

PHYSICS-INSPIRED ALGORITHMS

Nisheeth Vishnoi

Yale

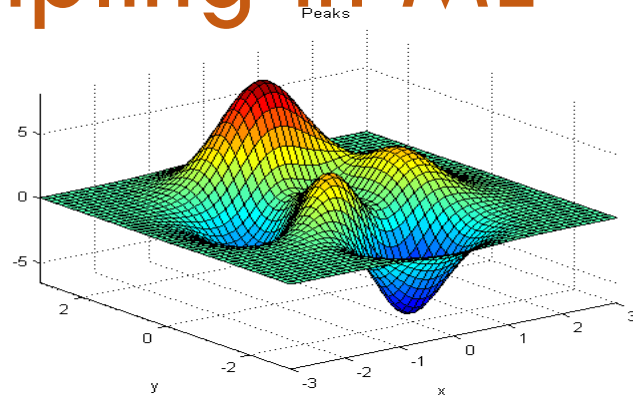
Optimization and Sampling in ML

Given access to $f: \mathbb{R}^d \rightarrow \mathbb{R}$

Optimize $\min_{\theta} f(\theta)$

Sample θ with prob. $\propto e^{-f(\theta)}$

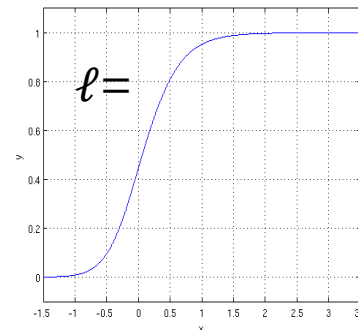
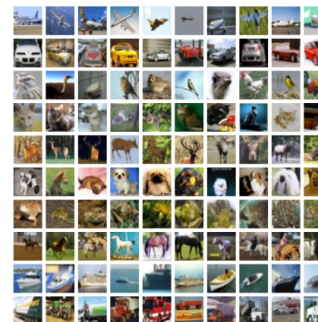
Typically, harder than optimization – but **robust**



Availability of **large, real-world datasets** has given rise to complex objective functions in high dimensions

$$f(\theta) = -f(\theta) = -\sum_i y_i \log \ell(\theta^T x_i) + (1 - y_i) \log \ell(-\theta^T x_i)$$

airplane
automobile
bird
cat
deer
dog
frog
horse
ship
truck



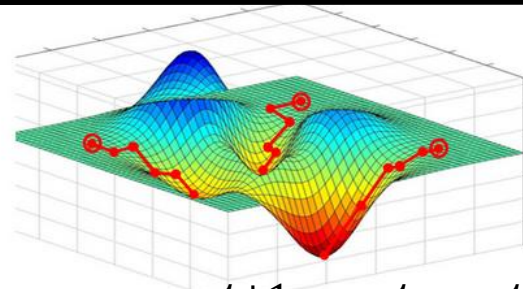
Two facets:

1. Develop methods

(associate a physical meaning and search for the right equations of motion)

2. Prove guarantees, tune parameters

(search for potential functions, “beyond worst case” assumptions on data)



$$\theta^{t+1} = \theta^t + \eta^t G_f(\theta^t)$$

Physics viewpoint has been helpful in both!



HAMILTONIAN DYNAMICS & SAMPLING from CONTINUOUS DISTRIBUTIONS

Sampling from Continuous Distributions

Given access to $f: \mathbb{R}^d \rightarrow \mathbb{R}$

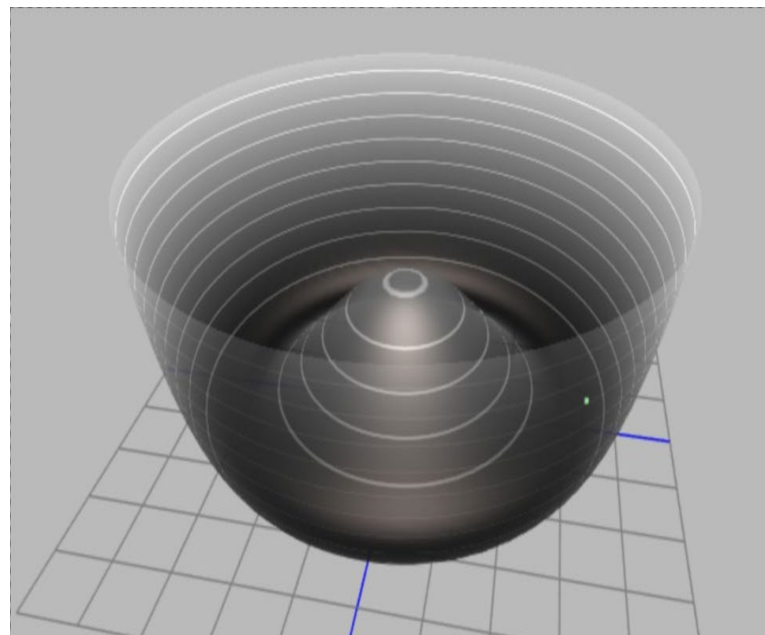
Sample θ with prob. $\pi(\theta) \propto e^{-f(\theta)}$

*Statistics, TCS, Optimization (vis annealing),
Bayesian inference, Molecular dynamics ..*

Iterative methods: MCMC+Metropolis

Propose: $\theta^{k+1} = \theta^k + G_f(\theta^k)$

Accept/Reject



Number of gradient (or function) evaluations to sample from smooth, strongly logconcave π (for smoothness/convexity = $\Theta(1)$):

- Random Walk Metropolis: d^2 [Gelman et al. '97]
- Unadjusted Langevin: d [Durmus, Moulines, '16]
- Underdamped Langevin: $d^{1/2}$ [Cheng et al. '17]

Beyond $d^{1/2}$?

Hamiltonian Monte Carlo

[Duane et al. '87] **No Accept/Reject step!**

Define: $H(\theta, v) = f(\theta) + \frac{1}{2} \|v\|^2$

In step i , sample $V_i \sim N(0, I_d)$

Obtain Θ_{i+1} by simulating **Hamiltonian Dynamics** starting at (Θ_i, V_i) for time T

Fact: Invariant distribution $\propto e^{-f(\theta)} e^{-\frac{1}{2}\|v\|^2}$



$$\frac{d\theta(t)}{dt} = v(t)$$

$$\frac{dv(t)}{dt} = -\nabla f(\theta(t))$$

2nd-order Leapfrog integrator

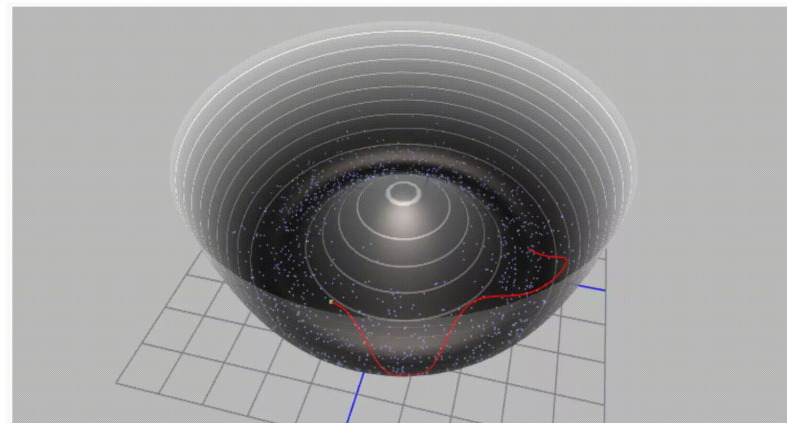
Let $(\theta_0, v_0) = (\Theta_i, V_i)$

For $j = 0, \dots, \frac{T}{\eta}-1$, do

$$\theta_{j+1} = \theta_j + \eta v_j - \frac{1}{2} \eta^2 \nabla f(\theta_j)$$

$$v_{j+1} = v_j - \eta \nabla f(\theta_j) - \frac{1}{2} \eta^2 \frac{\nabla f(\theta_{j+1}) - \nabla f(\theta_j)}{\eta}$$

