

# Nonlocal transport in InAs/GaSb composite quantum well: A candidate for 2D topological insulator

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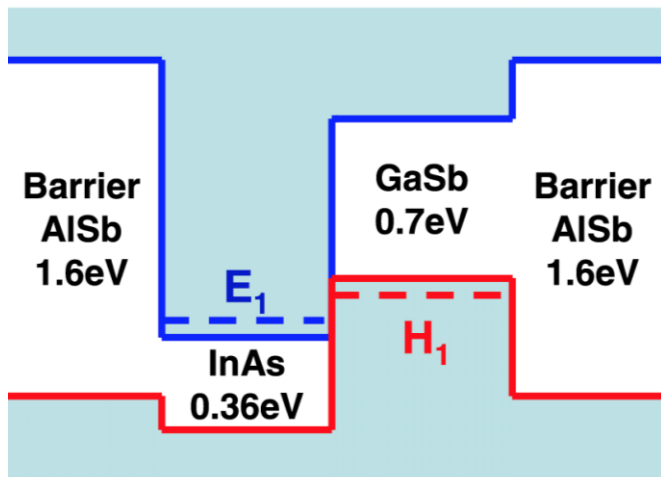
T. Tschirky

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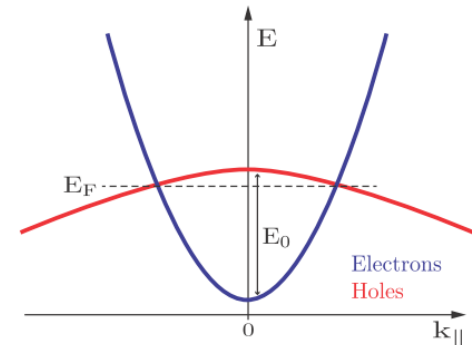
**Solid State Physics Laboratory, ETH Zurich**

ISPCM - 2017, ICTS, Bangalore

# InAs/GaSb Composite Quantum Well

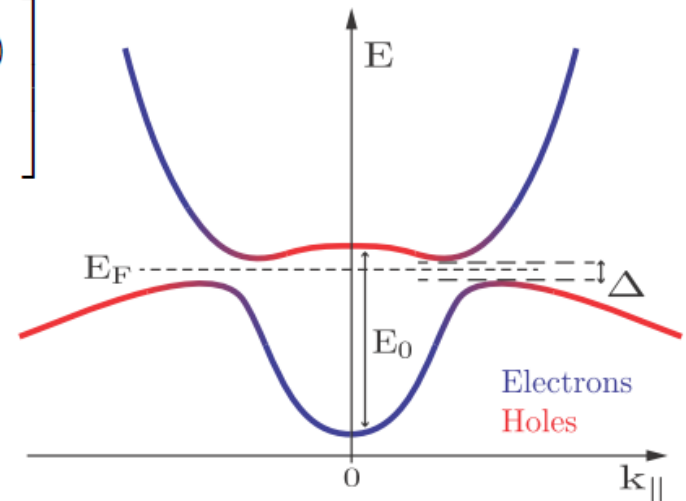


- Peculiar type-III band alignment
- Electron and holes in separate layers
- Hybridization gap due to coupling between layers

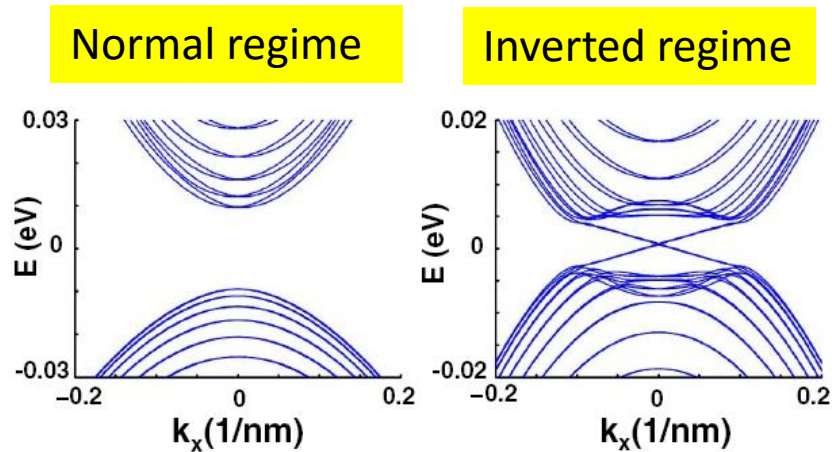


**Inverted**

$$H = \begin{bmatrix} E_e & V(k) \\ V(k)^* & E_h \end{bmatrix}$$



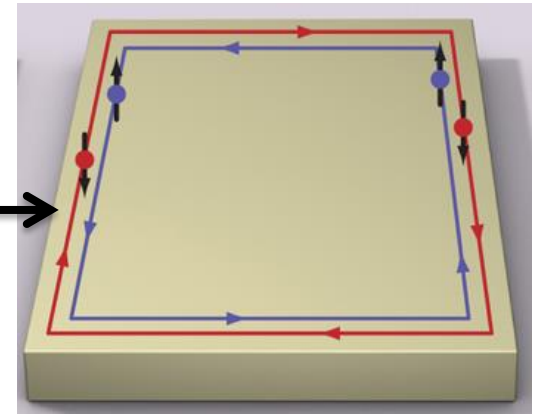
# Prediction of Quantum Spin Hall Effect



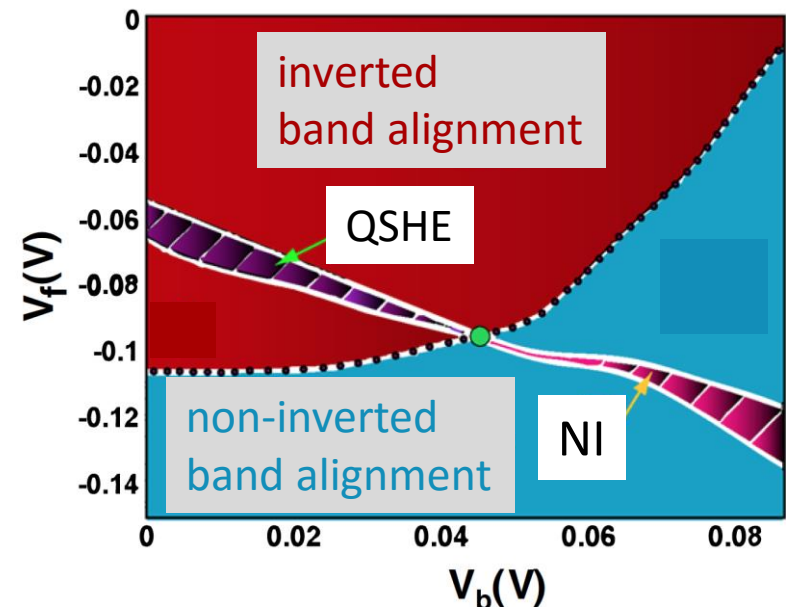
Liu *et al.* Phys. Rev. Lett. 100 236601 (2008)

- Band inversion and Hybridization gap
- Spin-orbit coupling
- Time reversal symmetry
- gap-modes at the sample edge form a Kramers-pair in k-space.

