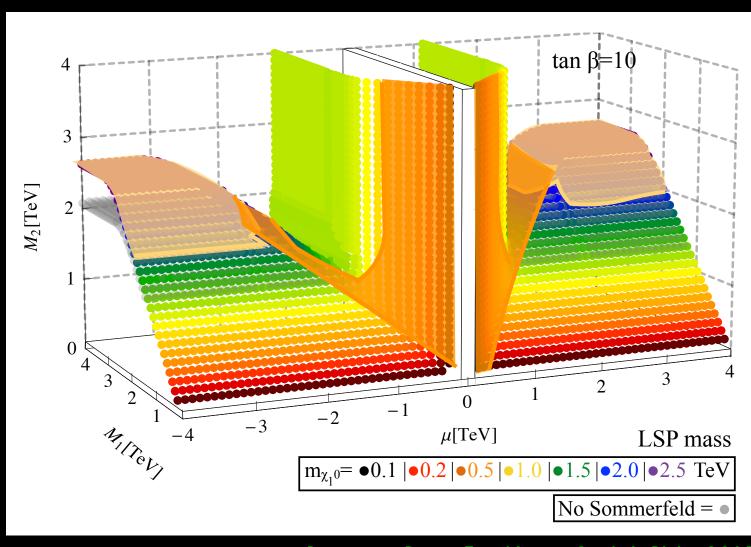
The CMSSM* is in dire straights, but there are many supersymmetric models

*Constrained Minimal Supersymmetric Standard Model



Supersymmetric dark matter

Neutralino dark matter with decoupled (heavy) sfermions



Excluded by LEP, HESS, LUX

All can be tested by LZ, CTA, and a I 00-TeV pp collider

Scalar phantom dark matter

"Gauge singlet scalar dark matter"

"Singlet scalar dark matter"

"Scalar singlet dark matter"

"Scalar Higgs-portal dark matter"

"The minimal model of dark matter"

Minimalist dark matter

do not confuse with minimal dark matter

Gauge singlet scalar field S stabilized by a \mathbb{Z}_2 symmetry $(S \rightarrow -S)$

$$\mathcal{L} = \frac{1}{2} \partial^{\mu} S \partial_{\mu} S + \frac{1}{2} \mu_S^2 S^2 - \frac{\lambda_S}{4} S^4 - \lambda_{HS} H^{\dagger} H S^2$$

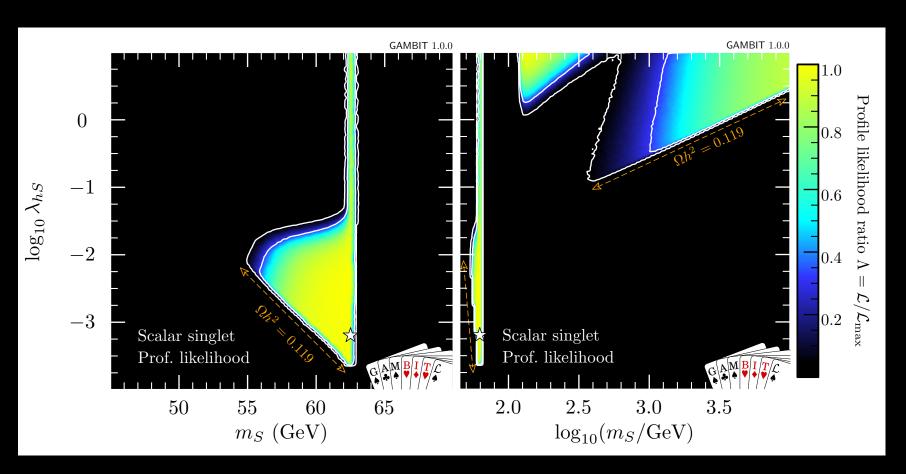
Silveira, Zee 1985; Andreas, Hambye, Tytgat 2008; Djouadi, Falkowksi, Mambrini, Quevillon 2012; Cline, Scott, Kainulainen, Weniger 2013; Hamada, Kawana 2015; Feng, Profumo, Ubaldi 2015; Athron et al 2017; etc

"Scalar phantom" is the original 1985 name

Scalar phantom dark matter

Global likelihood fits including collider, direct, indirect searches and cosmology.

The discovery of the Higgs boson at the LHC places the strongest constraints.



Anapole dark matter

The anapole moment is a C and P violating, but CP-conserving, electromagnetic moment

Zeldovich 1957

First measured experimentally in Cesium atoms

Wood et al 1997

Anapole dark matter

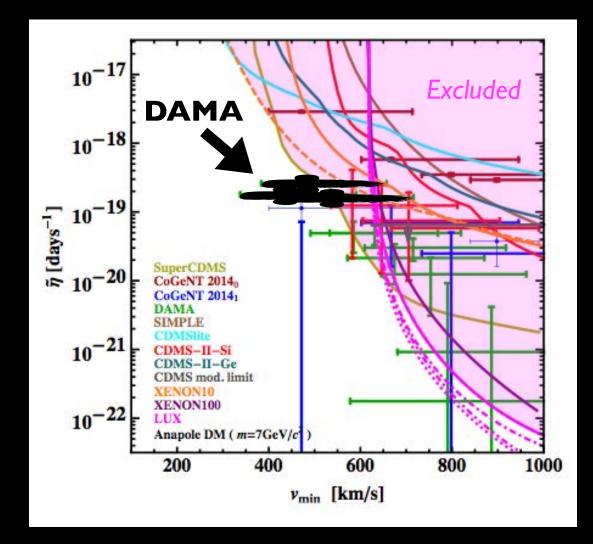
spin-1/2 Majorana fermion

$$\mathcal{L} = \frac{g}{2\Lambda^2} \bar{\chi} \gamma^{\mu} \gamma^5 \chi \, \partial^{\nu} F_{\mu\nu}$$

$$H = -\frac{g}{\Lambda^2} \, \vec{\sigma} \cdot \vec{\nabla} \times \vec{B}$$

The DAMA modulation is compatible with other searches for some WIMP velocity distribution

Del Nobile, Gelmini, Gondolo, Huh 2014



Other candidates

Self-interacting dark matter

Proposed as solution to 'cusp vs core' and 'too big to fail' puzzles in collisionless dark matter simulations.

Spergel, Steinhardt 1999

Tested on cluster and galaxy collisions

```
\sigma/m < 0.7~{\rm cm^2/g} mass loss in Bullet Cluster Randall et al 2008 \sigma/m < 0.47~{\rm cm^2/g} 72 cluster collisions Harvey et al 2015 \sigma/m = (1.7\pm0.7)\times10^{-4}~{\rm cm^2/g} in Abell 3827 (?) Massey et al 2015
```

Several particle models exist

```
Light mediator Feng, Kaplinghat, Yu; Buckley, Fox 2009; Tulin, Yu, Zurek 2013

Hidden vector dark matter (HVDM) Hambye 2008

Dark matter is 3 gauge bosons of a hidden SU(2) group spontaneously broken by a hidden Higgs doublet coupled to the SM Higgs.

Allowed in some regimes Bernal et al 2015
```

Asymmetric dark matter

- Dark matter in a hidden mirror sector ("dark sector")
- Dark matter asymmetry similar to baryon asymmetry, generated by similar mechanisms

$$n_{\chi} \approx n_{\rm p}$$

• Dark matter mass is a few times the proton mass

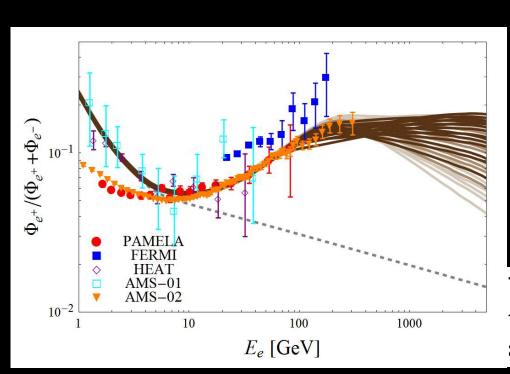
$$\Omega_{\chi} \approx \frac{m_{\chi}}{m_{p}} \, \Omega_{\rm p} \approx (\text{a few}) \, \Omega_{\rm p}$$