$$\Theta_{i+1} = \theta_{\frac{T}{\eta}}$$

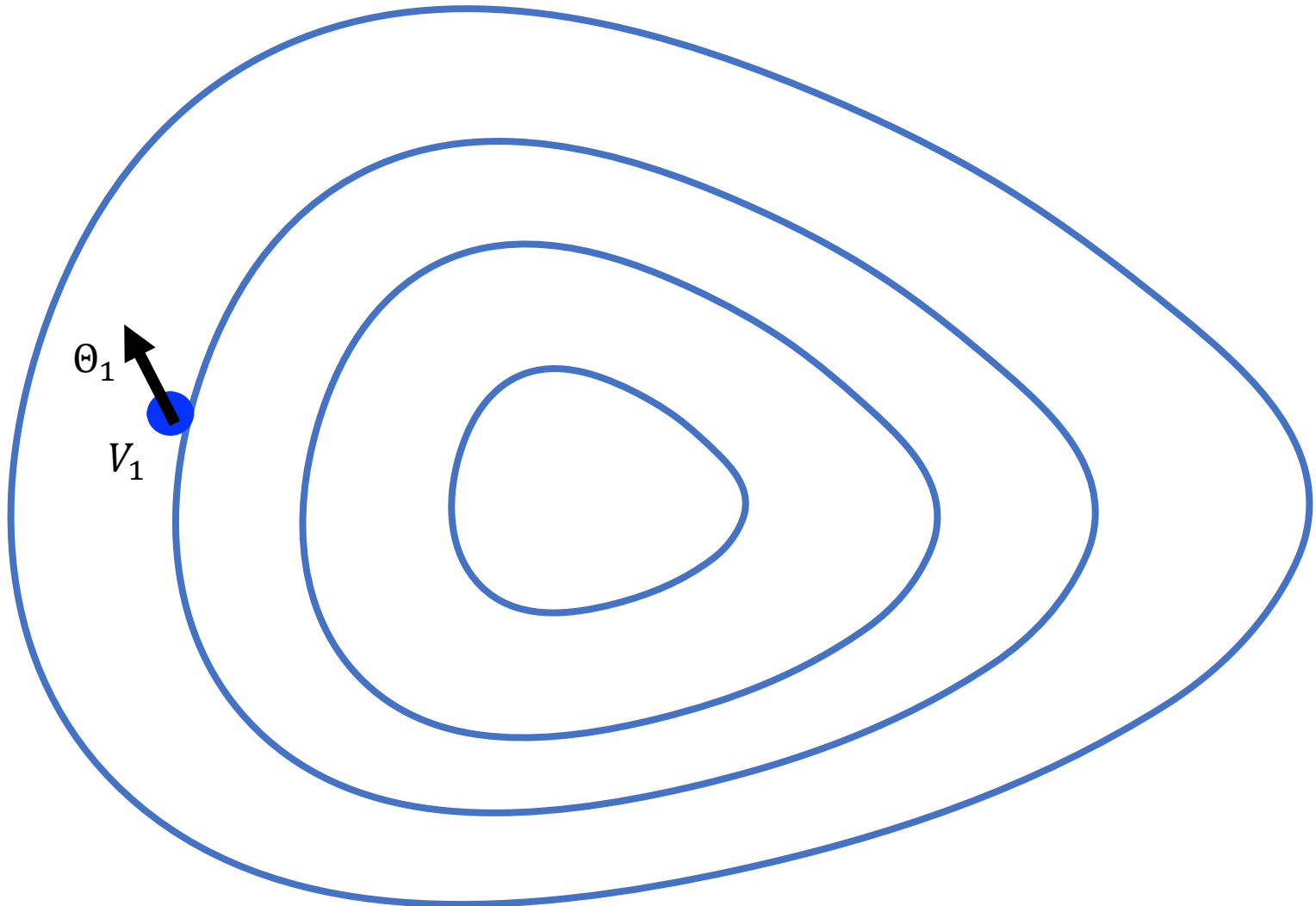


Widely deployed in practice – convergence bounds/tuning parameters?

(Informal) Conjecture: [Creutz, 1988]

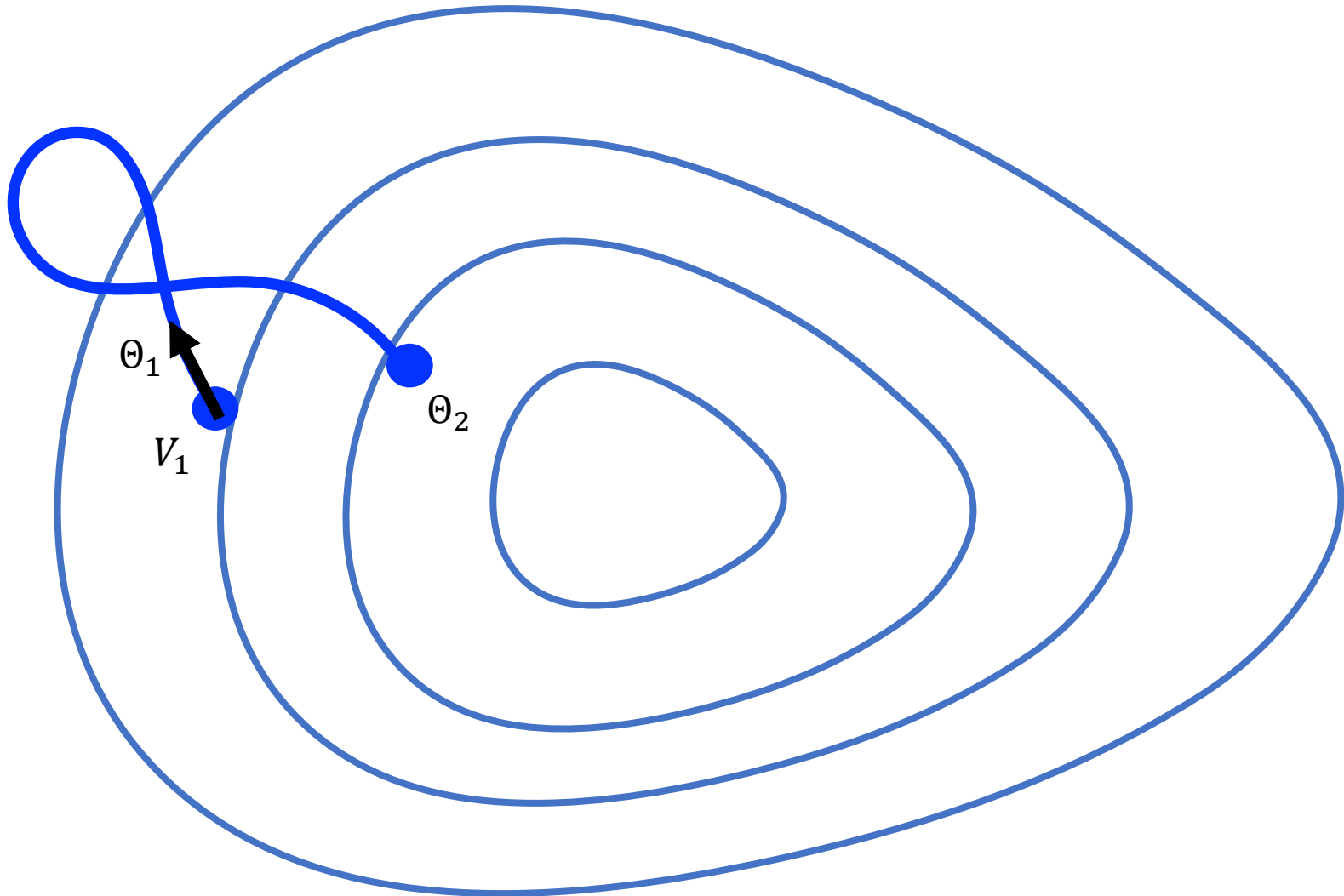
$d^{1/4}$ gradient evaluations are sufficient for 2nd-order HMC to sample from $O(1)$ -smooth, $O(1)$ -strongly convex π with bounded higher-order derivatives

