

Analysis of an instability in stratified fluid flow

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Buoyancy driven fluid flow, Bangalore

June 19 2017

Acknowledgements

The following is joint work with H Segur (U Colorado-Boulder) and B Deconinck (U Washington-Seattle). We build off the work of A Fokas (Cambridge), B Pelloni (U Reading), N Sheils (U Minnesota), P Treharne (U Sydney) and many others.

Boussinesq equations

- Assume a two-dimensional, viscous, incompressible fluid flow
- Density perturbations only relevant when multiplied by g (acceleration due to gravity)
- Fluid is bounded by two solid walls at $z = 0$ and $z = -h$ with no-slip boundary condition (Couette type)
- Mean density profile linear in z

$$\begin{aligned}\partial_t u + u\partial_x u + w\partial_z u &= -\frac{1}{\rho_0} \partial_x p + \nu(\partial_{xx} u + \partial_{zz} u), \\ \partial_t w + u\partial_x w + w\partial_z w &= -\frac{1}{\rho_0} \partial_z p - g\frac{\rho}{\rho_0} + \nu(\partial_{xx} w + \partial_{zz} w), \\ \partial_x u + \partial_z w &= 0, \\ \partial_t \rho + u\partial_x \rho + w\partial_z \rho &= wN + \kappa(\partial_{xx} \rho + \partial_{zz} \rho)\end{aligned}$$

Boussinesq equations

Exact solution

$$\begin{aligned}u &= U(z) = \frac{zU_0}{h} + U_0, \\w &= 0, \\ \rho &= 0\end{aligned}$$

$$\begin{aligned}\partial_t u + u\partial_x u + w\partial_z u &= -\frac{1}{\rho_0} \partial_x p + \nu(\partial_{xx} u + \partial_{zz} u), \\ \partial_t w + u\partial_x w + w\partial_z w &= -\frac{1}{\rho_0} \partial_z p - g\frac{\rho}{\rho_0} + \nu(\partial_{xx} w + \partial_{zz} w), \\ \partial_x u + \partial_z w &= 0, \\ \partial_t \rho + u\partial_x \rho + w\partial_z \rho &= wN + \kappa(\partial_{xx} \rho + \partial_{zz} \rho)\end{aligned}$$

Linear stability

Linearise about shear flow

$$u = U(z) + \tilde{u}, \quad w = \tilde{w}, \quad \rho = \tilde{\rho}$$

Introduce a stream function ψ : $\tilde{u} = -\partial_z \psi$, $\tilde{w} = \partial_x \psi$

Hat's denote Fourier transform with horizontal wave number k

$$\begin{aligned} -\partial_{tz} \hat{\psi} - ikU \partial_z \hat{\psi} + ik\hat{\psi} \frac{U_0}{h} &= -\frac{ik}{\rho_0} \hat{p} - \nu(-k^2 \partial_z \hat{\psi} + \partial_{zzz} \hat{\psi}), \\ ik\partial_t \hat{\psi} + (ik)^2 U \hat{\psi} &= -\frac{1}{\rho_0} \partial_z \hat{p} - g \frac{\hat{\rho}}{\rho_0} + ik\nu(-k^2 \hat{\psi} + \partial_{zz} \hat{\psi}), \\ \partial_t \hat{\rho} + ikU \hat{\rho} &= ik\hat{\psi} N + \kappa(-k^2 \hat{\rho} + \partial_{zz} \hat{\rho}) \end{aligned}$$

Linear stability

Orr-Sommerfeld equation

Combine momentum equations for single function $\hat{\psi}$ coupled with density equation

$$\hat{\psi} = \partial_z \hat{\psi} = 0, \quad z = 0, -h$$

$$\hat{\rho} = 0, \quad z = 0, -h$$

$$U = \frac{zU_0}{h} + U_0$$

$$\partial_t(-\partial_{zz} + k^2)\hat{\psi} + ikU(-\partial_{zz} + k^2)\hat{\psi} = gik \frac{\hat{\rho}}{\rho_0} - \nu(-\partial_{zz} + k^2)^2 \hat{\psi}$$

$$\partial_t \hat{\rho} + ikU \hat{\rho} = ik\hat{\psi}N + \kappa(-k^2 \hat{\rho} + \partial_{zz} \hat{\rho})$$

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Proceed with eigenvalue analysis: $\hat{\psi} = e^{\lambda t} f(z)$ $\hat{\rho} = e^{\lambda t} g(z)$

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Non-normal operator \Rightarrow cannot preclude short time growth even when all eigenvalues are stable

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An alternative viewpoint

Caution: new ideas ahead!

