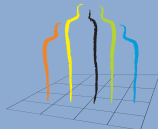


Transport Processes in Melts under External Fields



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On the Yielding of Colloidal (and Other) Glass Formers

Thomas Voigtmann

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J. A. Krumhansl Symposium, Bangalore, February 2012



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- S. Egelhaaf (D'dorf), M. Ballauff, M. Siebenbürger (HZ Berlin)



Outline

- Dynamical Yield Stress
- Startup: Creep and Micro-Rheology
- Residual Stresses

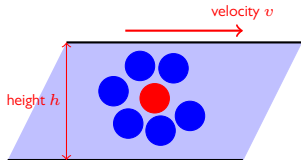


Introduction



Rheology of Dense Fluids

shear flow of dense fluids:



- external **flow rate** $\dot{\gamma} \sim v/h$ [1/s]
 - large structural **relaxation time** τ [s]
- \Rightarrow large effect when $\dot{\gamma}\tau \gg 1$
- (glassy) kinetic arrest: $\tau \rightarrow \infty$

- **apply** perturbation (shear $\dot{\gamma}$) \Rightarrow **measure response** (stress σ)

$$\mathcal{F}[\dot{\gamma}] = \sigma \quad \text{constitutive equation}$$

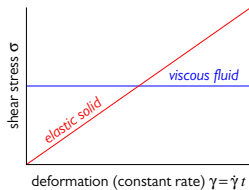
- \mathcal{F} is a *model* of the material
- *linear response, steady state*: Newtonian liquid $\sigma = \dot{\gamma} \times \eta$



Visco-Elasticity: Maxwell's Model

- Newtonian fluid: $\eta = \text{const.} \Rightarrow \sigma \propto \dot{\gamma}$
- Hookian elastic solid: $\sigma \propto \gamma$
- dense fluids: ??
- Maxwell: combine $\sigma \sim \gamma$ and $\sigma \sim \dot{\gamma}$

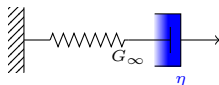
$$\dot{\gamma} = \dot{\sigma}/G_{\infty} + \sigma/\eta$$



- “spring-and-dashpot” model:
Hookian spring constant G_{∞}
- differential equation solved by

$$\sigma(t) = \int_{-\infty}^t \dot{\gamma}(t') G_{\infty} e^{-(t-t')/\tau} dt'$$

output input model



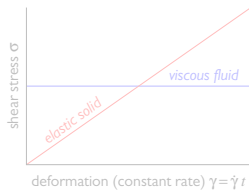
$$\eta = G_{\infty} \tau$$



Visco-Elasticity: Maxwell's Model

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$$\dot{\gamma} = \dot{\sigma}/G_{\infty} + \sigma/\eta$$



- “spring-and-dashpot” model:
Hookian spring constant G_{∞}



- Maxwell's constitutive equation

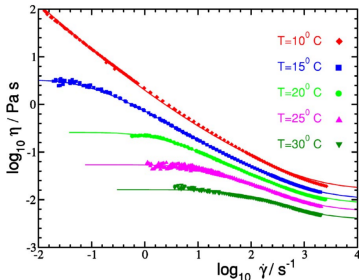
$$\sigma(t) = \int_{-\infty}^t \dot{\gamma}(t') G_{\infty} e^{-(t-t')/\tau} dt'$$

output
input
model

$$\eta = G_{\infty} \tau$$

Nonlinear Rheology: Shear Thinning

apply (steady) shear \Rightarrow dramatic decrease in apparent viscosity



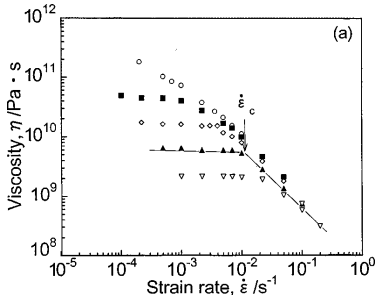
thermosensitive colloids

[Fuchs and Ballauff, J Chem Phys (2005)]

- non-linear response
 - linear response: $\eta \sim \text{const.}$
 - $\eta \rightarrow \infty$: glass
- $\eta \sim 1/\dot{\gamma}$
- applications: painting, coating, lubrication, ...
- “universal”: metallic melts, geophysics, soft matter, ...

