Understanding Spin-1 Kagomé Antiferromagnet Through Hida Model

Brijesh Kumar (JNU, New Delhi)

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[work done with **Pratyay Ghosh**]

PRB 97, 014413 (2018)

PRB **93**, 014427 (2016)

05 December 2018 @ *ICTS* Bengaluru (2nd Asia-Pacific Workshop on Quantum Magnetism)

Acknowledgement

Funding:

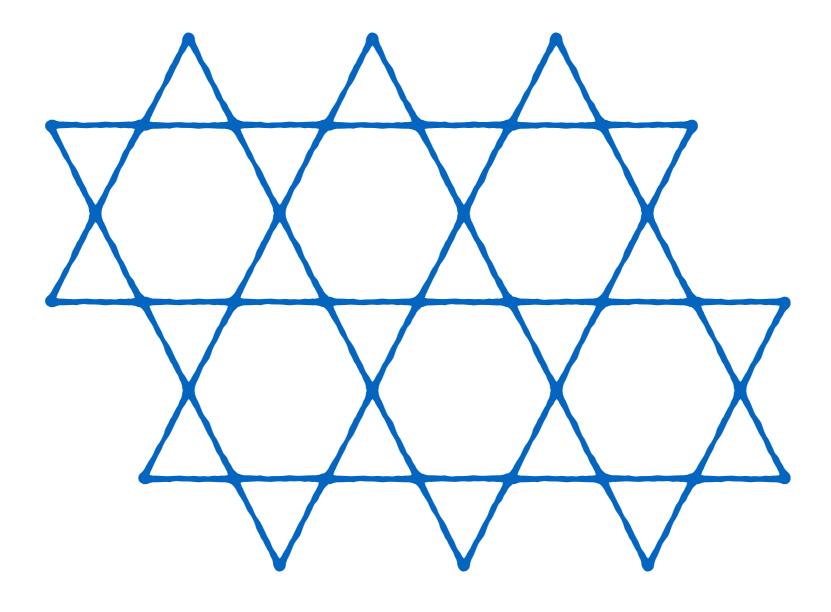
UGC (through UPE-II project)

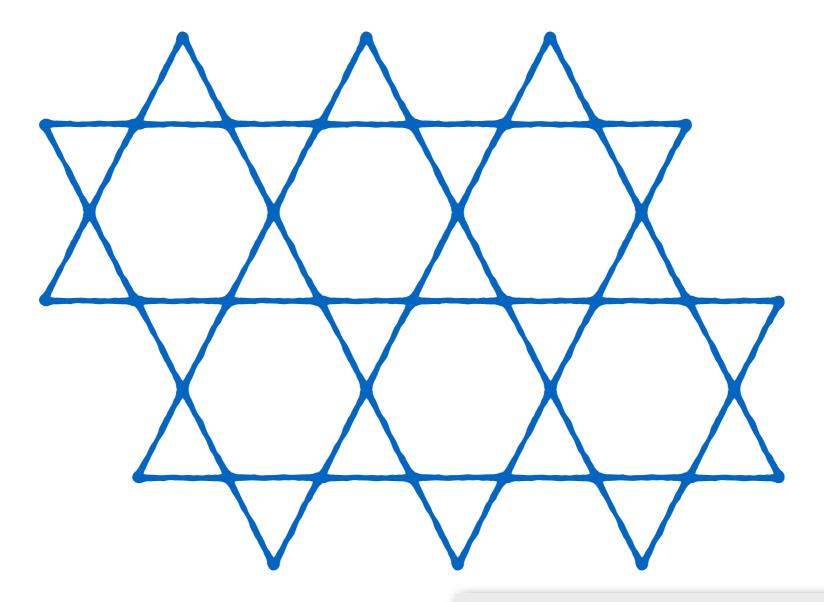
DST (through PURSE and FIST programs)

CSIR (for students' fellowships)

Computing facility:

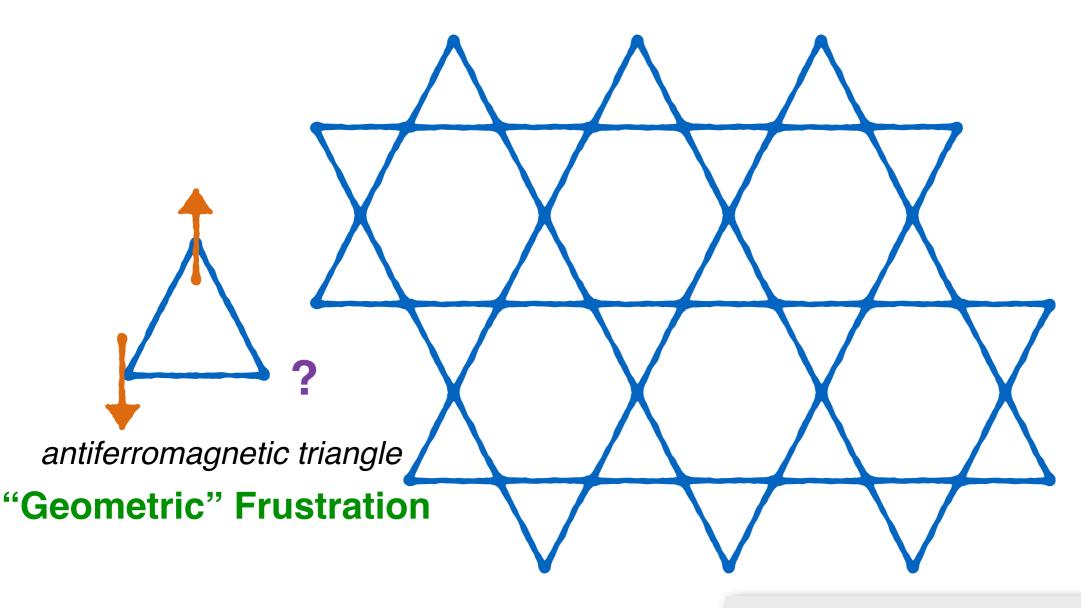
IUAC (New Delhi), and SPS (JNU)





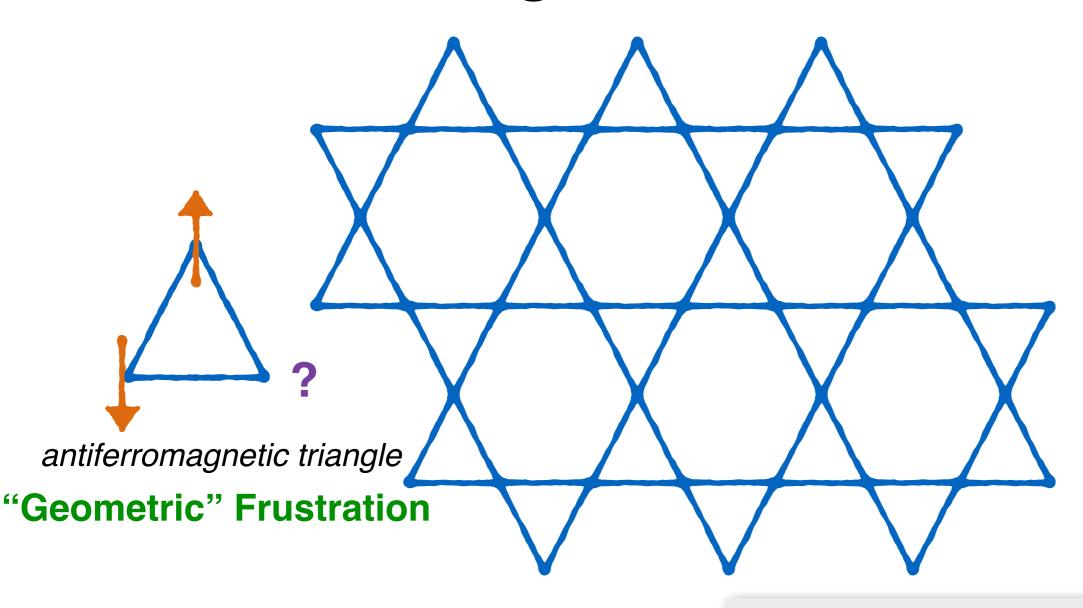
Heisenberg Antiferromagnet

$$H = \sum_{\langle i,j \rangle} J_{ij} \vec{S}_i \cdot \vec{S}_j$$



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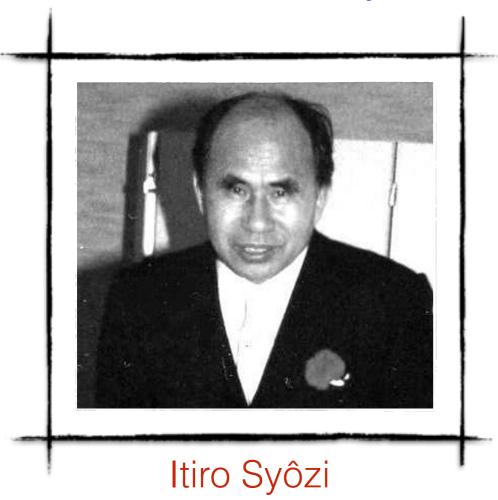


Heisenberg Antiferromagnet

"Highly" Frustrated Spin Model

$$H = \sum_{\langle i,j \rangle} J_{ij} \vec{S}_i \cdot \vec{S}_j$$

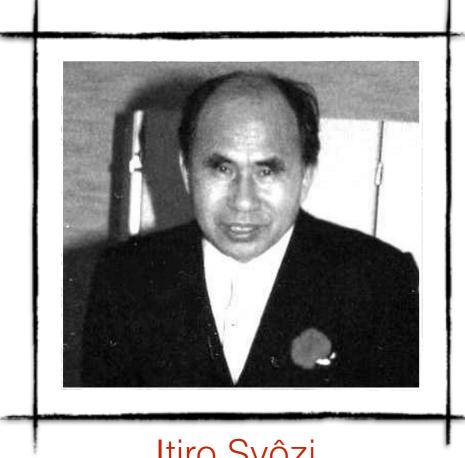
Appears first, through star-triangle relation, in the study of the Ising model on a "decorated" honeycomb lattice



Prog. Theor. Phys. 6, 306 (1951)

"... But in the antiferromagnetic case, there is no phase change."

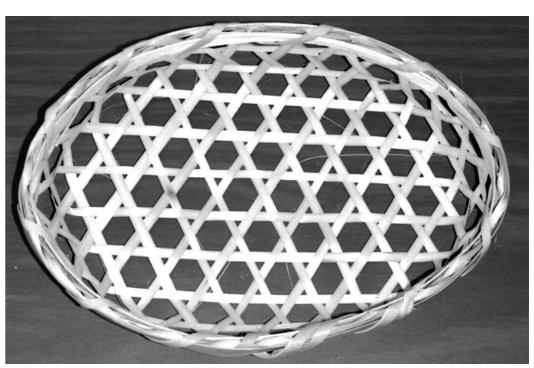
Appears first, through star-triangle relation, in the study of the Ising model on a "decorated" honeycomb lattice



Itiro Syôzi

Prog. Theor. Phys. 6, 306 (1951)

so named by Kodi Husimi...



Kagome = basket with holes

Kagome: The story of the basketweave lattice M. Mekata, Physics Today 56, 12 (2003)

> "... But in the antiferromagnetic case, there is no phase change."

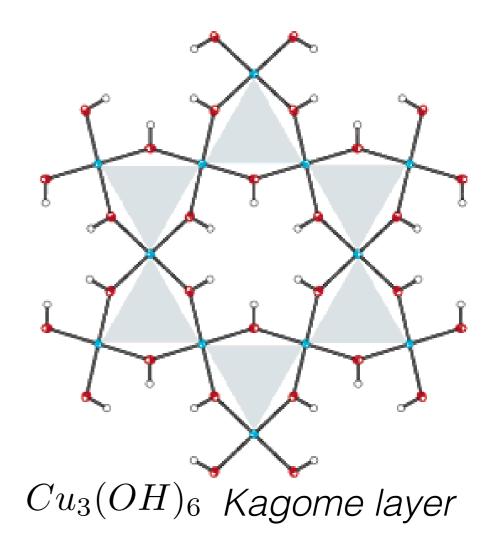
Kagome Antiferromagnets

Herbertsmithite

$$ZnCu_3(OH)_6Cl_2$$

nearly perfect quantum spin-1/2 Kagome Heisenberg Antiferromagnet

Shores et al, JACS (2005)



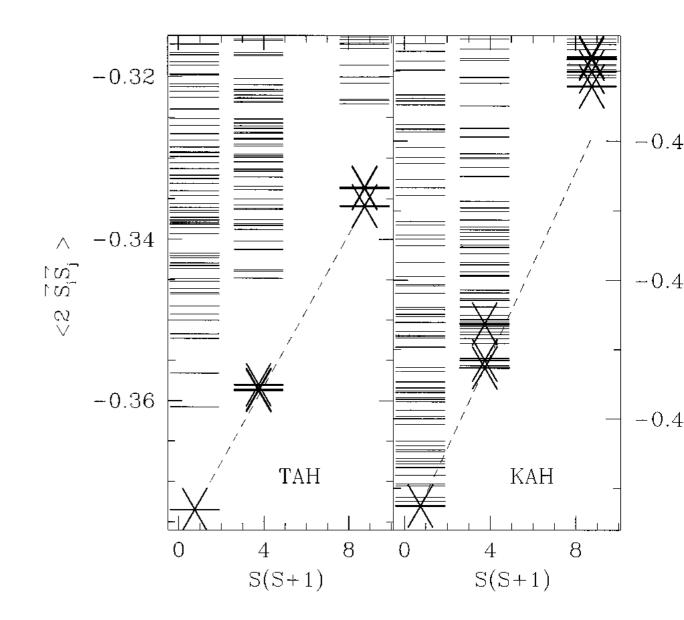
Absence of magnetic order in the ground state...

(KHA = Kagome Heisenberg Antiferromagnet)

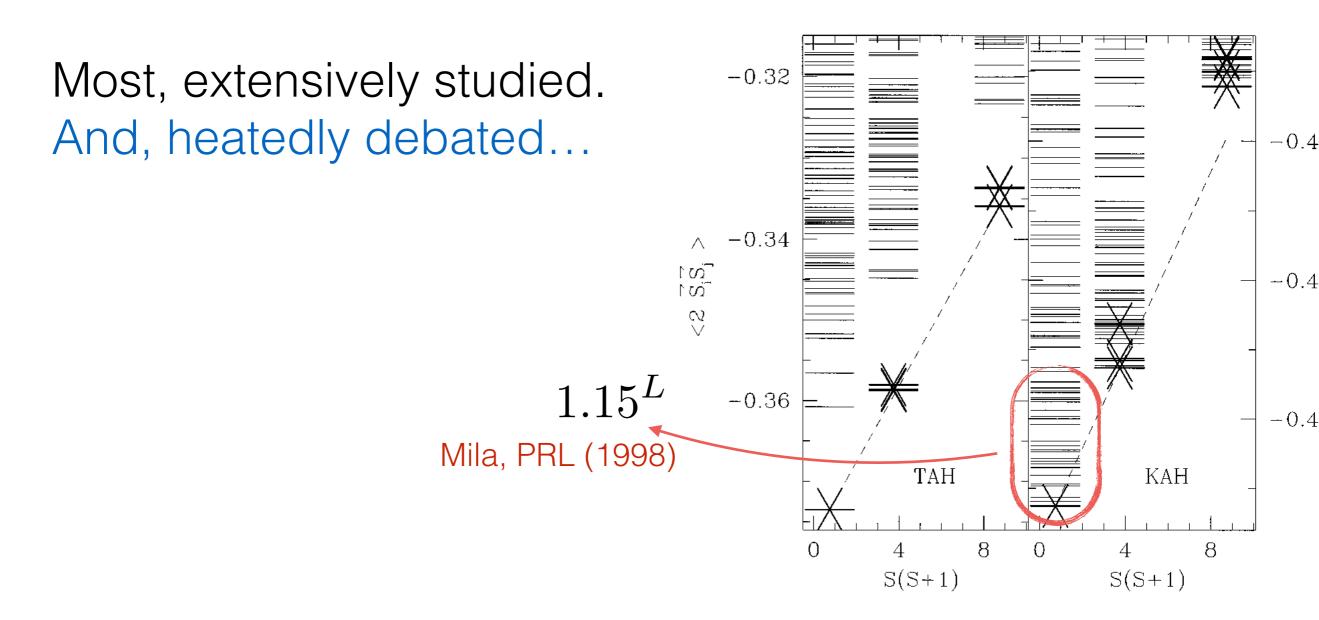
Most, extensively studied.