HMC



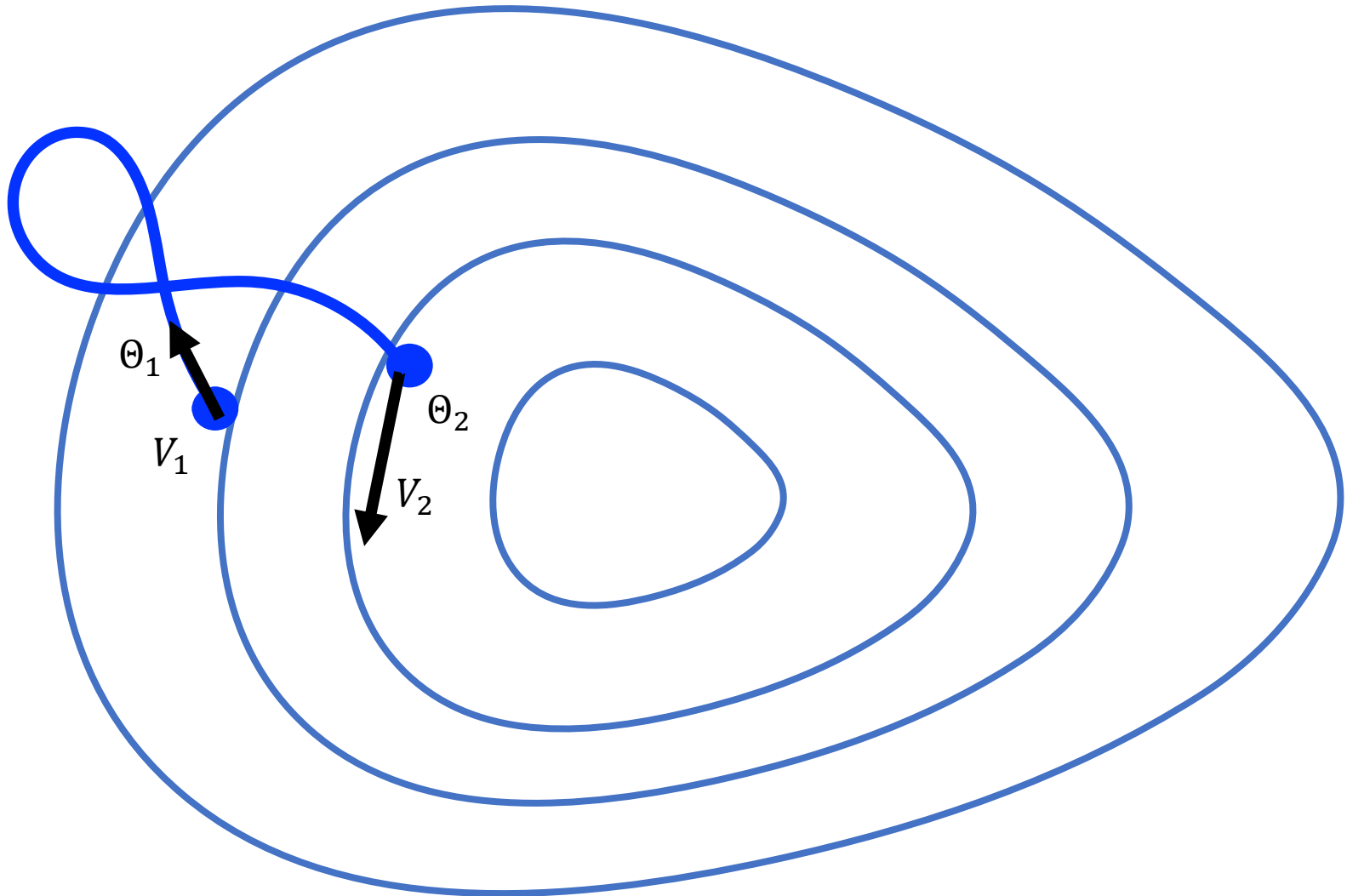
Step 1: sample $V_1 \sim N(0, I_d)$

HMC



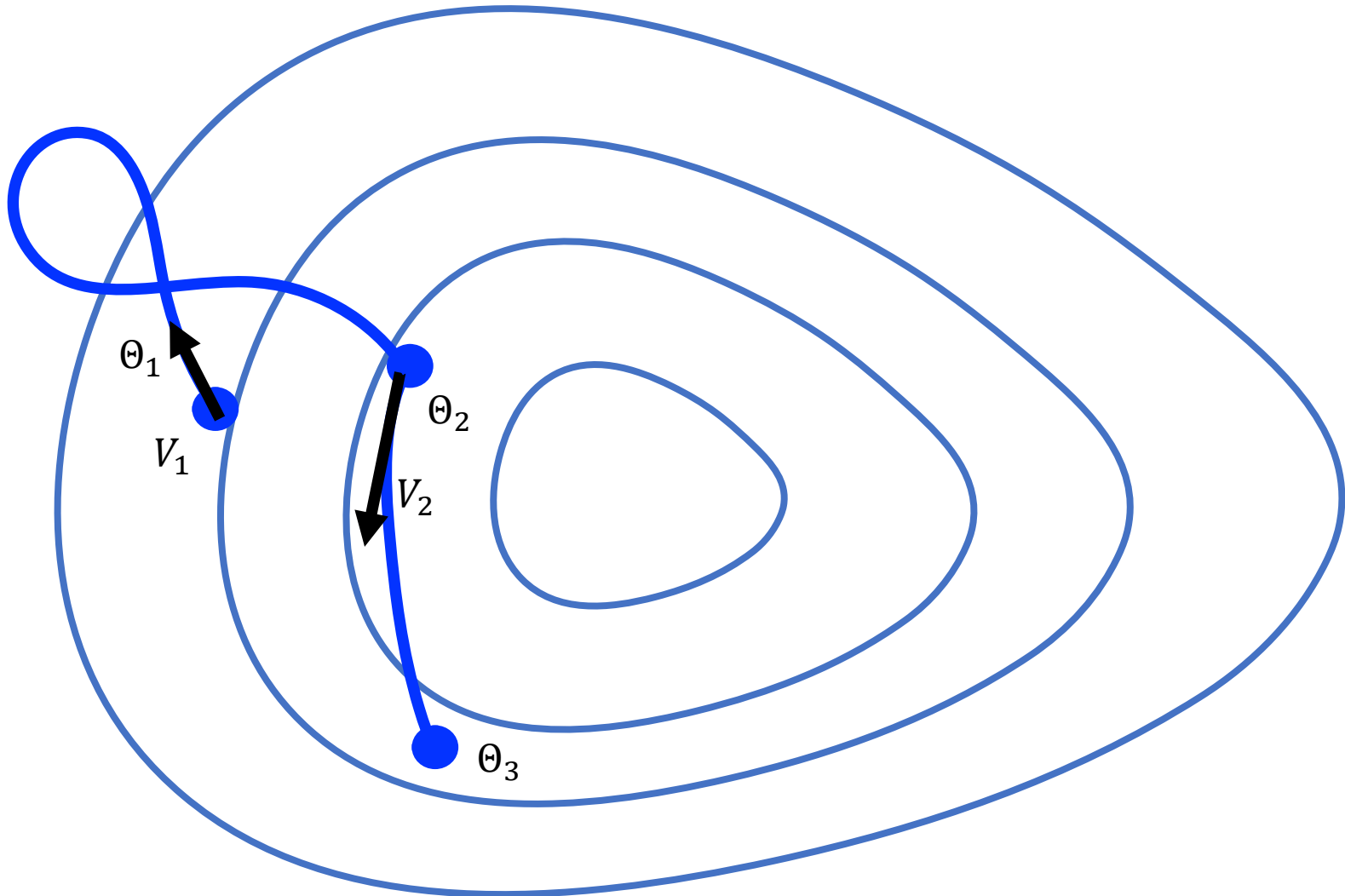
Step 2: Compute Hamiltonian trajectory for fixed time T

HMC



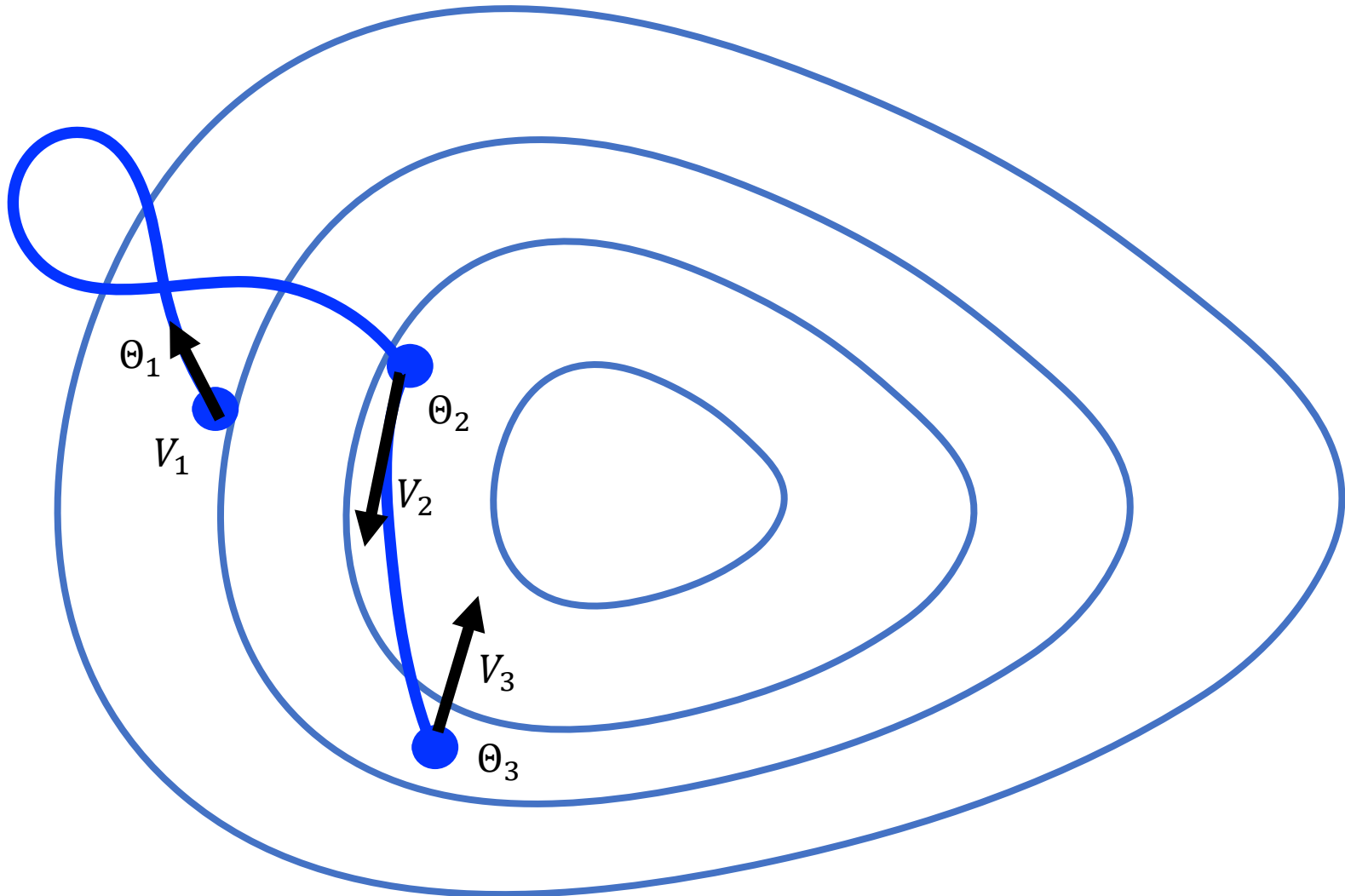
Step 3: Throw out old momentum and sample new independent momentum

