

High-dimensional statistics and
large-SpeNta expansion

OR

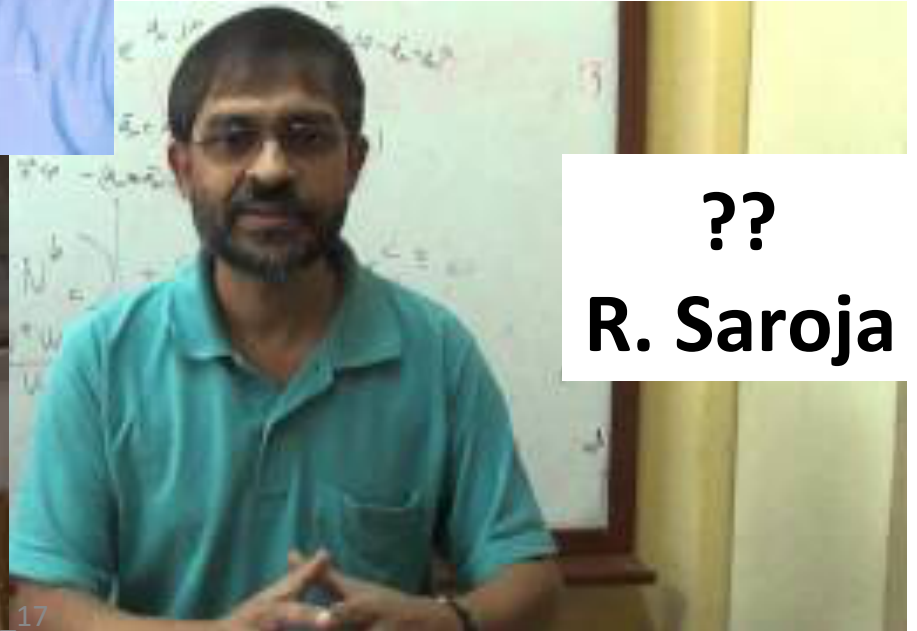
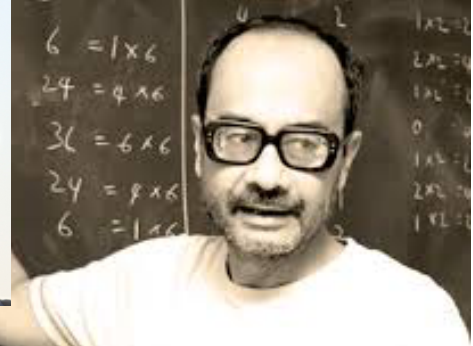
One can go only so far from home

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Rutgers University

TIFR String Excitations: 1987-88



??
R. Saroja

Fall 1989



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NEW CRITICAL BEHAVIOR IN $d = 0$ LARGE- N MATRIX MODELS

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Received 9 January 1990

The non-perturbative formulation of 2-dimensional quantum gravity in terms of the large- N limit of matrix models is studied to include the effects of higher order curvature terms. This leads to matrix models whose potential contains a symmetry breaking term of the form $\text{Tr } \phi A \phi A$, where A is a given fixed matrix. This is studied in $d = 0$ dimensions and effectively induces additional terms of the form $(\text{Tr } \phi^4)^2$ in the one matrix potential. An exact solution to leading order of the potential $V(\phi) = 1/2 \text{Tr } \phi^2 + g/N \text{Tr } \phi^4 + g'/N^2 (\text{Tr } \phi^2)^2$ is presented leading to 3 phases: $\gamma = -1/2$ (smooth surfaces), $\gamma = 1/2$ (branched polymer) and $\gamma = 1/3$ (intermediate phase). Including a $\text{Tr } \phi^6$ term in the potential gives rise to an additional phase with $\gamma = 1/4$. It is conjectured that for the general polynomial potential there are phases with $\gamma = 1/n$, $n = 2, 3, \dots$. The $\gamma > 0$ phases may correspond to $c > 1$ matter coupled to 2-dimensional gravity.

1. Introduction

It is generally believed that many fundamental issues in string theory are non-perturbative in nature. In this respect the standard first quantized approach seems to be of little use since it is always restricted to two-dimensional surfaces and hence to the perturbation expansion. Nevertheless, it has been realized recently that a deeper understanding of the integration over two-dimensional metrics in the first-

, USA

clidean quantum
surfaces of

Where was Spenta? (late eighties)



Spring 1991



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CLASSICAL SOLUTIONS OF 2-DIMENSIONAL STRING THEORY

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We present an exact one-parameter family of solutions to the classical graviton-dilaton system in two dimensions. The solution can be identified as a black hole. We present the solution both in a Schwarzschild-like gauge and in the target space conformal gauge. We discuss possible relations with matrix models.

