

Revisiting the Szilard engine

Abhishek Dhar

International centre for theoretical sciences (TIFR), Bangalore

www.icts.res.in

Deepak Bhat (ICTS)

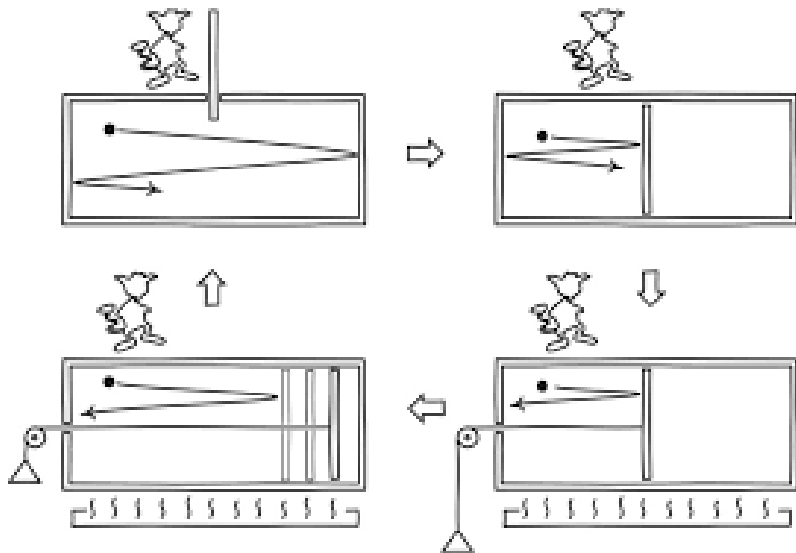
Anupam Kundu (ICTS)

Sanjib Sabhapandit (RRI)

ISPCM

Feb 17-19, 2017

The Szilard engine



The Szilard engine- Apparent violation of second law

- Work done by one-particle gas:

$$W = \int_{V/2}^V dV \frac{k_B T}{V} = k_B \ln 2$$

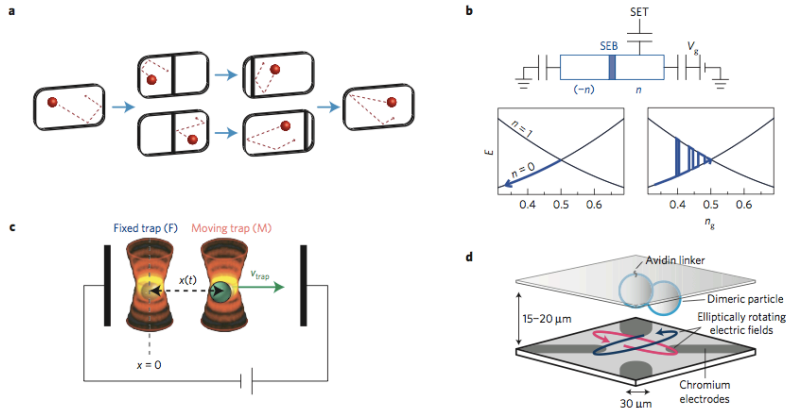
- Wall is thermal so that particle remains at constant temperature and heat flows into the system during the expansion.
- Energy conservation implies \rightarrow Heat absorbed BY system $Q = W = k_B T \ln 2$.
- Cyclic process. System unchanged at end of process. Change in entropy of the universe

$$\Delta S = -k_B \ln 2 < 0$$

Violation of second law.

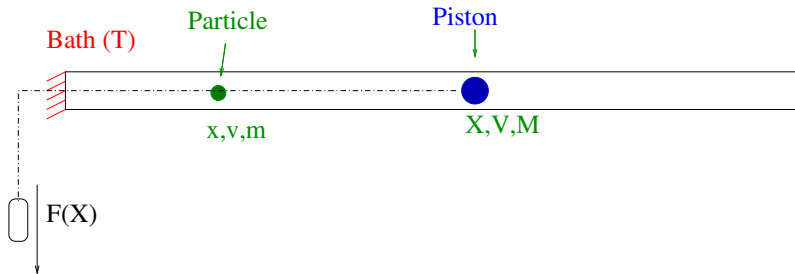
- Popular resolution: measurement leads to information of 1 bit which corresponds to generation of entropy $k_B \ln 2$.

The Szilard engine



Review article: Thermodynamics of information: Nature Physics (2015).

- Forget philosophical aspects....look at practical aspects!
- A microscopic understanding of the motion of the piston — In particular we expect that fluctuations are important in this “small” system and we want to understand their role in the motion of the piston.
- Focus on the following one-dimensional version



- $M \gg m$, X, V slow degrees, x, v fast degrees.