HMC



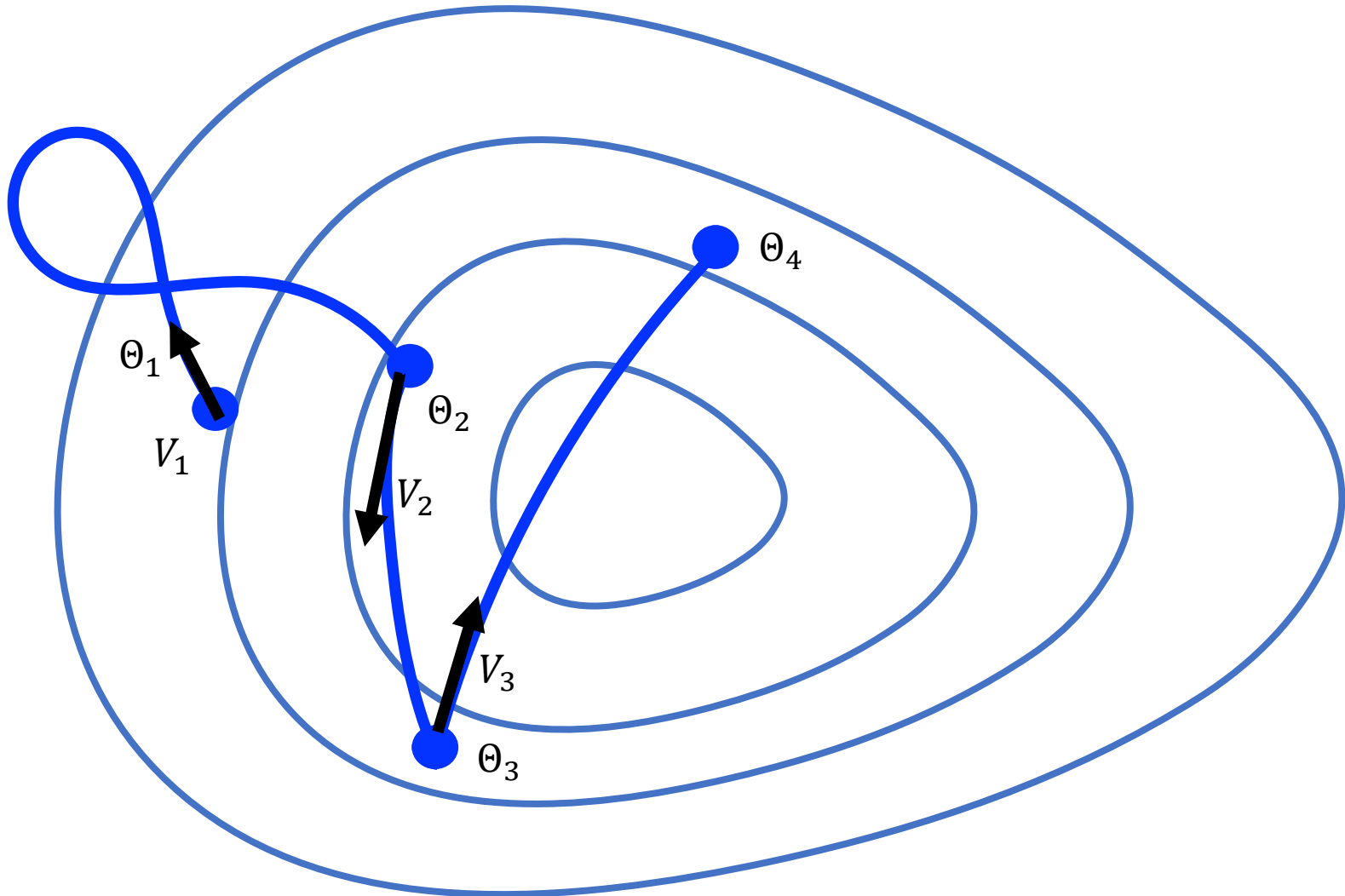
steps 4,5,...: iteratively repeat steps 1 and 2

HMC

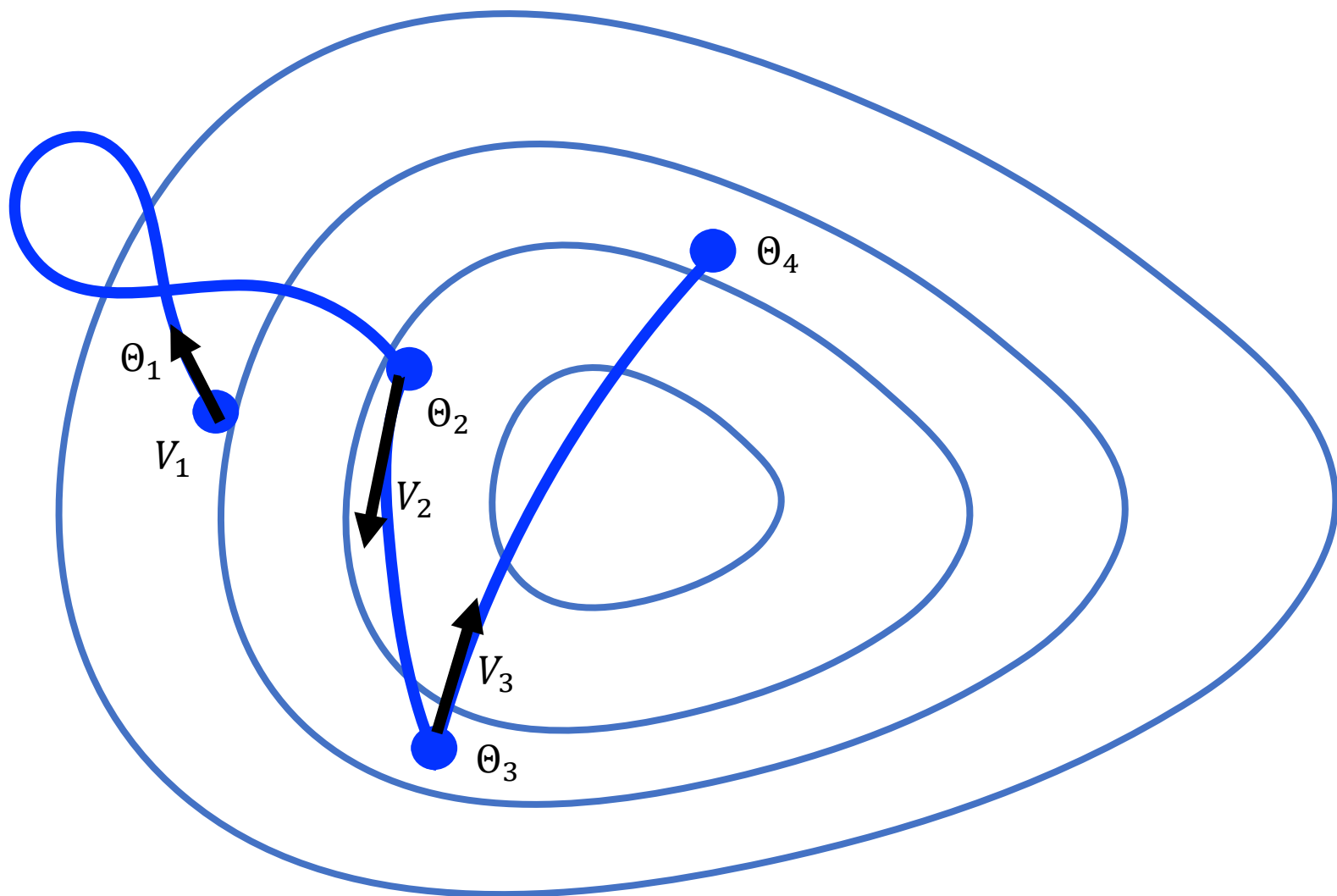


Steps 4,5,...: iteratively repeat steps 1 and 2

HMC



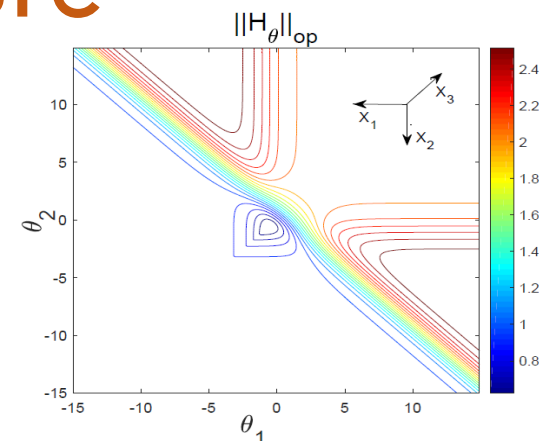
Steps 4,5,...: iteratively repeat steps 1 and 2