Nonlinear Rheology: Shear Thinning

apply (steady) shear \Rightarrow dramatic decrease in apparent viscosity



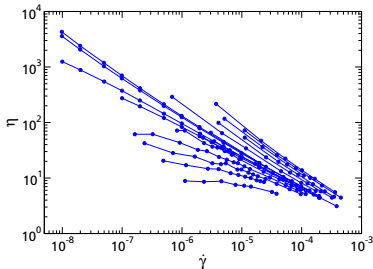
Pd₄₀Ni₁₀Cu₃₀P₂₀, various temperatures

[Kato et al., JAP (1998)]

- non-linear response
 - linear response: $\eta \sim \text{const.}$
 - $\eta \rightarrow \infty$: glass
- $\eta \sim 1/\dot{\gamma}$
- applications: painting, coating, lubrication, ...
- “universal”: metallic melts, geophysics, soft matter, ...

Nonlinear Rheology: Shear Thinning

apply (steady) shear \Rightarrow dramatic decrease in apparent viscosity



granular simulation

[Olsson/Teitel, PRL (2007)]

- non-linear response
 - linear response: $\eta \sim \text{const.}$
 - $\eta \rightarrow \infty$: glass
- $\eta \sim 1/\dot{\gamma}$
- applications: painting, coating, lubrication, ...
- “universal”: metallic melts, geophysics, soft matter, ... (?!?)



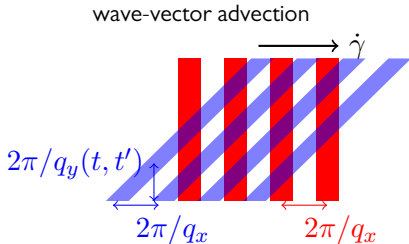
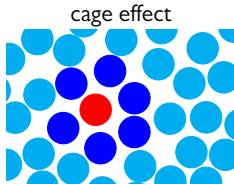
Rheo-Mode-Coupling Theory, Schematically

nonlinear schematic model – strain history $\gamma_{tt'} = \int_{t'}^t \dot{\gamma}(\tau) d\tau$

$$\underbrace{\sigma(t)}_{\text{output}} \sim \int_{-\infty}^t dt' \dot{\gamma}(t') G(t, t', [\dot{\gamma}]) \stackrel{\text{MCT}}{\approx} \int_{-\infty}^t dt' v_{\sigma} \underbrace{\dot{\gamma}(t')}_{\text{model}} \underbrace{\phi^2(t, t', [\gamma])}_{\text{input}}$$

$$\partial_t \phi(t, t') + \phi(t, t') + \int_{t'}^t m(t, t'', t') \partial_{t''} \phi(t'', t') dt'' = 0$$

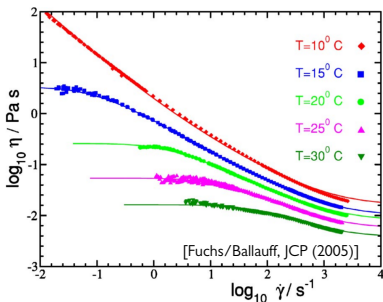
$$m(t, t'', t') = h[\gamma_{tt'}] h[\gamma_{tt''}] (v_1 \phi(t, t'') + v_2 \phi(t, t'')^2)$$



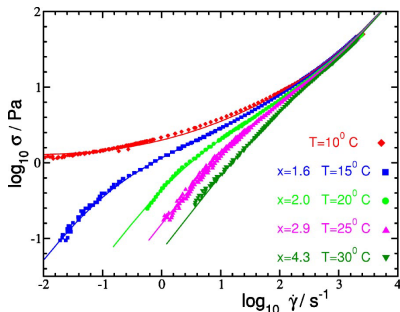


Dynamical Yield Stress

Dynamical Yield Stress



thermosens. colloids: $\eta(\varrho, \dot{\gamma})$



flow curves in steady state

- $\sigma(\dot{\gamma} \rightarrow 0) = \sigma_y > 0$ in the (idealized) glass: **dynamic yield stress**
- $\dot{\gamma} \rightarrow 0$ is singular; $\sigma = \dot{\gamma} \int_{-\infty}^t G(t-t', \dot{\gamma}) dt'$



