

Holographic Entanglement Entropy

Chandrasekhar Lecture 2
ICTS, December 11, 2014

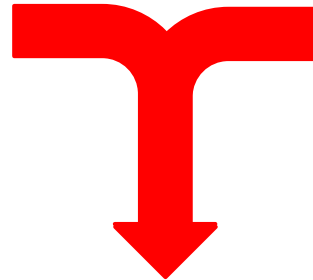
Entanglement Entropy

- condensed matter
- quantum information
- black hole microphysics

AdS/CFT correspondence

(gauge/gravity duality)

- string theory
- quantum gravity



Holographic Entanglement Entropy

- proposal by Ryu & Takayanagi (2006)

(New Horizons Prize, 2014)

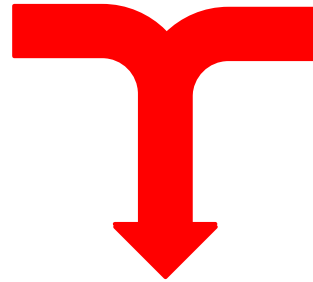
Entanglement Entropy

- condensed matter
- quantum information
- black hole microphysics

AdS/CFT correspondence

(gauge/gravity duality)

- string theory
- quantum gravity



Holographic Entanglement Entropy

- proposal by Ryu & Takayanagi (2006)

(New Horizons Prize, 2014)

Message:

- holographic entanglement entropy is part of a dialogue between string theory and a variety of fields (eg, condensed matter)
- offers potential to learn new lessons both about entanglement in quantum field theory and about quantum gravity

Overview:

1) Entanglement entropy in QFT

2) Primer on AdS/CFT correspondence

3) Introduction to Holographic Entanglement Entropy

4) Derivation of Holographic Entanglement Entropy
(for spherical entangling surfaces)

5) Conclusions

Quantum Entanglement

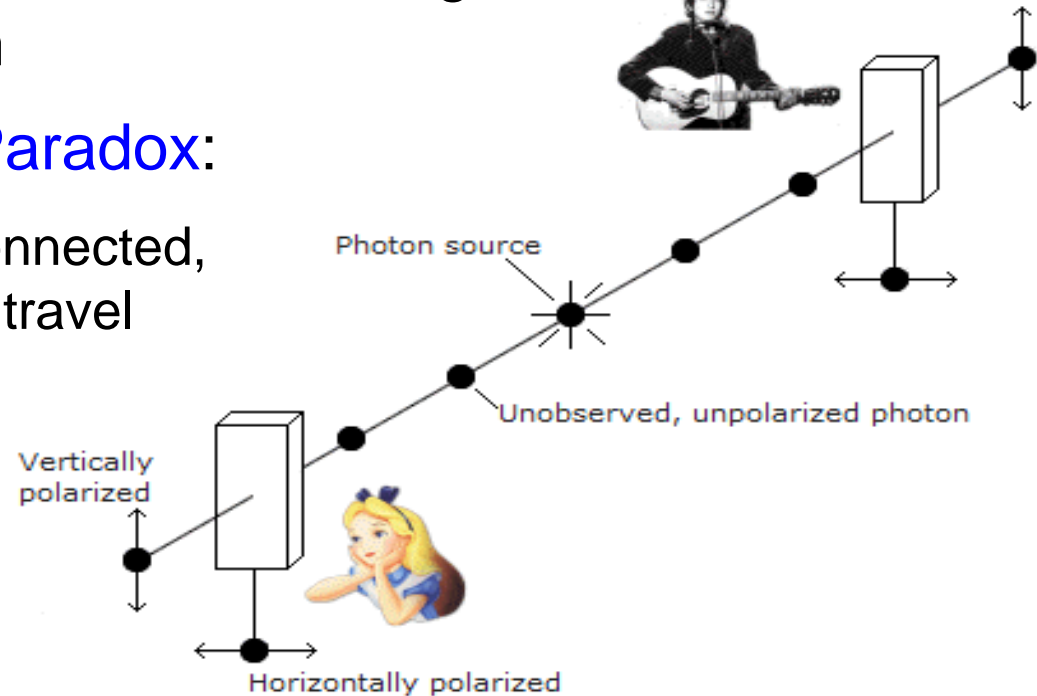
- different subsystems are correlated through global state of full system

Einstein-Podolsky-Rosen Paradox:

- properties of pair of photons connected, no matter how far apart they travel

“*spukhafte Fernwirkung*” = spooky action at a distance

$$|\psi\rangle = \frac{1}{\sqrt{2}} \left(|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle \right)$$



Quantum Information: entanglement becomes a resource for (ultra)fast computations and (ultra)secure communications

Condensed Matter: key to “exotic” phases and phenomena, e.g., quantum Hall fluids, unconventional superconductors, quantum spin fluids,

Entanglement Entropy

- general diagnostic: divide quantum system into two parts and use entropy as measure of correlations between subsystems
- procedure:
 - divide system into two subsystems, eg, A and B
 - integrate out degrees of freedom in subsystem B
 - remaining dof in A are described by a density matrix ρ_A
 - calculate **von Neumann entropy**: $S_{EE} = -Tr [\rho_A \log \rho_A]$
- simple examples found with spin systems:

$$\psi = \alpha |\uparrow\downarrow\rangle + \beta |\downarrow\uparrow\rangle \quad \text{with} \quad |\alpha|^2 + |\beta|^2 = 1$$

- “trace” over 2nd spin: $\rho_A = |\alpha|^2 |\uparrow\rangle\langle\uparrow| + |\beta|^2 |\downarrow\rangle\langle\downarrow|$

$$\longrightarrow S_{EE} = -|\alpha|^2 \log |\alpha|^2 - |\beta|^2 \log |\beta|^2$$

Entanglement Entropy

- general diagnostic: divide quantum system into two parts and use entropy as measure of correlations between subsystems
- procedure:
 - divide system into two subsystems, eg, A and B
 - integrate out degrees of freedom in subsystem B
 - remaining dof in A are described by a density matrix ρ_A
 - calculate **von Neumann entropy**: $S_{EE} = -Tr [\rho_A \log \rho_A]$
- simple examples found with spin systems:

$$\psi = \alpha |\uparrow\downarrow\rangle + \beta |\downarrow\downarrow\rangle \quad \text{with} \quad |\alpha|^2 + |\beta|^2 = 1$$

- “trace” over 2nd spin: $\rho_A = |\alpha|^2 |\uparrow\rangle\langle\uparrow| + |\beta|^2 |\downarrow\rangle\langle\downarrow| + \alpha\beta^\dagger |\uparrow\rangle\langle\downarrow| + \beta\alpha^\dagger |\downarrow\rangle\langle\uparrow|$

$$\longrightarrow S_{EE} = - \sum \lambda_n \log \lambda_n = 0 \quad (\lambda_1 = 0, \lambda_2 = 1)$$

Entanglement Entropy

- general diagnostic: divide quantum system into two parts and use entropy as measure of correlations between subsystems

- procedure:

- divide system into two subsystems, eg, A and B
- integrate out degrees of freedom in subsystem B
- remaining dof in A are described by a density matrix ρ_A
- calculate **von Neumann entropy**: $S_{EE} = -Tr [\rho_A \log \rho_A]$

- simple examples found with spin systems:

$$\psi = \alpha |\uparrow\downarrow\rangle + \beta |\downarrow\downarrow\rangle = (\alpha |\uparrow\rangle + \beta |\downarrow\rangle) \otimes |\downarrow\rangle$$

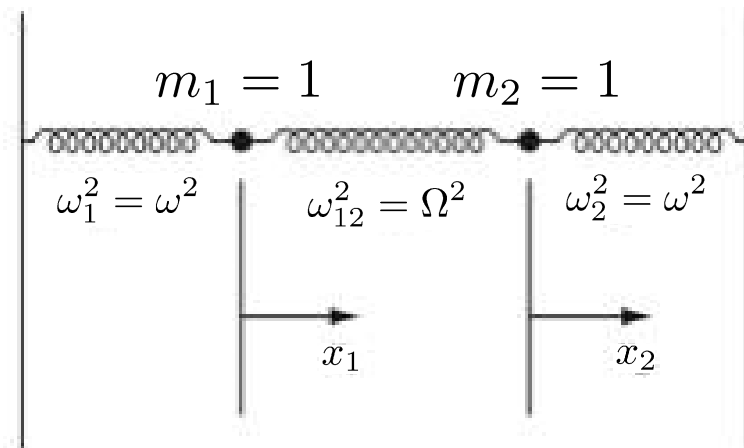
- “trace” over 2nd spin: $\rho_A = |\alpha|^2 |\uparrow\rangle\langle\uparrow| + |\beta|^2 |\downarrow\rangle\langle\downarrow| + \alpha\beta^\dagger |\uparrow\rangle\langle\downarrow| + \beta\alpha^\dagger |\downarrow\rangle\langle\uparrow|$

$$\longrightarrow S_{EE} = - \sum \lambda_n \log \lambda_n = 0 \quad (\lambda_1 = 0, \lambda_2 = 1)$$

Entanglement Entropy

- another example: two **coupled** harmonic oscillators

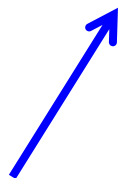
$$H = \frac{1}{2} \left[p_1^2 + p_2^2 + \omega^2 (x_1^2 + x_2^2) + \Omega^2 (x_1 - x_2)^2 \right]$$



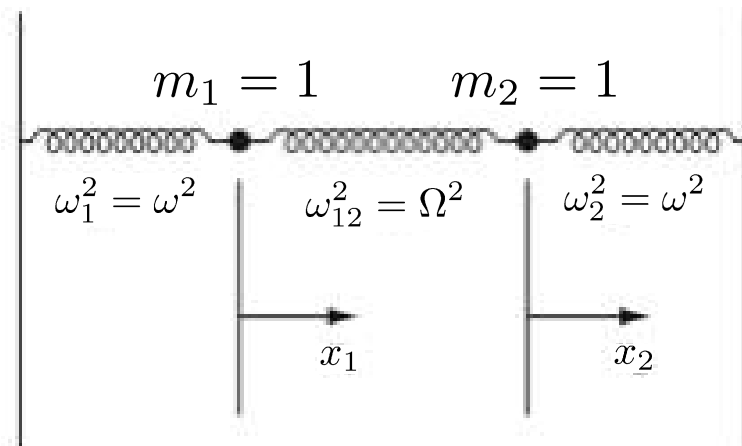
Entanglement Entropy

- another example: two **coupled** harmonic oscillators

$$\begin{aligned}
 H &= \frac{1}{2} \left[p_1^2 + p_2^2 + \omega^2 (x_1^2 + x_2^2) + \Omega^2 (x_1 - x_2)^2 \right] \\
 &= \frac{1}{2} \left[p_+^2 + \omega_+^2 x_+^2 + p_-^2 + \omega_-^2 x_-^2 \right] \quad \text{with } x_{\pm} = \frac{1}{\sqrt{2}} (x_1 \pm x_2), \\
 &\hspace{15em} \omega_+^2 = \omega^2, \quad \omega_-^2 = \omega^2 + 2\Omega^2
 \end{aligned}$$



find normal modes; problem reduces to two independent SHO's



Entanglement Entropy

- another example: two coupled harmonic oscillators

$$\begin{aligned}
 H &= \frac{1}{2} \left[p_1^2 + p_2^2 + \omega^2 (x_1^2 + x_2^2) + \Omega^2 (x_1 - x_2)^2 \right] \\
 &= \frac{1}{2} \left[p_+^2 + \omega_+^2 x_+^2 + p_-^2 + \omega_-^2 x_-^2 \right] \quad \text{with } x_{\pm} = \frac{1}{\sqrt{2}} (x_1 \pm x_2), \\
 &\qquad\qquad\qquad \omega_+^2 = \omega^2, \quad \omega_-^2 = \omega^2 + 2\Omega^2
 \end{aligned}$$

- ground-state wave-function:

$$\Psi_0(x_+, x_-) = \Psi_0(x_+) \Psi_0(x_-) = \frac{1}{\sqrt{\pi}(\omega_+\omega_-)^{1/4}} \exp \left[-\frac{1}{4} (\omega_+ x_+^2 + \omega_- x_-^2) \right]$$

$$\Psi_0(x_1, x_2) = \frac{1}{\sqrt{\pi}(\omega_+\omega_-)^{1/4}} \exp \left[-\frac{1}{8} (\omega_+ (x_1 + x_2)^2 + \omega_- (x_1 - x_2)^2) \right]$$

Entanglement Entropy

- another example: two coupled harmonic oscillators

$$\begin{aligned}
 H &= \frac{1}{2} \left[p_1^2 + p_2^2 + \omega^2 (x_1^2 + x_2^2) + \Omega^2 (x_1 - x_2)^2 \right] \\
 &= \frac{1}{2} \left[p_+^2 + \omega_+^2 x_+^2 + p_-^2 + \omega_-^2 x_-^2 \right] \quad \text{with } x_{\pm} = \frac{1}{\sqrt{2}} (x_1 \pm x_2), \\
 &\qquad\qquad\qquad \omega_+^2 = \omega^2, \quad \omega_-^2 = \omega^2 + 2\Omega^2
 \end{aligned}$$

- ground-state wave-function:

$$\Psi_0(x_1, x_2) = \frac{1}{\sqrt{\pi}(\omega_+ \omega_-)^{1/4}} \exp \left[-\frac{1}{8} (\omega_+ (x_1 + x_2)^2 + \omega_- (x_1 - x_2)^2) \right]$$

- trace over position of oscillator 1:

$$\begin{aligned}
 \rho(x_2, x'_2) &= \int_{-\infty}^{+\infty} dx_1 \Psi_0(x_1, x_2) \Psi_0^\dagger(x_1, x'_2) \\
 &= \frac{(\gamma - \beta)^{1/2}}{\sqrt{\pi}} \exp \left[-\frac{\gamma}{2} (x_2^2 + x_2'^2) + \beta x_2 x_2' \right]
 \end{aligned}$$

$$\text{with } \gamma = \frac{1}{4}(\omega_+ + \omega_-) + \frac{\omega_+ \omega_-}{\omega_+ + \omega_-} \quad \text{and} \quad \beta = \frac{1}{4}(\omega_+ + \omega_-) - \frac{\omega_+ \omega_-}{\omega_+ + \omega_-}$$

Entanglement Entropy

- another example: two coupled harmonic oscillators

$$\begin{aligned}
 H &= \frac{1}{2} \left[p_1^2 + p_2^2 + \omega^2 (x_1^2 + x_2^2) + \Omega^2 (x_1 - x_2)^2 \right] \\
 &= \frac{1}{2} \left[p_+^2 + \omega_+^2 x_+^2 + p_-^2 + \omega_-^2 x_-^2 \right] \quad \text{with } x_{\pm} = \frac{1}{\sqrt{2}} (x_1 \pm x_2),
 \end{aligned}$$

- ground-state wave-function: $\omega_+^2 = \omega^2, \quad \omega_-^2 = \omega^2 + 2\Omega^2$

$$\Psi_0(x_1, x_2) = \frac{1}{\sqrt{\pi}(\omega_+ \omega_-)^{1/4}} \exp \left[-\frac{1}{8} (\omega_+ (x_1 + x_2)^2 + \omega_- (x_1 - x_2)^2) \right]$$

- trace over position of oscillator 1:

$$\rho(x_2, x'_2) = \int_{-\infty}^{+\infty} dx_1 \Psi_0(x_1, x_2) \Psi_0^\dagger(x_1, x'_2)$$

$$= \frac{(\gamma - \beta)^{1/2}}{\sqrt{\pi}} \exp \left[-\frac{\gamma}{2} (x_2^2 + x_2'^2) + \beta x_2 x_2' \right]$$

$$\text{with } \gamma = \frac{1}{4}(\omega_+ + \omega_-) + \frac{\omega_+ \omega_-}{\omega_+ + \omega_-} \quad \text{and} \quad \beta = \frac{1}{4}(\omega_+ + \omega_-) - \frac{\omega_+ \omega_-}{\omega_+ + \omega_-}$$

$$\underline{\Omega/\omega \ll 1} \quad \sim \omega$$

$$\sim \Omega^4/8\omega^3$$

Entanglement Entropy

- another example: two coupled harmonic oscillators

$$\begin{aligned}
 H &= \frac{1}{2} \left[p_1^2 + p_2^2 + \omega^2 (x_1^2 + x_2^2) + \Omega^2 (x_1 - x_2)^2 \right] \\
 &= \frac{1}{2} \left[p_+^2 + \omega_+^2 x_+^2 + p_-^2 + \omega_-^2 x_-^2 \right] \quad \text{with } x_{\pm} = \frac{1}{\sqrt{2}} (x_1 \pm x_2),
 \end{aligned}$$

- ground-state wave-function: $\omega_+^2 = \omega^2, \quad \omega_-^2 = \omega^2 + 2\Omega^2$

$$\Psi_0(x_1, x_2) = \frac{1}{\sqrt{\pi}(\omega_+ \omega_-)^{1/4}} \exp \left[-\frac{1}{8} (\omega_+ (x_1 + x_2)^2 + \omega_- (x_1 - x_2)^2) \right]$$

- trace over position of oscillator 1:

$$\begin{aligned}
 \rho(x_2, x'_2) &= \int_{-\infty}^{+\infty} dx_1 \Psi_0(x_1, x_2) \Psi_0^\dagger(x_1, x'_2) \\
 &= \frac{(\gamma - \beta)^{1/2}}{\sqrt{\pi}} \exp \left[-\frac{\gamma}{2} (x_2^2 + x_2'^2) + \beta x_2 x_2' \right]
 \end{aligned}$$

entanglement
↙

$$\text{with } \gamma = \frac{1}{4}(\omega_+ + \omega_-) + \frac{\omega_+ \omega_-}{\omega_+ + \omega_-} \quad \text{and} \quad \beta = \frac{1}{4}(\omega_+ + \omega_-) - \frac{\omega_+ \omega_-}{\omega_+ + \omega_-}$$

$$\underline{\Omega/\omega \ll 1} \quad \sim \omega$$

$$\sim \Omega^4 / 8\omega^3$$

Entanglement Entropy

- another example: two coupled harmonic oscillators

$$\rho(x_2, x'_2) = \frac{(\gamma - \beta)^{1/2}}{\sqrt{\pi}} \exp \left[-\frac{\gamma}{2} (x_2^2 + x'_2{}^2) + \beta x_2 x'_2 \right]$$

- entanglement entropy: $S_{EE} = -\text{Tr} [\rho \log \rho]$?

