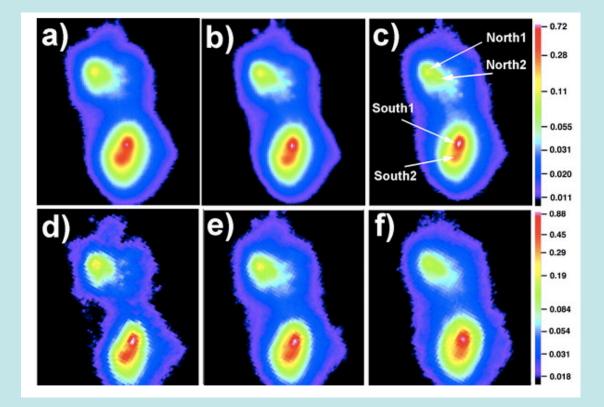
## The formation of the first black holes



Priyamvada Natarajan Departments of Astronomy and Physics, Yale

ICGC, Goa, India December 19, 2011

## Talk outline

- The context what do we know for certain about the accretion history of BHs?
- Observational evidence for BHs
- How did the first BHs form?
- How did they grow?
- Discriminating Pop III seeds from massive seeds
- Evidence for massive black hole seeds from high and low redshift
- Future prospects

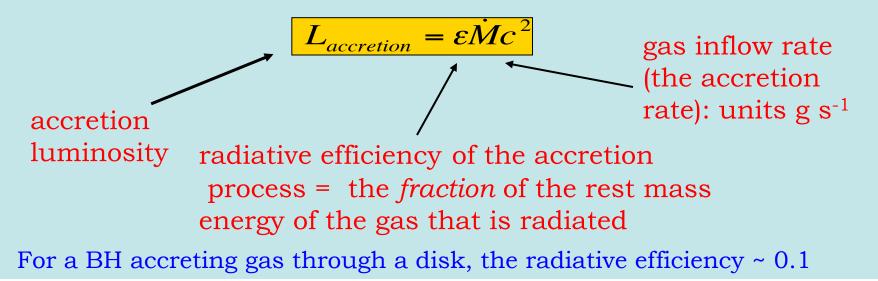
<u>Collaborators</u>: Giuseppe Lodato, Ezequiel Treister, Marta Volonteri, Andrew Davis

#### BH accretion: what we do know

- The critical luminosity when radiation pressure balances gravity is the Eddington limit
- Depends on mass of accreting object and opacity of the surrounding gas

$$f_{rad} = f_{grav}; \frac{\kappa L}{4\pi cr^2} = \frac{GM}{r^2}; L_{Edd} = \frac{4\pi cGM}{\kappa_T} = 1.25 \times 10^{45} \left(\frac{M}{10^7 M_{sun}}\right) \text{ erg s}^{-1}$$

Gas flowing toward black holes produces radiation as the gravitational potential energy is released



## <u>Characteristic growth timescale</u>

- Bright quasars must have  $M > 10^8$  Msun
- Eddington limit caps growth rate of mass

$$\frac{dM}{dt} = \frac{L_{acc}}{\eta c^2} < \frac{4\pi G M m_p}{\eta c \sigma_T}$$

$$M \le M_0 e^{\tau}$$

$$\tau = \frac{\eta c \sigma_T}{4\pi G m_p} \approx 5 \times 10^7 \, yr$$

Salpeter timescale

## Energy output from accreting BHs

• Accretion=> Gravitational Energy=> EM radiation If all energy is thermalized, Black body with temperature

$$T_b = \left\{\frac{L_{acc}}{4\pi R^2 \sigma}\right\}^{\frac{1}{4}}$$

For SMBHs:

$$M \approx 10^8 M_{sun}; R = \frac{2GM}{c^2} \approx 3 \times 10^{13} cm$$

$$T_b \approx 10^7 \left(\frac{M_{sun}}{M}\right)^{\frac{1}{4}} K \approx 10^5 K$$

Expect accreting BHs to be UV, X-ray and possibly gamma-ray emitters

### Approach to building accretion histories

• <u>Globally averaged constraint</u> (Soltan 1982) census of the total energy emitted by a population of accreting BHs

$$E = \iint E(L,t)dLdt = \frac{4\pi}{c}\int (1+z)dz \int n(S,z)SdS$$

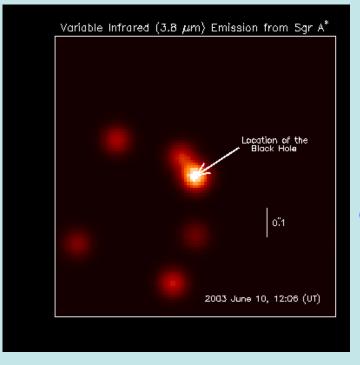
$$\rho_{BH} = \int_{0}^{t_{age}} dt \int_{0}^{\infty} d(\ln L) L \phi(L, t)$$

 <u>Continuity argument</u> for mass accumulation as a function of time (Small & Blandford 1992; Haehnelt, PN, Rees 1998, Haiman & Loeb 1999......Volonteri et al. 2005; Merloni & Heinz 2008)

attempt to relate the observed evolution of quasars to physical models of AGN

## Astrophysical environment of BHs

Empirically, two modes of accretion in galactic nuclei, X-ray binaries



UCLA Galactic Center group

• radiatively inefficient accretion e.g. Galactic Center  $\dot{m} = \dot{M} / \dot{M}_{Edd} \sim 10^{-7}$ (Shcherbakov et al. 2010; Narayan & Yi) Ubiquitous: fed by stellar winds,

"hot"  $T_{ion} \sim T_{vir}$  quasi-spherical flow

• thin disk accretion: radiatively efficient (e  $\sim$  0.1) AGN mode

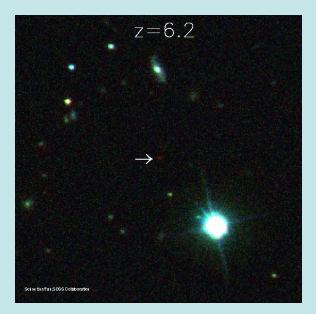
Feeding mechanism unclear, rare in the local Universe

Typical merger environment *only* if thin disks **cause** or are otherwise preferentially correlated with BH coalescences

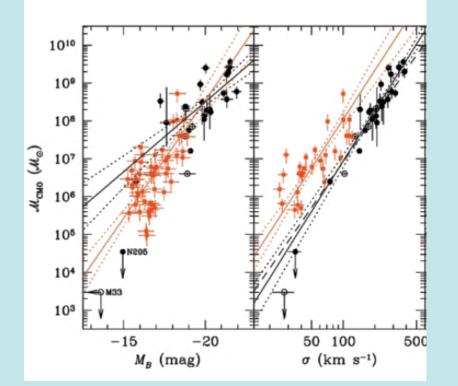
## Observations from high and low redshift

<u>Abundance & LF of</u> <u>High redshift quasars</u>

> Age of the Universe 2 Gyr!