A Nonlinear Maxwell Model

- *shear accelerates dynamics*: relaxation time $\sim 1/\dot{\gamma}$

$$\tau^{-1} \mapsto \tau^{-1} + \dot{\gamma}$$

- nonlinear Maxwell model

$$G_{\infty} \dot{\gamma} = \dot{\sigma} + \sigma/\tau + \sigma \dot{\gamma} / \gamma_c$$

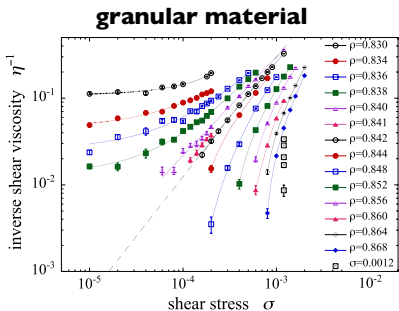
(plus a high-shear Newtonian viscosity...)

$$\sigma = G_{\infty} \tau_0 \dot{\gamma} + G_{\infty} \tau \frac{\dot{\gamma}}{1 + \dot{\gamma} \tau / \gamma_c}$$

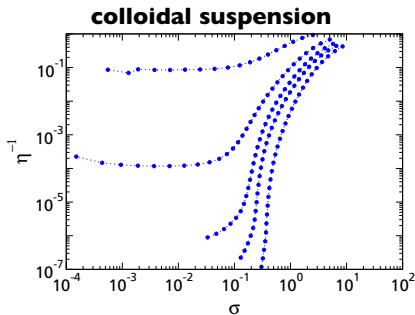
\Rightarrow as $\tau \rightarrow \infty$: *critical dynamical yield stress*

$$\sigma_y = G_{\infty} \gamma_c > 0$$

Flow Curves



[Olsson/Teitel, PRL (2007)]

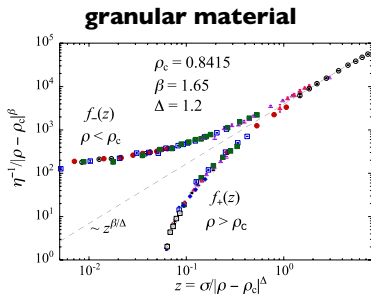


[data: M. Siebenbürger]

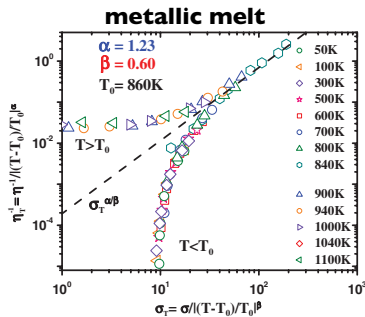
- scenarios for dynamical yielding: colloidal vs. granular
- remember: different protocols, $k_B T$ finite vs. zero



Flow Curves: Scaling Proposal



[Olsson/Teitel, PRL (2007)]



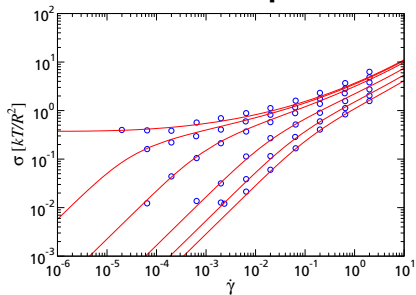
[Guan/Chen/Egami, PRL (2010)]

- (granular) point J as a critical point
- scaling: suggests $\sigma(T \rightarrow T_c) \sim \dot{\gamma}^x \Rightarrow$ hence: $\sigma_y^c = 0$