And, heatedly debated...

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Lecheminant el al, PRB (1997)



Lecheminant el al, PRB (1997)

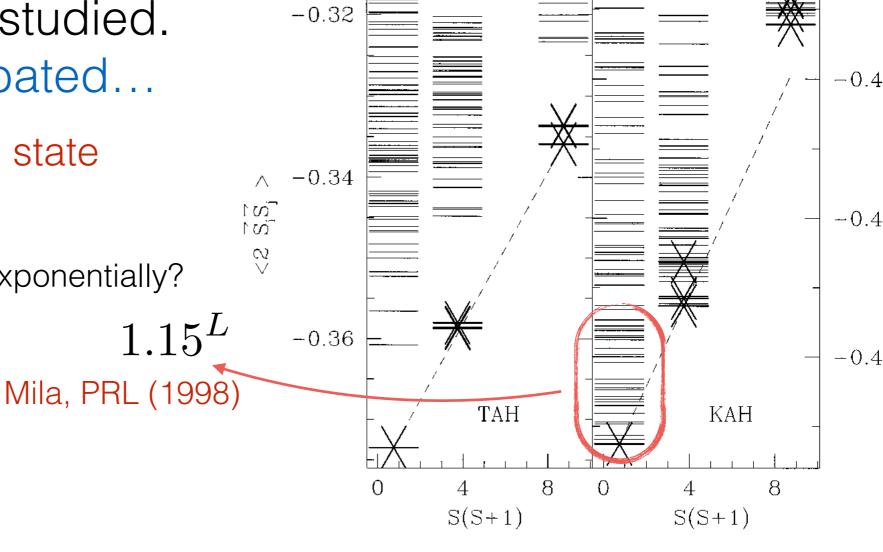
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Quantum spin liquid state

$$\langle \mathbf{S}_i \rangle = 0$$

do all correlations decay exponentially?



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Quantum spin liquid state

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 1.15^L Mila, PRL (1998)

-0.32-0.4-0.34-0.4-0.36 TAH KAH 8 S(S+1)S(S+1)

Gapped RVB? S. Yan et al (2011)

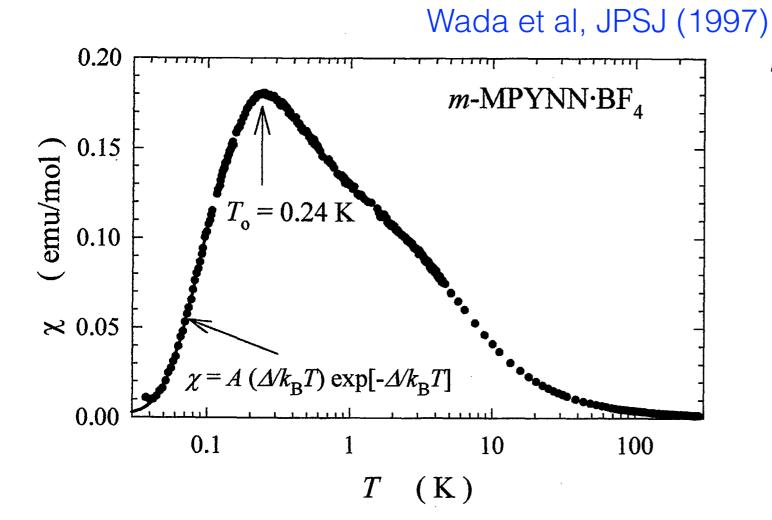
Gapless Algebraic? Y. Iqbal el al (2013)

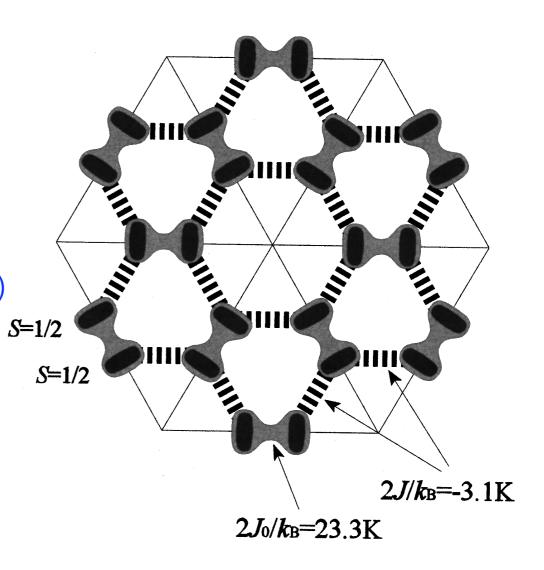
Lecheminant el al, PRB (1997)

Valence bond crystal? A. Ralko el al (2018)

 $m - MPYNN.BF_4$

(m-N-methylperidium-a-nitronyl-nitroxide)



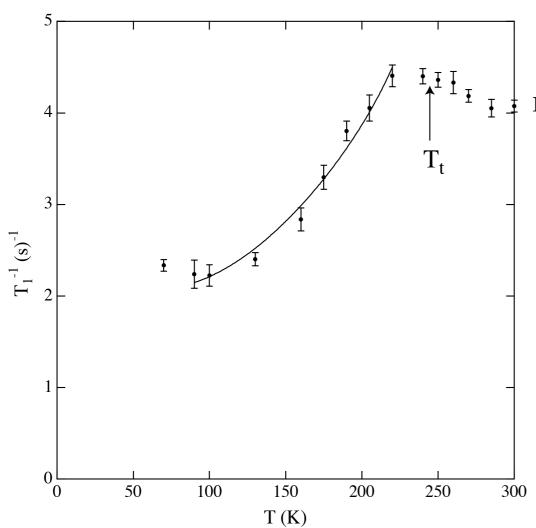


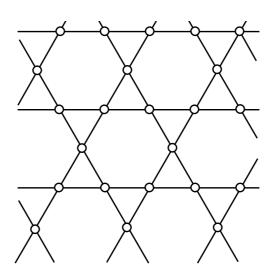
Watanabe et al, PRB (1998)

exhibit spin gap behaviour

 NaV_6O_{11}

Kato et al, JPSJ (2001)





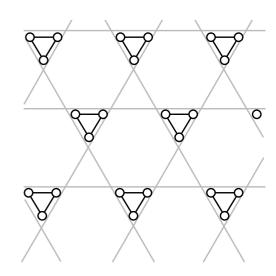
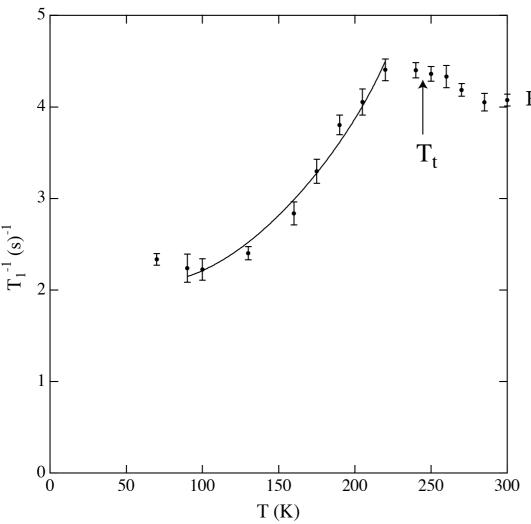


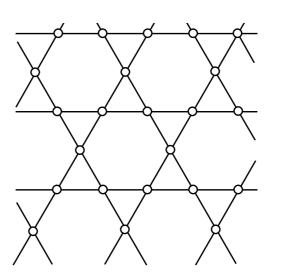
Fig. 1. (a) Crystal structure of NaV₆O₁₁ at room temperature. Dashed lines express a unit cell. (b) V(1) kagomé lattice at $T > T_t$. (c) Distorted V(1) kagome lattice at $T < T_t$.

exhibit spin gap behaviour

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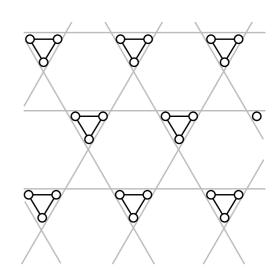


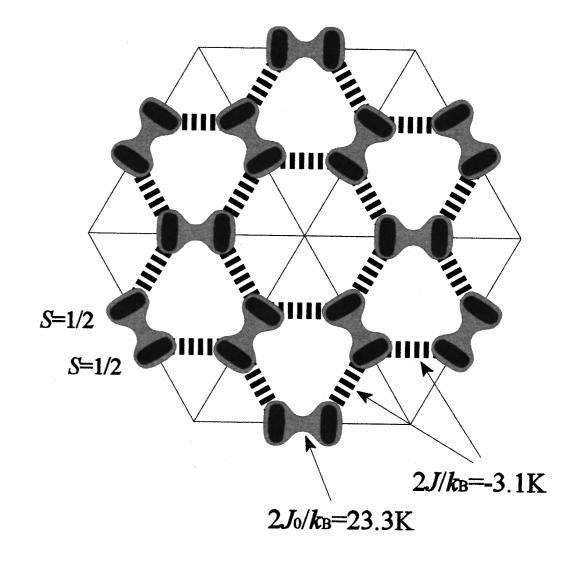
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Another spin-1 Kagome compound

 $KV_3Ge_2O_9$

Hara et al, JPSJ (2012)

[Kazuo Hida, JPSJ (2000)]

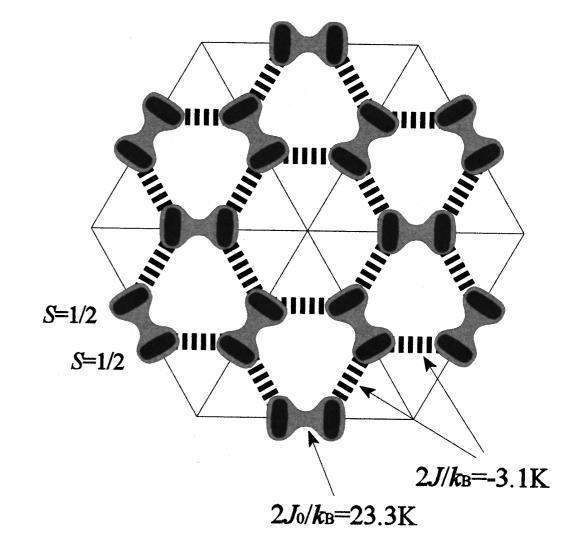


 $m - MPYNN.BF_4$

[Kazuo Hida, JPSJ (2000)]

Minimal Model

(quantum spin-1/2 on honeycomb lattice)

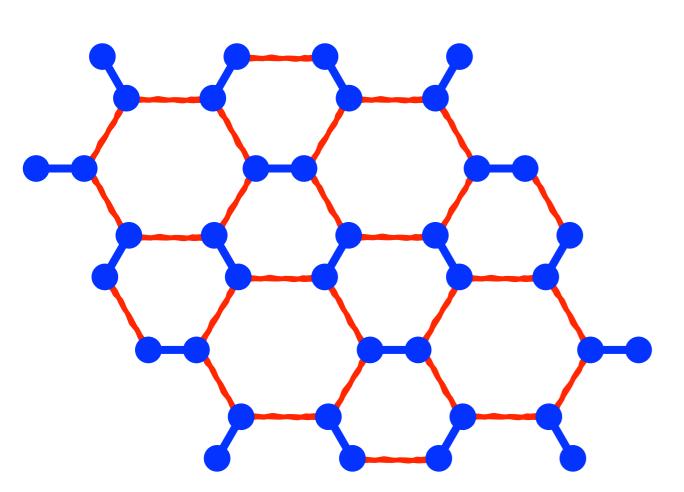


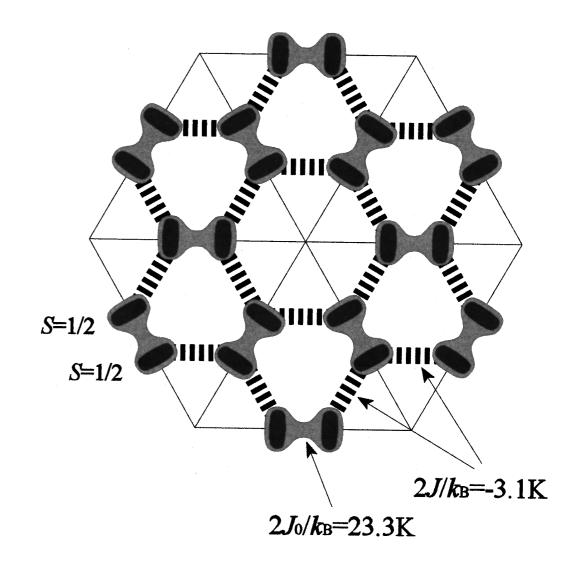
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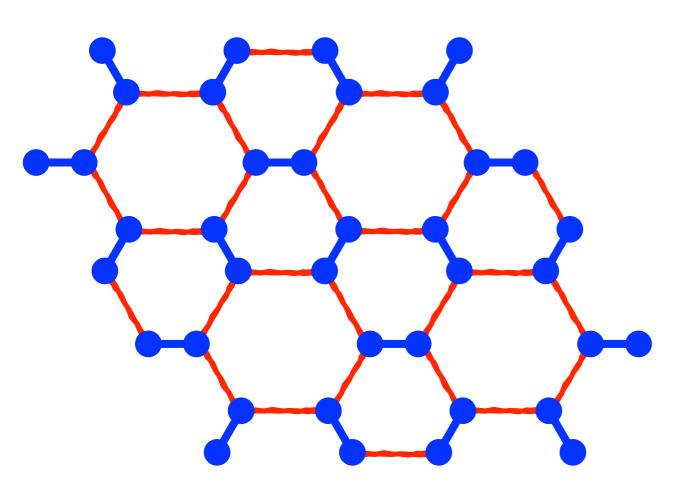
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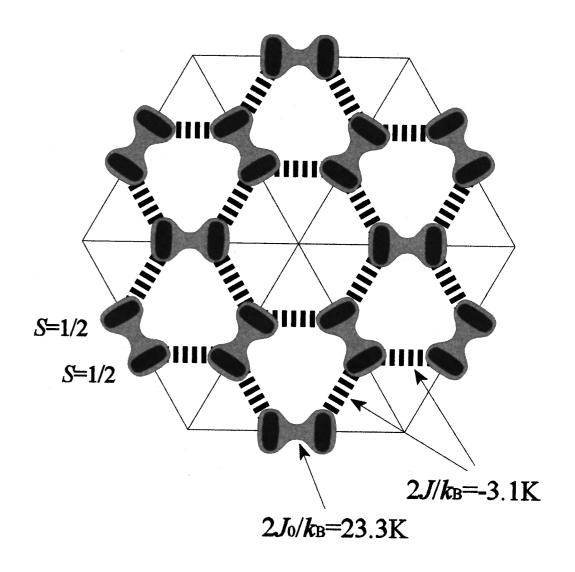
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antiferromagnetic hexagons (J_A) coupled ferromagnetically (J_F)





