



Jet Deceleration in Radio Galaxies and the FR Dichotomy

Manel Perucho

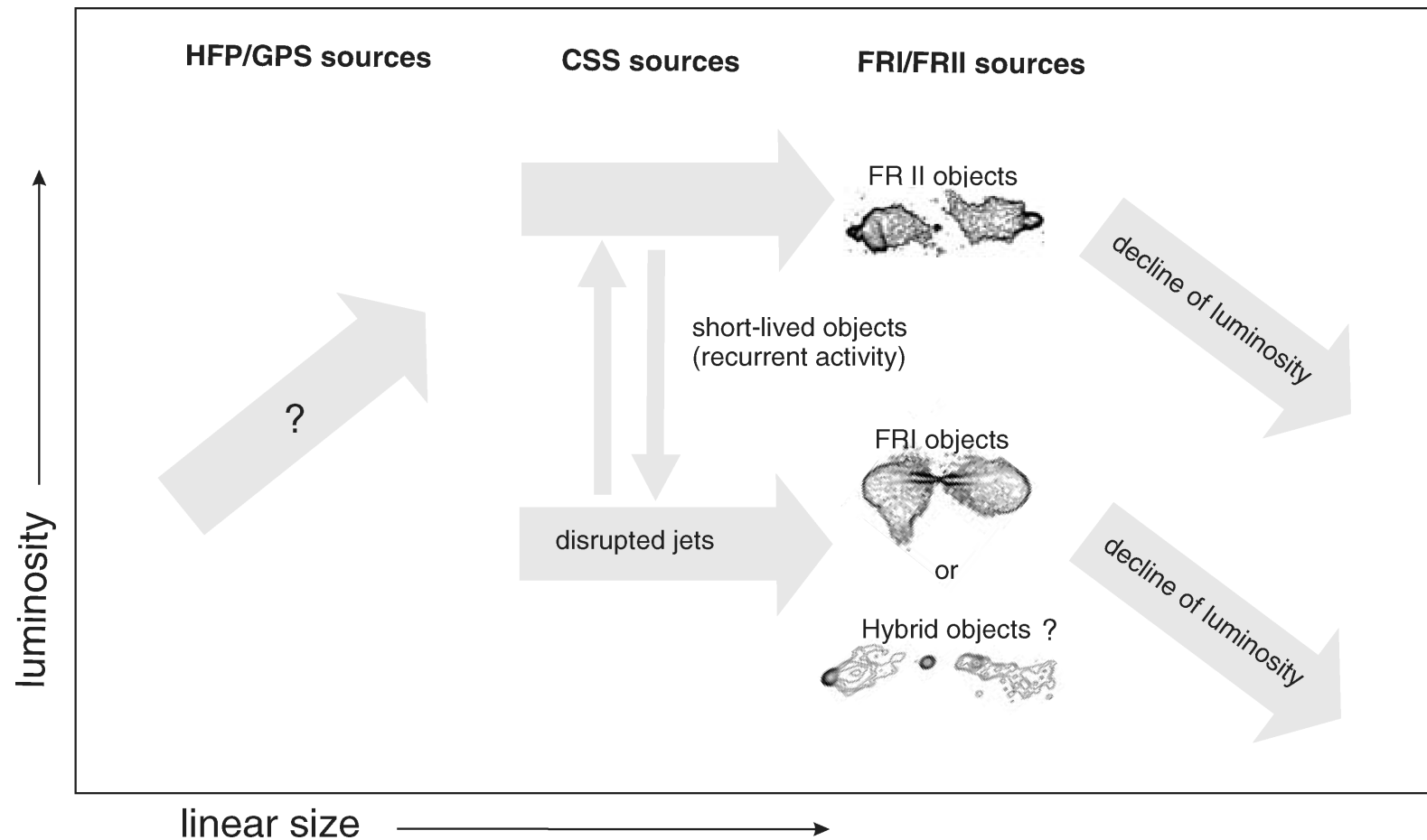
Computational Astrophysics and Cosmology group
DAA/OAUV

Universitat de València

Extragalactic Relativistic Jets: Cause and Effect.

Bangalore, October 16th 2015.

Radio-source evolution: PD diagram

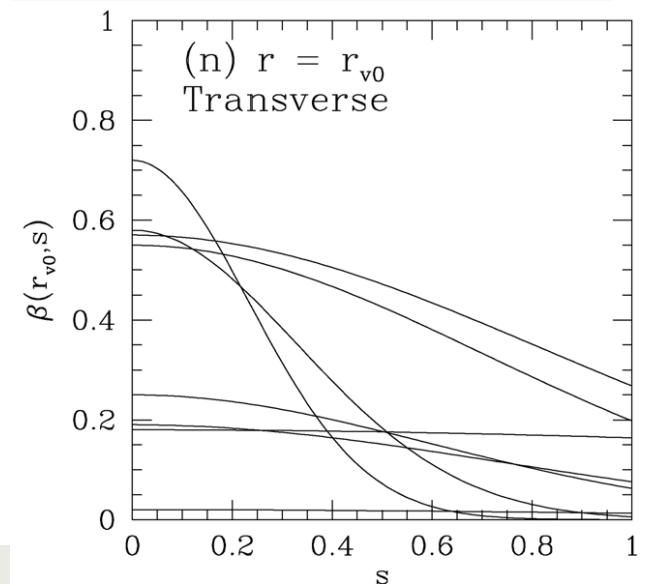
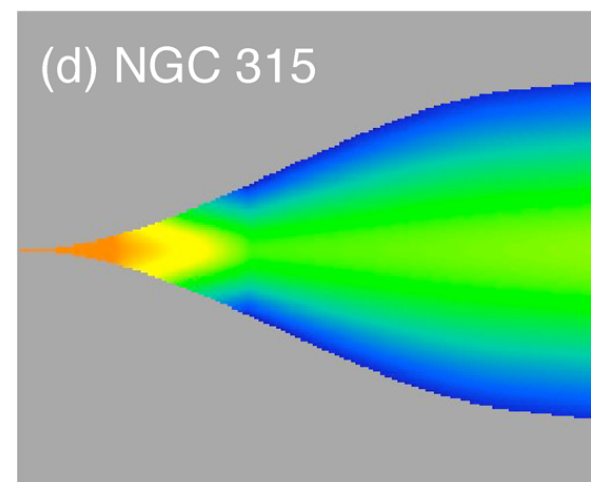
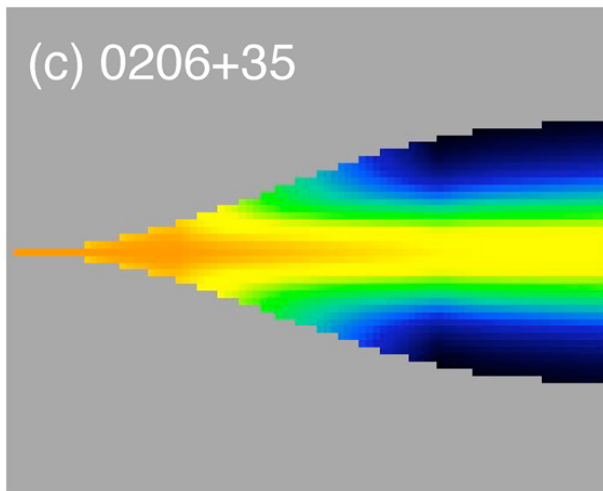
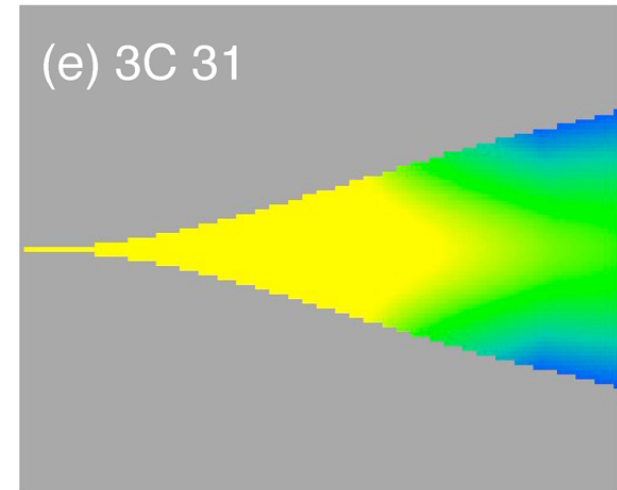
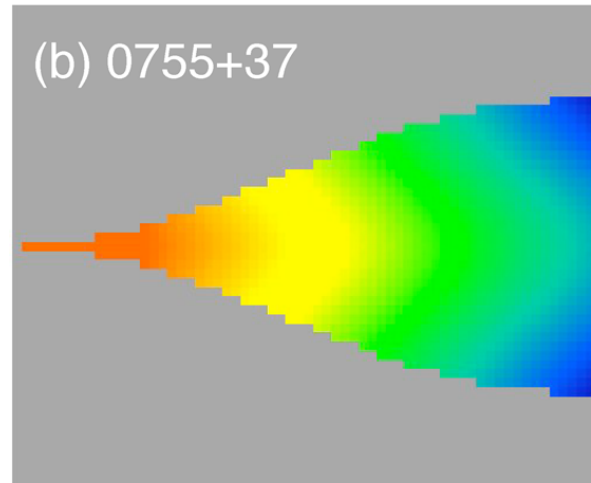
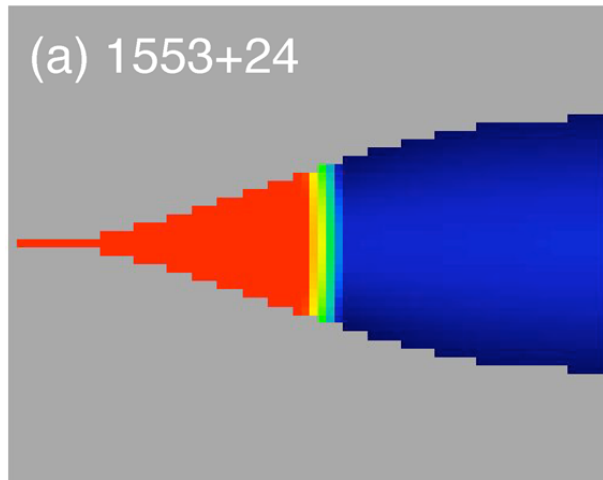


From Kunert-Bajraszewska et al. (2010)

Deceleration of FRI jets

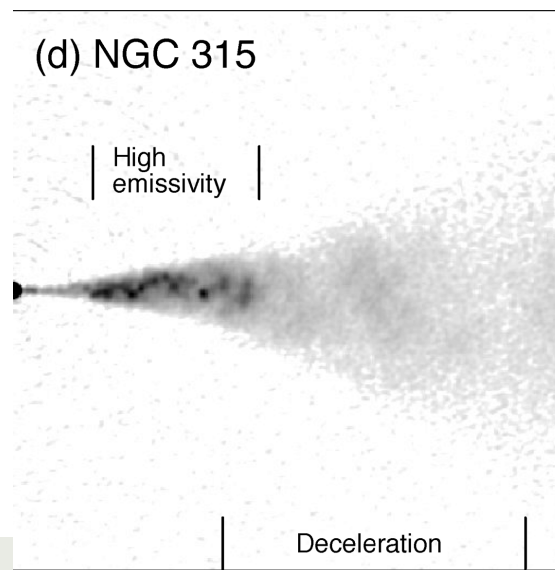
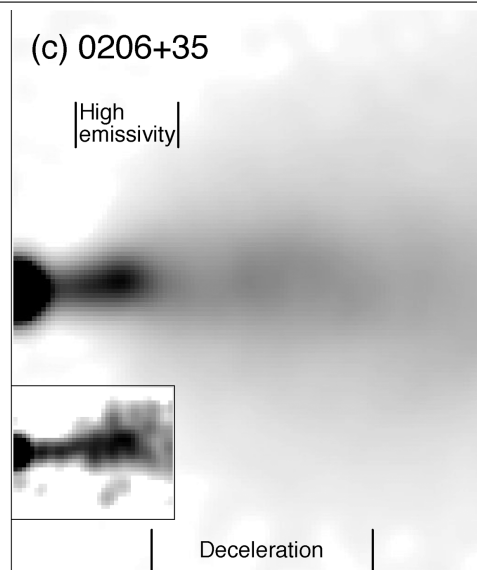
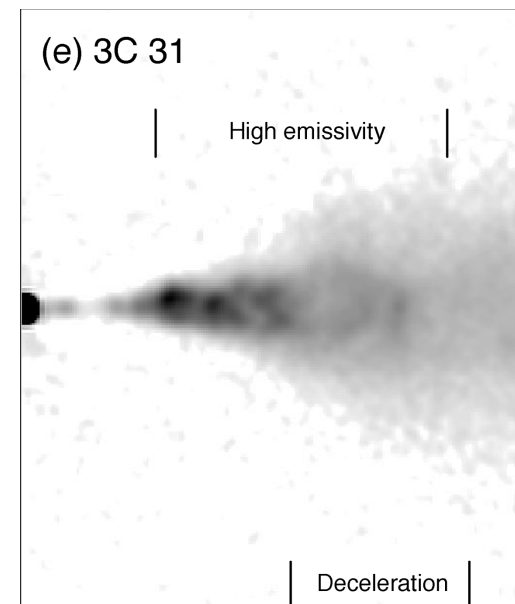
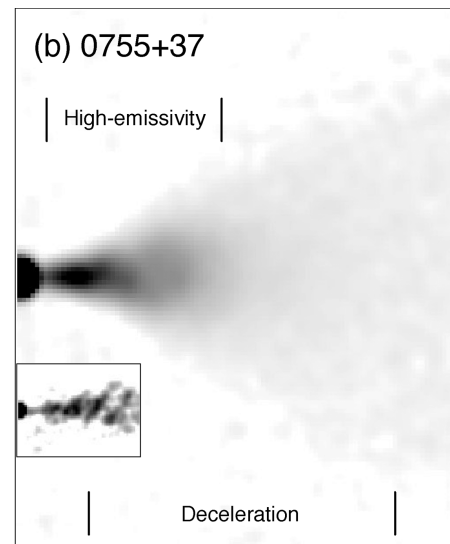
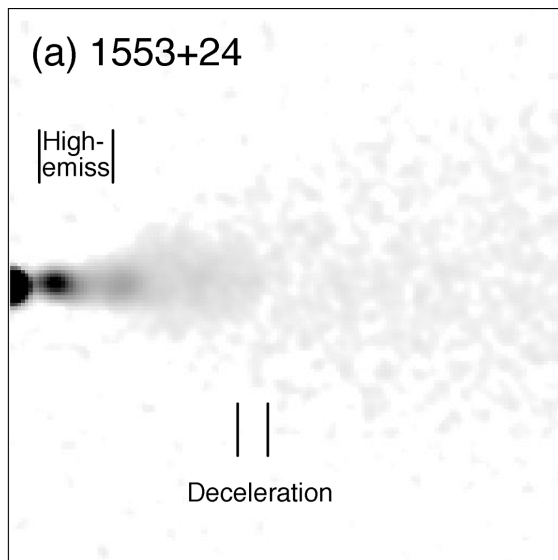