Helical edge channels



Electrically tunable

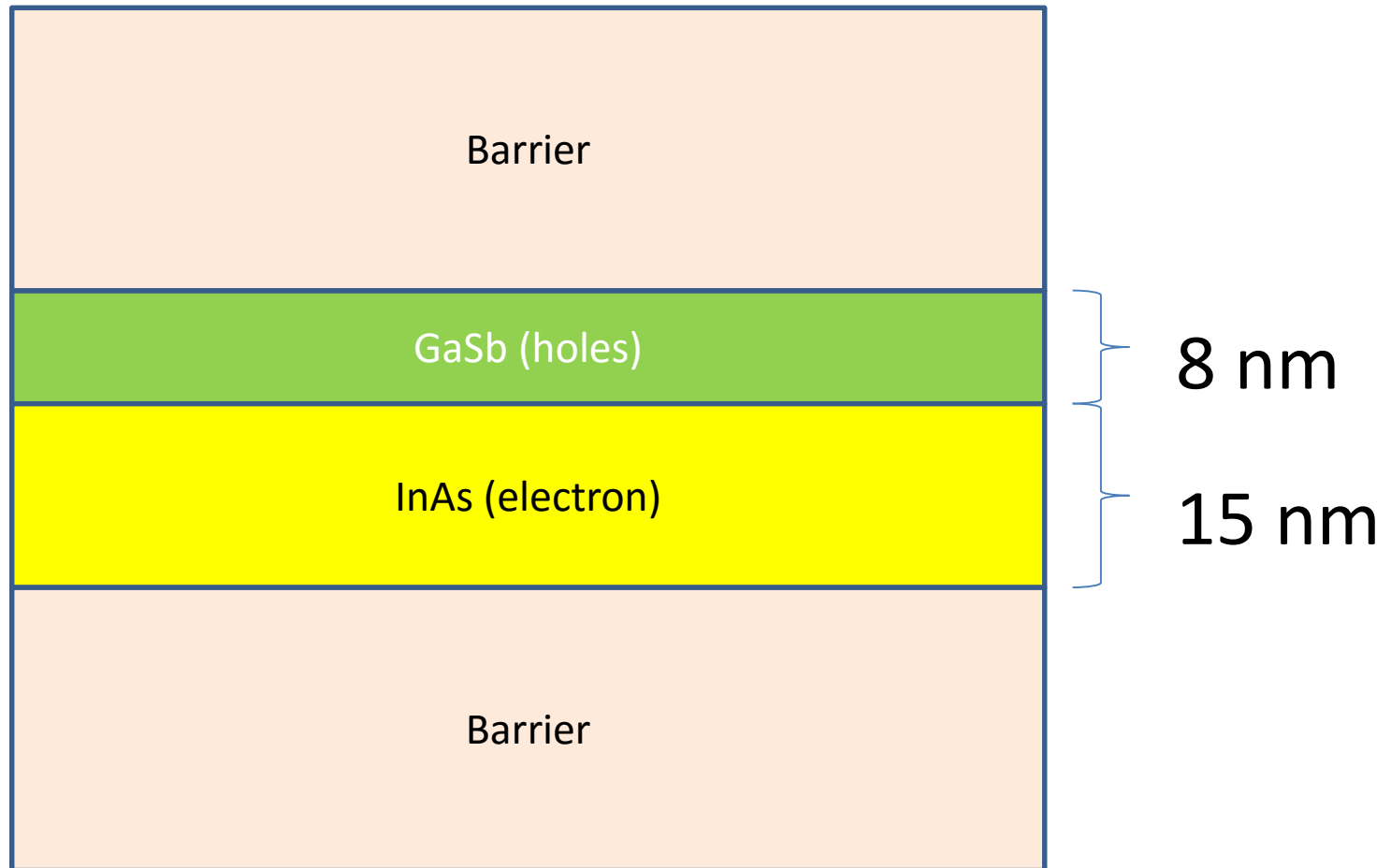


Liu *et al.* PRL 100 236601 (2008)

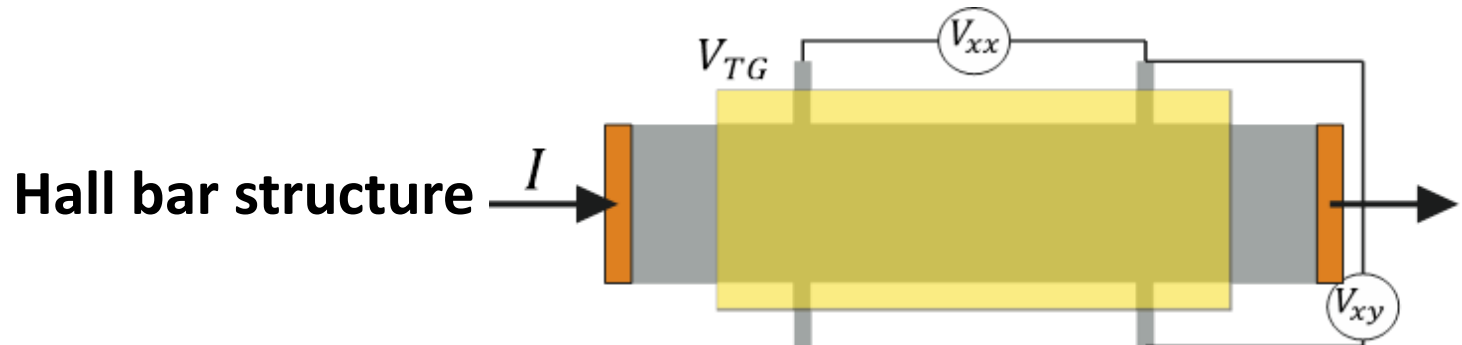
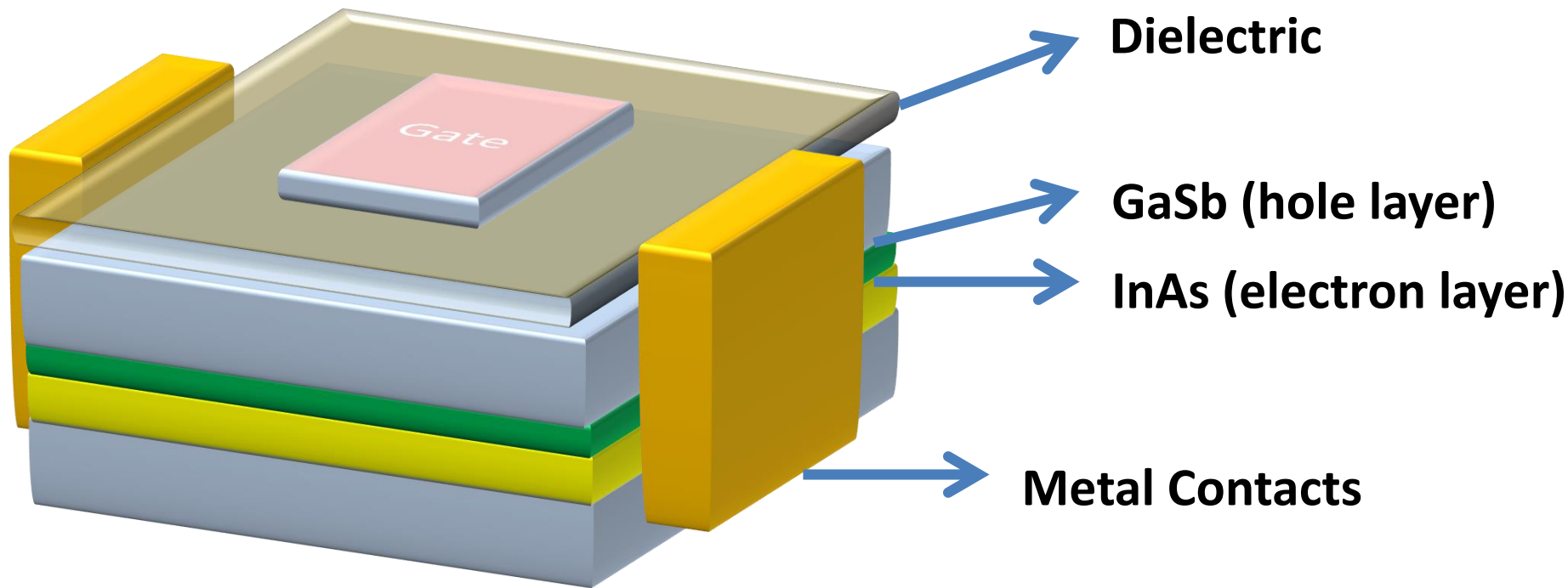
# Questions.....

- **Electron-hole transport and their hybridization?**
- **What is effect of disorder?**
- **What is effect of magnetic field?**
- **Can we observe Quantum spin Hall edge states?**