- How much progress can be made without resorting to numerics?
- Step one: reformulate problem (H Segur)
- Step two: simplify the problem
- Step three: analyse the problem

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Disclaimer!

Reformulation

recall

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Reformulation

$\hat{w} - \hat{p} - \hat{\rho}$ equations

The two momentum equations are equivalent iff

$$(-\partial_{zz} + k^2)\hat{p} = 2\rho_0(ik)^2\hat{\psi}\frac{U_0}{h} - g\partial_z\rho$$

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To keep things clean, re-introduce $\hat{w} = ik\hat{\psi}$

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$$U = \frac{zU_0}{h} + U_0$$

Boundary conditions

$$\hat{w} = 0, \partial_z \hat{w} = 0, \quad z = 0, -h$$

$$\hat{\rho} = 0, \quad z = 0, -h$$

$$(-\partial_{zz} + k^2)\hat{p} = 2\rho_0 ik \frac{U_0}{h} \hat{w} - g\partial_z \rho$$

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No boundary conditions for \hat{p} !

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General solution to second order ODE

$$\hat{p} = \rho_0 \int_{-h}^0 G(z-y) \left(2ik \frac{U_0}{h} \hat{w} - g \frac{\partial_z \rho}{\rho_0} \right) dy + A(t) \rho_0 e^{-|k|z} + B(t) \rho_0 e^{|k|z}$$

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No boundary conditions to apply \Rightarrow retain A and B as unknowns.

$$(-\partial_{zz} + k^2)\hat{p} = 2\rho_0 ik \frac{U_0}{h} \hat{w} - g \partial_z \rho$$

$$\partial_t \hat{w} + ikU \hat{w} = -\frac{1}{\rho_0} \partial_z \hat{p} - g \frac{\hat{\rho}}{\rho_0} + \nu(-k^2 \hat{w} + \partial_{zz} \hat{w}),$$

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$$\partial_t \begin{pmatrix} \hat{w} \\ \hat{\rho} \end{pmatrix} = \begin{pmatrix} -ikU & 0 \\ 0 & -ikU \end{pmatrix} \begin{pmatrix} \hat{w} \\ \hat{\rho} \end{pmatrix} + \mathcal{L} \begin{pmatrix} \hat{w} \\ \hat{\rho} \end{pmatrix}$$

A toy problem

Consider

$$\partial_t q = i f(x) q + L q,$$

where L is some operator (e.g. $L = \partial_{xx}$) posed on finite interval $(0, 1)$ with homogeneous BC's

- If eigenvalues of L are known, then perturbation theory can be used to find eigenvalues for $i f(x) + L$
- Perturbation theory usually requires $f(x)$ be small
- Energy integral

$$\partial_t \int \frac{|q|^2}{2} dx = \text{Re} \int q^* L q dx$$

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- Energy integral

$$\partial_t \int \frac{|q|^2}{2} dx = \operatorname{Re} \int q^* L q dx$$

$$\|q\| = \|q_0\| \exp \left(\int_0^t \operatorname{Re} \frac{\int q^* L q dx}{\|q\|^2} \right)$$

where

$$\|q\|^2 = \int_0^1 |q|^2 dx$$

A toy problem

$$\begin{aligned}\partial_t q &= if(x)q + Lq, \\ \Rightarrow \|q\| &= \|q_0\| \exp\left(\int_0^t \operatorname{Re} \frac{\int q^* Lq dx}{\|q\|^2} dt\right)\end{aligned}$$