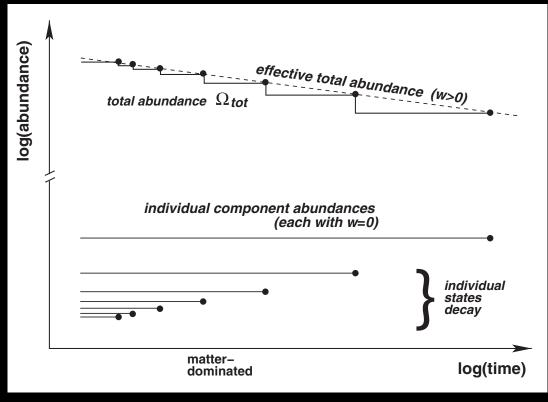
Dynamical dark matter

Dienes, Thomas 2011, 2012 Dienes, Kumar, Thomas 2012, 2013

A vast ensemble of fields decaying from one to another

Example: Kaluza-Klein tower of axions in extra-dimensions

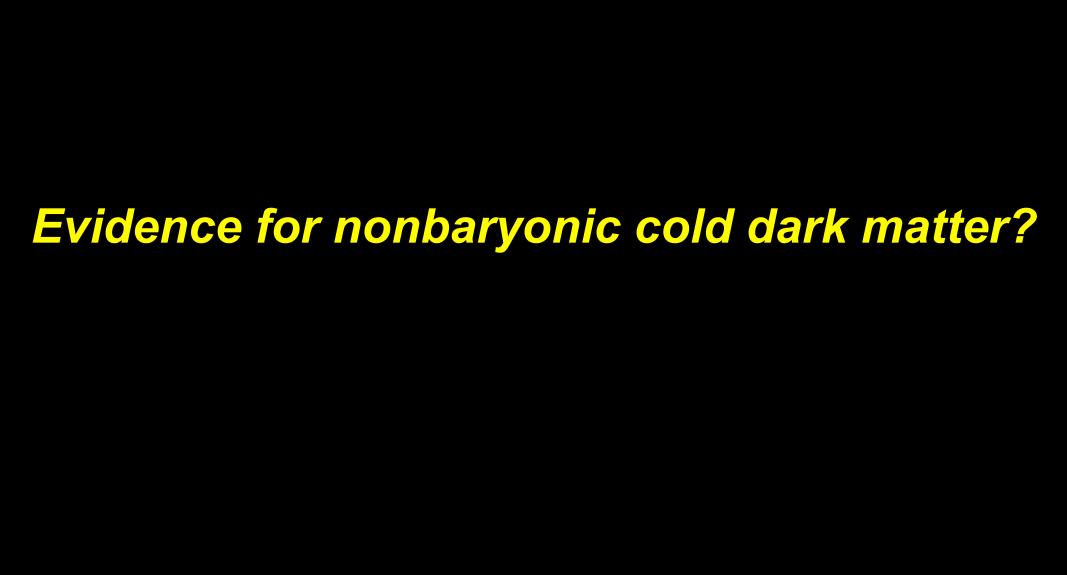




Phenomenology obtained through scaling laws

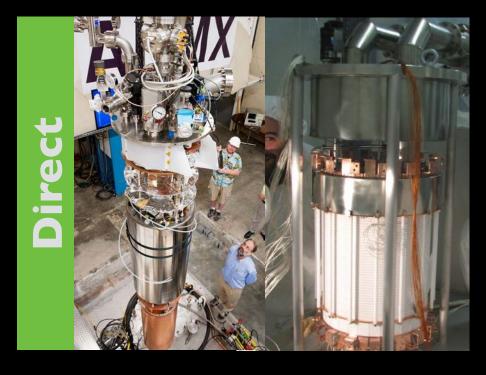
$$m_n = m_0 + n^{\delta} \Delta m,$$
$$\rho_n \sim m_n^{\alpha}, \, \tau_n \sim m_n^{-\gamma}$$

This model can fit the positron excess and has no cut off.



Searches for particle dark matter









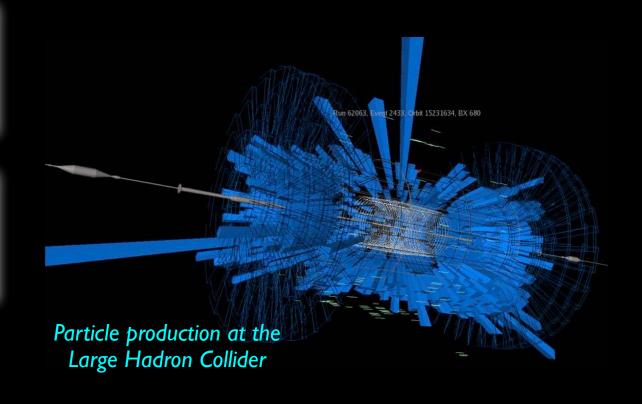
Searches for particle dark matter: colliders

Dark matter particles (or their "cousins") are produced in high-energy collisions

Dark matter particles are produced and escape detection (missing energy)

Charged/colored "cousins" of the dark matter particle are produced

LEP ALEPH, DELPHI, OPAL, ... Tevatron CDF, D0, ... LHC ATLAS, CMS, ...



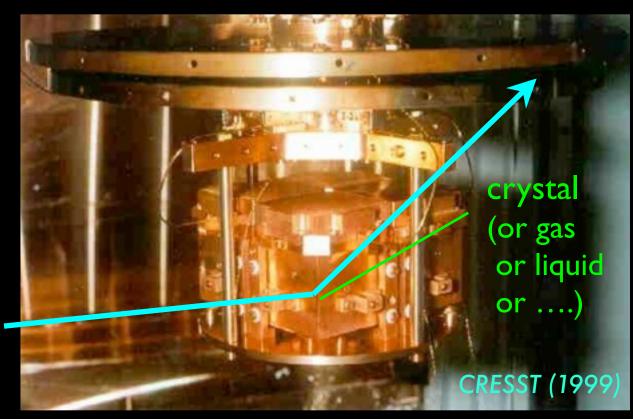
Searches for particle dark matter: direct

Dark matter particles that arrive on Earth interact with a laboratory detector

WIMP scattering off nuclei or electrons

Goodman, Witten 1985

Dark matter particle



Low-background underground detector

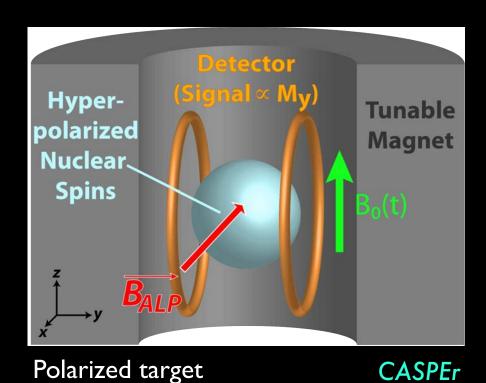
DAMA, SuperCDMS, XENON, LUX, XMASS, PICO, CoGeNT, DEAP, DRIFT, ANAIS, CRESST, LZ, DARWIN, DM-ICE, NEWAGE, ...

Searches for particle dark matter: direct

Dark matter particles that arrive on Earth interact with a laboratory detector

Axion conversion to photons

Primakoff Conversion Sikivie 1983 T_N B₀ Dark Single real matter photon particle Virtual photon Magnetic field Resonating cavity **ADMX** Coherent axion field acts as electromagnetic current



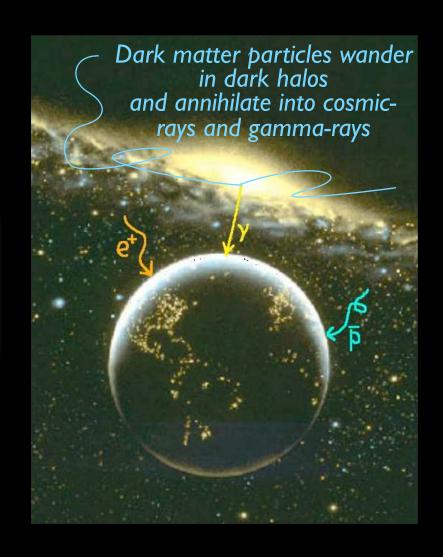
ADMX

Searches for particle dark matter: indirect

Dark matter particles transform into ordinary particles, which are then detected or inferred

Gunn, Lee, Lerche, Schramm, Steigman 1978; Stecker 1978

Gamma-rays, positrons, antiprotons from our galaxy and beyond



cosmic-rays
PAMELA
AMS
...
gamma-rays
MAGIC
HESS
VERITAS
Fermi-LAT
HAWK
CTA
...

Searches for particle dark matter: indirect

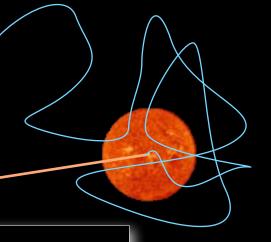
Dark matter particles transform into ordinary particles, which are then detected or inferred



IceCube ANTARES

•••

Dark matter particles sink into the Sun/Earth where they transform into neutrinos



Neutrinos from the Sun

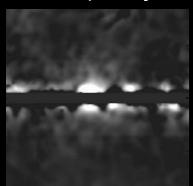
Press, Spergel 1985; Silk, Olive, Srednicki 1985

Neutrinos from the Earth

Freese 1986; Krauss, Srednicki, Wilczek 1986

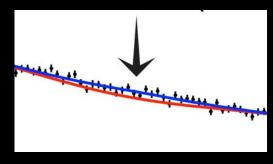
Signals from dark matter?

GeV γ -rays



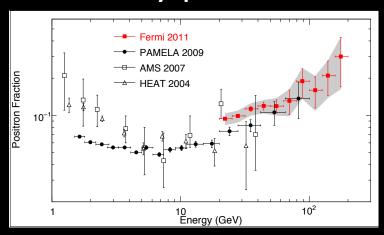
Hooper et al 2009-14

3.5 keV X-ray line



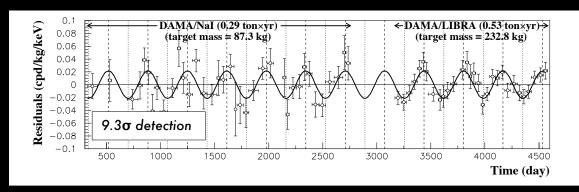
Bulbul et al 2014

Cosmic-ray positrons



Adriani et al 2009; Ackerman et al 2011; Aguilar et al 2013

Annual modulation in direct detection



Bernabei et al 1997-now

Gamma-ray GeV excess

Gamma-rays from WIMP annihilation

factor

annihilation

$$\frac{d^2\phi}{d\Omega dE} = \frac{\langle \sigma v \rangle}{8\pi m_{\chi}^2} \frac{dN_{\gamma}}{dE} \times \int_{\text{l.o.s}} \rho^2 ds$$

Galactic DM Halo

- good S/N
- difficult background
- angular information

Galactic Center

- brightest DM source
- bright background

DM clumps

- no baryons
- bright enough?
- boost overall signal

Extragalactic

- nearly isotropicvisible near Galactic Poles
- angular information
- galaxy clusters?