Most recent census from SDSS and 2dF Fan+ 2007; Croom+ 2004



#### Observational estimates of masses of central massive objects

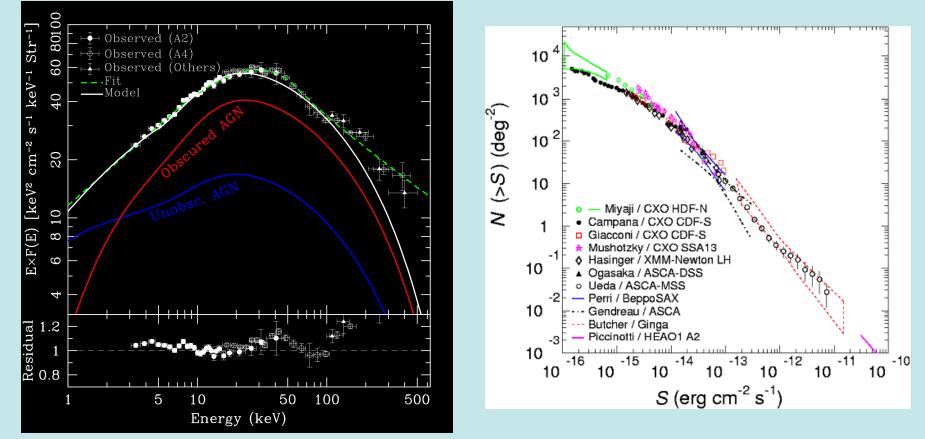
observed correlation between bulge luminosity and BH mass => BH mass and vel. Disp

Ferrarese+ 2006; Ferrarese & Merritt 2002; Tremaine+ 2002; Kaspi+ 2005

## The X-Ray Background (XRB) spectrum and

#### number counts

- The XRB is the integrated emission from all AGNs in the Universe, the energy density peaks at  $\widetilde{}$  30 keV
- The shape of the XRB indicates that most of the AGNs are obscured, harder energies needed to detect the contribution of these obscured sources



Comastri+ 1995; Ueda+ 2003; Treister & Urry 2005; Merloni+ 2004; PN & Treister 2008

#### <u>Co-evolution of galaxy and super massive black holes</u> <u>in galactic centres</u>

#### Star forming history vs BH mass vs stellar mass accretion history 0 EIIIIIII Ζ 0.3 0.5 1.0 2.0 5.0 2.0×10<sup>-</sup> black hole growth (a) Log do/dt [M<sub>®</sub> yr<sup>-1</sup> Mpc<sup>-3</sup>] dp<sub>see</sub>/dt+8.0e−4 [Fardal et al.] Ϋ́ dpm/dt+8.0e−4 [Hopkins & Beacom] .odW 1.5×10<sup>-4</sup> 1.0×10<sup>-4</sup> [M⊙ yr 5.0×10<sup>-5</sup> ΒН ·Σ -5 log 2 12 10 8 6 4 SFR (IR) - Chary & Elbaz 2001 BH Accretion Rate (this paper) t [Gyr] $^{-6}$ rescaled (x 4000. BH Acc. Rate (Barger et al. 2001 -7 [] 1..... 2 0 1 3 4

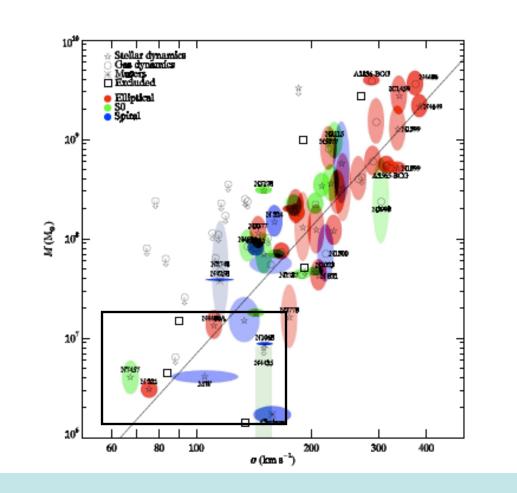
#### e.g., Marconi & Hunt 2003; Shankar, Weinberg & Miralda-Escude 2007

Marconi+ 2004

 $\mathbf{Z}$ 

5

## The $M_{bh}$ -sigma relation



Most recent estimates: Gultekin+ 2010; Greene & Ho 2009, Greene+ 2011; Reines+ 2011

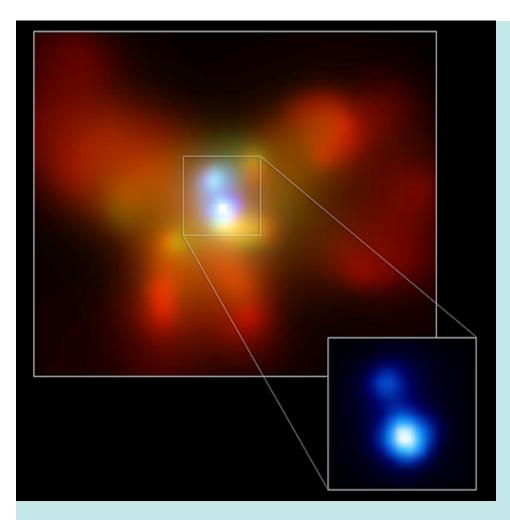
## What is the origin of M-sigma?

- Merging history of DM halos naturally provides correlation
- Slope and scatter contain information on black hole seeds and feedback processes

#### WHAT ARE THESE SIGNATURES?

Does an initial correlation between seed masses and halo masses survive the merging and accretion history of mass build up over cosmic time?