- $\operatorname{Num}(L) = \{z \in \mathbb{C} \mid z = \int \phi^* L\phi dx \ \forall \|\phi\| = 1\}$
- If $L\phi + if(x)\phi = \lambda\phi$ then $\lambda \in \operatorname{Num}(L)$
- $\operatorname{Num}(P + Q) \subseteq \operatorname{Num}(P) + \operatorname{Num}(Q)$
- If $\operatorname{Re} \operatorname{Num}(L) < 0$ then q is asymptotically stable \Rightarrow all solutions decay to zero
- $\|q\| \leq \|q_0\| e^{Mt}$, $M = \sup\{\operatorname{Re} \operatorname{Num}(L)\}$
- If all eigenvalues of $L + if(x)$ are stable, then M gives the short time growth rate and we can estimate greatest possible algebraic growth rate
- Long time behavior is dictated by eigenvalues

A toy problem

Moral of the story...

Consider

$$\partial_t q = if(x)q + Lq,$$

where L is some operator (e.g. $L = \partial_{xx}$) posed on finite interval $(0, 1)$ with homogeneous BC's

$$\|q\| = \|q_0\| \exp\left(\int_0^t \operatorname{Re} \frac{\int q^* Lq dx}{\|q\|^2}\right)$$

•

$$\partial_t \begin{pmatrix} \hat{w} \\ \hat{\rho} \end{pmatrix} = \begin{pmatrix} -ikU & 0 \\ 0 & -ikU \end{pmatrix} \begin{pmatrix} \hat{w} \\ \hat{\rho} \end{pmatrix} + \mathcal{L} \begin{pmatrix} \hat{w} \\ \hat{\rho} \end{pmatrix}$$

is the system equivalent of the toy problem and all the claims follow.

The effective boundary-value problem

$$\hat{w} = \partial_z \hat{w} = \hat{\rho} = 0, \quad z = 0, -h$$

$$\begin{aligned} \partial_t \hat{w} = & -\partial_z \int_{-h}^0 G(z-y) \left(2ik \frac{U_0}{h} \hat{w} - g \frac{\partial_z \rho}{\rho_0} \right) dy - A|k|e^{-|k|z} + B|k|e^{|k|z} \\ & - g \frac{\hat{\rho}}{\rho_0} + \nu(-k^2 \hat{w} + \partial_{zz} \hat{w}), \end{aligned}$$

$$\partial_t \hat{\rho} = \hat{w}N + \kappa(-k^2 \hat{\rho} + \partial_{zz} \hat{\rho})$$

- We now seek as much as information about the BVP as possible.

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- We now seek as much as information about the BVP as possible.
Claim: this problem can be solved explicitly!
- New method to solve boundary value problems: Unified Transform Method (UTM)
- Encompasses the standard methods.
- It is more general than the standard methods: UTM handles non self-adjoint operators (*etc.*)
- Readily obtain eigenvalues, estimate asymptotic behavior, compute numerical range.
- Efficient evaluation of the solution: contour parametrization leading to stable numerical methods, steepest descent, the residue theorem, *etc.*

Unified Transform Method

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- Developed by Fokas and collaborators (see SIAM Review Deconinck, Trogdon & Vasan 2014)
- Recently extended to problems with nonlocal operators (Vasan 2016), systems of equations (Deconinck et al 2017), interface problems (Sheils, Pelloni and others)
- Solution given as

$$\int_{\Gamma_k} \frac{e^{ikx - \omega_j t}}{\Delta(k)} \hat{q}_0(k) dk + \int_{\Gamma_\nu} \frac{e^{ikx - \omega_j t}}{\Delta(k)} \hat{q}_0(\nu) dk$$

with $\Gamma_{\nu, k}$ as contours in complex plane. ω_j is dispersion relation

- $\Delta(k_m) = 0$ give rise to eigenfunctions and $\omega_j(k_m)$ is the eigenvalue

Concluding remarks

- Fluid shear problems in rectangular domains give rise to

$$\partial_t \bar{s} = -ikU\bar{s} + \mathcal{K}_U \bar{s} + \mathcal{L}\bar{s}$$

\bar{s} are the fluid variables of interest, \mathcal{K}_U is a compact operator and \mathcal{L} is an operator amenable to UTM.

- Using perturbation theory and numerical range we can estimate growth rates in terms of \mathcal{L} and shear
- Currently we're studying interfacial waves with density and viscosity jump with a background time-dependent shear profile

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Thank you!