

# Lattice study of gauge/gravity duality in two-dimensional $N=(8,8)$ SYM

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“Nonperturbative and numerical approaches to quantum gravity, string theory and holography”  
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Giguere and D.K., JHEP 1505, 082 (2015).

Giguere and D.K., in preparation

# Talk plan

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1. Duality and Lattice studies
2. Sugino lattice action and SUSY WTI
3. Numerical test of Duality in 2D  $N=(8,8)$  SYM
4. Summary



# 1. Duality and Lattice studies

# Gauge/gravity duality

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- AdS/CFT [Maldacena, 1997] and more general duality conjecture states

strongly coupled gauge theory

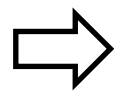
= classical gravity on a curved space

- Application of duality

(1) Non-perturbative formulation of string theory

(2) BH information loss paradox

(3) Holographic QCD, superconductor...



precision test of the gauge/gravity duality  
by lattice simulations is important.

# Maximally supersymmetric Yang-Mills

- Action in  $p+1$ -dimensions

$$S = \frac{N}{\lambda} \int d^{p+1} x \operatorname{tr} \left\{ \frac{1}{2} F_{\mu\nu}^2 + \frac{1}{2} (D_\mu X_i)^2 - \frac{1}{4} [X_i, X_j]^2 - \frac{i}{2} \psi_\alpha \gamma_\mu D_\mu \psi_\alpha + \frac{1}{2} \psi_\alpha \gamma_i [X_i, \psi_\alpha] \right\}$$

$A_\mu$  : gauge field       $\mu = 0, \dots, p$

$X_i$  : scalar field       $i = p+1, \dots, 9$

$\psi_\alpha$  : fermion       $\alpha = 1, \dots, 16$

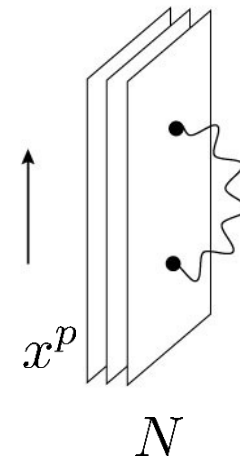
$\lambda = g^2 N$  't Hooft coupling constant

$$\gamma_0 = i$$

$$\{\gamma_m, \gamma_n\} = \delta_{mn}$$

- dual to black  $p$ -brane solution in SUGRA

[Maldacena 1997, Itzhaki et al. 1998]



# Numerical tests of duality in $p+1$ dimensions

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- $p=0$  (BFSS model)

momentum sharp cut-off:

- Hanada et al (2007-2014)

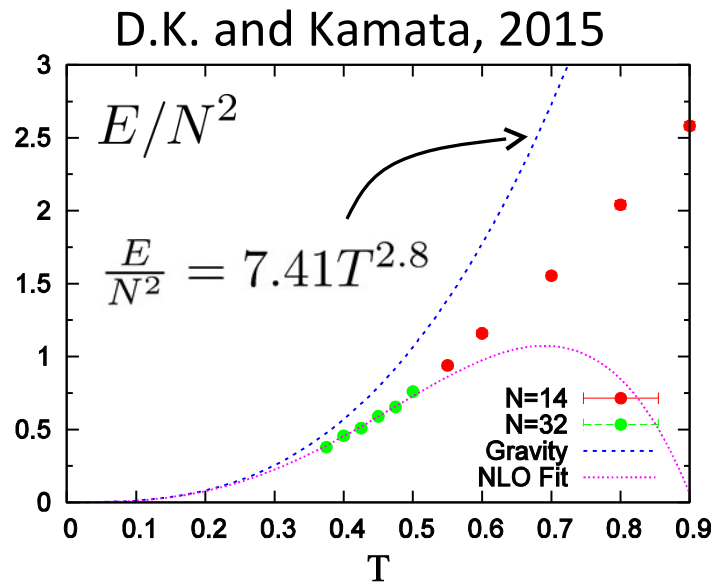
lattice:

Hanada's lecture

- Catterall and Wiseman (2007-2010)
- D.K. and Kamata (2012-2015)
- Filev and O'Connor (2015)
- Berkowitz, Rinaldi, Hanada, Ishiki, Shimasaki, Vranas (2016)

...