# Layer by layer growth by MBE

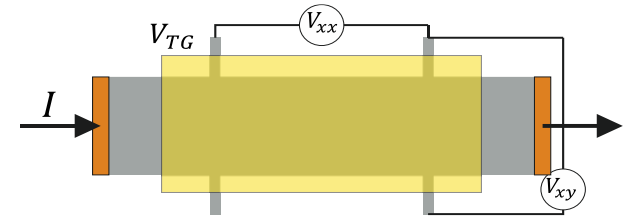


# Basic device structure

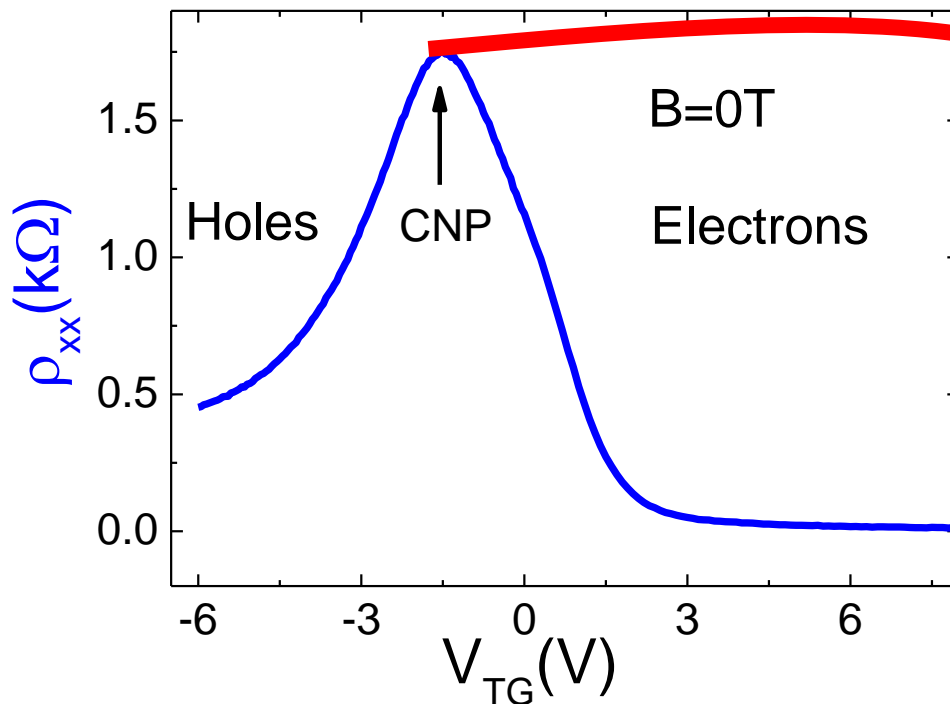


# Electron-hole transport: Pseudographene!!

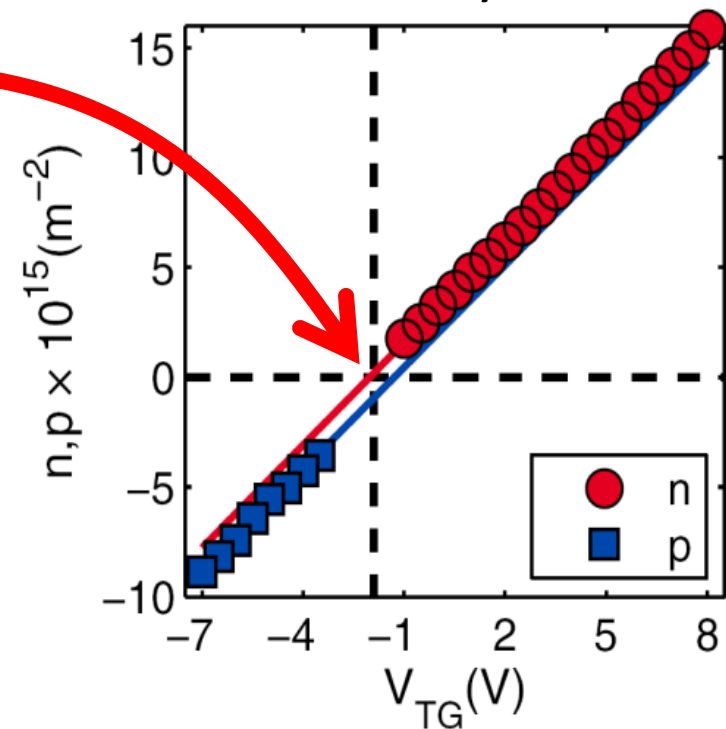
25  $\mu\text{m}$  x 50  $\mu\text{m}$  Hall bar



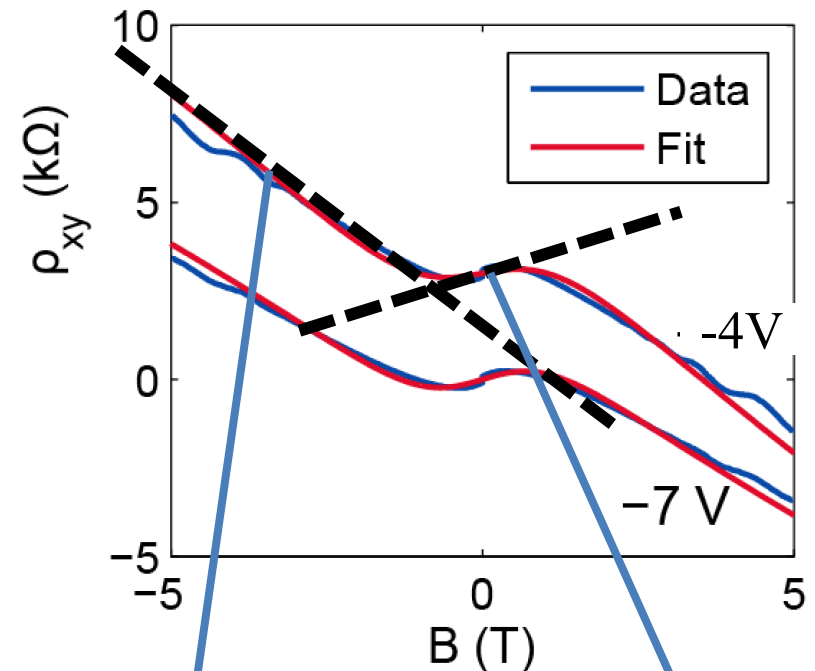
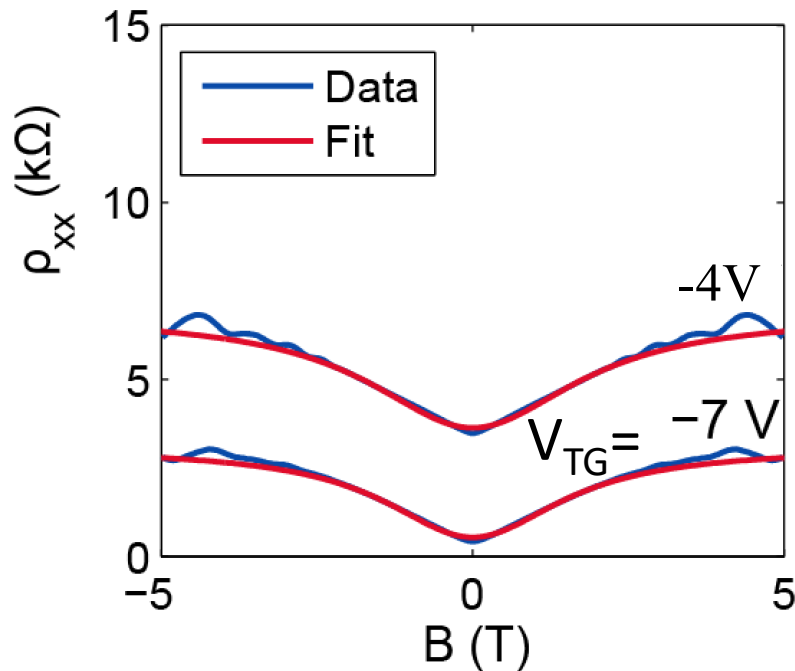
Longitudinal resistivity



Hall Density



# Low-field magnetoresistivity in perpendicular magnetic field



Classical two-band model works  
Electrons and holes coexist

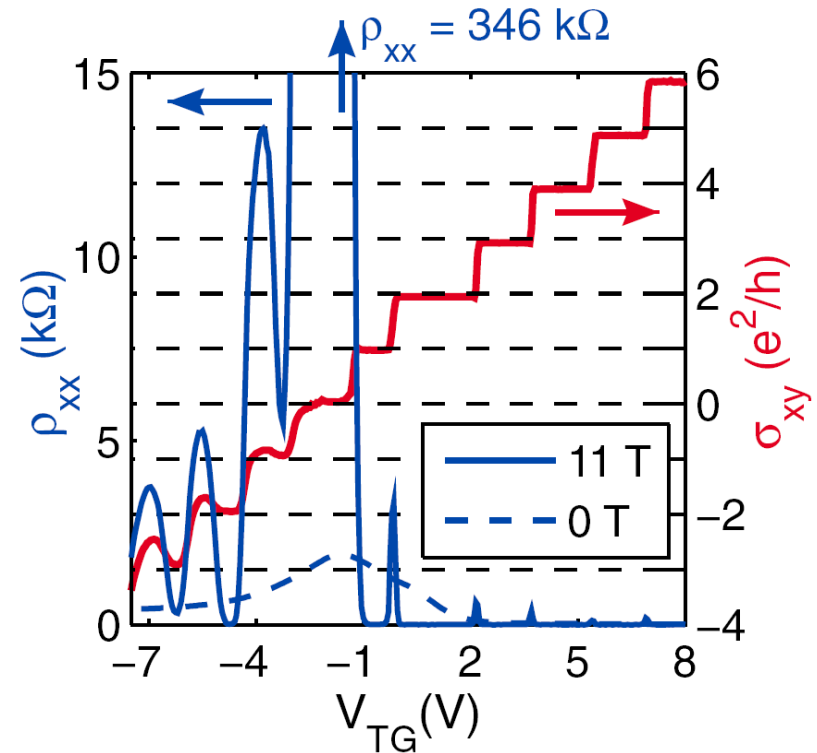
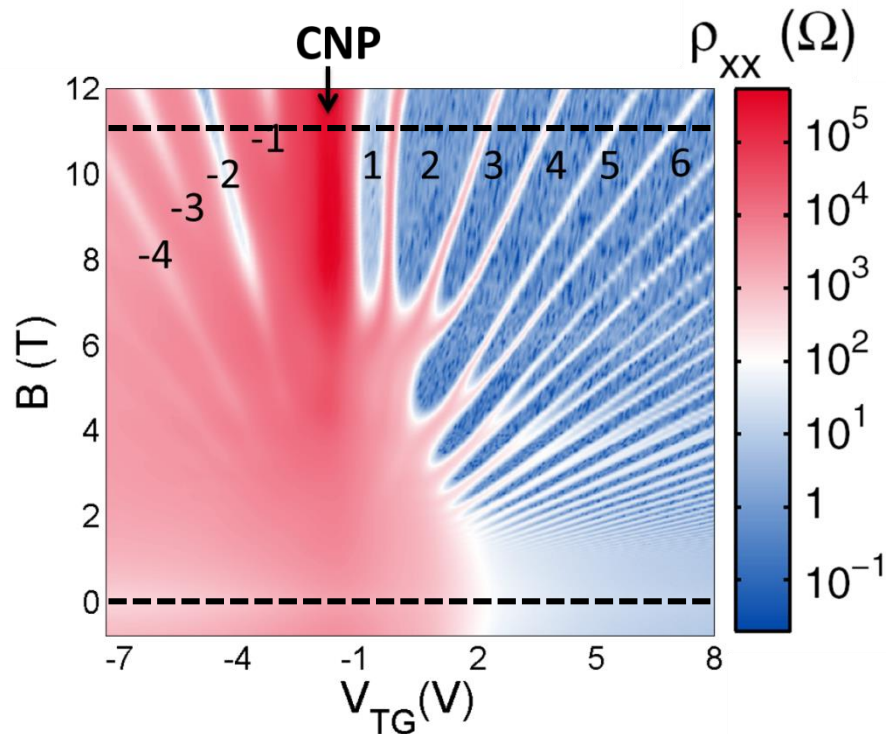
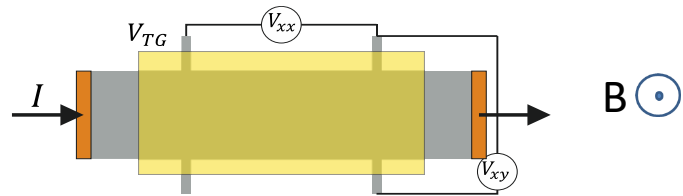
E Zaremba, Phys. Rev. B 45, 14143 (1992).