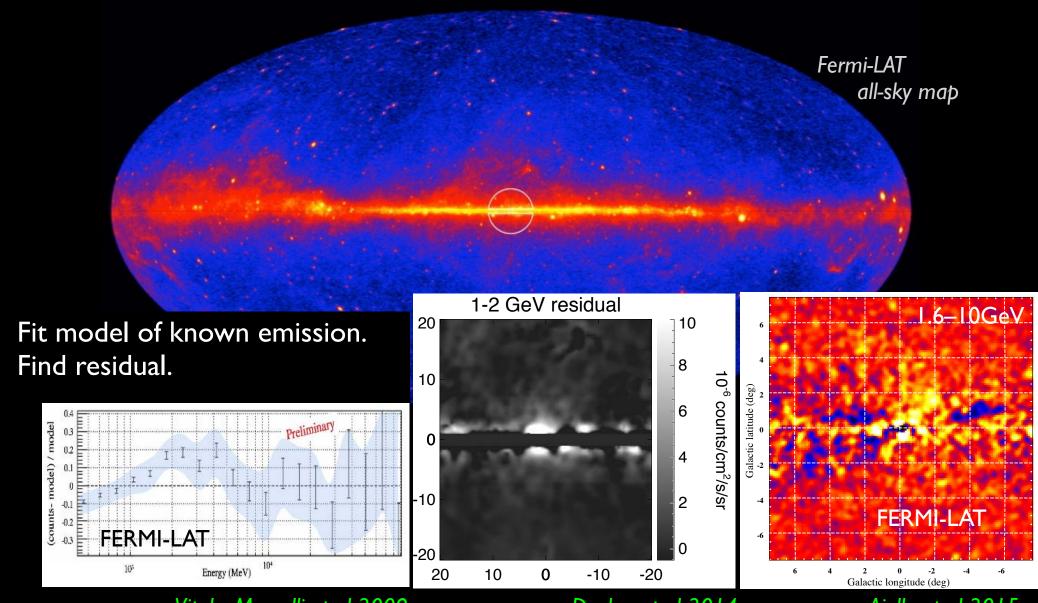
Dwarf Spheroidal Galaxies - harbor small number of stars

- otherwise dark (no γ -ray emission)

Kuhlen, Diemand, Madau 2007

1 GeV y-ray excess at Galactic Center

Goodenough, Hooper; Vitale, Morselli et al 2009; Hooper, Goodenough; Boyarsky, Malyshev, Ruchayskiy; Hooper, Linden 2011; Abazajian, Kaplinghat 2012; Gordon, Macias 2013; Abazajian, Canac, Horiuchi, Kaplinghat; Daylan et al; Calore, Cholis, Weniger 2014



Vitale, Morselli et al 2009

Daylan et al 2014

Ajello et al 2015

1 GeV γ-ray excess at Galactic Center

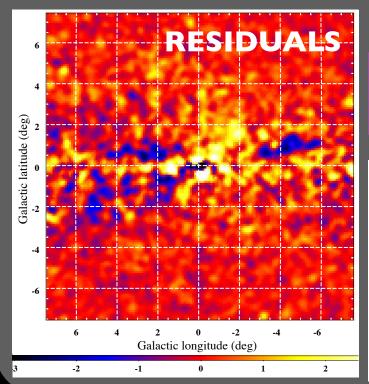
Goodenough, Hooper; Vitale, Morselli et al 2009; Hooper, Goodenough; Boyarsky, Malyshev, Ruchayskiy;

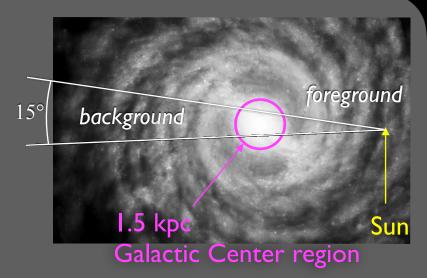
Hoober Lindon 2011-Abazaiian Kablinghat 2012-Cordon Maciae 2012-Abazaiian Canac Horiuchi

The official Fermi-LAT analysis

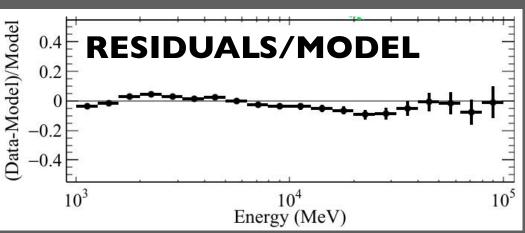
Ajello et al 2016

- **–** Four foreground+background cosmic-ray models:
 - Intensity-scaled or index-scaled pulsars sources
 - Intensity-scaled or index-scaled OB-stars sources
- Point sources



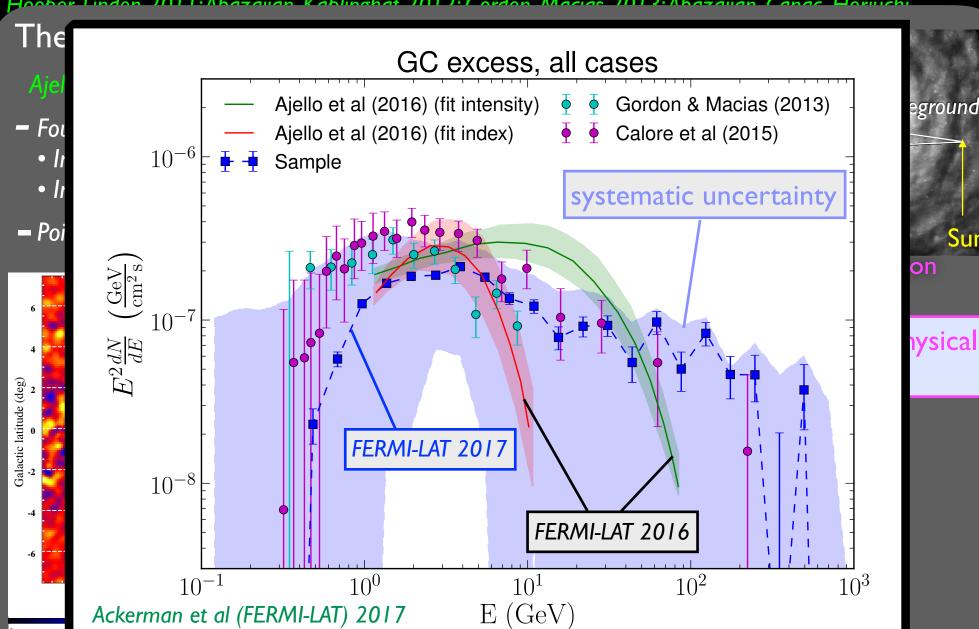


An extended residual is present. A precise physical interpretation of its origin is premature.



1 GeV γ-ray excess at Galactic Center

Goodenough, Hooper; Vitale, Morselli et al 2009; Hooper, Goodenough; Boyarsky, Malyshev, Ruchayskiy;



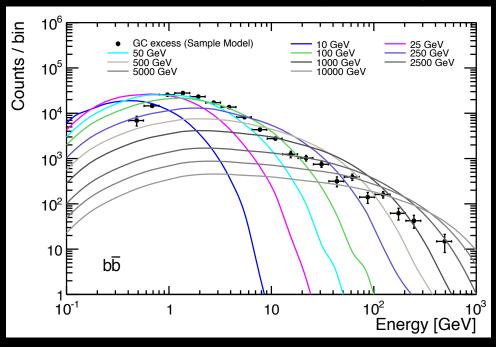
1 GeV γ-ray excess at Galactic Center

Dark matter annihilation

Goodenough, Hooper 2009; Hooper, Goodenough; Hooper, Linden 2011; Abazajian, Kaplinghat 2012; Abazajian, Canac, Horiuchi, Kaplinghat; Daylan et al; Calore, Cholis, Weniger 2014;

Possible for specific WIMP and dark halo models

Ackerman et al (FERMI-LAT) 2017



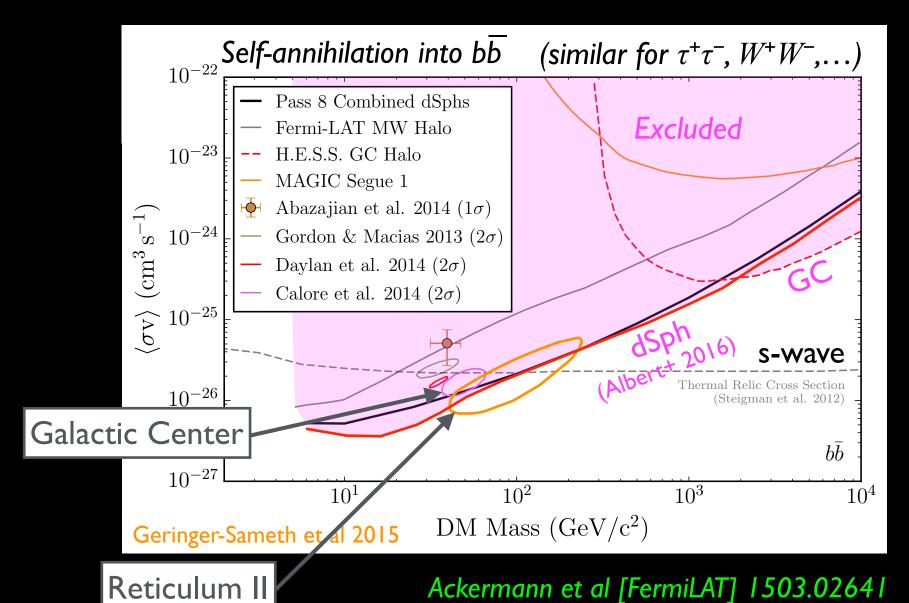
Millisecond pulsars

Wang et al 2005; Abazajian 2011; Gordon, Macias 2013; Hooper et al 2013; Yuan, Zhang 2014; Calore et al 2014; Cholis et al 2014; Petrovic et al 2014; Lee et al 2014; Bartels et al 2014

Favored by wavelet analysis and nonpoissonian point spread function Lee, Lisanti, Safdi 2014; Bartels, Krishnamurthy, Weniger 2015

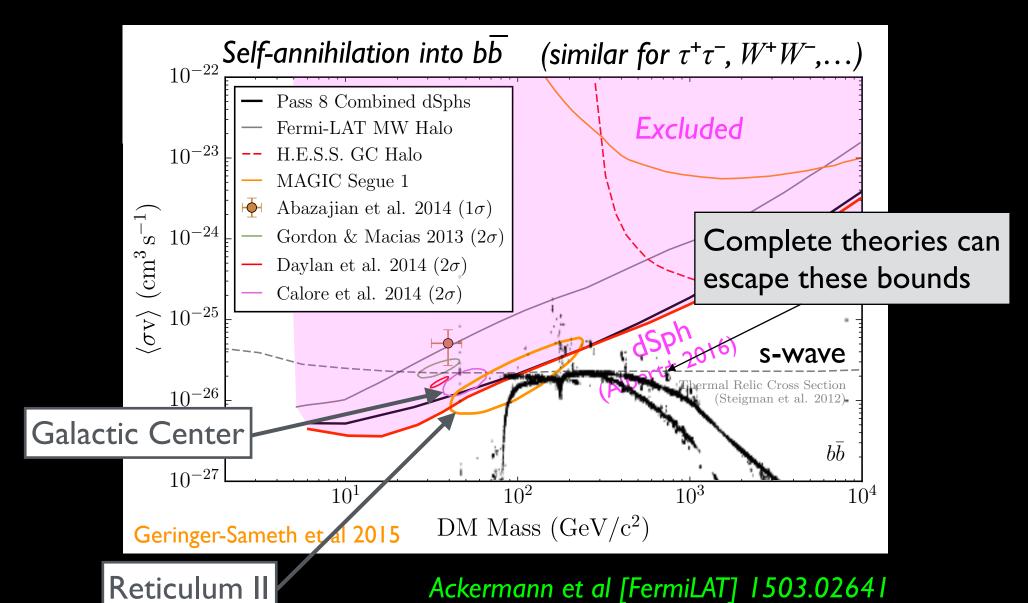
Gamma-rays from dark matter

Upper limits on the WIMP annihilation cross section from dwarf spheroidal galaxies and Galactic Center