Laing & Bridle 2014



Deceleration of FRI jets

Laing & Bridle 2014



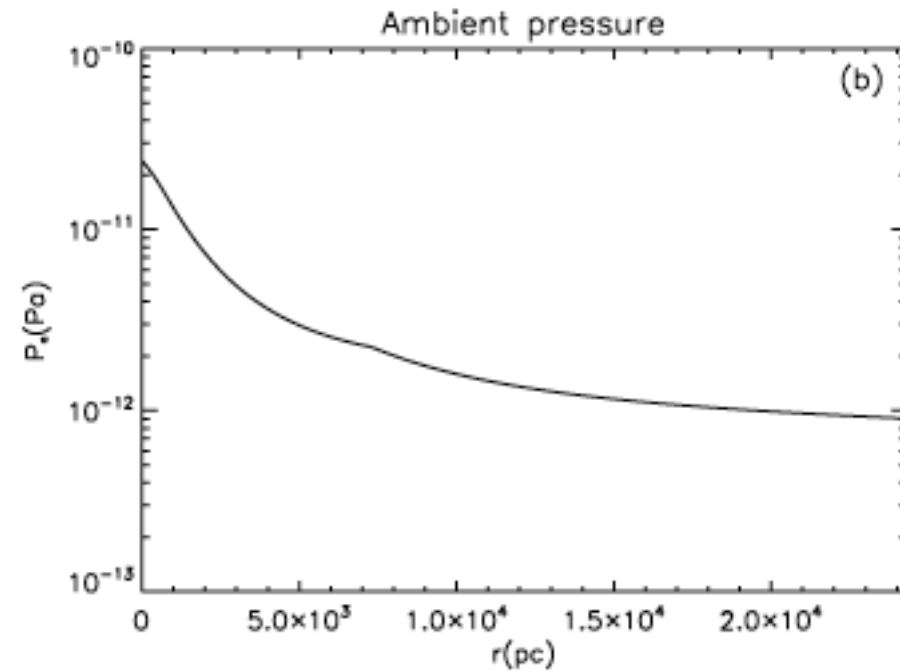
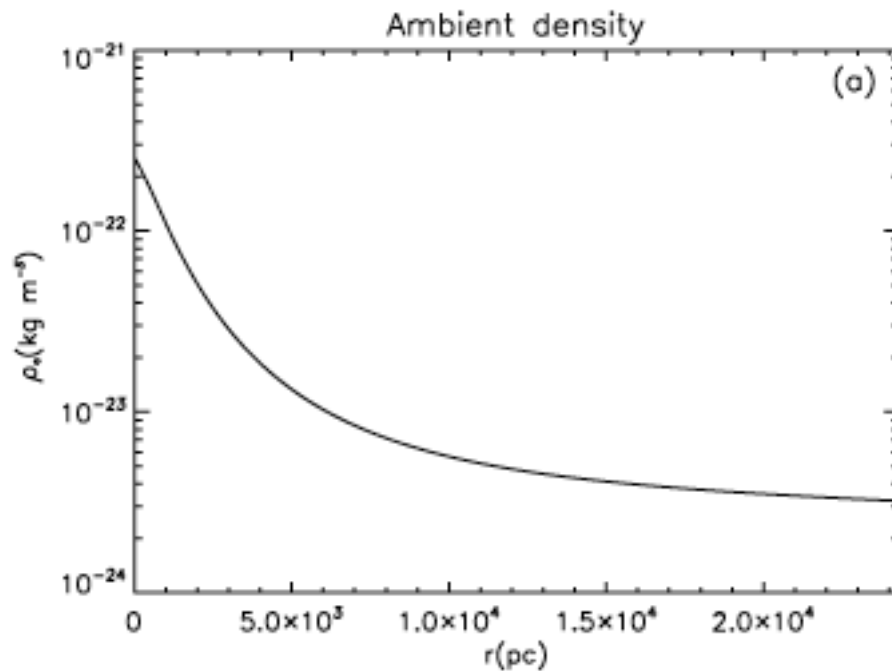
Kharb et al. (2012) find X-ray emission from these regions In 15/21 FRI's (nine X-ray jets).

Deceleration: role of the ambient medium

$$n_{\text{ext}}(r) = n_c \left(1 + \frac{r^2}{r_c^2}\right)^{-3\beta_{\text{atm},c}/2} + n_g \left(1 + \frac{r^2}{r_g^2}\right)^{-3\beta_{\text{atm},g}/2}$$

$$T = T_c + (T_g - T_c) \frac{r}{r_m} \quad (r < r_m)$$

$$T = T_g \quad (r \geq r_m),$$

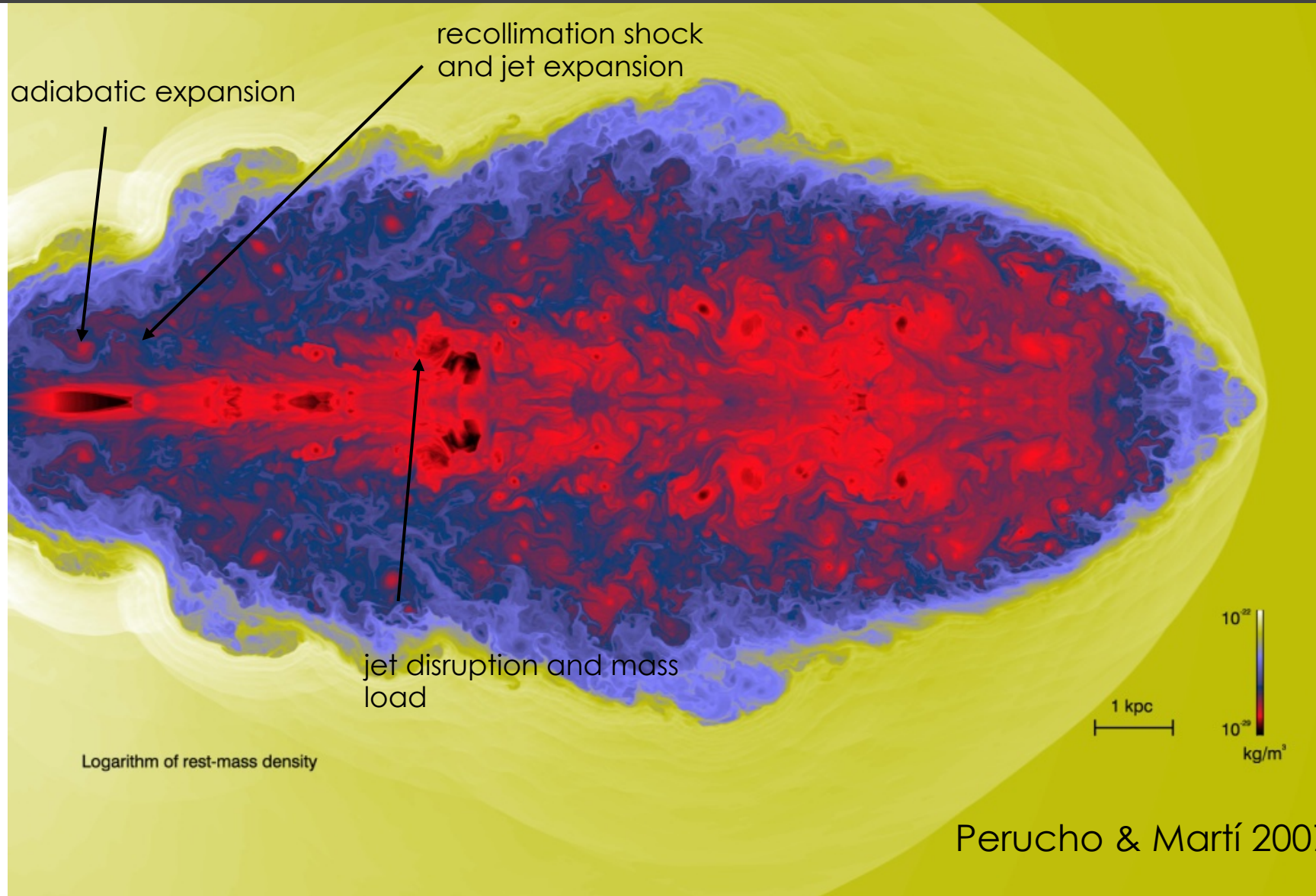


Perucho
& Martí
(2007)

Component	Central density	Form factor	Core radius	Temperature
Galaxy	$n_c = 1.8 \times 10^5 \text{ m}^{-3}$	$\beta_{\text{atm},c} = 0.73$	$r_c = 1.2 \text{ kpc}$	$T_c = 4.9 \times 10^6 \text{ K}$
Group	$n_g = 1.9 \times 10^3 \text{ m}^{-3}$	$\beta_{\text{atm},g} = 0.38$	$r_g = 52 \text{ kpc}$	$T_g = 1.7 \times 10^7 \text{ K}$

From Hardcastle et al. (2002)

Deceleration: role of the ambient medium

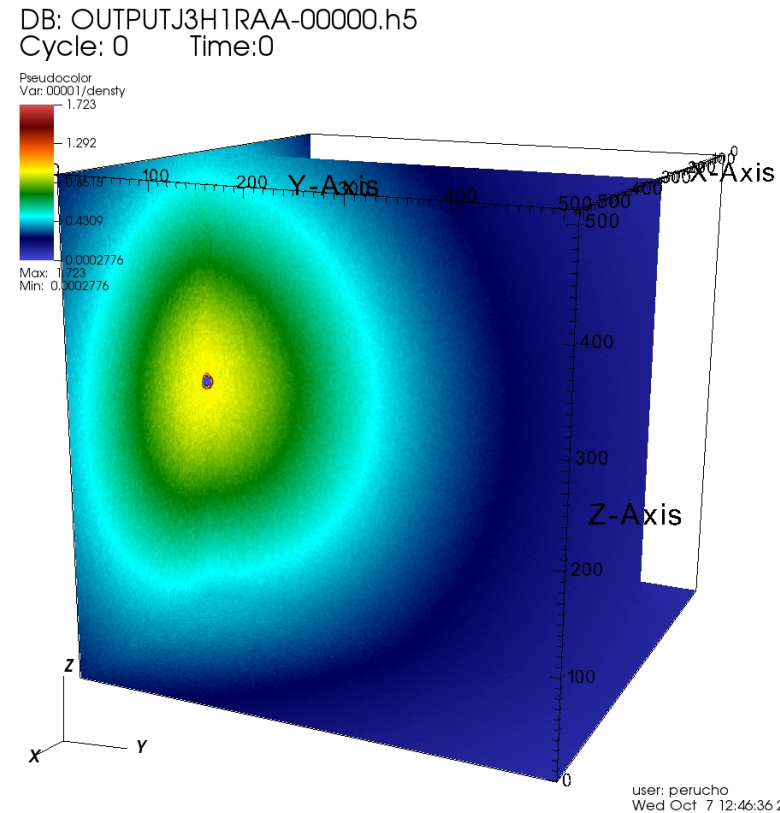
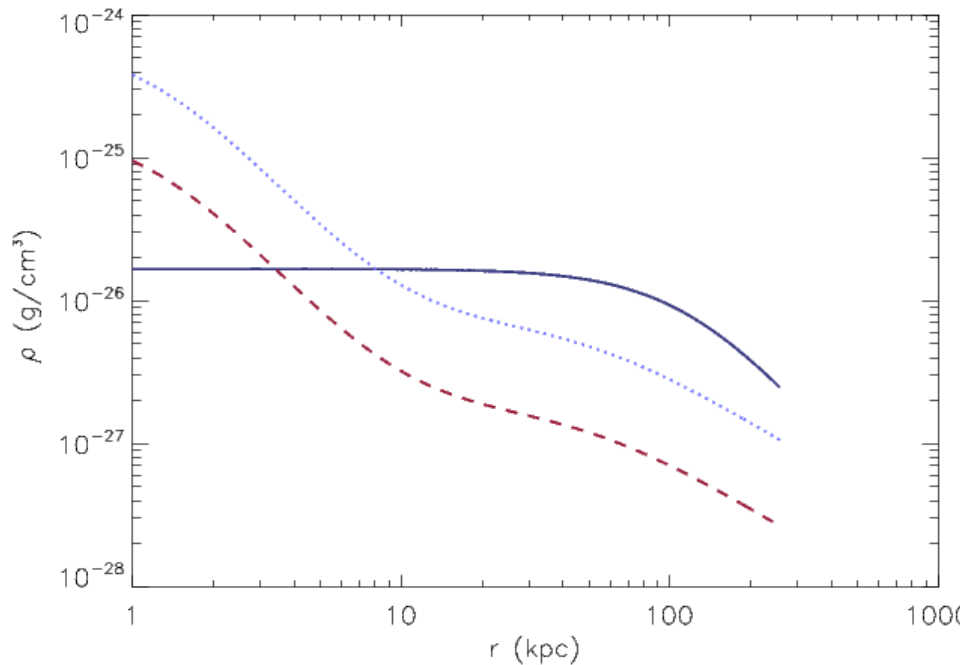