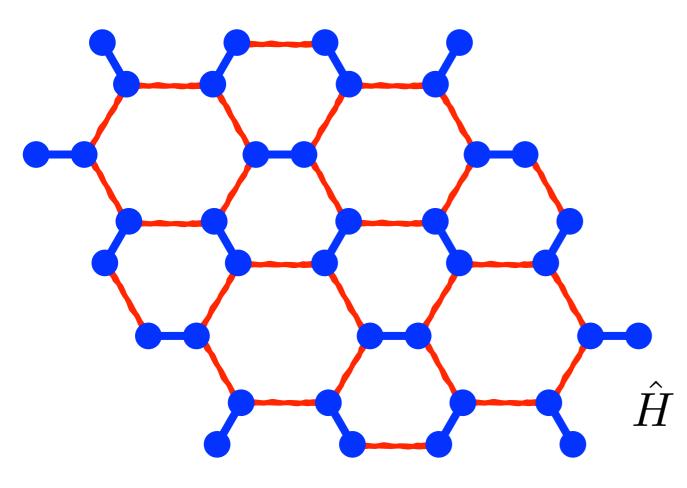
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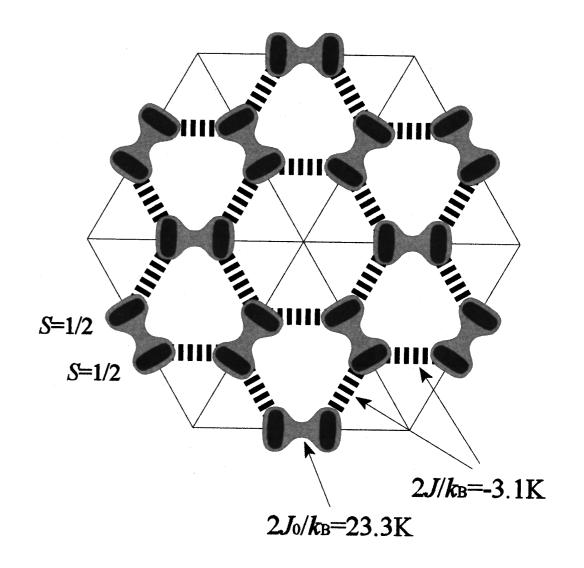
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$$\hat{H} = J_F \sum_{\langle i,j \rangle} \mathbf{S}_i \cdot \mathbf{S}_j + J_A \sum_{\langle i,j \rangle} \mathbf{S}_i \cdot \mathbf{S}_j$$

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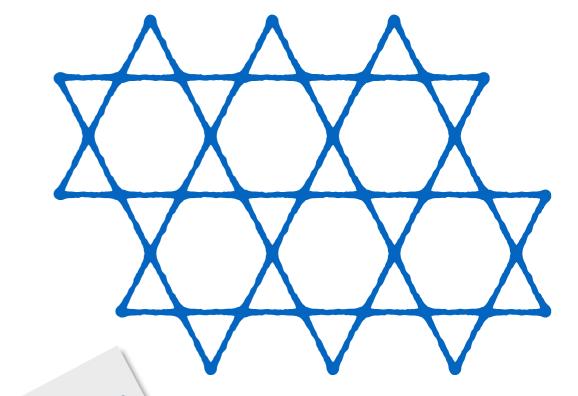
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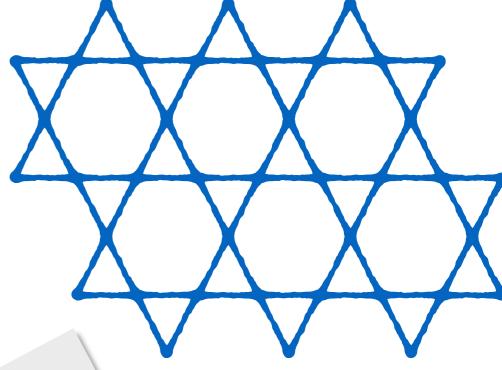
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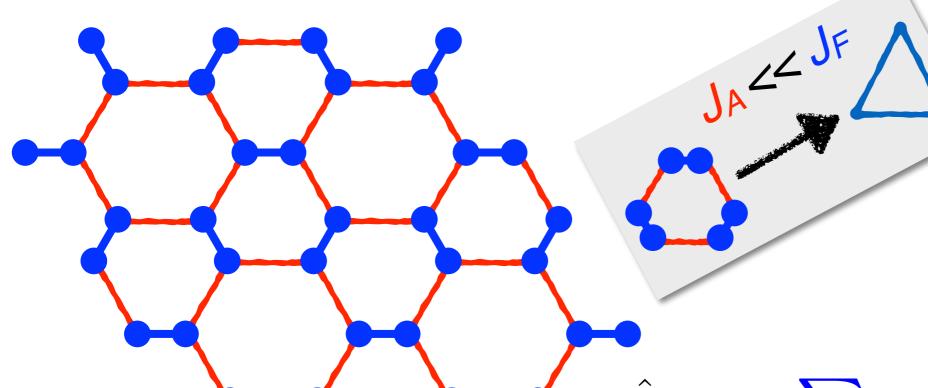
(spin-1 KHA)

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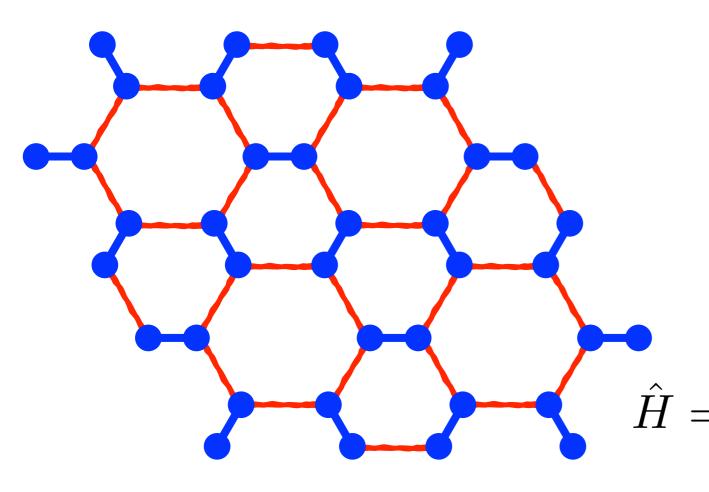
"Hida Model"

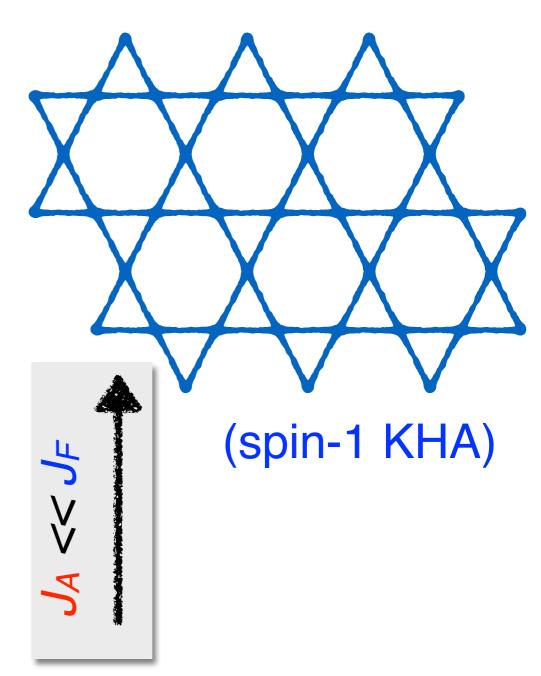


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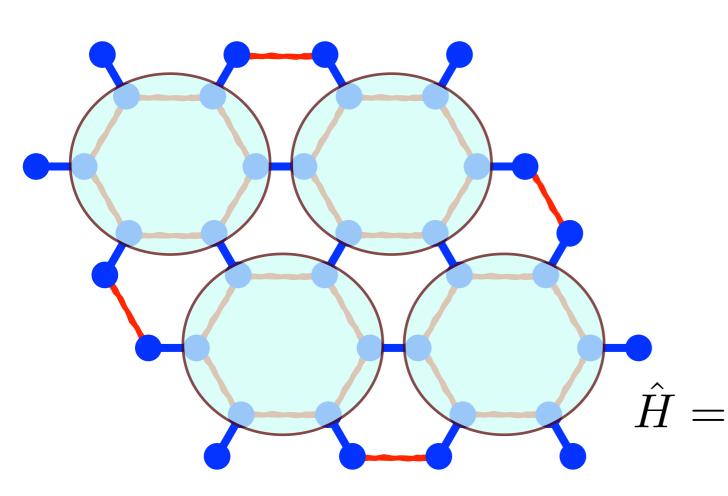


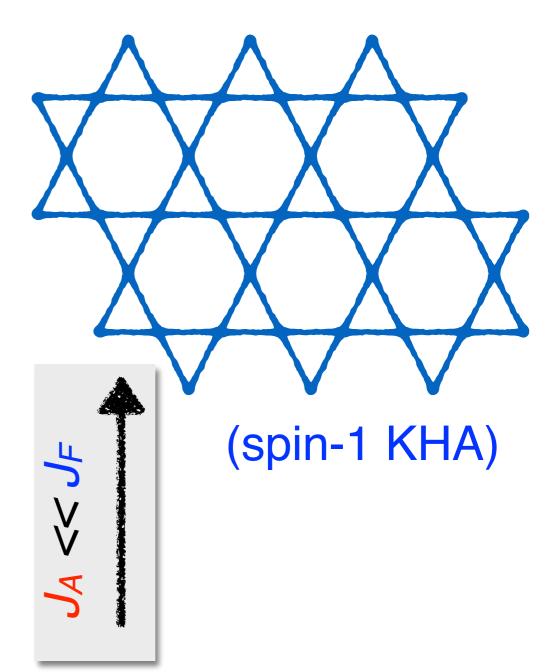
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[Kazuo Hida, JPSJ (2000)]

Hida's ground state for spin-1 KHA

hexagonal singlets



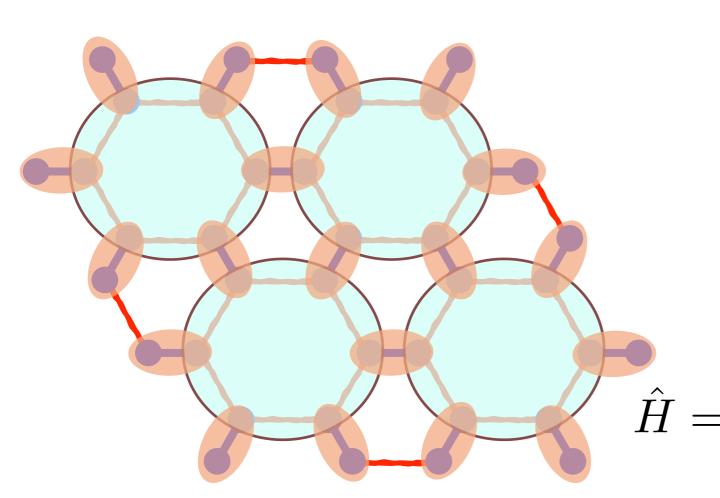


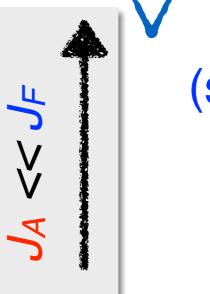
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[Kazuo Hida, JPSJ (2000)]

Hida's ground state for spin-1 KHA

hexagonal singlets with symmetrized FM bonds





(spin-1 KHA)

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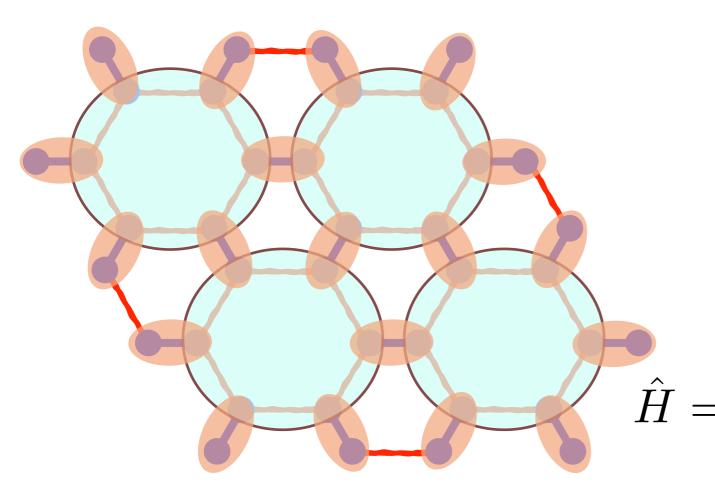
[Kazuo Hida, JPSJ (2000)]

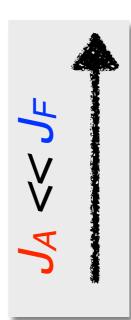
Hida's ground state for spin-1 KHA

hexagonal singlets with symmetrized FM bonds

HSS state

(AKLT type)





(spin-1 KHA)

$$\hat{H} = J_F \sum_{\langle i,j \rangle} \mathbf{S}_i \cdot \mathbf{S}_j + J_A \sum_{\langle i,j \rangle} \mathbf{S}_i \cdot \mathbf{S}_j$$

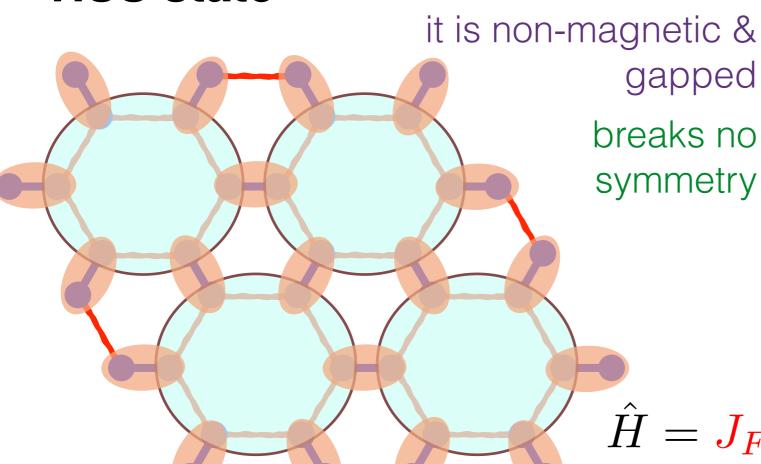
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Hida's ground state for spin-1 KHA

hexagonal singlets with symmetrized FM bonds

HSS state

(AKLT type)



√A << √F

(spin-1 KHA)

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[Kazuo Hida, JPSJ (2000)]

Hida's ground state for spin-1 KHA

hexagonal singlets with symmetrized FM bonds

HSS state

(AKLT type)

it is non-magnetic & gapped

breaks no symmetry

(spin-1 KHA)

However, the question of the ground state of spin-1 KHA has been revisited recently.