Confirming Creutz's Conjecture

[Mangoubi-V. NeurIPS '18] Strongly convex f + regularity conditions

HMC with Leapfrog Integrator requires (roughly) $d^{1/4}$ gradient evaluations



Bit more formally: Suppose that

1. $\frac{1}{10}I \preceq \nabla^2 f(\theta) \preceq 10I$
2. $\nabla^2 f$ satisfies a Lipschitz condition for $L_\infty, r > 0$ and $x_1, \dots, x_r \in \mathbb{S}^d$:

$$\|(\nabla^2 f(\theta_1) - \nabla^2 f(\theta_2))v\|_2 \leq L_\infty \|X^\top(\theta_1 - \theta_2)\|_\infty \times \|X^\top v\|_\infty,$$

where $X := [x_1, \dots, x_r]$

Then Leapfrog HMC requires $\tilde{O}(\max(d^{\frac{1}{4}}, \sqrt{L_\infty}) \varepsilon^{-1/2})$ gradient calls to obtain a sample ε -close (in Wasserstein-2 metric) to π

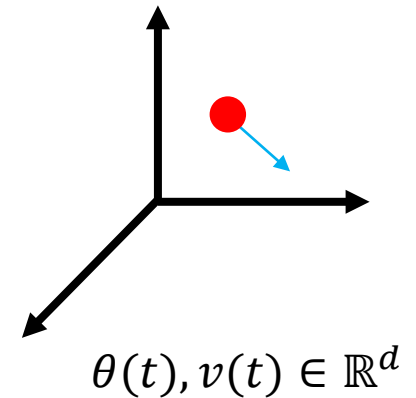
Application of our result: *Fast sampling from logistic “ridge” regression*

$$f(\theta) = \|\theta\|^2 - \sum_i y_i \log \ell(\theta^\top x_i) + (1 - y_i) \log \ell(-\theta^\top x_i)$$
$$L_\infty = \sqrt{C}, \text{ where coherence } C := \max_{i \in [r]} \sum_{j=1}^r |x_i^\top x_j|$$

Hamiltonian Dynamics

Setting:

- Particle with position $\theta(t)$ and momentum/velocity $v(t)$
- Moves according to classical physics laws in a potential well f



Hamiltonian: $H(\theta, v) = f(\theta) + \frac{1}{2} \|v\|^2$

Properties:

- **Time Reversible**
- **Preserves Hamiltonian (Energy):**

Hamilton's Equations:

- **Momentum:** $\frac{d\theta(t)}{dt} = \frac{\partial H}{\partial v} = v(t)$

$$\frac{dH}{dt} = \sum_i \frac{d\theta_i}{dt} \frac{\partial H}{\partial \theta_i} + \frac{dv_i}{dt} \frac{\partial H}{\partial v_i} = 0$$

- **Force:** $\frac{dv(t)}{dt} = -\frac{\partial H}{\partial \theta} = -\nabla f(\theta(t))$

Preserves Volume:

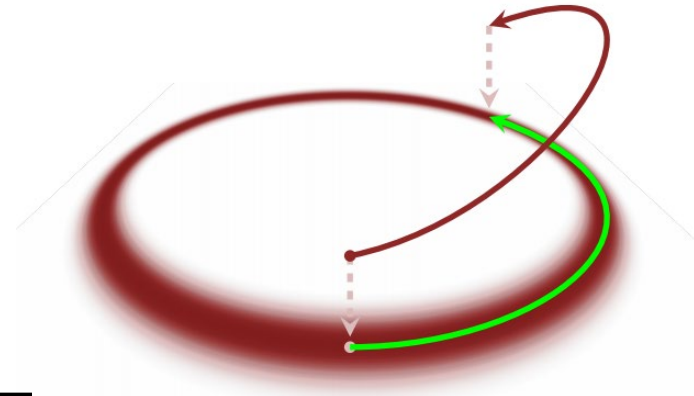
Vector field F in $\mathbb{R}^d \times \mathbb{R}^d$ at (θ, v)

$$\frac{d\theta}{dt}, \frac{dv}{dt}$$

Check: $\text{div } F = \sum_i \frac{\partial}{\partial \theta_i} \frac{d\theta_i}{dt} + \frac{\partial}{\partial v_i} \frac{dv_i}{dt} = 0$

Correctness of continuous-time HMC

Correct: Time reversible,
energy-preserving, volume
preserving (in “phase space”)



Proof: Two steps in the HMC chain. Sufficient to that $e^{-H(\cdot, \cdot)}$ is invariant

Refresh Velocity: Only v is changing, independent of θ and sampled from the right marginal. Hence, $e^{-H(\theta, v)} = e^{-f(\theta)} e^{-\frac{1}{2}\|v\|^2}$ is invariant

Simulate Hamiltonian dynamics:

Partition the phase space into infinitesimal cubes and let C be one cube and (θ, v) be a point in C . The probability of being in C is proportional to $e^{-H(\theta, v)} \times \text{vol}(C)$

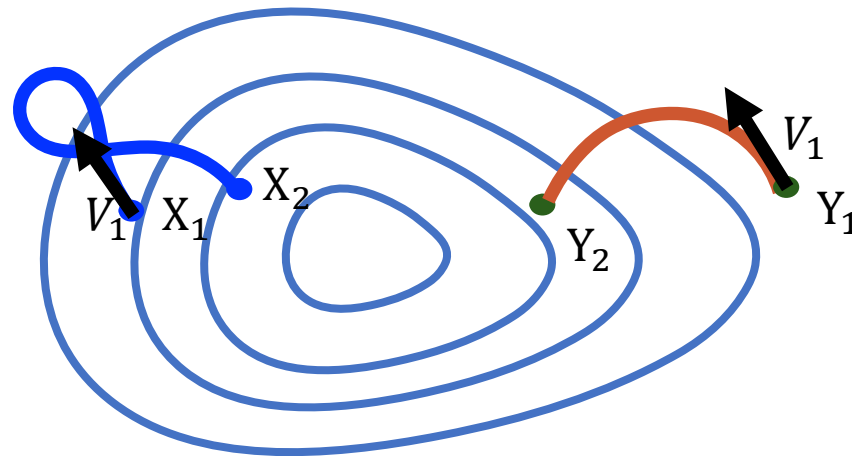
Since the Hamiltonian flow conserves the Hamiltonian (energy) and the volume, the probability of being in the image of C is also conserved (uses *time-reversibility* of Hamiltonian dynamics)

Coupling Bounds for Idealized HMC

- Two chains X_1, X_2, \dots and Y_1, Y_2, \dots with same transition kernel
- Choose a coupling so that $\|X_{i+1} - Y_{i+1}\| \leq c \|X_i - Y_i\|$ for $c < 1$

Momentum (velocity) allows HMC to take long steps ...

Can we couple the momentum of the two (ideal) HMC chains in a way that leads to large (dimension independent) contractions over these long steps?



Hamiltonian trajectories contract for strongly convex potentials; c independent of d

Exercise: $f(\theta) = \sum_i c_i \theta_i^2, \frac{1}{10} \leq c_i \leq 10$

Example: Coupled Pendulums



- pendulums kicked with the same initial velocity
- distance between pendulums contracts for a long time
- difference between velocities increases during this time

Leapfrog Integrator

2nd-order Leapfrog integrator

For $j = 0, \dots, \frac{T}{\eta}-1$, do

$$\theta_{j+1} = \theta_j + \eta v_j - \frac{1}{2}\eta^2 \nabla f(\theta_j)$$

$$v_{j+1} = v_j - \eta \nabla f(\theta_j) - \frac{1}{2}\eta^2 \frac{\nabla f(\theta_{j+1}) - \nabla f(\theta_j)}{\eta} \approx \eta^2 \nabla^2 f(\theta_j) v_j = \eta^2 H(\theta_j) v_j$$

- ***Symplectic integrator***: Approximately conserves target measure
 - volume is conserved in phase space
 - a *perturbed* Hamiltonian is conserved
- Only one gradient call per iteration

Bound numerical error for a given discretization η ?

