## Emergent behavior in self-organized dynamics: from consensus to hydrodynamic flocking

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• "Living agents": flocks of birds; schools of fish; colonies of ants, bacteria, cells









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• "Non-living agents": sensor-based robots, UAVs, micro-motors, nematic fluids, ...









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   attraction-repulsion, phase transition, ...
- Key issue: the notion of 'local neighborhood'

#### Outline

- Rules of engagement: alignment
  - Krause model for opinion dynamics
  - Sensor-based motion the rendezvous problem
  - Vicsek model for flocking; phase transition
  - Cucker-Smale models for flocking near and far from equilibrium
- $t \to \infty$ : The emergence of consensus, parties, leaders, ...
  - Large time behavior consensus, flocking, ...
  - Synchronization Kuramoto model
  - Taking tendency into account emergence of leaders
  - A general perspective
- 3  $N \to \infty$ : Social hydrodynamics
  - Kinetic description
  - From kinetic to hydrodynamic description of flocking
  - Hydrodynamic alignment smooth solutions must flock
  - Critical thresholds in flocking hydrodynamics

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- ullet Neighborhood  $\mathcal{N}_i$  geometric neighborhood dictated by  $\phi$

## Example#1: Krause model<sup>2</sup> for opinion dynamics

• State space — vectors of "opinions":  $\{\mathbf{x}_i(t)\}_{i=1}^N$ 

<sup>&</sup>lt;sup>2</sup>U. Krause (1997,2000), R. Hegselmann & U.Krause: Opinion dynamics and bounded confidence models, analysis and simulation (2002,2004)

# Example#1: Krause model<sup>2</sup> for opinion dynamics

- State space vectors of "opinions":  $\{\mathbf{x}_i(t)\}_{i=1}^N$
- ullet Krause-Hegselmann model (1997) interaction through **local** averaging:

$$\mathbf{x}_i(t+\Delta t) = \frac{1}{N_i} \sum_{|\mathbf{x}_i-\mathbf{x}_j| \leq R} \mathbf{x}_j(t), \qquad N_i := \#\{\mathbf{x}_j : |\mathbf{x}_j-\mathbf{x}_i| < R\}$$

"Environmental averaging":

$$\mathbf{x}_i(t+\Delta t) = \sum_{j\in\mathcal{N}_i} a_{ij}\mathbf{x}_j(t)$$
  $\sum_j a_{ij} = 1$ 

Act on difference of opinions:

$$a_{ij} = \frac{\phi(|\mathbf{x}_i - \mathbf{x}_j|)}{\deg_i} \qquad \qquad \phi(r) = \mathbb{1}_{[0,R)}(r)$$

•  $\deg_i = \sum_{k \in \mathcal{N}_i} \phi(|\mathbf{x}_i - \mathbf{x}_k|) \rightsquigarrow N_i$ 

— the degree of influence on  $agent_i$ 

\* Typical question in collective dynamics: does "averaging" lead to "consensus"?

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- State space vectors of positions:  $\{\mathbf{p}_i(t)\}_{i=1}^N \leadsto \{\mathbf{x}_i(t)\}_{i=1}^N$
- Mean-shift gradient flow<sup>3</sup> (non parametric clustering):

$$\mathbf{x} (t + \Delta t) = \frac{\sum_{j} \phi(|\mathbf{x} (t) - \mathbf{x}_{j}|)\mathbf{x}_{j}}{\sum_{j} \phi(|\mathbf{x} (t) - \mathbf{x}_{j}|)} \leftarrow \text{local mean}$$

$$= \mathbf{x} + \frac{\Delta t}{\deg(\mathbf{x})} \sum_{j \in \mathcal{N}_{i}} \phi(|\mathbf{x} - \mathbf{x}_{j}|)(\mathbf{x}_{j} - \mathbf{x})_{|\mathbf{x} = \mathbf{x}_{i}(t)} \leftarrow \text{alignment}$$

- Local means are shifted in direction of maximal increase of density
- Robotic agents: the "rendezvous problem"  $^{3b-3d}$   $\mathbf{x}_i(t) \mathbf{x}_j(t) \xrightarrow{t \to \infty} 0$ ?