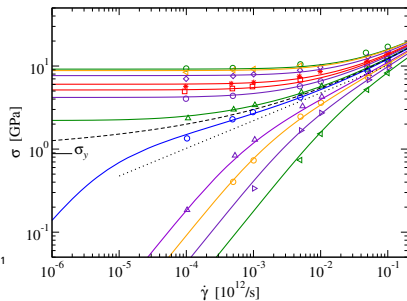
Flow Curves: Finite Yield Stress

Brownian hard spheres



[data: Henrich et al., Phil Trans Roy Soc (2009)]

metallic melt

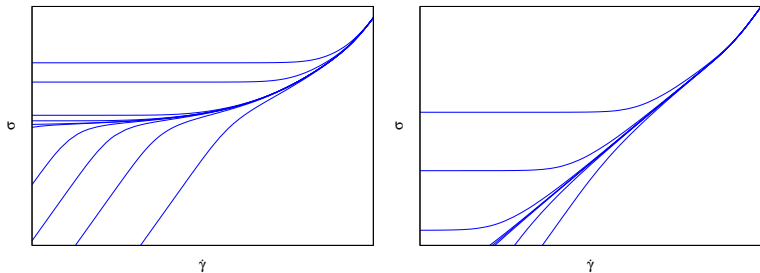


[Guan/Chen/Egami, PRL (2010)]

- fits using schematic model of mode-coupling theory (MCT)
- prediction $\sigma_y^c = \mathcal{O}(0.1 k_B T / R^3)$ – *apparent* power laws



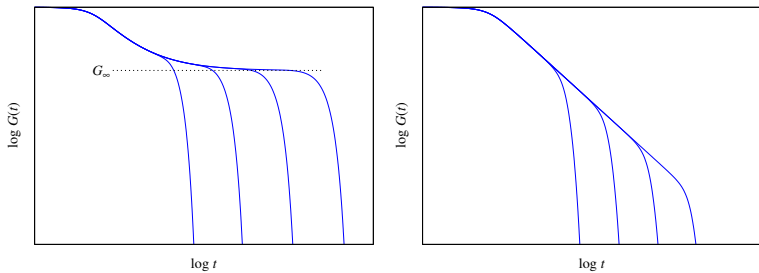
Different Scenarios for Flow Curves



- “discontinuous” vs. “continuous” yield-stress scenario
- different yielding mechanisms: local cages vs. avalanches
- energy densities $k_B T / R^3$ vs. overlap energies



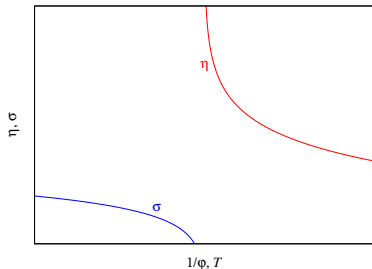
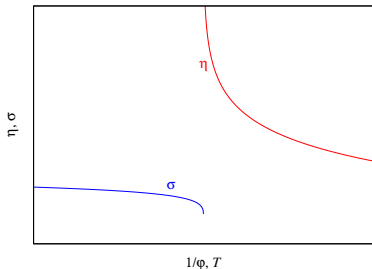
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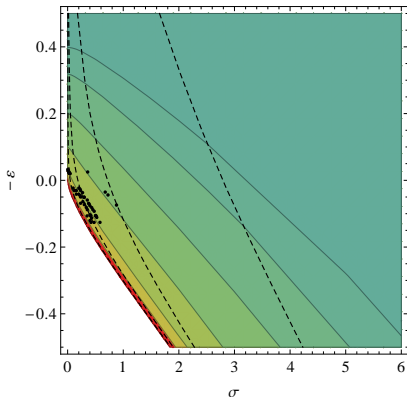
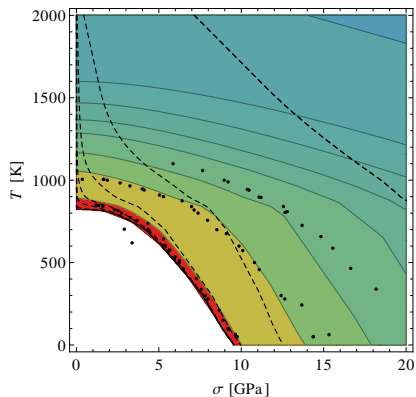
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Jamming Diagram

iso-viscosity lines in the (T, σ) and $(1/\rho, \sigma)$ plane:



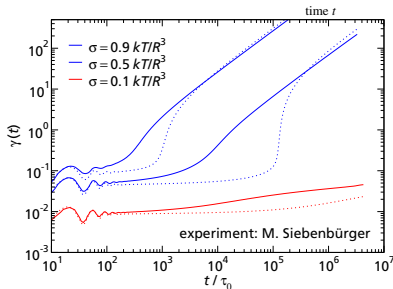
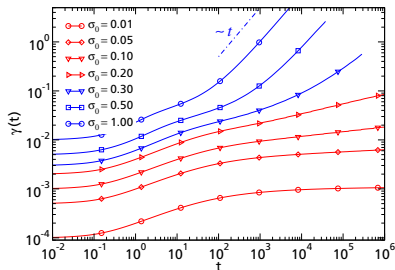
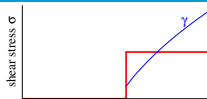
(white: "jammed")



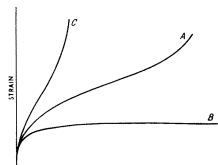
Startup: Creep and Micro-Rheology



- deformation $\gamma(t)$ after sudden step stress

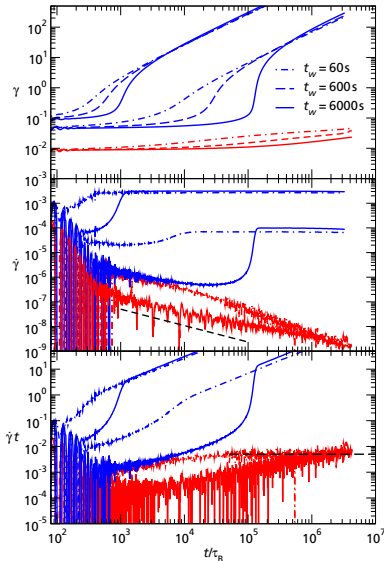


- nonequilibrium transition:
plastic deformation / flow
- static yield stress σ_c
- anomalous flow behavior (creep)?

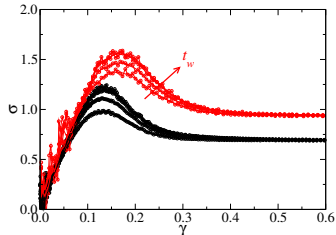


[Cottrell, "The Time Laws of Creep"]

Creep Continued



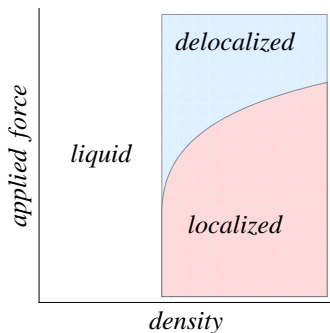
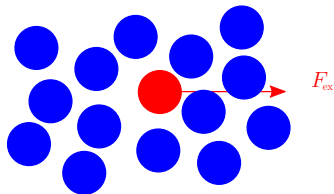
- creep laws (hard matter):
 - logarithmic, $\dot{\gamma}(t)t \sim \text{const.}$
 - Andrade, $\dot{\gamma}(t) \sim t^{-\alpha}$, $\alpha \approx 2/3$
 - secondary, $\dot{\gamma}(t) \sim \text{const.}$
- “viscosity thinning”, $\dot{\gamma}(t) \sim t^{1+x}$
 - related to stress overshoot?
 - aging-time dependent!





Static Yielding: A Force Threshold

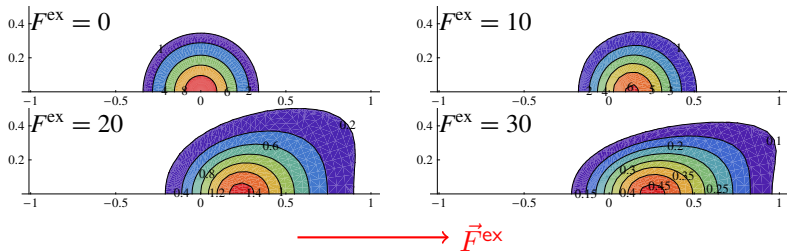
- steady external shear \Rightarrow glass molten (always)
 - steady external force \Rightarrow yielding transition σ_c
 - microscopic analog?
 - yielding of individual “cages” by local external force
- \Rightarrow microrheology





