

## Gaussian relaxational models in the ordered phase (12 points)

Model A/B ( $\alpha = 0, 2$ ) coupled Langevin equations:  $\alpha = 1, \dots, n$

$$\frac{\partial S^\alpha(\vec{x}, t)}{\partial t} = -D(i\vec{v})^\alpha \left[ (E - \vec{v}^2 + \frac{u}{6} \sum_{\beta=1}^n S^\beta(\vec{x}, t)^2) S^\alpha(\vec{x}, t) - h^\alpha(\vec{x}, t) \right] + \xi^\alpha(\vec{x}, t)$$

$$\langle S^\alpha(\vec{x}, t) \rangle = 0, \quad \langle S^\alpha(\vec{x}, t) S^\beta(\vec{x}, t') \rangle = 2Dh_B T (i\vec{v}_x)^\alpha \delta(\vec{x} - \vec{x}') \delta(t - t') S^{\alpha\beta}$$

$T < T_c$  ( $T < 0$ ): parametrize  $S^\alpha = \pi^\alpha$ ,  $\alpha = 1, \dots, n-1$

$$S^n = \sqrt{\frac{6|T|}{u}} + \sigma$$

in terms of the fluctuating fields  $\pi^\alpha, \sigma$ :

$$\begin{aligned} \frac{\partial \pi^\alpha(\vec{x}, t)}{\partial t} = & -D(i\vec{v})^\alpha \left[ (-|v| - \vec{v}^2 + \frac{u}{6} \sqrt{\frac{6|T|}{u}} + 2\sqrt{\frac{6|T|}{u}} \sigma(\vec{x}, t)) \pi^\alpha(\vec{x}, t) \right. \\ & \left. + \sum_{\beta=1}^{n-1} \pi^\beta(\vec{x}, t)^2 \right] + \xi^\alpha(\vec{x}, t) \end{aligned}$$

$$\text{linearize: } \approx D(i\vec{v})^\alpha \left[ \vec{v}^2 \pi^\alpha(\vec{x}, t) + h^\alpha(\vec{x}, t) \right] + \xi^\alpha(\vec{x}, t)$$

$$\begin{aligned} \frac{\partial \sigma(\vec{x}, t)}{\partial t} = & -D(i\vec{v})^\alpha \left[ (-|v| - \vec{v}^2 + \frac{u}{6} \sqrt{\frac{6|T|}{u}} + 2\sqrt{\frac{6|T|}{u}} \sigma(\vec{x}, t) + \sigma(\vec{x}, t)^2 \right. \\ & \left. + \sum_{\beta=1}^{n-1} \pi^\beta(\vec{x}, t)^2) \left( \sqrt{\frac{6|T|}{u}} + \sigma(\vec{x}, t) \right) - h^n(\vec{x}, t) \right] + \xi^n(\vec{x}, t) \end{aligned}$$

$$\text{linearize: } \approx -D(i\vec{v})^\alpha \left[ (2|v| - \vec{v}^2) \sigma(\vec{x}, t) - h^n(\vec{x}, t) \right] + \xi^n(\vec{x}, t)$$

$$\text{Fourier transform: } (-i\omega + Dq^2 + \alpha) \pi^\alpha(\vec{q}, \omega) = Dq^\alpha h^\alpha(\vec{q}, \omega) + \xi^\alpha(\vec{q}, \omega)$$

$$[-i\omega + Dq^\alpha (2|v| + q^2)] \sigma(\vec{q}, \omega) = Dq^\alpha h^n(\vec{q}, \omega) + \xi^n(\vec{q}, \omega)$$

$$\text{with } \langle \xi^\alpha(\vec{q}, \omega) \rangle = 0, \quad \langle \xi^\alpha(\vec{q}, \omega) \xi^\beta(\vec{q}', \omega') \rangle = 2Dh_B T q^\alpha (2\pi)^{d+1} \delta(\vec{q} + \vec{q}') \delta(\omega - \omega') S^{\alpha\beta}$$

→ dynamic response functions in Gaussian approximation:

$$X_0^{AB}(\vec{q}, \omega) = \frac{\partial \langle \pi^\alpha(\vec{q}, \omega) \rangle}{\partial h^\beta(\vec{q}, \omega)} = X_0^T(\vec{q}, \omega) S^{\alpha\beta},$$

$$X_0^T(\vec{q}, \omega) = \frac{1}{-\frac{i\omega}{Dq^2} + q^2}$$

transverse response function

"massless":

$n-1$  Goldstone modes

$$X_0^{nn}(\vec{q}, \omega) = \frac{\partial \langle \sigma(\vec{q}, \omega) \rangle}{\partial h^n(\vec{q}, \omega)} = X_0^L(\vec{q}, \omega),$$

$$X_0^L(\vec{q}, \omega) = \frac{1}{-\frac{i\omega}{Dq^2} + 2|v| + q^2}$$

longitudinal response function

and  $\chi_o^T(\vec{q}, t) = Dq^a e^{-Dq^2 + \alpha t} \theta(t)$  always "critical"

$$\chi_o^L(\vec{q}, t) = Dq^a e^{-Dq^2(2|\tau| + q^2)t} \theta(t)$$

dynamic correlation functions in Gaussian approximation:

$$C_o^T(\vec{q}, \omega) = \frac{2Dk_B T q^a}{\omega^2 + (Dq^2 + \alpha)^2}, \quad C_o^L(\vec{q}, \omega) = \frac{2Dk_B T q^a}{\omega^2 + [Dq^2(2|\tau| + q^2)]^2}$$