Our Lipschitz Hessian Condition

Suppose “Euclidean Lipschitz” Hessian

$$\|(H(\theta_1) - H(\theta_2))v\|_2 \leq L_2 \cdot \|\theta_1 - \theta_2\|_2 \cdot \|v\|_2$$

Turns out that numerical error: $\|\eta(H(\theta + \eta v) - H(\theta))v\|_2 \leq L_2 \cdot \eta^2 \cdot \|v\|_2^2$

Here $v \sim N(0, I_d)$, so $\|v\|_2 \approx \sqrt{d}$, so $\eta \approx 1/\sqrt{d}$ leading to no better than \sqrt{d} bound!

Idea: Use a different norm ..

Infinity Lipschitz Hessian

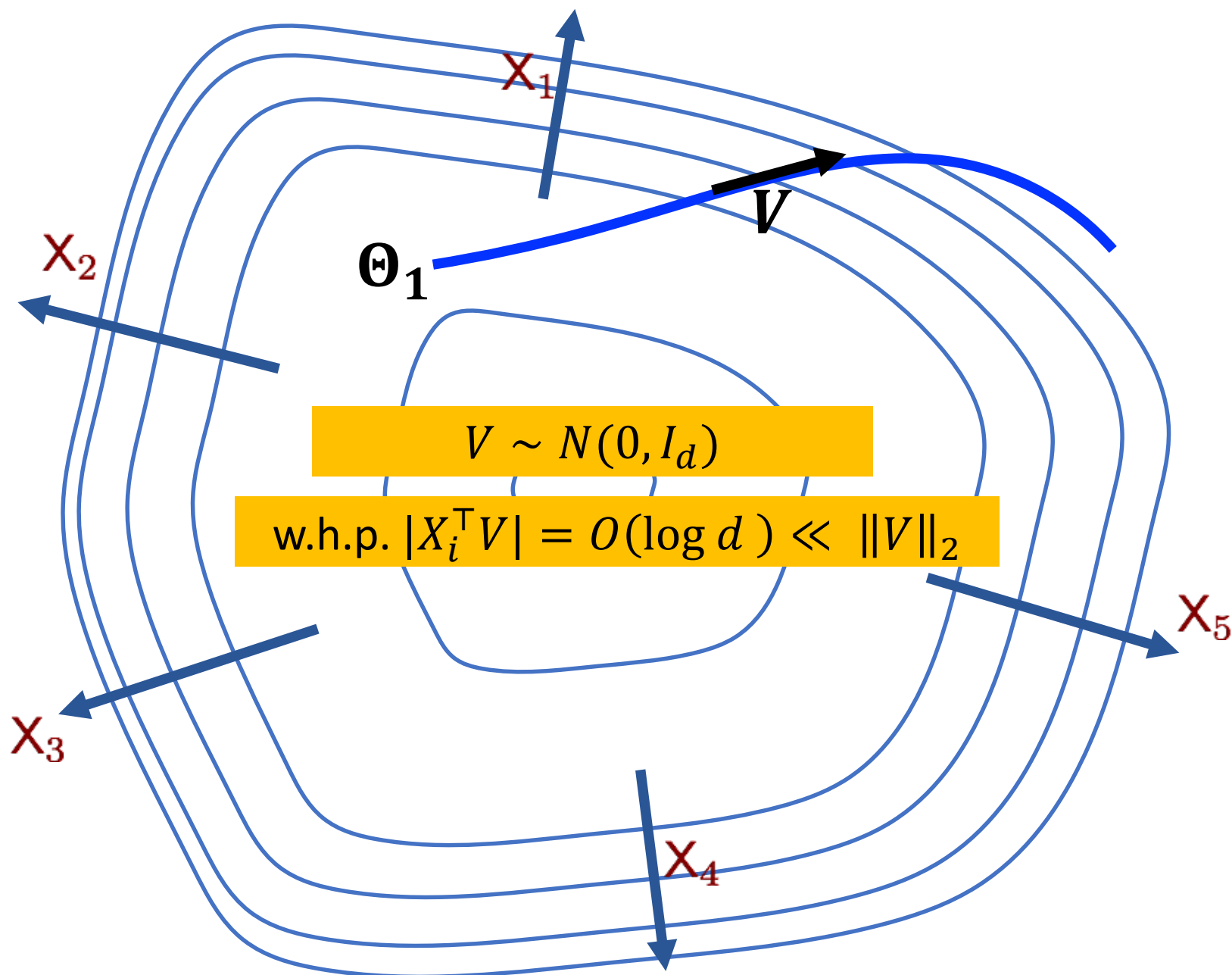
$$\|(H(\theta_1) - H(\theta_2))v\|_2 \leq L_\infty \cdot \|\theta_1 - \theta_2\|_\infty \cdot \|v\|_\infty$$

Positive: $\|v\|_\infty \approx \sqrt{\log d}$, **Negative:** L_∞ is large unless f is separable

Idea: transform the norm to align with the “data vectors”

We use: $\|X^\top v\|_\infty$ where $X := [x_1, \dots, x_r]$

Intuition



Concluding the Proof (for $d = r$)

We bound (inductively on j) the errors $\|\theta_j - \theta(\eta j)\|_2$ and $\|v_j - v(\eta j)\|_2$ by $O(\eta j \varepsilon)$, where $(\theta(t), v(t))$ is the continuous solution to Hamilton's equations with initial conditions in that phase. Since $\eta j \leq T = O(1)$, $O(\eta j \varepsilon) = O(\varepsilon)$

- The error in the quadratic term of the velocity update is roughly
$$\|(\eta^2 H(\theta + \eta v_j) - \eta^2 H(\theta))v_j\|_2 \leq \eta^3 L_\infty \sqrt{d} \|X^\top v_j\|_\infty^2$$
- The invariance property of Hamiltonian mechanics implies v_j is roughly $N(0, I_d)$ at every point on the exact trajectory if HMC has a warm start
- Thus, $\|X^\top v_j\|_\infty = O(\log(d))$ w.h.p., since by inductive assumption $\|v_j - v(\eta j)\|_2 = O(\eta j \varepsilon) = O(1)$
- After T/η iterations, the errors sum to $\tilde{O}(\eta^2 L_\infty \sqrt{d})$. Choosing η to have error ε , # of gradients is $T/\eta = \tilde{\Theta}(\varepsilon^{-1/2} d^{1/4} L_\infty^{1/2})$

LANGEVIN DYNAMICS,
SIMULATED ANNEALING
AND
NOISY CONVEX OPTIMIZATION

Optimizing using Noisy Oracles

Input: Noisy approximation \hat{F} to a convex function $F: \mathbb{R}^d \rightarrow \mathbb{R}$ with global minimum θ^*

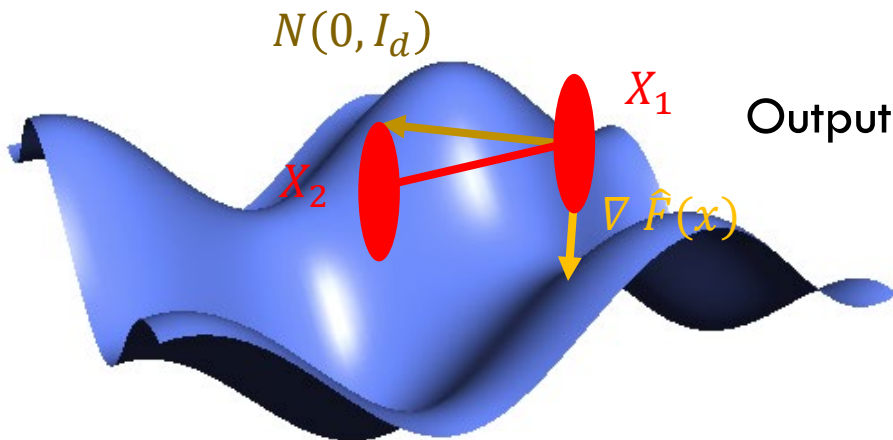
Goal: Find \hat{x} , s.t. $F(\hat{\theta}) - F(\theta^*) < \varepsilon$ for given $\varepsilon > 0$

Applications:

- Optimizing F when an accurate value of F is expensive to compute
- Optimizing non-convex functions which are close to a convex function

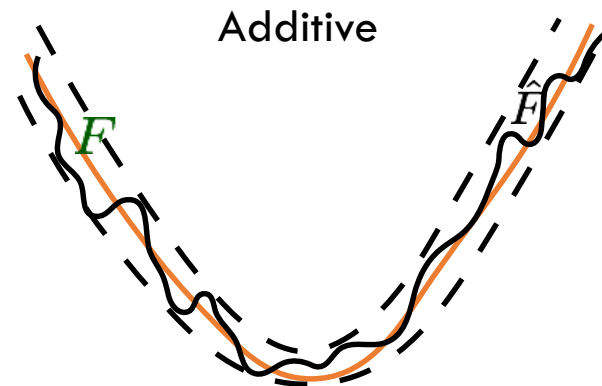
Algorithm: Langevin Dynamics

(Arises in statistical physics)