Entanglement Entropy

- another example: two coupled harmonic oscillators

$$\rho(x_2, x'_2) = \frac{(\gamma - \beta)^{1/2}}{\sqrt{\pi}} \exp \left[-\frac{\gamma}{2} (x_2^2 + x'_2{}^2) + \beta x_2 x'_2 \right]$$

- entanglement entropy: $S_{EE} = -Tr [\rho \log \rho]$?

$$\longrightarrow S_{EE} = - \sum \lambda_n \log \lambda_n \quad \text{where} \quad \int_{-\infty}^{+\infty} dx' \rho(x, x') f_n(x') = \lambda_n f_n(x)$$

- find: $\lambda_n = (1 - \xi) \xi^n$ with $f_n(x) = H_n(\alpha^{1/2} x) \exp [-\alpha x^2/2]$,

$$\xi = \beta/(\gamma + \alpha) \quad \text{and} \quad \alpha = (\omega_+ \omega_-)^{1/2}$$

$$\longrightarrow S_{EE} = -\log(1 - \xi) - \frac{\xi}{1 - \xi} \log \xi$$

Entanglement Entropy

- another example: two coupled harmonic oscillators

$$\rho(x_2, x'_2) = \frac{(\gamma - \beta)^{1/2}}{\sqrt{\pi}} \exp \left[-\frac{\gamma}{2} (x_2^2 + x'_2{}^2) + \beta x_2 x'_2 \right]$$

- entanglement entropy: $S_{EE} = -\text{Tr} [\rho \log \rho]$?

$$\longrightarrow S_{EE} = -\sum \lambda_n \log \lambda_n \quad \text{where} \quad \int_{-\infty}^{+\infty} dx' \rho(x, x') f_n(x') = \lambda_n f_n(x)$$

- find: $\lambda_n = (1 - \xi) \xi^n$ with $f_n(x) = H_n(\alpha^{1/2} x) \exp[-\alpha x^2/2]$,

$$\xi = \beta/(\gamma + \alpha) \quad \text{and} \quad \alpha = (\omega_+ \omega_-)^{1/2}$$


$$\longrightarrow S_{EE} = -\log(1 - \xi) - \frac{\xi}{1 - \xi} \log \xi$$

- notice: as $\Omega/\omega \rightarrow 0$, $\xi \sim (\Omega/2\omega)^4 \rightarrow 0$ and $S_{EE} \sim \frac{1}{4} \left(\frac{\Omega}{\omega} \right)^4 \log(2\omega/\Omega)$
 $\rightarrow 0$

Entanglement Entropy

- another example: two coupled harmonic oscillators

$$\rho(x_2, x'_2) = \frac{(\gamma - \beta)^{1/2}}{\sqrt{\pi}} \exp \left[-\frac{\gamma}{2} (x_2^2 + x_2'^2) + \beta x_2 x_2' \right]$$

- entanglement entropy: $S_{EE} = -\log(1 - \xi) - \frac{\xi}{1 - \xi} \log \xi$
- note entanglement arises because I insisted on working with original basis x_1, x_2  guided by physical narrative
- no entanglement in normal mode basis x_+, x_-

Entanglement Entropy

- another example: two coupled harmonic oscillators

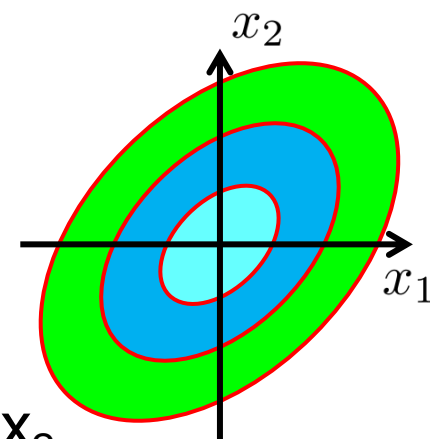
$$\rho(x_2, x'_2) = \frac{(\gamma - \beta)^{1/2}}{\sqrt{\pi}} \exp \left[-\frac{\gamma}{2} (x_2^2 + x'_2{}^2) + \beta x_2 x'_2 \right]$$

- entanglement entropy: $S_{EE} = -\log(1 - \xi) - \frac{\xi}{1 - \xi} \log \xi$
- note entanglement arises because I insisted on working with original basis x_1, x_2 \longrightarrow **guided by physical narrative**
- no entanglement in normal mode basis x_+, x_-
- can interpret eq's as two-dimensional SHO

$$V(x_1, x_2) = \frac{1}{2}\omega^2 (x_1^2 + x_2^2) + \frac{1}{2}\Omega^2 (x_1 - x_2)^2$$

\longrightarrow x_1 and x_2 take less precedence

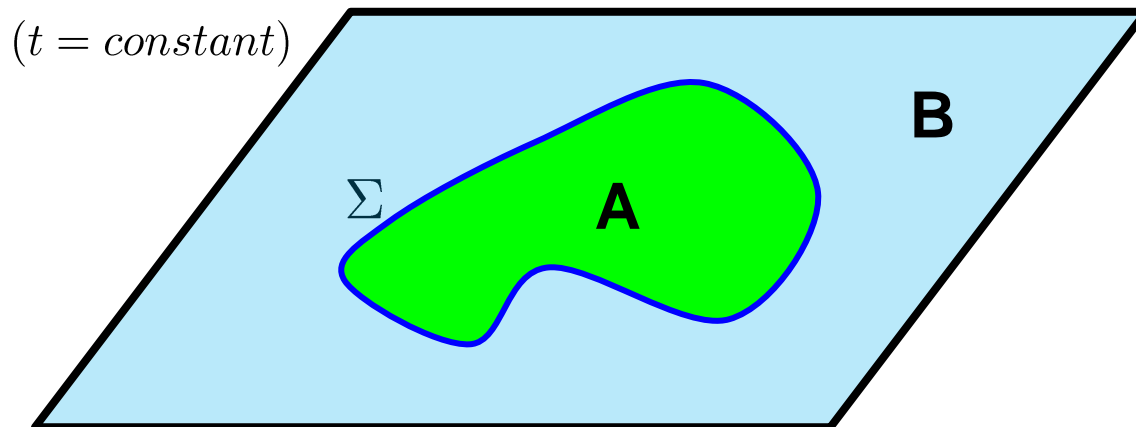
\longrightarrow $S_{EE} \neq 0$ integrating out any axis except $x_1 = \pm x_2$



Entanglement Entropy 2:

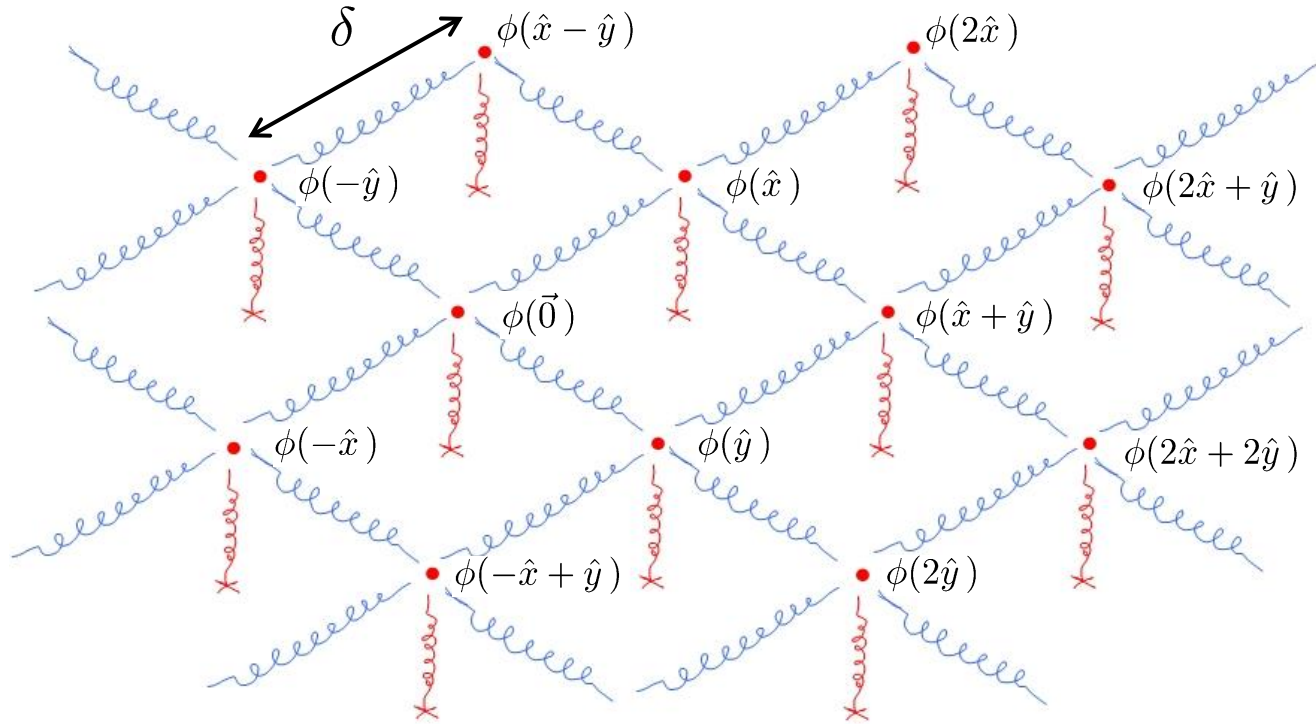
Quantum Field Theory??

- in the context of holographic entanglement entropy, S_{EE} is applied in the context of **quantum field theory**
 - in **QFT**, typically introduce a (smooth) boundary **or entangling surface** Σ which divides the space into two separate regions
 - integrate out degrees of freedom in “outside” region
 - remaining dof are described by a density matrix ρ_A
- calculate **von Neumann entropy** $S_{EE} = -Tr [\rho_A \log \rho_A]$



- free scalar field (in d dimensions) = many coupled oscillators

$$H = \frac{1}{2} \sum_{\vec{n}} \left[\frac{p(\vec{n})^2}{\delta^{d-1}} + \delta^{d-1} \left(\frac{1}{\delta^2} \sum_i (\phi(\vec{n}) - \phi(\vec{n} - \hat{x}_i))^2 + \mu^2 \phi(\vec{n})^2 \right) \right]$$

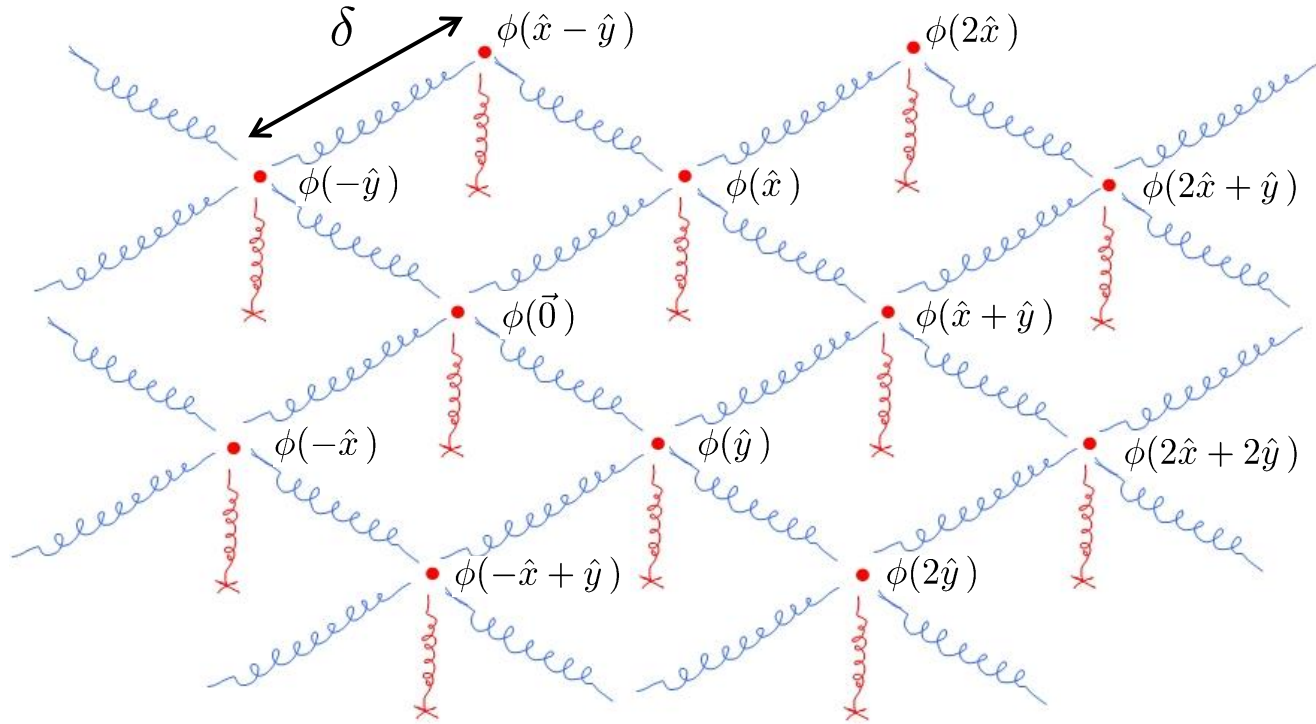


with $\omega^2 = \mu^2$, $\Omega^2 = 1/\delta^2$ and reinstated $m = \delta^{d-1}$.

→ physical narrative: have a lattice of oscillators with
 $\delta =$ lattice spacing = short distance cut-off

- free scalar field (in d dimensions) = many coupled oscillators

$$H = \frac{1}{2} \sum_{\vec{n}} \left[\frac{p(\vec{n})^2}{\delta^{d-1}} + \delta^{d-1} \left(\frac{1}{\delta^2} \sum_i (\phi(\vec{n}) - \phi(\vec{n} - \hat{x}_i))^2 + \mu^2 \phi(\vec{n})^2 \right) \right]$$



- consider continuum limit: $\delta \rightarrow 0$


$$\phi(\vec{n}) \rightarrow \phi(\vec{x}), \quad p(\vec{n})/\delta^{d-1} \rightarrow \pi(\vec{x}), \quad \delta^{d-1} \sum_{\vec{n}} \rightarrow \int d^{d-1}x$$

$$H = \frac{1}{2} \int d^{d-1}x \left[|\pi(x)|^2 + |\vec{\nabla} \phi(x)|^2 + \mu^2 |\phi(x)|^2 \right]$$

- free scalar field (in d dimensions) = many coupled oscillators

$$H = \frac{1}{2} \int d^{d-1}x \left[|\pi(x)|^2 + |\vec{\nabla}\phi(x)|^2 + \mu^2|\phi(x)|^2 \right]$$


$$= \frac{1}{2} \int d^{d-1}k \left[|\pi(k)|^2 + k^2|\phi(k)|^2 + \mu^2|\phi(k)|^2 \right]$$


$$\phi(\vec{k}) = \int d^{d-1}x e^{i\vec{k}\cdot\vec{x}} \phi(x)$$

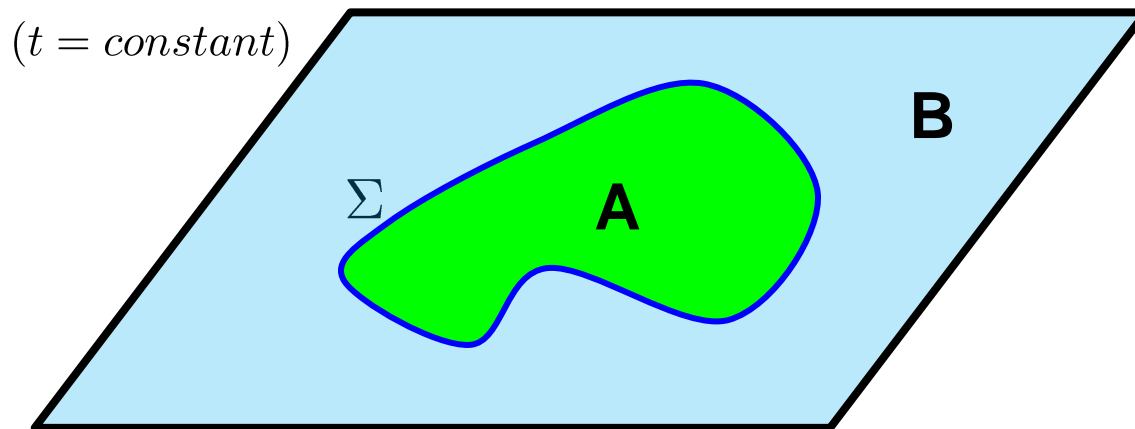
- normal modes given by Fourier transform
- much of “textbook QFT” is perturbing lots of coupled SHO’s

- free scalar field (in d dimensions) = many coupled oscillators

$$\begin{aligned}
 H &= \frac{1}{2} \int d^{d-1}x \left[|\pi(x)|^2 + |\vec{\nabla}\phi(x)|^2 + \mu^2|\phi(x)|^2 \right] \\
 &= \frac{1}{2} \int d^{d-1}k \left[|\pi(k)|^2 + k^2|\phi(k)|^2 + \mu^2|\phi(k)|^2 \right]
 \end{aligned}$$


 $\phi(\vec{k}) = \int d^{d-1}x e^{i\vec{k}\cdot\vec{x}} \phi(x)$


- normal modes given by Fourier transform
- much of “textbook QFT” is perturbing lots of coupled SHO’s



- proposed entanglement calculation done in position basis

- free scalar field (in d dimensions) = many coupled oscillators

$$\begin{aligned}
 H &= \frac{1}{2} \int d^{d-1}x \left[|\pi(x)|^2 + |\vec{\nabla}\phi(x)|^2 + \mu^2|\phi(x)|^2 \right] \\
 &= \frac{1}{2} \int d^{d-1}k \left[|\pi(k)|^2 + k^2|\phi(k)|^2 + \mu^2|\phi(k)|^2 \right]
 \end{aligned}$$



$$\phi(\vec{k}) = \int d^{d-1}x e^{i\vec{k}\cdot\vec{x}} \phi(x)$$

- normal modes given by Fourier transform
- much of “textbook QFT” is perturbing lots of coupled SHO’s
- proposed entanglement calculation done in position basis


Ground state wave-functional:

$$\Psi_0(\phi(k)) \sim \exp \left[-\frac{1}{2} \int \frac{d^{d-1}k}{(2\pi)^{d-1}} \sqrt{k^2 + \mu^2} \phi^\dagger(k) \phi(k) \right]$$