# Internal energy of Black hole (p=0)



$$\frac{E}{N^2} = 7.41T^{2.8} - CT^p + DT^{p+6/5}$$

	$C$	$p$	$D$
SUGRA	unknown	4.6	unknown
SYM			
Hanada et al. (2007)	$5.55 \pm 0.07$	$4.58 \pm 0.03$	—
D.K. – Kamata (2015)	$9.0 \pm 2.6$	$4.74 \pm 0.35$	—
Berkowitz et al. (2016)	$10.2 \pm 2.4$	$4.6 \pm 0.3$	$6.2 \pm 2.6$

# Lattice studies for $p \geq 1$ cases

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- $p=1$

Catterall, Joseph and Wiseman (2008-2010)

Giguere and D.K. (2015-) ← This talk

David Schaich's talk

- $p=2$

Raghav G. Jha's talk

- $p=3$  (AdS/CFT)

Catterall, Giedt, Schaich, Damgaard,  
Damgaard and Degrand (2014-)



# This work

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2D  $N=(8,8)$  SYM  
at finite  $T$  and large  $N$

$\stackrel{=}{?}$

Thermodynamics of  
black string

- Sugino lattice action [Sugino, 2005]

↪ exact two supersymmetries

- Standard Monte-Carlo method (with quenching phase of fermion pfaffian)

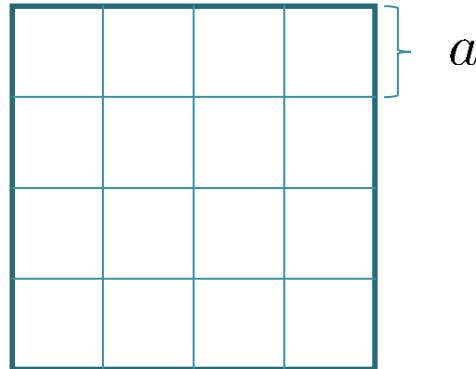


## 2. Sugino lattice action and SUSY WTI

# Lattice gauge theory

Wilson 1974

- Lattice



- Fields live on sites or links
- $\partial_\mu \rightarrow$  difference operator  $\nabla_\mu$

- Lattice gauge field

continuum limit

$$P_{\mu\nu}(x) \sim 1$$

$$U_\mu(x, y) = e^{iaA_\mu(x, y)} \in G$$

$$S_G = \frac{N}{g^2} \sum_{x, \mu < \nu} \left\{ 1 - \frac{1}{N} \text{Re tr} (P_{\mu\nu}(x)) \right\} \longrightarrow \int d^4x \frac{1}{4} \text{tr} (F_{\mu\nu}^2(x))$$

$$P_{\mu\nu}(x) = U_\mu(x) U_\nu(x + a\hat{\mu}) U_\mu^{-1}(x + a\hat{\nu}) U_\nu^{-1}(x)$$

# Lattice supersymmetry

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- SUSY invariance in continuum theory

$$\partial(fg) = (\partial f)g + f(\partial g)$$

$$\Rightarrow \delta S_{cont.} = \int d^4x \partial_\mu X_\mu(x) = 0$$

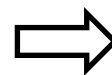
- SUSY is broken by the lattice cut-off

Dondi-Nicolai, Kato-Sakamoto-So, Bergner

$$\nabla(fg) = (\nabla f)g + f(\nabla g) + O(a)$$

$$\delta S_{lat.} = a^4 \sum_x \{ \partial_\mu X_\mu(x) + O(a) \} = O(a)$$

+ UV divergences  
 $O(1/a)$



SUSY is broken  
at quantum level

# Lattice SYM with exact SUSY

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- 2003-2004:

Cohen-Kaplan-Katz-Unsal

Sugino

Catterall

Catterall's lecture

- The other lattice formulations of SYM and applications:

Giedt, Endres, Kawamoto, D'Adda, Kanamori, Kikukawa, Suzuki,

Damgaard, Matsuura, Hanada, Joseph, O'Connor, Schaich, Jha,

... and D.K.

- c.f. Lattice Wess-Zumino model and SUSY QM:

Dondi, Nicolai, Sakai, Sakamoto, Wipf, Wenger,

Bergner, Wozar, Baumgartner,...