holes

electrons



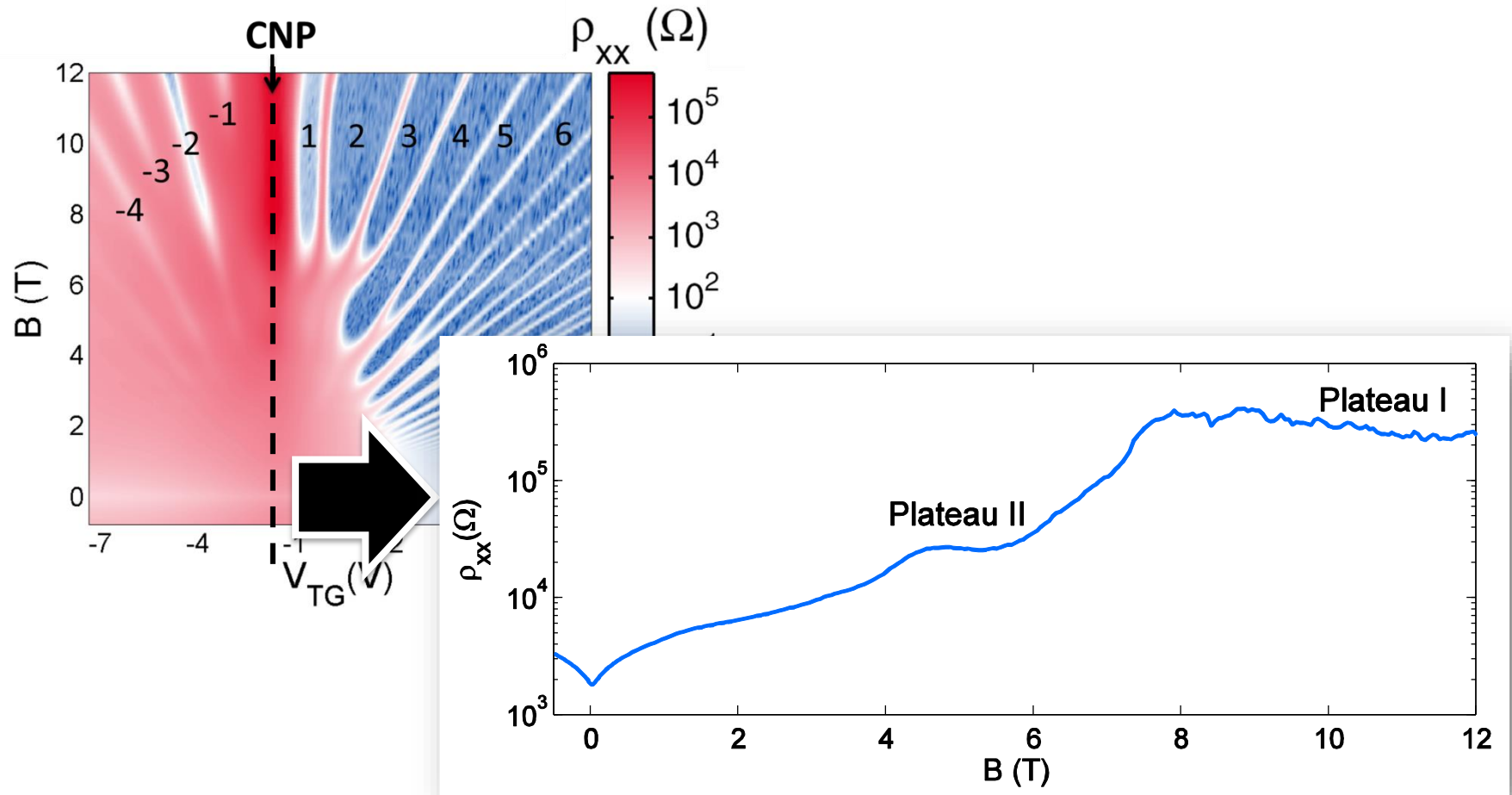
# Quantum Hall Effect of Electrons and Holes



- Resistance peak and  $\nu = 0$  plateau at the CNP : more than 10 times higher than resistance quantum ( $\hbar/e^2 = 25.8 \text{ k}\Omega$ )

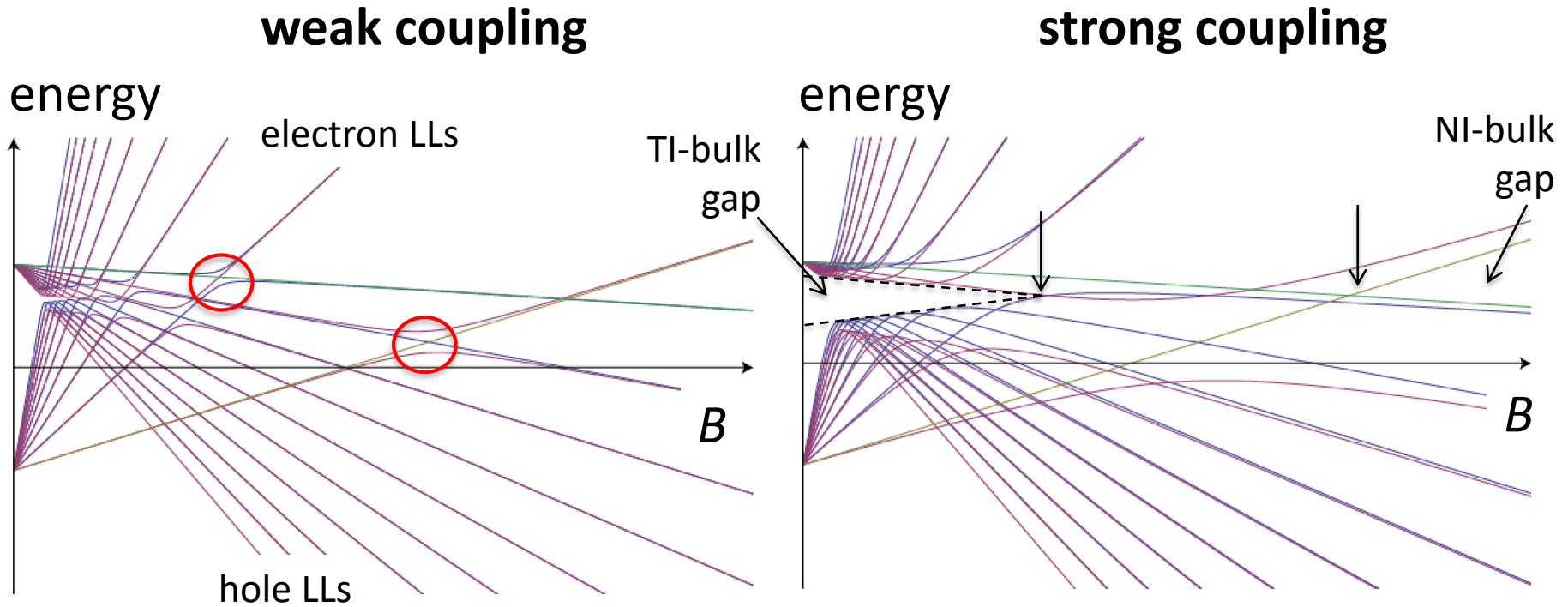
# Resistance Peak at CNP

(charge neutrality point)



- Resistance peak and  $\nu = 0$  plateau at the CNP
- Plateaus in  $\rho_{xx}$  at the CNP as a function of magnetic field