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<sup>&</sup>lt;sup>3b</sup>On limited visibility — Ji & Egerstedt (2007); Bellaiche & Bruckstein (2015);

<sup>&</sup>lt;sup>3c</sup>Time delay — Olfati-Saber & Murray (2004), Somarakis & Baras (2013-);

 $<sup>^{3</sup>d}$ Connectivity of mobile networks — Zalvanos, Pappas, et. al (2007-2011)

## Example #3: Vicsek model<sup>4</sup> — alignment of orientations

• Fix a speed s. Averaging of orientations  $\{\mathbf{p}_i(t)\}_{i=1}^N \leadsto \{\omega_i(t)\}_{i=1}^N \in \mathbb{S}^{\mathsf{d}-1}$ 

$$\omega_i(t+\Delta t) = \left(s\sum_{j\in\mathcal{N}_i} a_{ij}\,\omega_j(t) + \mathsf{noise}
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ullet 2D additive Noise= uniform in angle in [- au, au]

$$\bullet \ a_{ij} = \frac{\phi(|\mathbf{x}_i - \mathbf{x}_j|)}{\deg_i} \quad \leadsto \quad \mathbf{v}_i(t) := \frac{s}{\deg_i} \sum_{j: |\mathbf{x}_i - \mathbf{x}_j| < R} \phi(|\mathbf{x}_i - \mathbf{x}_j|) \omega_j(t)$$

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Vicsek model as an alignment dynamics

$$\frac{\mathrm{d}}{\mathrm{d}t}\mathbf{v}_{i}(t) = \alpha \sum_{i \in \mathcal{N}_{i}} a_{ij} \left(\mathbf{v}_{j} - \frac{\langle \mathbf{v}_{i}, \mathbf{v}_{j} \rangle}{|\mathbf{v}_{i}|^{2}} \mathbf{v}_{i}\right)$$

• Beyond environmental averaging:

Projection  $\frac{\langle \mathbf{v}_i, \mathbf{v}_j \rangle}{|\mathbf{v}_i|^2}$  enforces fixed speed,  $|\mathbf{v}_i(t)| = s$ , and noise  $(\tau)$ 

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• State space — vectors of velocities:  $\{\mathbf{p}_i(t)\}_{i=1}^N \leadsto \{\mathbf{v}_i(t)\}_{i=1}^N$ 

<sup>&</sup>lt;sup>5</sup>Emergent Behavior in Flocks (2007)

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Emergent behavior in self-organized dynamics

 $<sup>^5</sup>$ Emergent Behavior in Flocks (2007)  $^{5a}a_{ii}:=1-\sum_{j
eq i}a_{ij}$ 

# Example#4: Cucker-Smale model<sup>5</sup>— velocity alignment

- State space vectors of velocities:  $\{\mathbf{p}_i(t)\}_{i=1}^N \leadsto \{\mathbf{v}_i(t)\}_{i=1}^N$
- Environmental averaging of velocities  $\{\mathbf{v}_i(t)\}_{i=1}^N \in \mathbb{R}^d \leadsto \text{alignment}^{6a}$

$$\frac{\mathsf{d}}{\mathsf{d}t}\mathbf{v}_i(t) = \alpha \Big(\sum_{j \in \mathcal{N}_i} a_{ij}\mathbf{v}_j(t) - \mathbf{v}_i(t)\Big) = \alpha \sum_{j \in \mathcal{N}_i} a_{ij} \Big(\mathbf{v}_j(t) - \mathbf{v}_i(t)\Big)$$

• A second-order model:

$$a_{ij}(\mathbf{x}(t)) = rac{1}{\deg_i}\phi(|\mathbf{x}_i(t) - \mathbf{x}_j(t)|), \quad rac{\mathsf{d}}{\mathsf{d}t}\mathbf{x}_i(t) = \mathbf{v}_i(t)$$

ullet Global vs. local models: scaling by graph degree  $\deg_i = \sum_k \phi(|\mathbf{x}_i - \mathbf{x}_k|)$ 

 $<sup>^5</sup>$ Emergent Behavior in Flocks (2007)  $^{5a}a_{ii}:=1-\sum_{j
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- $\star$  Local models when diameter (Supp $\{\phi\}$ )  $\ll$  max  $|\mathbf{x}_i(t) \mathbf{x}_j(t)|$ :

Short-range interactions involve 'nearby' neighbors  $\deg_i \rightsquigarrow N_i$ 

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### Example#5: far from equilibrium

• (with S. Motsch<sup>6</sup>):  $a_{ij}(\mathbf{x}(t)) = \frac{1}{\deg_i} \phi(|\mathbf{x}_i(t) - \mathbf{x}_j(t)|)$ 