- simple “product” in terms of normal modes

- free scalar field (in d dimensions) = many coupled oscillators

$$\begin{aligned}
 H &= \frac{1}{2} \int d^{d-1}x \left[|\pi(x)|^2 + |\vec{\nabla}\phi(x)|^2 + \mu^2|\phi(x)|^2 \right] \\
 &= \frac{1}{2} \int d^{d-1}k \left[|\pi(k)|^2 + k^2|\phi(k)|^2 + \mu^2|\phi(k)|^2 \right]
 \end{aligned}$$



$$\phi(\vec{k}) = \int d^{d-1}x e^{i\vec{k}\cdot\vec{x}} \phi(x)$$

- normal modes given by Fourier transform
- much of “textbook QFT” is perturbing lots of coupled SHO’s
- proposed entanglement calculation done in position basis

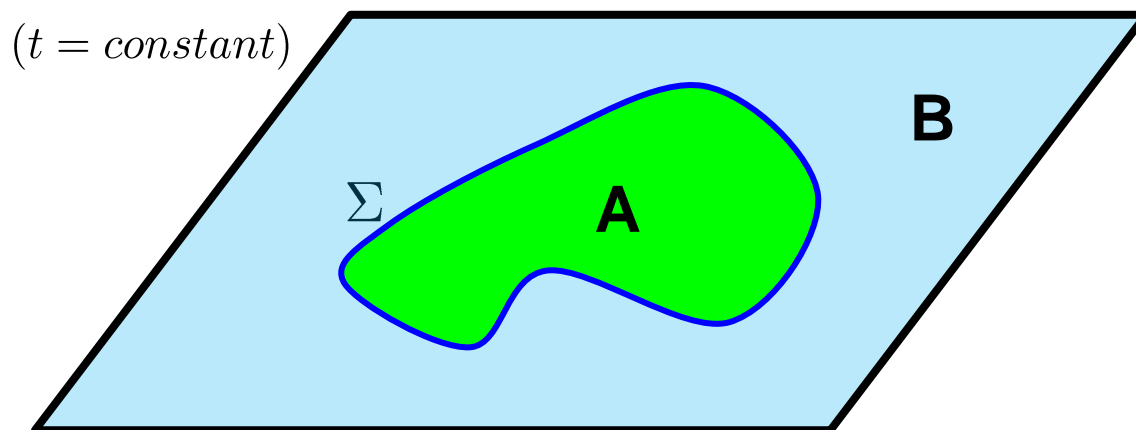
Ground state wave-functional:

$$\Psi_0(\phi(x)) \sim \exp \left[-\frac{1}{2} \int d^{d-1}x_1 \int d^{d-1}x_2 \phi(x_1) W(x_1, x_2) \phi(x_2) \right]$$

with $W(x_1, x_2) = \int \frac{d^{d-1}k}{(2\pi)^{d-1}} \sqrt{k^2 + \mu^2} e^{-i\vec{k}\cdot(\vec{x}_1 - \vec{x}_2)}$

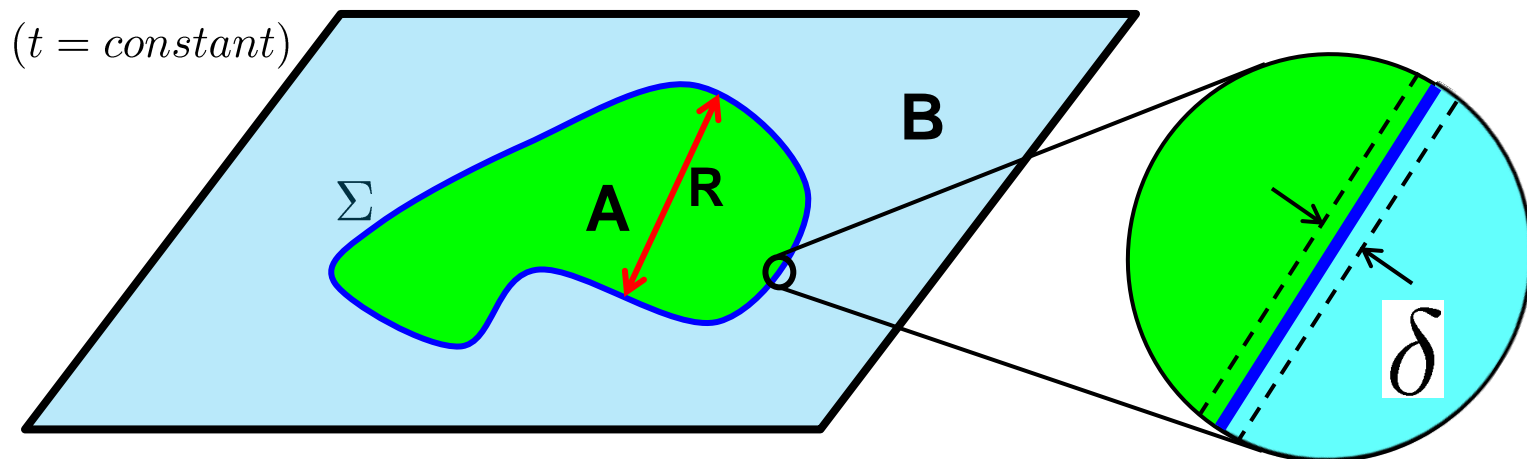
Entanglement Entropy 2:

- in the context of holographic entanglement entropy, S_{EE} is applied in the context of **quantum field theory**
 - in **QFT**, typically introduce a (smooth) boundary **or entangling surface** Σ which divides the space into two separate regions
 - integrate out degrees of freedom in “outside” region
 - remaining dof are described by a density matrix ρ_A
- calculate **von Neumann entropy** $S_{EE} = -Tr [\rho_A \log \rho_A]$



Entanglement Entropy 2:

→ calculate von Neumann entropy: $S_{EE} = -Tr [\rho_A \log \rho_A]$



- result is **UV divergent!** dominated by short-distance correlations
- must regulate calculation: $\delta = \text{short-distance cut-off}$

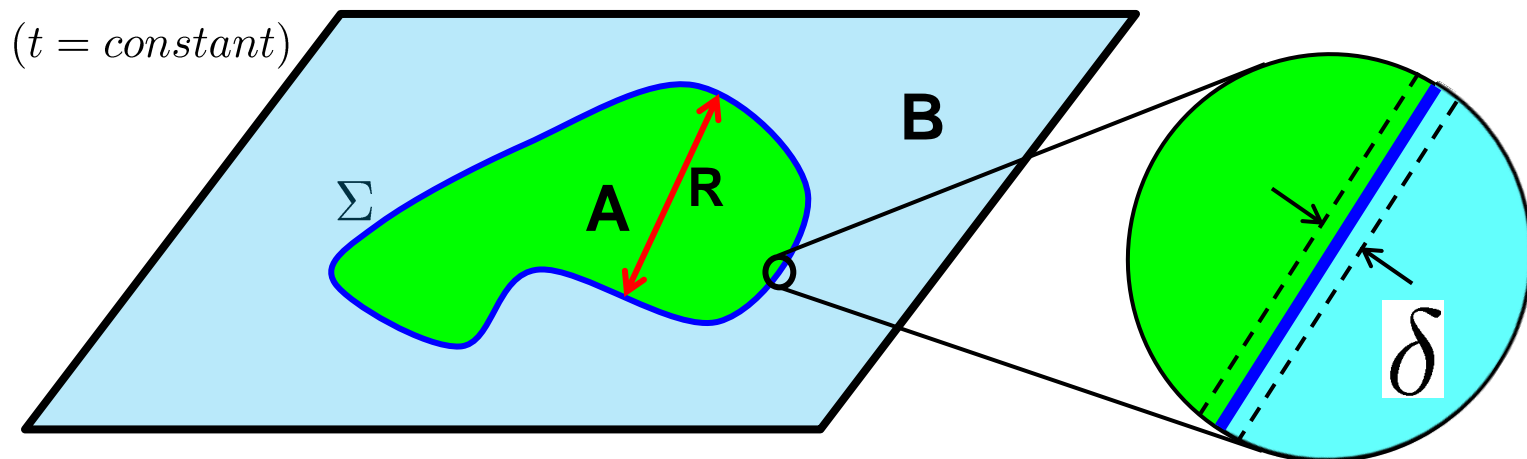
$$S = c_0 \frac{R^{d-2}}{\delta^{d-2}} + c_2 \frac{R^{d-4}}{\delta^{d-4}} + \dots \quad d = \text{spacetime dimension}$$

- careful analysis reveals geometric structure, eg,

$$S = \tilde{c}_0 \frac{\mathcal{A}_\Sigma}{\delta^{d-2}} + \tilde{c}_2 \frac{\int_\Sigma d^{d-2} \sigma \text{ "R" }}{\delta^{d-4}} + \dots$$

Entanglement Entropy 2:

→ calculate von Neumann entropy: $S_{EE} = -\text{Tr} [\rho_A \log \rho_A]$



- result is **UV divergent!** dominated by short-distance correlations
- must regulate calculation: $\delta = \text{short-distance cut-off}$

$$S = c_0 \frac{R^{d-2}}{\delta^{d-2}} + c_2 \frac{R^{d-4}}{\delta^{d-4}} + \dots \quad d = \text{spacetime dimension}$$

- leading coefficients sensitive to details of regulator, eg, $\delta \rightarrow 2\delta$
- find **universal** information characterizing underlying QFT in subleading terms, eg, $S = \dots + c_d \log(R/\delta) + \dots$

More general comments on **Entanglement Entropy**:

- nonlocal quantity which is (at best) very difficult to measure
→ no accepted experimental procedure
- still useful diagnostic in condensed matter theory and in quantum information theory

More general comments on **Entanglement Entropy**:

Where did [hep-th](#) find “Entanglement Entropy”? **→ black holes**

Bekenstein & Hawking: “quantum” black holes are thermal systems leaking blackbody radiation

→ event horizon carries entropy!!

$$S_{BH} = \mathcal{A}/4G_N$$



- **Sorkin '84**: leading term in EE obeys “area law”: $S = c_0 \frac{\mathcal{A}_\Sigma}{\delta^{d-2}} + \dots$
→ suggestive of Bekenstein-Hawking formula if $\delta \simeq \ell_P$

(Sorkin '84; Bombelli, Koul, Lee & Sorkin; Srednicki; Frolov & Novikov; . . .)

More general comments on **Entanglement Entropy**:

Where did [hep-th](#) find “Entanglement Entropy”? → **black holes**

Bekenstein & Hawking: “quantum” black holes are thermal systems leaking blackbody radiation

→ event horizon carries entropy!!

$$S_{BH} = \mathcal{A}/4G_N$$



- **Sorkin '84**: leading term in EE obeys “area law”: $S = c_0 \frac{\mathcal{A}_\Sigma}{\delta^{d-2}} + \dots$
→ suggestive of Bekenstein-Hawking formula if $\delta \simeq \ell_P$
(Sorkin '84; Bombelli, Koul, Lee & Sorkin; Srednicki; Frolov & Novikov; . . .)
- **issues?**: leading singularity not universal; regulator dependent
- active topic in 90's but story left unresolved ([returned to recently](#))

More general comments on **Entanglement Entropy**:

Where did **hep-th** find “Entanglement Entropy”? **→ black holes**

Bekenstein & Hawking: “quantum” black holes are thermal systems leaking blackbody radiation

→ event horizon carries entropy!!

$$S_{BH} = \mathcal{A}/4G_N$$



- **Sorkin '84**: leading term in EE obeys “area law”: $S = c_0 \frac{\mathcal{A}_\Sigma}{\delta^{d-2}} + \dots$

→ suggestive of Bekenstein-Hawking formula if $\delta \simeq \ell_P$

(Sorkin `84; Bombelli, Koul, Lee & Sorkin; Srednicki; Frolov & Novikov; . . .)

- **issues?**: leading singularity not universal; regulator dependent
- active topic in 90's but story left unresolved (**returned to recently**)



- recently considered in **AdS/CFT correspondence**

(Ryu & Takayanagi `06)

Overview:

1) Entanglement entropy in QFT

2) Primer on AdS/CFT correspondence

3) Introduction to Holographic Entanglement Entropy

4) Derivation of Holographic Entanglement Entropy
(for spherical entangling surfaces) and Beyond

5) Conclusions

AdS/CFT Correspondence:

Bulk:

- quantum gravity
- negative cosmological constant
- **d+1** spacetime dimensions

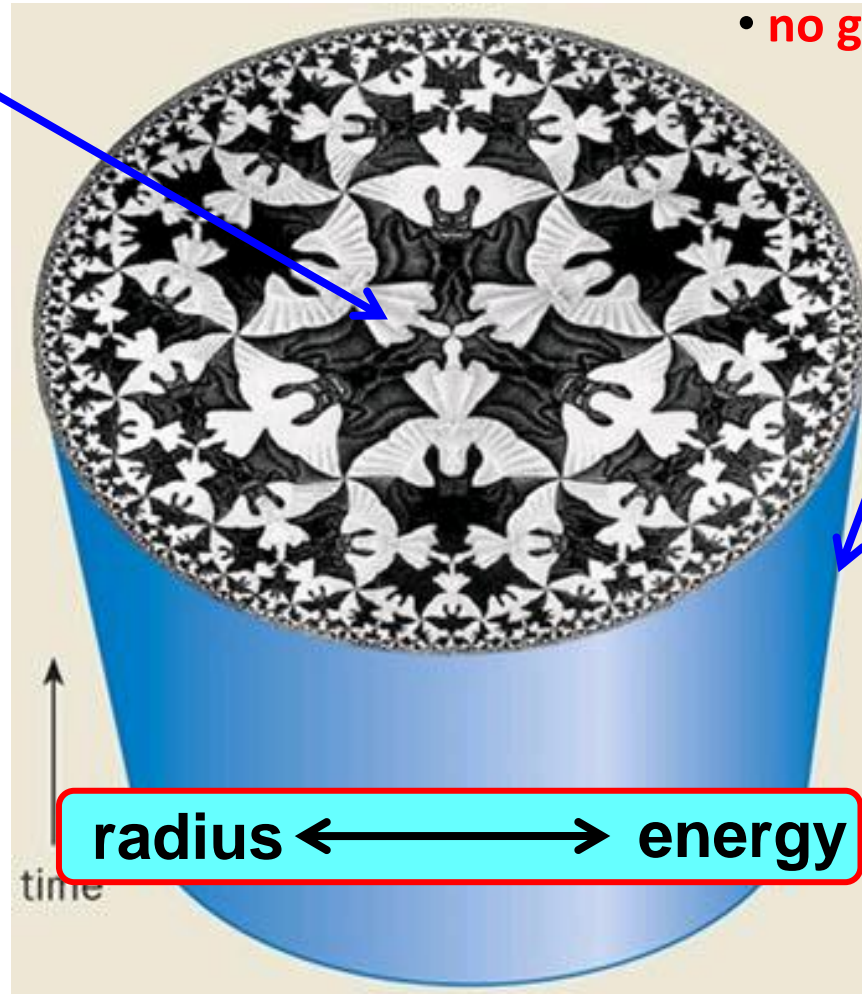
Boundary:

- quantum field theory
- no scale (at quantum level)
- **d** spacetime dimensions
- **no gravity!**

Holography



anti-de Sitter
space



conformal
field theory

radius \longleftrightarrow energy

(Maldacena '97)

AdS/CFT correspondence:

anti-de Sitter space

quantum gravity

- negative cosmological constant
- d+1 spacetime dimensions



conformal field theory

quantum field theory

- no scale (at quantum level)
- d spacetime dimensions
- **no gravity!**



holography



classical gravity

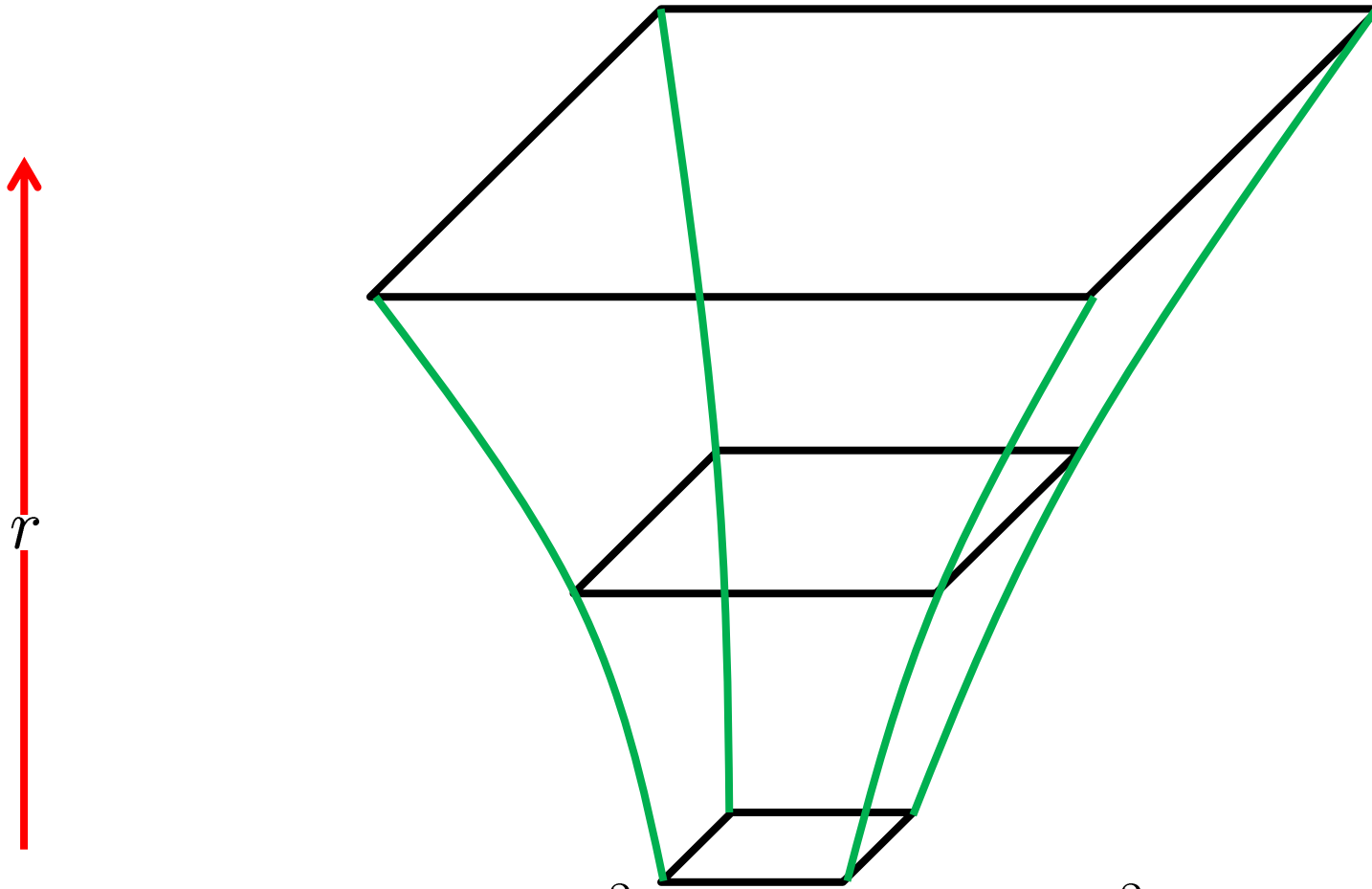
with small curvatures

→ large central charge ($N_c \rightarrow \infty$)