# Q-exact form of continuum N=(8,8) action

- Q-transformation

$$Q_{\pm}A_{\mu} = \psi_{\pm\mu}, \quad Q_{\pm}\psi_{\pm\mu} = iD_{\mu}\phi_{\pm}, \dots$$

$$\Rightarrow \quad Q_{\pm}^2 = i\delta_{\phi_{\pm}}, \quad \{Q_{+}, Q_{-}\} = -i\delta_C$$

$\delta_{\omega}$  : gauge transformation

- Q-exact action

$$S = Q_{+}Q_{-} \frac{N}{2\lambda} \int d^2x \operatorname{tr} \left\{ -2iB_i G_i - \frac{2}{3}\epsilon_{ijk} B_i B_j B_k \right. \\ \left. -\psi_{+\mu}\psi_{-\mu} - \chi_{+i}\chi_{-i} - \frac{1}{4}\eta_{+}\eta_{-} \right\}$$

$$G_0 = D_0 X_3 + D_1 X_2$$

$$G_1 = D_1 X_3 - D_0 X_2$$

$$G_2 = i[X_2, X_3] + F_{01}$$

# Sugino lattice action

[Sugino,2005]

- Q-exact lattice action

$$S = Q_+ Q_- \frac{N}{2\lambda} \sum_x \text{tr} \left\{ -2iB_i G_i - \frac{2}{3} \epsilon_{ijk} B_i B_j B_k - \psi_{+\mu} \psi_{-\mu} - \chi_{+i} \chi_{-i} - \frac{1}{4} \eta_+ \eta_- \right\}$$

$$G_0 = \nabla_0^+ X_3 + \nabla_1^+ X_2$$

$$G_1 = \nabla_1^- X_3 - \nabla_0^- X_2$$

$$G_2 = i[X_2, X_3] + F_{01}^{lat}$$

lattice field strength

$\nabla_\mu^\pm$  : forward/ backward difference operators

- Can define lattice Q-transformation satisfying

$$Q_\pm^2 = i\delta_{\phi_\pm}, \quad \{Q_+, Q_-\} = -i\delta_C$$

⇒ Two exact SUSYs

# Choice of gauge action

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- A simple choice (double plaquette gauge action)

$$F_{01}^{lat}(x) = -\frac{i}{2} \left( P_{01}(x) - P_{01}^\dagger(x) - \frac{1}{N} \text{tr}(P_{01}(x) - P_{01}^\dagger(x)) \right)$$

- However, extra vacuum configurations exist:

$$S_B|_{X_i=0} = \frac{N}{2\lambda_0} \sum_{t,x} \text{tr} \left\{ -\frac{1}{4} \left( P_{01} - P_{01}^\dagger - \frac{1}{N} \text{tr}(P_{01} - P_{01}^\dagger) \right)^2 \right\}$$