Deceleration: ambient medium and instabilities

dashed-red line: ambient density profile in the 2D simulations.

$L_j = 10^{46}$ erg/s, $\rho_j = 10^{-4} \rho_{a,0}$, $\rho_{e,j} = 0.8 m_p / \text{cm}^3$,
 $v_j = 0.92 c$, $T_a = 5 \cdot 10^7$ K.

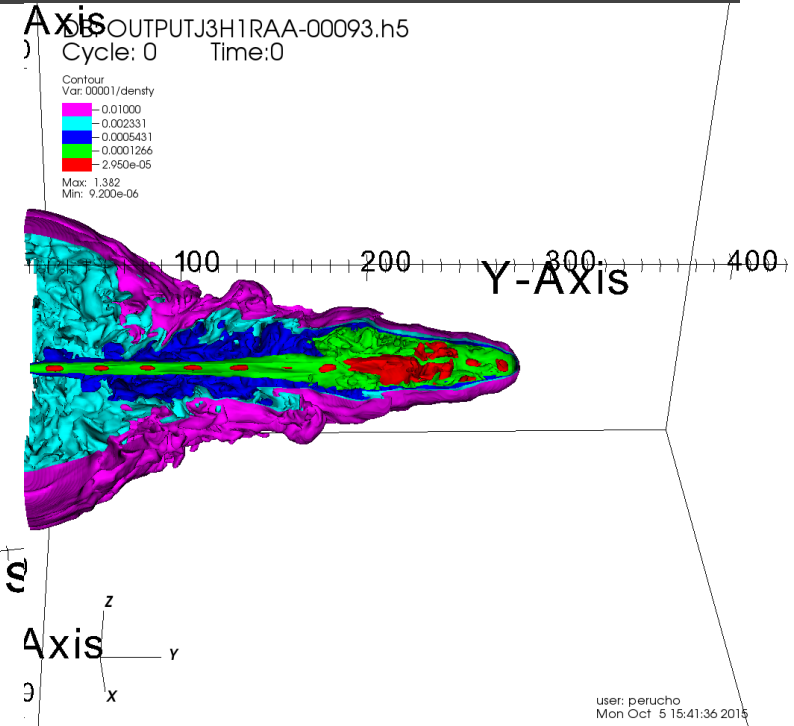
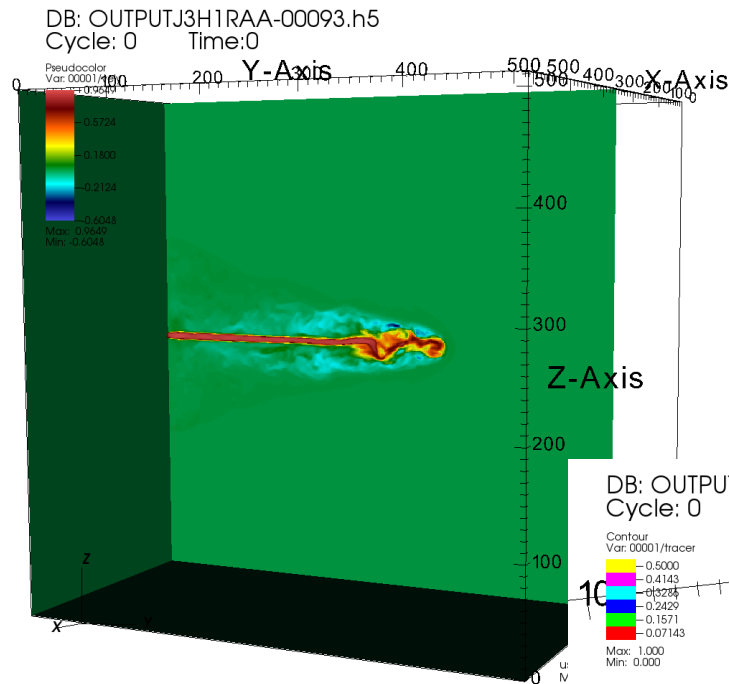
solid-blue line: ambient density profile in the 3D simulation.



Perucho et al., in preparation

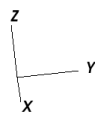
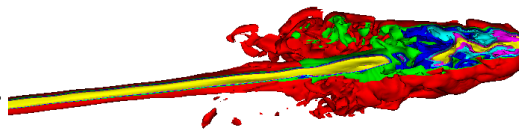
Ratpenat is a HRSC, **3D RHD** code that combines MPI and OMP parallelization. It includes a **relativistic equation of state**, which allows us to introduce different populations of particles.

Deceleration: ambient medium and instabilities



user: perucho
Mon Oct 5 15:41:36 2015

The jet develops pinching and (induced) helical instabilities and it is disrupted beyond the density core.

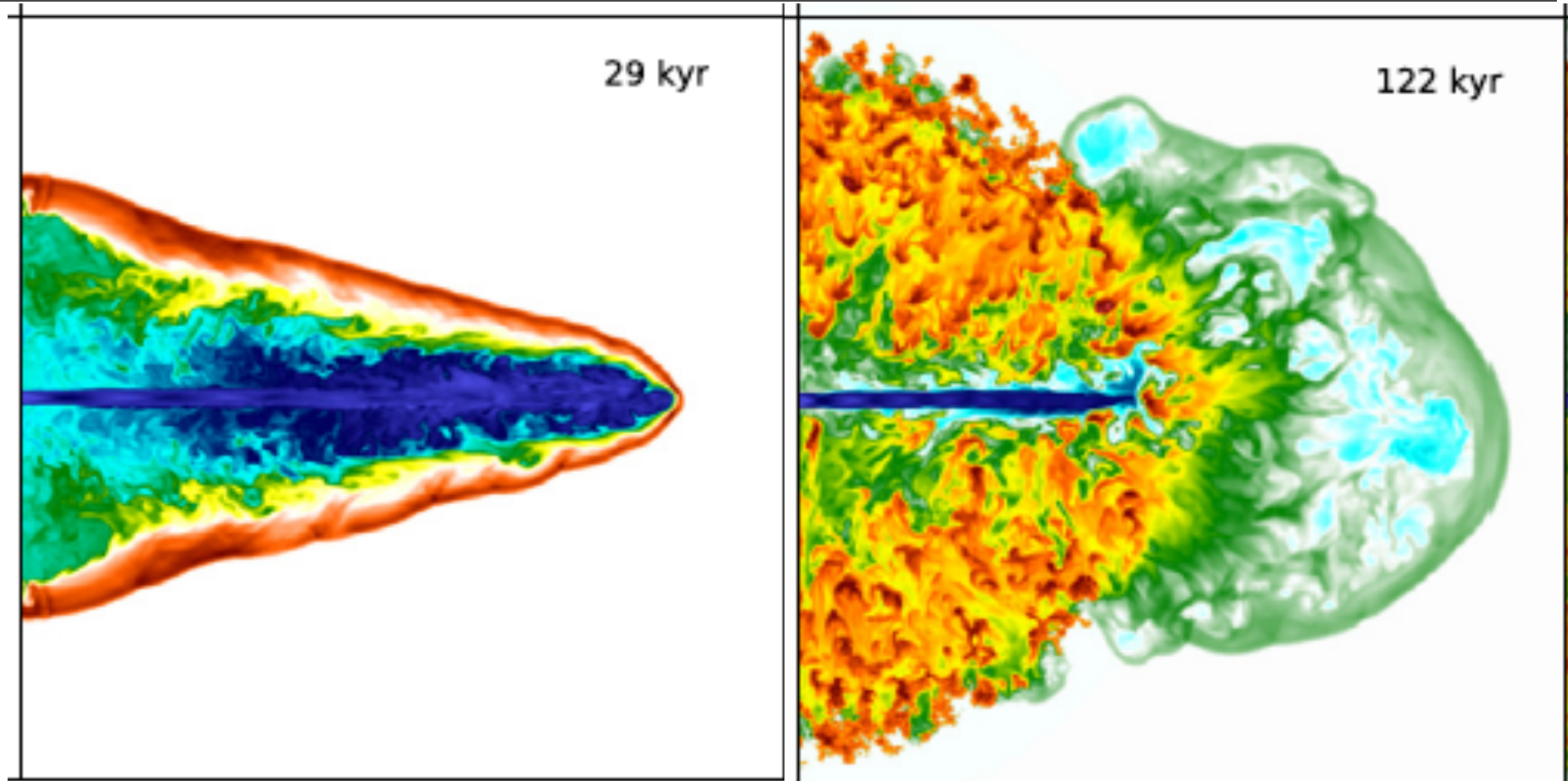


Deceleration of the jet head and development of fat lobes is expected following this disruption.

See Rossi et al. (2008) for the case of a homogeneous ambient medium.

user: perucho
Mon Oct 5 17:13:39 2015

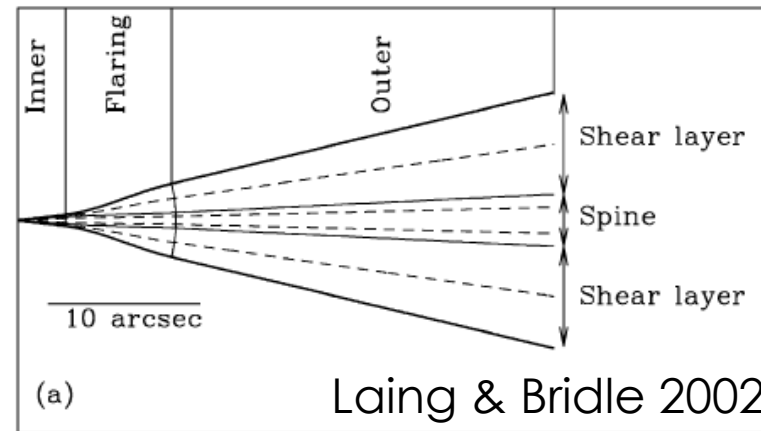
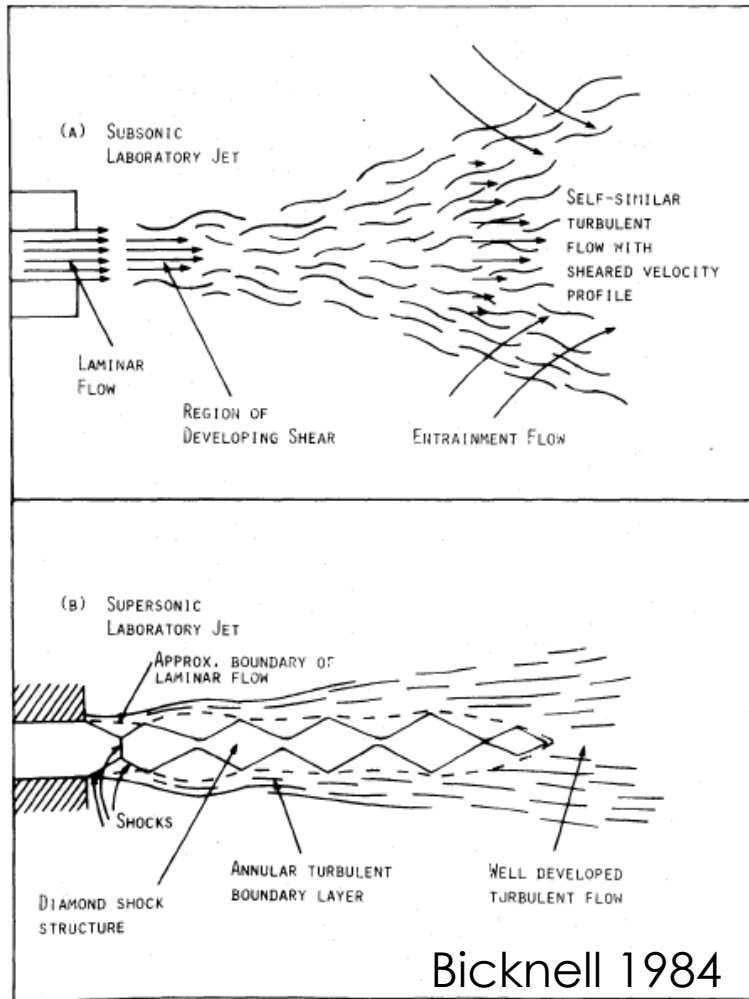
Deceleration: inhomogeneous media