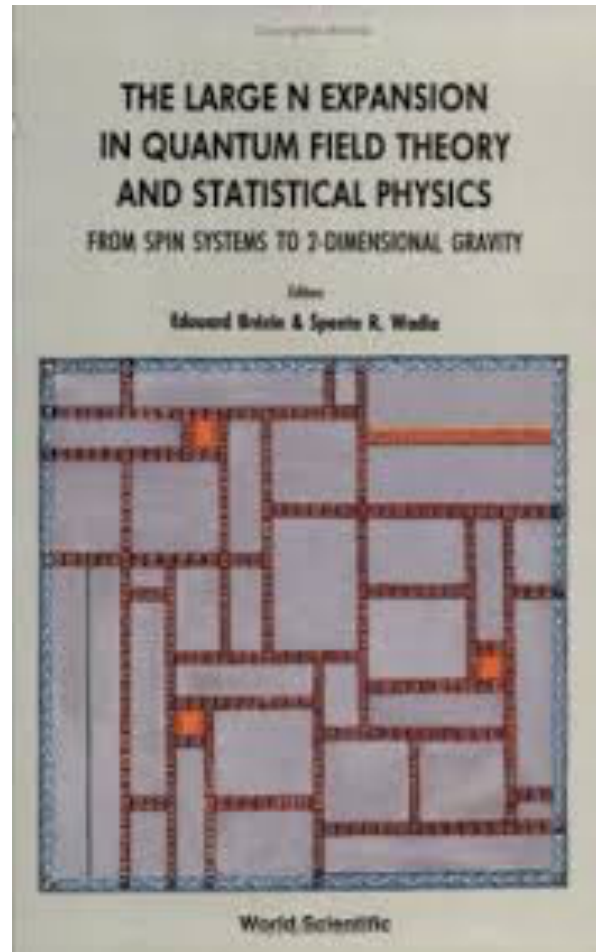
1. Introduction

Recently there has been considerable interest in the problem of quantized 2-dimensional gravity coupled to a scalar field, which in the absence of gravity would be a $c = 1$ system. There are presently two approaches to this problem — the matrix model^{1,3-9} and the continuum Liouville theory.¹⁰⁻¹² There are encouraging results which indicate that the weak coupling tree level S -matrices computed in the matrix model can be obtained by Liouville theory calculation.⁸ One of the interesting results of the continuum Liouville theory is that this system has a 2-dimensional target space, in that the general couplings are functions of two variables one of them being the conformal mode of the 2-dimensional metric.² The reasoning that led to this proposal was similar to that which leads to the σ -model approach to critical

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²By a Liouville theory here we mean a string in 2-dimensional target space with the following backgrounds: $G_{\mu\nu} = \eta_{\mu\nu}$ (metric), $\Phi = Q\eta/2$ (dilaton) and $T = a \exp(Q\eta/2)$ (tachyon), where η is the Liouville mode. However, the 2-D gravity theory coupled to matter can admit of more general backgrounds having more complicated dependences on η . The reasoning that led to this proposal² is similar to the one that led to the σ -model approach in critical string theory.

Large N Expansion



Recurring theme we heard today

- Emphasis on physical intuition
- Emphasis on unity of physics

Large N expansion as a fixed point

Kondo Physics:

Sengupta, Georges (1995)

Parcollet, Georges, Kotliar, Sengupta (1998)

Singular Value Distribution of Random Matrices:

Sengupta, Mitra (1999)

Information Theory of MIMO Systems:

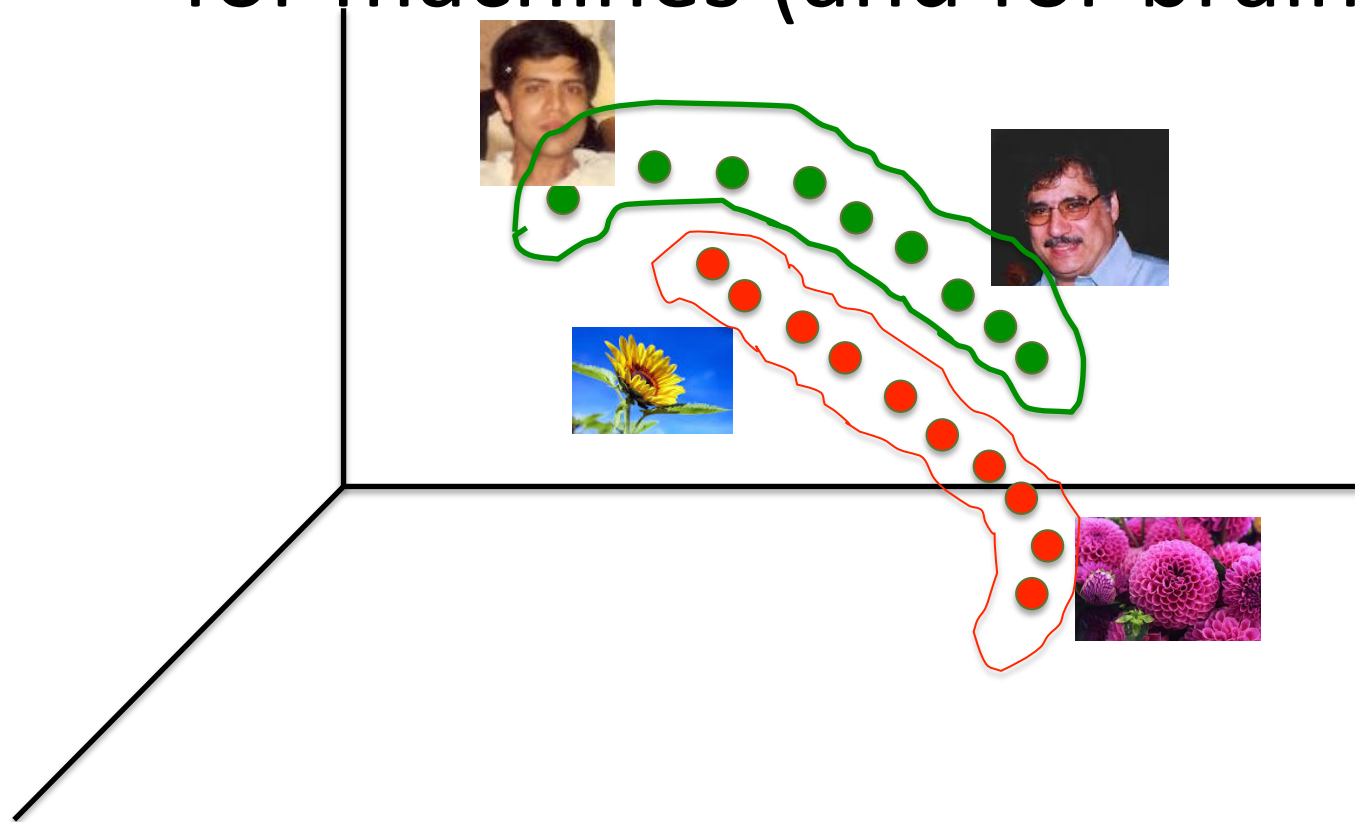
Moustakas, Baranger, Balents, Sengupta, Simon (2000)