→ strong coupling ($\lambda \rightarrow \infty$)

anti-de Sitter space:

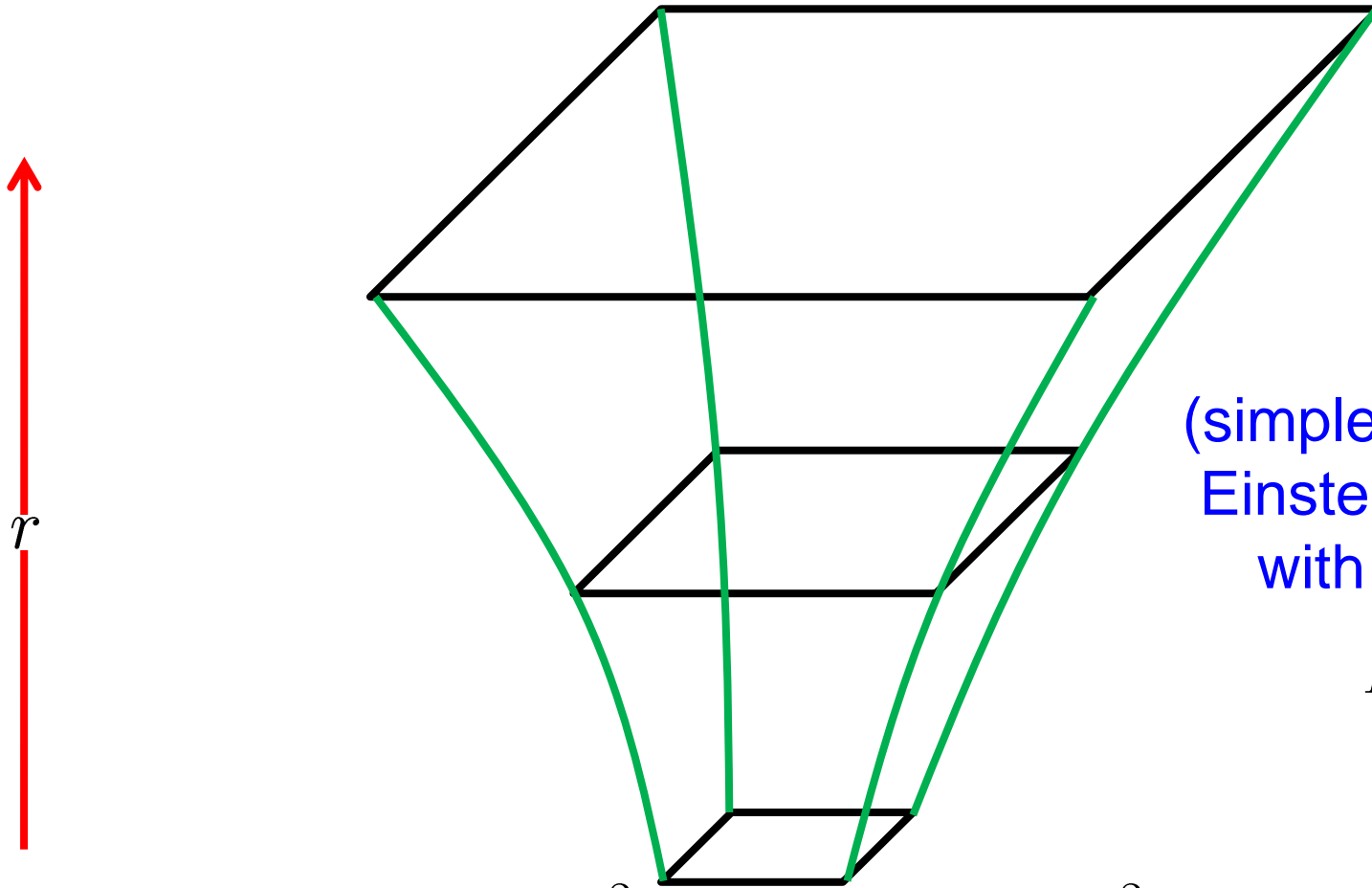
maximally symmetric geometry with negative curvature



$$ds^2 = \frac{r^2}{L^2} (-dt^2 + d\vec{x}^2) + \frac{L^2}{r^2} dr^2$$

anti-de Sitter space:

maximally symmetric geometry with negative curvature



(simplest) solution of
Einstein's equations
with negative Λ :

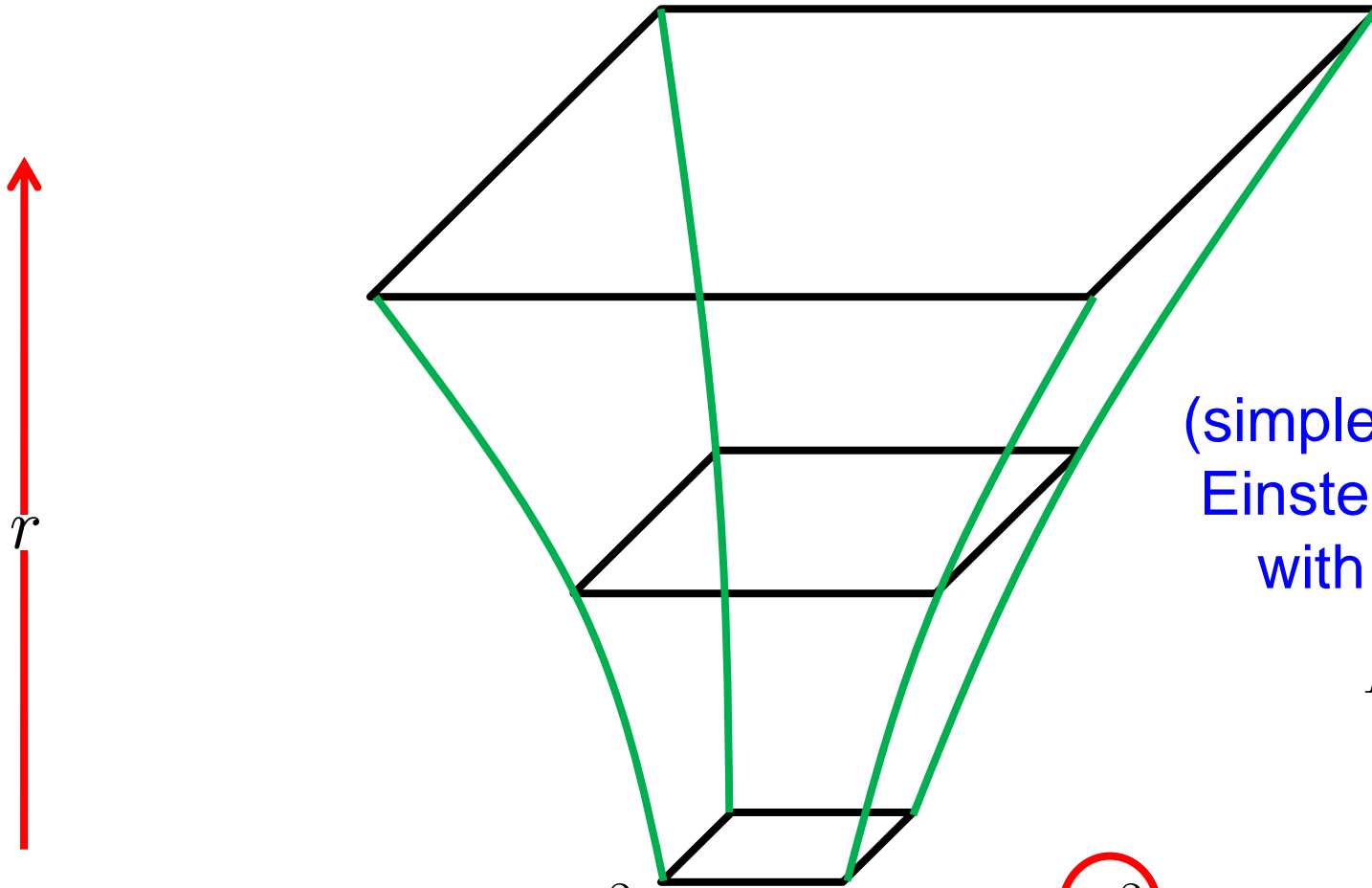
$$R_{ab} = -\frac{d}{L^2} g_{ab}$$

$$ds^2 = \frac{r^2}{L^2} (-dt^2 + d\vec{x}^2) + \frac{L^2}{r^2} dr^2$$

anti-de Sitter space:

$$R \sim -\frac{1}{L^2}$$

maximally symmetric geometry with negative curvature



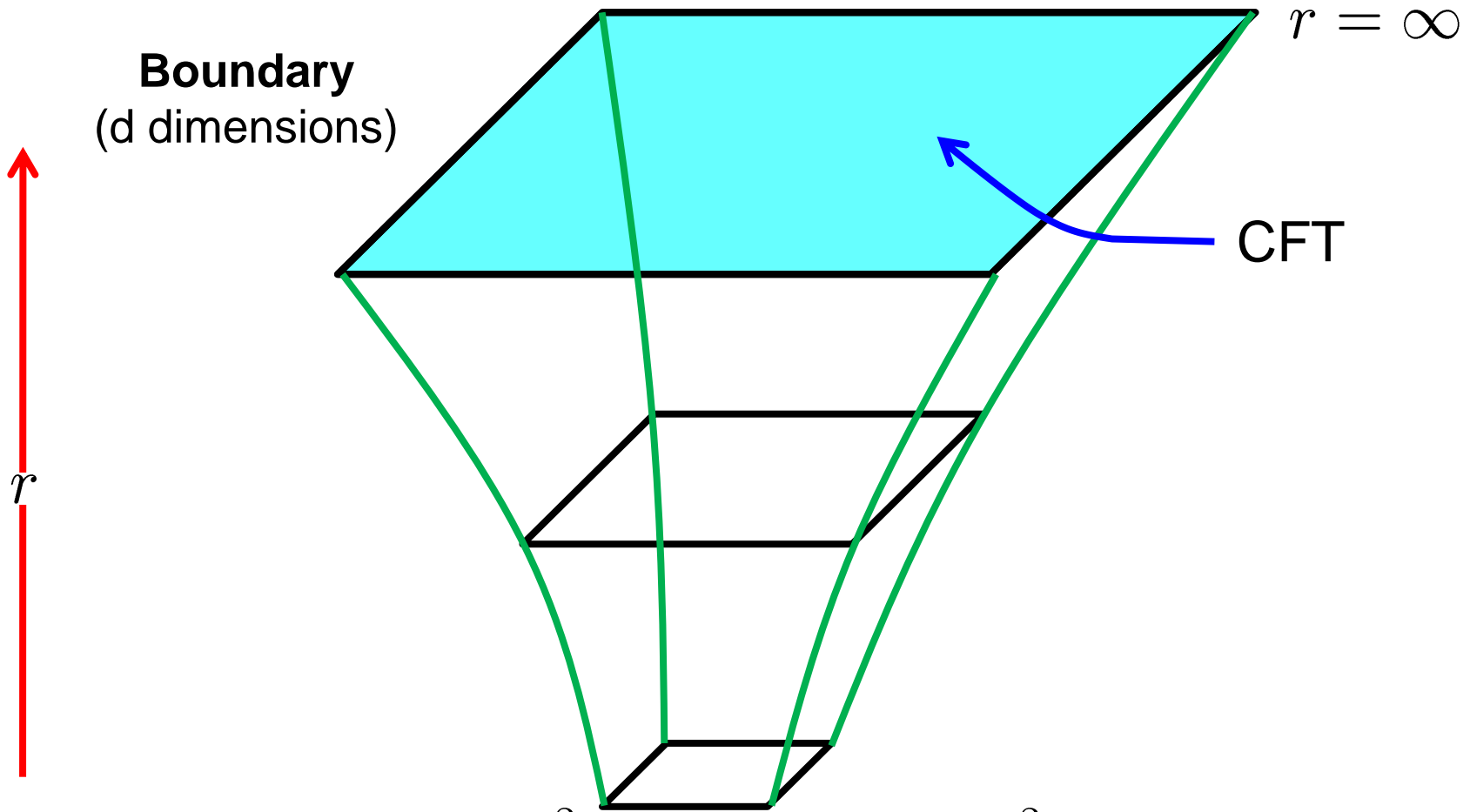
(simplest) solution of Einstein's equations with negative Λ :

$$R_{ab} = -\frac{d}{L^2}$$

$$ds^2 = \frac{r^2}{L^2} (-dt^2 + d\vec{x}^2) + \frac{L^2}{r^2} dr^2$$

anti-de Sitter space:

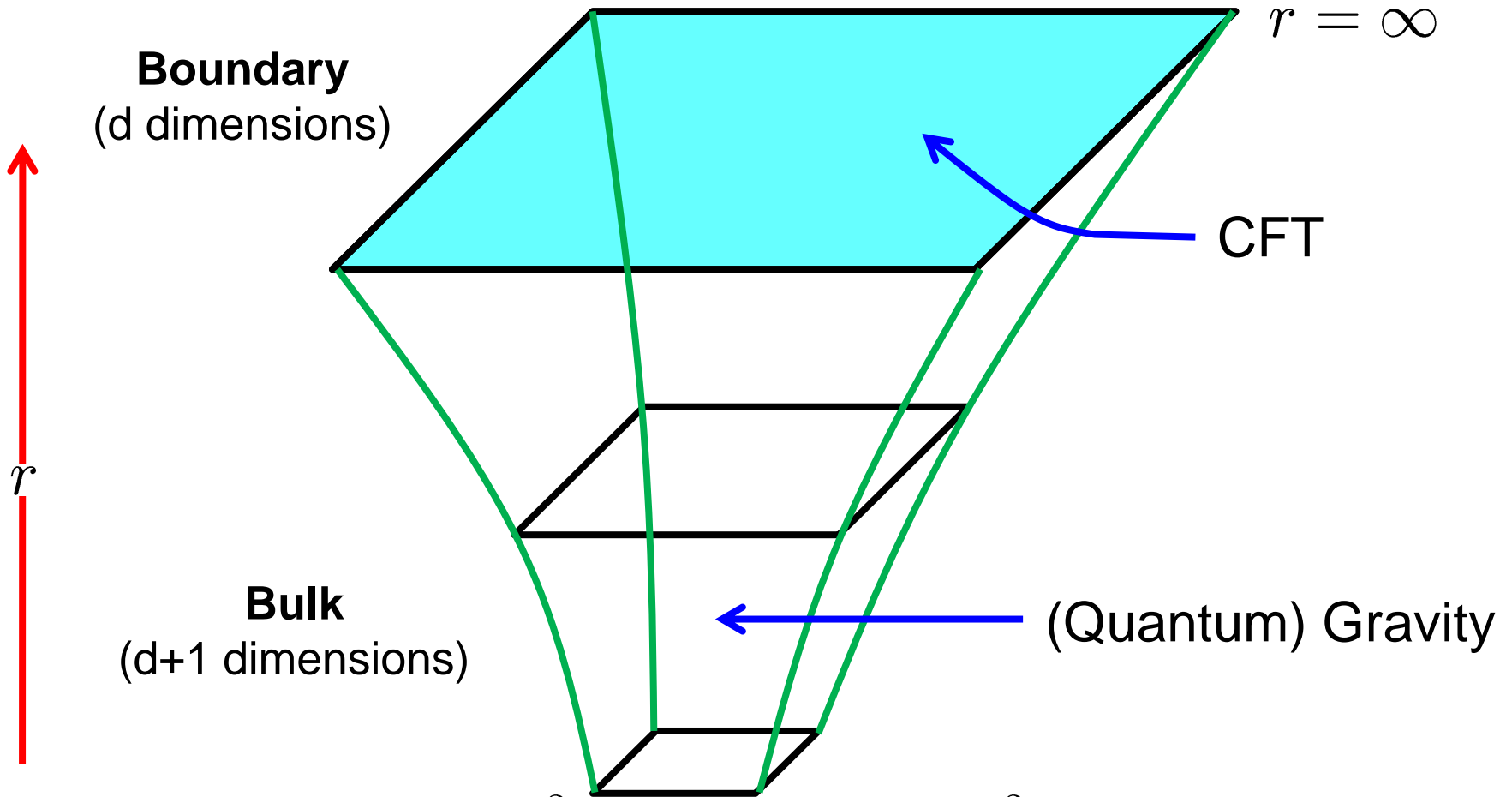
maximally symmetric geometry with negative curvature



$$ds^2 = \frac{r^2}{L^2} (-dt^2 + d\vec{x}^2) + \frac{L^2}{r^2} dr^2$$

anti-de Sitter space:

maximally symmetric geometry with negative curvature

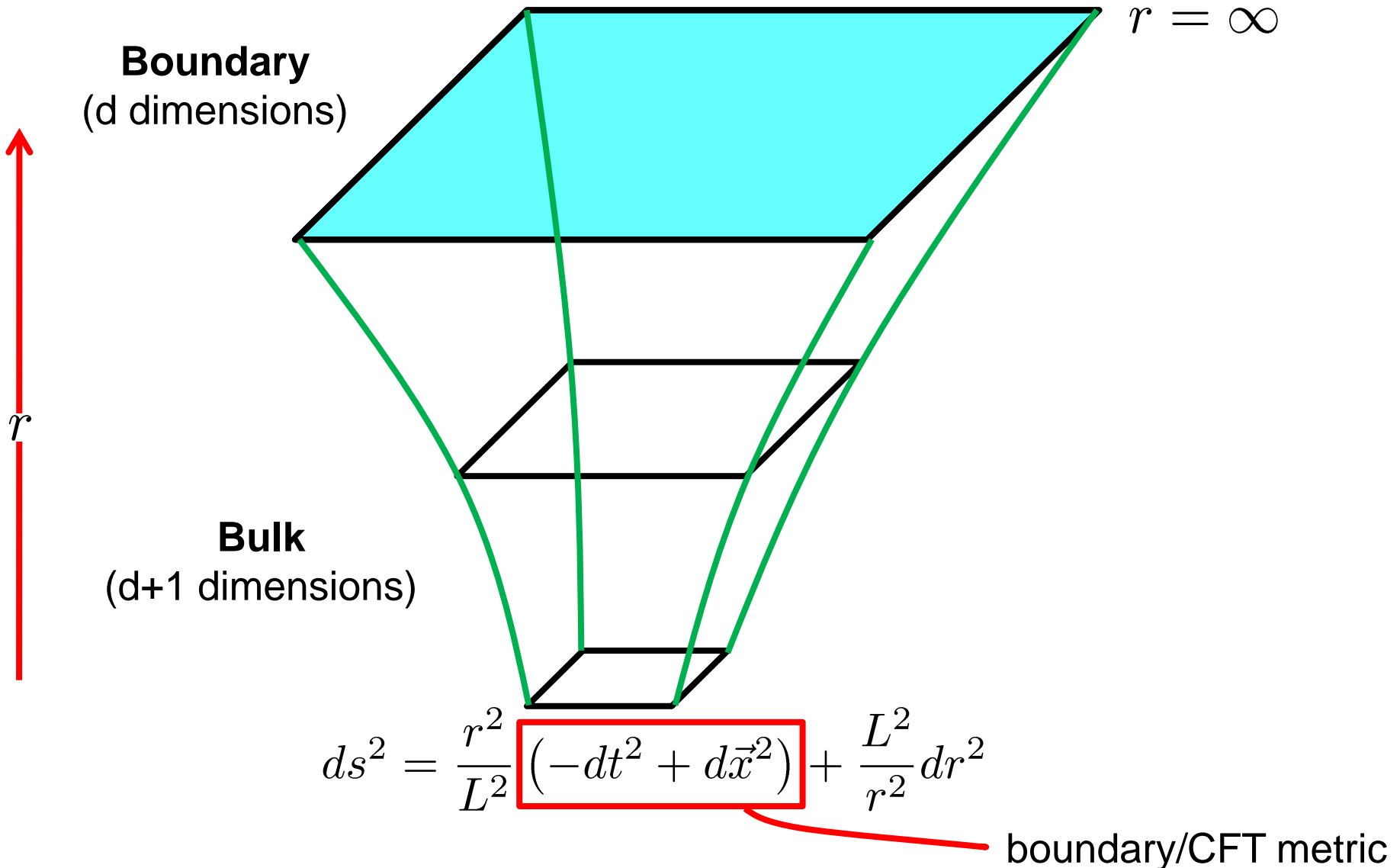


$$ds^2 = \frac{r^2}{L^2} (-dt^2 + d\vec{x}^2) + \frac{L^2}{r^2} dr^2$$

anti-de Sitter space:

$$R \sim -\frac{1}{L^2}$$

maximally symmetric geometry with negative curvature



Boundary
(d dimensions)

$$r = \infty$$



r

Bulk
(d+1 dimensions)

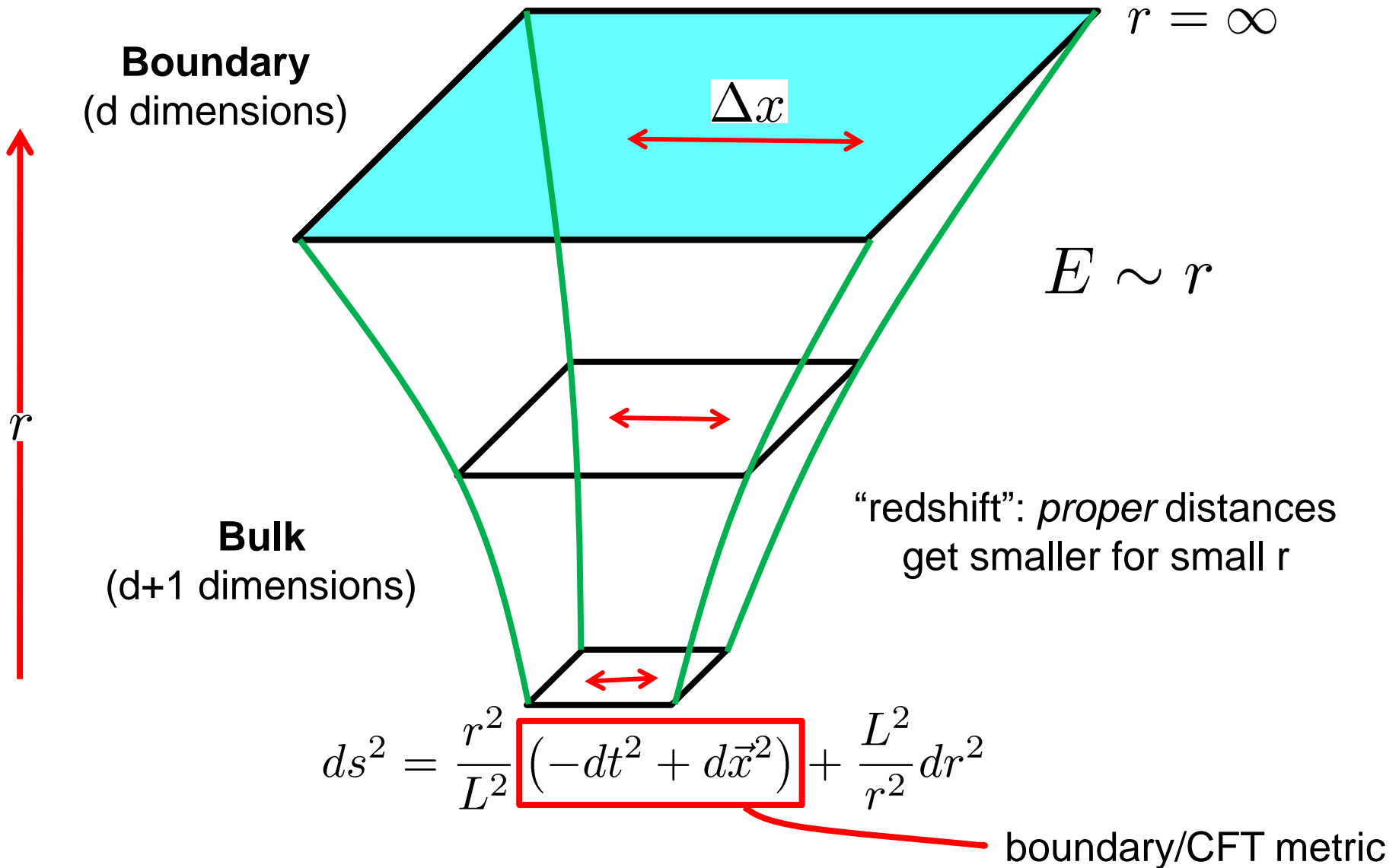
$$ds^2 = \frac{r^2}{L^2} (-dt^2 + d\vec{x}^2) + \frac{L^2}{r^2} dr^2$$

boundary/CFT metric

anti-de Sitter space:

$$R \sim -\frac{1}{L^2}$$

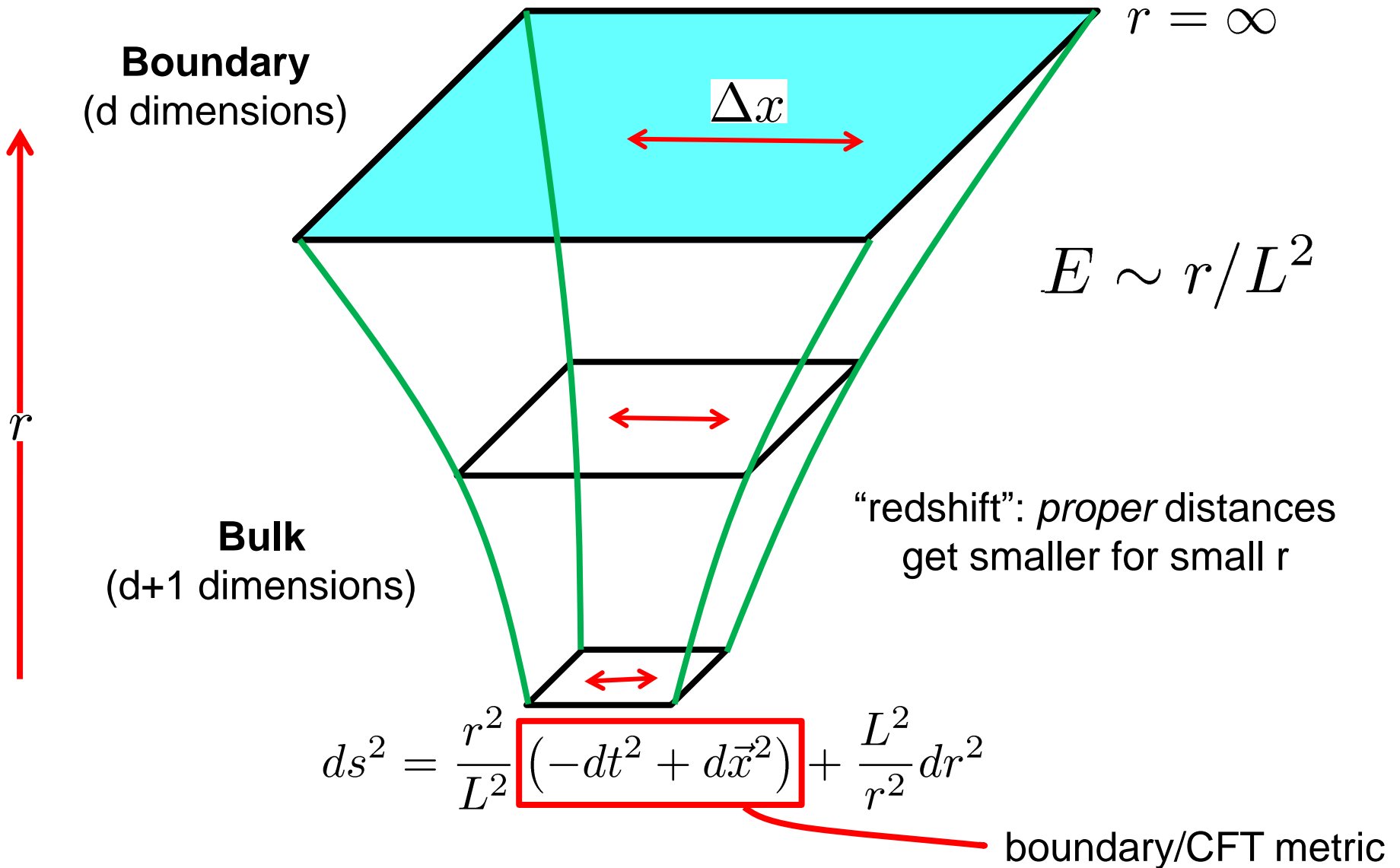
maximally symmetric geometry with negative curvature



anti-de Sitter space:

$$R \sim -\frac{1}{L^2}$$

maximally symmetric geometry with negative curvature



$$ds^2 = \frac{r^2}{L^2} (-dt^2 + d\vec{x}^2) + \frac{L^2}{r^2} dr^2$$

boundary/CFT metric

AdS/CFT Correspondence:

Bulk:

- quantum gravity
- negative cosmological constant
- **d+1** spacetime dimensions

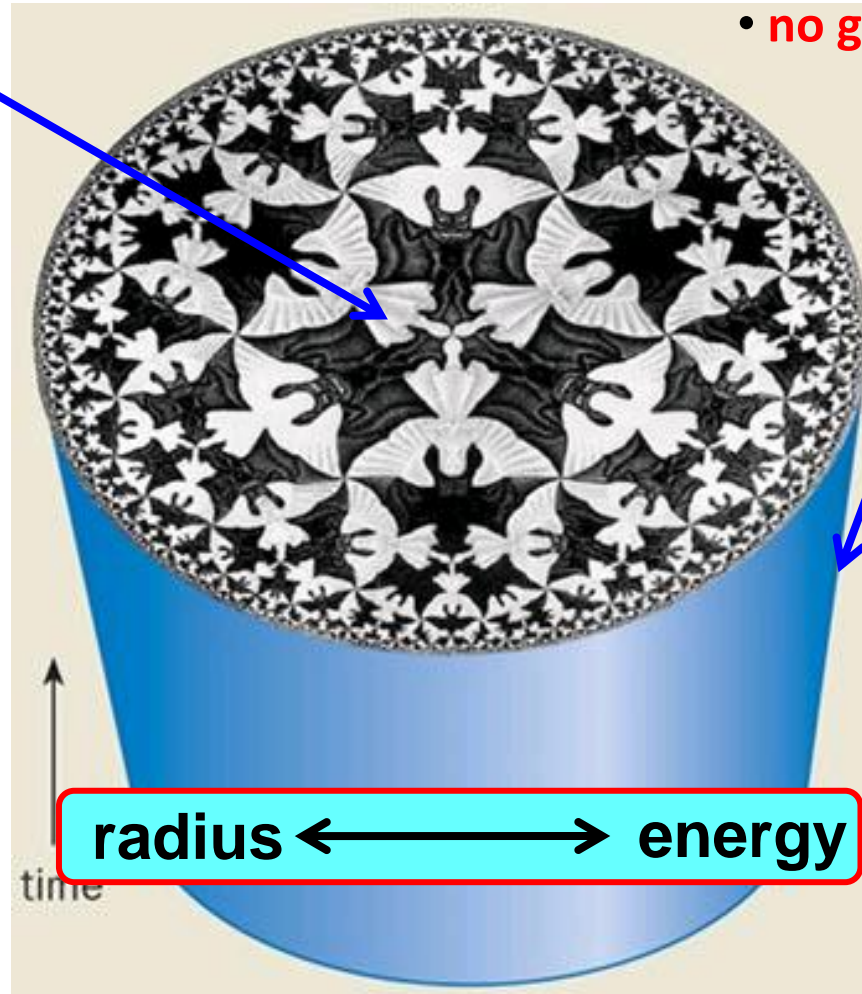
Boundary:

- quantum field theory
- no scale (at quantum level)
- **d** spacetime dimensions
- **no gravity!**

Holography



anti-de Sitter
space



conformal
field theory

radius \longleftrightarrow energy

(Maldacena '97)

Overview:

1) Entanglement entropy in QFT

2) Primer on AdS/CFT correspondence

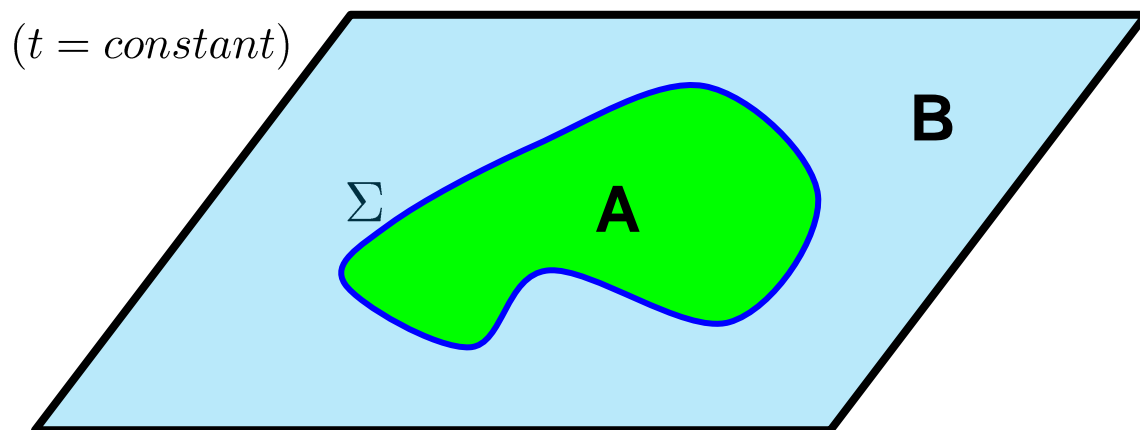
3) Introduction to Holographic Entanglement Entropy

4) Derivation of Holographic Entanglement Entropy
(for spherical entangling surfaces) and Beyond

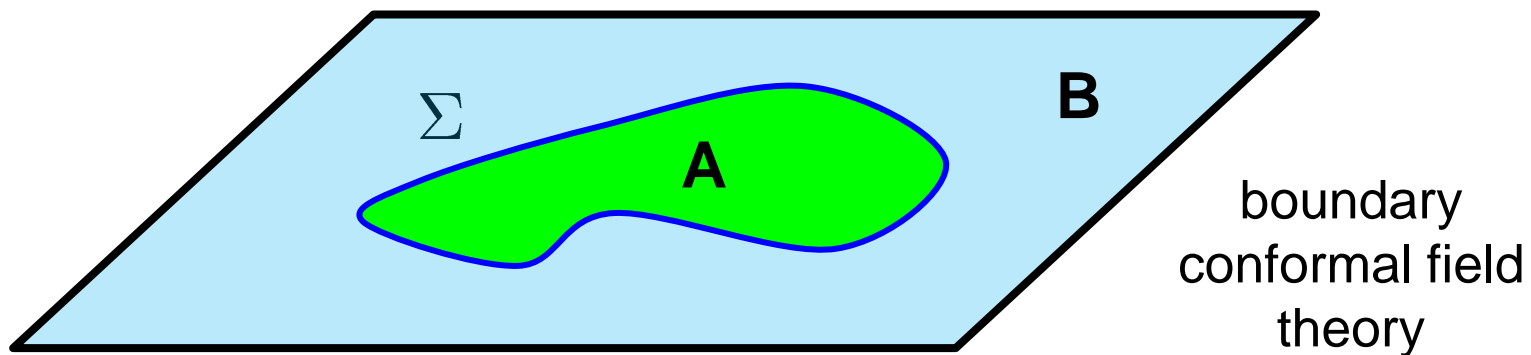
5) Conclusions

Entanglement Entropy

- general tool; divide quantum system into two parts and use entropy as measure of correlations between subsystems
 - in QFT, typically introduce a (smooth) boundary or **entangling surface** Σ which divides the space into two separate regions
 - integrate out degrees of freedom in “outside” region
 - remaining dof are described by a density matrix ρ_A
- calculate **von Neumann entropy**: $S_{EE} = -Tr [\rho_A \log \rho_A]$