Wagner & Bicknell (2011)
Wagner et al. (2012)

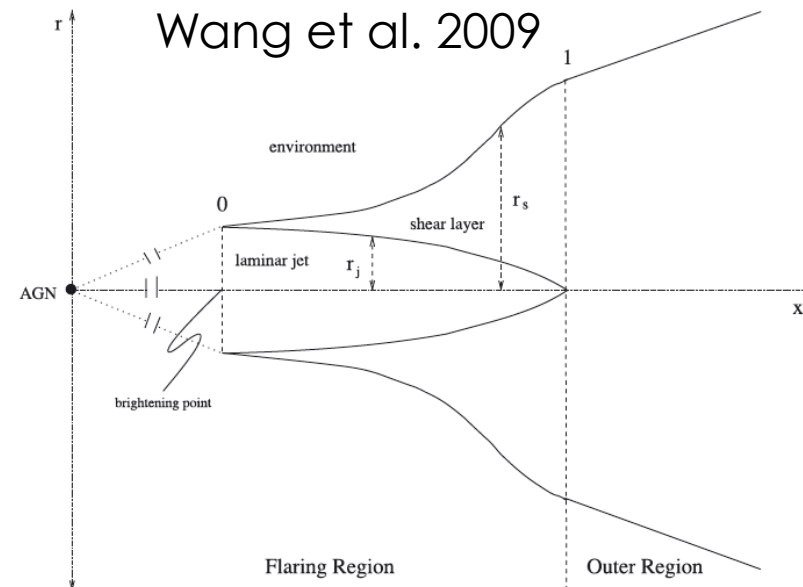
Homogeneous versus clumpy medium

Deceleration: instabilities

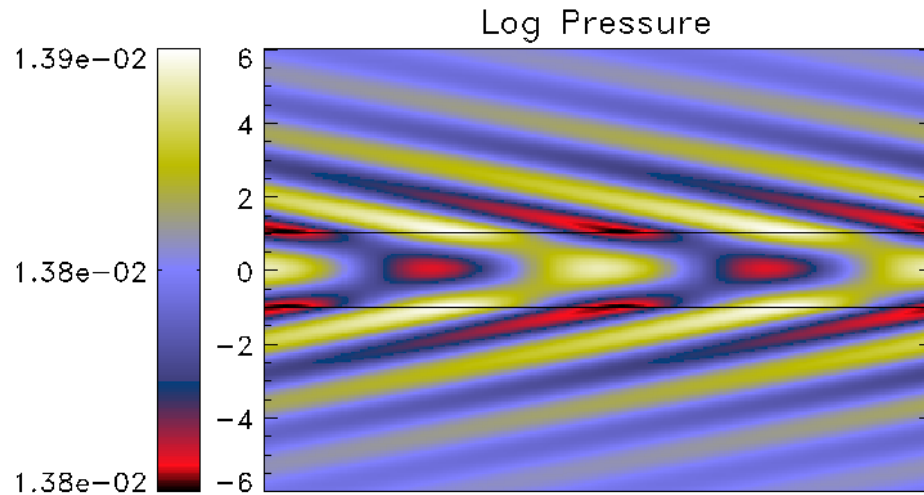


Laing & Bridle 2002, 2014

Shear-layer loading by turbulent mixing.

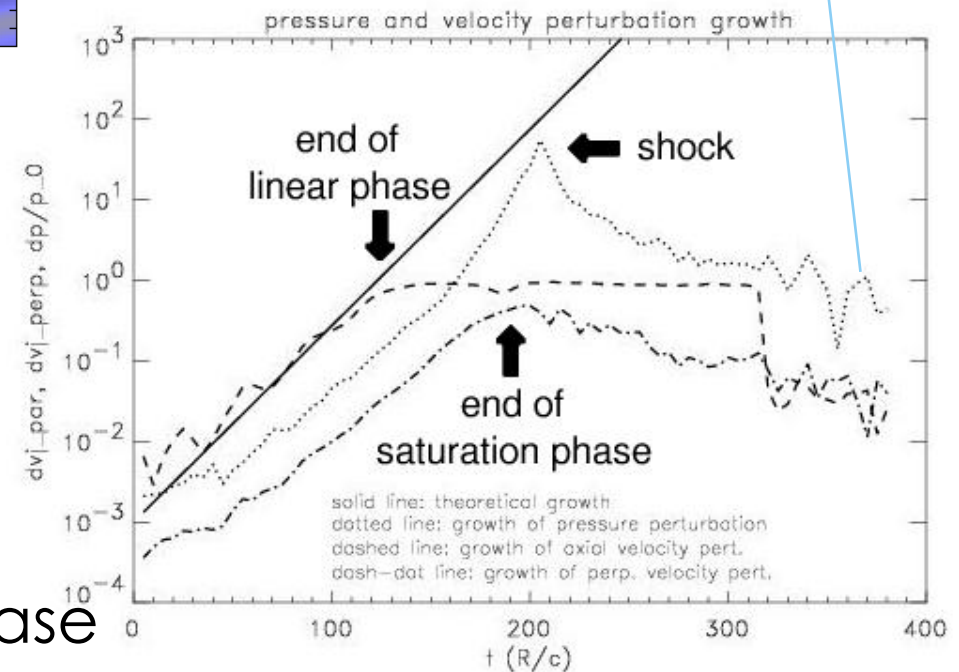
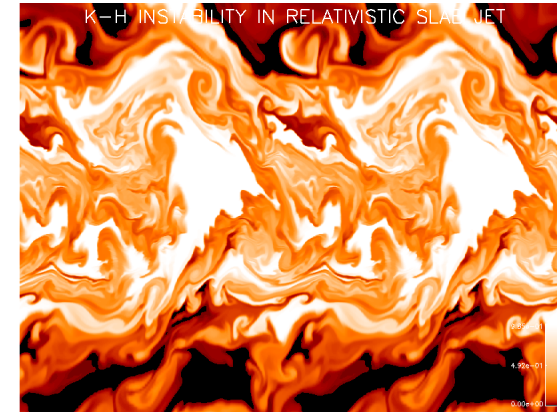


Deceleration: instabilities



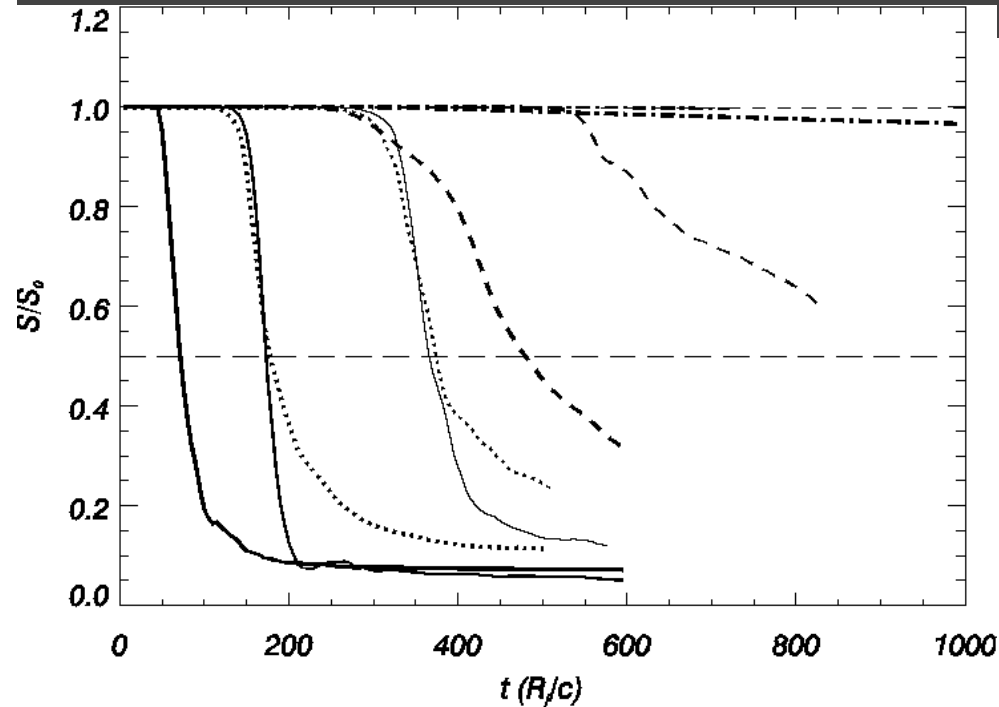
Initial model

- Parameters:
 - Lorentz factor.
 - Rest-mass density contrast.
 - Specific internal energy.
 - Pressure equilibrium.



Linear phase

Deceleration: instabilities

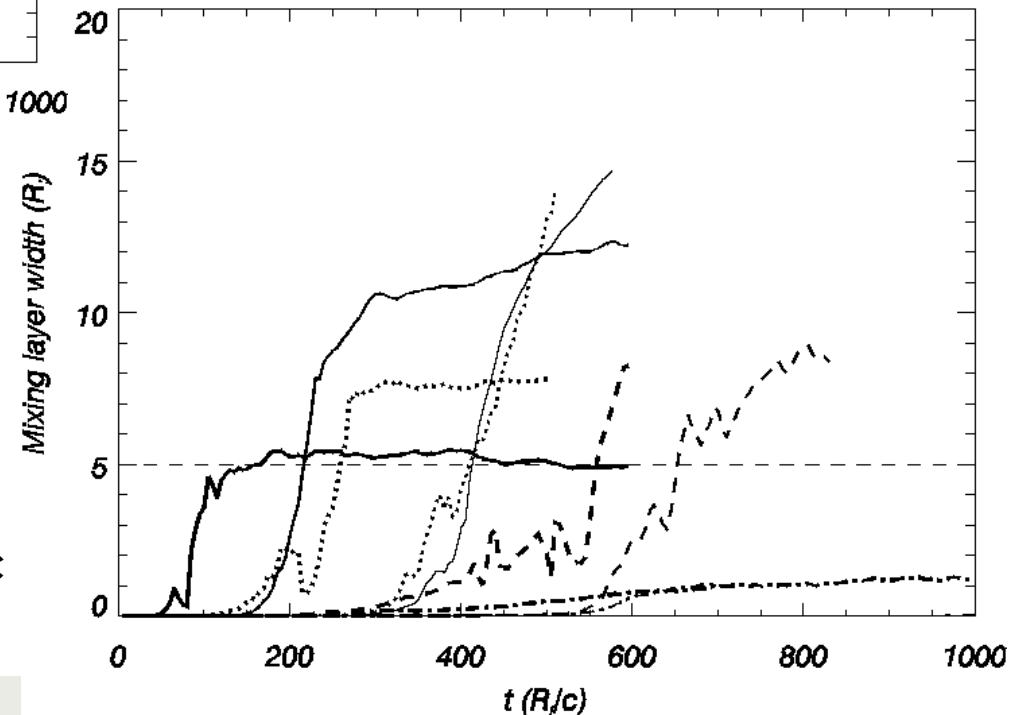


2D simulations of temporal evolution of KH instabilities in relativistic flows.