[Kazuo Hida, JPSJ (2000)]

Hida's ground state for spin-1 KHA

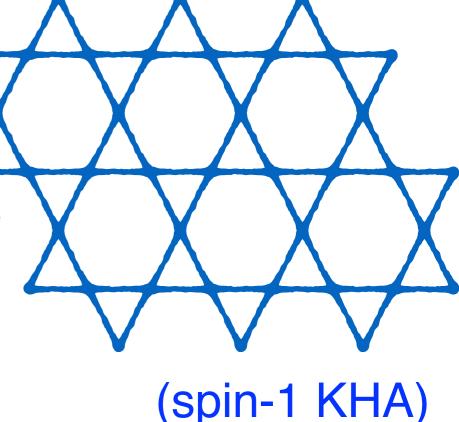
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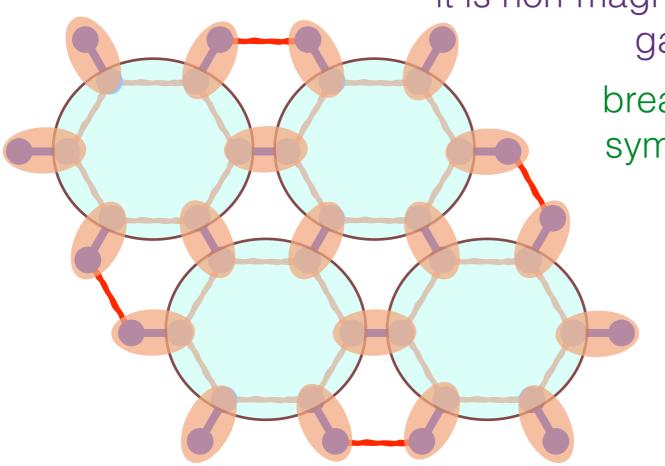
it is non-magnetic & gapped

> breaks no symmetry



However, the question of the ground state of spin-1 KHA has been revisited recently.

and, energetically better candidate for the ground state has been found.



Spontaneously Trimerised Singlet (TS) state

Liu et al, PRB-Rapid (2015)

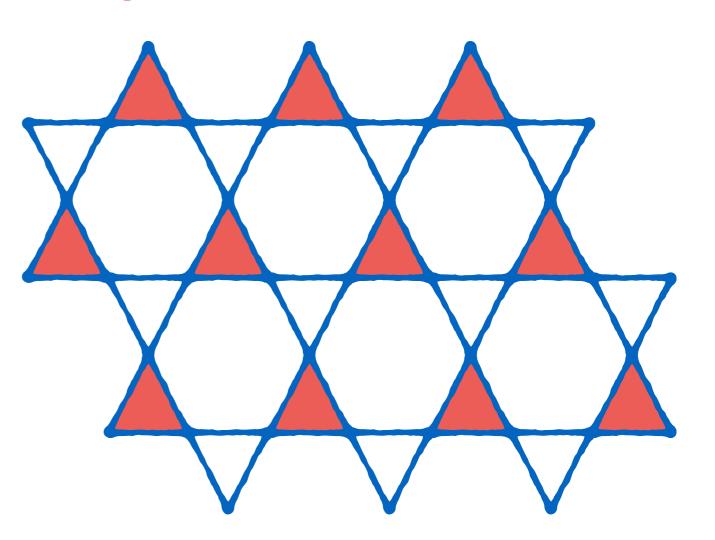
Picot and Poilblanc, PRL (2015)

Changlani and Läuchli, PRB (2015)

Ghosh, Verma and Kumar, PRB (2016)

Rajiv Singh et al (2016)

and others...



[Z. Cai et al, J. Phys. Condens. Matter. (2009)]

A Model on Trimerised Kagome Lattice

$$\hat{H} = J \sum_{\Delta} \vec{S}_i \cdot \vec{S}_j$$

$$+J' \sum_{\nabla} \vec{S}_i \cdot \vec{S}_j$$

A Model on Trimerised Kagome Lattice

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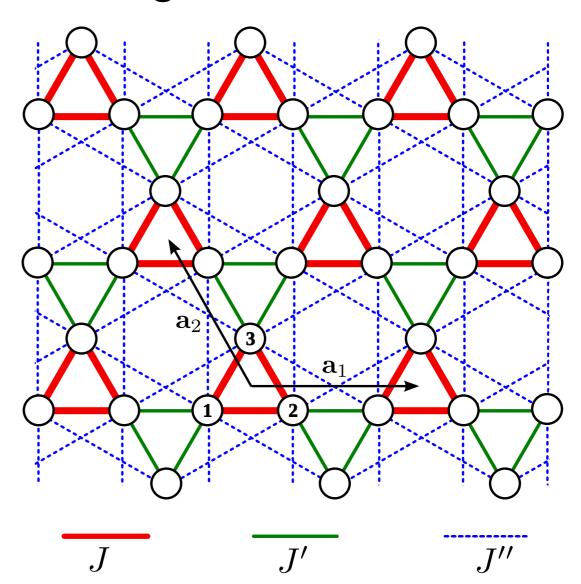
$$+J'' \sum_{\langle \langle i,j \rangle \rangle} \vec{S}_i \cdot \vec{S}_j$$

A Model on Trimerised Kagome Lattice

$$\hat{H} = J \sum_{\Delta} \vec{S}_i \cdot \vec{S}_j$$

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$$+J'' \sum_{\langle \langle i,j \rangle \rangle} \vec{S}_i \cdot \vec{S}_j$$



Plaquette Triplon Analysis

eigenstates of a spin-1 triangle

$$\hat{H}_{\triangle} = J(\vec{S}_1 \cdot \vec{S}_2 + \vec{S}_2 \cdot \vec{S}_3 + \vec{S}_3 \cdot \vec{S}_1)$$



9 triplet states -2J

plaquette operators

$$|s\rangle := \hat{s}^{\dagger} |\phi\rangle,$$

 $|t_{mv}\rangle := \hat{t}_{mv}^{\dagger} |\phi\rangle.$

$$\hat{s}^{\dagger}\hat{s} + \sum_{m,\nu} \hat{t}_{m\nu}^{\dagger}\hat{t}_{m\nu} = 1.$$
 (constraint)

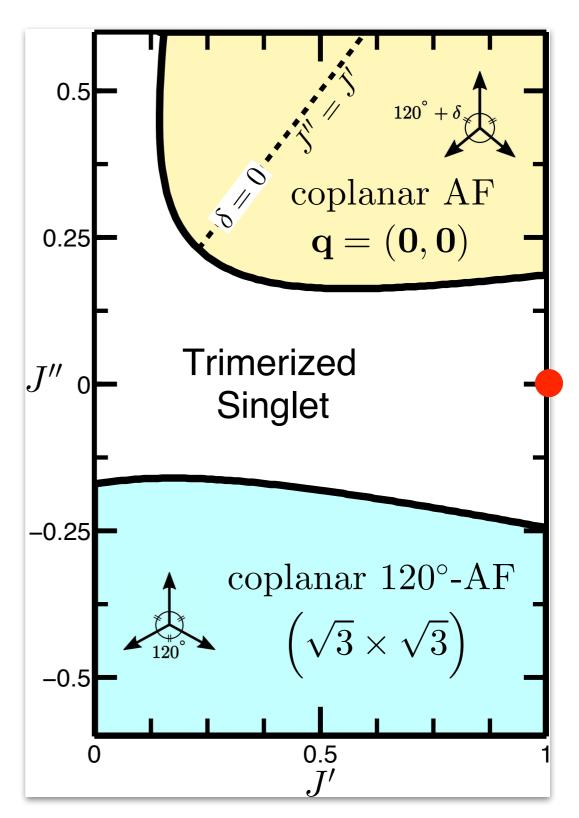
akin to the bond-operators for dimerised AF problems [Sachdev & Bhatt (1990); BK (2010)]

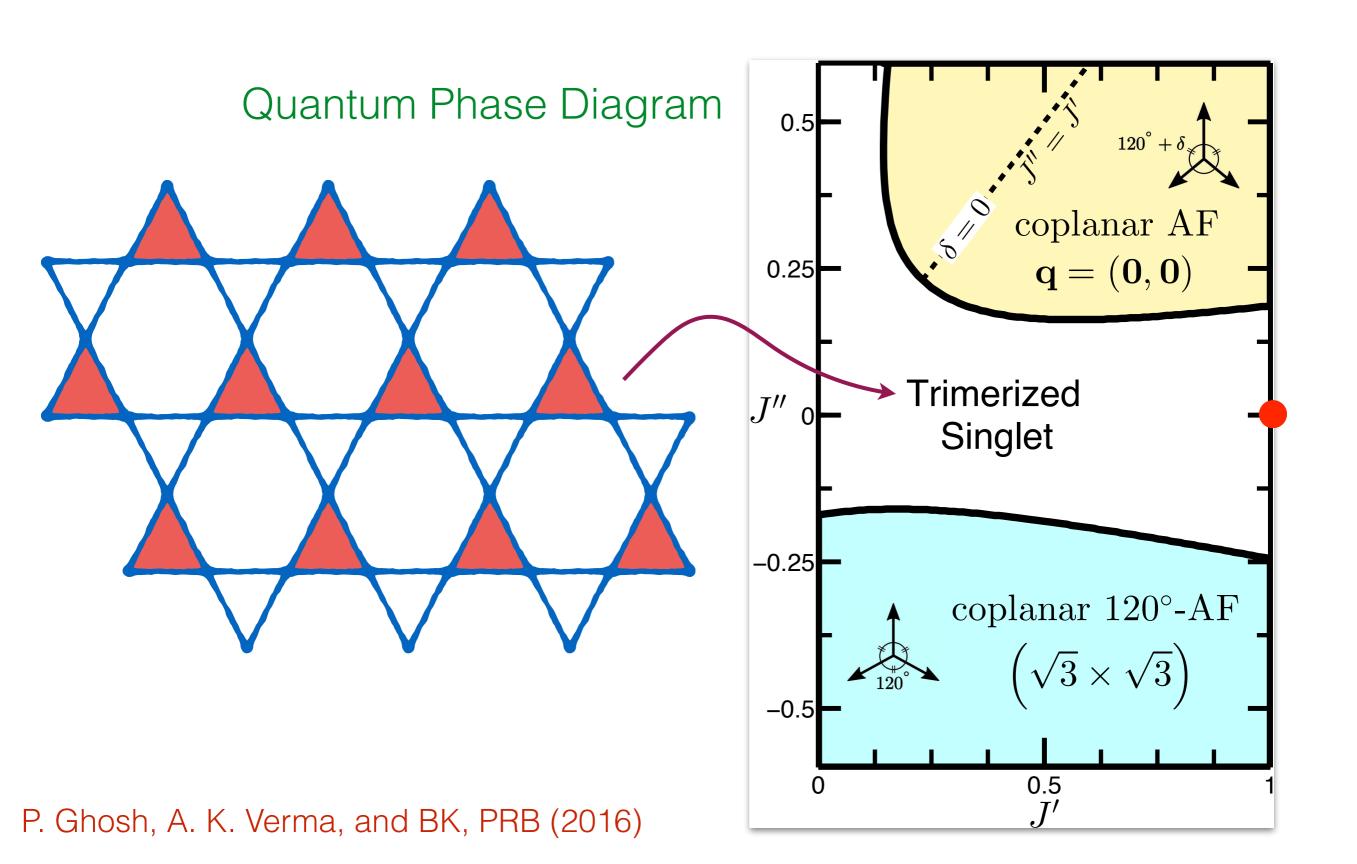
Quantum Phase Diagram

$$\hat{H} = J \sum_{\Delta} \vec{S}_i \cdot \vec{S}_j$$

$$+J' \sum_{\nabla} \vec{S}_i \cdot \vec{S}_j$$

$$+J'' \sum_{\langle \langle i,j \rangle \rangle} \vec{S}_i \cdot \vec{S}_j$$





Magnetic Order with wave vector q

$$\mathbf{m}_{j}(\mathbf{r}) = m_{j}(\cos{[\varphi_{j} - \mathbf{q} \cdot \mathbf{r}]}, \sin{[\varphi_{j} - \mathbf{q} \cdot \mathbf{r}]}, 0)$$

$$(m_1,\varphi_1) = \left(2\bar{s}\sqrt{\frac{n_c}{3(1+\zeta)}},\frac{\pi}{2}\right),\,$$

$$(m_2, \varphi_2) = \left(m_1 \frac{\sqrt{1+3\zeta}}{2}, \varphi_1 + \frac{2\pi}{3} + \delta\right)$$

$$(m_3,\varphi_3)=\left(m_2,\varphi_1-\frac{2\pi}{3}-\delta\right).$$

$$(m_3, \varphi_3) = \left(m_2, \varphi_1 - \frac{2\pi}{3} - \delta\right).$$

$$\delta = \tan^{-1} \left[\sqrt{3}(1 - \sqrt{\zeta})/(1 + 3\sqrt{\zeta})\right]$$

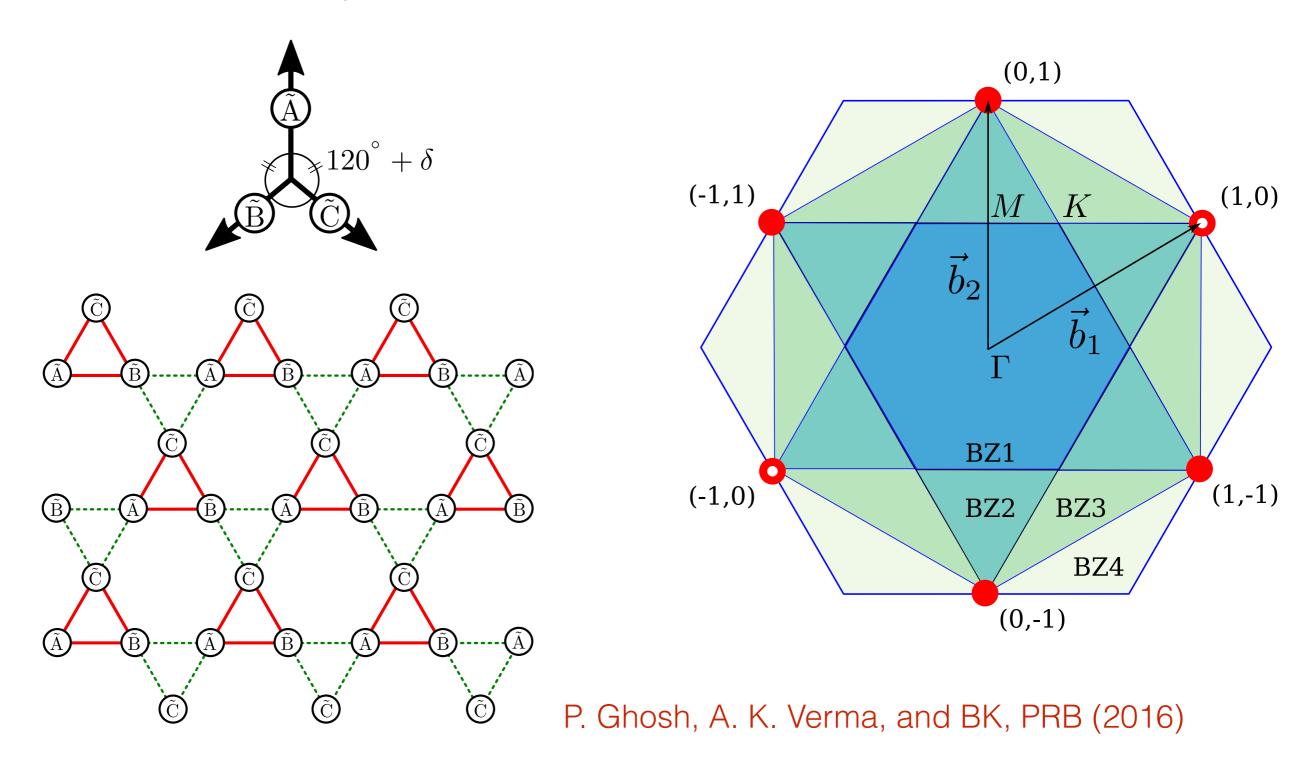
$$\zeta = n_{c1}/n_{c\bar{1}}$$

Magnetic Order in the Two Phases $\sqrt{3} \times \sqrt{3}$

Magnetic Order in the Two Phases Novel q=(0,0) order $\sqrt{3} \times \sqrt{3}$ $120^{\circ} + \delta$