Moustakas, Simon, Sengupta (2003)

Algorithmic Phase Transitions in Compressed Sensing:

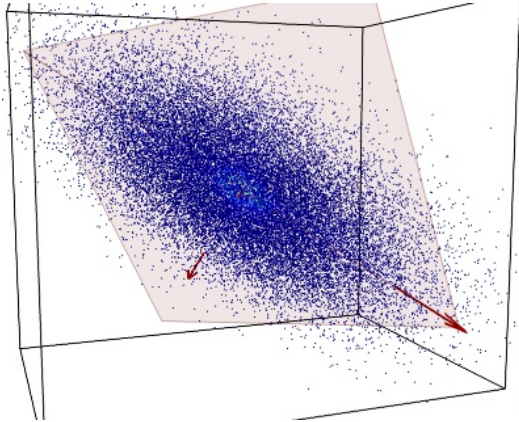
Ramezanali, Mitra, Sengupta (2015)

High dimensional pattern recognition: for machines (and for brains)



Curse of dimensionality

Dimension Reduction 101: Principal Components Analysis (PCA)



Diagonalize covariance matrix with
Element given by:

$$C_{ij} = \sum_a (X_{ia} - \bar{X}_i)(X_{ja} - \bar{X}_j)$$

PCA->

Singular Value Decomposition (SVD)

If we define $\hat{X}_{ia} = X_{ia} - \bar{X}_i$ then $C = \hat{X}\hat{X}^T$

Diagonalizing covariance matrix C is the same thing as doing a singular value decomposition of deviations \hat{X} , and ordering singular vectors according to the size of singular values.

$$\begin{pmatrix} \hat{X} \\ \begin{pmatrix} x_{11} & x_{12} & \dots & x_{1n} \\ x_{21} & x_{22} & \dots & \\ \vdots & \vdots & \ddots & \\ x_{m1} & & & x_{mn} \end{pmatrix} \\ m \times n \end{pmatrix} \approx \begin{pmatrix} U \\ \begin{pmatrix} u_{11} & \dots & u_{1r} \\ \vdots & \ddots & \\ u_{m1} & & u_{mr} \end{pmatrix} \\ m \times r \end{pmatrix} \begin{pmatrix} S \\ \begin{pmatrix} s_{11} & 0 & \dots \\ 0 & \ddots & \\ \vdots & & s_{rr} \end{pmatrix} \\ r \times r \end{pmatrix} \begin{pmatrix} V^T \\ \begin{pmatrix} v_{11} & \dots & v_{1n} \\ \vdots & \ddots & \\ v_{r1} & & v_{rn} \end{pmatrix} \\ r \times n \end{pmatrix}$$

SVD for Denoising/Compressing Images

$$X_{\text{recon}} = \bar{X} + US_K V^\dagger$$



Figure 2: 1st, 50th, and 120th image approximations of the image for $\sigma = 0\%$



Figure 3: 1st, 14th, and 120th image approximations of the image for $\sigma = 5\%$



Figure 4: 1st, 9th, and 120th image approximations of the image for $\sigma = 10\%$



Figure 5: 1st, 5th, and 120th image approximations of the image for $\sigma = 15\%$

$$S_K = \begin{pmatrix} s_{11} & \dots & \dots & 0 & \dots & \dots & 0 \\ \vdots & \ddots & \vdots & \vdots & \dots & \vdots & \vdots \\ 0 & \dots & s_{KK} & \vdots & \dots & \vdots & \vdots \\ 0 & \dots & \dots & 0 & \dots & 0 & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots & \vdots \\ 0 & \dots & \dots & 0 & \dots & 0 & 0 \end{pmatrix}$$

Q: What K is optimal?

A: Departure from singular values of pure noise.