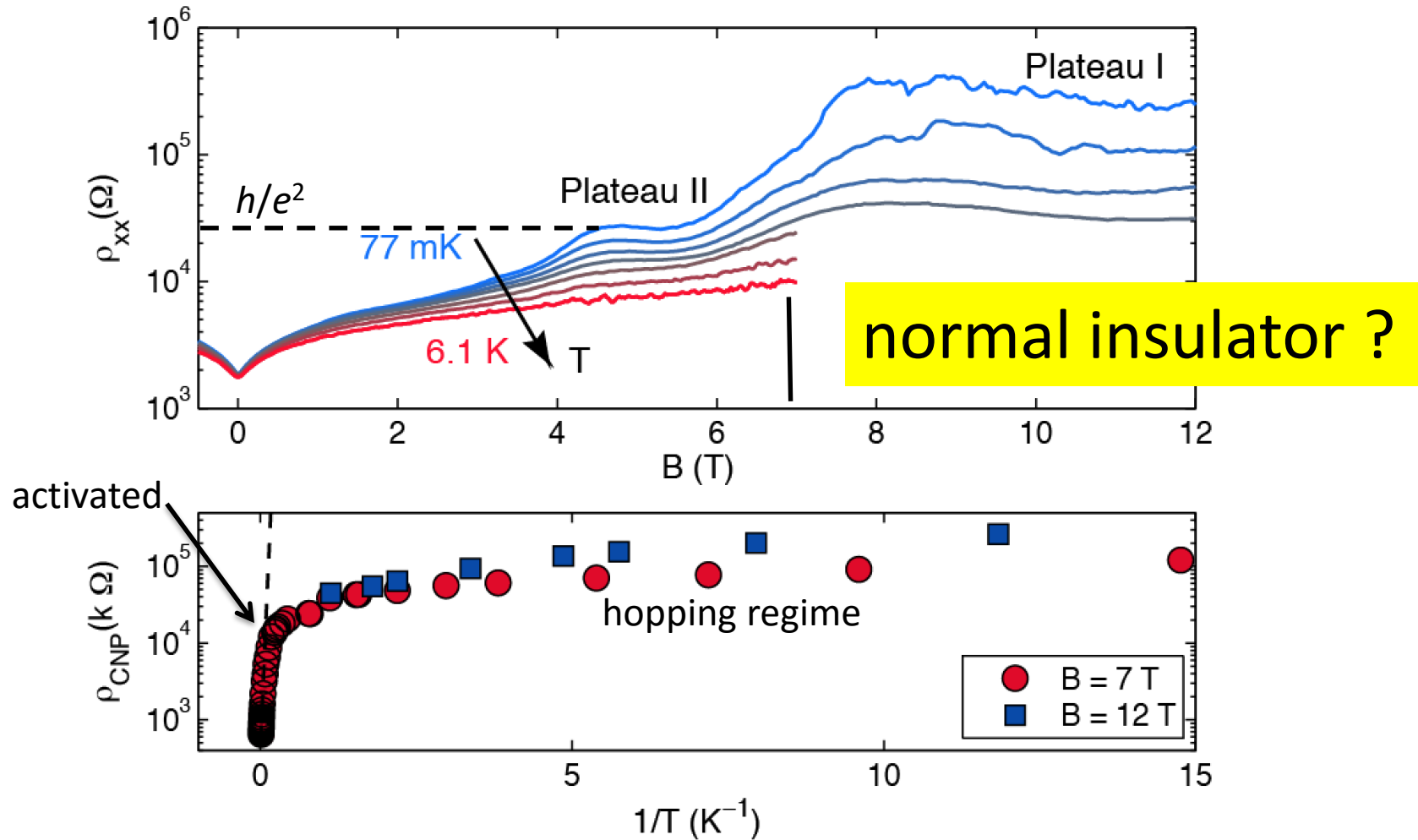
# Landau-level coupling



Interlayer tunneling and spin-orbit interaction  
couple electron- and hole-Landau levels

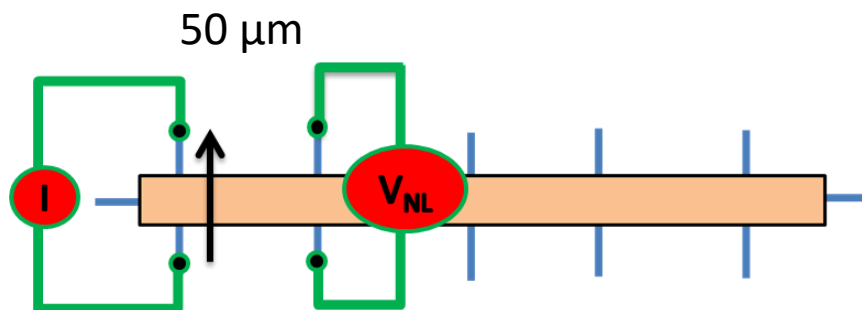
e.g. Wagner *et al*, Superlattices and Microstructures **21**, 95 (1997).  
D. I. Pikulin *et al*, Phys. Rev. B **89**, 161403 (2014).

# Resistance Peak at CNP

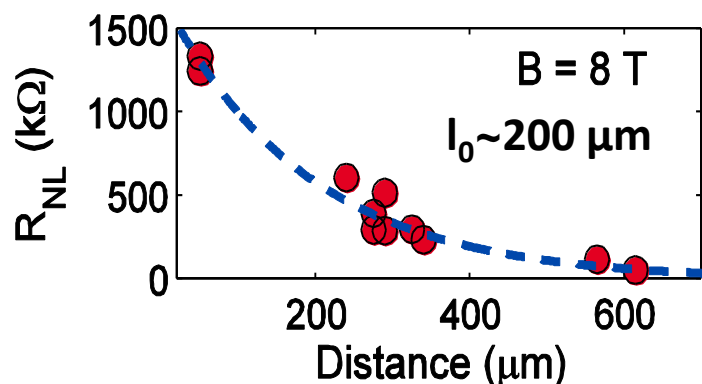


Magnetic field perpendicular to the plane of the quantum well

# Non-local Transport at the CNP



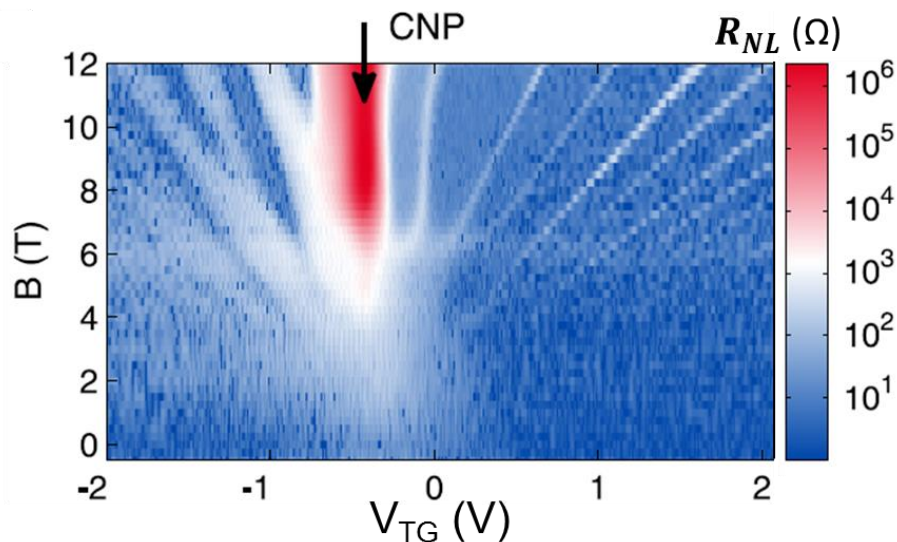
$$R_{NL} = \frac{V_{NL}}{I}$$



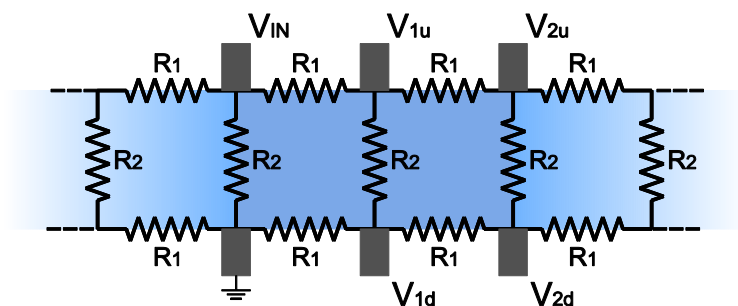
- The non-local resistance follows an exponential decay:

$$R_{NL}(x) = R_0 e^{-x/l_0}$$

- Diffusive transport :  $l_0 = W/\pi = 8 \mu\text{m}$



- Simple resistive network model

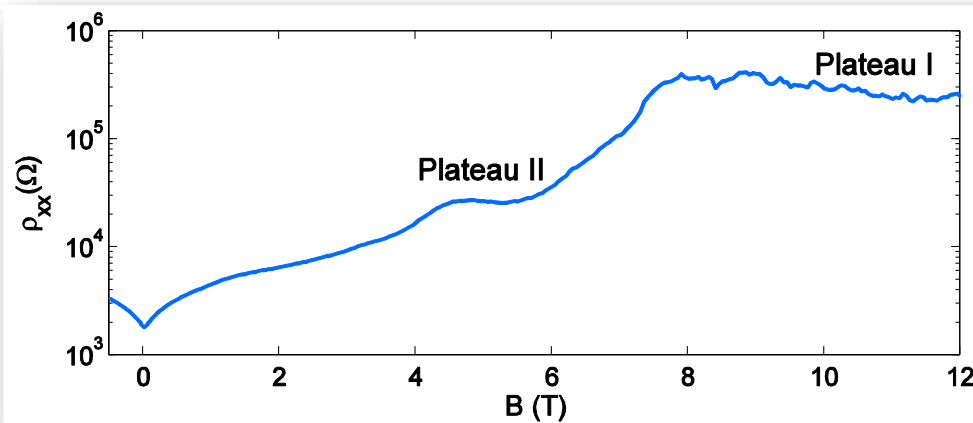
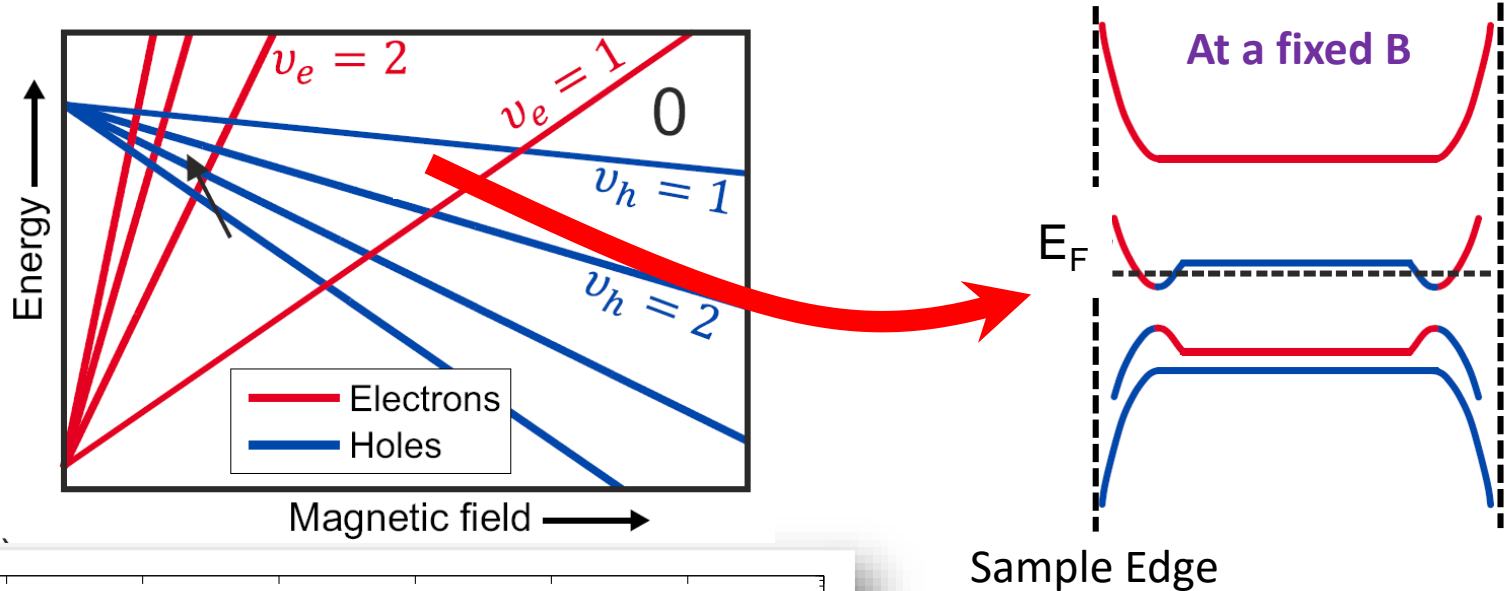