P. Ghosh, A. K. Verma, and BK, PRB (2016)

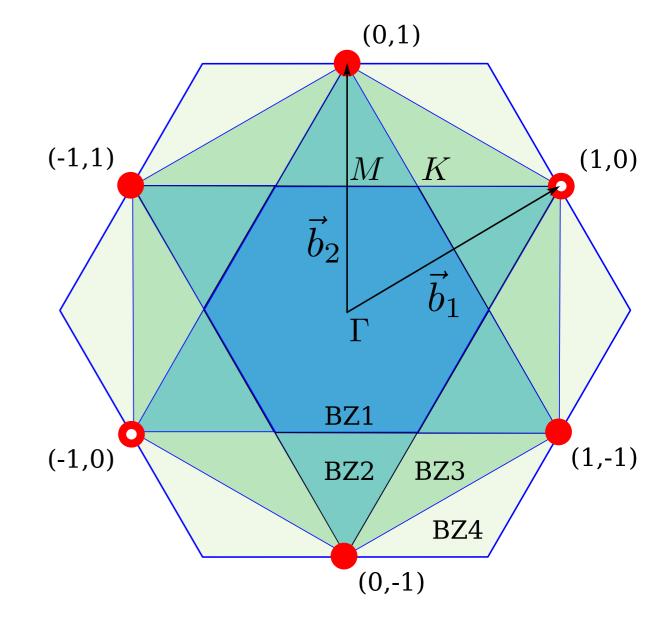
Structure factor of q=(0,0) order with $\delta \neq 0$



Structure factor of q=(0,0) order with $\delta \neq 0$

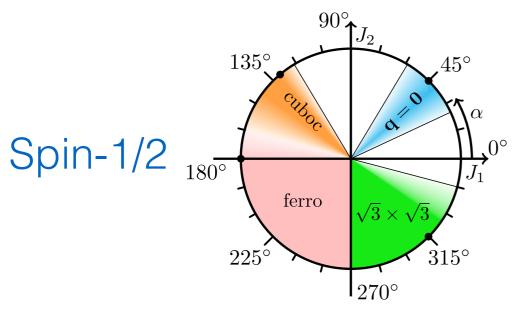
experimental measure of δ

$$\frac{I_{(1,0)}}{I_{(0,1)}} = \frac{4}{1+3\zeta} = 4\sin^2\left(\frac{\pi}{6} + \delta\right)$$

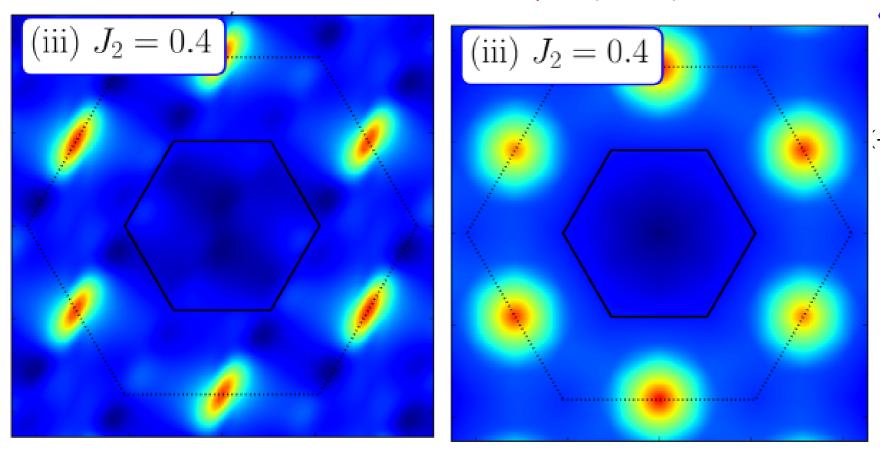


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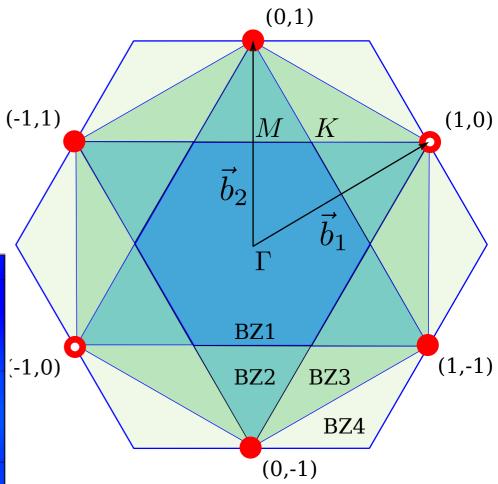
Structure factor for q=(0,0) order



FRG; Suttner et al, PRB-Rapid (2014)



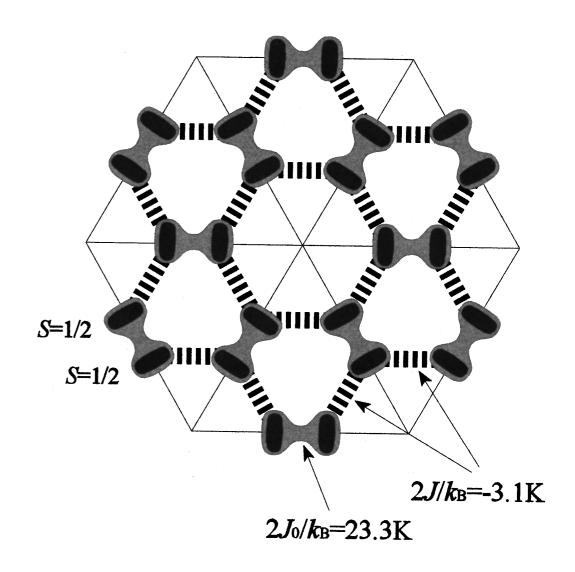
Spin-1: δ≠0



P. Ghosh, A. K. Verma, and BK, PRB (2016)

DMRG; Kolley et al, PRB (2015)

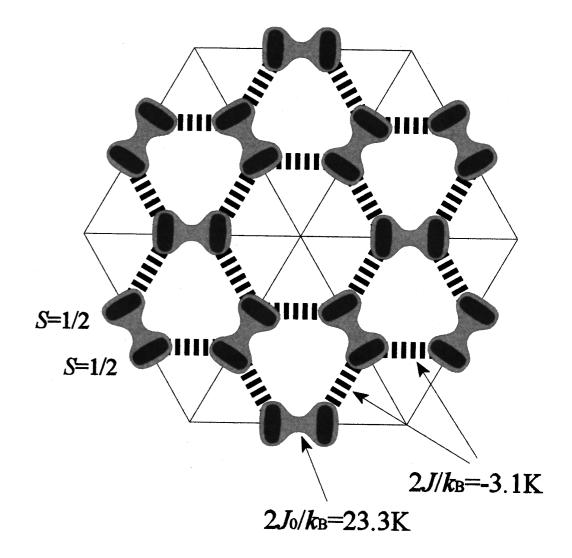
What state do they realize?



 $m - MPYNN.BF_4$

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- Trimerized singlet state, or
- the HSS state of Hida?



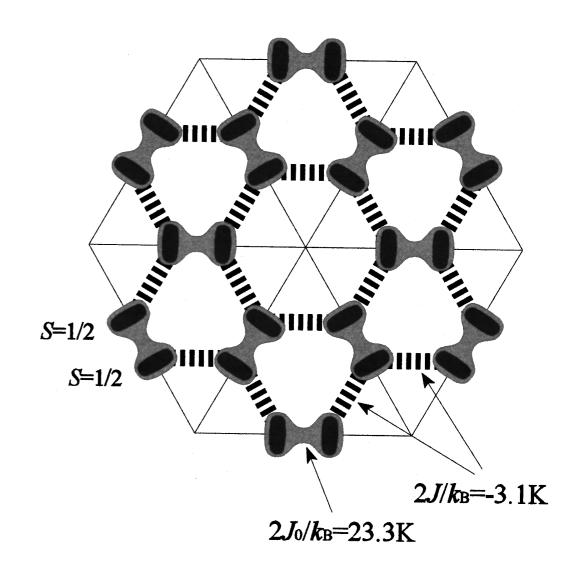
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TABLE I. The exchange interactions for m-MPYNN⁺ · X^+ · $\frac{1}{3}$ (acetone) from Refs. [16,17].

X	I	BF ₄	$(BF_4)_{0.72}I_{0.28}$	ClO ₄
J_A	1.6 K	3.11 K	1.20 K	0.19 K
$J_F \ J_F/J_A$	- 10.2 K - 6.375	- 23.26 K - 7.479	− 11.3 K − 9.416	- 10.5 K - 55.263



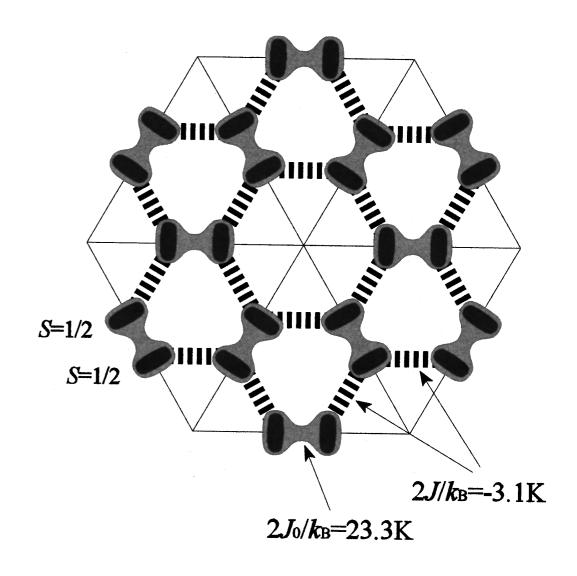
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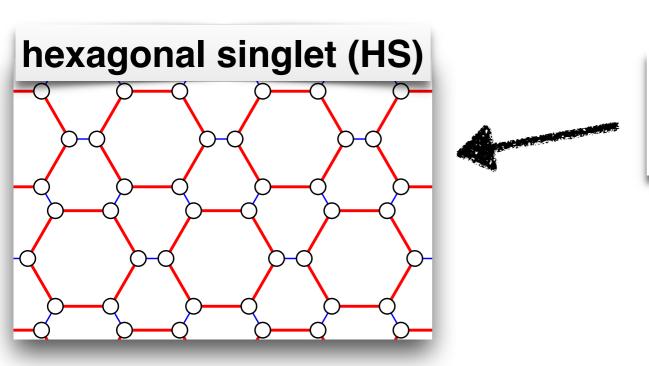
How does the ground state of the Hida model evolve with J_F/J_A ?

How does the ground state evolve with J_F/J_A ?

Candidate states.

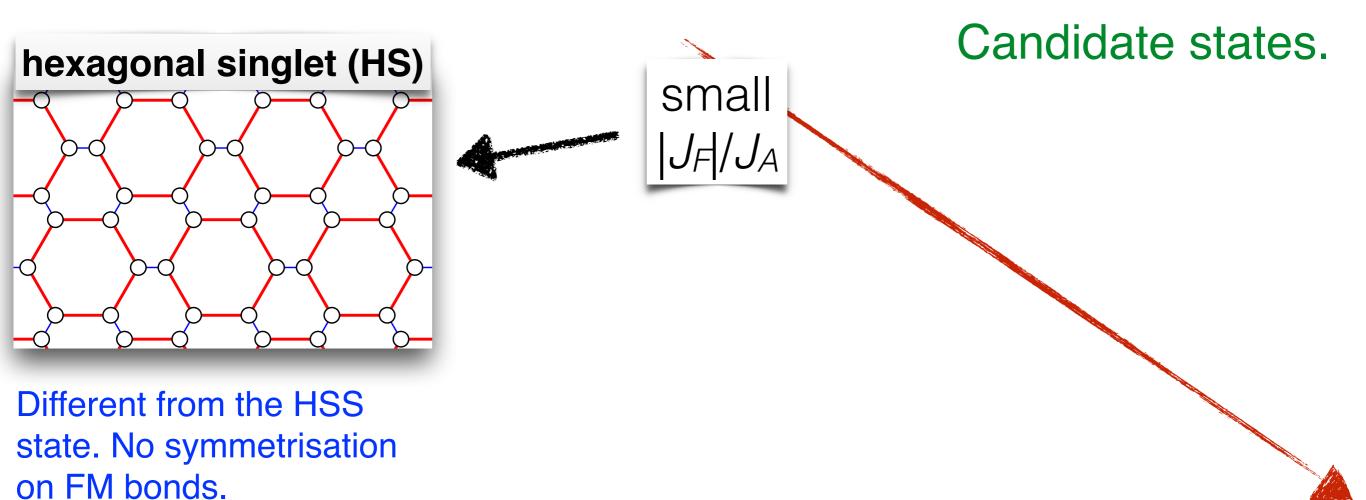
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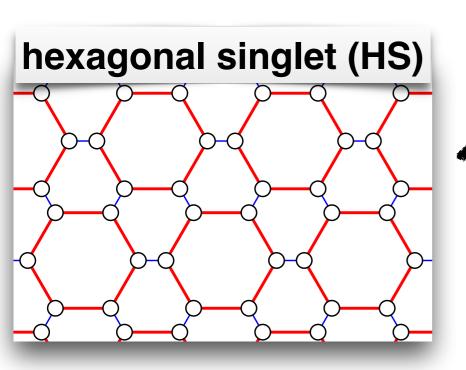
 $|J_F|/J_A$