Local Melting of the Glass

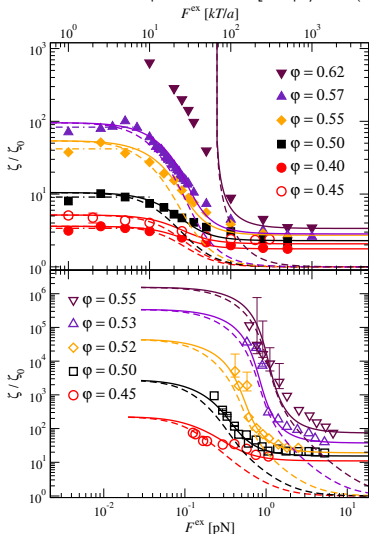
- $F^{\text{ex}} < F_c^{\text{ex}}$: **localized** probe
- distorted probe probability density $\phi^S(\vec{r}, t \rightarrow \infty)$



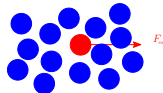
- $F^{\text{ex}} > F_c^{\text{ex}}$: **delocalized** probe
- $F_c^{\text{ex}} \gg k_B T / \sigma$: **cages**, not thermal forces

Microscopic Yielding

sim.: A M Puertas / exp.: Habdas et al. [Europhys Lett (2004)]

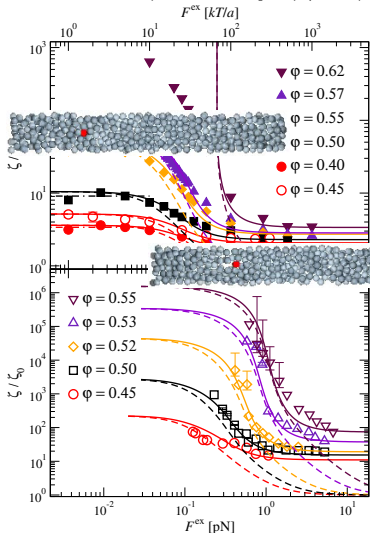


- depinning signature at $F \approx F_c$:
measures typical **cage strength**
- fits: schematic model (MCT)
 - modes $\parallel \vec{F}^{\text{ex}}, \perp \vec{F}^{\text{ex}}$
 - high-force plateau:
fluctuations \perp force
 - strong influence of hydrodynamic interactions
- MCT power laws \curvearrowright
 $\langle v \rangle_\infty \sim (F - F_c)^{1/a-1}$

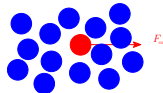


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fluctuations \perp force
 - strong influence of hydrodynamic interactions
- MCT power laws \curvearrowright
 $\langle v \rangle_\infty \sim (F - F_c)^{1/a-1}$





Residual Stresses



Constitutive Equations

common Ansatz: nonlinear *Boltzmann superposition principle*

$$\sigma(t) = \int_{-\infty}^t \gamma_{tt'} \psi(t - t', \gamma_{tt'}) dt' \quad (\text{BKZ})$$

⇒ prediction: **single-step-strain response** contains it all

⇒ $\sigma(\infty) = 0$ whenever $\sum_i \gamma_i = 0$ in n -step-strain

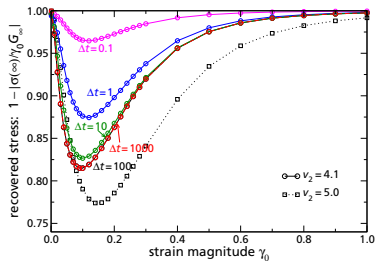
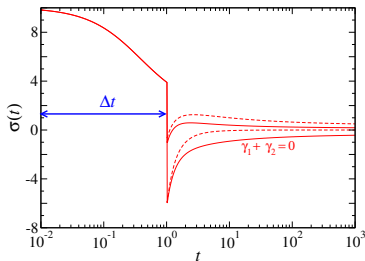
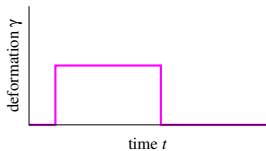
● test: double step strain