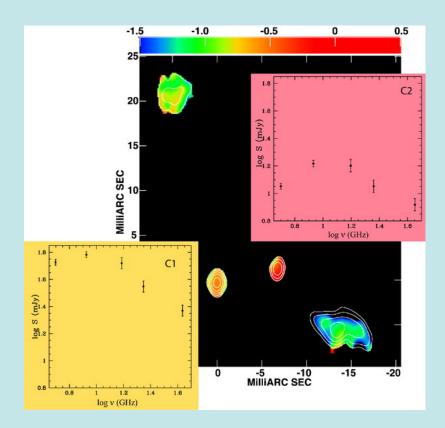
Volonteri & PN 08; Peng 07



#### NGC 6240: d ~ kpc Komossa+ 2003

## **Observationally**

Handful of identified pre-merger binary black holes



0402+379: d = 7pc (Rodriguez+ 2006)

## <u>Challenges posed by high-z quasars</u>

- Properties of high redshift quasars (10 found at z > 6 in SDSS, 20 in CFHQS; record z = 7.08 Mortlock et al. 2011)
- Luminosities suggest  $M_{bh} = L_{obs}/L_{Edd} \sim 10^{9-10}$  Msun assembled and in place when t = 0.77 Gyr
- And these are just the tip of the iceberg!
- Number of e-foldings required for a Pop III seed  $\widetilde{\phantom{a}}$  25 in the time available
- Need to be accreting steadily at Eddington or super-Eddington the entire time

## **Conditions for BH seeds**

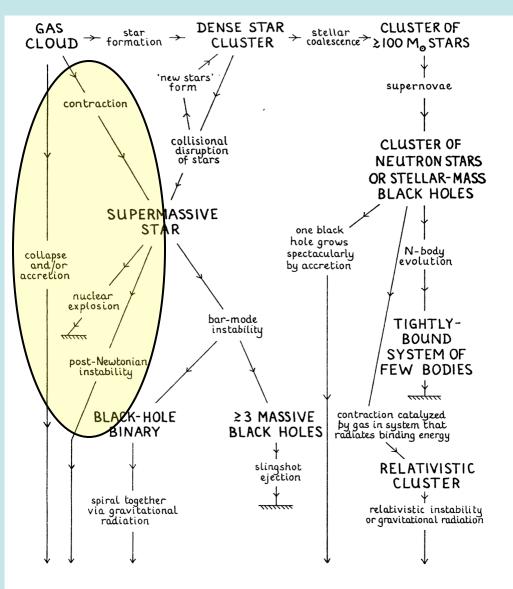
#### STELLAR SEEDS

- \*uninterrupted near-Eddington accretion onto  $^{\sim}100~M_{\rm sun}$  seed
- \* continuous gas supply
- $^{\ast}$  need to avoid radiative feedback depressing accretion rate
- $^{\ast}\,must$  avoid ejection from halos and loosing BHs

#### DIRECT COLLAPSE

- \*rapid formation from direct collapse of gas or via supermassive star, quasi-star, or ultra-dense star cluster
- \*gas must be driven in rapidly (deep potential)
- \*transfer angular momentum and must avoid fragmentation

### Pathways to making BH seeds



### massive black hole

Figure 1 Schematic diagram [reproduced from Rees (106)] showing possible routes for runaway evolution in active galactic nuclei.



#### First black holes in pre-galactic halos z ~20-30

 $M_{\rm BH} \sim 100 - 500 \ M_{\rm sun}$ 

Pop III remnants

- Simulations suggest that the first stars were massive 100 500 M<sub>sun</sub> (Bromm+ 2002; Abel+ 2002; Abel+ 2000; Alvarez+ 2008)
- Metal free Pop III stars with M > 260  $M_{sun}$  leave remnant BHs with  $M_{seed}$  > 100  $M_{sun}$  (Fryer, Woosley & Heger)

 $\rm M_{BH} \sim 10^3$  -  $10^6 \rm \ M_{sun}$ 

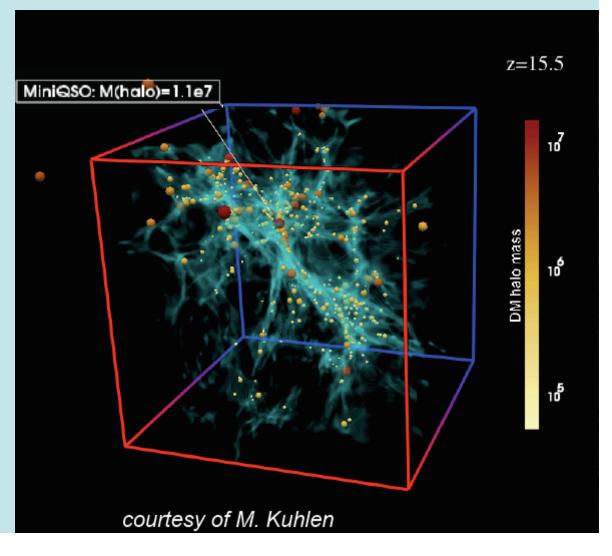
- Viscous transport efficient angular momentum transfer,
- formation of central concentration (Eisenstein & Loeb 1995; Koushiappas+ 2004)
- + proper dynamical treatment of disk stability (Lodato & PN 2006, 2007)

Supermassive star (Haehnelt & Rees 1993)

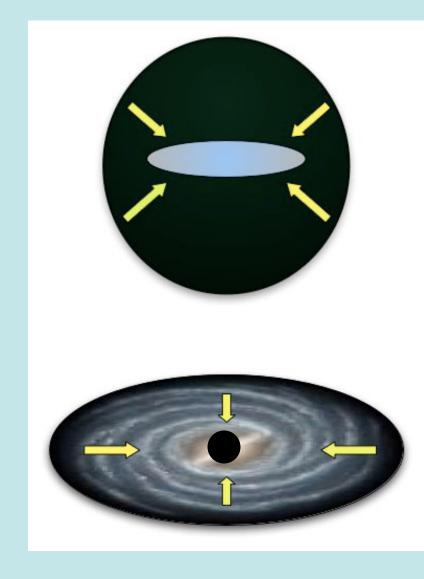
Bar unstable self-gravitating gas + large quasi-star (Begelman, Volonteri & Rees 2006)