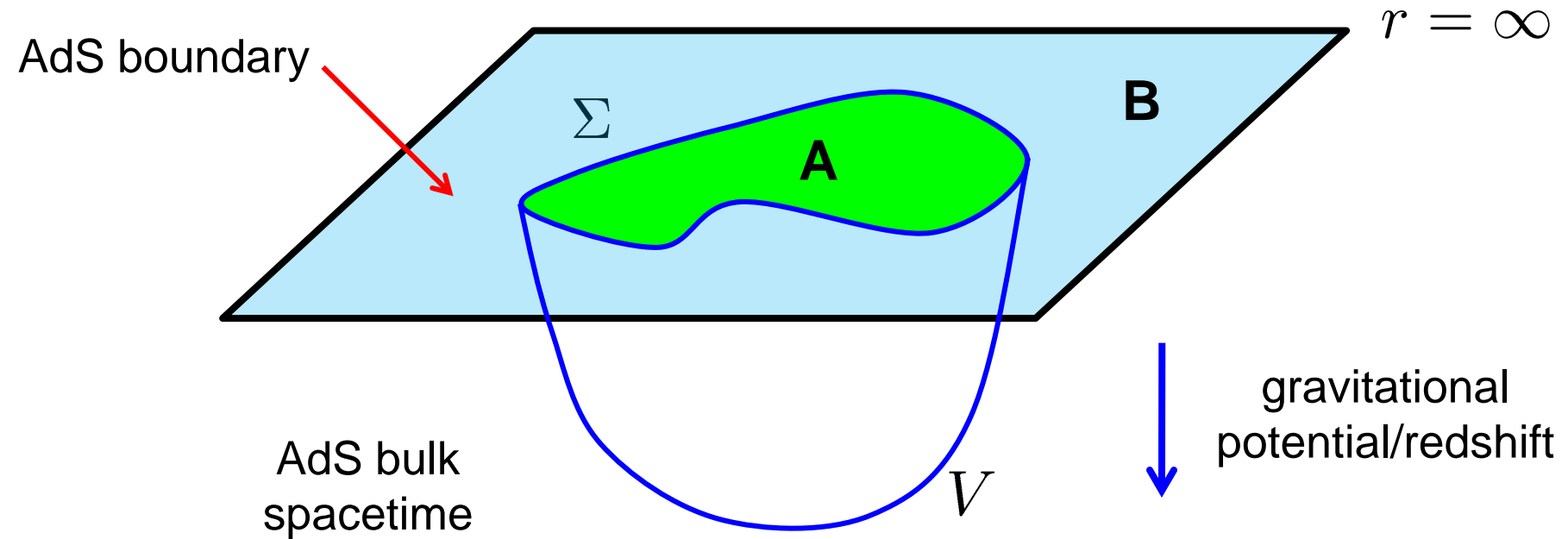


Holographic Entanglement Entropy:



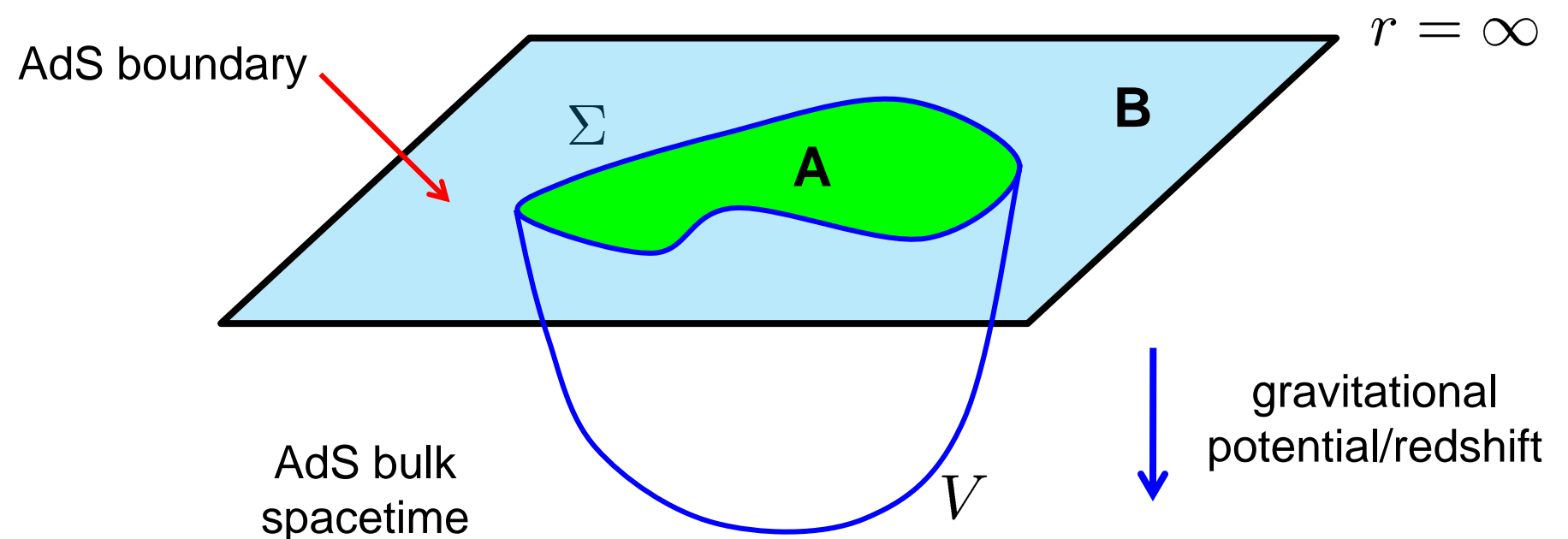
$$S(A) = ??$$

Holographic Entanglement Entropy:



$$S(A) = \min_{\partial V = \Sigma} \frac{A_V}{4G_N}$$

Holographic Entanglement Entropy:

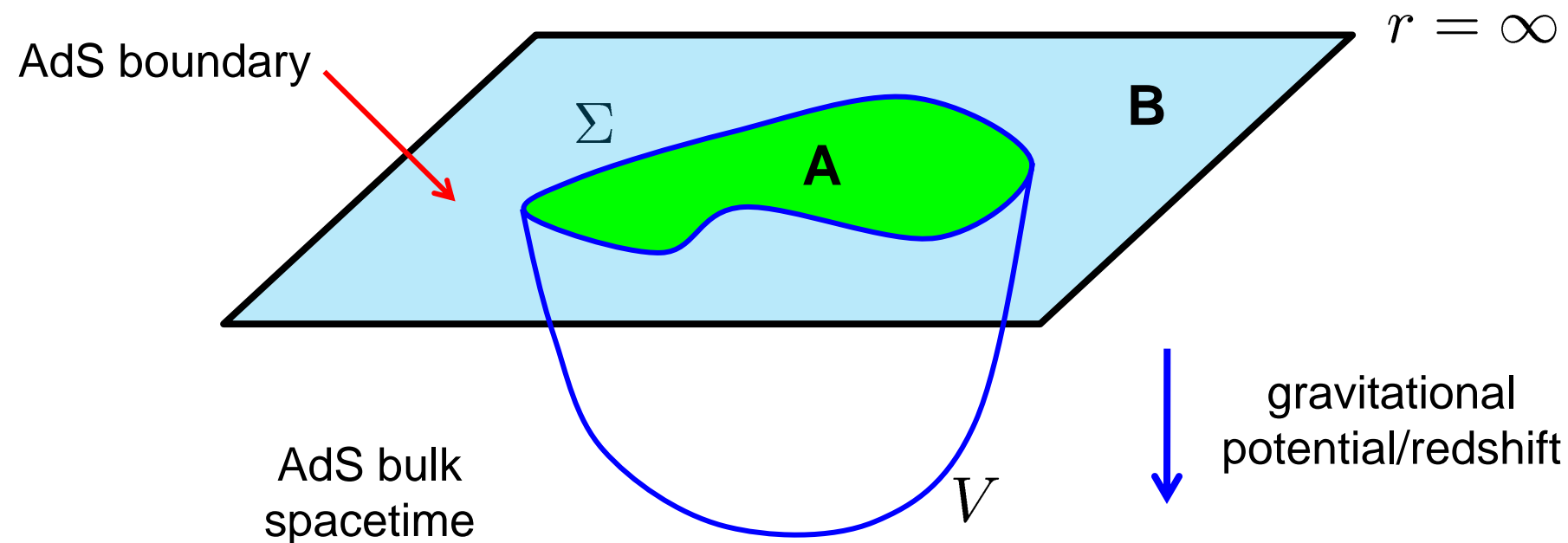


$$S(A) = \text{ext}_{V \sim A} \frac{A_V}{4G_N}$$

$(d - 1)$ dimensional

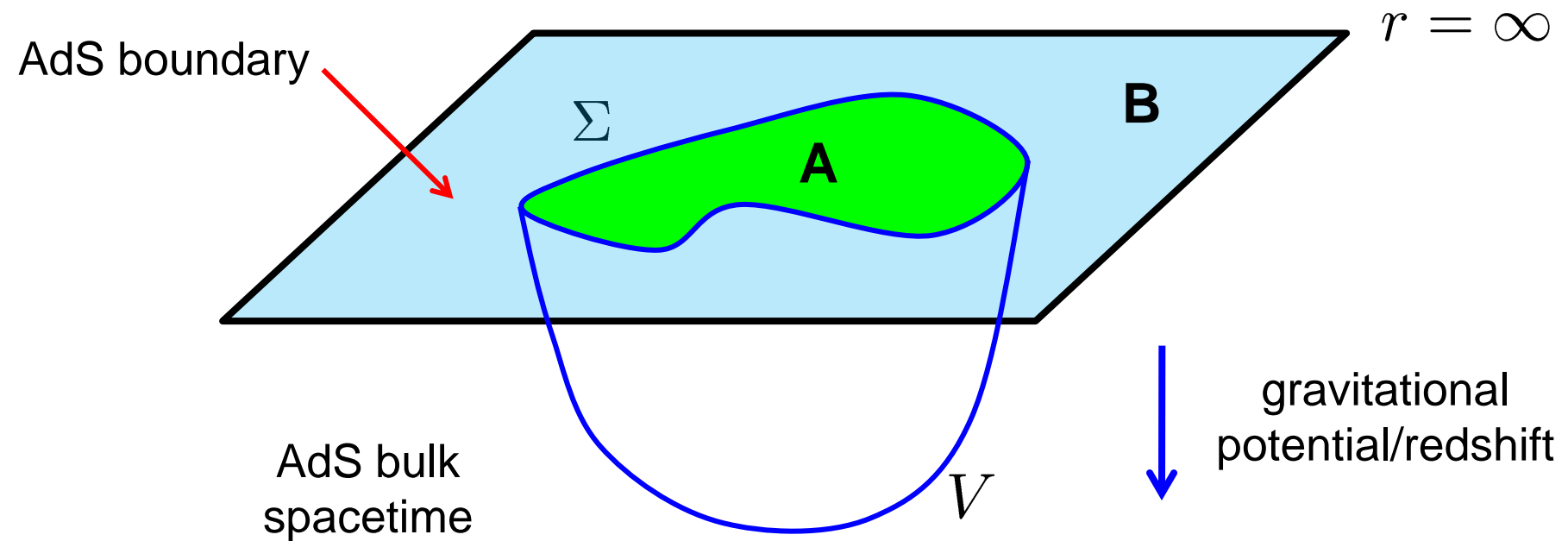
looks like
BH entropy!

Holographic Entanglement Entropy:



$$S(A) = \text{ext}_{V \sim A} \frac{A_V}{4G_N} = \infty!!$$

Holographic Entanglement Entropy:

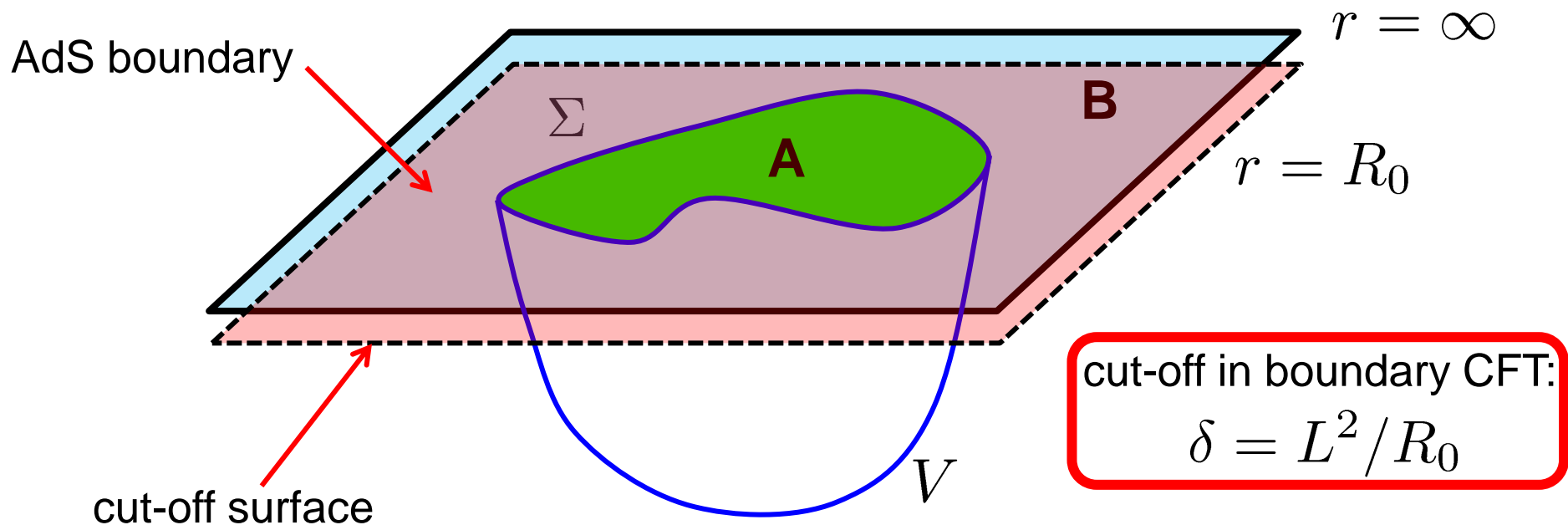


$$S(A) = \text{ext}_{V \sim A} \frac{A_V}{4G_N} = \infty!!$$

- “UV divergence” because area integral extends to $r = \infty$

(a feature! **not** a bug!)

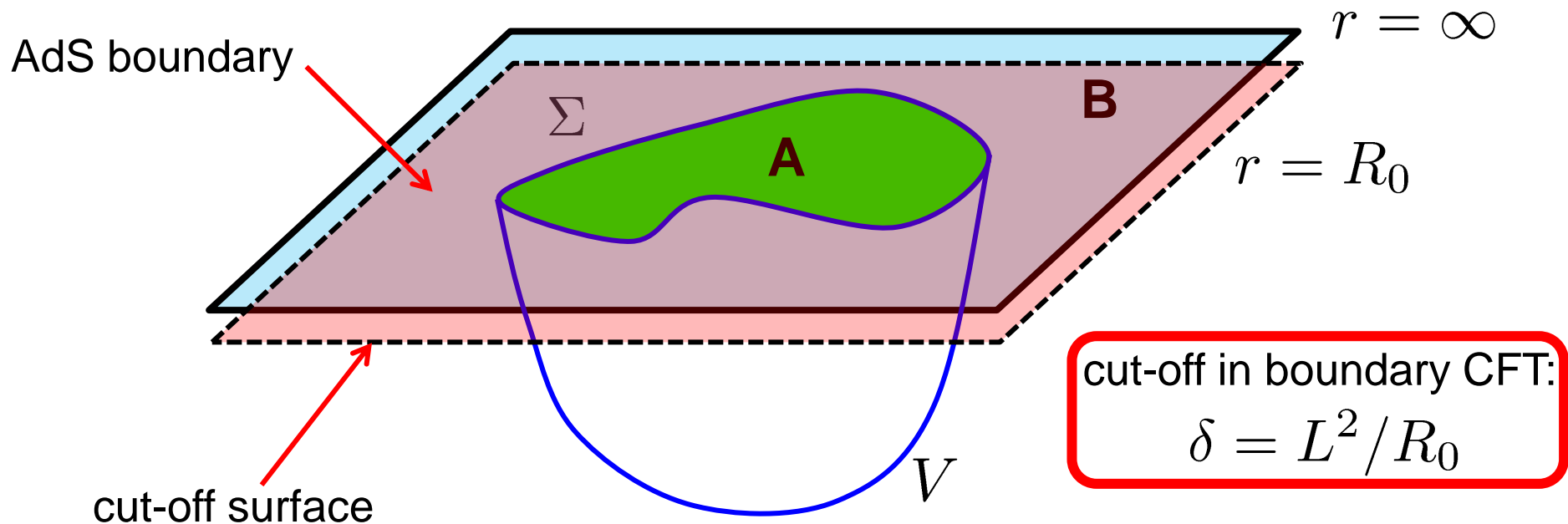
Holographic Entanglement Entropy:



$$S(A) = \text{ext}_{\partial V = \Sigma} \frac{A_V}{4G_N}$$

- “UV divergence” because area integral extends to $r = \infty$
- finite result by stopping radial integral at large radius: $r = R_0$
 → short-distance cut-off in boundary theory: $\delta = L^2 / R_0$

Holographic Entanglement Entropy:



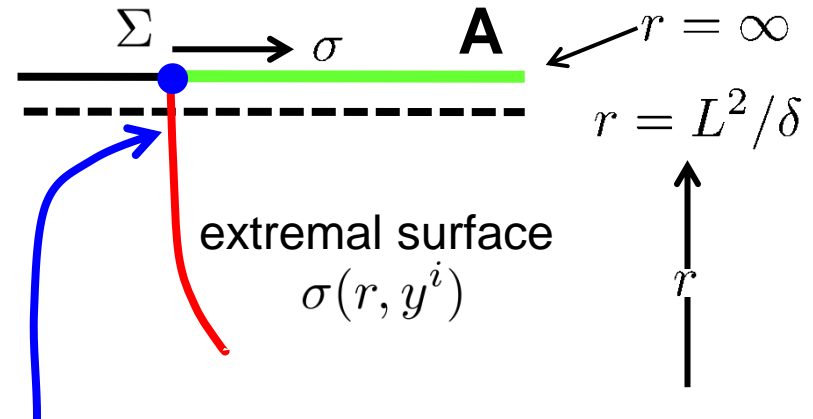
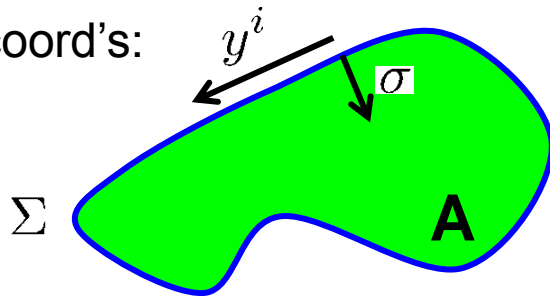
$$S(A) = \text{ext}_{\partial V = \Sigma} \frac{A_V}{4G_N} \simeq \frac{L^{d-1}}{G_N} \frac{\mathcal{A}_\Sigma}{\delta^{d-2}} + \dots$$