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<sup>&</sup>lt;sup>6</sup>A new model for self-organized dynamics and its flocking behavior (2011)

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$$\star \deg_i (i \in G_1) \approx N_1 \phi_0 \quad \leadsto \quad \frac{\mathsf{d}}{\mathsf{d} t} \mathbf{v}_i(t) = \frac{\alpha}{N_1 \phi_0} \sum_{G_1} \phi_{ij} \left( \mathbf{v}_j(t) - \mathbf{v}_i(t) \right)$$

10

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•  $A_{\phi} = \{a_{ij} = \frac{\phi_{ij}}{\deg_i}\}$  is <u>not</u> symmetric . . .

but involves the symmetric graph Laplacian  $L_A:=I-D^{1/2}A_\phi D^{-1/2}$ 

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#### Outline

- Rules of engagement: alignment
  - Krause model for opinion dynamics
  - Sensor-based motion the rendezvous problem
  - Vicsek model for flocking; phase transition
  - Cucker-Smale models for flocking near and far from equilibrium
- $oldsymbol{2} t 
  ightarrow \infty$ : The emergence of consensus, parties, leaders, ...
  - Large time behavior consensus, flocking, . . .
  - Synchronization Kuramoto model
  - Taking tendency into account emergence of leaders
  - A general perspective
- - Kinetic description
  - From kinetic to hydrodynamic description of flocking
  - Hydrodynamic alignment smooth solutions must flock
  - Critical thresholds in flocking hydrodynamics

Graph 
$$G = (V, E)$$
: vertices  $V = \{\mathbf{p}_i\} \subset \mathbb{R}^N$ ; edges  $E_{\phi} = \{\mathbf{e}_{ij}\} \subset \mathbb{R}^N \times \mathbb{R}^N$  grad  $\nabla_{\phi}(\mathbf{p})_{ij} := \sqrt{\phi_{ij}}(\mathbf{p}_i - \mathbf{p}_j)$ ; divergence  $\operatorname{div}_{\phi}(\mathbf{e})_i := \sum_i \sqrt{\phi_{ij}}(\mathbf{e}_{ij} - \mathbf{e}_{ji})$ 

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- Consensus depends on Fiedler  $\#\mu(t) = \lambda_2(L_{A(\mathbf{x}(t))}) > 0$   $(L_A := I A_\phi)$
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- \* Interplay between dynamics on graph and graph driven by the dynamics

## Cucker-Smale vs. Motsch-Tadmor model (global influence)

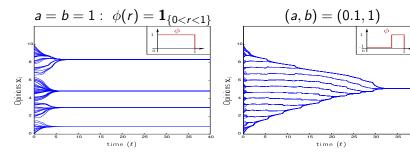
• Global influence:  $\phi_{ij} = \phi(|\mathbf{x}_i - \mathbf{x}_j|)$ 

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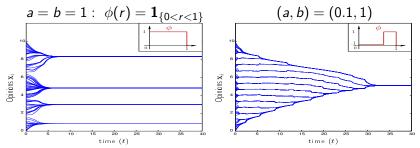
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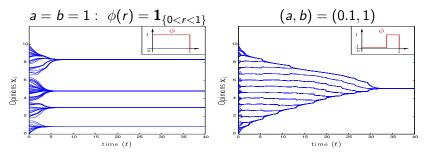
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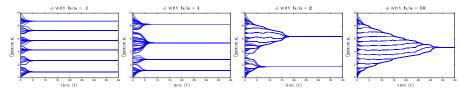
ullet 100 uniformly distributed opinions:  $\phi(r) = a {f 1}_{\{r \leqslant rac{1}{\sqrt{2}}\}} + b {f 1}_{\{rac{1}{\sqrt{2}} \leqslant r < 1\}}$ 



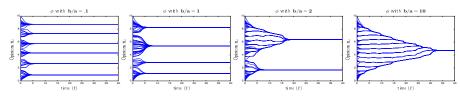
• Homophilious dynamics: align with those that think alike  $(a \gg b)$ 



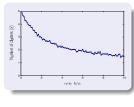
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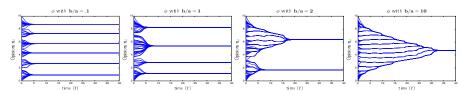
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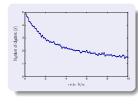
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- $\star$  Heterophilious dynamics enhances consensus  $\phi(\cdot)_{|0 < r < R}$  is increasing  $\leadsto \lambda_2(L_{A(\mathbf{x}(t))}) > \eta > 0$ ?