## First black holes

- Unclear if PopIII remnants will do as seeds (Alvarez+ 2008!)
- Other ways of making more massive BH seeds need to be considered



## BH seed formation at high z



Baryons inside DM halo collapse and form a rotating pre-galactic disc

Disc becomes gravitationally unstable and accretes to the center

Lodato & PN 2006, 2007

## Sequence of events

DM halo mass M,  $T_{vir}$ , no metals, gas mass  $f_bM$  hot disc ~4000 K, cold disc ~ 400 K

 $T_{vir}$  >  $T_{gas}$  gas collapses and forms rotationally supported disc, disc subject to grav. instabilities, onset when Toomre  $Q_{crit} \sim 1 - 3$ 

Disc evolution tussle between accretion and fragmentation

Bars lead to redistribution of J, feed matter to center continues till central mass stabilizes the disc

Accumulated central mass depends on spin, halo mass,  $T_{gas}/T_{vir}$ , max. spin for which the disc is grav. Unstable, provides upper bound on  $M_{BH}$ 

## **Evolution of self-gravitating discs**

- Controlled mainly by two parameters:
  - Q>1 non-self-gravitating
  - *Q*~1 subject to grav. instabilities

$$t_{\rm cool}\Omega = \frac{t_{\rm cool}}{t_{\rm dyn}}$$

 $Q = \frac{c_{\rm s}\kappa}{\pi G\Sigma}$ 

- $t_{cool}/t_{dyn} < 1$  fragmentation
- $t_{cool}/t_{dyn} > 1$  self-regulation
- Zero-metallicity: long cooling times --> fragmentation can be avoided
- When disc self-regulated, steady ang. mom. transport
- Max. mass accretion rate

$$\dot{M}_{\rm max} = 2\alpha_{
m crit} rac{c_{
m s}^3}{G}$$
  $\dot{M}_{
m max} \approx 10^{-2} M_{\odot}/{
m yr}$  No H<sub>2</sub>  
 $\dot{M}_{
m max} \approx 3 \ 10^{-4} M_{\odot}/{
m yr}$  H<sub>2</sub> present

 $\alpha_{crit} \sim 0.06$  (Rice, Lodato & Armitage 2005)

(Gammie 2001, Lodato & Rice 2004, 2006, Mayer et al., Durisen et al., etc....)

For large M, the internal torques needed to redistribute J too large to be sustained, causes disc to fragment when  $T_{vir} > T_{max}$ 

Happens for critical value of  $\alpha \sim 0.06$  in Keplerian discs

Fragmentation is rapid, timescale local dynamical time stops when enough mass is converted into stars to make disc stable, no J losses, no mass funnelled to center

Key property is T<sub>gas</sub>, atomic or molecular H cooling 2 extreme cases: fragmentation quenches accretion and fragmentation not taken into account in all previous treatments

$$\frac{dn}{dM_{bh}}(M_{bh};z) = \int_{M(T_{\min})}^{M(T_{\max})} \frac{dn}{dM}(M;z)p[\lambda(M_{bh},M)] | \frac{d\lambda}{dM_{BH}} | dM$$

n(M) of DM halos Sheth & Tormen 1999; Gammie+ 2001; Rice & Lodato 2005

## <u>Conditions to assemble a central mass rapidly</u>

- Small baryon angular momentum (spin parameter  $\lambda < 0.05$ ) ---> disc becomes grav. unstable
- Avoid fragmentation:
  - Need long cooling times: zero-metallicity, atomic hydrogen dominated cooling (high redshift)
  - Small infall rates from the halo (relatively low DM halo mass)
- Final state known: the disc will be exactly marginally stable  $(Q=Q_c)$  and excess mass will have been accreted at the center

$$M_{\mathbf{BH}} = m_{\mathrm{d}} M \left[ 1 - \sqrt{\frac{8\lambda}{m_{\mathrm{d}}Q_{\mathrm{c}}} \left(\frac{j_{\mathrm{d}}}{m_{\mathrm{d}}}\right) \left(\frac{T_{\mathrm{gas}}}{T_{\mathrm{vir}}}\right)^{1/2}} \right]$$

## **Fragmentation criteria**

3 interesting regimes:

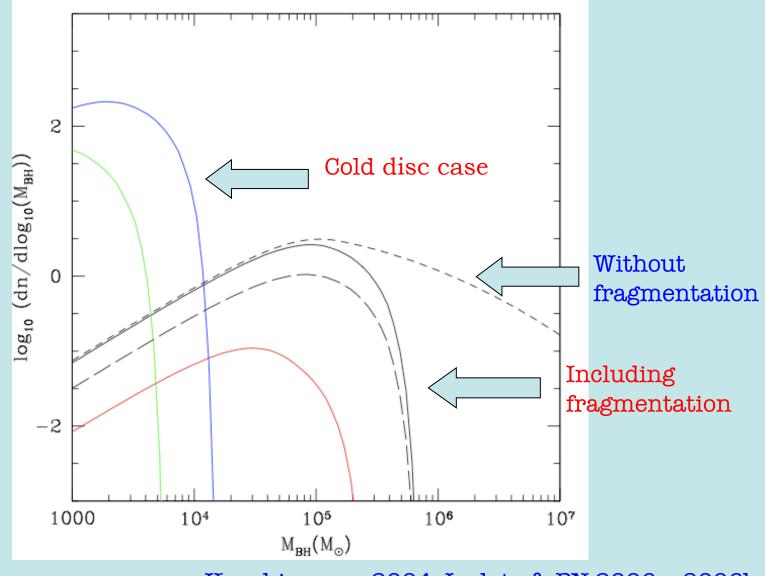
 $T_{vir}/T_{gas} > 3$  will fragment and form stars and no central mass concentrations

 $2 < T_{vir}/T_{gas} < 3$  will fragment and form stars and central mass concentrations

 $T_{vir}/T_{gas} < 2$  will not fragment to form stars, will accrete gas into central mass concentrations that will form BHs