Gamma-rays from dark matter

Upper limits on the WIMP annihilation cross section from dwarf spheroidal galaxies and Galactic Center

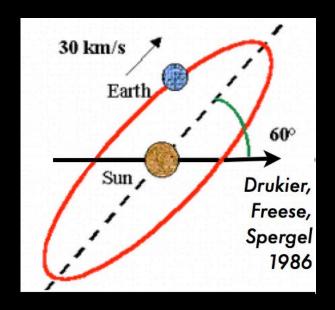


DAMA annual modulation

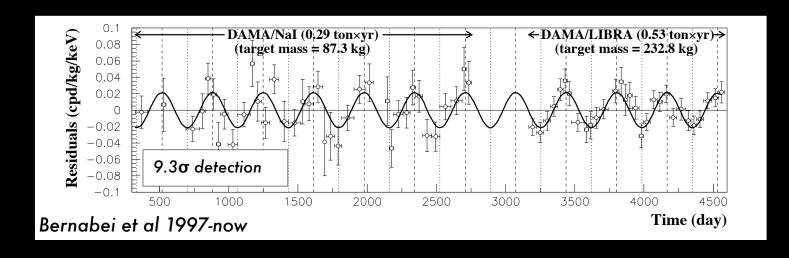
Annual modulation in direct detection

 The revolution of the Earth around the Sun modulates the WIMP event rate

Drukier, Freese, Spergel 1986



DAMA observes such kind of modulation



DAMA annual modulation

Model Independent Annual Modulation Result

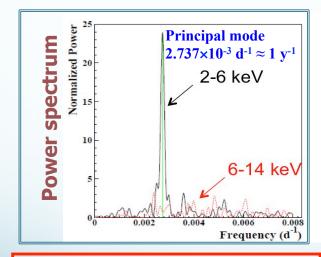
DAMA/NaI + DAMA/LIBRA-phase1 Total exposure: 487526 kg×day = 1.33 ton×yr

EPJC 56(2008)333, EPJC 67(2010)39, EPJC 73(2013)2648

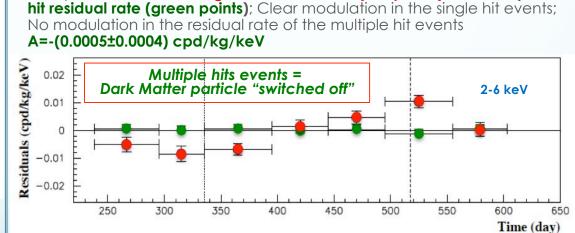
The measured modulation amplitudes (A), period (T) and phase (t_0) from the single-hit residual rate vs time

	A(cpd/kg/keV)	$T=2\pi/\omega$ (yr)	to (day)	C.L.
DAMA/NaI+DAMA/LIBRA-phase1				
(2-4) keV	0.0190 ±0.0020	0.996 ±0.0002	134 ± 6	9.50
(2-5) keV	0.0140 ±0.0015	0.996 ±0.0002	140 ± 6	9.3σ
(2-6) keV	0.0112 ±0.0012	0.998 ±0.0002	144 ± 7	9.3σ

 $A\cos[\omega(t-t_0)]$



No systematics or side reaction able to account for the measured modulation amplitude and to satisfy all the peculiarities of the signature



Comparison between single hit residual rate (red points) and multiple

This result offers an additional strong support for the presence of DM particles in the galactic halo further excluding any side effect either from hardware or from software procedures or from background

The data favor the presence of a modulated behaviour with all the proper features for DM particles in the galactic halo at about 9.2 σ C.L.

DAMA annual modulation

Model Independent Annual Modulation Result

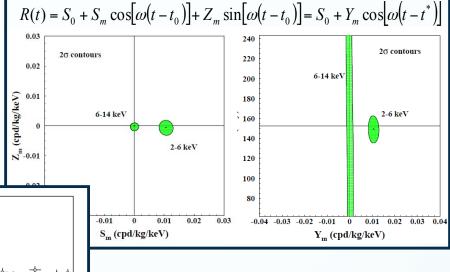
DAMA/NaI + DAMA/LIBRA-phase1 Total exposure: 487526 kg×day = 1.33 ton×yr

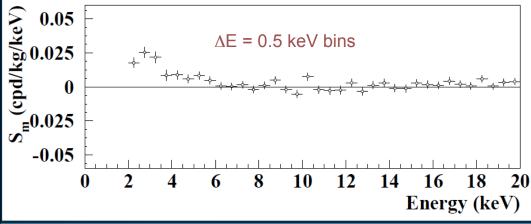
EPJC 56(2008)333, EPJC 67(2010)39, EPJC 73(2013)2648

- No modulation above 6 keV
- No modulation in the whole energy spectrum
- No modulation in the 2-6 keV multiple-hit events

$$R(t) = S_0 + S_m \cos \left[\omega \left(t - t_0\right)\right]$$

here $T = 2\pi/\omega = 1$ yr and $t_0 = 152.5$ day

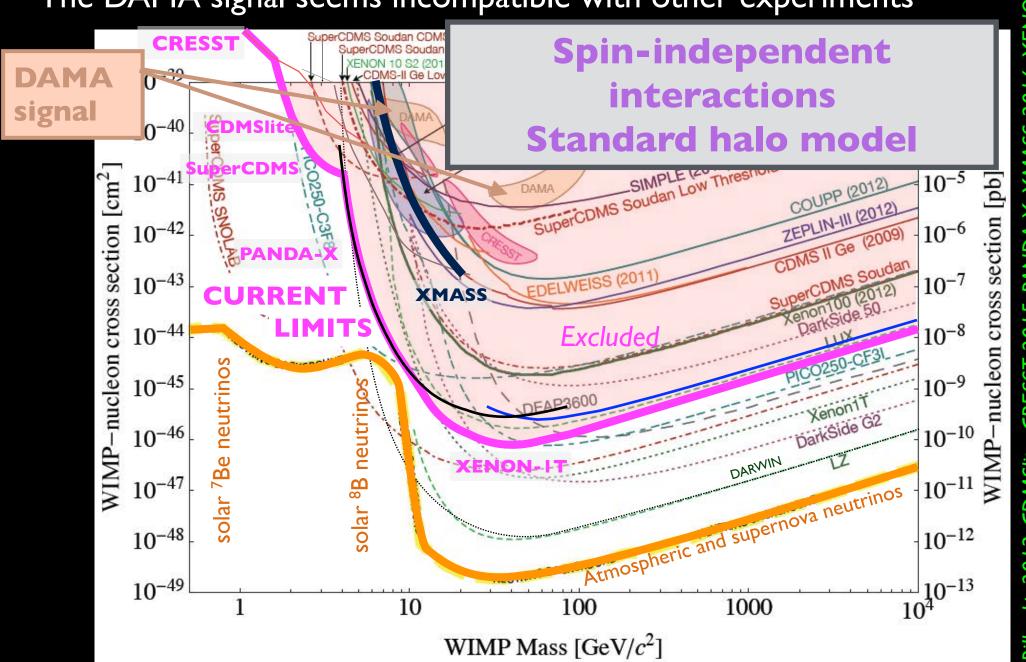




No systematics or side processes able to quantitatively account for the measured modulation amplitude and to simultaneously satisfy the many peculiarities of the signature are available.

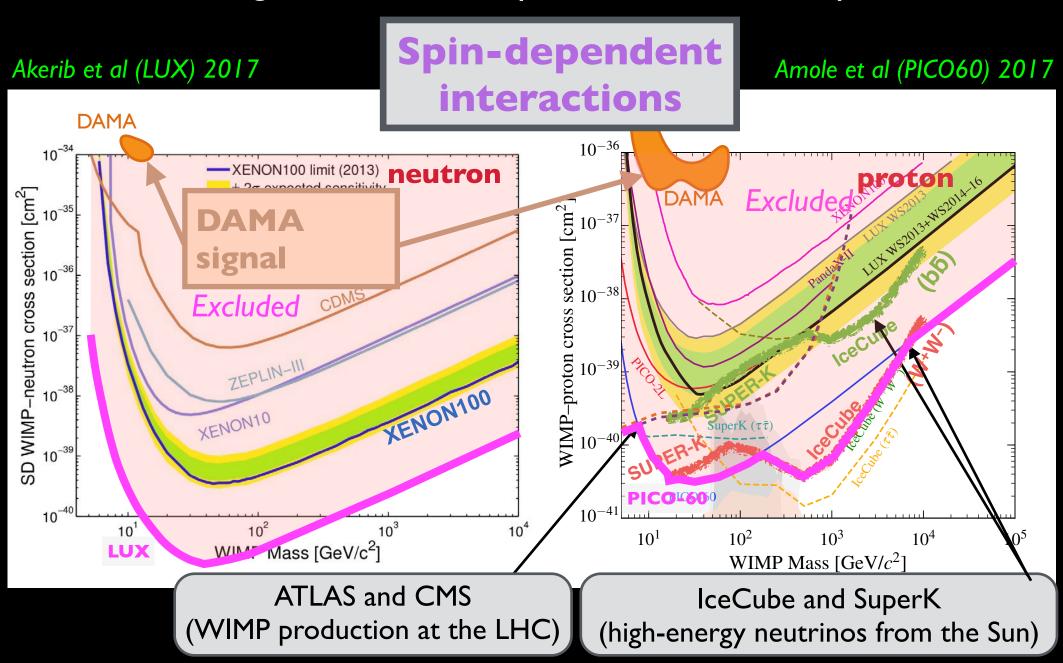
Direct evidence for dark matter particles?

