A New Perspective on Holographic Entanglement

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1 How should one think about the minimal surface?

In semiclassical gravity, surface areas are related to entropies

Bekenstein-Hawking ['74]:

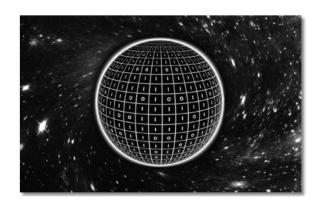
For black hole,

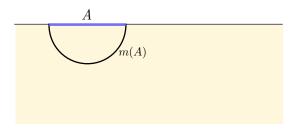
$$S = \frac{1}{4G_{\rm N}} \operatorname{area}(\operatorname{horizon})$$

Why?

Possible answer:

Microstate bits "live" on horizon, 1 bit / 4 Planck areas





Ryu-Takayanagi ['06]: For region in holographic field theory (classical Einstein gravity, static state)

$$S(A) = \frac{1}{4G_N} \operatorname{area}(m(A))$$

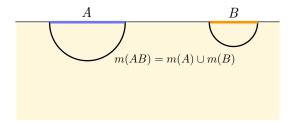
m(A) = bulk minimal surface homologous to A

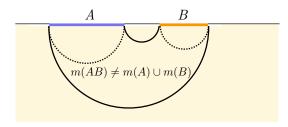
Do microstate bits of A "live" on m(A)?

Unlike horizon, m(A) is not a special place; by choosing A, we can put m(A) almost anywhere

Puzzles:

ullet Under continuous changes in boundary region, minimal surface can jump Example: Union of separated regions A,B





• Information-theoretic quantities are given by differences of areas of surfaces passing through different parts of bulk:

Conditional entropy: H(A|B) = S(AB) - S(B)

 $\text{Mutual information:} \qquad I(A:B) = S(A) + S(B) - S(AB)$

Conditional mutual information: I(A:B|C) = S(AB) + S(BC) - S(ABC) - S(C)

$$H(A|B) = S(AB) - S(B)$$
 $I(A:B) = S(A) + S(B) - S(AB)$

Information-theoretic meaning (heuristically):

Classical: H(A|B) = # of (independent) bits belonging purely to A I(A:B) = # shared with B $\longleftarrow S(AB) \longrightarrow A$ 1011010010111101000110010001100111 B $\longleftarrow H(A|B) \longrightarrow \longleftarrow I(A:B) \longrightarrow A$

Quantum: Entangled (Bell) pair contributes 2 to I(A:B), -1 to H(A|B); can lead to H(A|B) < 0

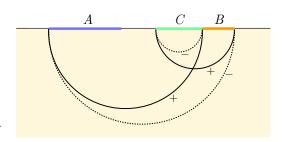
$$= \frac{1}{\sqrt{2}} \left(\begin{vmatrix} 1 \\ 1 \end{vmatrix} \right) + \begin{vmatrix} 0 \\ 0 \end{vmatrix} \right) \qquad A \qquad \underbrace{ \qquad \qquad S(B) \longrightarrow}_{\qquad \qquad \qquad } I(A:B) \longrightarrow \underbrace{ \qquad \qquad }_{\qquad \qquad \qquad } I(A:B) \longrightarrow \underbrace{ \qquad \qquad }_{\qquad \qquad \qquad } I(A:B) \longrightarrow \underbrace{ \qquad \qquad }_{\qquad \qquad \qquad } I(A:B) \longrightarrow \underbrace{ \qquad \qquad }_{\qquad \qquad \qquad } I(A:B) \longrightarrow \underbrace{ \qquad \qquad }_{\qquad \qquad \qquad } I(A:B) \longrightarrow \underbrace{ \qquad \qquad }_{\qquad \qquad \qquad } I(A:B) \longrightarrow \underbrace{ \qquad \qquad }_{\qquad \qquad \qquad } I(A:B) \longrightarrow \underbrace{ \qquad \qquad }_{\qquad \qquad \qquad } I(A:B) \longrightarrow \underbrace{ \qquad \qquad }_{\qquad \qquad \qquad } I(A:B) \longrightarrow \underbrace{ \qquad \qquad }_{\qquad \qquad \qquad } I(A:B) \longrightarrow \underbrace{ \qquad \qquad }_{\qquad \qquad \qquad } I(A:B) \longrightarrow \underbrace{ \qquad \qquad }_{\qquad \qquad \qquad } I(A:B) \longrightarrow \underbrace{ \qquad \qquad }_{\qquad \qquad \qquad } I(A:B) \longrightarrow \underbrace{ \qquad \qquad }_{\qquad \qquad \qquad } I(A:B) \longrightarrow \underbrace{ \qquad \qquad }_{\qquad \qquad \qquad } I(A:B) \longrightarrow \underbrace{ \qquad \qquad }_{\qquad \qquad \qquad } I(A:B) \longrightarrow \underbrace{ \qquad \qquad }_{\qquad \qquad \qquad } I(A:B) \longrightarrow \underbrace{ \qquad \qquad }_{\qquad \qquad \qquad } I(A:B) \longrightarrow \underbrace{ \qquad \qquad }_{\qquad \qquad \qquad } I(A:B) \longrightarrow \underbrace{ \qquad \qquad }_{\qquad \qquad \qquad } I(A:B) \longrightarrow \underbrace{ \qquad \qquad }_{\qquad \qquad } I(A:$$