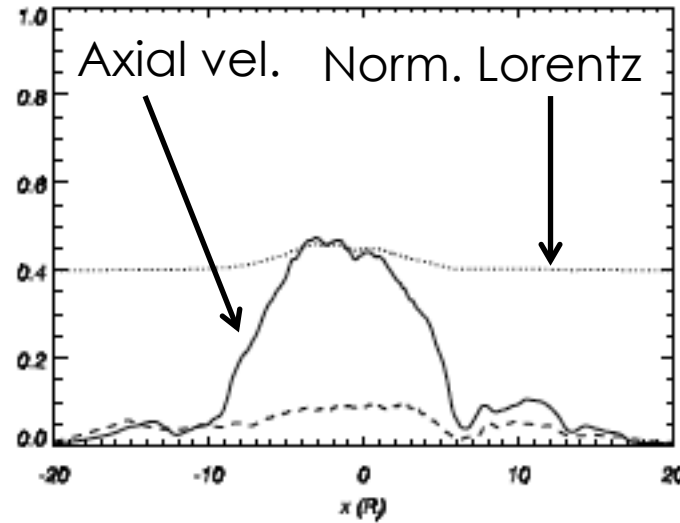
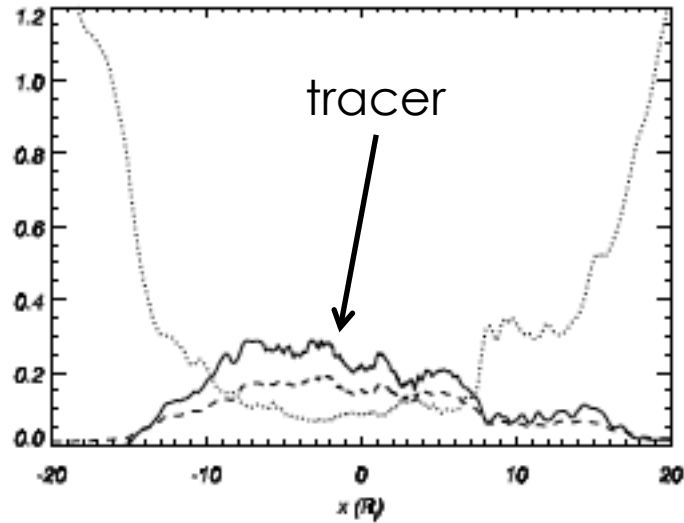
Deceleration distance 300 pc – 3 kpc

Perucho, Martí, Hanasz (2005),
Perucho et al. (2010):

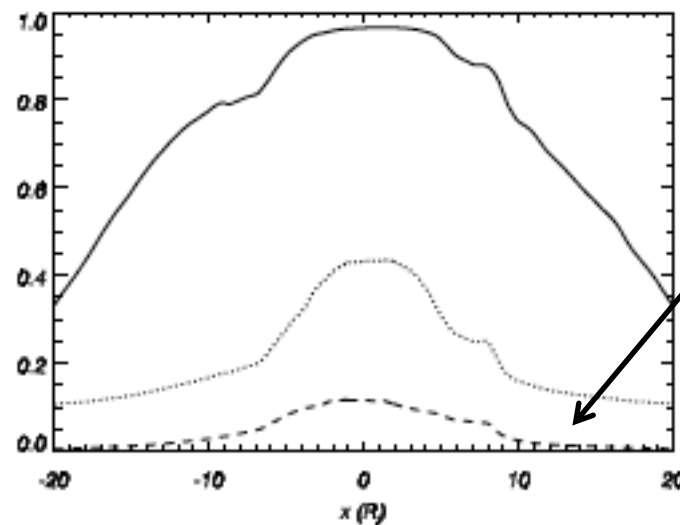
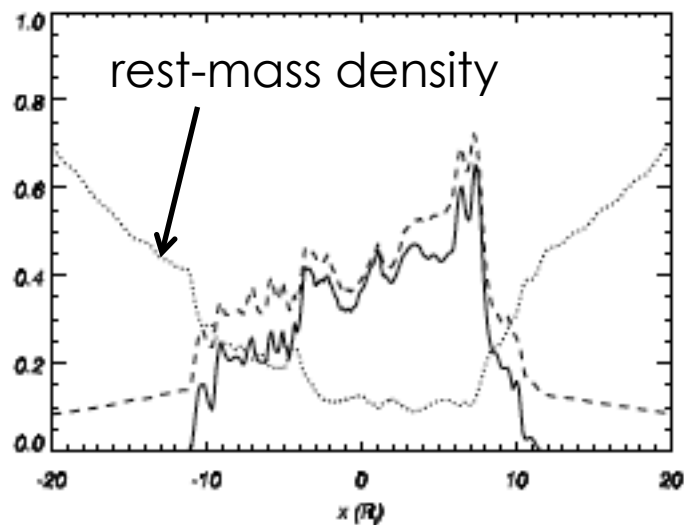
KH instability. The disruption process can be slow if the growing modes have small wavelengths.



Deceleration: instabilities



Profiles across the jet in 2D simulations.

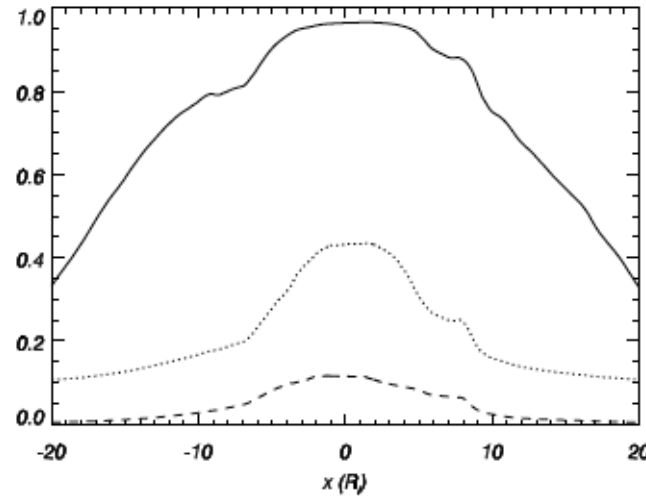
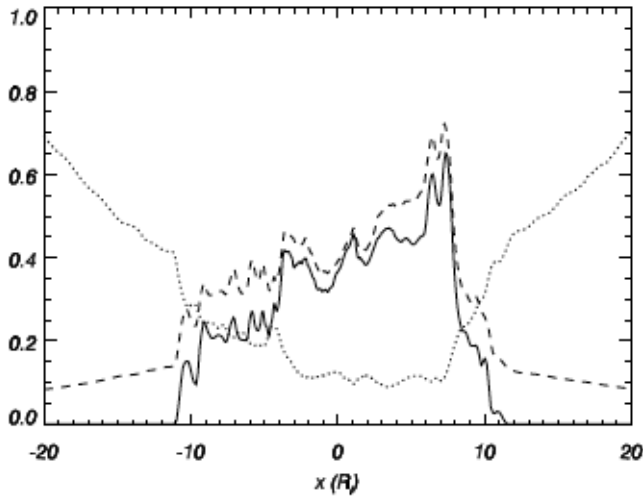


Norm. Momentum In the axial direction

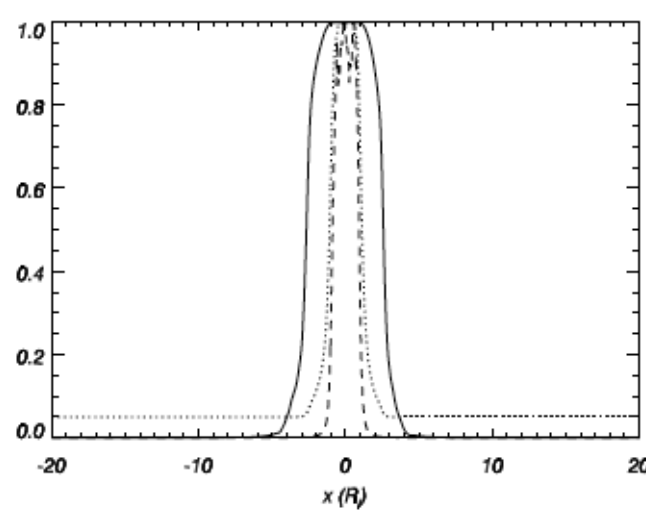
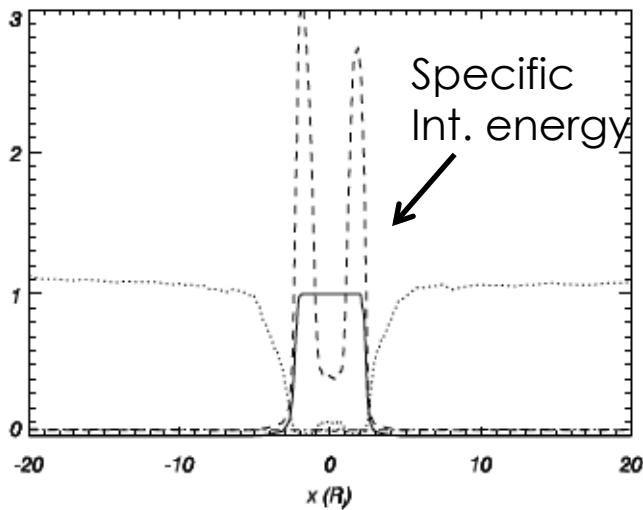
Left panel (thermodynamical quantities): solid line, tracer; dotted line, rest mass density; dashed line, specific internal energy.

Right panels solid line, longitudinal velocity; dotted line, lorentz factor normalized to the initial value in the jet; dashed line, longitudinal momentum normalized to the initial value in the jet.

Deceleration: instabilities



Profiles across the jet in 2D simulations.

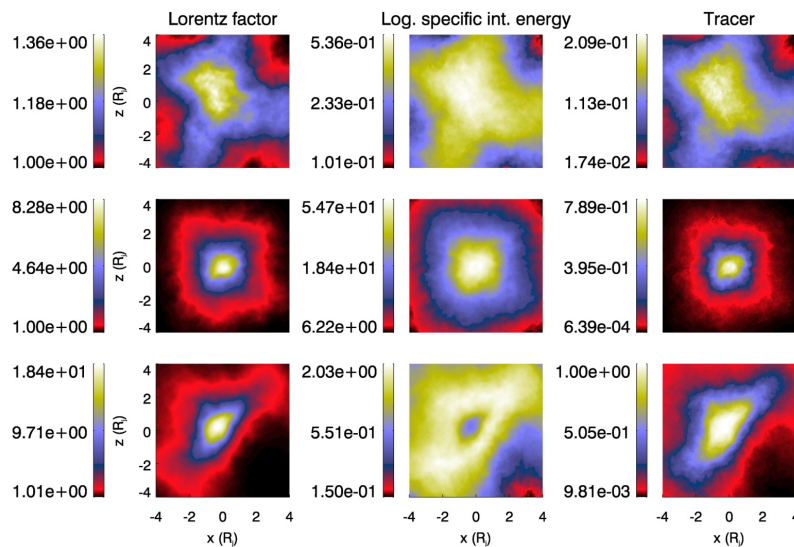
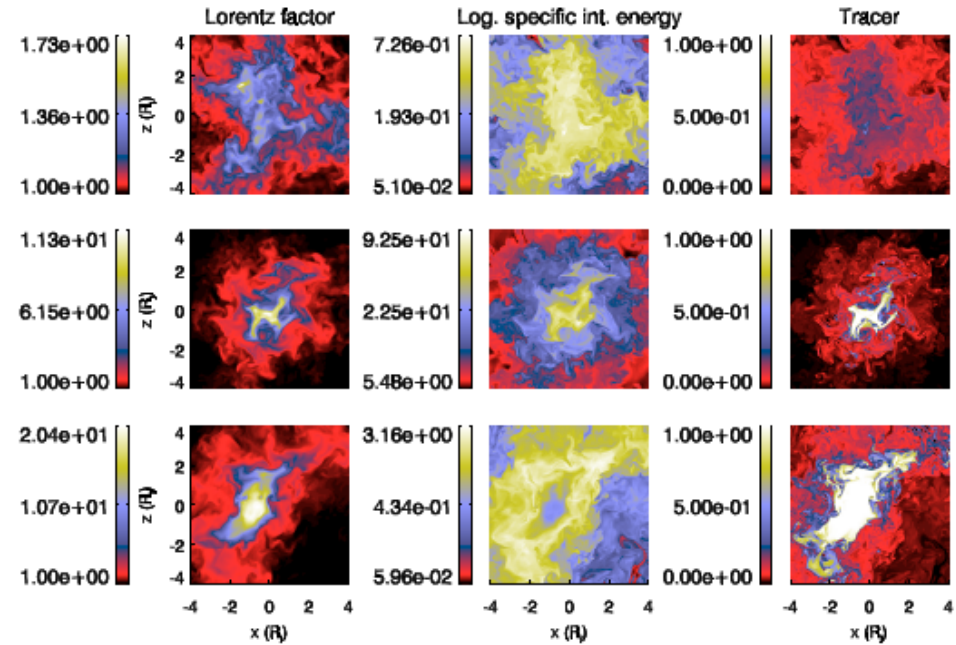
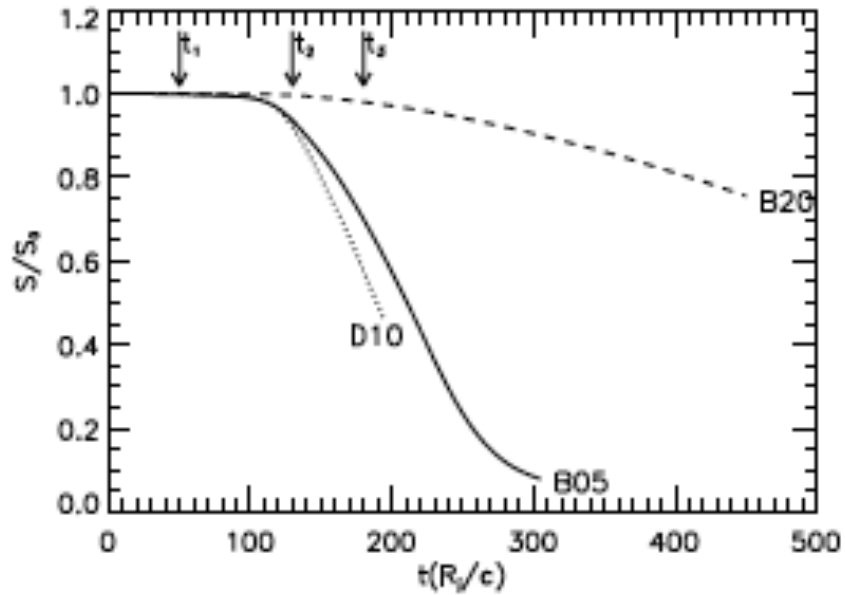