$$S_B = 0 \text{ at } P_{01} = +1, \quad \textcircled{P_{01} = -1} \quad SU(2)$$

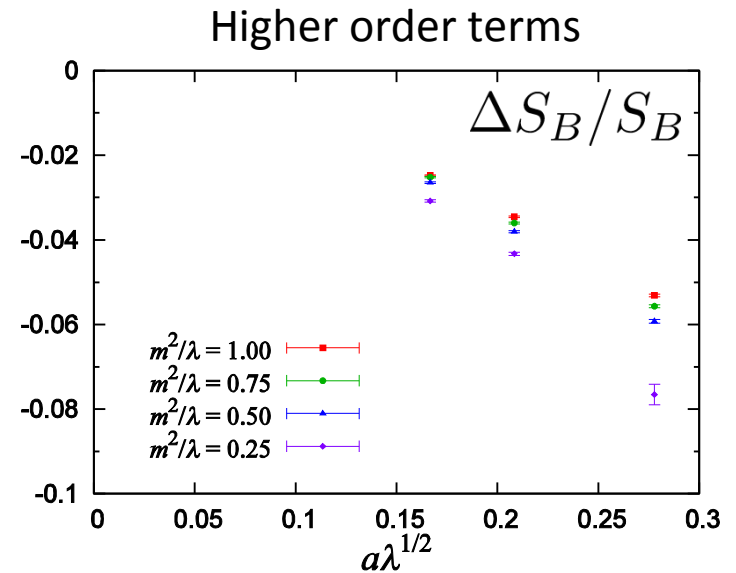
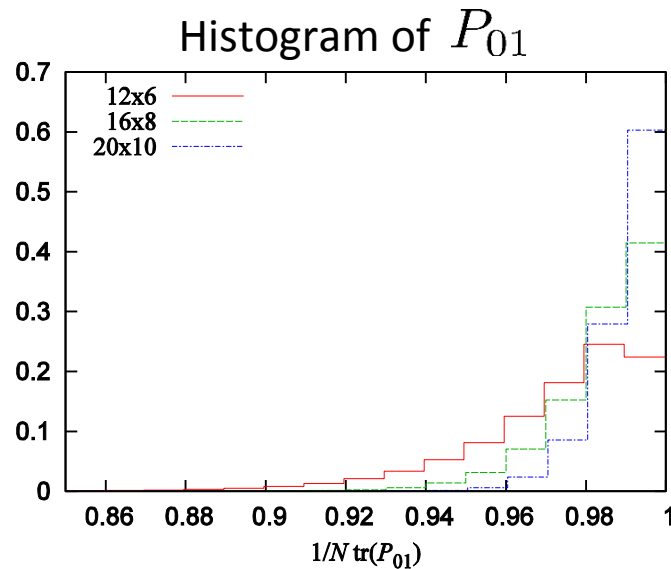
c.f. admissible gauge action [Sugino, 2004]

tangent-type action [Matsuura-Sugino, 2014]



# Test of boson part

Eric Giguere and D.K., 2015



$$\lambda_0 = \lambda a^2 \rightarrow 0$$

$$P_{01}(x) \rightarrow 1, \quad \Delta S_B \rightarrow 0.$$

Sugino lattice action of  $N=(8,8)$  SYM works  
for double plaquette action

# SUSY Ward-Takahashi identity

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- SUSY breaking source
  - (1) finite T (anti-periodic boundary)
  - (2) finite a  $\longrightarrow$  disappear in the continuum limit?

- SUSY WTI with supercurrent  $J_\mu(x)$  holds in continuum SYM even at finite temperature

(c.f. Kananori-Suzuki, 2007)

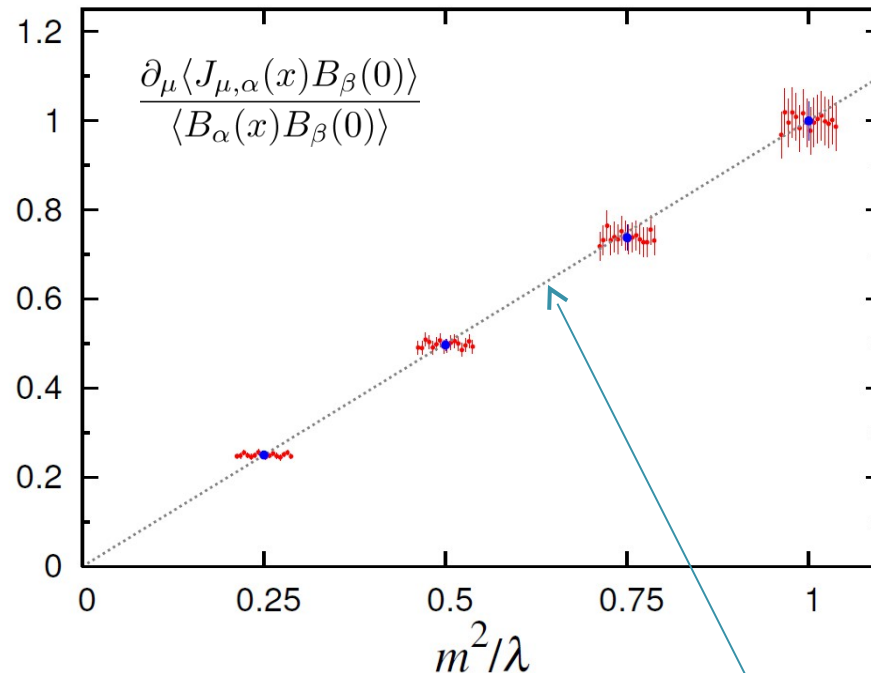
$$\langle \partial_\mu J_\mu(x) \mathcal{O}(y) \rangle = 0 \quad \text{for } x \neq y$$