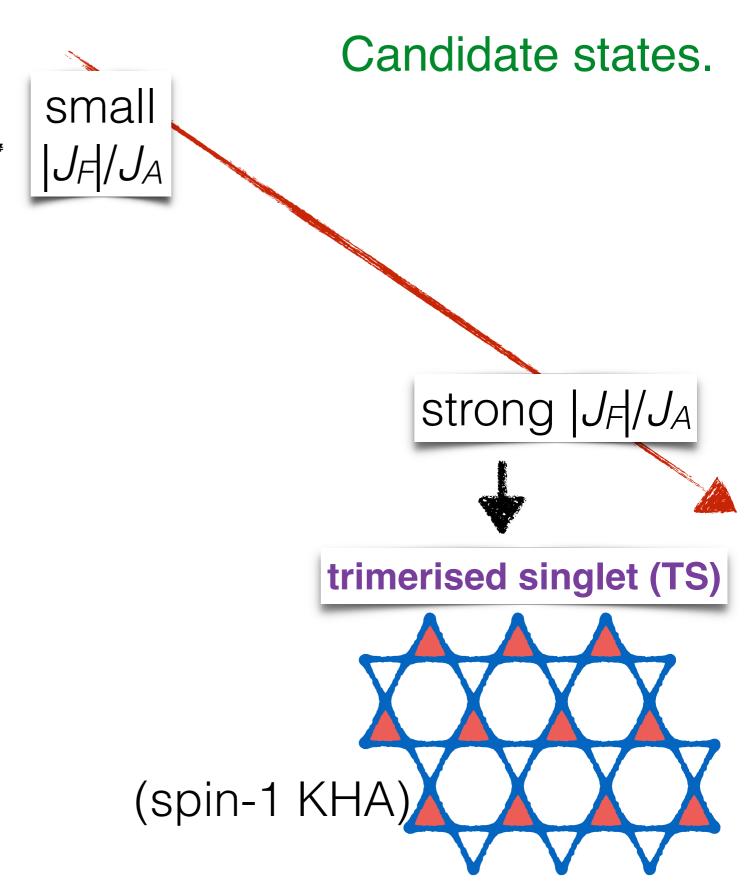
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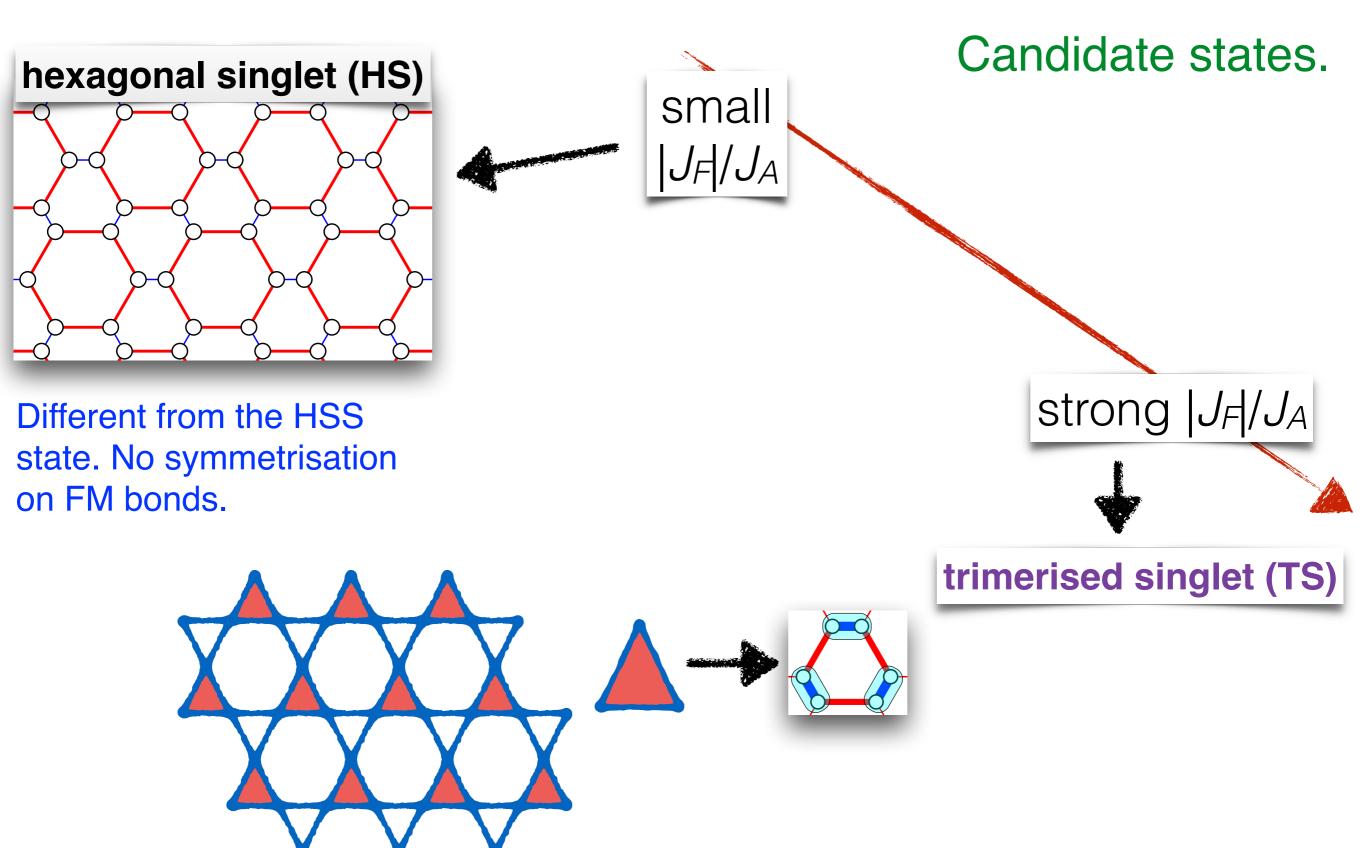
Different from the HSS state. No symmetrisation on FM bonds.

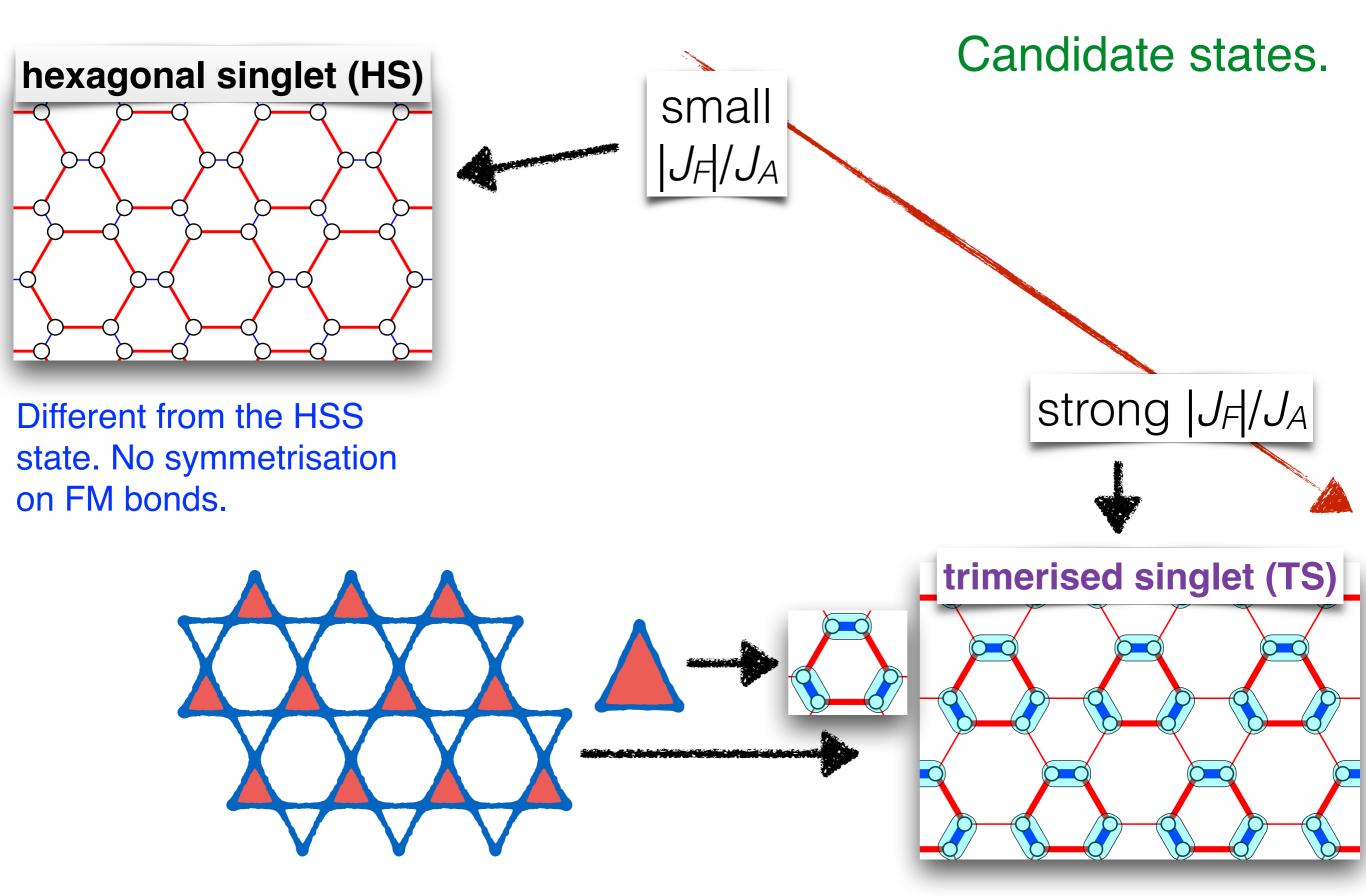




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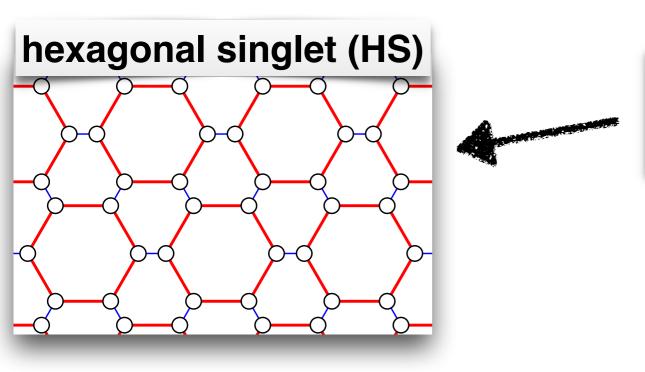




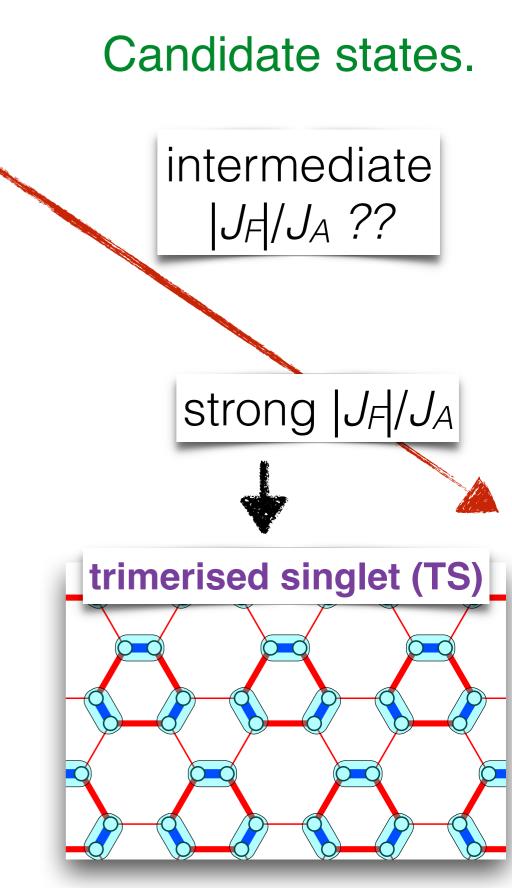


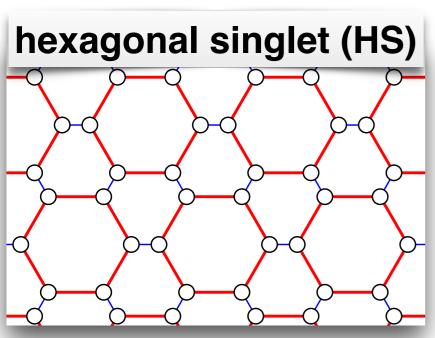
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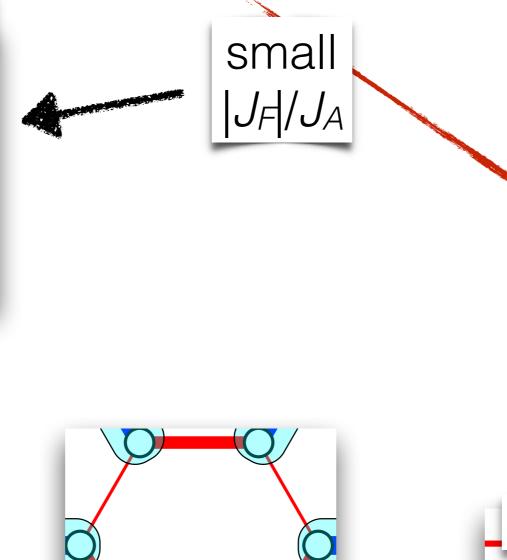


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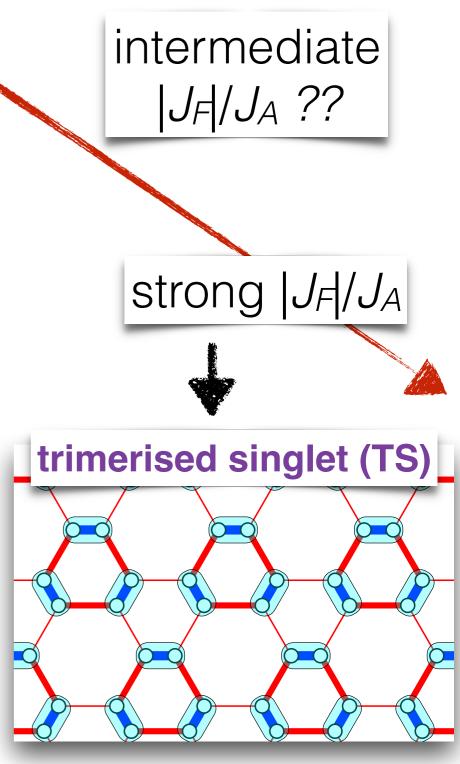




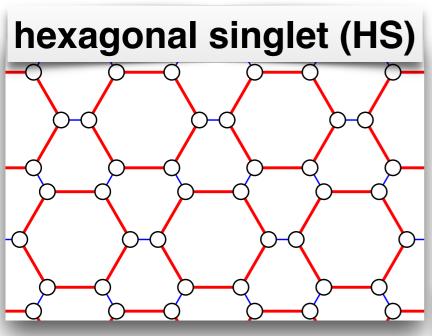
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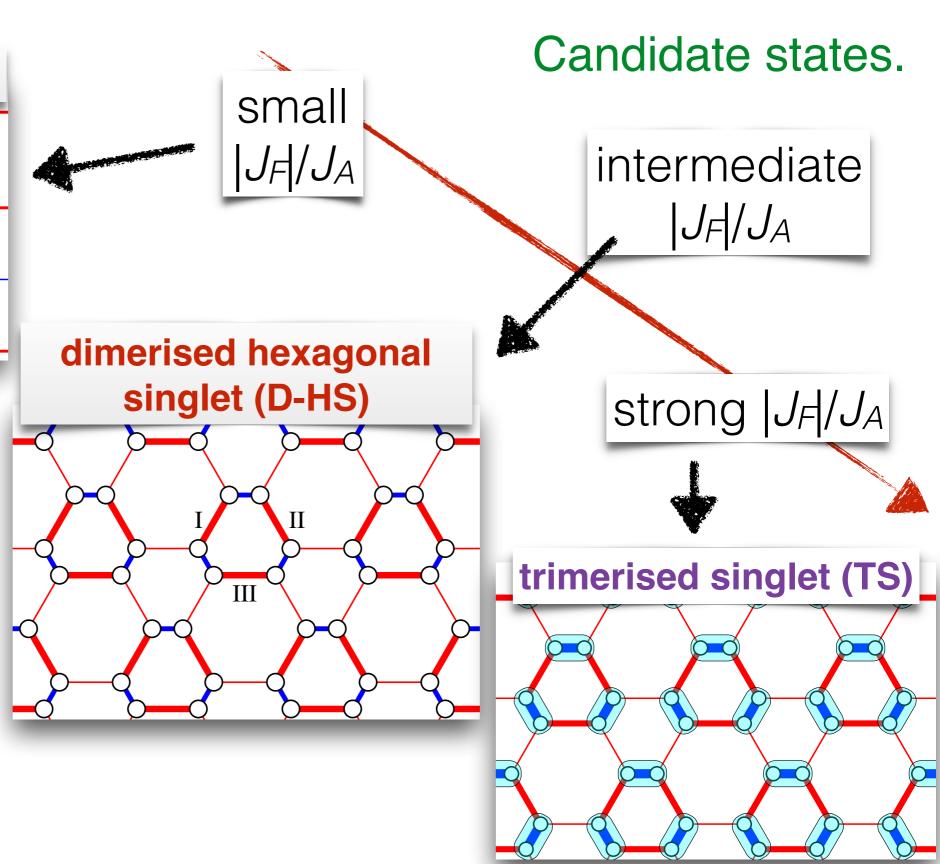
Antiferromagnetic hexagons are dimerized.