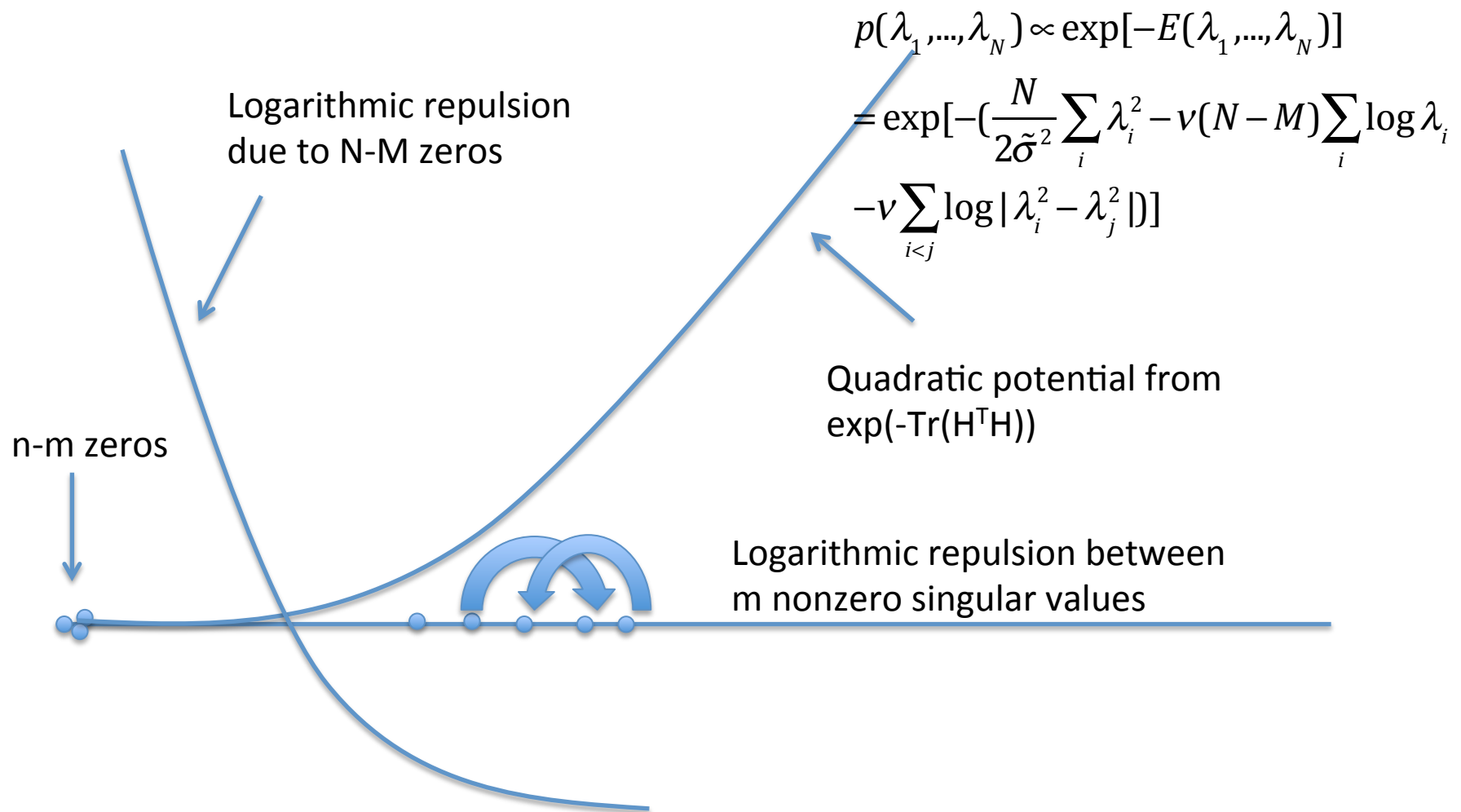
From Poon, Ng and Sridharan, 2009

Describe the distribution in a different coordinate system

$$H = U \Lambda V^\dagger = U \begin{pmatrix} \lambda_1 & 0 & \dots & 0 & 0 & \dots & 0 \\ 0 & \lambda_2 & \dots & 0 & 0 & \dots & 0 \\ \vdots & \vdots & \ddots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \dots & \lambda_M & 0 & \dots & 0 \end{pmatrix} V^\dagger$$

$$\begin{aligned} P(H) \prod_{ia} dH_{ia} &= C \exp\left[-\frac{N}{2\tilde{\sigma}^2} \text{tr}(H^2)\right] \prod_{ia} dH_{ia} \\ &\rightarrow C e^{-\frac{N}{2\tilde{\sigma}^2} \sum_i \lambda_i^2} \prod_{i < j} |\lambda_i^2 - \lambda_j^2|^\nu \prod_i \lambda_i^{\nu(N-M+1)-1} \left(\prod_i d\lambda_i \right) d\mu(U) d\mu(V) \end{aligned}$$

Coulomb of Singular Values



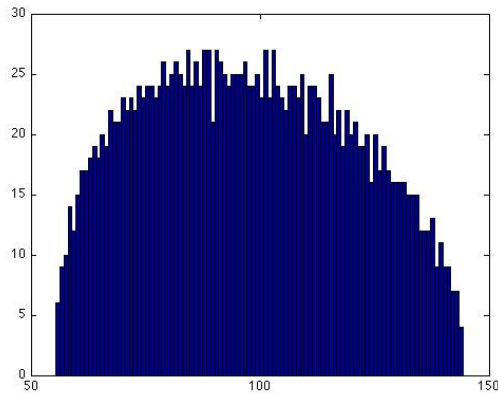
SVD of 'Just Noise'

$$\rho(\lambda) = \frac{a}{\lambda} \sqrt{(\lambda_{\max}^2 - \lambda^2)(\lambda^2 - \lambda_{\min}^2)}$$