Oh+ 2003; Lodato & PN 2006

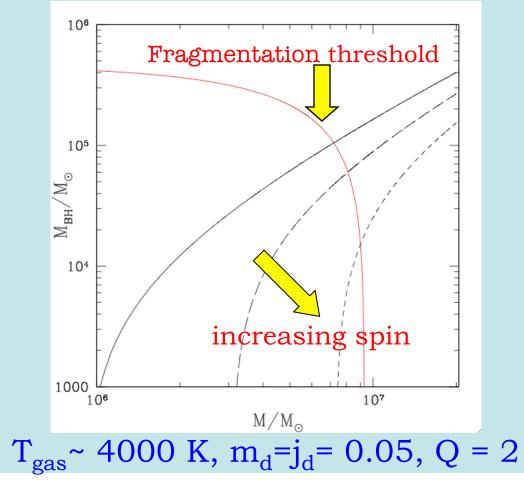
#### Mass function of seed BHs at z = 15



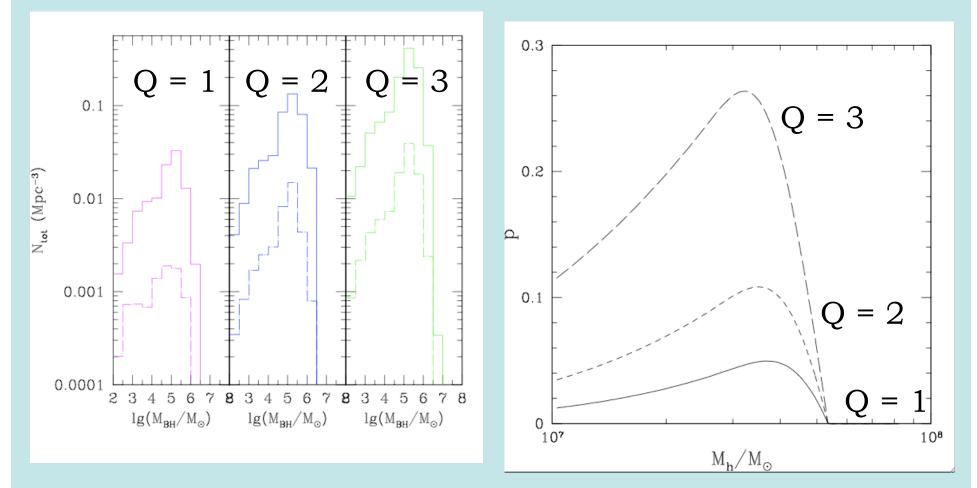
Koushiappas+ 2004; Lodato & PN 2006a, 2006b

## Key features of the model

- Different relationship between  $M_{\rm BH}$  and spin parameter
- High spin halos do not host BHs at early times as disc is not massive enough, stable to grav. instabilities
- Massive, low spin halos host the most massive seeds



### Model features

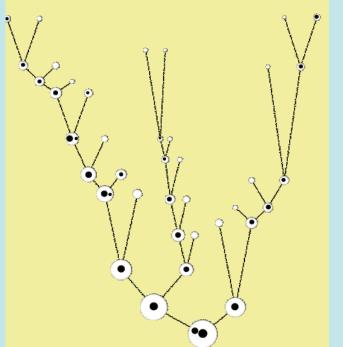


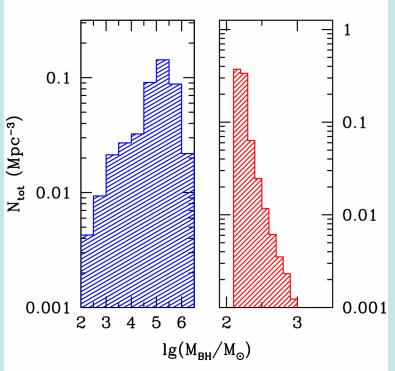
The mass function of seeds for 3 different BH formation efficiency models

Probability of hosting a BH seed of any mass

## <u>Merger induced accretion + CDM merger tree +</u>

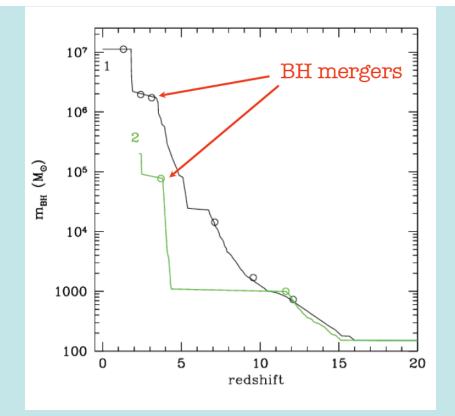






- Plant with initial mass function of seeds
- Generate Monte Carlo merger tree in LCDM
- MBH seed formation ceases at z  $\widetilde{}$  12
- Propogate different BH seed models
- Every major merger (mass ratii 1:10 or greater) induces accretion of gas
- + 2 models: (I) gas mass accreted scales with  $v_{\rm c}{}^5$  and (II) BH mass simply doubles

Hernquist+ 2007; Kauffmann & Haehnelt 2000; Croton+ 2006; Bower+ 2006; Kimm+ 2008; Volonteri, Haardt & Madau 2003; Miloslavjevic & Merritt 2001; Armitage & PN 2004; Hopkins+ 2007; 2008; 2009



- MBH mergers are rare events, as they require the merger of two galaxies each with a central BH
- Not only all MBHs experience a major merger during their lifetime, only  $^{\sim}$  40 50 %
- Dynamical and gravitational interactions can displace MBHs
- Mergers detectable with signatures EM and GW signatures, predict event rates for LISA

Sesana+ 2004; Dotti+ 2006; PN& Volonteri 2008

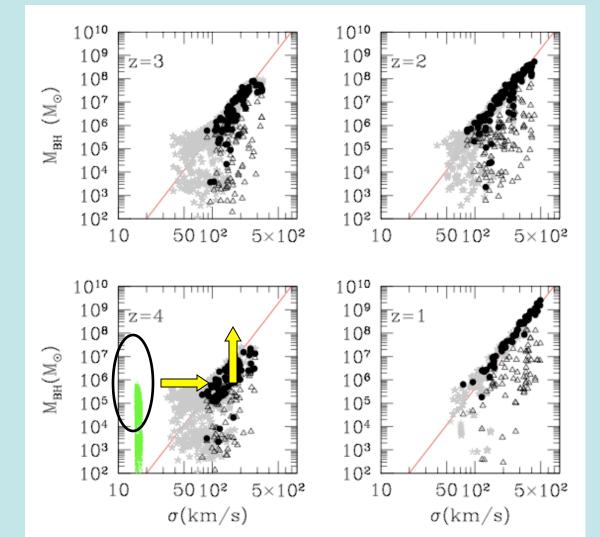
#### Including a model for mergers

- Black holes grow during major mergers
- A fixed fraction of the available cold gas driven to the center and accreted at arbitrarily high rates
- Amount of cold gas is set by cooling and stellar feedback
- Growth during the optically bright quasar and the obscured phase
- Simple light curve with luminosity independent/dependent lifetime