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### 2D heterophilious dynamics

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"Regarding networks and homophily, birds of a feather flock together, but people are influenced by those they like. Both these processes result in the same outcome (similar people together in groups), but there is no standard accepted way of separating these two processes".

Introduction to Mathematical Sociology, P. Bonacich & P. Lu

<sup>&</sup>lt;sup>7</sup>Garnier, Papanicolaou, Yang, Consensus convergence with stochastic effects (2017)

Emergent behavior in self-organized dynamics

16

• Empirical distribution  $\rho^{N}(t, \mathbf{x}) := \frac{1}{N} \sum_{i} \delta_{\mathbf{x} - \mathbf{x}_{i}(t)} \cong \overline{\rho} + \frac{\rho_{1}}{\rho_{1}}(t, \mathbf{x})$ 

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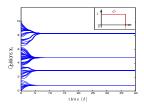
<sup>&</sup>lt;sup>7</sup>Garnier, Papanicolaou, Yang, Consensus convergence with stochastic effects (2017)

Emergent behavior in self-organized dynamics

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- Empirical distribution  $ho^N(t,\mathbf{x}) := \frac{1}{N} \sum_i \delta_{\mathbf{x}-\mathbf{x}_i(t)} \cong \overline{\rho} + \frac{\rho_1(t,\mathbf{x})}{\rho_1(t,\mathbf{x})}$ Fluctuations<sup>7</sup>:  $|\widehat{\rho_1}(t,k)| \sim e^{\gamma_{\max}t}$ ,  $\gamma_{\max} := \max_{k>0} \left[ 2k \int_0^1 \phi(s) \sin(ks) s \, \mathrm{d}s \right]$
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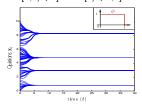
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# Fluctuations and mean distance of interacting clusters

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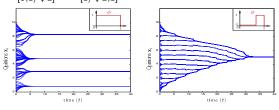


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mean distance  $\simeq 0.7$ 

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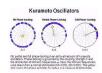
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### Example #6: Synchronization

• Kuramoto model<sup>8</sup>  $\{\mathbf{p}_i\} \leadsto \text{phases } \{\theta_i\}$  or frequencies  $\{\omega_i = \dot{\theta}_i\}$ 

$$\frac{\mathsf{d}}{\mathsf{d}t}\theta_i(t) = \Omega_i + \frac{K}{N} \sum_j \sin(\theta_j - \theta_i) \rightsquigarrow \Omega_i + \frac{K}{\mathsf{deg}_i} \sum_j a_{ij}(\theta_j - \theta_i), \ a_{ij} = \frac{\sin(\theta_j - \theta_i)}{\theta_j - \theta_i}$$

- $|\Omega_i| < \alpha r$ : steady states; more oscillators recruited into synchronized clusters  $(r \approx 1)$  as  $\alpha > \alpha_c$  increases
- $|\Omega_i| > \alpha r$  : no synchronization  $(r \approx 0)$  is possible for  $\alpha < \alpha_c$



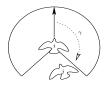
Stewart Heitmann

<sup>&</sup>lt;sup>8</sup>Kuramoto, Lecture Notes Phys. (1975, 1984)...Acebron et. al. RevModPhys (2005) <sup>8b</sup>Ha et. al (2010 –), Gerard-Varet, Dietert, Fernandez, Giacomin (2016)

### Example #7: tendency to follow in "thinking agents"



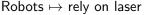
Robots  $\mapsto$  rely on laser

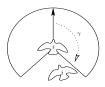


"Living agents"  $\mapsto$  rely on vision

### Example #7: tendency to follow in "thinking agents"







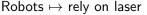
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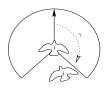
• Influenced by the **projection** of those moving ahead tracing using pheromones rather than vision . . . following the footsteps of – education . . .