The DAMA signal seems incompatible with other experiments



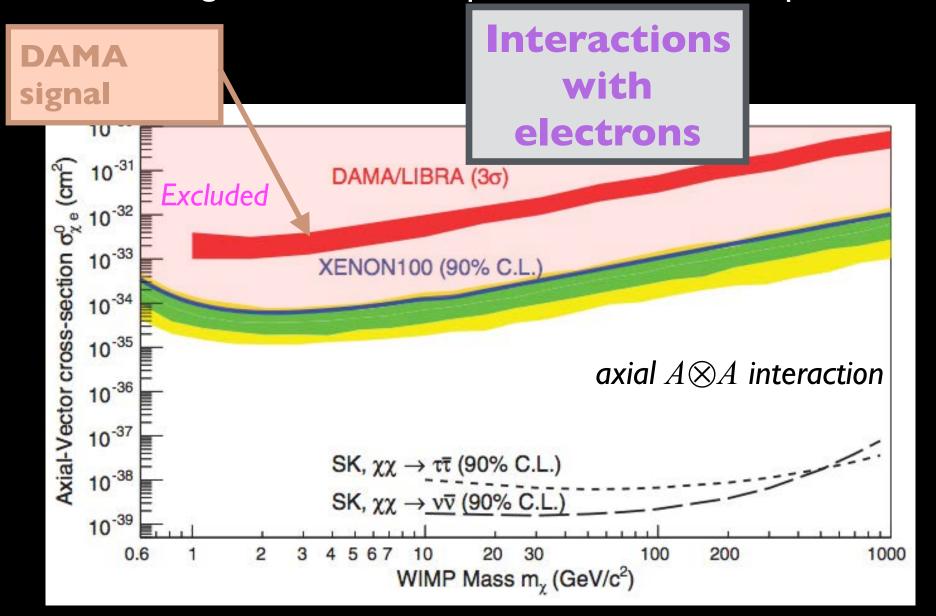
Direct evidence for dark matter particles?

The DAMA signal seems incompatible with other experiments



Direct evidence for dark matter particles?

The DAMA signal seems incompatible with other experiments



Recent trends: "make no assumptions"

"Make no assumptions"

All particle physics models

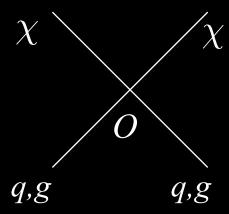
- Consider all possible interactions between dark matter and standard model particles
- This program has been carried out in some limits (e.g., non-relativistic conditions, heavy mediators)

All astrophysical models

- Halo-independent methods of analysis have been developed
- Ideally they require no assumption on the astrophysical density and velocity distributions of dark matter particles

Effective operators

if mediator mass » exchanged energy



Four-particle effective operator

There are many possible operators.

Interference is important although often, but not always, neglected.

Long(ish) distance interactions are not included.

Effective operators: LHC & direct detection

Name	Operator	Coefficient
D1	$ar{\chi}\chiar{q}q$	m_q/M_*^3
D2	$\bar{\chi}\gamma^5\chi \bar{q}q$	im_q/M_*^3
D3	$\bar{\chi}\chi\bar{q}\gamma^5q$	im_q/M_*^3
D4	$\bar{\chi}\gamma^5\chi\bar{q}\gamma^5q$	m_q/M_*^3
D5	$\bar{\chi}\gamma^{\mu}\chi\bar{q}\gamma_{\mu}q$	$1/M_*^2$
D6	$\bar{\chi}\gamma^{\mu}\gamma^5\chi\bar{q}\gamma_{\mu}q$	$1/M_*^2$
D7	$\bar{\chi}\gamma^{\mu}\chi\bar{q}\gamma_{\mu}\gamma^5q$	$1/M_*^2$
D8	$\left \bar{\chi}\gamma^{\mu}\gamma^5\chi\bar{q}\gamma_{\mu}\gamma^5q\right $	$1/M_*^2$
D9	$\bar{\chi}\sigma^{\mu\nu}\chi\bar{q}\sigma_{\mu\nu}q$	$1/M_*^2$
D10	$\left \bar{\chi}\sigma_{\mu\nu}\gamma^5\chi\bar{q}\sigma_{\alpha\beta}q\right $	i/M_*^2
D11	$\bar{\chi}\chi G_{\mu\nu}G^{\mu\nu}$	$\alpha_s/4M_*^3$
D12	$\bar{\chi}\gamma^5\chi G_{\mu\nu}G^{\mu\nu}$	$i\alpha_s/4M_*^3$
D13	$\bar{\chi}\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$i\alpha_s/4M_*^3$
D14	$\bar{\chi}\gamma^5\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$\alpha_s/4M_*^3$

Name	Operator	Coefficient
C1	$\chi^\dagger \chi ar q q$	m_q/M_*^2
C2	$\chi^{\dagger}\chi \bar{q}\gamma^5 q$	im_q/M_*^2
C3	$\chi^{\dagger}\partial_{\mu}\chi \bar{q}\gamma^{\mu}q$	$1/M_{*}^{2}$
C4	$\chi^{\dagger} \partial_{\mu} \chi \bar{q} \gamma^{\mu} \gamma^5 q$	$1/M_{*}^{2}$
C5	$\chi^{\dagger}\chi G_{\mu\nu}G^{\mu\nu}$	$\alpha_s/4M_*^2$
C6	$\chi^{\dagger} \chi G_{\mu\nu} \tilde{G}^{\mu\nu}$	$i\alpha_s/4M_*^2$
R1	$\chi^2 \bar{q} q$	$m_q/2M_*^2$
R2	$\chi^2 \bar{q} \gamma^5 q$	$im_q/2M_*^2$
R3	$\chi^2 G_{\mu\nu} G^{\mu\nu}$	$\alpha_s/8M_*^2$
R4	$\chi^2 G_{\mu\nu} \tilde{G}^{\mu\nu}$	$i\alpha_s/8M_*^2$

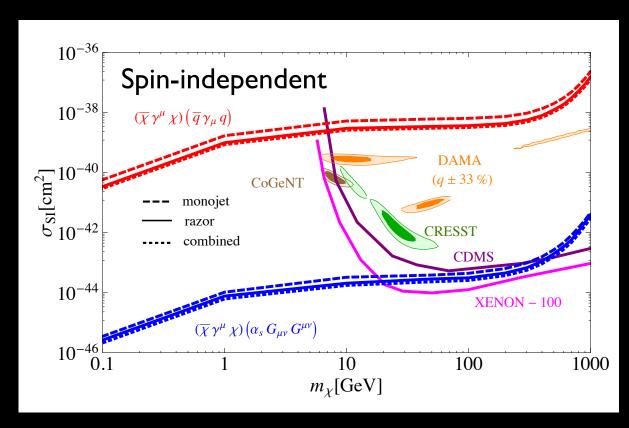
Table of effective operators relevant for the collider/direct detection connection

Goodman, Ibe, Rajaraman, Shepherd, Tait, Yu 2010

Effective operators: LHC & direct detection

LHC limits on WIMP-quark and WIMP-gluon interactions are competitive with direct searches

Beltran et al., Agrawal et al., Goodman et al., Bai et al., 2010; Goodman et al., Rajaraman et al. Fox et al., 2011; Cheung et al., Fitzptrick et al., March-Russel et al., Fox et al., 2012......



These bounds do not apply to SUSY, etc.

Complete theories contain sums of operators (interference) and not-so-heavy mediators (Higgs)

Fox, Harnik, Primulando, Yu 2012

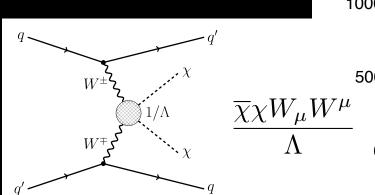
Effective operators: LHC & indirect detection

Limits on WIMP couplings to vector bosons (γ , W, g, ...)

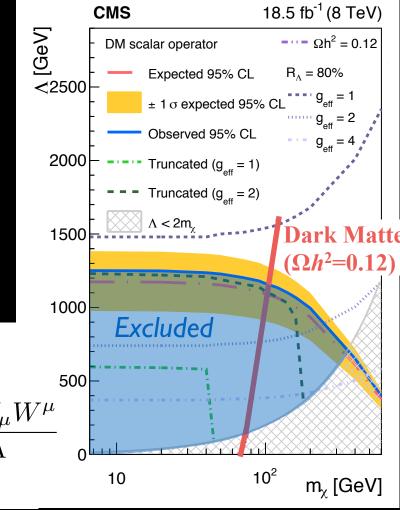
Name	Expression	Norm.	Vertices	Sub-Procs.	Ann.
$\overline{dim = 5:}$					
D5a	$\bar{\chi}\chi V^{a\mu}V^a_\mu$	Λ^{-1}	4pt	ZZ,WW	v^2
D5b	$\bar{\chi}i\gamma_5\chi V^{a\mu}V_{\mu}^a$	Λ^{-1}	4pt	ZZ,WW	1
D5c	$\bar{\chi}\sigma_{\mu\nu}t^a\chi V^{a\mu\nu}$	Λ^{-1}	3/4pt	A,Z,WW	1
D5d	$\bar{\chi}\sigma_{\mu\nu}t^a\chi\widetilde{V}^{a\mu\nu}$	Λ^{-1}	3/4pt	A,Z,WW	$1 (VV), v^2 (f\bar{f})$
dim = 6:					
D6a	$\bar{\chi}\gamma_{\mu}t^{a}D_{\nu}\chi V^{a\mu\nu}$	Λ^{-2}	3/4pt	A,Z,WW	1
D6b	$\bar{\chi}\gamma_{\mu}\gamma_{5}t^{a}D_{\nu}\chi V^{a\mu\nu}$	Λ^{-2}	3/4pt	A,Z,WW	$1 (VV), v^2 (f\bar{f})$
dim = 7:					
D7a	$\bar{\chi}\chi V^{\mu\nu}V_{\mu\nu}$	Λ^{-3}	4pt	AA,AZ,ZZ,WW	v^2
D7b	$\bar{\chi}i\gamma_5\chi V^{\mu\nu}V_{\mu\nu}$	Λ^{-3}	4pt	AA,AZ,ZZ,WW	1
D7c	$\bar{\chi}\chi V^{\mu u}\widetilde{V}_{\mu u}$	Λ^{-3}	4pt	AA,AZ,ZZ,WW	v^2
D7d	$\bar{\chi}i\gamma_5\chi V^{\mu u}\widetilde{V}_{\mu u}$	Λ^{-3}	4pt	AA,AZ,ZZ,WW	1