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Deceleration: instabilities

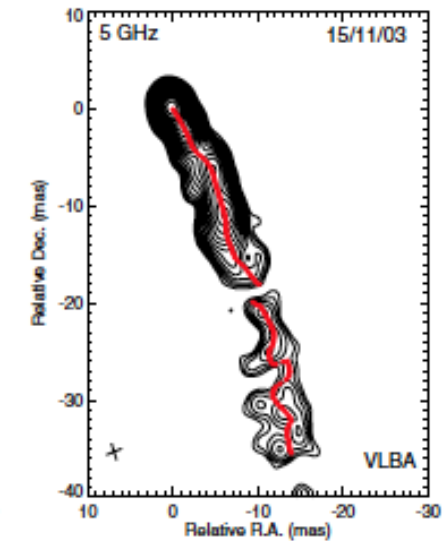
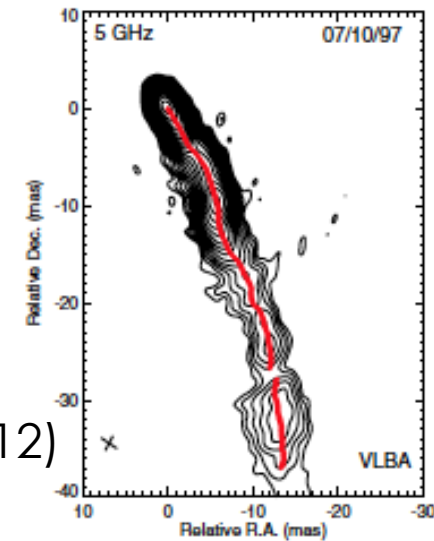
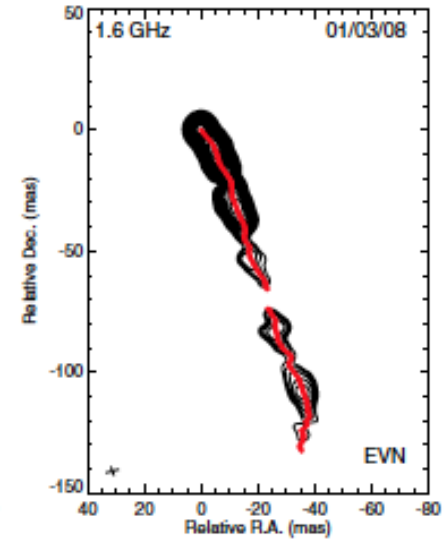
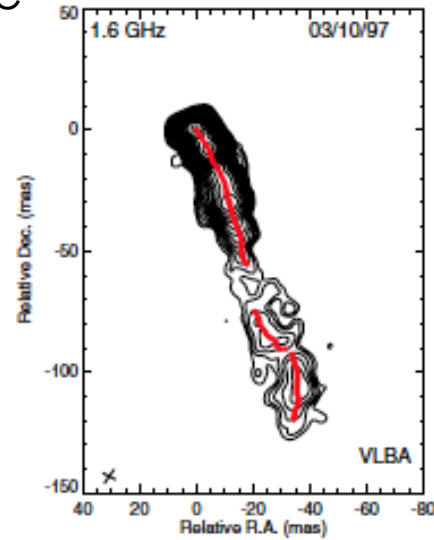
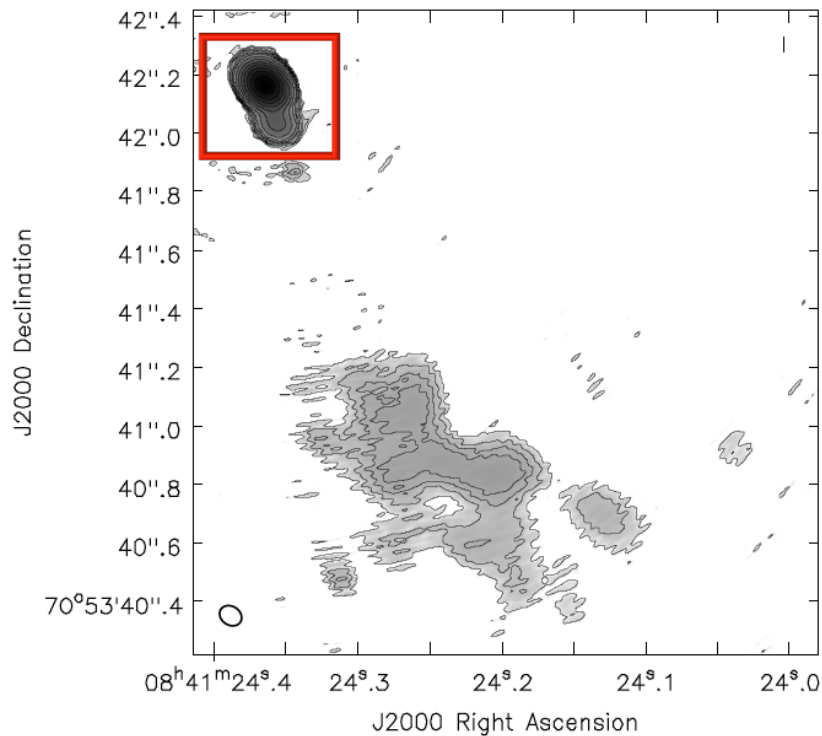


cold
 hot
 cold
 Lorentz factor

Perucho et al. 2005, 2010: KH instability. The disruption process can be slow if the growing modes have small wavelengths.

Deceleration: instabilities

0836+710: The jet is possibly disrupted by the growth of instabilities.

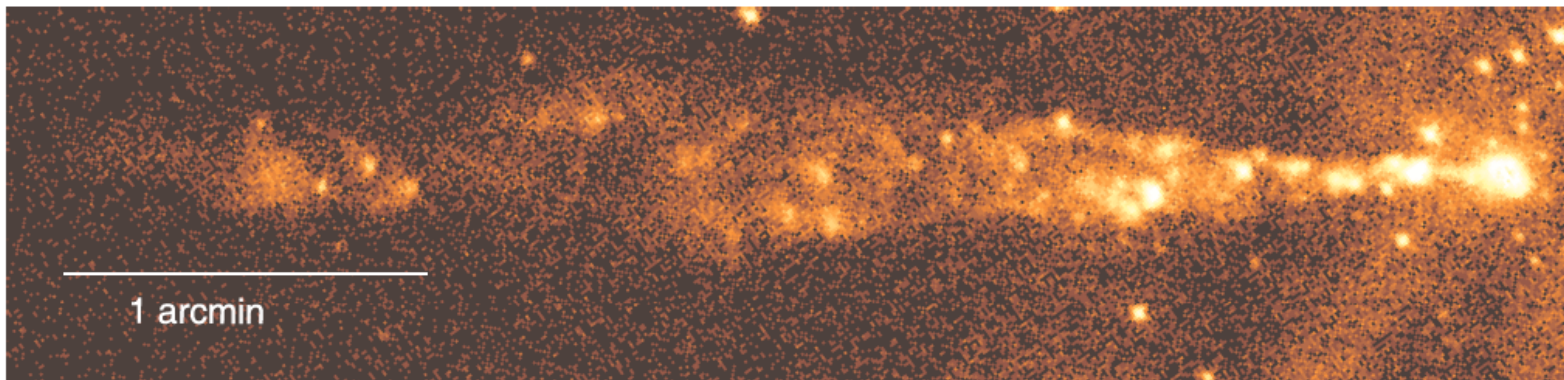
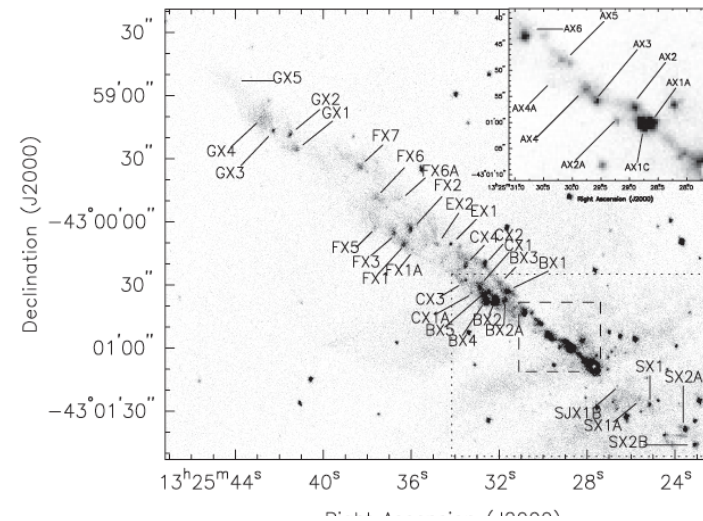
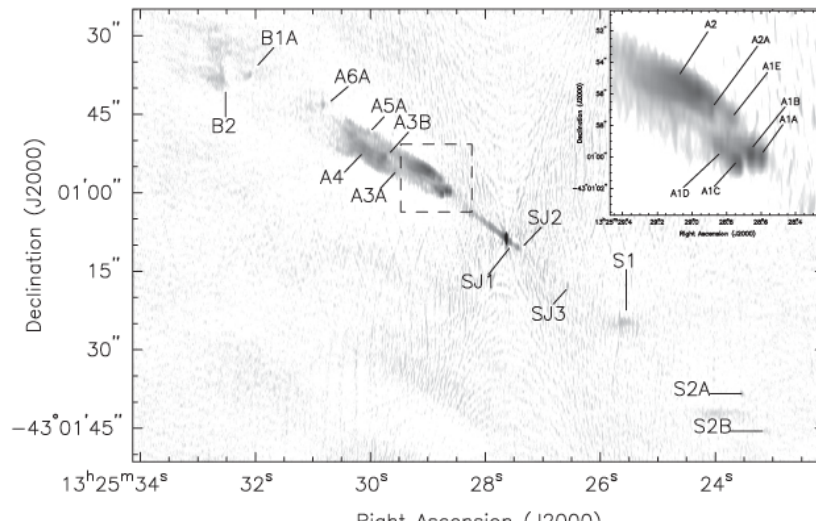


MP, Kovalev, Lobanov, Hardee, Agudo (2012)

MP, Martí-Vidal, Lobanov, Hardee (2012)

Deceleration: mass load by stellar winds

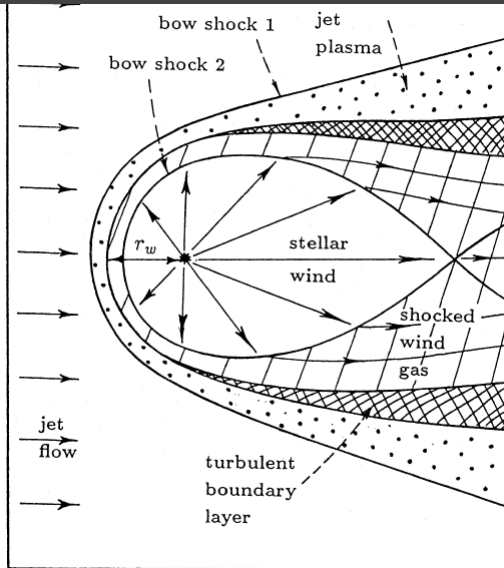
Worrall et al. 2008, Goodger et al. 2010, Wykes et al. 2013, 2015, Müller et al. 2014.
Possible interaction with obstacles in Centaurus A: Clouds, O/B type stars?



SEE TALK BY
S. WYKES

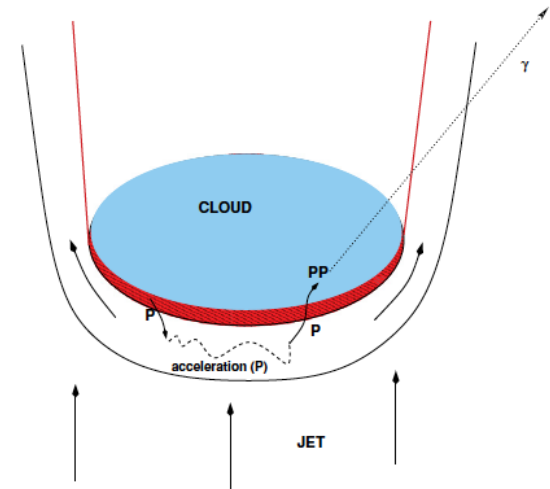
FIG. 1.—A 0.8–3 keV unsmoothed rotated image of the six *Chandra* exposures comprising the Cen A VLP, with 0.492×0.492 arcsec² pixels.