$$\frac{R_1}{R_2} = \frac{1}{100}$$

$R_1$ : Edge channel resistance  
 $R_2$ : Bulk resistance

# A qualitative model

- At finite field, electron and hole Landau levels can hybridize
- The formation of energy gaps is masked by disorder potential



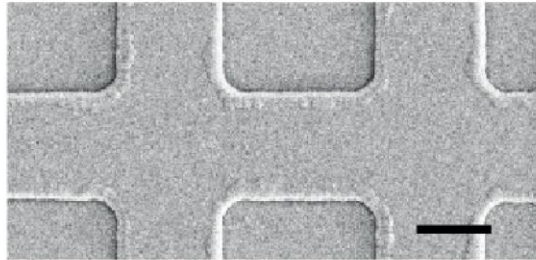
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Part II:

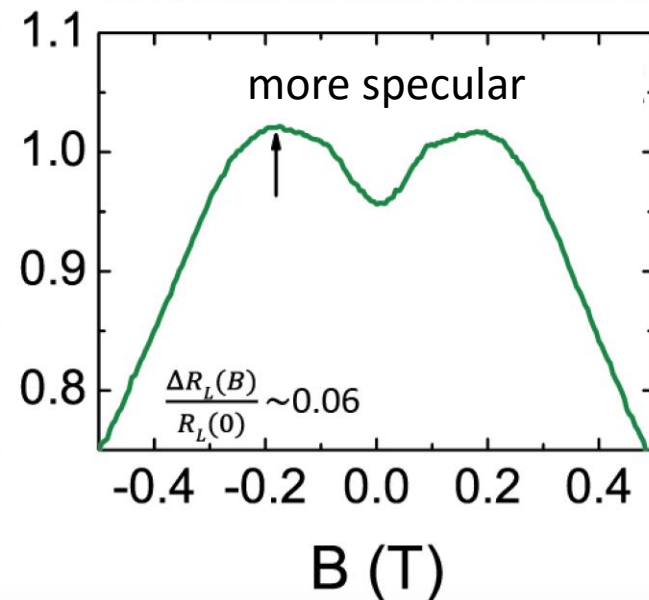
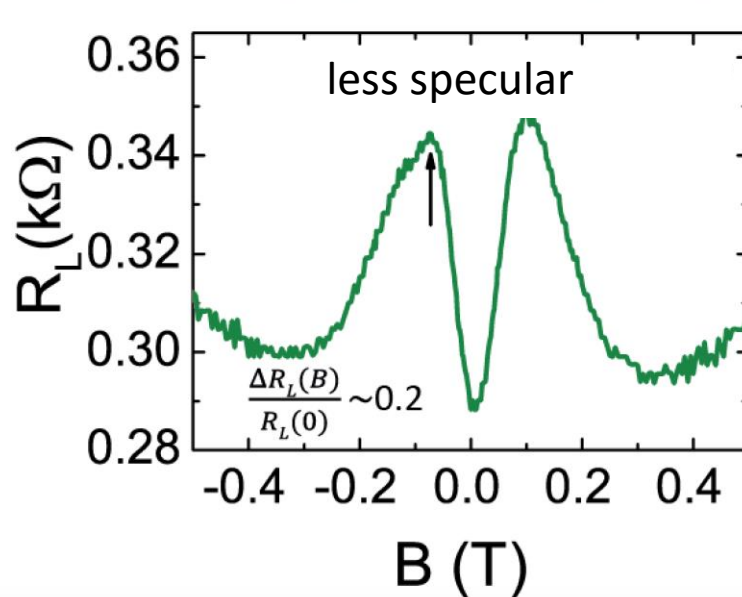
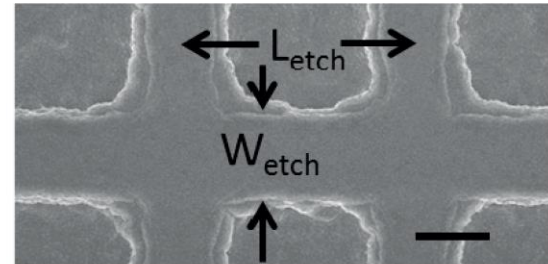
Towards quantum spin Hall effect

# Fabrication of mesoscopic Hall bars

Dry etched



Wet etched

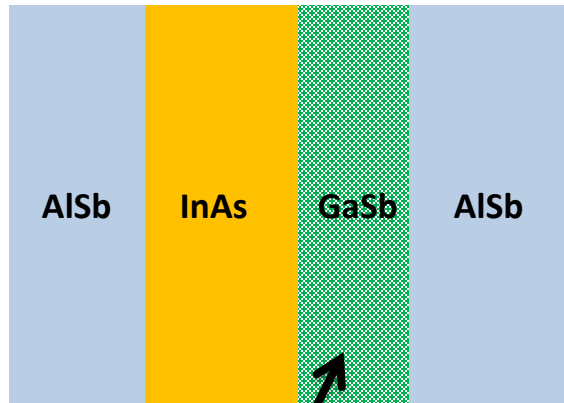


devices passivated by 200 nm of  $\text{Si}_3\text{N}_4$



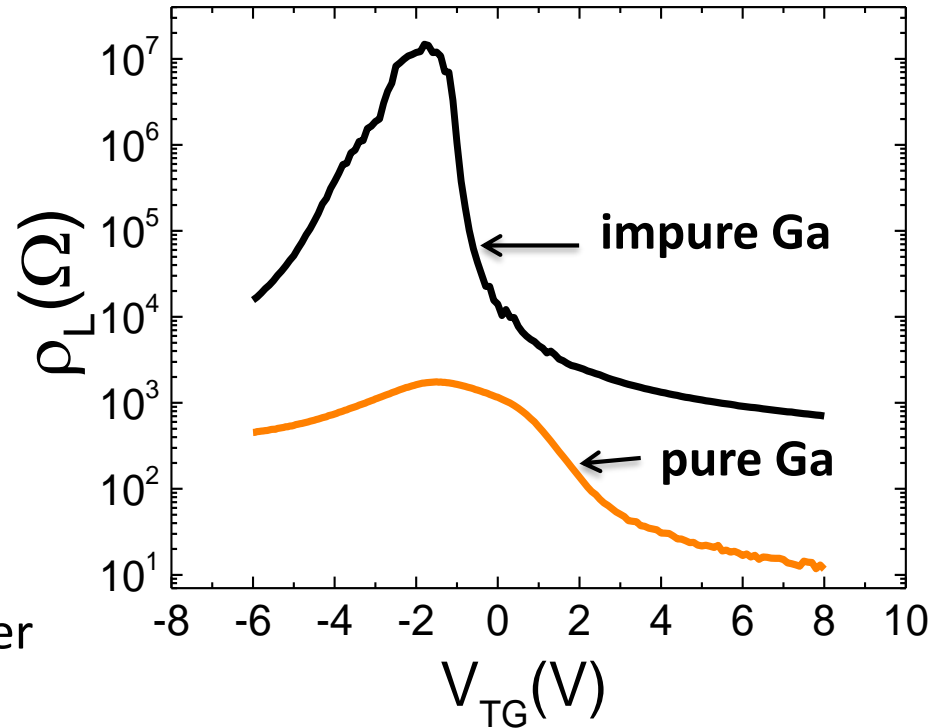
# Reducing bulk conduction in InAs/GaSb

- Small hybridization gap (few meV).
- Wave functions do not 100% overlap:  
Residual e and h