where

$$\lambda_{\min, \max} = \sigma \sqrt{(M+N) \mp 2\sqrt{MN}}$$

Marcenko
And
Pastur,
1967



Histogram of singular values
For a 2000X10000 random matrix

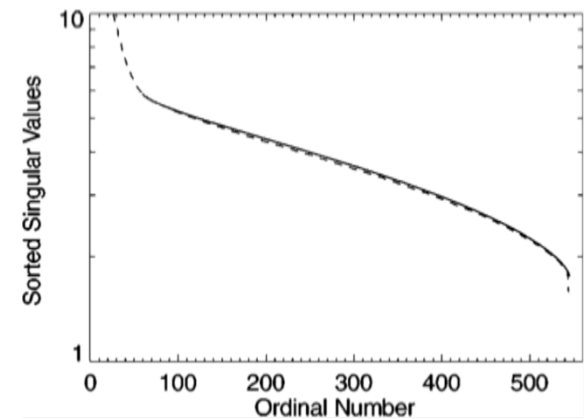
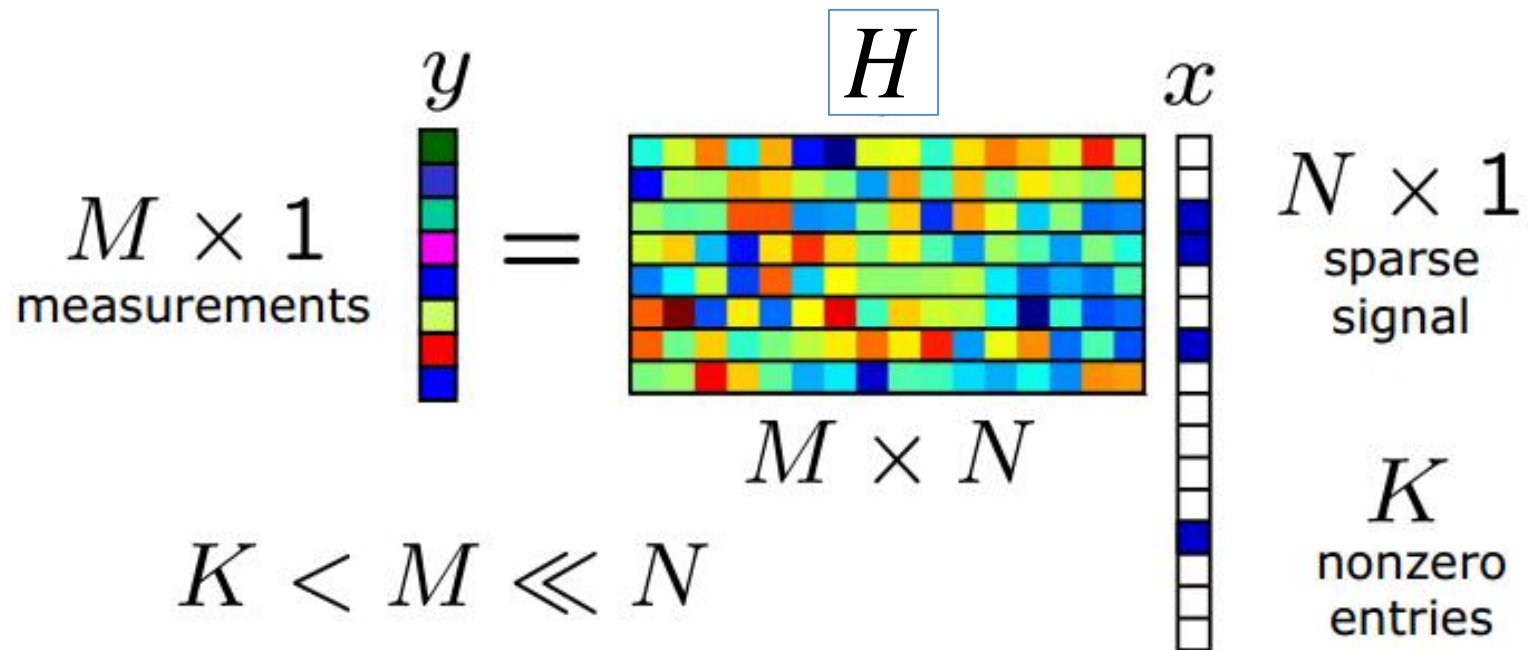


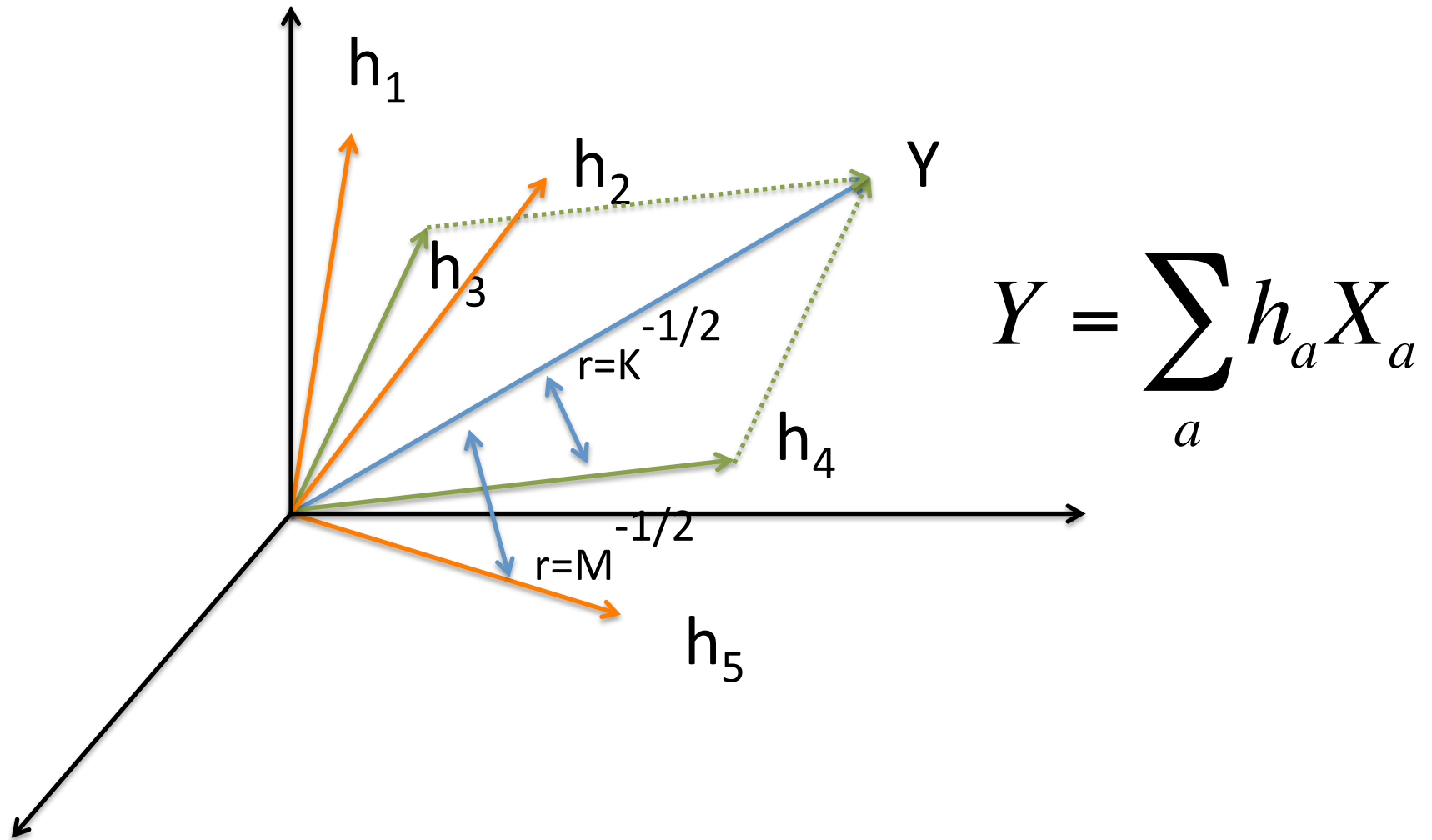
FIG. 3. Comparison of singular values from a SVD of an fMRI data set with the theoretical formula for a noise only matrix.