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$$\frac{\mathrm{d}}{\mathrm{d}t}\mathbf{p}_{i}(t) = \frac{1}{\mathrm{deg}_{i}}\sum_{i\in\mathcal{N}_{i}}\phi(|\mathbf{x}_{i}-\mathbf{x}_{j}|)\left(\frac{\langle\mathbf{p}_{i},\mathbf{p}_{j}\rangle}{|\mathbf{p}_{j}|^{2}}\mathbf{p}_{j}-\mathbf{p}_{i}\right),$$

### Tendency and emergence of leaders - 1st-order model

$$\frac{d}{dt}\mathbf{x}_{i}(t) = \sum_{j \in \mathcal{N}_{i}} a_{ij}\mathbf{x}_{j} - \mathbf{x}_{i}, \quad a_{ij} = \frac{1}{\deg_{i}}\phi(|\mathbf{x} - \mathbf{x}_{j}|) \frac{\langle \mathbf{x}_{i}, \mathbf{x}_{j} \rangle}{|\mathbf{x}_{j}|^{2}}$$

Figure: Random initial conditions,  $\phi = \mathbb{1}_{[0,1]}$ . Snapshots: t = 0, 0.3, 0.5, 2, 5, 70.

#### Self-organized dynamics

- Biology The role of empirical data
  - Flocks, swarms, colonies, ... how are they formed?
  - Since there is no Newton's law what are the rules of engagement?
  - ★ Agents <u>are</u> different are these observed patterns observed patterns system specific?

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- Mathematics Agent-based models; non-local PDEs
   Agent-based → kinetic models → macroscopic models
   ★ Numerical and analytical studies of 'social hydrodynamics'

#### Outline

- Rules of engagement: alignment
  - Krause model for opinion dynamics
  - Sensor-based motion the rendezvous problem
  - Vicsek model for flocking; phase transition
  - Cucker-Smale models for flocking near and far from equilibrium
- 2  $t \to \infty$ : The emergence of consensus, parties, leaders, ...
  - Large time behavior consensus, flocking, ...
  - Synchronization Kuramoto model
  - Taking tendency into account emergence of leaders
  - A general perspective
- - Kinetic description
  - From kinetic to hydrodynamic description of flocking
  - Hydrodynamic alignment smooth solutions must flock
  - Critical thresholds in flocking hydrodynamics

# Second limit — behavior of large crowds $^9$ $N ightarrow \infty$

• Empirical distribution  $f^N := \frac{1}{N} \sum_j \delta_{\mathbf{x} - \mathbf{x}_j(t)} \otimes \delta_{\mathbf{v} - \mathbf{v}_j(t)} \longrightarrow f(t, \mathbf{x}, \mathbf{v})$ 

Vlasov eq. for distribution  $f(t, \mathbf{x}, \mathbf{v})$ 

$$f_t + \mathbf{v} \cdot \nabla_{\mathbf{x}} f + \alpha \nabla_{\mathbf{v}} \cdot Q(f, f) = 0$$

<sup>&</sup>lt;sup>9</sup>Ha & ET (2008); Canizo, Carrillo, Rosado (2009); Carrillo-Fornasier-Rosado-Toscani (2009), Motsch-ET (2014)

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$$Q(f,f) = \frac{1}{\deg(t,\mathbf{x})} \int_{\mathbb{R}^{2d}} \frac{\phi(|\mathbf{x} - \mathbf{y}|)(\mathbf{w} - \mathbf{v})f(t,\mathbf{x},\mathbf{v})f(t,\mathbf{y},\mathbf{w})d\mathbf{y}d\mathbf{w} + \dots$$

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Flocking  $(K = 1)$   $f \rightsquigarrow \rho(t,\mathbf{x}) \delta(\mathbf{v} - \mathbf{u}(t,\mathbf{x})) \dots$ 

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• Flocking (K = 1)  $f \rightsquigarrow \rho(t, \mathbf{x})\delta(\mathbf{v} - \mathbf{u}(t, \mathbf{x}))...$ 

.... recovered in terms of moments —

$$\begin{bmatrix} \text{density ......} & \rho \\ \text{momentum ...} & \rho \mathbf{u} \end{bmatrix} = \int \begin{bmatrix} 1 \\ \mathbf{v} \end{bmatrix} f(t, \mathbf{x}, \mathbf{v}) d\mathbf{v}$$

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$$\begin{cases} \text{ mass : } & \partial_t \rho + \nabla_{\mathbf{x}} \cdot (\rho \mathbf{u}) = 0 \\ \text{ momentum : } & \partial_t (\rho \mathbf{u}) + \nabla_{\mathbf{x}} \cdot (\rho \mathbf{u} \otimes \mathbf{u} + P(\mathbf{f})) = \rho \, \mathcal{A}_{\rho}(\mathbf{u}) \end{cases}$$

<sup>&</sup>lt;sup>10</sup>R. Shvydkoy & ET, Eulerian dynamics with commutator forcing (2016)