Growth dominated by accretion + at late times by merging of BHs

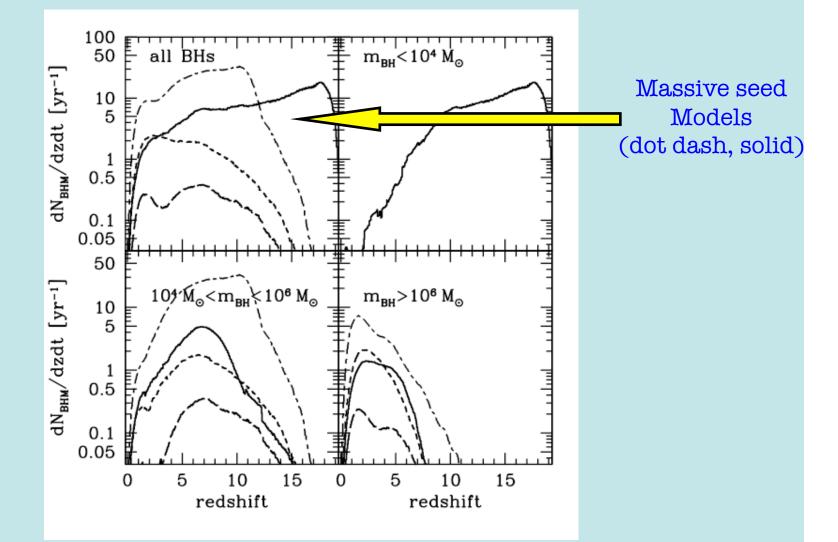
Haehnelt, PN & Rees 1998; Haiman & Loeb 1999; PN 2003; Yu & Tremaine 2002; Kauffmann & Haehnelt 2000,2003; Wyithe & Loeb 2004; Granato+ 2004; Shanker+ 2005; PN 2006; Kauffmann & Haehnelt 2000, 2003

# Merger history of a present day $10^{13}$ M<sub>sun</sub> halo (merging BHs)



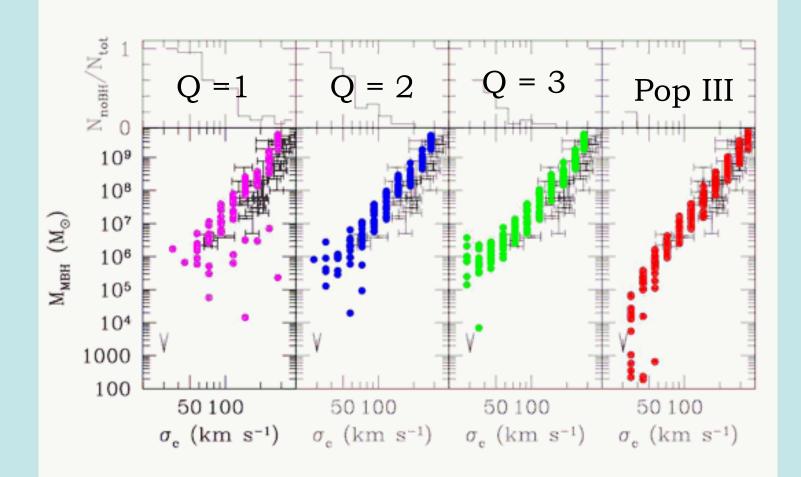
initially over-massive, massive seeds in 3.5-4 sigma peaks

## Number of MBHB mergers/yr



Sesana+ 2007; Volonteri & PN 2010

#### Key prediction at the low mass end



Volonteri, Lodato & PN 2007

## Observations at the low mass end

DRAFT VERSION JULY 25, 2011 Preprint typeset using  ${\rm P}T_{\rm E}{\rm X}$  style emulate apj v. 6/22/04

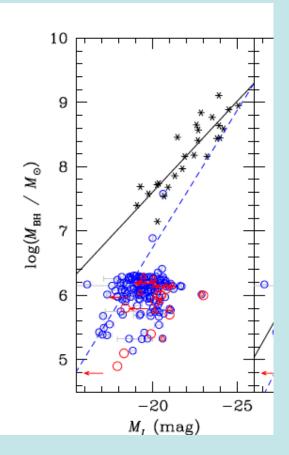
#### BLACK HOLE MASS AND BULGE LUMINOSITY FOR LOW-MASS BLACK HOLES

YAN-FEI JIANG<sup>1</sup>, JENNY E. GREENE<sup>1</sup> AND LUIS C. Ho<sup>2</sup> Draft version July 25, 2011

#### ABSTRACT

We study the scaling between bulge magnitude and central black hole (BH) mass in galaxies with virial BH masses  $\leq 10^6 M_{\odot}$ . Based on careful image decomposition of a snapshot Hubble Space Telescope I-band survey, we found that these BHs are found predominantly in galaxies with pseudobulges. Here we show that the  $M_{\rm BH} - L_{\rm bulge}$  relation for the pseudobulges at low mass is significantly different from classical bulges with BH masses  $\geq 10^7 M_{\odot}$ . Specifically, bulges span a much wider range of bulge luminosity, and on average the luminosity is larger, at fixed  $M_{\rm BH}$ . The trend holds both for the active galaxies from Bentz et al. and the inactive sample of Gültekin et al. and cannot be explained by differences in stellar populations, as it persists when we use dynamical bulge masses. Put another way, the ratio between bulge and BH mass is much larger than ~ 1000 for our sample. This is consistent with recent suggestions that  $M_{\rm BH}$  does not scale with the pseudobulge luminosity. The low-mass scaling relations appear to flatten, consistent with predictions from Volonteri & Natarajan for massive seed BHs.