- Add mass term to the action for avoiding the problem of flat direction:  $S \rightarrow S + \frac{N}{\lambda} \int d^2x \frac{m^2}{2} \text{tr}(X_i^2)$

$$\langle \partial_\mu J_\mu(x) \mathcal{O}(y) \rangle = \frac{m^2}{\lambda} \langle B(x) \mathcal{O}(y) \rangle \quad \text{for } x \neq y$$

# Numerical test of Ward-Takahashi identity

Eric Giguere and D.K., 2015




SU(2)  
T=1

$$\frac{\partial_\mu \langle J_{\mu,\alpha}(x) B_\beta(0) \rangle}{\langle B_\alpha(x) B_\beta(0) \rangle} = \frac{m^2}{\lambda}$$

SUSY WTI holds for all  $(J_\mu(x))_\alpha, \alpha = 1, 2, \dots, 16$ .

All of the SUSYs are restored in the continuum limit.



### 3. Numerical test of Duality in 2D $N=(8,8)$ SYM

# SYM at finite temperature

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- $\lambda$  is only dimensionful parameter

$$S = \frac{N}{\lambda} Q_+ Q_- (\dots) \quad \dim[\lambda] = 2$$

- Scale-transformation  $\dim[a] = -1$

$$\beta \equiv \frac{1}{T} = \beta_0 a, \quad L = L_0 a, \quad \lambda = \lambda_0 a^{-2}$$

$$\begin{aligned} a \frac{\partial}{\partial a} \log Z &= a \frac{\partial \beta}{\partial a} \frac{\partial}{\partial \beta} \log Z + a \frac{\partial L}{\partial a} \frac{\partial}{\partial L} \log Z = -\beta L (\epsilon - p), \\ &= 2 \langle S \rangle \end{aligned}$$

$$\Rightarrow \epsilon - p = -\frac{2}{\beta L} \langle S \rangle \xrightarrow{T \rightarrow 0} 0$$

# Thermodynamics of black string

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- SUGRA

$$S = -\frac{1}{2\kappa^2} \int d^{10}x \sqrt{g} \left[ R - \frac{1}{2}(\partial_\mu \phi)^2 - \frac{1}{12}e^\phi F_{\mu\nu\rho}^2 \right],$$

- Black hole thermodynamics [Klebanov-Tseytlin, 1996]

$$E = \frac{2}{3}c_0LN^2T^3$$

$$S = c_0LN^2T^2$$

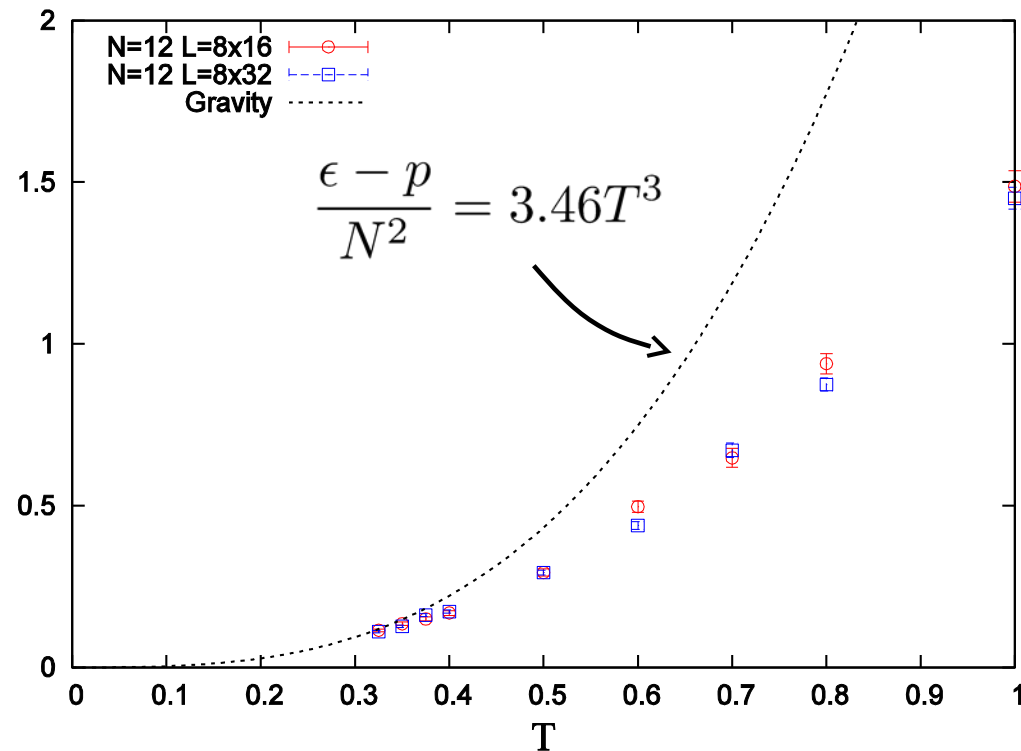
$$c_0 = \frac{2^4\pi^{5/2}}{3^3} = 10.37\dots$$