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$$\begin{split} \bullet \text{ Commutator form}^{10} \colon & \mathcal{A}_{\rho}(\mathbf{u}) := \frac{\alpha}{\deg} [L_{\phi}, \mathbf{u}](\rho), \\ & [L_{\phi}, \mathbf{u}](\rho) = L_{\phi}(\rho \mathbf{u}) - L_{\phi}(\rho) \mathbf{u}, \qquad L_{\phi}(g) := \int_{\mathbb{R}^d} \phi(\mathbf{x}, \mathbf{y}) (g(\mathbf{y}) - g(\mathbf{x})) \mathrm{d}\mathbf{y} \end{split}$$

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$$\text{Stress tensor}^{10b} \qquad P_{ij}(\mathbf{f}) = \int_{\mathbb{R}^d} (\mathbf{v}_i - \mathbf{u}_i)(\mathbf{v}_j - \mathbf{u}_j)\mathbf{f}(\mathbf{t}, \mathbf{x}, \mathbf{v})d\mathbf{v} \equiv 0$$

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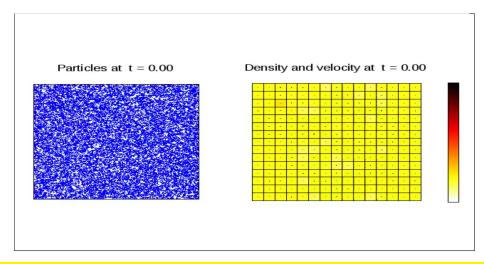
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Transport+Alignment: 
$$\mathbf{u}_t + (\mathbf{u} \cdot \nabla_{\mathbf{x}})\mathbf{u} = \frac{\mathcal{A}_{\rho}(\mathbf{u})}{\mathcal{A}_{\rho}(\mathbf{u})} = \frac{\alpha}{\deg}[L_{\phi}, \mathbf{u}](\rho)$$

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### Hydrodynamic vs. agent-base description

Vicsek model: agent-base model vs. hydrodynamic description



## Flocking behavior – bounded kernels

• Classical solutions must flock





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# Flocking behavior – bounded kernels





• Classical solutions must flock

$$\begin{array}{lcl} \rho_t + \nabla_{\mathbf{x}} \cdot (\rho \mathbf{u}) & = & 0, & \underline{\text{compactly supported } \rho_0} \\ \mathbf{u}_t + (\mathbf{u} \cdot \nabla_{\mathbf{x}}) \mathbf{u} & = & \mathcal{A}_{\rho}(\mathbf{u}), & (\mathcal{A}_{\rho}(\mathbf{u}) = \frac{\alpha}{\deg} [L_{\phi}, \mathbf{u}](\rho)) \end{array}$$

# Flocking behavior - bounded kernels





• Classical solutions must flock

$$\rho_t + \nabla_{\mathbf{x}} \cdot (\rho \mathbf{u}) = 0, \qquad \underbrace{\text{compactly supported}}_{t} \rho_0$$

$$\mathbf{u}_t + (\mathbf{u} \cdot \nabla_{\mathbf{x}}) \mathbf{u} = \underbrace{\mathcal{A}_{\rho}}_{t}(\mathbf{u}), \qquad (\underbrace{\mathcal{A}_{\rho}}_{t}(\mathbf{u}) = \frac{\alpha}{\deg}[L_{\phi}, \mathbf{u}](\rho))$$

$$\underbrace{\text{Theorem}^{11}}_{t}. \text{ Set diameter } [\mathbf{u}(t)]_{\infty} := \sup_{\mathbf{x}, \mathbf{y} \in \text{Supp } \rho(t, \cdot)} |\mathbf{u}(t, \mathbf{x}) - \mathbf{u}(t, \mathbf{y})|$$

$$\underbrace{\text{If } \mathbf{u} \in C^1 \ \, \rightsquigarrow \ \, \frac{d}{dt}[\mathbf{u}(t)] \leqslant -\alpha \, \mu(t) \, [\mathbf{u}(t)] :$$

$$\mu_{\infty}(t) \text{ is coefficient of ergodicity } \geqslant \min_{\mathbf{x}, \mathbf{y} \in \text{Supp } \rho(t, \cdot)} \phi(|\mathbf{x} - \mathbf{y}|)$$

$$\underbrace{\text{Then } \underline{\mathbf{global}}}_{t} \text{ interaction } \int_{0}^{\infty} \phi(s) ds = \infty \ \, \rightsquigarrow \ \, \text{unconditional flocking}^{11b} :$$

$$\underbrace{\mathbf{u}(t, \cdot) \xrightarrow{t \to \infty}_{t} \overline{\mathbf{u}}}_{t}$$