Cotta, Hewett, Le, Rizzo 2013

Limits on non-standardmodel dijets with vector boson fusion topology



The CMS Collaboration, 2016



Effective operators: non relativistic

Nonrelativistic WIMP-nucleon contact operators classified

Barger et al 2008, Fan et al 2010, Fitzpatrick et al 2012, Dent et al 2015

$$\mathcal{O}_{1} = 1_{\chi} 1_{N} \qquad \qquad \mathcal{O}_{7} = \vec{S}_{N} \cdot \vec{v}_{\chi N}^{\perp}$$

$$\mathcal{O}_{3} = -i \vec{S}_{N} \cdot \left(\frac{\vec{q}}{m_{N}} \times \vec{v}_{\chi N}^{\perp}\right) \qquad \mathcal{O}_{8} = \vec{S}_{\chi} \cdot \vec{v}_{\chi N}^{\perp}$$

$$\mathcal{O}_{4} = \vec{S}_{\chi} \cdot \vec{S}_{N} \qquad \qquad \mathcal{O}_{9} = -i \vec{S}_{\chi} \cdot \left(\vec{S}_{N} \times \frac{\vec{q}}{m_{N}}\right)$$

$$\mathcal{O}_{5} = -i \vec{S}_{\chi} \cdot \left(\frac{\vec{q}}{m_{N}} \times \vec{v}_{\chi N}^{\perp}\right) \qquad \qquad \mathcal{O}_{10} = -i \vec{S}_{N} \cdot \frac{\vec{q}}{m_{N}}$$

$$\mathcal{O}_{6} = \left(\vec{S}_{\chi} \cdot \frac{\vec{q}}{m_{N}}\right) \left(\vec{S}_{N} \cdot \frac{\vec{q}}{m_{N}}\right) \qquad \qquad \mathcal{O}_{11} = -i \vec{S}_{\chi} \cdot \frac{\vec{q}}{m_{N}}$$

and more in Barger et al. 2008, Fan et al. 2010, Dent et al 2015

To leading order in q and v, only O_1 and O_4 appear, which are the spin-independent and spin-dependent terms, respectively.

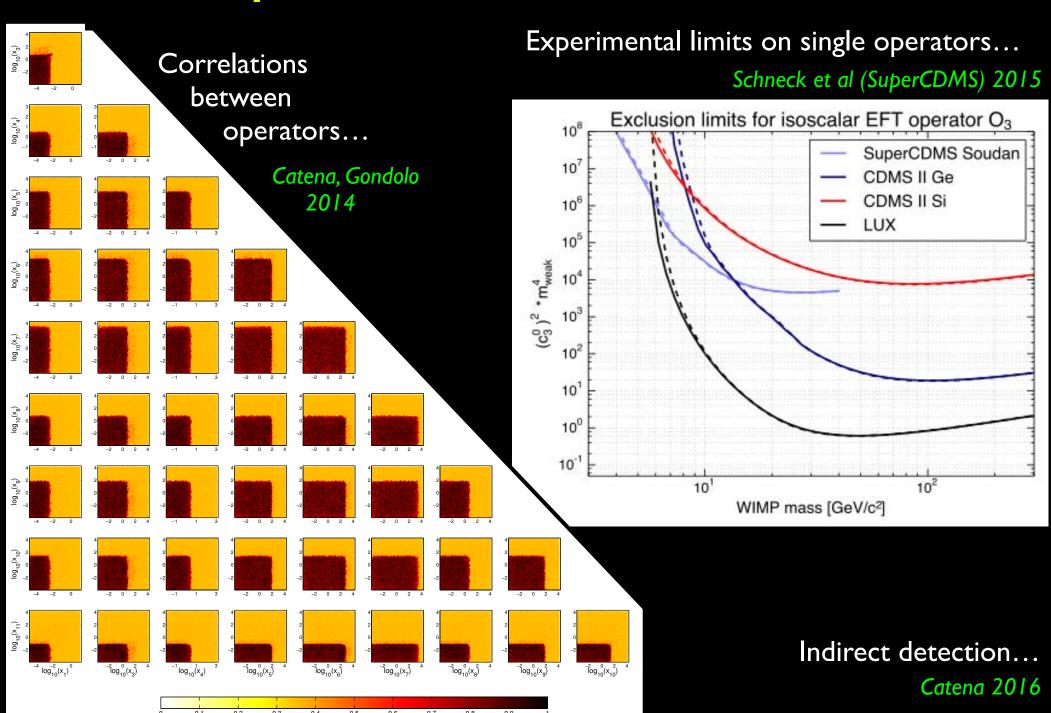
Effective operators: non relativistic

Operators appearing in "simplified" WIMP models Dent et al 2015

	$s_{\chi}=0$ X^{0}	s _χ =0 X ⁺	$ \begin{array}{c c} s_{\chi}=1/2 \\ X^{0} \end{array} $	$s_{\chi}=1/2$ X^{+}	$s_{\chi}=1$ X^{0}	$s_{\chi}=1$ X^{+}
S	1	1	1,11	1,11	1	1
V	_		1,8,9	1,8,9	4,5,6,8,9,11,17	
Т		_	_	4,10,11,12		4,11,12,18
A	10		4,7,9	4,7,9	4,6,9,10,14,18	
P	10	10	6,10	6,10	10	10

Entries denote operator index. X^0 (X^+) is neutral (charged) mediator.

Effective operators: non relativistic



Effective operators: nucleon matrix elements

To connect high-energy theories to the nonrelativistic contact operators one must obtain WIMP-nucleon interactions from WIMP-quark and WIMP-gluon interactions.

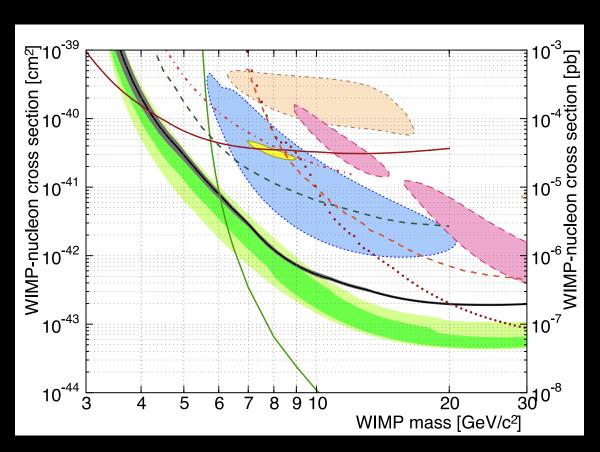
Nucleon matrix elements of quark and gluon currents

$$\langle N|\overline{q}\Gamma_{i}q|N\rangle = \sum_{j} f_{ij}^{(q,N)} \overline{N}\Gamma_{j}N \qquad (q = u, d, s)$$
$$\langle N|G_{\mu\nu}^{a}G_{\lambda\rho}^{a}h_{i}^{\mu\nu\lambda\rho}|N\rangle = \sum_{j} f_{ij}^{(g,N)} \overline{N}\Gamma_{j}N$$

See e.g. Kaplan, Manohar 1988; Cheng 1989; Drees, Nojiri 1993; Adam+ 1995; Aoki+ 1997; Mallot 1999; Pospelov&Ritz, Leinweber+ 2004; Doi+ 2009; Alekseev+ 2010; Bacchetta+, Bali+, Hisano+ 2012; Anselmino+, Dienes+, Fuyuto+ 2013; Hill&Solon 2014; Agrawal+, Bhattacharya+, Hisano+ 2015, ...

Systematic analysis under way: some large uncertainties, some unknown matrix elements,

$$\begin{pmatrix} \text{event} \\ \text{rate} \end{pmatrix} = \begin{pmatrix} \text{detector} \\ \text{response} \end{pmatrix} \times \begin{pmatrix} \text{particle} \\ \text{physics} \end{pmatrix} \times \begin{pmatrix} \text{astrophysics} \end{pmatrix}$$
FIXED



Standard Halo Model

truncated Maxwellian

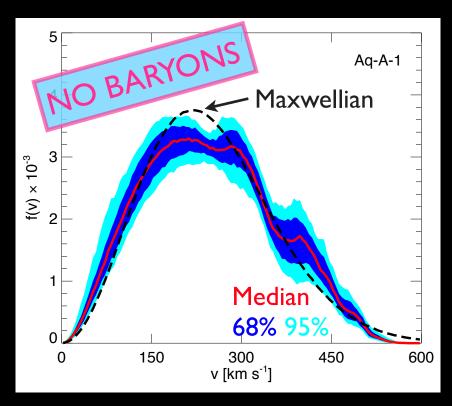
$$f(\vec{v}) = C e^{-|\vec{v} + \vec{v}_{\text{obs}}|/\bar{v}_0^2} \Theta(v - v_{\text{esc}})$$



The spherical cow of direct WIMP searches

Gelmini

We know very little about the dark matter velocity distribution near the Sun

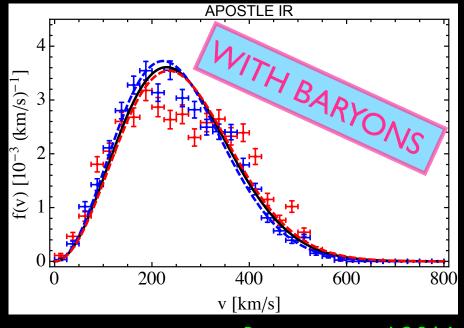


Vogelsberger et al 2009

Cosmological N-Body simulations including baryons are challenging but underway



Odenkirchen et al 2002 (SDSS)