$$C_o^T(\vec{q}, t) = \frac{k_B T}{q^2} e^{-Dq^2 + \alpha |t|}$$

$$C_o^L(\vec{q}, t) = \frac{k_B T}{2|\tau| + q^2} e^{-Dq^2(2|\tau| + q^2)|t|}$$

note static limits:  $\chi_o^T(\vec{q}) = \chi_o^T(\vec{q}, \omega=0) = \frac{1}{q}$  "massless"

$$\chi_o^L(\vec{q}) = \frac{1}{2|\tau| + q^2}$$
 "massive"

$$C_o^T(\vec{q}) = C_o^T(\vec{q}, t=0) = \frac{k_B T}{q^2}, \quad C_o^L(\vec{q}) = \frac{k_B T}{2|\tau| + q^2}$$

coupling of longitudinal fluctuations to Goldstone modes

induces coexistence singularities:

$$\chi^L(\vec{q}, \omega) \sim \frac{1}{\left[ -i\omega + \frac{1}{Dq^a} + q^2 + h^n \right]^{\frac{d-4}{2}}}$$



## Reversible mode-coupling terms and equilibrium condition

(8 points)

equilibrium condition for reversible stationary probability current:

divergence-free in space of random variables:

$$0 \stackrel{!}{=} \int d^d x \sum_{\alpha} \frac{\delta}{\delta S^{\alpha}(\vec{x})} \left( F_{\text{rev}}^{\alpha}[S^{\alpha}](\vec{x}) e^{-H[S^{\alpha}](\vec{x})/k_B T} \right) = \\ = \int d^d x \sum_{\alpha} \left[ \frac{\delta F_{\text{rev}}^{\alpha}[S^{\alpha}](\vec{x})}{\delta S^{\alpha}(\vec{x})} - \frac{F_{\text{rev}}^{\alpha}[S^{\alpha}](\vec{x})}{k_B T} \frac{\delta H[S^{\alpha}]}{\delta S^{\alpha}(\vec{x})} \right] e^{-H[S^{\alpha}](\vec{x})/k_B T}$$

ansatz:

$$F_{\text{rev}}^{\alpha}[S^{\alpha}](\vec{x}) = - \int d^d x' \sum_{\alpha'} \left[ Q^{\alpha \alpha'}(\vec{x}, \vec{x}') \frac{\delta H[S^{\alpha}]}{\delta S^{\alpha'}(\vec{x}')} - k_B T \frac{\delta Q^{\alpha \alpha'}(\vec{x}, \vec{x}')}{\delta S^{\alpha'}(\vec{x}')} \right]$$

with antisymmetric  $Q^{\alpha \alpha'}(\vec{x}, \vec{x}') = -Q^{\alpha' \alpha}(\vec{x}', \vec{x})$

$$\rightarrow \sum_{\alpha} \frac{\delta F_{\text{rev}}^{\alpha}[S^{\alpha}](\vec{x})}{\delta S^{\alpha}(\vec{x})} = - \int d^d x' \sum_{\alpha, \alpha'} \left[ \frac{\delta Q^{\alpha \alpha'}(\vec{x}, \vec{x}')}{\delta S^{\alpha}(\vec{x})} \frac{\delta H[S^{\alpha}]}{\delta S^{\alpha'}(\vec{x}')} \right. \\ \left. + Q^{\alpha \alpha'}(\vec{x}, \vec{x}') \frac{\delta^2 H[S^{\alpha}]}{\delta S^{\alpha}(\vec{x}) \delta S^{\alpha'}(\vec{x}')} - k_B T \underbrace{\frac{\delta^2 Q^{\alpha \alpha'}(\vec{x}, \vec{x}')}{\delta S^{\alpha}(\vec{x}) \delta S^{\alpha'}(\vec{x}')}}_{=0 \text{ from antisymmetry}} \right]$$

after  $\int d^d x'$  integration: vanish because of antisymmetry  
 $\alpha \leftrightarrow \alpha', \vec{x} \leftrightarrow \vec{x}'$

This leaves in equilibrium condition:

$$- \int d^d x \int d^d x' \sum_{\alpha, \alpha'} \left[ \frac{\delta Q^{\alpha \alpha'}(\vec{x}, \vec{x}')}{\delta S^{\alpha}(\vec{x})} \frac{\delta H[S^{\alpha}]}{\delta S^{\alpha'}(\vec{x}')} - \frac{Q^{\alpha \alpha'}(\vec{x}, \vec{x}')}{k_B T} \frac{\delta H[S^{\alpha}]}{\delta S^{\alpha}(\vec{x})} \frac{\delta H[S^{\alpha}]}{\delta S^{\alpha'}(\vec{x}')} \right. \\ \left. + Q^{\alpha \alpha'}(\vec{x}, \vec{x}') \frac{\delta^2 H[S^{\alpha}]}{\delta S^{\alpha}(\vec{x}) \delta S^{\alpha'}(\vec{x}')} - k_B T \underbrace{\frac{\delta^2 Q^{\alpha \alpha'}(\vec{x}, \vec{x}')}{\delta S^{\alpha}(\vec{x}) \delta S^{\alpha'}(\vec{x}')}}_{=0 \text{ from antisymmetry}} \right] e^{-H[S^{\alpha}]/k_B T} = 0 \checkmark$$

combine to  $0$  after  $\alpha \leftrightarrow \alpha', \vec{x} \leftrightarrow \vec{x}'$   
in second term