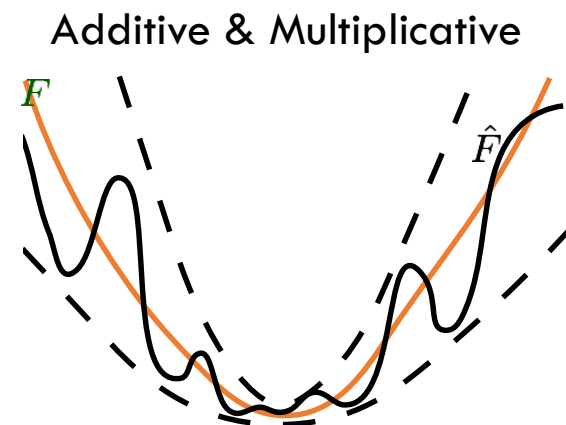


Outputs $\theta \propto e^{-\hat{F}(\theta)/T}$

$$X_{i+1} = X_i - \eta \nabla \hat{F}(X_i) + \sqrt{2\eta T} N(0, I_d)$$



[Applegate-Kannan '91, Zhang et al. '17]



[Belloni-Liang-Narayanan-Rakhlin '15]

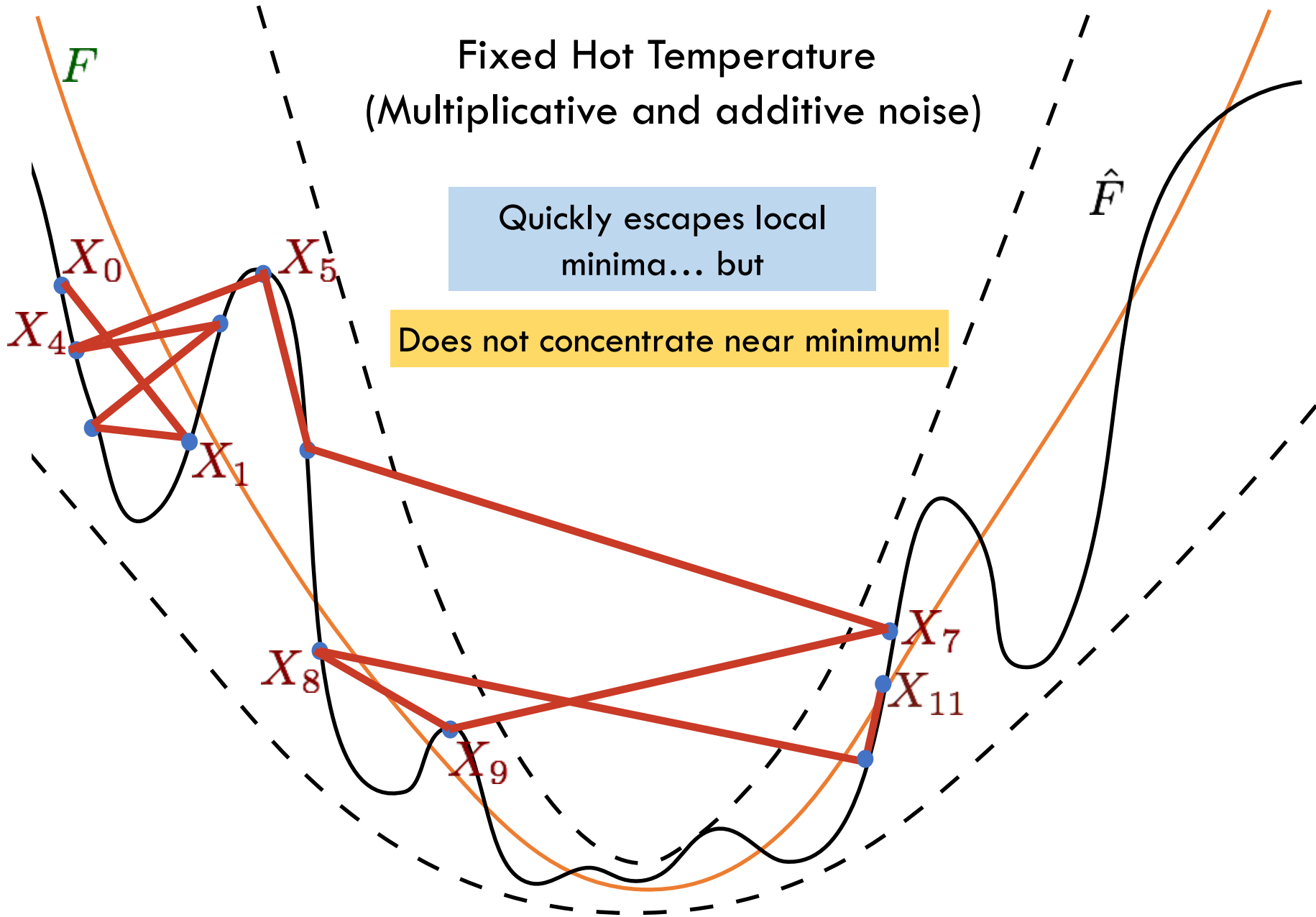
High T : escape local minima quickly

Low T : limiting distribution concentrates near minimum value

Fixed Hot Temperature
(Multiplicative and additive noise)

Quickly escapes local
minima... but

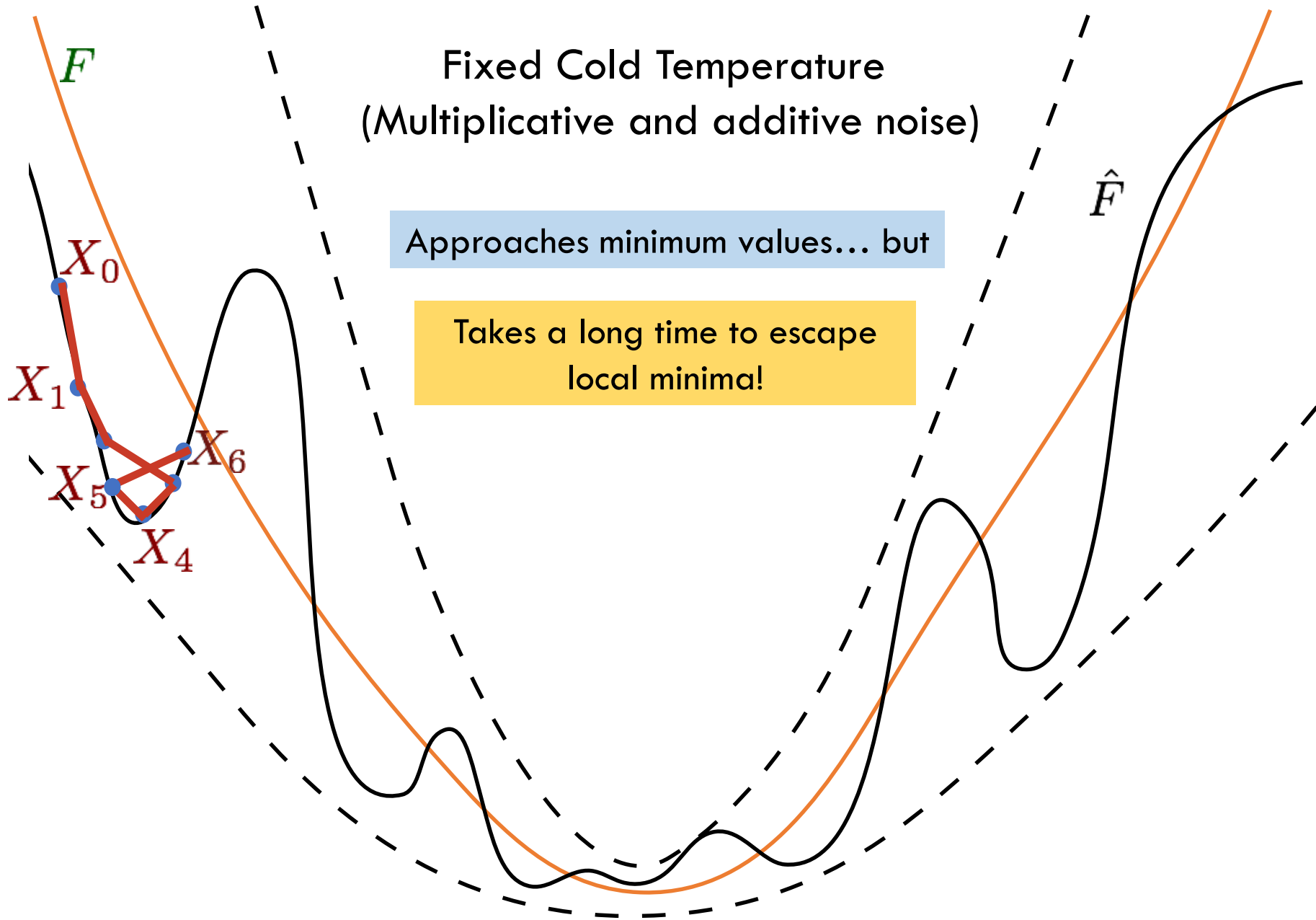
Does not concentrate near minimum!