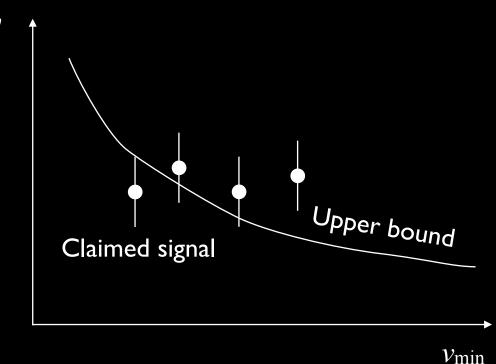
Bozorgnia et al 2016

Rescaled astrophysics factor common to all experiments

$$\eta(v_{\min}) = \frac{\rho_{\chi}}{m_{\chi}} \int_{v_{\min}}^{\infty} \frac{f(\mathbf{v})}{v} d^3v$$

"Velocity integral"

Proxy for dark matter flux

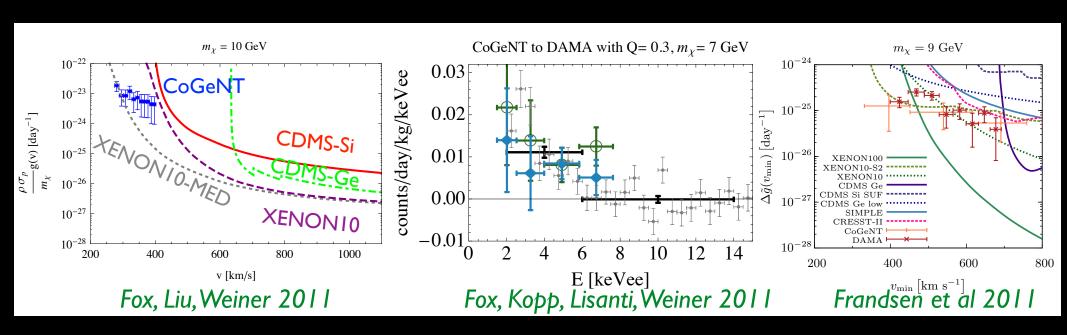


Minimum WIMP speed to impart recoil energy E_R

Find velocity integral from one experiment and use it for another.

Fox, Liu, Weiner 2011

$$\frac{dR}{dE_R} = \frac{A^2 F^2(E_R)}{2\mu_{\chi p}^2} \,\tilde{\eta}(v_{\min}) \qquad \qquad \qquad \qquad \tilde{\eta}(v_{\min}) = \frac{2\mu_{\chi p}^2}{A^2 F^2(E_R)} \,\frac{dR}{dE_R}$$

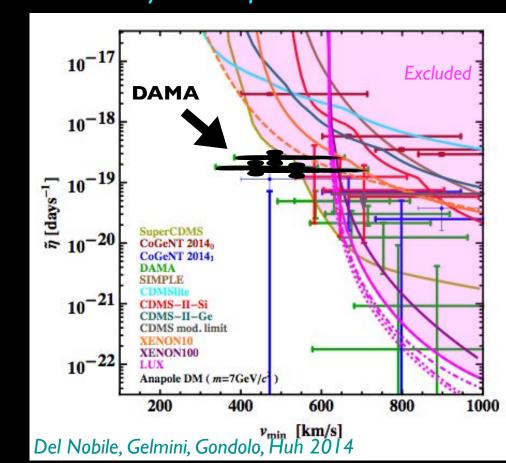


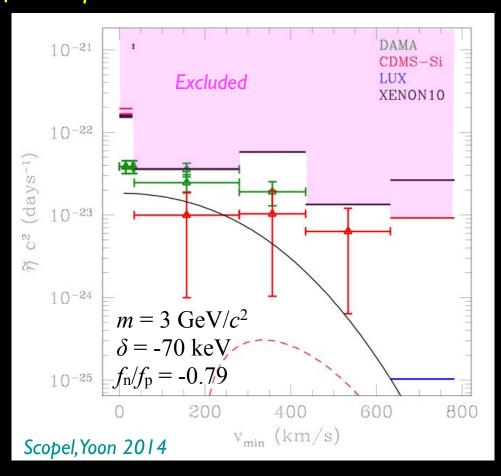
Needs unique relation between measured energy and minimum WIMP speed, available for **single-target detectors with excellent energy resolution**. For composite targets, lucky event pattern in CRESST allowed inversion.

In general, for composite targets and finite energy resolution, one can still find weighted averages of the velocity integral.

Gondolo, Gelmini 2012

DAMA may be compatible with null searches for anapole and exothermic dark matter.

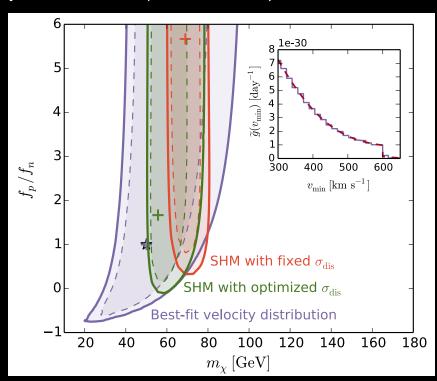




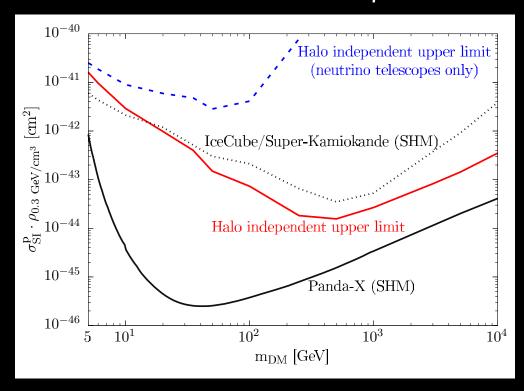
Allows for any velocity and energy dependent cross section, and indirect searches through neutrinos from the Sun/Earth.

Alternatively, one has sampled **discretized velocity distributions** to find bounds from direct and indirect experiments (neutrinos from the Sun).

Likelihood for particle-physics parameters (mock data)



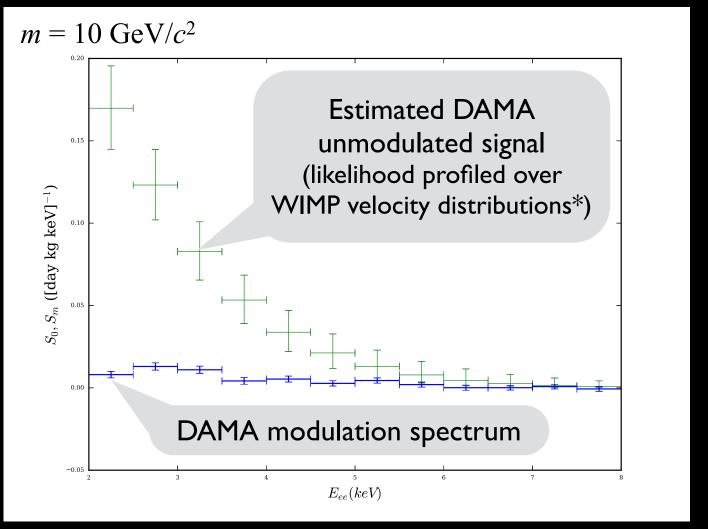
Bounds on cross section from direct detection and neutrinos from the Sun



New powerful techniques inspired by probability theory (problem of moments and Chebyshev inequalities) have just been introduced.

First application: halo-independent DAMA unmodulated signal

Gondolo, Scopel 2017



$S_{0,i}$	$S_{0,i} + B_i$
$0.17^{+0.026}_{-0.025}$	1.029
$0.12^{+0.022}_{-0.021}$	1.228
$0.083^{+0.01}_{-0.01}$	$\frac{18}{7}$ 1.294
$0.053^{+0.01}_{-0.01}$	$\frac{15}{4}$ 1.140
$0.034^{+0.01}_{-0.01}$	$\frac{13}{12}$ 0.956

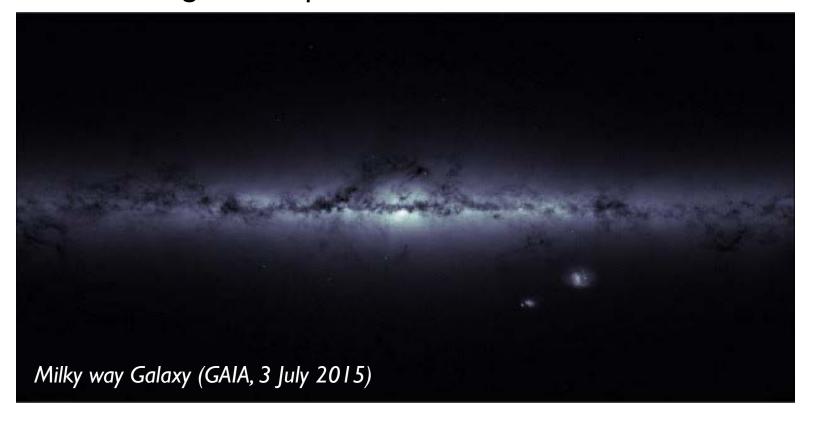
The unmodulated signal is compatible with the background

(*isotropic in galactic rest frame)

In the next future

In the next future..... Small-scale structure

GAIA is measuring the 3D position of about one billion stars.

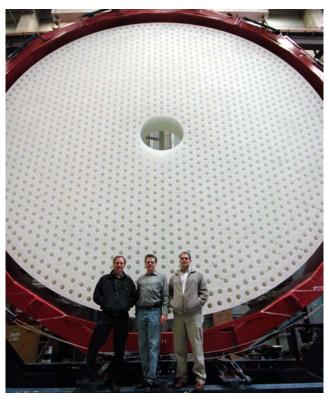


GAIA may detect dark matter substructures in the galactic halo. Feldmann, Spolyar 2014

In the next future..... Large-scale structure

The Dark Energy Survey (DES) and the Large Synoptic Survey Telescope (LSST) will map the large-scale distribution of dark matter and its growth with time.