I(A:B|C) = S(AB) + S(BC) - S(ABC) - S(C) =correlation between A & B, conditioned on C What do differences between areas of surfaces, passing through different parts of bulk, have to do with these measures of information?

• RT obeys strong subadditivity [Headrick-Takayanagi '07]

$$I(A:BC) \ge I(A:C)$$

What does proof (by cutting & gluing minimal surfaces) have to do with information-theoretic meaning of SSA (monotonicity of correlations)?



To try to answer these questions, I will present a new formulation of RT

- Does not refer to minimal surfaces (demoted to a calculational device)
- Suggests a new way to think about the holographic principle & about the connection between spacetime geometry and information

2 Reformulation of RT

Consider a Riemannian manifold with boundary

Flow: vector field v obeying $\nabla \cdot v = 0$, $|v| \le 1$

Think of flow as a set of oriented threads (flow lines) beginning & ending on boundary, transverse density $=|v|\leq 1$

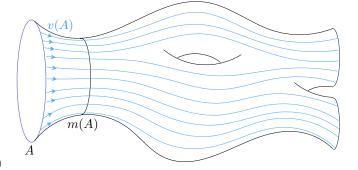
Let A be a subset of boundary

Max flow-min cut theorem (originally on graphs; Riemannian version: [Federer '74, Strang '83, Nozawa '90]):

$$\max_{v} \int_{A} v = \min_{m \sim A} \operatorname{area}(m)$$

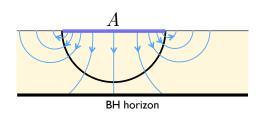
Note:

- $\hbox{ Max flow is highly non-unique} \\ \hbox{ (except on } m(A), \hbox{ where } v=\hbox{ unit normal)} \\ \hbox{ Let } v(A) \hbox{ denote } \textit{any } \max \hbox{ flow}$
- Finding max flow is a convex optimization problem



RT version 2.0:

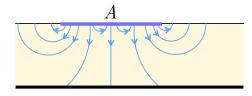
$$S(A) = \max_{v} \int_{A} v$$
 $(4G_{
m N}=1)$ $= \max \# ext{ of threads beginning on } A$



Threads can end on A^c or horizon

Each thread has cross section of 4 Planck areas & is identified with 1 (independent) bit of A

Automatically incorporates homology & global minimization conditions of RT



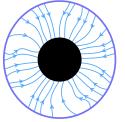
Threads are "floppy": lots of freedom to move them around in bulk & move where they attach to ${\cal A}$

Also lots of room near boundary to add extra threads that begin & end on A (don't contribute to S(A))

Role of minimal surface: bottleneck, where threads are maximally packed, hence counted by area

Holographic principle: entropy \propto area because bits are carried by one-dimensional objects

Bekenstein-Hawking:



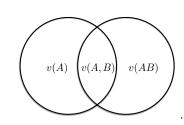
3 Threads & information

Now we address conceptual puzzles with RT raised before

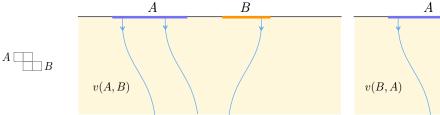
First: even when m(A) jumps, v(A) changes continuously with A

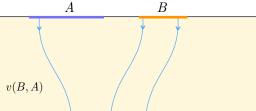
Next, consider two regions A, B

We can maximize flux through A or B, not in general both But we ${\it can}$ always maximize through A and AB (nesting property) Call such a flow v(A,B)



Example 1:
$$S(A) = S(B) = 2$$
, $S(AB) = 3 \Rightarrow I(A : B) = 1$, $H(A|B) = 1$



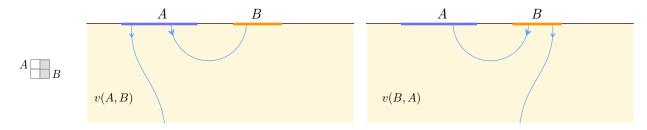


Lesson 1:

- ullet Threads that are stuck on A represent bits unique to A
- ullet Threads that can be moved between $A\ \&\ B$ represent correlated pairs of bits

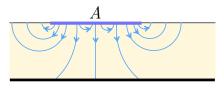
Example 2:
$$S(A) = S(B) = 2$$
, $S(AB) = 1 \Rightarrow I(A : B) = 3$, $H(A|B) = -1 \Rightarrow$ entanglement!