Sengupta and Mitra,
PRE 1999

Compressed Sensing



Explaining Data



$$Y = \sum_a h_a X_a$$

Convex relaxation

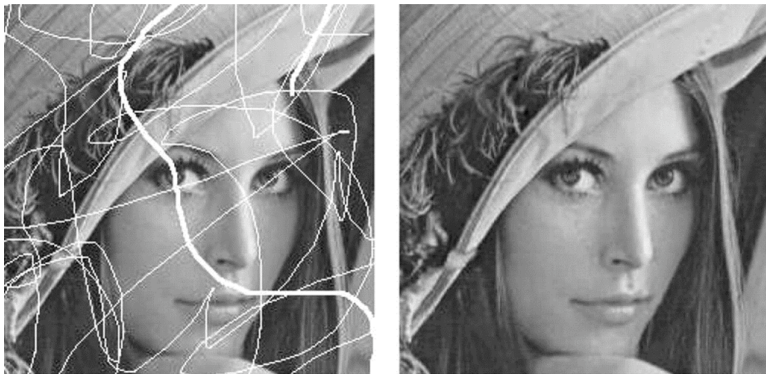
Compressed sensing:

$y \approx Wx_0$, with $y \in R^M, x_0 \in R^N, M < N$,

but $\|x_0\|_0 \leq K < M$.

Find $\hat{x} = \arg \min_x \|x\|_1$ subject to $y = Wx$.

Is $\hat{x} \approx x_0$ when W is a random matrix?

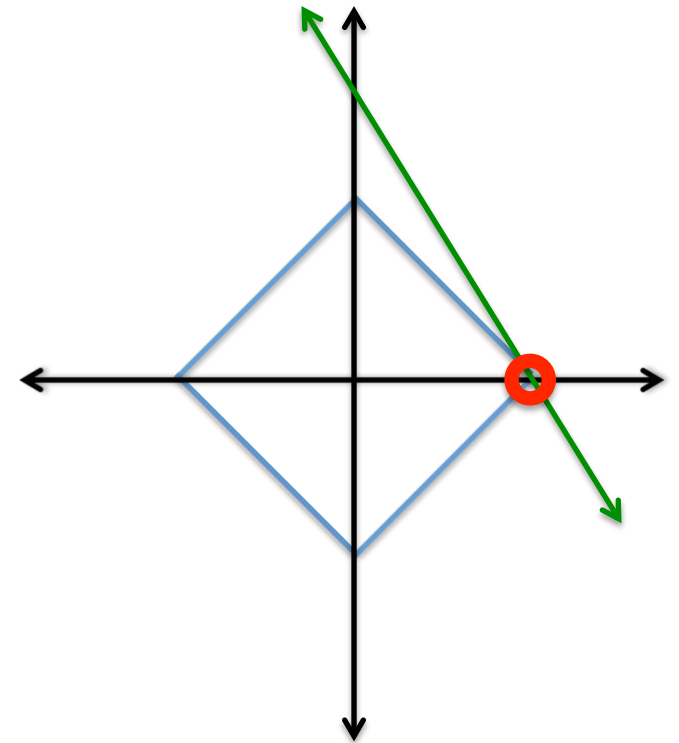


Sparsity enforced by l_1 norm

L0-L1 relaxation:

Minimize $\|X\|_1$

Subject to $Y = HX$



Works of Candes, Tao
Donoho, from around 2004

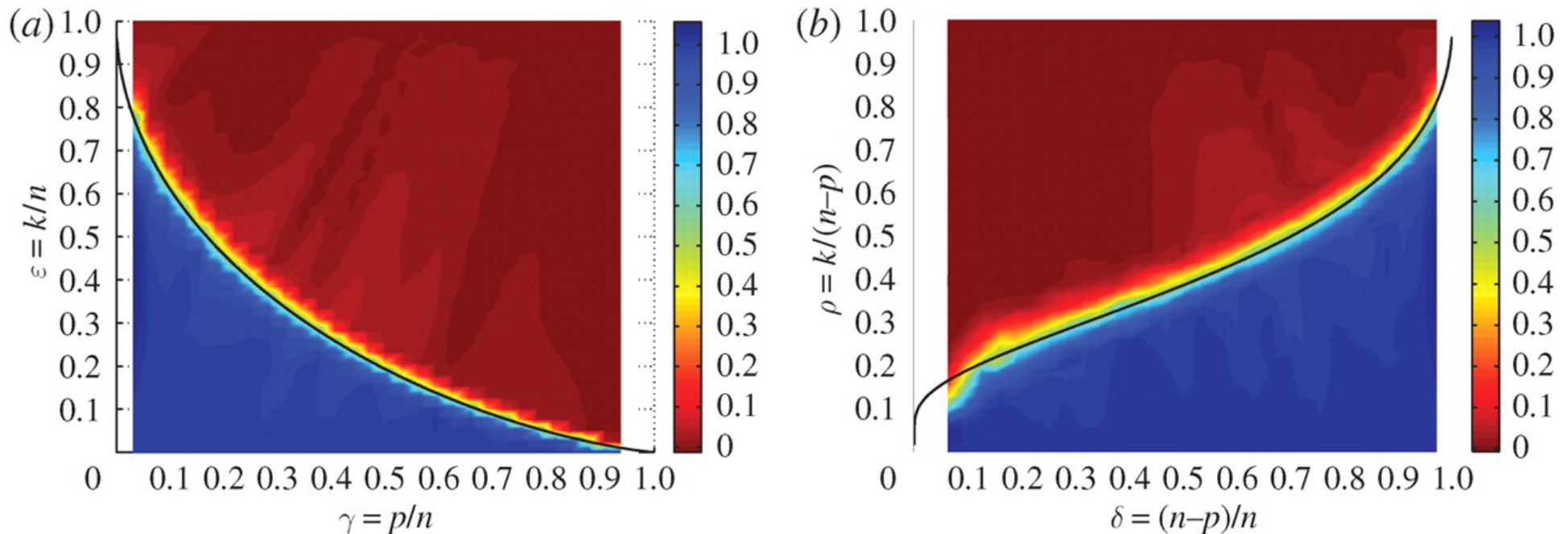
Large Dimensions

The transition happens with

$N, M, K \rightarrow \infty$, holding

$M / N = \alpha, K / N = \rho$ fixed.