Introducing impure **Ga** in the **GaSb** layer

**D1: 25  $\mu\text{m}$  x 50  $\mu\text{m}$  Hall bar**



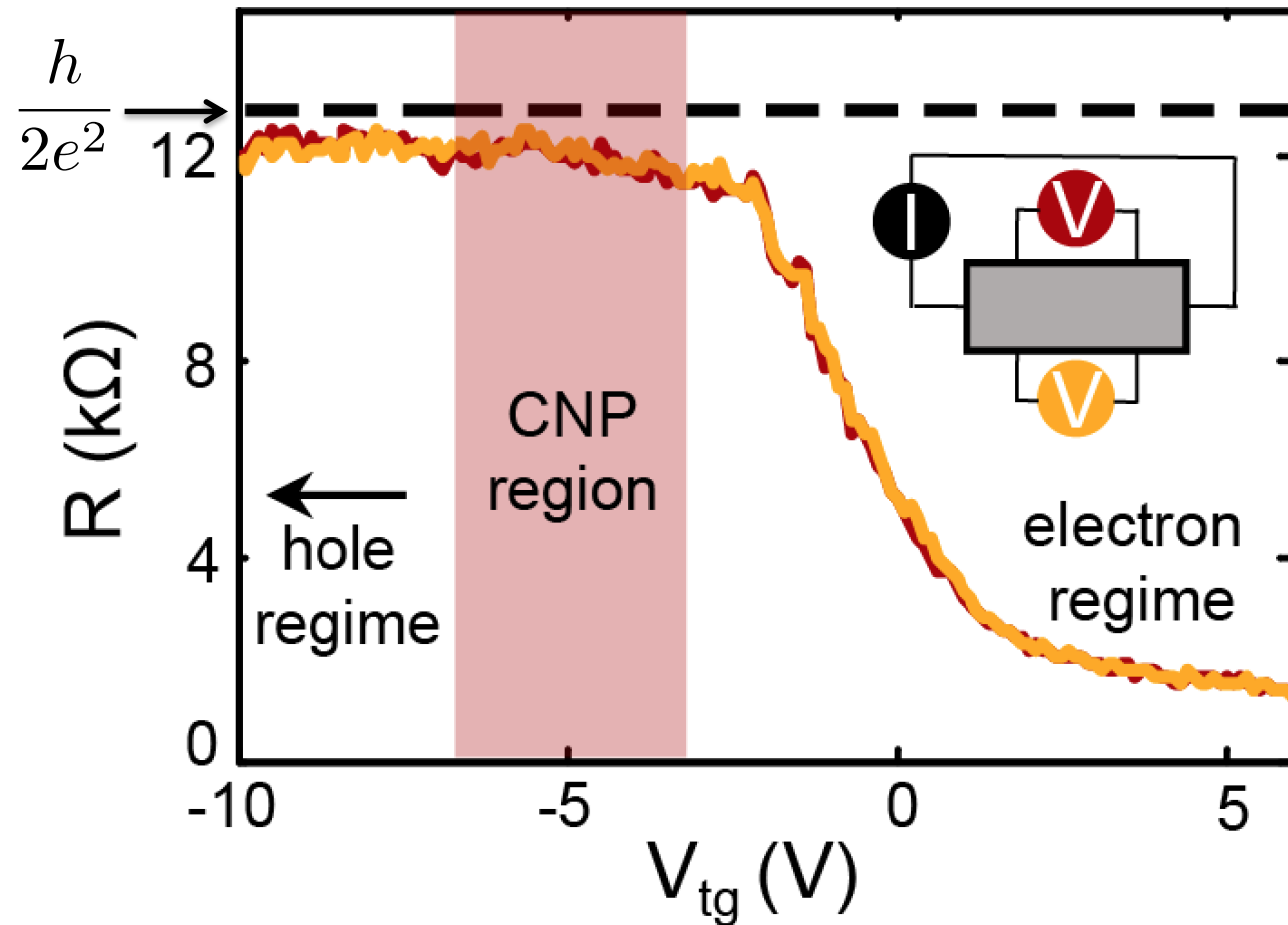
Charpentier *et al.* APL 103, 112102 (2013)

Du *et al.* PRL 114, 096802 (2015)

Dong-Hui Xu *et al.*, PRB 89, 195104 (2014)

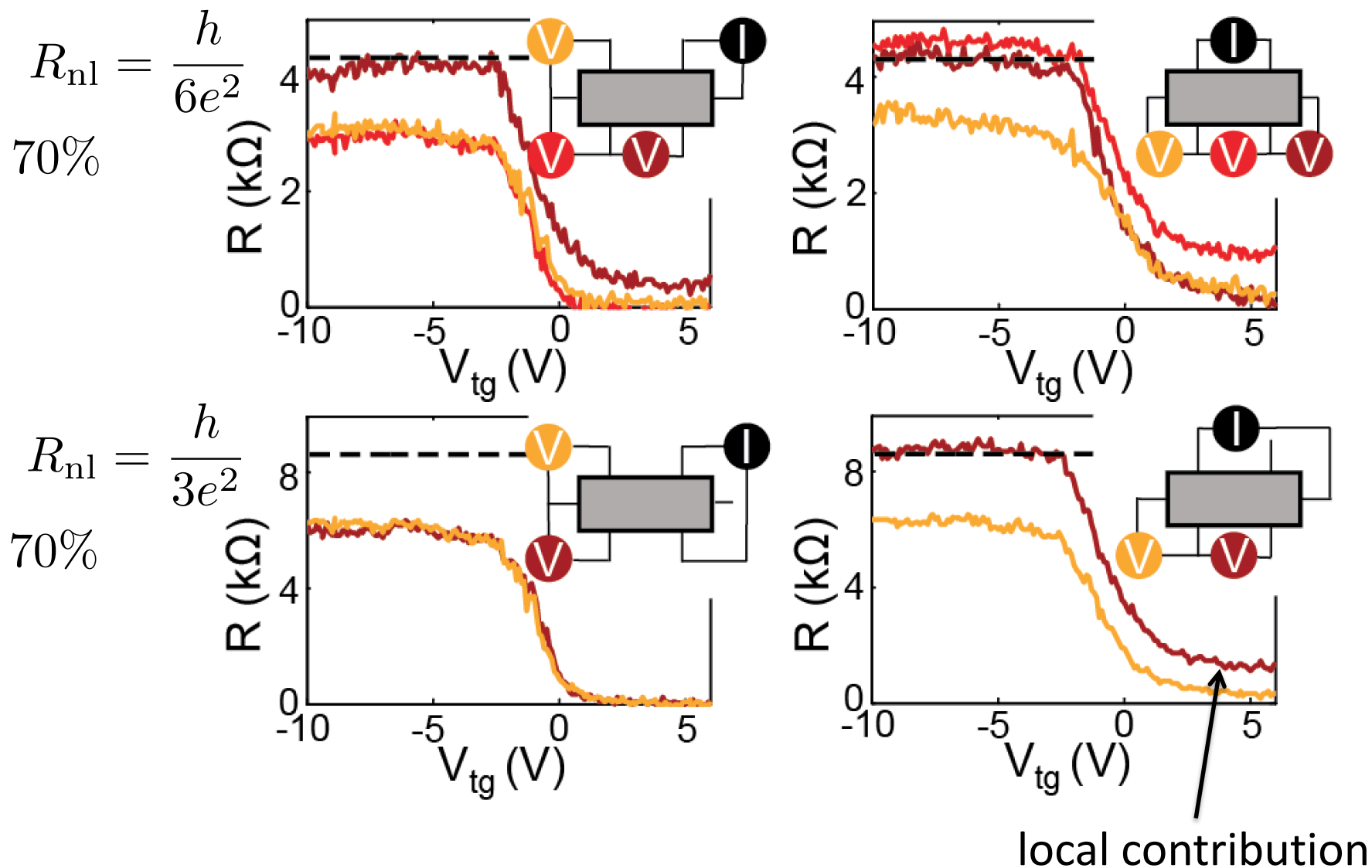
Suzuki *et al.* PRB 87, 235311 (2013)

# Small devices – four-terminal resistance



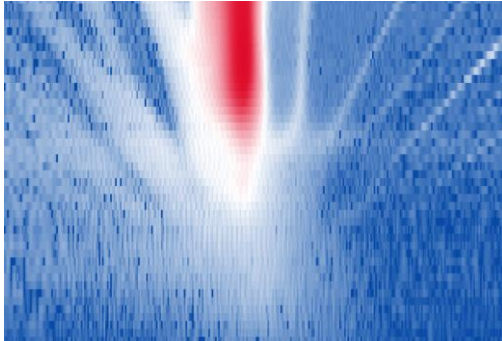
➔ Plateau value close to expected  $h/2e^2$ : helical edge transport?