- CLAIM — Langevin equation of motion for the piston (heavy mass)

$$M \frac{dV}{dt} = -\frac{\partial U}{\partial x} + \frac{k_B T}{X} - \frac{\alpha}{X} V + \left(\frac{2\alpha k_B T}{X} \right)^{1/2} \eta(t),$$

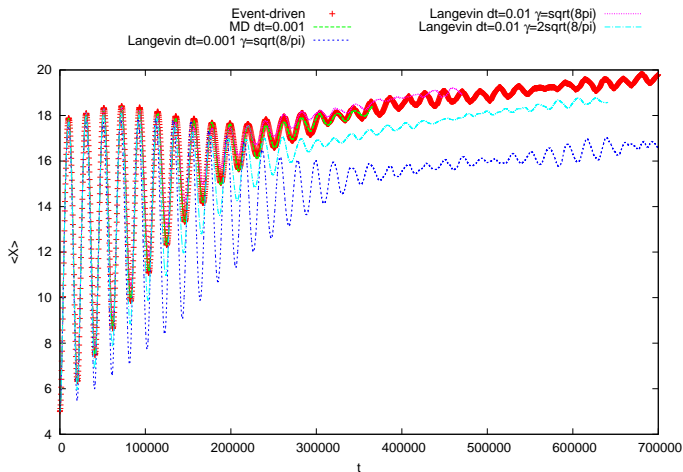
$$\text{where } \alpha = c \left(\frac{8k_B T m}{\pi} \right)^{1/2},$$

and $\eta(t)$ is Gaussian white noise.

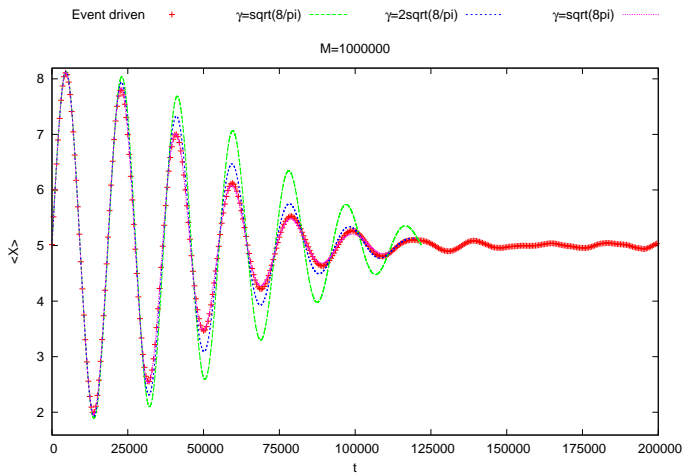
- Naive theory gives $c = 1$.
- Simulations give $c = \pi$.
- Derivation from the master equation for $P(x, X, v, V, t)$.
 - Eliminate fast variables (x, v) .
 - Write effective Fokker-Planck equation for $P(X, V, t)$.
 - Read off the corresponding Langevin equation.
- OTHER results:** Energetics of the Szilard engine, optimum protocols.

Comparison of direct simulations with solution of Langevin equations.

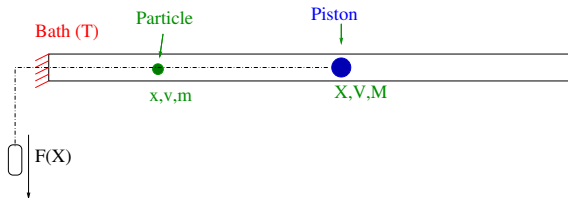
Example (I)— Szilard-engine case



Example (II) — Three particle case



Master equation approach



$$\begin{aligned} \frac{\partial P(x, X, v, V, t)}{\partial t} = & -v \frac{\partial P}{\partial x} - V \frac{\partial P}{\partial X} + \frac{U'(X)}{M} \frac{\partial P}{\partial V} \\ + & \delta(x) \int \left\{ [-v' \theta(-v') \phi(v)] P(x, X, v', V, t) - [-v \theta(-v) \phi(v')] P(x, X, v, V, t) \right\} dv' \\ + & \delta(X - x) \left\{ \theta(V - v)(V - v) P \left[x, X, \frac{(\mu - 1)v + 2V}{1 + \mu}, \frac{(1 - \mu)V + 2\mu v}{1 + \mu}, t \right] \right. \\ & \left. - \theta(v - V)(v - V) P(x, X, v, V, t) \right\}, \end{aligned}$$

where $\phi(v) = \sqrt{\beta m v} e^{-\beta m v^2 / 2}$ — distribution of particles emerging from a thermal bath and $\mu = m/M$.

GREEN — collision of small particle with thermal wall.

BLUE — collision of small particle with piston.

Master equation approach

- Introduce new $O(1)$ variables $u = \sqrt{\beta m}v$, $y = \sqrt{\beta M}V$, $\tau = t/\sqrt{\beta m}$ and let us denote the small parameter in the problem as $\epsilon = \sqrt{m/M}$.
- Small ϵ -expansion of the master equation results in:

$$\frac{\partial P(x, X, u, y, \tau)}{\partial \tau} = \mathcal{L}^{(0)}P + \epsilon \mathcal{L}^{(1)}P + \epsilon^2 \mathcal{L}^{(2)}P.$$

- Then write

$$P(x, X, u, U, \tau) = \Phi(x, u, \tau | X, U, \tau) \Psi(X, U, \tau)$$

$$\Phi = \Phi_0 + \epsilon \Phi_1 + \epsilon^2 \Phi_2$$

$$\Psi = \Psi_0 + \epsilon \Psi_1 + \epsilon^2 \Psi_2$$

- Write equations at each order, use solutions to establish equation at next order.
- Finally write Fokker-Planck equation for $\Psi(X, U, \tau)$.

.....Solution in time-Laplace transform possible.....inversion problem remains!

MORE DETAILS — Deepak Bhat's Poster.