$$E = TS - PV \quad \Rightarrow \quad \epsilon - p = \frac{1}{3}c_0N^2T^3$$

# Lattice results of e-p

Giguere and D.K., in preparation

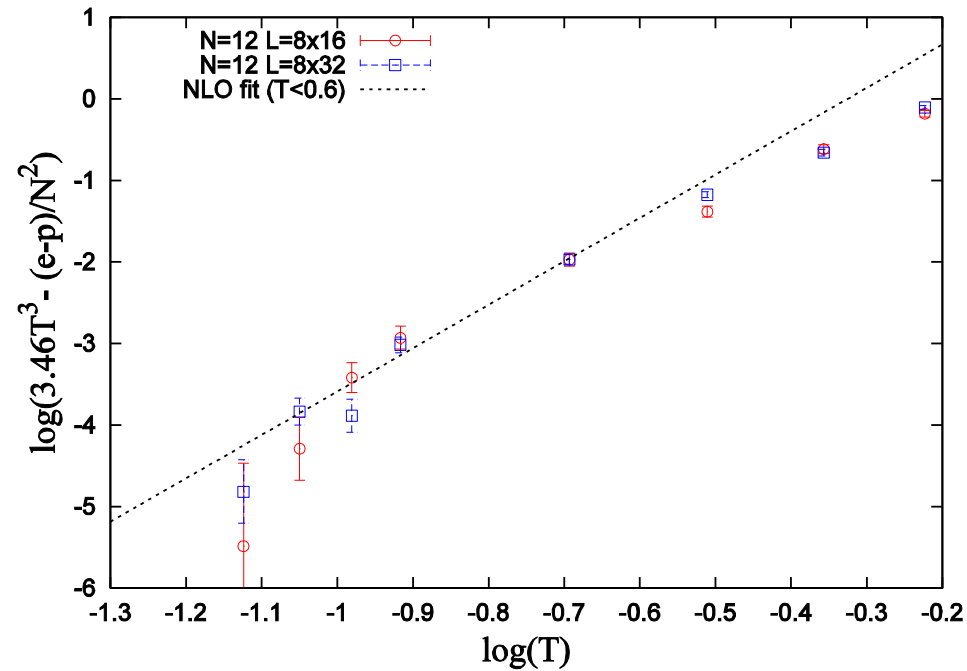
$$\frac{\epsilon - p}{N^2}$$



Gauge/gravity duality holds, at least, qualitatively.

# $\alpha'$ -correction

Giguere and D.K., in preparation



$$\epsilon - p = 3.46T^3 - dT^\beta + \dots$$

$d$

$\beta$

SUGRA	unknown	5
SYM (Our result)	$5.7 \pm 1.3$	$5.32 \pm 0.28$





## 4. Summary

# Summary

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- ☑ SUSY Ward-Takahashi identity

  - Sugino lattice action works well for 2d SYM with 16 SUSY beyond PT.

- ☑ e-p of black string

  - Lattice results reproduce the prediction, and NLO result shows that the duality is true.

- ☐ Take continuum limit and large N limit of the results

- ☐ Need a method solving sign problem

Thank you.