↑
"Area Law"

Area law contribution:

recall AdS metric: $ds^2 = \frac{r^2}{L^2} (-dt^2 + d\vec{x}^2) + \frac{L^2}{r^2} dr^2$

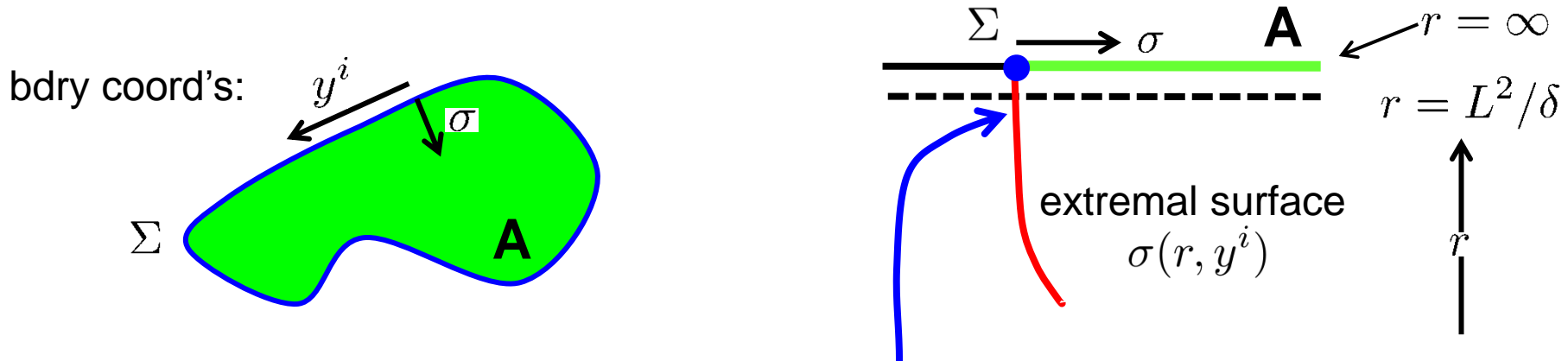
bdry coord's:



to leading order, surface falls straight down, ie, $\sigma \simeq 0$

Area law contribution:

recall AdS metric:
$$ds^2 = \frac{r^2}{L^2} (-dt^2 + d\vec{x}^2) + \frac{L^2}{r^2} dr^2$$



to leading order, surface falls straight down, ie, $\sigma \simeq 0$

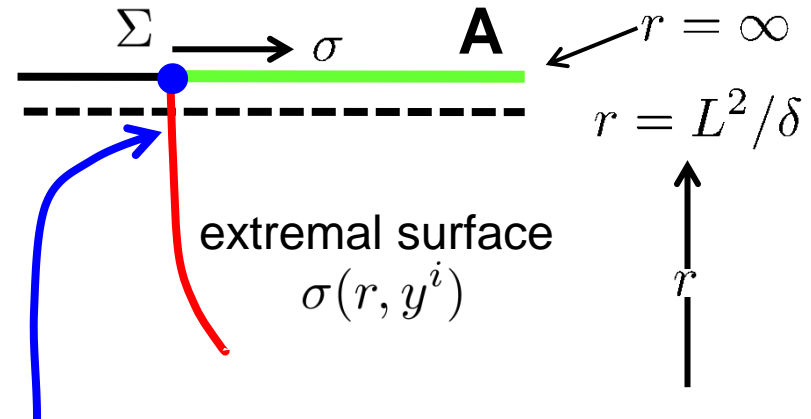
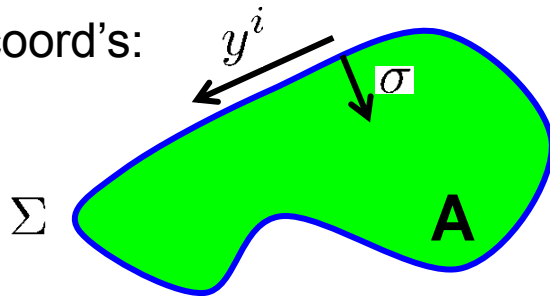
$$S(A) = \frac{A_V}{4G_N} = \frac{1}{4G_N} \int d^{d-1} \sigma \sqrt{h}$$

$$= \frac{1}{4G_N} \int^{R_0} dr \left(\frac{L}{r} \right) \times \int d^{d-2} y \sqrt{h_y} \left(\frac{r}{L} \right)^{d-2}$$

Area law contribution:

recall AdS metric:
$$ds^2 = \frac{r^2}{L^2} (-dt^2 + d\vec{x}^2) + \frac{L^2}{r^2} dr^2$$

bdry coord's:



to leading order, surface falls straight down, ie, $\sigma \simeq 0$

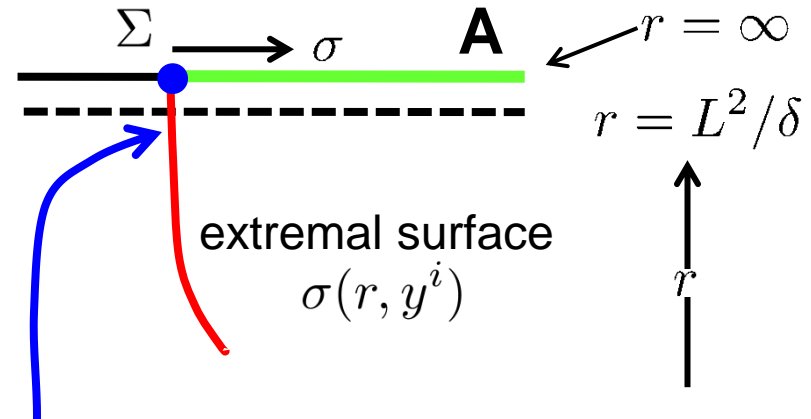
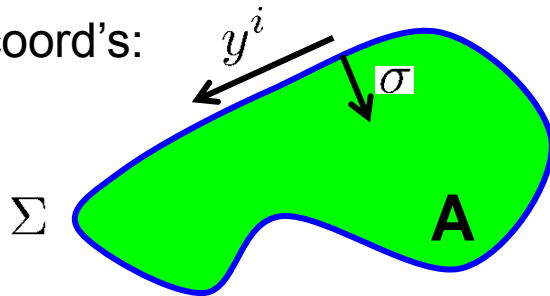
$$S(A) = \frac{A_V}{4G_N} = \frac{1}{4G_N} \int d^{d-1} \sigma \sqrt{h}$$

$$= \frac{1}{4G_N} \mathcal{A}_\Sigma \int^{L^2/\delta} dr \left(\frac{r}{L} \right)^{d-3} \simeq \frac{L^{d-1}}{4(d-2)G_N} \frac{\mathcal{A}_\Sigma}{\delta^{d-2}}$$

Area law contribution:

recall AdS metric:
$$ds^2 = \frac{r^2}{L^2} (-dt^2 + d\vec{x}^2) + \frac{L^2}{r^2} dr^2$$

bdry coord's:



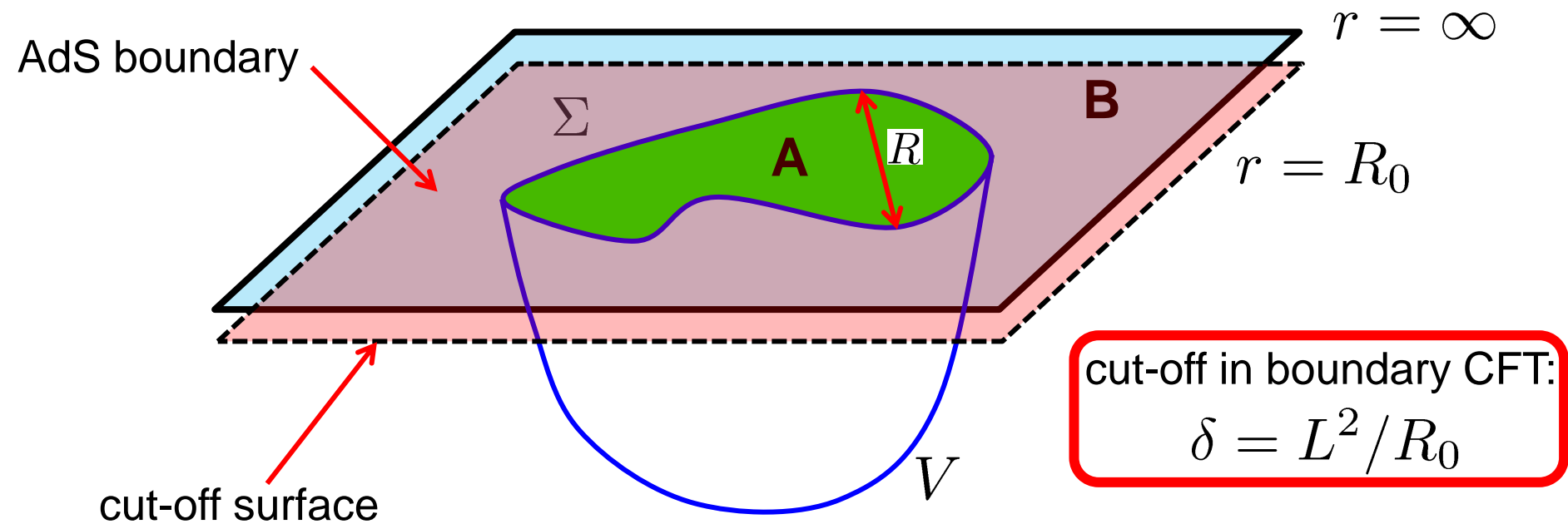
to leading order, surface falls straight down, ie, $\sigma \simeq 0$

$$S(A) = \frac{A_V}{4G_N} = \frac{1}{4G_N} \int d^{d-1} \sigma \sqrt{h}$$

$$= \frac{1}{4G_N} \mathcal{A}_\Sigma \int^{L^2/\delta} dr \left(\frac{r}{L} \right)^{d-3} \sim \frac{L^{d-1}}{\ell_{Planck}^{d-1}} \frac{\mathcal{A}_\Sigma}{\delta^{d-2}}$$

central charge
(counts dof)

Holographic Entanglement Entropy:



general expression (as desired):

$$S(A) \simeq c_0 (R/\delta)^{d-2} + c_1 (R/\delta)^{d-4} + \dots$$

$$\left\{ \begin{array}{l} + c_{d-2} \log(R/\delta) + \dots \quad (\text{d even}) \\ + \underbrace{c_{d-2} + \dots}_{\text{universal contributions}} \quad (\text{d odd}) \end{array} \right.$$

universal contributions

Holographic Entanglement Entropy:

$$S(A) = \text{ext}_{\partial V = \Sigma} \frac{A_V}{4G_N}$$

appeared as conjecture!

Holographic Entanglement Entropy:

$$S(A) = \text{ext}_{\partial V = \Sigma} \frac{A_V}{4G_N}$$

appeared as conjecture!

Extensive consistency tests:

1) leading contribution yields “area law”

$$S \simeq \frac{L^{d-1}}{G_N} \frac{\mathcal{A}_\Sigma}{\delta^{d-2}} + \dots$$

Holographic Entanglement Entropy:

$$S(A) = \text{ext}_{\partial V = \Sigma} \frac{A_V}{4G_N}$$

appeared as conjecture!

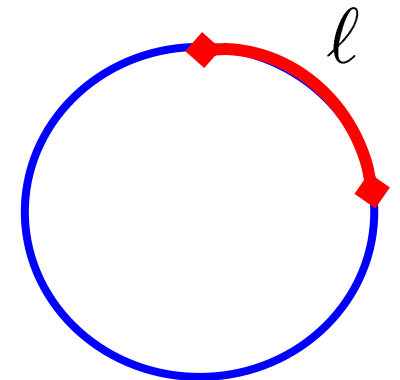
Extensive consistency tests:

1) leading contribution yields “area law” $S \simeq \frac{L^{d-1}}{G_N} \frac{\mathcal{A}_\Sigma}{\delta^{d-2}} + \dots$

2) recover known results of Calabrese & Cardy
for d=2 CFT

$$S = \frac{c}{3} \log \left(\frac{C}{\pi \delta} \sin \frac{\pi \ell}{C} \right)$$

(also result for thermal ensemble)



$C = \text{circumference}$

Holographic Entanglement Entropy:

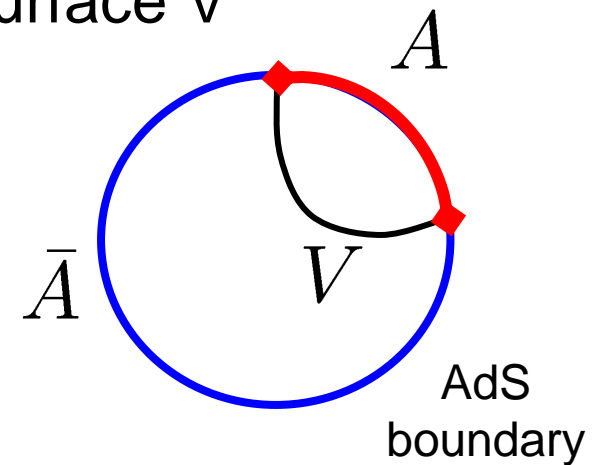
$$S(A) = \text{ext}_{\partial V = \Sigma} \frac{A_V}{4G_N}$$

appeared as conjecture!

Extensive consistency tests:

3) $S(A) = S(\bar{A})$ in a pure state

→ A and \bar{A} both yield same bulk surface V



Holographic Entanglement Entropy:

$$S(A) = \text{ext}_{\partial V = \Sigma} \frac{A_V}{4G_N}$$

appeared as conjecture!

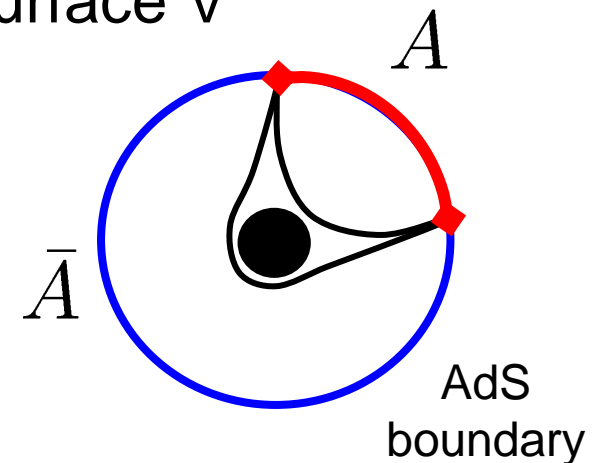
Extensive consistency tests:

3) $S(A) = S(\bar{A})$ in a pure state

→ A and \bar{A} both yield same bulk surface V

cf: thermal ensemble \neq pure state

horizon in bulk → $S(A) \neq S(\bar{A})$



AdS/CFT Dictionary: thermal bath ↔ black hole

Holographic Entanglement Entropy:

$$S(A) = \text{ext}_{\partial V = \Sigma} \frac{A_V}{4G_N}$$

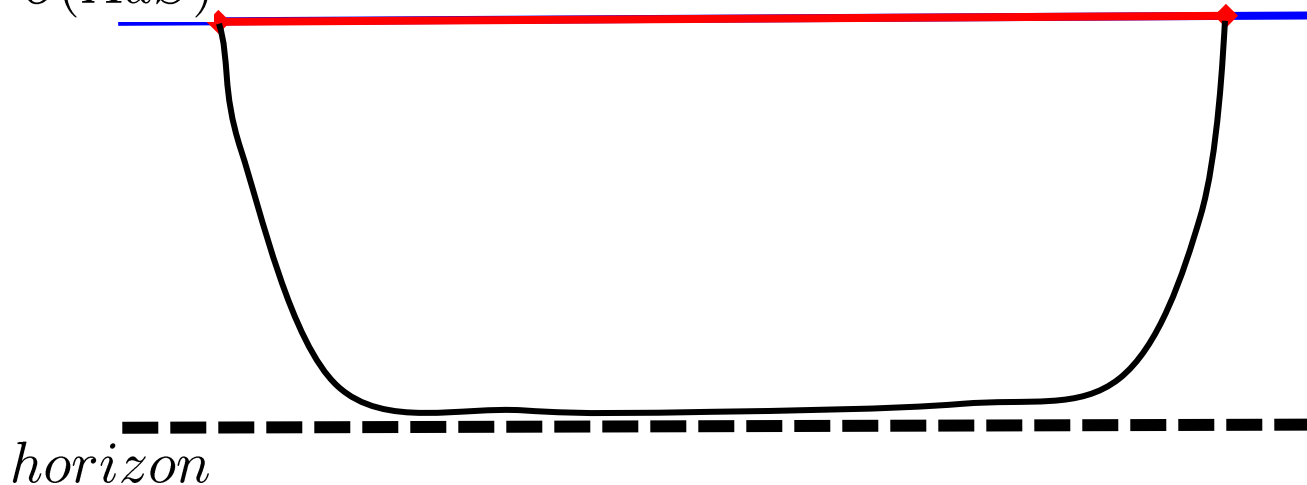
appeared as conjecture!