● note: MCT does *not* reduce to BKZ form

⇒ prediction of *residual stresses*



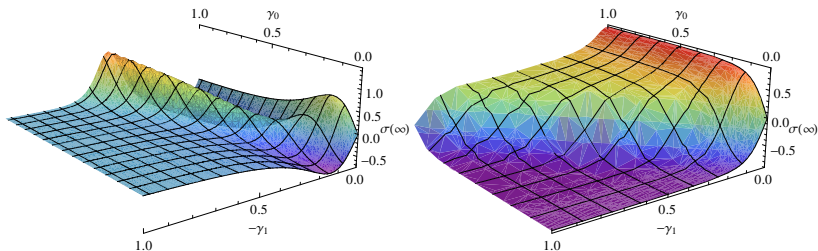
Stress Recovery After Reversing Strain



- glass: finite $\sigma(\infty)$ even for reversed strain
- recovered stress: $< 100\%$ due to **memory effects**



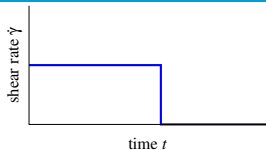
Double Step Strain: Protocol Dependence



- protocol dependence: sudden “affine” deformation vs. strain-rate-ramp
- response: “echo” vs. “wipe-out”



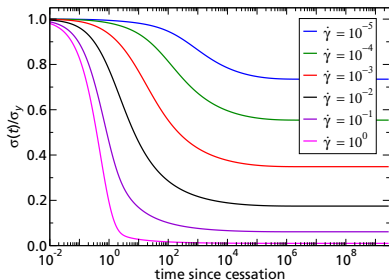
Residual Stresses: Switching off Steady Shear



ss eq

$t_w = 0$ t_{w1} t_{w2}

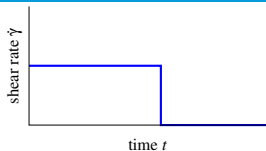
measure decay of shear stress $\sigma(t)$



- glass can sustain finite residual stress $\sigma(\infty)$
- strong shear “forgotten” quicker



Residual Stresses: Switching off Steady Shear

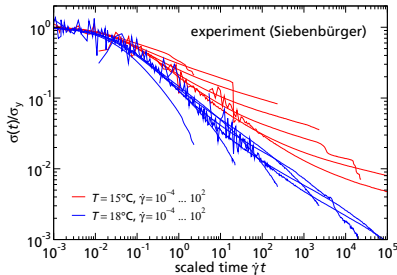
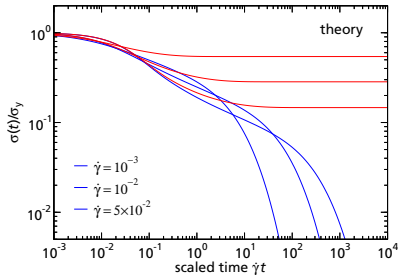


ss

eq



measure decay of shear stress $\sigma(t)$



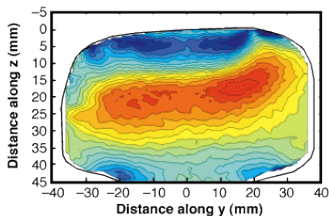
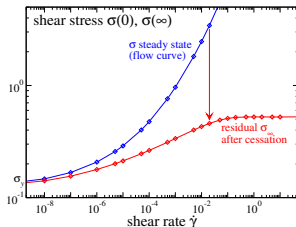
- glass can sustain finite residual stress $\sigma(\infty)$ – transition
- strong shear “forgotten” quicker

Frozen-In Stresses: History Dependence

- glass re-freezes after shear melting
⇒ frozen in stresses: memory of flow history

$$\sigma(t) = \dot{\gamma} \int_{-\infty}^0 dt' G(t > 0, t' < 0; \dot{\gamma})$$

- relevant for applications



residual stress in a worn railway rail



stress
birefringence

[Webster et al., MRS Forum (2002)]

[image: wikipedia.org]



Summary

- dynamical yield stress σ_y
 - goes to zero continuously as liquid is approached? \Rightarrow “granular case”
 - goes to constant? \Rightarrow “Maxwell case” ($G_\infty \neq 0$ in liquid)
- creep
 - analogy colloidal systems – metallic systems
 - “failure time” – overshoot in startup flow
- residual stresses
 - signature of energy dissipation due to relaxation processes
 - not obtained in standard BKZ-type constitutive equations