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$$\begin{array}{lcl} \rho_t + \nabla_{\mathbf{x}} \cdot (\rho \mathbf{u}) & = & 0, & \underline{\text{compactly supported } \rho_0} \\ \mathbf{u}_t + (\mathbf{u} \cdot \nabla_{\mathbf{x}}) \mathbf{u} & = & \mathcal{A}_{\rho}(\mathbf{u}), & (\mathcal{A}_{\rho}(\mathbf{u}) = \frac{\alpha}{\deg} [L_{\phi}, \mathbf{u}](\rho)) \end{array}$$

$$\underline{\text{If }} \mathbf{u} \in C^1 \implies \frac{\mathsf{d}}{\mathsf{d}t} [\mathbf{u}(t)] \leqslant -\alpha \, \mu(t) [\mathbf{u}(t)]:$$

$$\mu_{\infty}(t)$$
 is coefficient of ergodicity  $\geqslant \min_{\mathbf{r}, \mathbf{x}, \mathbf{y} \in \mathsf{Supp}\, \rho(t, \cdot)} \phi(|\mathbf{x} - \mathbf{y}|)$ 

Then global interaction  $\int_{-\infty}^{\infty} \phi(s) ds = \infty \implies \text{unconditional flocking}^{11b}$ :

$$\mbox{unconditional flocking} \dots \left\{ \begin{aligned} \mathbf{u}(t,\cdot) &\xrightarrow{t \to \infty} \overline{\mathbf{u}} \\ \rho(t,\mathbf{x}) - \rho_{\infty}(\mathbf{x} - t\overline{\mathbf{u}}) &\xrightarrow{t \to \infty} 0 \end{aligned} \right.$$

• When does  $\mathbf{u}(t,\cdot) \in C^1$ -solution exist? What about <u>local</u> interaction?

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# We have a much larger class in mind...

$$\mathbf{u}_t + \mathbf{u} \cdot 
abla \mathbf{u} = [L_{\phi}, \mathbf{u}](
ho) = \int_{\mathbb{R}^d} \phi(|\mathbf{x} - \mathbf{y}|) (\mathbf{u}(\mathbf{y}) - \mathbf{u}(\mathbf{x})) 
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- Singular  $\phi$ 's Fractional dissipation:  $L = -(-\Delta)^{\beta/2} \iff \phi_{\beta}(\mathbf{x}) = |\mathbf{x}|^{-(d+\beta)}$

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• "Local action" — Limiting case  $\beta=2$  — Navier-Stokes eqs.  $L=\Delta$ 

$$\rightsquigarrow$$
  $(\rho \mathbf{u})_t + \nabla(\rho \mathbf{u} \otimes \mathbf{u}) = \nabla(\rho^2 D \mathbf{u}), \quad D \mathbf{u} = \{\partial_i u_j\}$ 

 $\bullet$  For bounded  $\phi$ 's — under certain critical threshold conditions in configuration space

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- For bounded  $\phi$ 's under certain critical threshold conditions in configuration space
- For NS eqs<sup>12</sup> ( $\beta = 2$ ):  $\mathbf{u}_t + \mathbf{u} \cdot \nabla \mathbf{u} = \text{div}(\rho^2 \nabla \mathbf{u})$

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- $\star \beta = 1$  is critical ... in fact, the whole range  $0 < \beta < 1$  is critical!  $^{12c}$
- $\star$  In contrast to blow-up in fractional Burgers: regularity iff  $^{12d}$   $eta\geqslant 1$

$$u_t + uu_x = \int_{\mathbb{R}} \frac{u(y) - u(x)}{|x - y|^{1+\beta}} dy, \qquad \beta < 2$$

• The role of (i) no vacuum and (ii) the spectral gap in 2D dynamics

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<sup>&</sup>lt;sup>12d</sup>Caffarelli-Vasseur, Kiselev-Nazarov, Constantin-Vicol

#### Progress — mostly one- and two-dimensional models

ullet Critical threshold – 1D flocking hydrodynamics  $^{13}$ 





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• Flocking – fractional dissipation  $^{13b}$ :  $\phi_{\beta}(x) = |x|^{-(1+\beta)}$ 



ullet Spectral gap in 2D flocking hydrodynamics  $^{13c}$ 



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 $<sup>^{13</sup>b}$ Shvydkoy & ET Eulerian dynamics with commutator forcing. I, II and III (2017)  $^{13c}$ S. He & ET, Global regularity of 2D flocking hydrodynamics (2017)