Fixed Cold Temperature
(Multiplicative and additive noise)

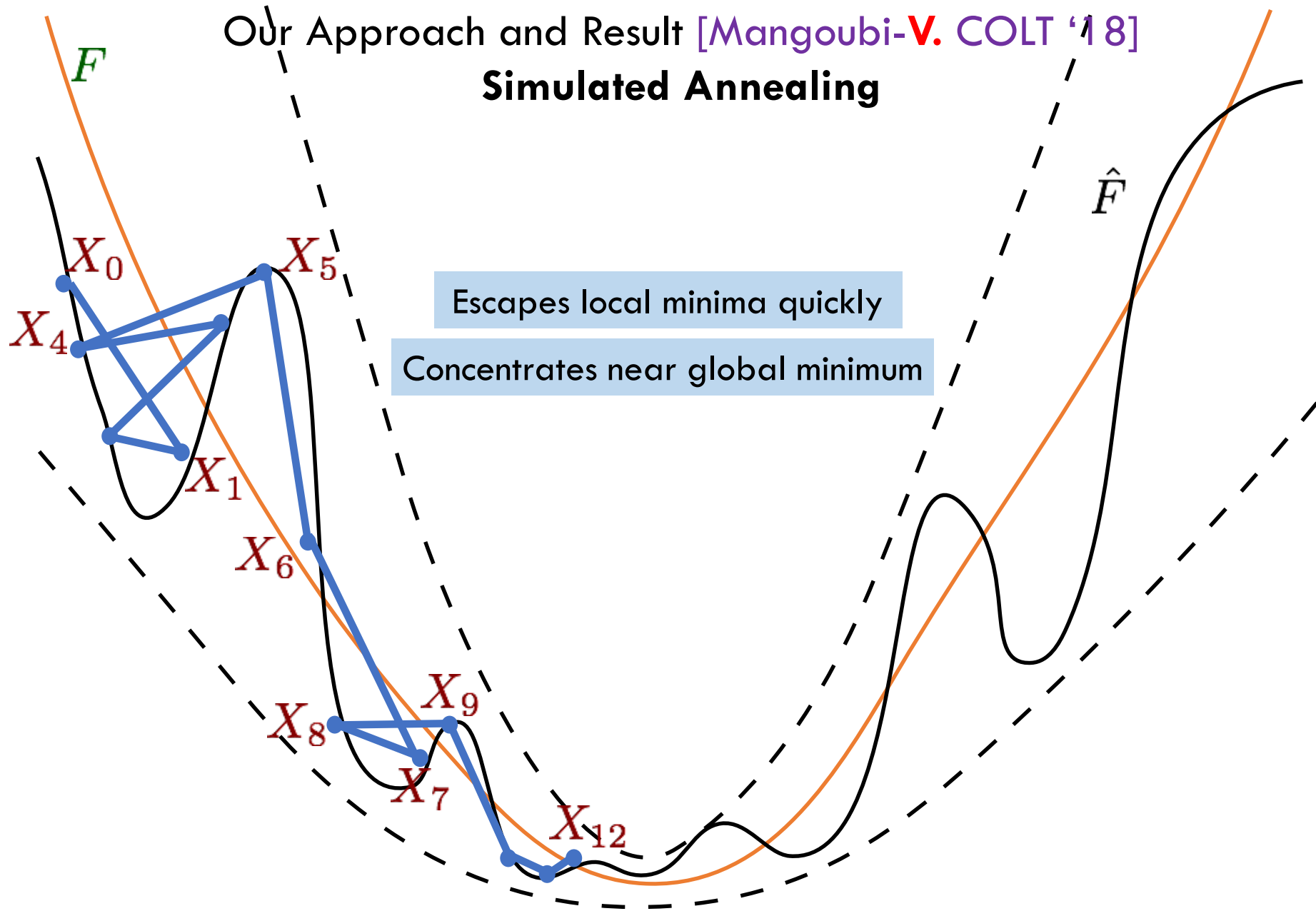
Approaches minimum values... but

Takes a long time to escape
local minima!



Our Approach and Result [Mangoubi-V. COLT '18]

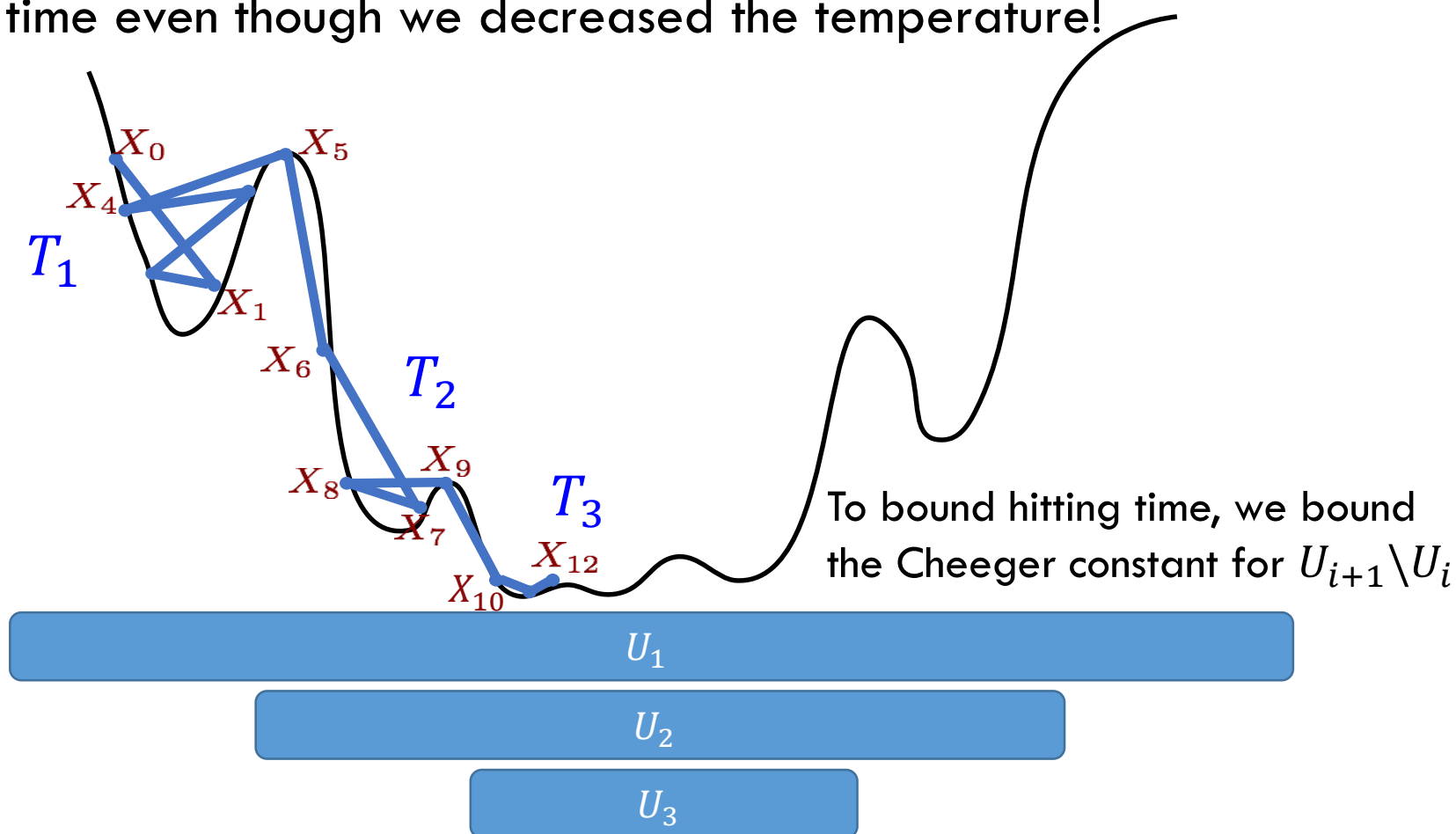
Simulated Annealing



Theorem: Suppose the approximation \hat{F} has additive noise $< O(\frac{\varepsilon}{d})$ and multiplicative noise $< O(\frac{1}{d})$. Then, our algorithm can minimize F to accuracy ε in time that is polynomial in d

Proof Strategy

1. At epoch $i + 1$: show that our algorithm remains inside U_i w.h.p.
2. Noise is lower than at previous i , since multiplicative noise decreases
3. A lower noise implies shallower local minima and therefore a fast hitting time even though we decreased the temperature!



Conclusion

- Physics and physical systems (have been) and can be a great source of ideas and insights for
 - **rigorous** algorithm design
 - **tuning** parameters
 - identifying the right **potential** functions
 - obtaining “**beyond worst case**” assumptions on functions/data
- **Other recent physics-inspired optimization and sampling algo.**
 - Online sampling with $O(1)$ update time [H.Lee-Mangoubi-**V.** NeurIPS '19]
 - Langevin dynamics for non-convex potentials [Mangoubi-**V.** COLT '19]
 - Sampling from polytopes [Mangoubi-**V.** FOCS '19]

Thanks! Questions?