Deceleration: mass load by stellar winds



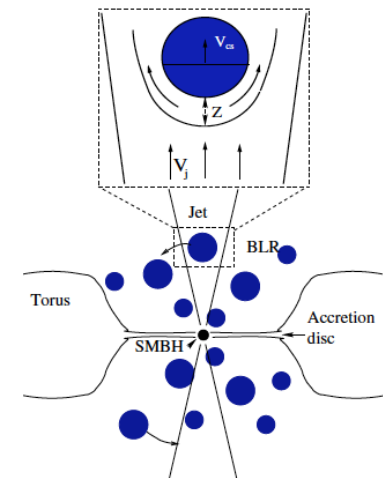
FRI deceleration due to mass-load of stellar wind gas?

Komissarov 1994 (RHD problem),
Bowman, Leahy & Komissarov 1996,
Hubbard & Blackman 2006



Interaction between jets and clouds invoked to explain

high-energy emission: Bednarek, Protheroe 1997
Araudo et al. 2010
Barkov et al. 2010, 2012
Wykes et al. 2013, 2015

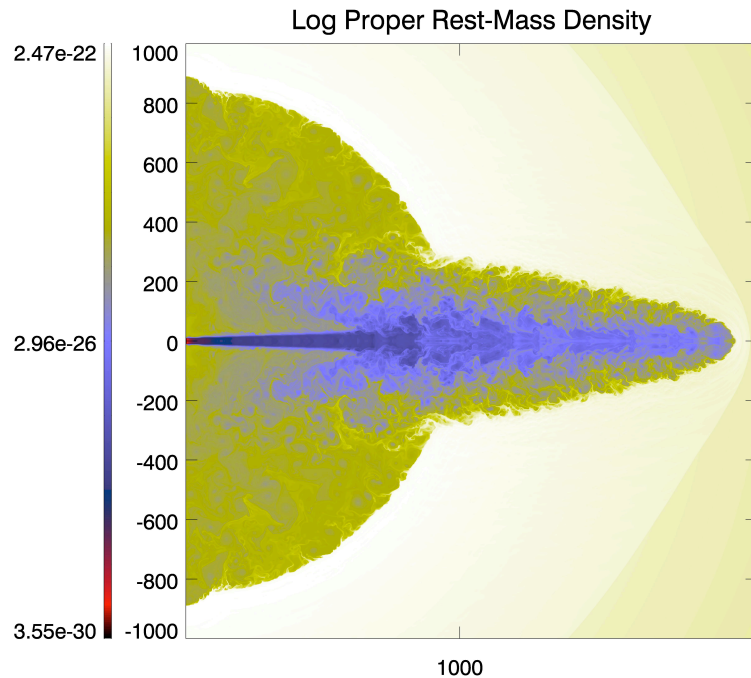


Deceleration: mass load by stellar winds

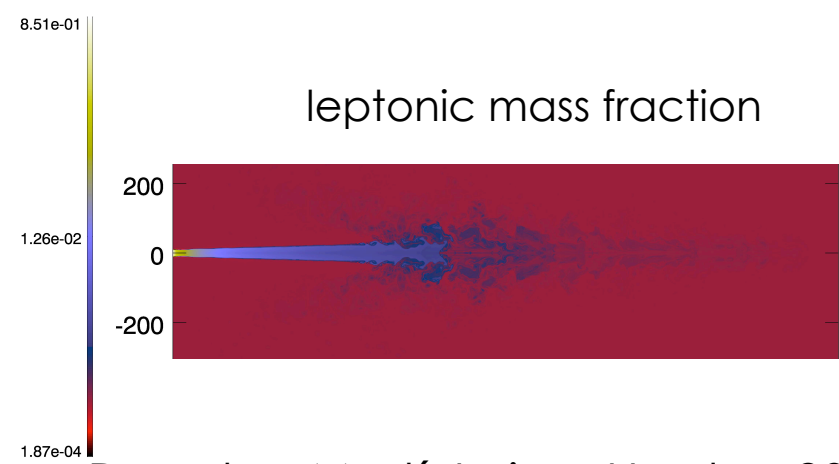
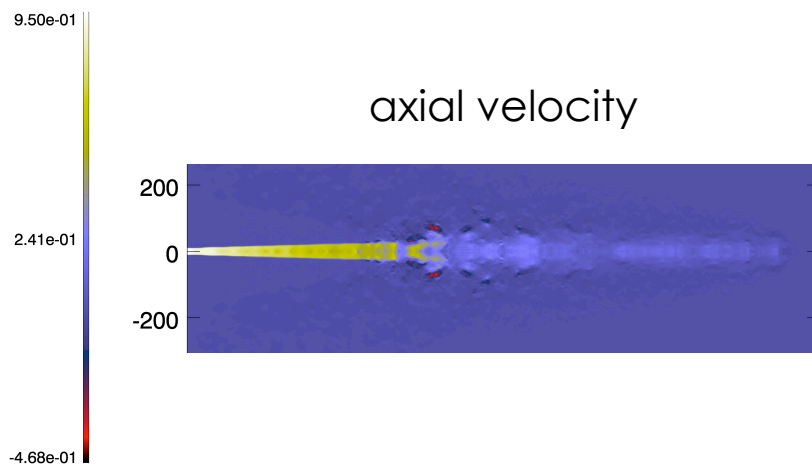
Following Komissarov (1994), Bowman et al. (1996), we performed simulations of FRI jets with a source term in mass accounting for mass-load from stellar wind.

Laing & Bridle (2002) studied this case in 3C31 (10^{44} erg/s): concluded it was not enough to decelerate the jet flow.

Model	Velocity [c]	Density [g/cm ³]	Temperature [K]	P_j/P_{am}	L_k [erg/s]	q_0 [gyr ⁻¹ pc ⁻³]	Jet length [kpc]	t_{sim} [Myrs]
Po	0.99	9.65×10^{-30}	3×10^9	17.6	10^{44}	4.95×10^{22}	2.2	-
Pr	0.99	9.65×10^{-30}	3×10^9	17.6	10^{44}	4.95×10^{22}	2.2	-
A0	0.95	3×10^{-33}	3×10^{11}	0.54	5×10^{41}	0	1.5	1.6
A	0.95	3×10^{-33}	3×10^{11}	0.54	5×10^{41}	4.95×10^{22}	2.1	24.0
B	0.95	3×10^{-34}	3×10^{12}	0.54	5×10^{41}	4.95×10^{22}	2.0	21.0
C	0.95	3×10^{-35}	3×10^{13}	0.54	5×10^{41}	4.95×10^{22}	1.8	19.0
D	0.95	3×10^{-35}	3×10^{13}	0.54	5×10^{41}	4.95×10^{21}	1.8	18.0



injection point at 80 pc – initial jet radius 10 pc



King density profile.

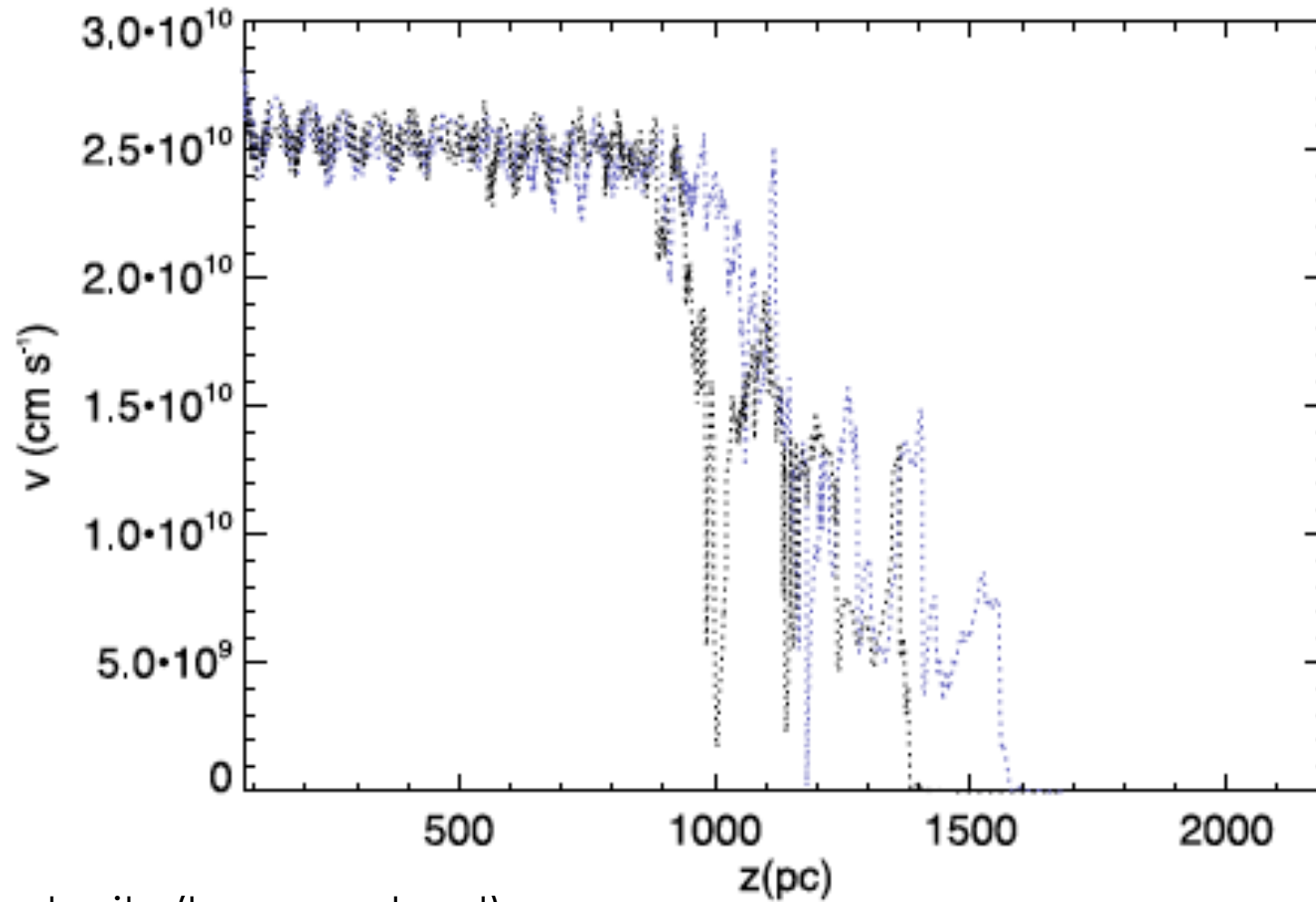
$$n_{ext} = n_c \left(1 + \frac{r^2}{r_c^2} \right)^{-3\beta_{atm,c}/2}$$

Deprojected Nuker profile (Lauer et al. 2007).

$$S_\rho = q_0 \left(\frac{r_b}{r} \right)^\gamma \left(1 + \left(\frac{r}{r_b} \right)^\alpha \right)^{(\gamma-\beta)/\alpha}$$