DES

LSST

In the next future..... DAMA's revenge?

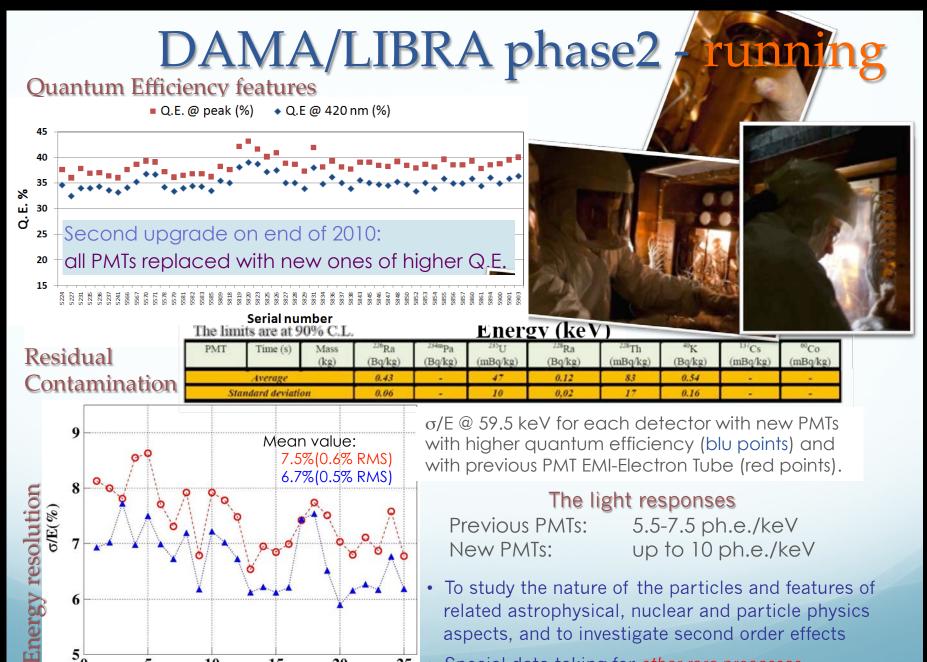
5

10

15

Detector number

20



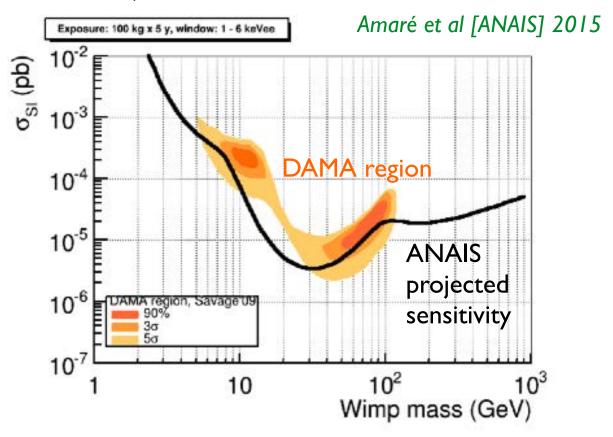
Special data taking for other rare processes

In the next future..... Direct check on DAMA

Experiments have been proposed that can directly check the DAMA modulation using the same target material

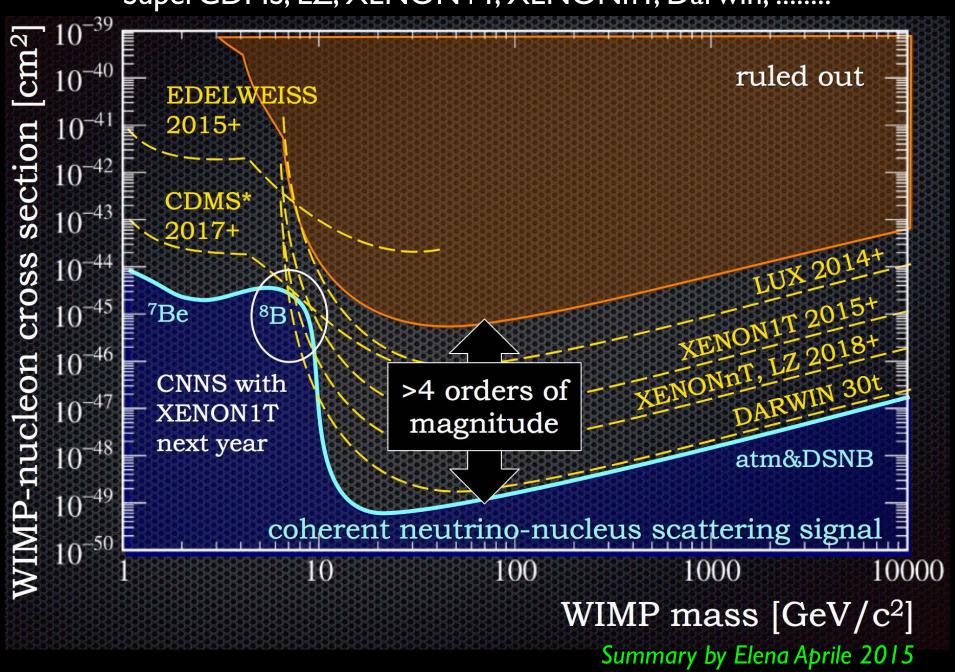
COSINE-100 (DM-ICE+KIMS-NaI)

ANAIS, SABRE, ...

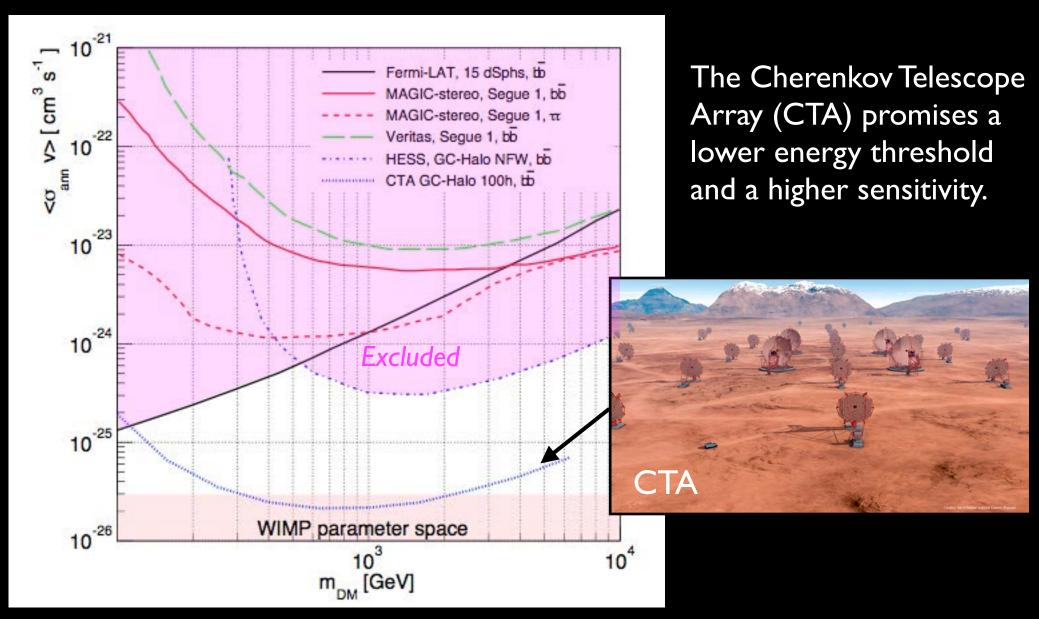


In the next future..... Giant direct detectors

SuperCDMS, LZ, XENON1T, XENONnT, Darwin,



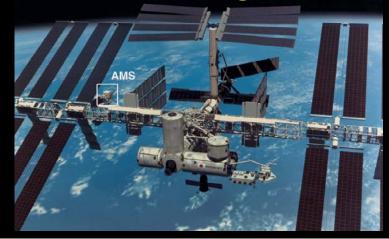
In the next future..... High-energy γ-rays

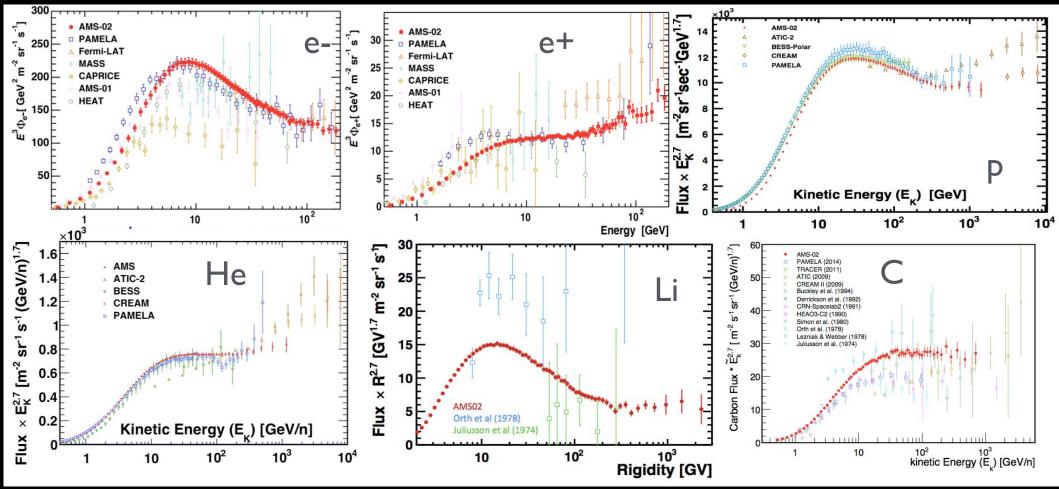


In the next future..... Precision cosmic rays

AMS (Alpha Magnetic Spectrometer)

Isotopic ratios measured to better than 1% precision up to Fe and ~100 GeV/nucleon allow for better Galactic cosmic ray models





Summary

- There is overwhelming astrophysical evidence for nonbaryonic cold dark matter
- The nature of cold dark matter is still unknown, and many candidates have been proposed
- There are some controversial claims of having detected signals from particle dark matter
- The next future will see measurements of the small- and large-scale structure of dark matter, as well as larger and more precise searches for dark matter signals