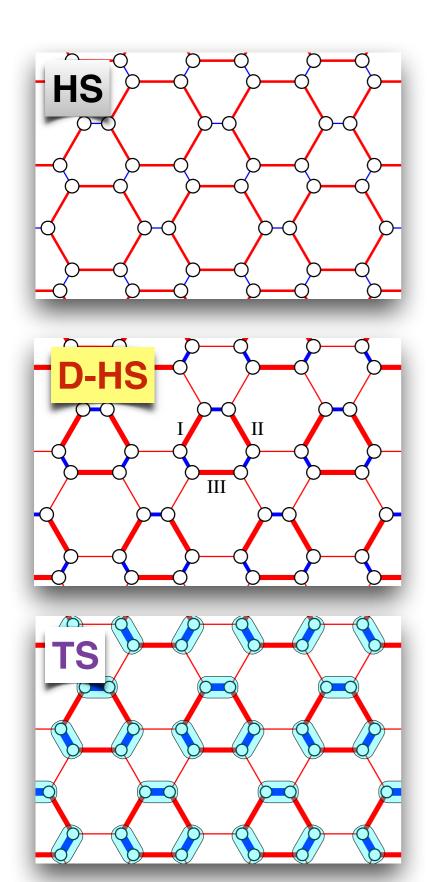


Candidate states.

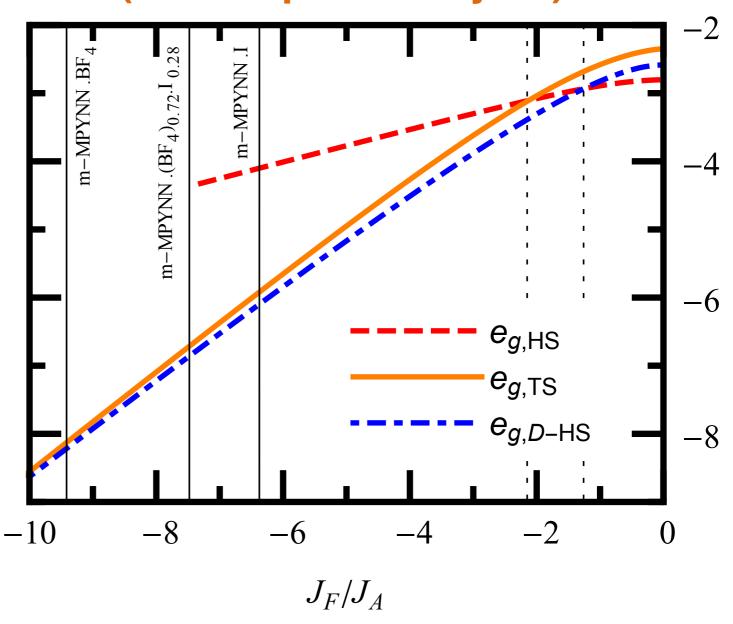


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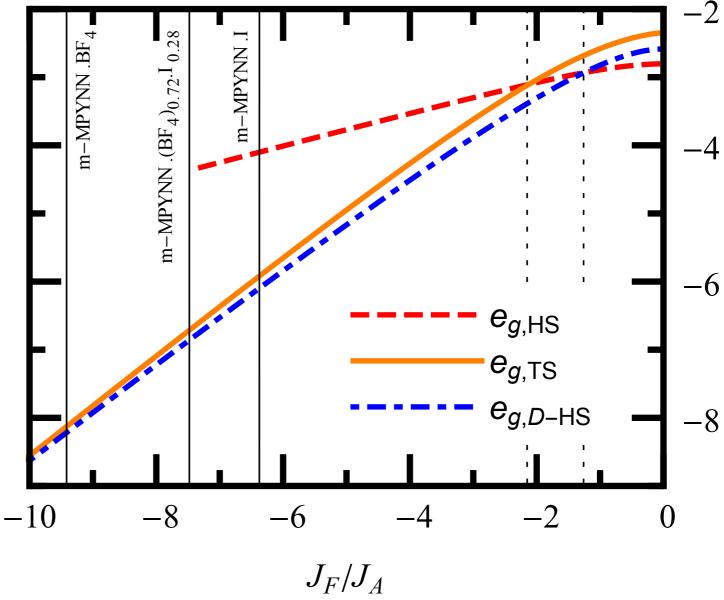


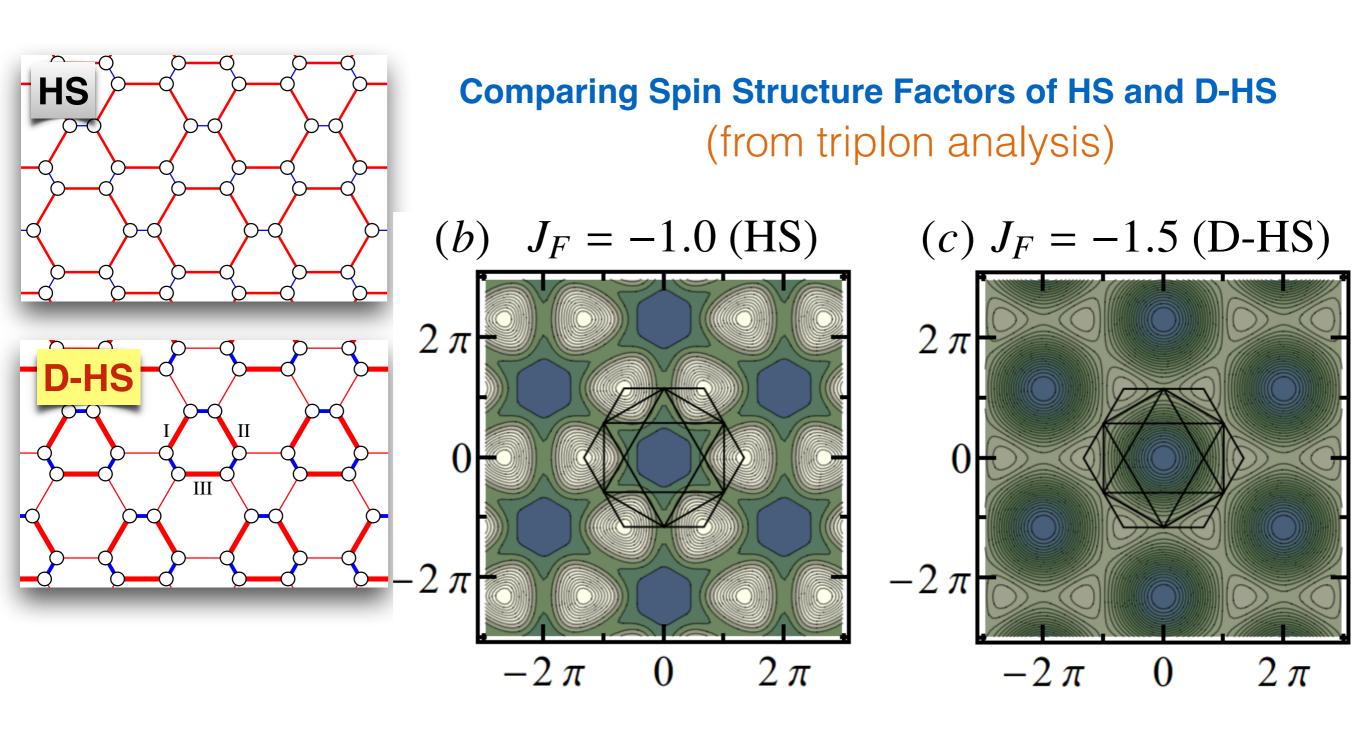
Comparison of their energies (from triplon analysis)

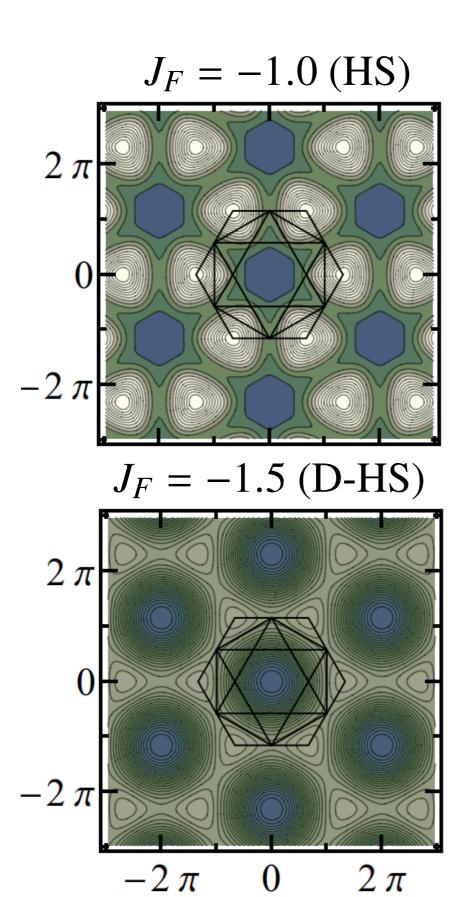


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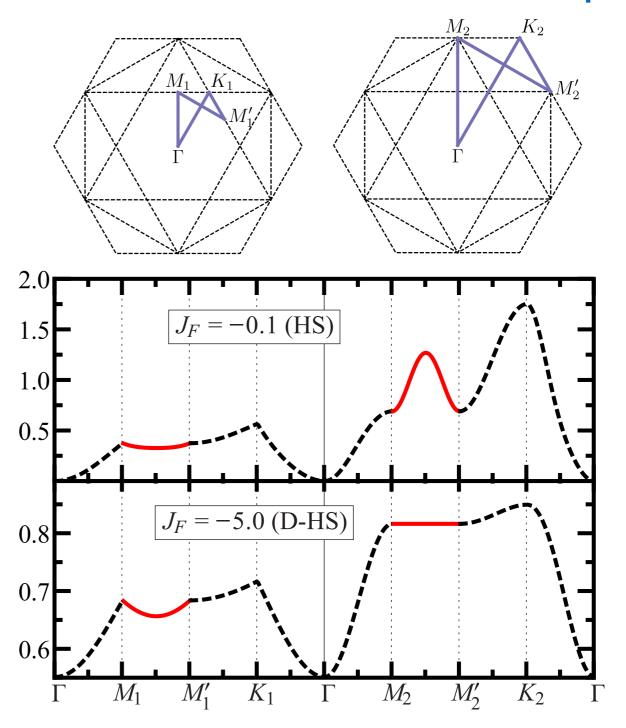
It suggests that the m-MPYNN.X family of organic salts realise the D-HS (TS) phase at low temperatures.





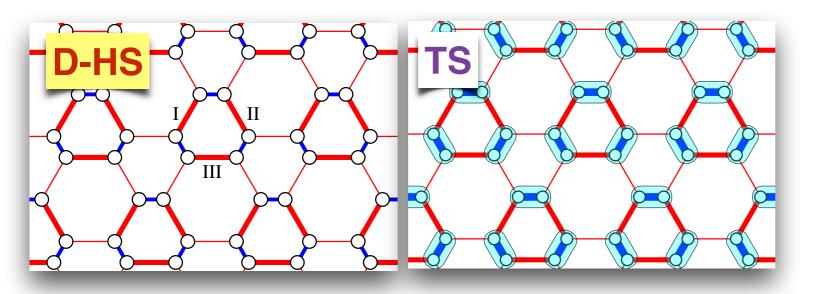


Static Structure Factors of HS and D-HS phases

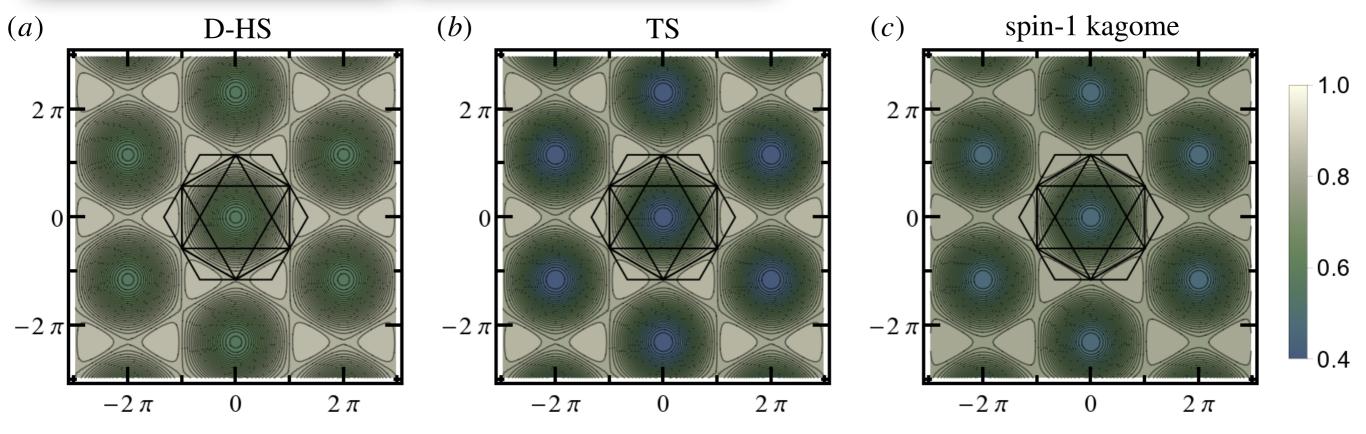


P. Ghosh and <u>BK</u>, PRB 97, 014413 (2018)

Comparing Spin Structure Factors of D-HS and TS



(from triplon analysis)



Schwinger boson representation

$$S_i^+ = a_i^{\dagger} b_i,$$
 $S_i^- = b_i^{\dagger} a_i,$
 $S_i^z = \frac{1}{2} (a_i^{\dagger} a_i - b_i^{\dagger} b_i)$
 $a_i^{\dagger} a_i + b_i^{\dagger} b_i = 2S \text{ (constraint)}$