# Non-local resistance – systematics



# Conclusion

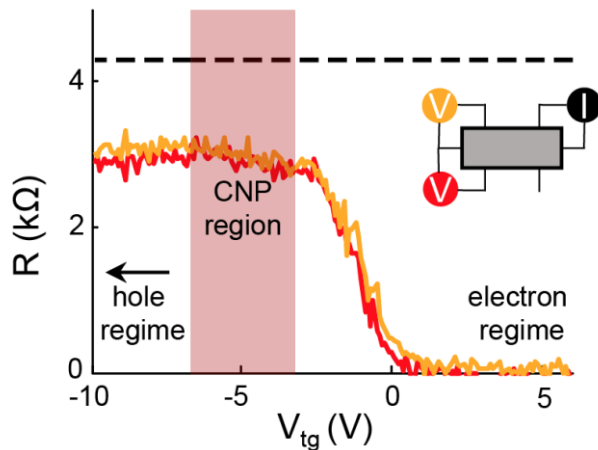
## Magnetotransport in high-mobility InAs/GaSb quantum wells



Giant non-local response at the CNP  
at high magnetic fields mimicking helical  
edge channels

F. Nichele, A N Pal *et al*, PRL **112**, (2014).

## Hunting topological insulator behavior in lower mobility samples



A glimpse of helical edge channel  
transport at zero magnetic field

Non-local resistivities are systematically low

S. Mueller, A N Pal *et al* PRB (R) 92, 081303(R) (2015)

# Thank you for your attention!

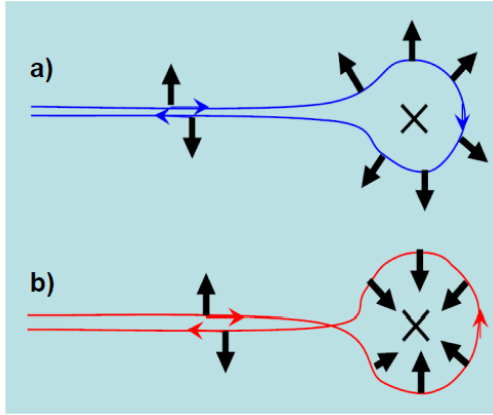


Figure 2.7 : Scattering of a quantum spin Hall edge state around a non-magnetic impurity. Due to time reversal symmetry the spin rotates by  $\pi$  in the clockwise direction and  $-\pi$  in the counterclockwise direction, with a total change in phase of  $2\pi$ . Upon  $2\pi$  spin rotation the wavefunction changes sign and the two paths interfere destructively [5].

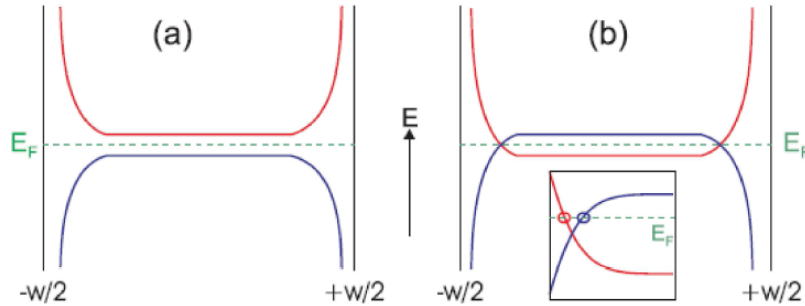
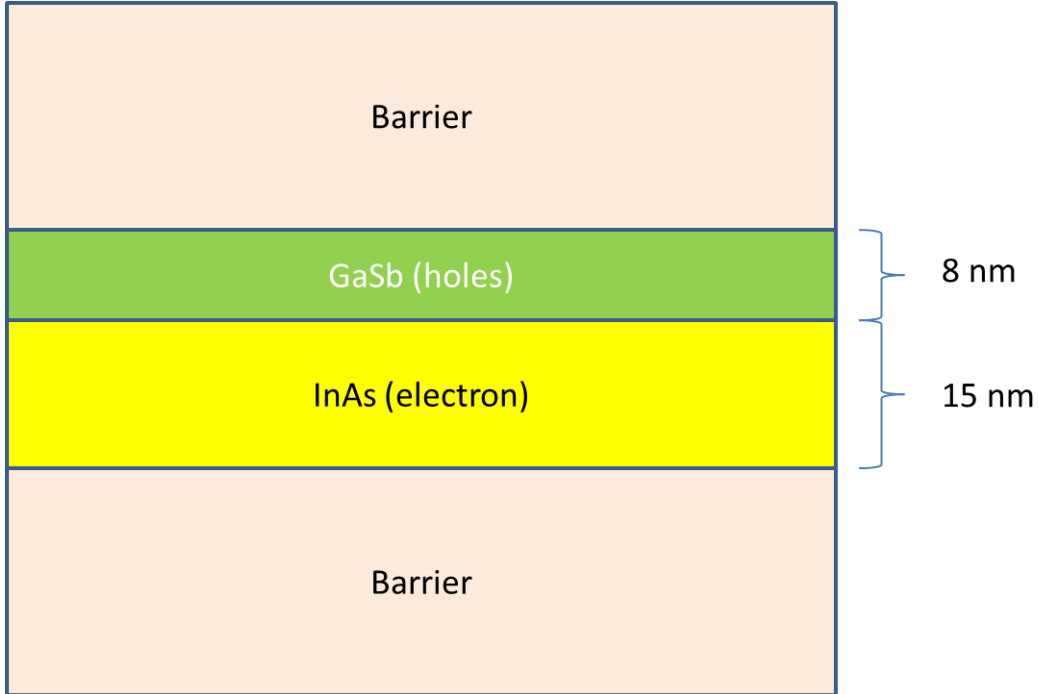
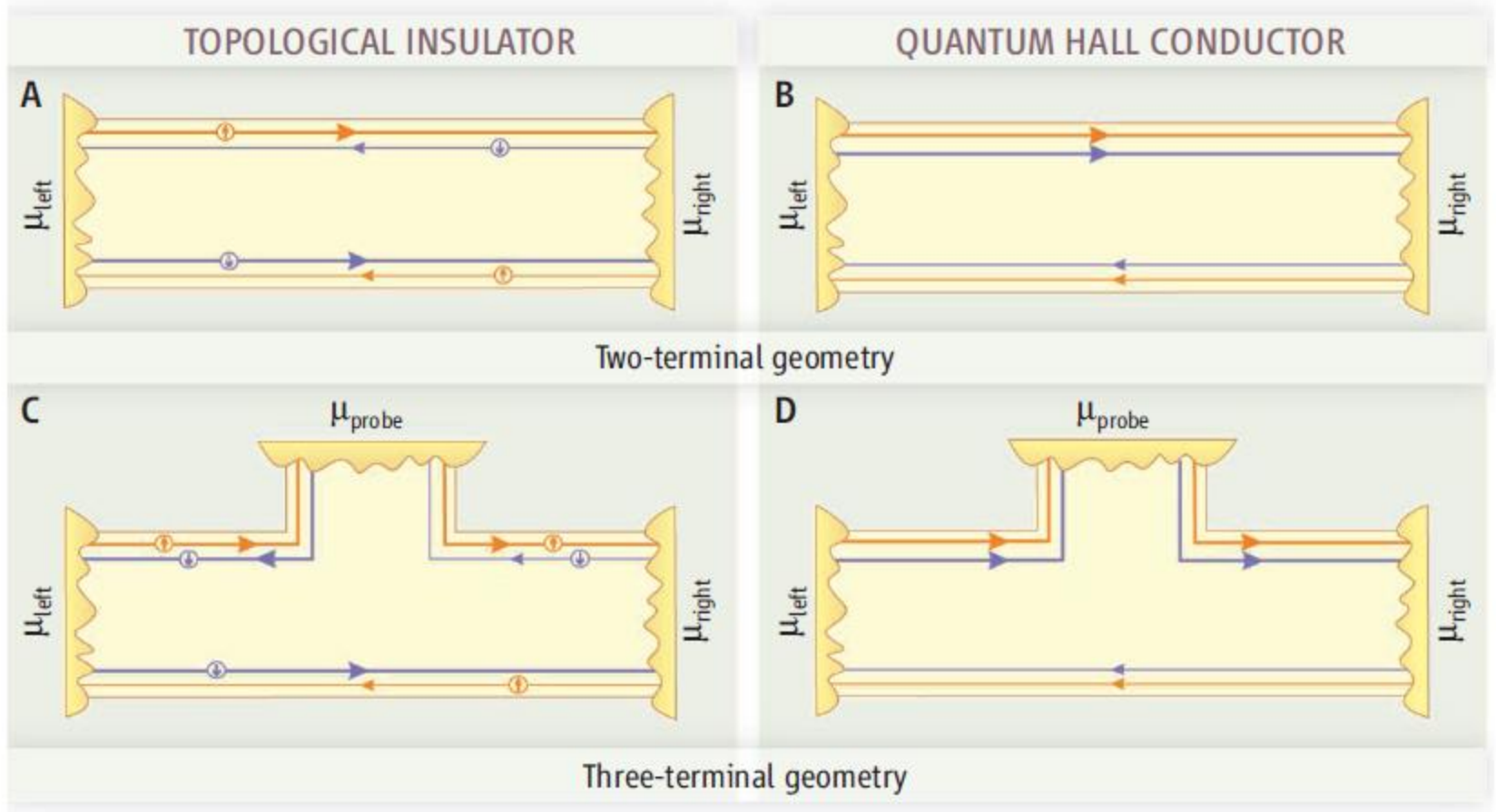


Figure 2.6 : Energy spectrum in real space for normal regime in a) and inverted regime in b). The confinement potential for electrons and holes has opposite direction. For the inverted structure and Fermi level  $E_F$  in the gap,  $E_F$  crosses energy states at edges, while for the normal regime  $E_F$  stays in the gap for the entire width of the sample [58].



# QHE and QSHE



Markus Büttiker, Science, 325, 278 (2013)