Subject headings: galaxies: active — galaxies: nuclei — methods: observational



#### Jiang, Greene & Ho 2011

#### Self regulation of SF and BH growth

Quasar driven wind sweeps up gas shell and expels it, inhibiting SF and limiting BH mass that depends on spin

$$M_{bh} > \alpha \frac{\sigma^{5} \kappa}{G^{2} c} = 8 \times 10^{8} \gamma (\frac{\sigma}{500 km/s})^{5} M_{\Theta}$$

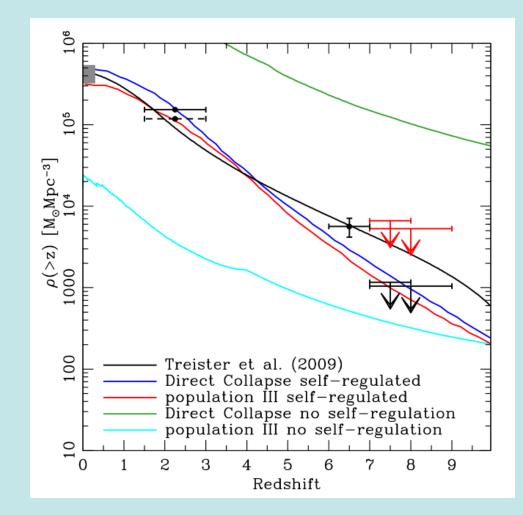
$$M_{bh} \approx 10^{8} (\frac{f_{kin}}{0.001})^{-1} j_{d}^{-5} (\frac{\lambda}{0.05})^{-5} (\frac{m_{d}}{0.1})^{5} (\frac{v_{halo}}{400 km/s})^{5} M_{\Theta}$$

$$Momentum-driven winds$$

$$L_{crit} = \frac{4 f_{g} c}{G} \sigma^{4} \qquad M_{\bullet,crit} / M_{sun} = 0.12 \eta_{Edd}^{-1} \left(\frac{f_{g}}{0.1}\right) \left(\frac{\sigma}{km/s}\right)^{4}$$

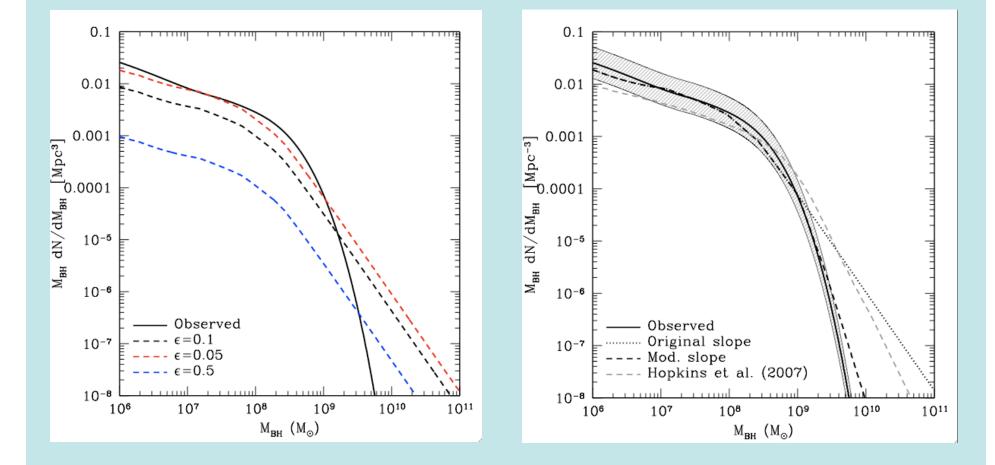
Haehnelt, PN & Rees 1998; Silk & Rees 1998; Wyithe & Loeb 2003; King 2004; Thompson, Murray & Quataert 2004; PN 2008

## **Evidence for self-regulation?**



Treister+ Nature, 2011

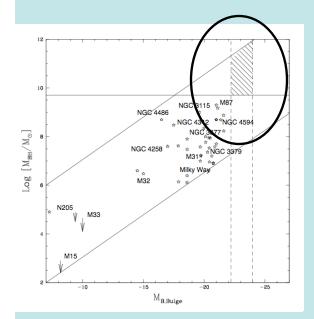
## <u>An upper limit to BH masses?</u>



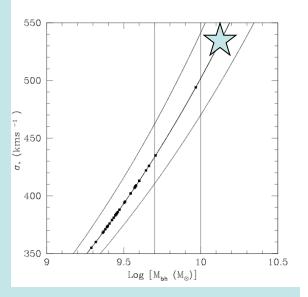
 $M_{upper-limit}$  (BH) ~ few times X 10<sup>10</sup> Msun for cD galaxies

PN & Treister 2008

## Predict existence of UMBHs





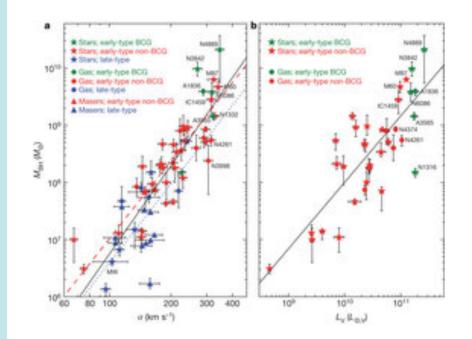


Expected in the centers of cDs

SDSS results of nearby cDs and bright ellipticals

Lauer+ 2005; 2006; Bernardi+ 2006

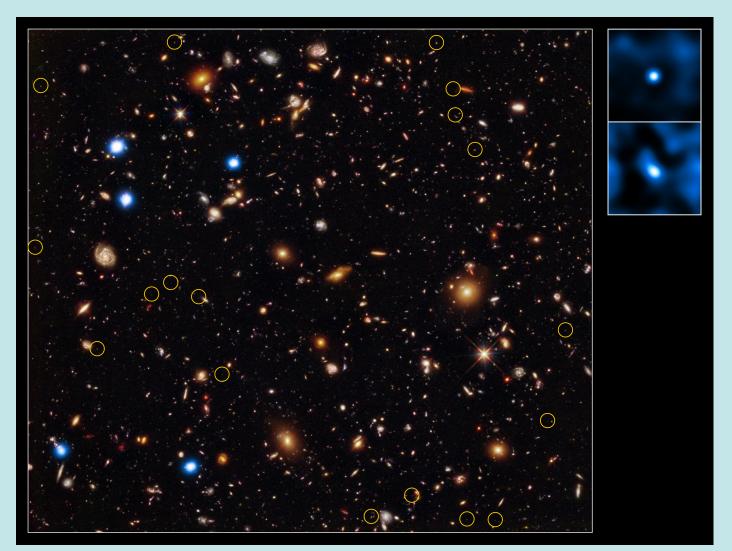
## <u>Ultra-massive BHs?</u>



**a**, Black-hole mass ( $M_{BH}$ ) versus stellar velocity dispersion ( $\sigma$ ) for 65 galaxies with direct dynamical measurements of  $M_{BH}$ . For galaxies with spatially resolved stellar