Phase Transition in Compressed Sensing by L0-L1 Relaxation

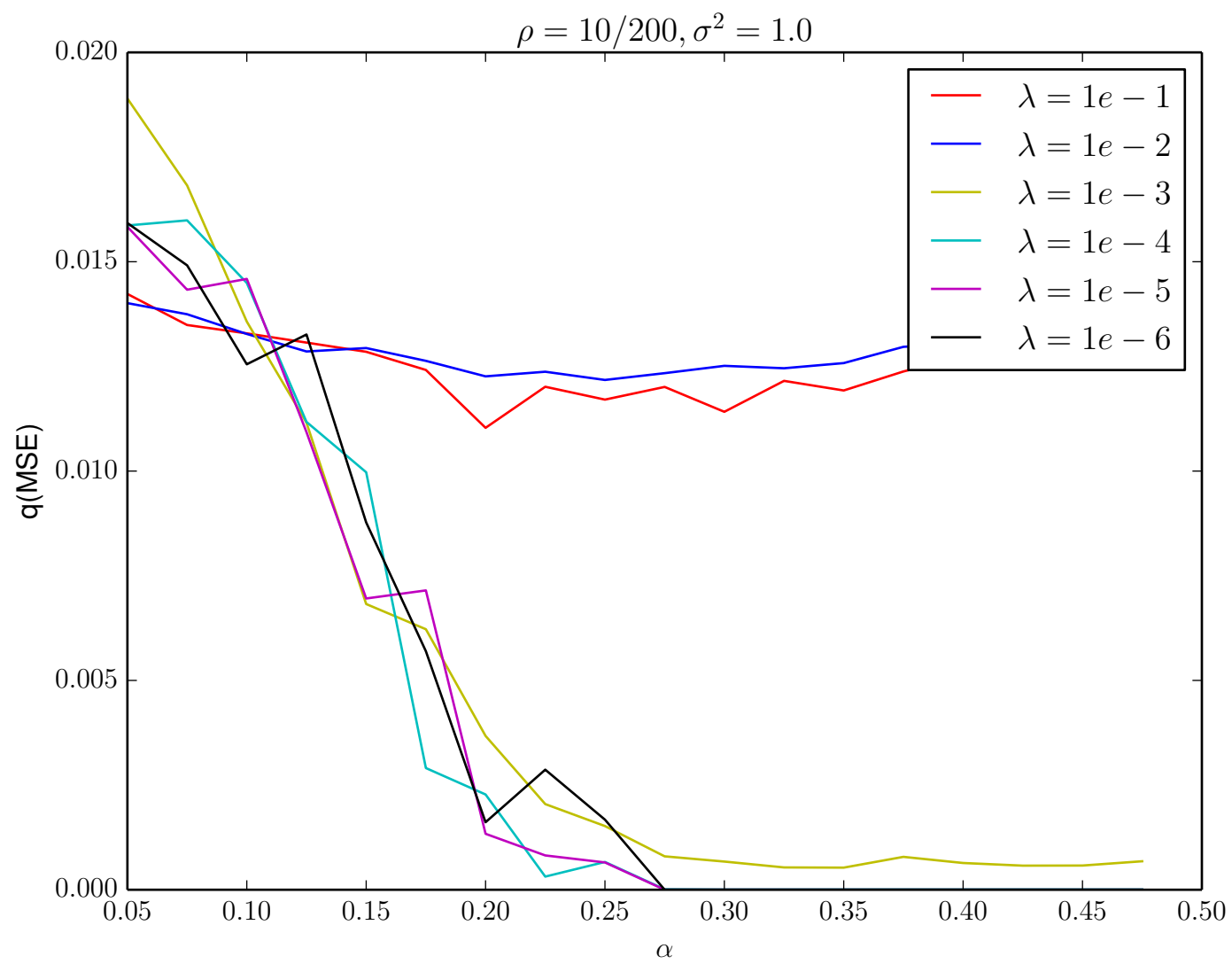


Donoho and Tanner, Phil. Trans. Roy. Soc 2009

Replica Method Analysis: Kabashima

Wadayama and Tanaka, 2009

Ganguly and Sompolinsky, 2010



Random-like deep network

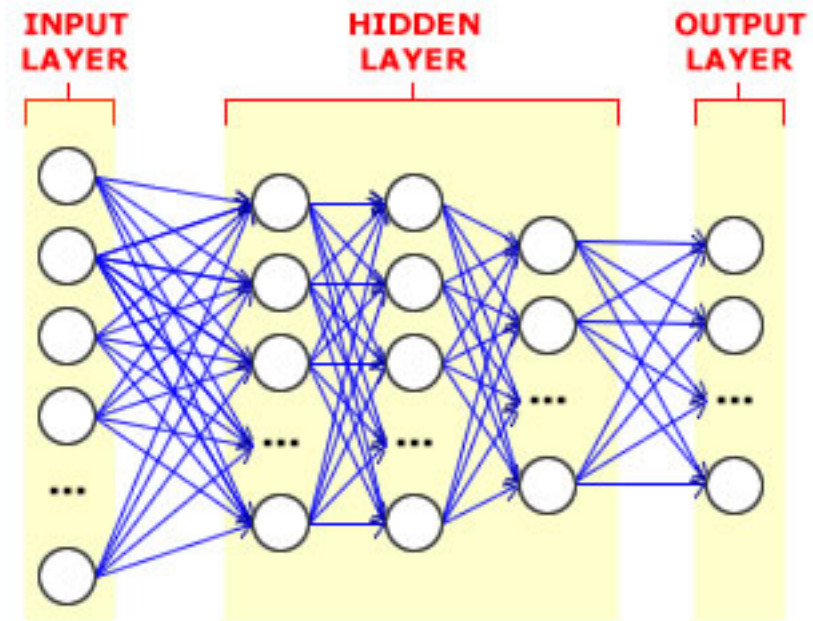
Deep nets

$$h_{M_1 \times 1}^{(1)} \approx \sigma(W_{M_1 \times N}^{(1)\dagger} x_{N \times 1} + b_1),$$

$$h_{M_2 \times 1}^{(2)} \approx \sigma(W_{M_2 \times M_1}^{(2)\dagger} h_{M_1 \times 1}^{(1)} + b_2)$$

...

$$h_{M_d \times 1}^{(d)} \approx \sigma(W_{M_d \times M_1}^{(d)\dagger} h_{M_1 \times 1}^{(d-1)} + b_d)$$



From <http://www.andrewsnoke.com>

'Pre-image' of the final layer activation

- Fully-connected layers of deep nets

$$h^{(1)} \approx \sigma(W^{(1)\dagger}x + b_1),$$

$$h^{(2)} \approx \sigma(W^{(2)\dagger}h^{(1)} + b_2)$$

...

$$h^{(d)} \approx \sigma(W^{(d)\dagger}h^{(d-1)} + b_d)$$

Call this map Φ_f :

$$h^{(d)} = \Phi_f(x)$$



What is the preimage of Φ_f ?

Reconstruction :

$$\tilde{x} = \arg \min_x l(\Phi_f(x), h^{(d)}) + R(x)$$



Mahendran and Vedaldi 2015

Alternative ‘generative’ model

Alternative: run the net backwards
randomly zeroing some coordinates

Arora, Liang and Ma 2016

$$\tilde{h}^{(d-1)} \approx \tau(W^{(d)} h^{(d)}) \odot n_{d-1}^{drop}$$

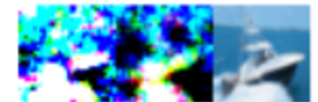
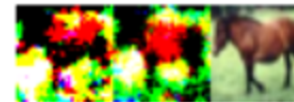
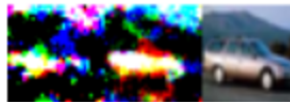
$$\tilde{h}^{(d-2)} \approx \tau(W^{(d-1)} \tilde{h}^{(d-1)}) \odot n_{d-2}^{drop}$$

...

$$\tilde{x} \approx \tau(W^{(1)} \tilde{h}^{(1)}) \odot n_0^{drop},$$

Define a distribution:

$$P(\tilde{x} | h^{(d)})$$



Are deep nets reversible when W 's are random?

Namely, is $h^{(d)} \approx \Phi_f(\tilde{x})$?

Reversibility in One-layer 'Deep' Net

Run the net backwards

randomly zeroing some coordinates

$$\tilde{x} = \gamma \text{ReLU}(Wh) \odot n^{drop}$$

$$\tilde{h} = \text{ReLU}(W^\dagger \tilde{x} + b)$$

Compute joint distribution $P(\tilde{h}_i, h_i)$

Turns out that $W^\dagger \tilde{x} + b$ is nearly Gaussian.

Use that to show $\tilde{h}_i \approx h_i$.

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