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Heisenberg exchange operator

$$\vec{S}_i \cdot \vec{S}_j = :F_{ij}^{\dagger} F_{ij} : -A_{ij}^{\dagger} A_{ij},$$

$$A_{ij} = \frac{1}{\sqrt{2}} (a_i b_j - b_i a_j)$$

$$F_{ij} = \frac{1}{\sqrt{2}} (a_i^{\dagger} a_j + b_i^{\dagger} b_j)$$

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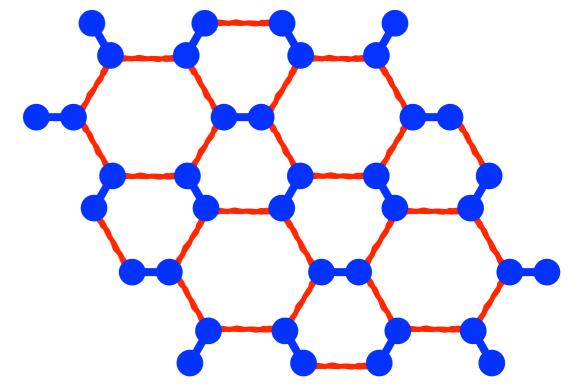
Mean-field approximation for Hida model:

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$$\alpha_F = \langle A_{ij} \rangle$$
 and $\phi_F = \langle F_{ij} \rangle$ (FM bonds)
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$$\alpha'_A = \langle A_{ij} \rangle \text{ and } \phi'_A = \langle F_{ij} \rangle$$

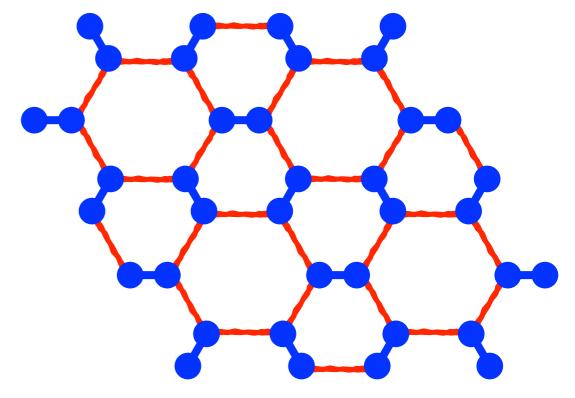
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(AFM bonds)

Heisenberg exchange operator

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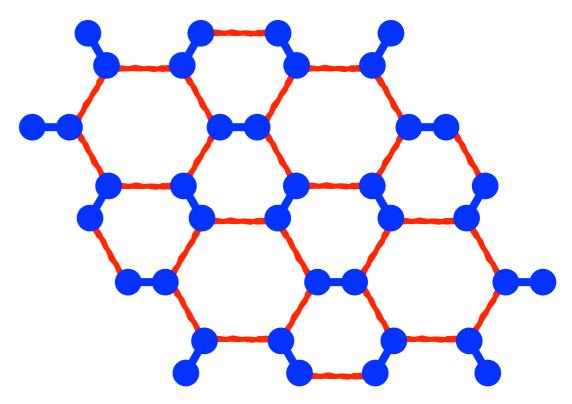
determine these mean-field parameters self-consistently.

Heisenberg exchange operator

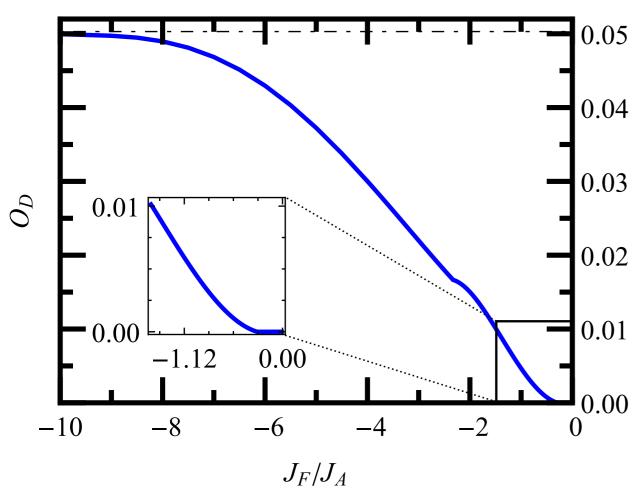
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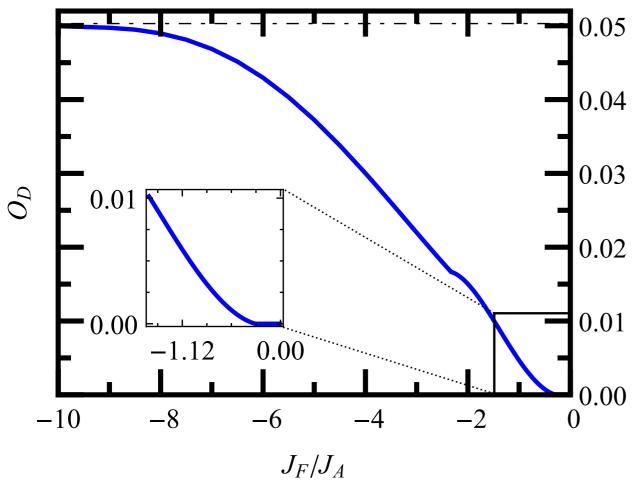
Two Quantum Phase Transitions



spontaneous dimerisation

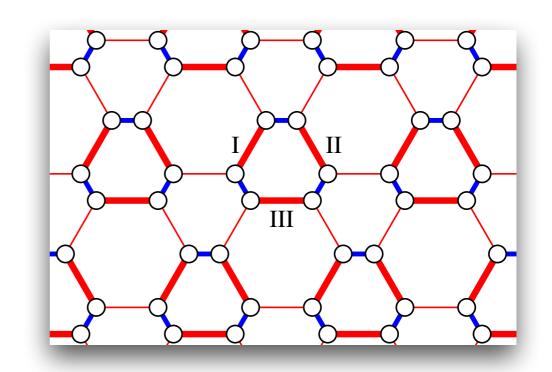
 $J_F/J_A \sim -0.28$

Two Quantum Phase Transitions



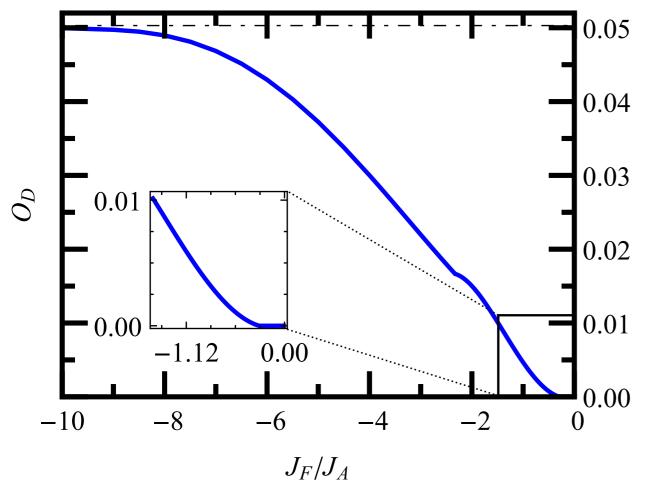
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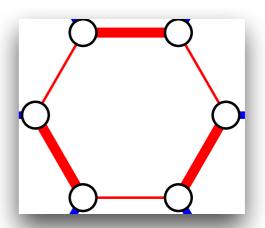
$$O_D = \frac{1}{3N_{\rm uc}} \left| \sum_{\langle i,j \rangle}^{\circ \circ} \langle \vec{S}_i \cdot \vec{S}_j \rangle - \sum_{\langle i,j \rangle}^{\circ \circ} \langle \vec{S}_i \cdot \vec{S}_j \rangle \right|$$

Two Quantum Phase Transitions



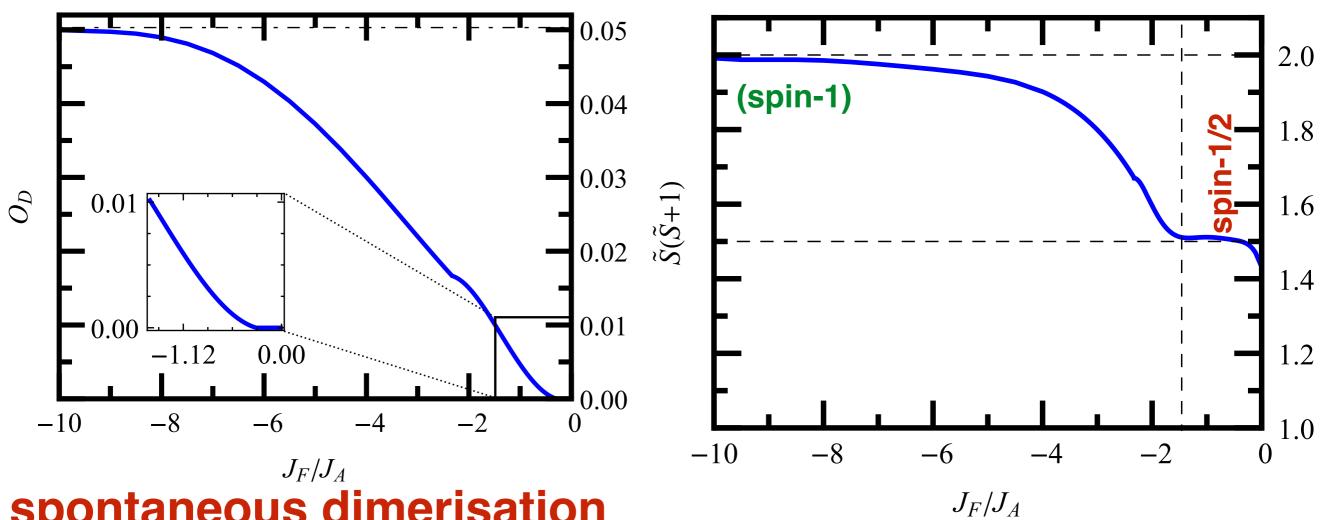
spontaneous dimerisation

 $J_F/J_A \sim -0.28$



$$O_D = \frac{1}{3N_{\rm uc}} \left| \sum_{\langle i,j \rangle}^{2} \langle \vec{S}_i \cdot \vec{S}_j \rangle - \sum_{\langle i,j \rangle}^{2} \langle \vec{S}_i \cdot \vec{S}_j \rangle \right|$$

Two Quantum Phase Transitions



spontaneous dimerisation

$$J_F/J_A \sim -0.28$$

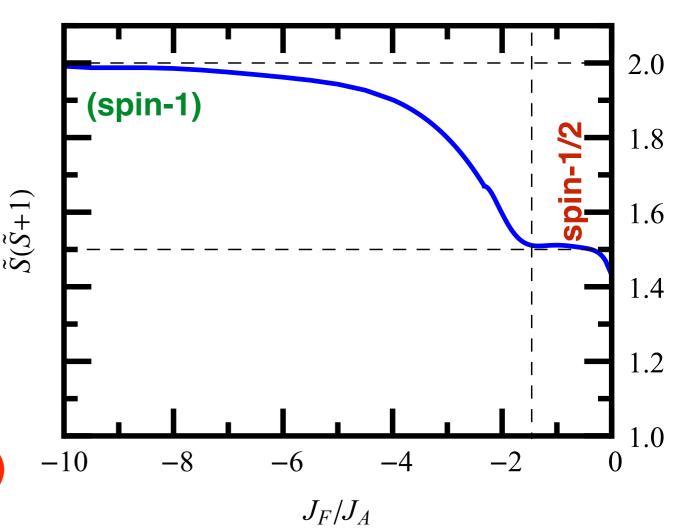
moment formation $J_F/J_A \sim -1.46$

Two Quantum Phase Transitions

$$\widetilde{S}(\widetilde{S}+1) = \frac{1}{3N_{\text{uc}}} \sum_{\langle i,j \rangle}^{\infty} \langle (\vec{S}_i + \vec{S}_j)^2 \rangle$$

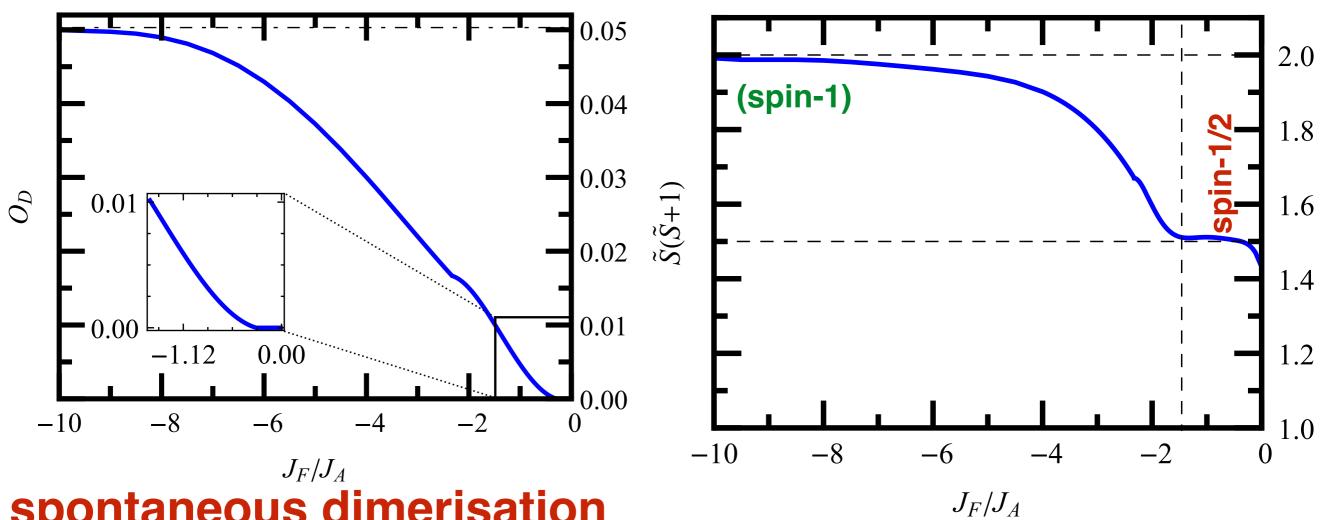
$$= \frac{1}{3N_{\text{uc}}} \sum_{\langle i,j \rangle}^{\infty} (\langle \vec{S}_i^2 \rangle + \langle \vec{S}_j^2 \rangle + 2\langle \vec{S}_i \cdot \vec{S}_j \rangle)$$

(average total spin per FM bond)



moment formation $J_F/J_A \sim -1.46$

Two Quantum Phase Transitions



spontaneous dimerisation

$$J_F/J_A \sim -0.28$$

moment formation $J_F/J_A \sim -1.46$

Summary

The spin-1/2 Hida model, applicable to a family of spin-gapped organic salts, exhibits two transition in the spin-gapped, singlet ground with respect to J_F/J_A .

It first undergoes **spontaneous dimerisation**" at $J_F/J_A \sim -0.3$

For $J_F/J_A \sim -1.46$, pair of spin-1/2 moments on every FM bond begin to behave as "bound" moment.

The spontaneously trimerised singlet phase of the spin-1 KHA is same as the dimerised-HS phase of the Hida model for strong $J_F/J_{A.}$

The m-MPYNN.X family of organic salts realise the D-HS (TS) phase at low temperatures.

Thank you.