#### McConnell+ 2011

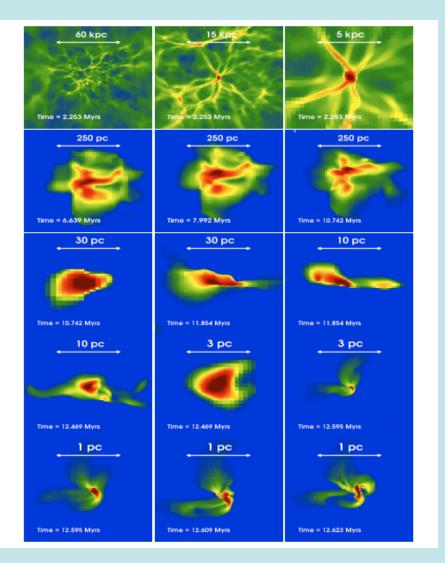
## Discovery of a new population of baby BHs at z = 6 - 8?



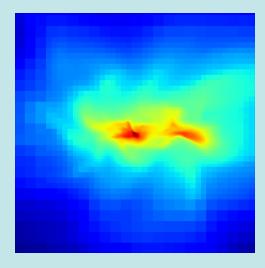
Treister, PN + *Nature*, June 2011; see also Cowie, Barger & Hasinger preprint; Fiore et al and Willott

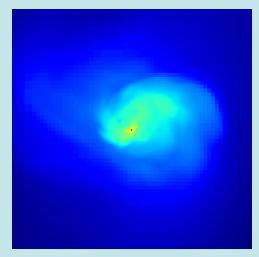


## **Massive BH seed formation**



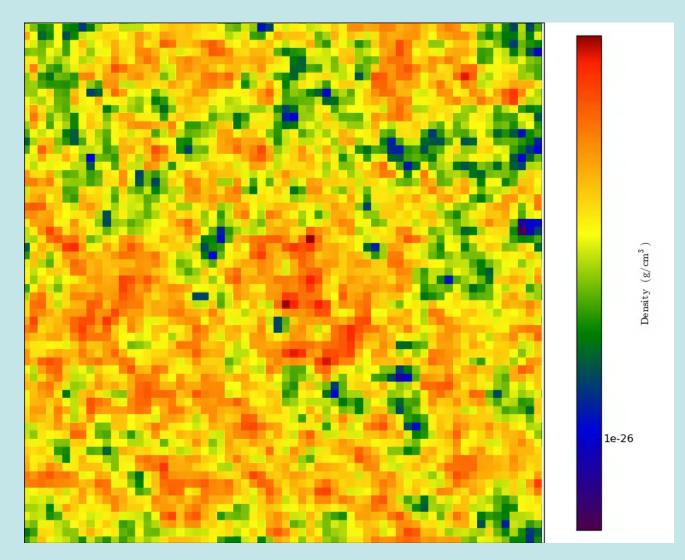




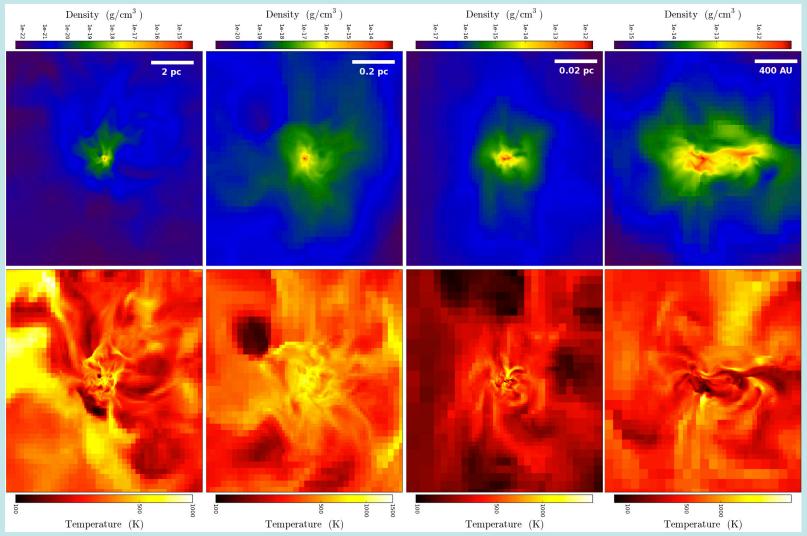


#### Smith, Davis & PN 2011

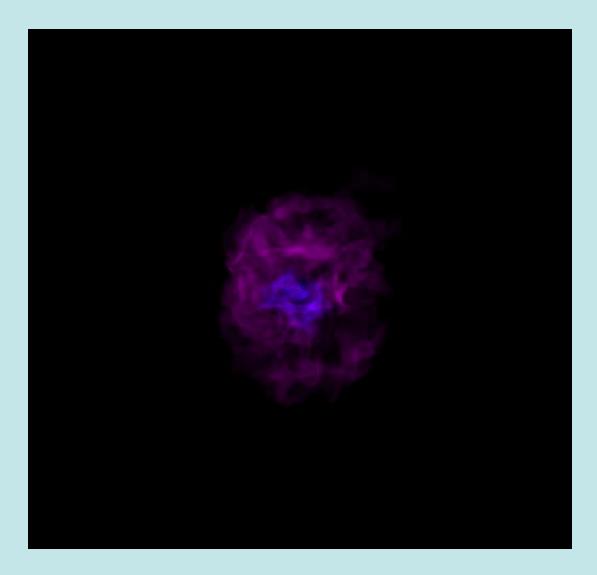
## Baryonic collapse in a high spin halo



## The fate of rotating fat disks



#### Davis, Smith, Turk, PN, in prep



## Work in progress....