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$$u_t + uu_x = \int \phi(|x - y|)(u(t, y) - u(t, x)\rho(t, y)dy$$

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• 1D alignment:

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- Riccati balanced by alignment:  $(d + \phi * \rho)' = -d(d + \phi * \rho)$
- Global smooth solution under a critical threshold condition iff  $u_0'$  "is not too negative":  $u_0'(x) + \phi * \rho_0(x) \geqslant 0$

<sup>&</sup>lt;sup>14</sup>Y.-P. Choi, J. Carrillo, E.T., C. Tan (2015)



$$(\mathsf{d} + \phi * \rho)' + \mathsf{d}(\mathsf{d} + \phi * \rho) = 0$$



$$(\mathbf{d} + \phi * \rho)' + \mathbf{d}(\overbrace{\mathbf{d} + \phi * \rho}^{\mathbf{e} := u_{\mathsf{x}} + \phi * \rho}) = 0$$



$$(d + \phi * \rho)' + d(d + \phi * \rho) = 0 \quad \Leftrightarrow \quad e' + u_x e = 0$$



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$$u_t + uu_x = [L_\beta, u](\rho), \quad L_\beta(\rho) = \text{p.v.} \int_{\mathbb{R}} \frac{\rho(y) - \rho(x)}{|x - y|^{1+\beta}} dy, \ \phi_\beta = |x|^{-(1+\beta)}$$



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$$e := u_x + L_\beta(\rho)$$
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$$\frac{e}{\rho}$$
 propagates along particle paths:  $(\frac{e}{\rho})' = 0$ 

$$\Rightarrow \frac{\mathsf{e}}{\rho} = \frac{\mathsf{e}_0}{\rho_0} := c_0$$
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- A bound on the density  $\rho_+(t) = \max_x \rho(x, t) < \infty$ :



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- A bound on the density  $\rho_+(t) = \max_x \rho(x, t) < \infty$ :

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ho_+$$

Quadratic growth  $-{
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ho_+\sim c_0
ho_+^2~$  vs. decay  $L_{eta}(
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ho_+,~-c_{eta}\ll-1$ 



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 'implies'  $u_{\mathsf{x}} > -\infty$  (at least along particle path)

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• In fact, substitution  $u_x = e - \Lambda^{\beta} \rho$  yields <u>diffusive</u> mass eq.

$$\rho_t + u\rho_x + e\rho = \rho\Lambda^{\beta}\rho, \qquad \Lambda^{\beta}g = \text{p.v.} \int_{\mathbb{R}} \frac{g(y) - g(x)}{|x - y|^{1+\beta}} dy$$

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ullet Global smooth solution 0 < eta < 2 with no CT; Singularity helps!

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• Global regularity (+ flocking) provided:  $\begin{cases} \operatorname{div}_{\mathbf{x}}(\mathbf{u}_0)(x) \geqslant -\phi * \rho_0(x), \\ \max_{x} |\eta(\mathcal{S}_0(x))| \leqslant \frac{1}{2}M_0\phi_{\infty} \end{cases}$ 

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Back to the fundamentals<sup>17</sup>



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   What about more realistic short-range interactions?
  - (with R. Shvydkoy): Singular kernels which are adapted to the density. Regular kernels: if the variation of the density  $\max \rho \min \rho$  is not too large relative to  $1 \widehat{\phi}(k)$  but independent of  $supp \, \phi$

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  - (with J. Morales) agents adjust to their own (internal) intensities but they synchronize the group orientations; unifying Cucker-Smale and Kuramoto



J. Morales

- A key point: in the discrete case requires propagation of connectivity; the density is bounded away from vacuum  $\int_0^\infty \rho_-(t)dt = \infty$
- Self-organized dynamics different type of "traits"
   Agents are equipped with traits that interact with the environment;
   there are 'genotype' traits that affect their readiness to interact.
  - (with J. Morales) agents adjust to their own (internal) intensities but they synchronize the group orientations; unifying Cucker-Smale and Kuramoto
- How different 'rules of engagement' dictate large time behavior?



KI-Net: Kinetic description of emerging challenges in multiscale problems of natural sciences





#### THANK YOU

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