One thread leaving A must go to B, and vice versa



Lesson 2:

ullet Threads that connect $A\ \&\ B$ (switching orientation) represent entangled pairs of bits



Apply lessons to single region:

- \bullet freedom to move beginning points around reflects correlations within A
- \bullet freedom to add threads that begin & end on A reflects entanglement within A

In equations:

$$\begin{split} H(A|B) &= S(AB) - S(B) \\ &= \int_{AB} v(AB) - \int_{B} v(B) \\ &= \int_{AB} v(B,A) - \int_{B} v(B,A) \\ &= \int_{A} v(B,A) \\ &= \min \text{ flux on } A \text{ (maximizing on } AB) \end{split}$$

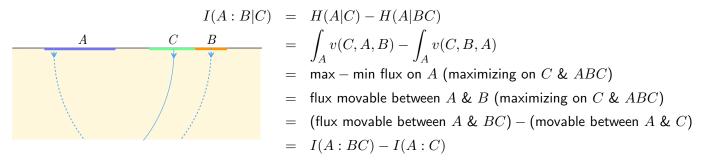
Mutual information:

$$\begin{split} I(A:B) &= S(A) - H(A|B) \\ &= \int_A v(A,B) - \int_A v(B,A) \\ &= \max - \min \text{ flux on } A \text{ (maximizing on } AB) \\ &= \text{ flux movable between } A \text{ and } B \text{ (maximizing on } AB) \end{split}$$

Max flow can be defined even when flux is infinite: flow that cannot be augmented Regulator-free definition of mutual information:

$$I(A:B) = \int_{A} (v(A,B) - v(B,A))$$

Conditional mutual information:



Strong subadditivity $(I(A:B|C) \ge 0)$ is clear

In each case, clear connection to information-theoretic meaning of quantity/property

Open problem: Use flows to prove "monogamy of mutual information" property of holographic EEs [Hayden-Headrick-Maloney '12]:

$$I(A:BC) \ge I(A:B) + I(A:C)$$

(and generalizations to more parties [Bao et al '15])

Flow-based proof may illuminate the information-theoretic meaning of these inequalities

4 Extensions

4.1 Emergent geometry

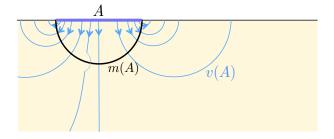
Metric \longrightarrow Set of thread configurations allowed by $\nabla \cdot v = 0$, $|v| \leq 1$ Set of thread configurations \longrightarrow metric: Given set $\{w\}$ of closed (d-1)-forms, find $g_{\mu\nu}$ with pointwise smallest det such that $|w| \leq 1 \ \forall w$

4.2 Quantum corrections

Faulkner-Lewkowycz-Maldacena ['13]:

Quantum (order $G_{
m N}^0$) correction to RT is from entanglement of bulk fields

May be reproduced by allowing threads to jump from one point to another (or tunnel through microscopic wormholes, à la ER = EPR [Maldacena-Susskind '13])



4.3 Covariant bit threads

With Veronika Hubeny (to appear)

Hubeny-Rangamani-Takayanagi ['07] covariant entanglement entropy formula:

$$S(A) = area(m(A))$$

m(A) = minimal extremal surface homologous to A

Need generalization of max flow-min cut theorem to Lorentzian setting

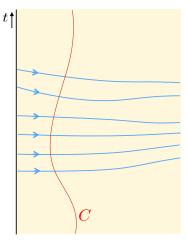
Define a flow as a vector field v (in full Lorentzian spacetime) obeying

- $\bullet \ \nabla \cdot v = 0$
- no flux into or out of singularities
- integrated norm bound: \forall timelike curve C,

$$\int_C ds \, |v_\perp| \le 1 \qquad (v_\perp = \text{projection of } v \text{ orthogonal to } C)$$

Any observer counts, over her lifetime, total of at most 1 thread / 4 Planck areas

(Can be expressed as local constraint)



Theorem (assuming NEC, using results of Wall ['12] & Headrick-Hubeny-Lawrence-Rangamani ['14]):

$$\max_{v} \int_{D(A)} v = \operatorname{area}(m(A)) \qquad \qquad D(A) = \text{boundary causal domain of } A$$

Linearizes problem of finding extremal surface area

HRT version 2.0:

$$S(A) = \max_{v} \int_{D(A)} v$$

To maximize flux, threads seek out m(A), automatically confining themselves to entanglement wedge

Threads can lie on common Cauchy slice (equivalent to Wall's ['12] maximin by standard max flow-min cut) or spread out in time