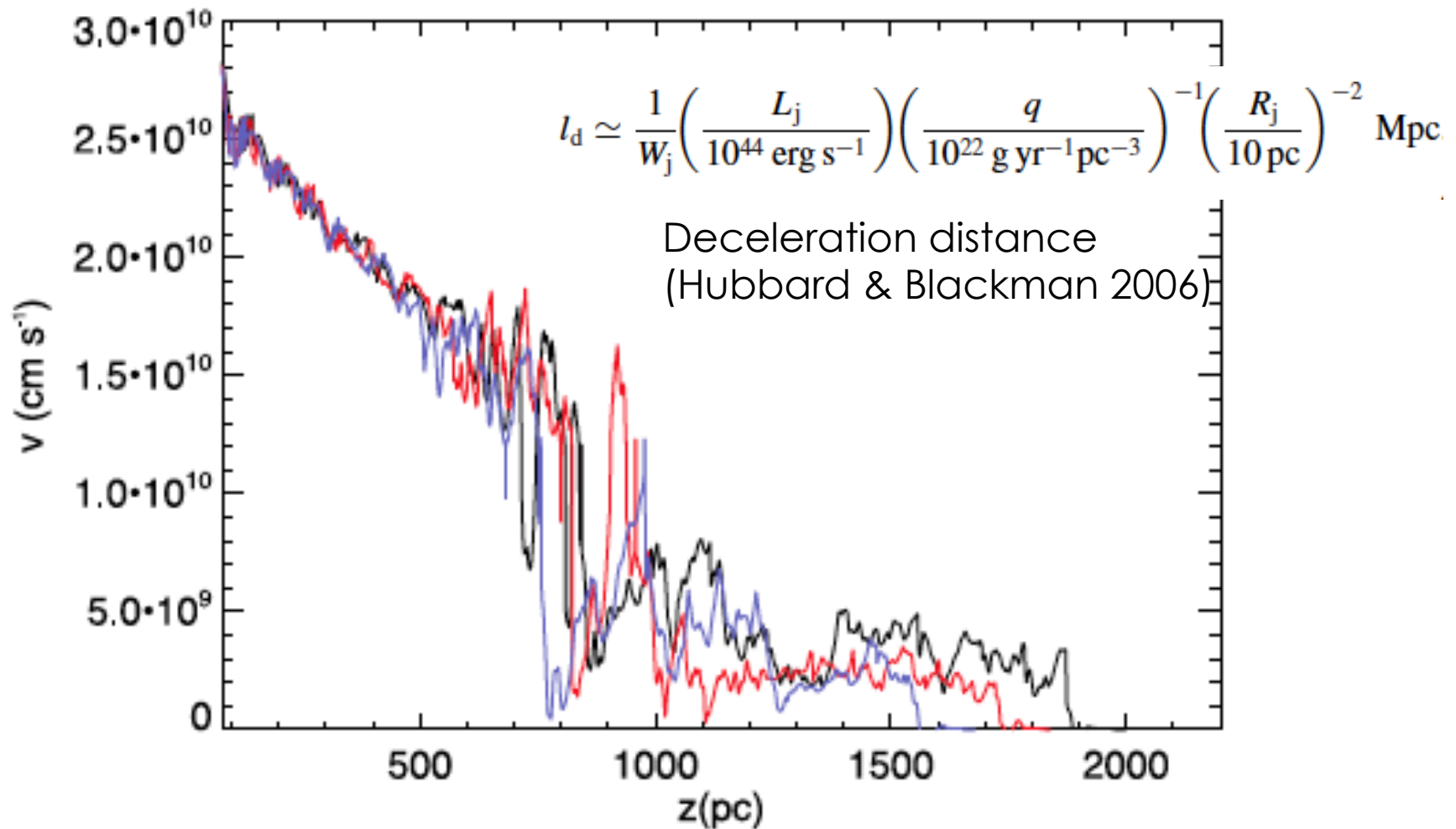
The stars are assumed to be all the same, with stellar mass losses 10^{-11} - $10^{-12} M_\odot \text{yr}^{-1}$.

Deceleration: mass load by stellar winds



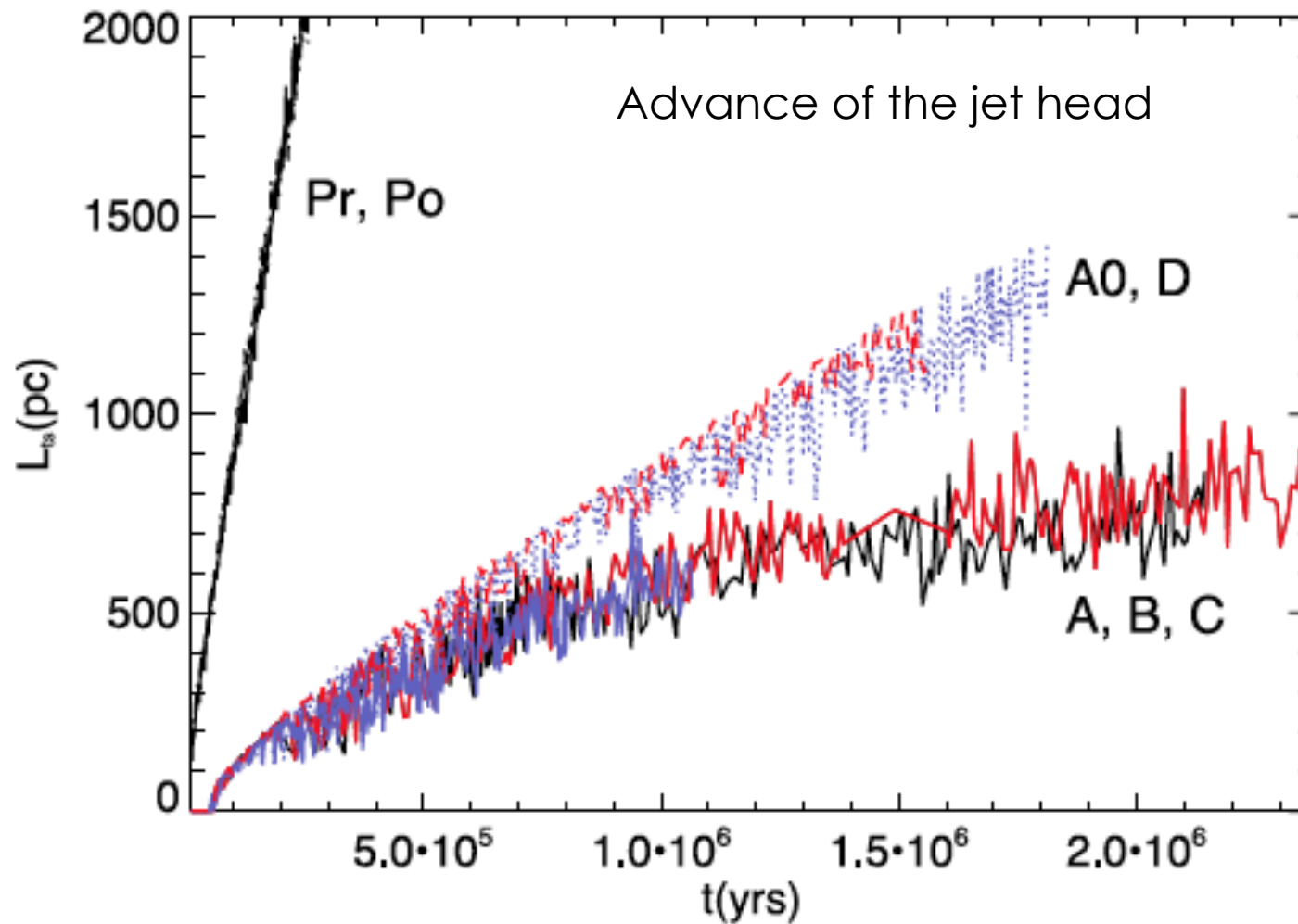
Jet velocity (less or no load)

Deceleration: mass load by stellar winds



Jet velocity (mass load)

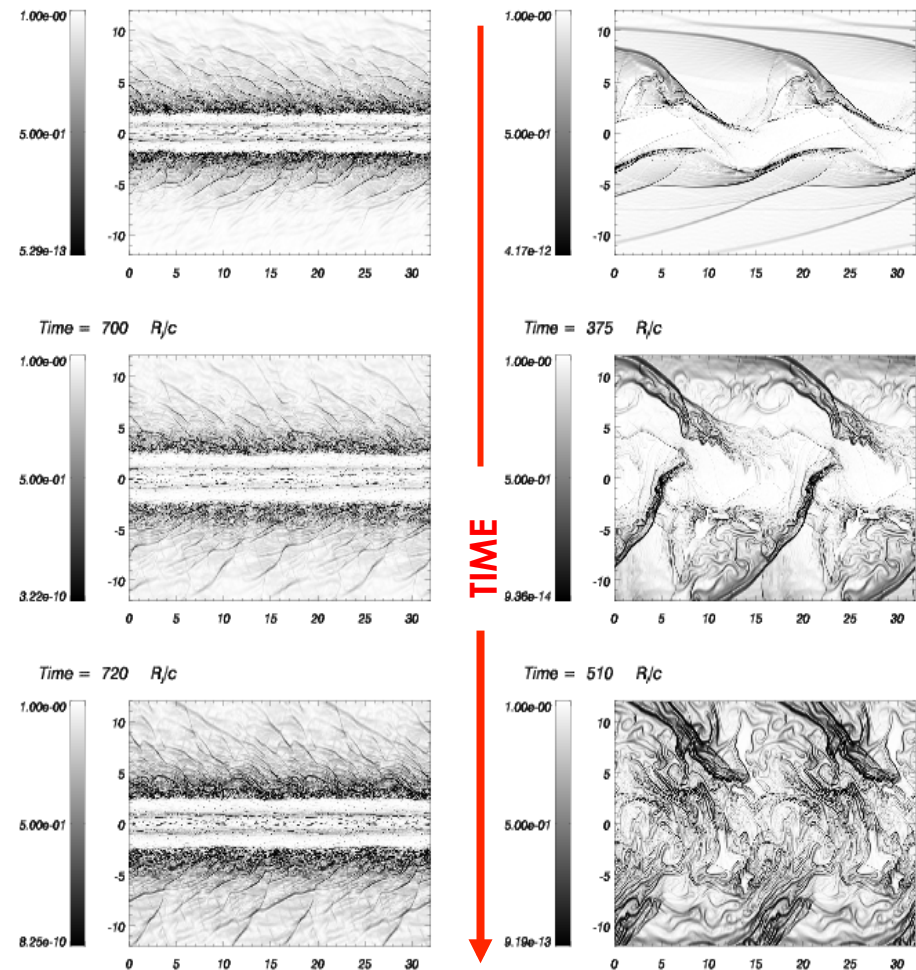
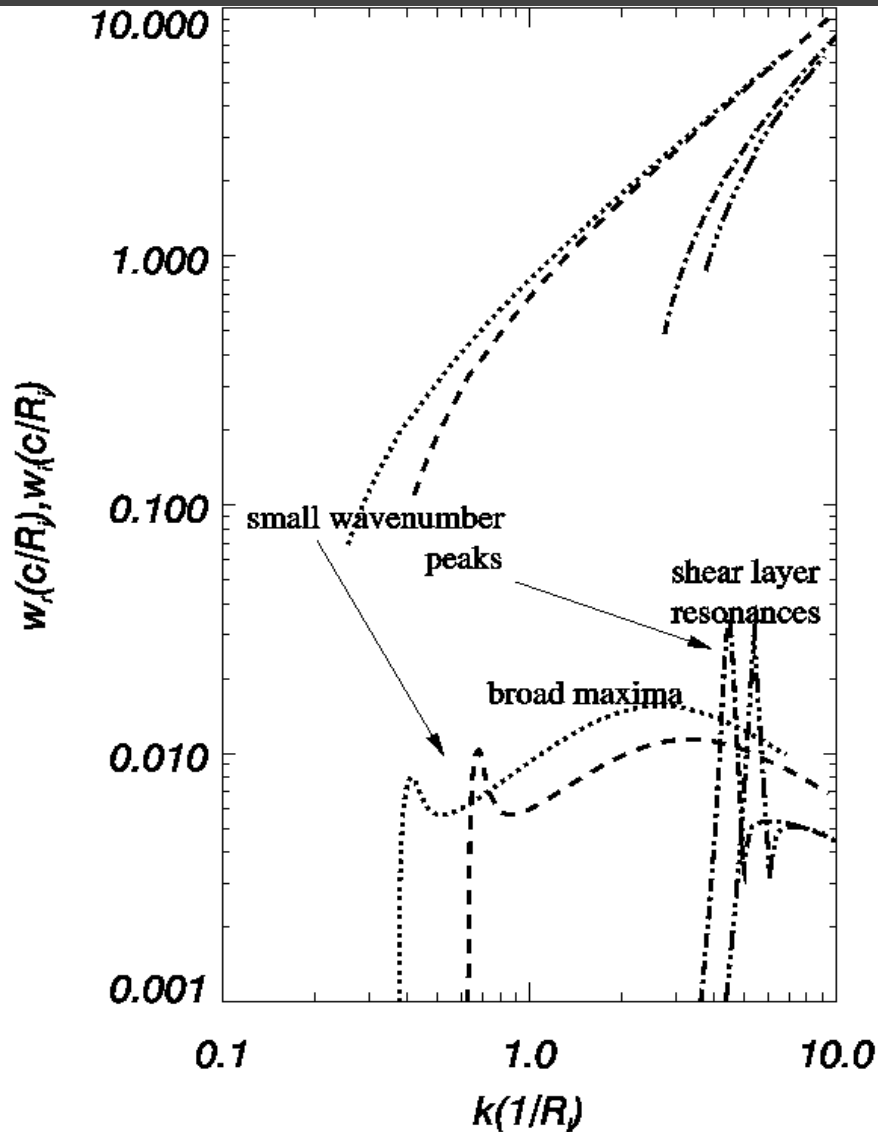
Deceleration: mass load by stellar winds



Conclusions

- Jet power and environmental properties (including stellar population) seem to determine the evolution of radio-sources.
- Jet deceleration:
 - Strong recollimation shocks in steep density gradients (not tremendously favoured by the observations).
 - Small scale instabilities in jets with $L_j \sim 10^{43-44}$ erg/s can cause mixing, dissipation, increase of internal energy and subsequent expansion within the expected FRI deceleration scales.
 - Long wavelength instabilities can cause the disruption of powerful jets in the long-term (possibly also the flaring of the inner spine in FRI's?).
 - Mass load by stellar winds and clouds can efficiently decelerate low-power AGN jets with $L_j \sim 10^{41-42}$ erg/s (Seyfert, LLAGN galaxies?), if the stellar population is old with weak stellar winds.

Deceleration: instabilities



Sheared jet ($d=0.2 R_j$)
Lorentz factor 20

Sheared jet ($d=0.2 R_j$)
Lorentz factor 5

Perucho et al. 2005, 2007

See also short λ saturation (Hardee 2011)