Extensive consistency tests:

4) for thermal bath: $S(A) \supset S_{therm} = \alpha T^{d-1} \times volume$

$\partial(AdS)$ A

$R \gg 1/T$



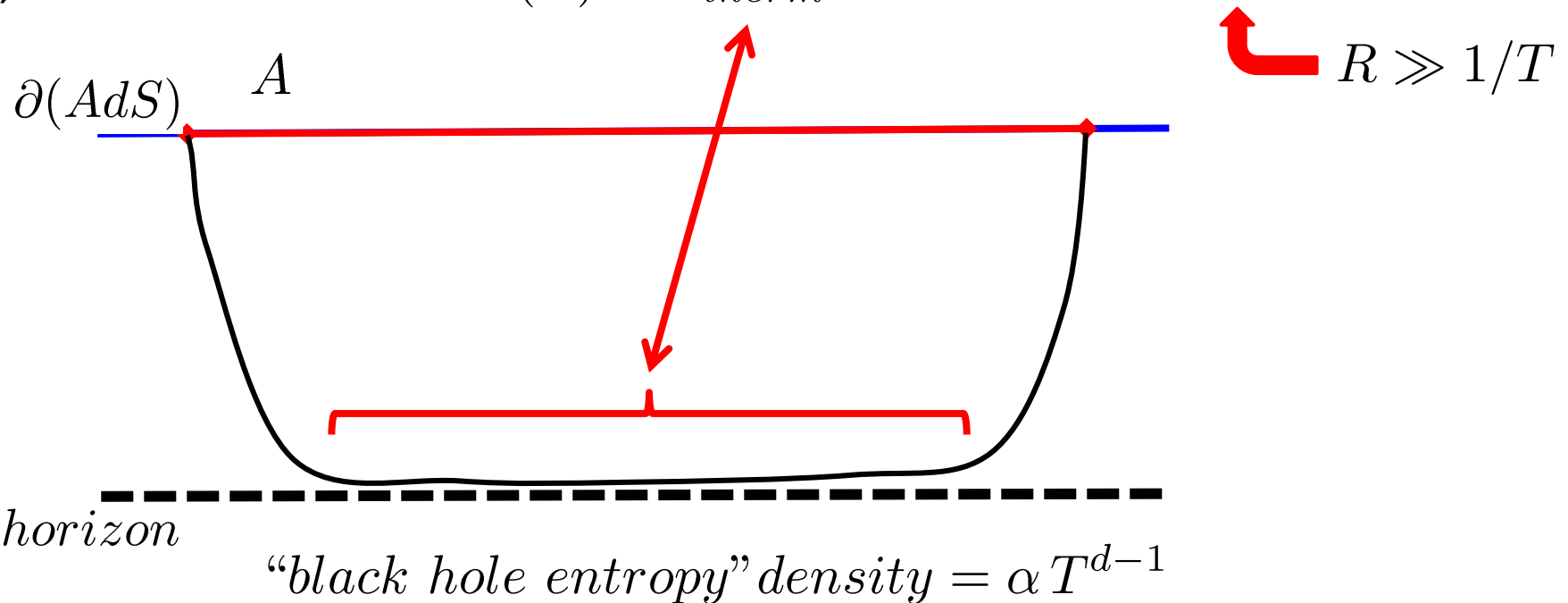
Holographic Entanglement Entropy:

$$S(A) = \text{ext}_{\partial V = \Sigma} \frac{A_V}{4G_N}$$

appeared as conjecture!

Extensive consistency tests:

4) for thermal bath: $S(A) \supset S_{therm} = \alpha T^{d-1} \times \text{volume}$



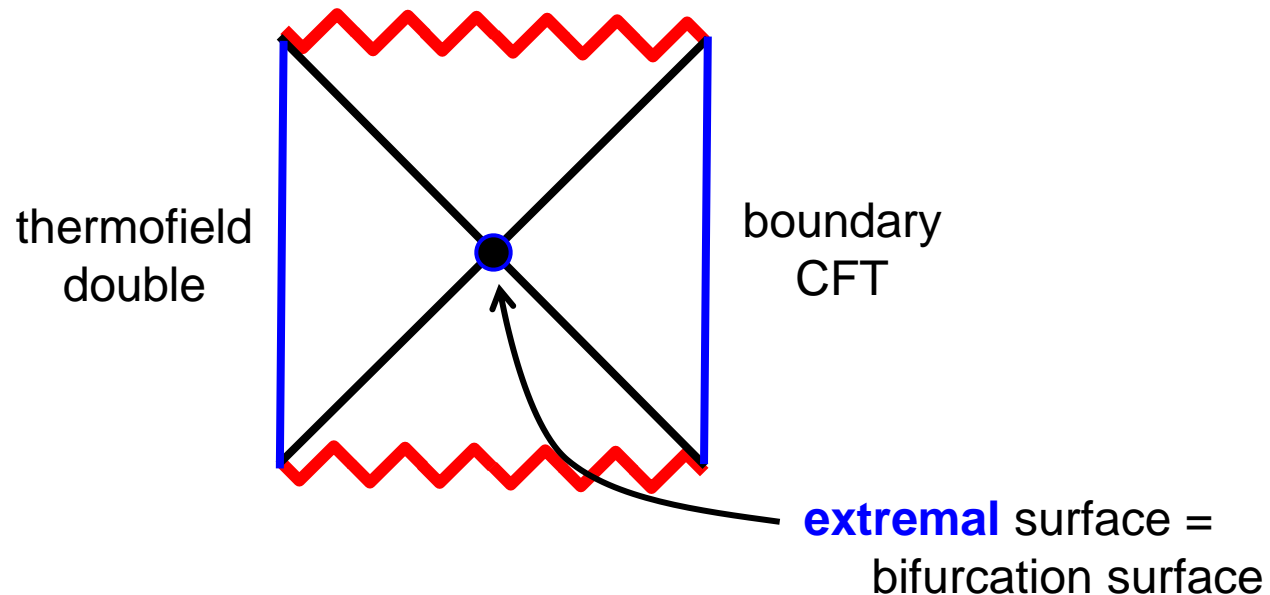
Holographic Entanglement Entropy:

$$S(A) = \text{ext}_{\partial V = \Sigma} \frac{A_V}{4G_N}$$

appeared as conjecture!

Extensive consistency tests:

4b) Entropy of eternal black hole =
entanglement entropy of boundary CFT & thermofield double
(Headrick)



Holographic Entanglement Entropy:

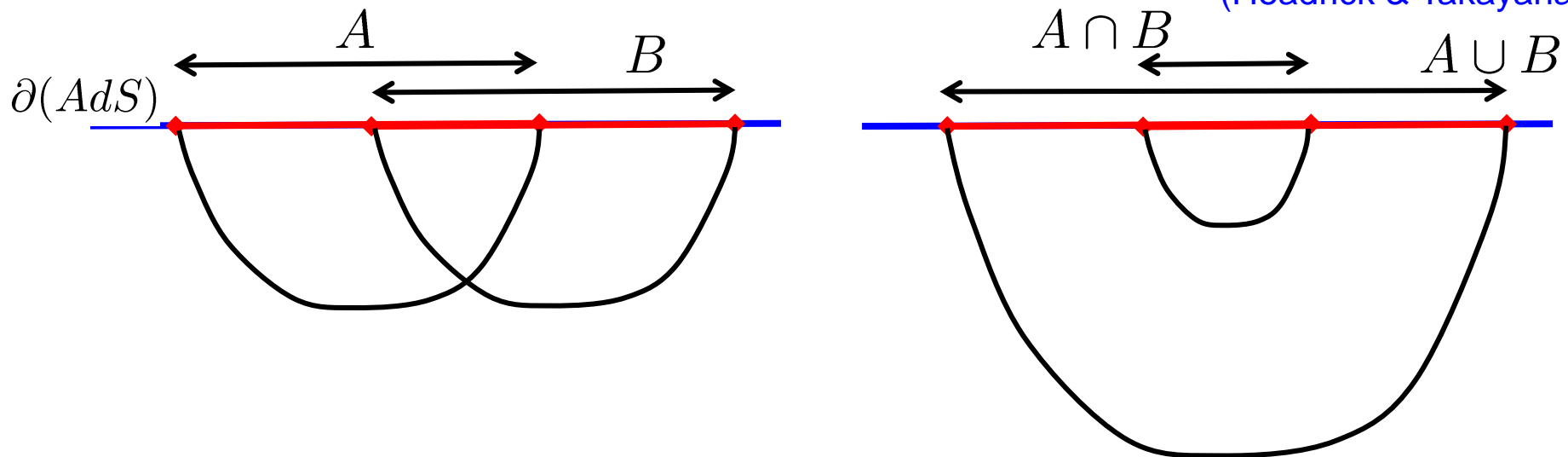
$$S(A) = \text{ext}_{\partial V = \Sigma} \frac{A_V}{4G_N}$$

appeared as conjecture!

Extensive consistency tests:

5) strong sub-additivity: $S(A) + S(B) \geq S(A \cup B) + S(A \cap B)$

(Headrick & Takayanagi)



Holographic Entanglement Entropy:

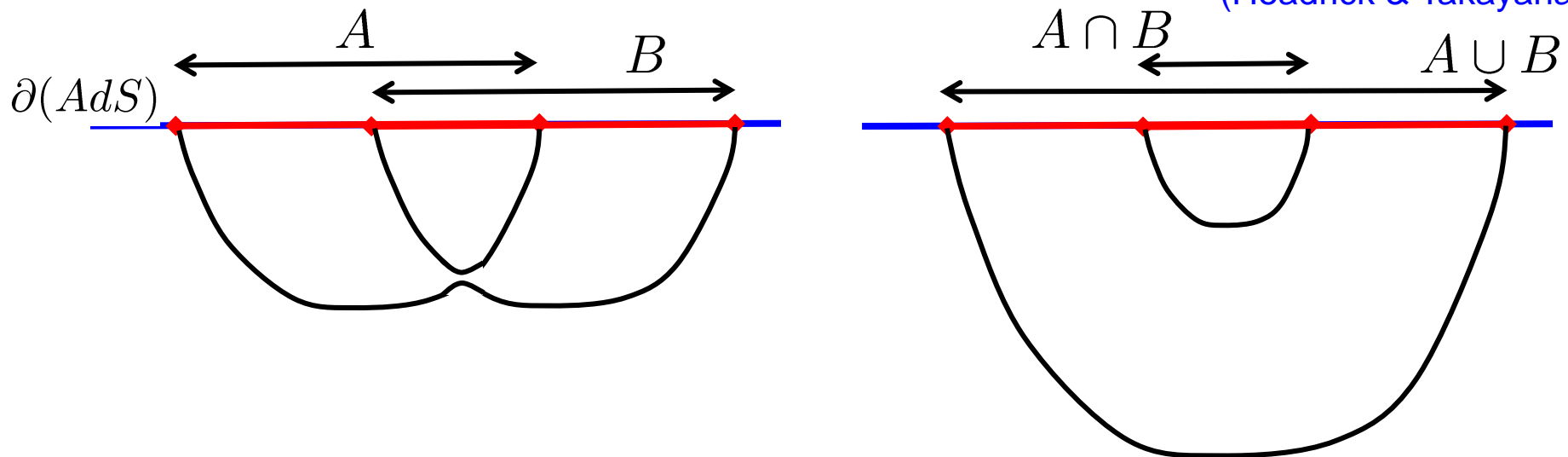
$$S(A) = \text{ext}_{\partial V = \Sigma} \frac{A_V}{4G_N}$$

appeared as conjecture!

Extensive consistency tests:

5) strong sub-additivity: $S(A) + S(B) \geq S(A \cup B) + S(A \cap B)$

(Headrick & Takayanagi)



Holographic Entanglement Entropy:

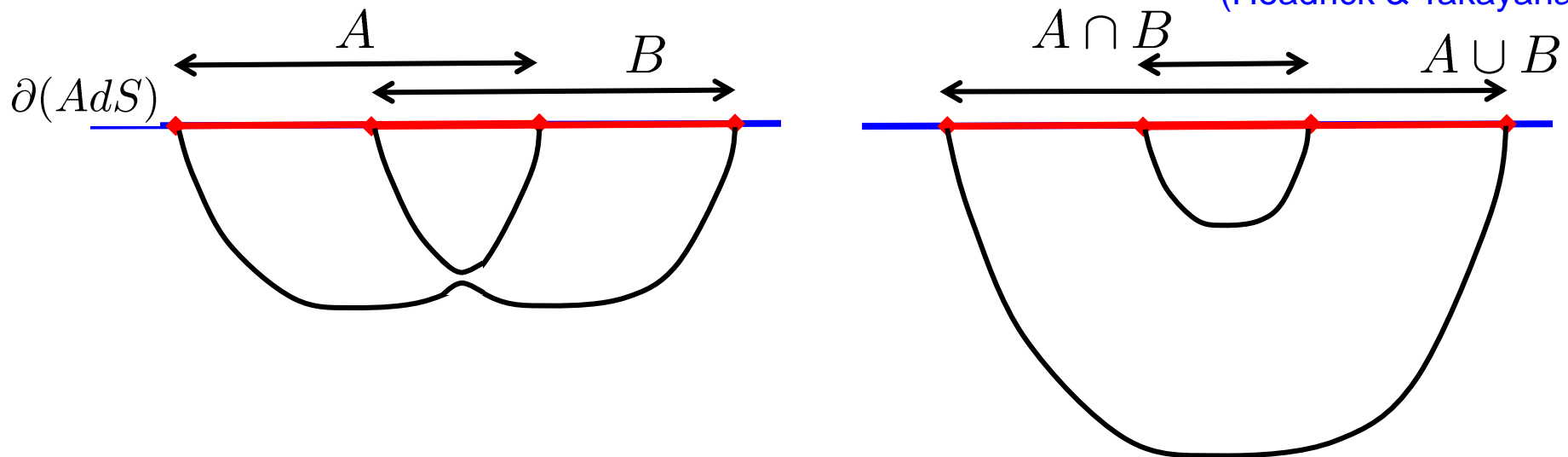
$$S(A) = \text{ext}_{\partial V = \Sigma} \frac{A_V}{4G_N}$$

appeared as conjecture!

Extensive consistency tests:

5) strong sub-additivity: $S(A) + S(B) \geq S(A \cup B) + S(A \cap B)$

(Headrick & Takayanagi)



[extended to dynamical setting: [Wall](#)]

[further monogamy relations: [Hayden, Headrick & Maloney](#)]

Holographic Entanglement Entropy:

$$S(A) = \text{ext}_{\partial V = \Sigma} \frac{A_V}{4G_N}$$

appeared as conjecture!

Extensive consistency tests:

6) for even d , connection of universal/logarithmic contribution in S_{EE} to central charges of boundary CFT, eg, in $d=4$

$$S_{uni} = \log(R/\delta) \frac{1}{2\pi} \int_{\Sigma} d^2x \sqrt{h} \left[\mathbf{c} \left(C^{ijkl} \tilde{g}_{ik}^{\perp} \tilde{g}_{jl}^{\perp} - K_a^{ib} K_b^{ia} + \frac{1}{2} K_a^{ia} K_b^{ib} \right) - \mathbf{a} \mathcal{R} \right]$$

(Hung, RM & Smolkin)

7) derivation of holographic EE for spherical entangling surfaces

(Casini, Huerta & RM, RM & Sinha)

Holographic Entanglement Entropy:

$$S(A) = \text{ext}_{\partial V = \Sigma} \frac{A_V}{4G_N}$$

appeared as conjecture!

Extensive consistency tests: **→ recent proof!!!**
(Lewkowycz & Maldacena)

Holographic Entanglement Entropy:

$$S(A) = \text{ext}_{\partial V = \Sigma} \frac{A_V}{4G_N} \quad \text{— appeared as conjecture! —}$$

Extensive consistency tests: \longrightarrow **recent proof!!!**

(Lewkowycz & Maldacena)

- generalization of Euclidean path integral calc's for S_{BH} , extended to “periodic” bulk solutions without Killing vector

- for AdS/CFT, translates replica trick for boundary CFT to bulk

$$\Delta\tau = 2\pi \rightarrow 2\pi n \quad \longrightarrow \quad \log Z(n) = \log \text{Tr} [\rho^n] = -I_{\text{grav}}(n)$$

$$\longrightarrow \quad S = -n \partial_n [\log Z(n) - n \log Z(1)] \Big|_{n=1}$$

- at $n \sim 1$, linearized gravity eom demand: $K^\alpha = h^{ij} K_{ij}^\alpha = 0$

\longrightarrow τ shrinks to zero on an extremal surface in bulk

- evaluating Einstein action yields $S = A/4G_N$ for extremal surface

Overview:

- 1) Entanglement entropy in QFT
- 2) Primer on AdS/CFT correspondence
- 3) Introduction to Holographic Entanglement Entropy
- 4) Derivation of Holographic Entanglement Entropy
(for spherical entangling surfaces)**
- 5) Conclusions

Why is deriving Holographic Entanglement Entropy hard?

$$S_{EE} = -Tr [\rho_A \log \rho_A]$$

- a “standard” approach relies on **replica trick**, first calculating **Renyi entropy** and then taking $n \rightarrow 1$ limit

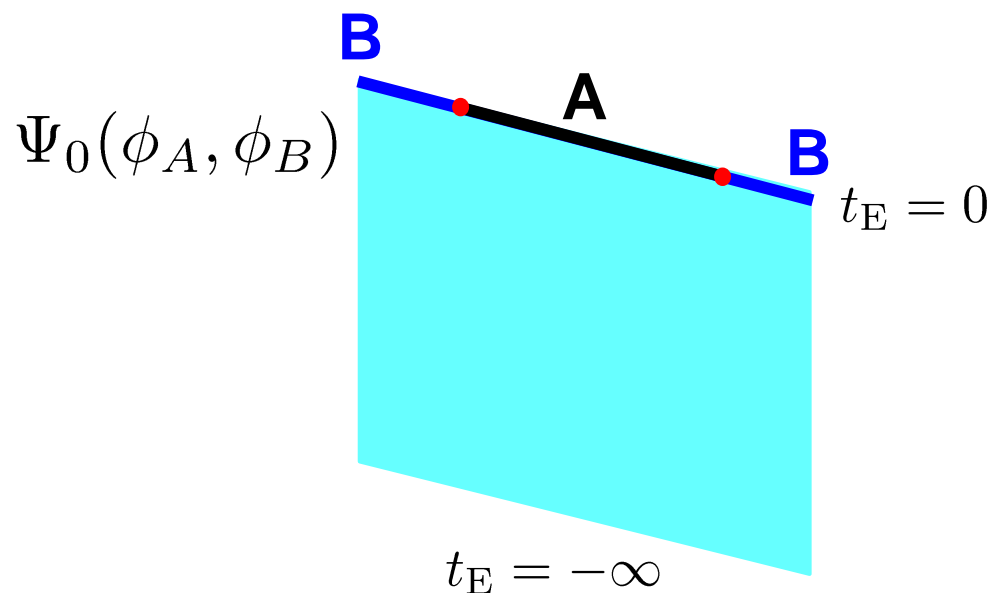
$$S_n = \frac{1}{1-n} \log Tr [\rho_A^n]$$

$$S_{EE} = \lim_{n \rightarrow 1} S_n$$

Replica trick:

0. analytically continue: $t_E = i t$

1. path integral representation
of ground state wave function



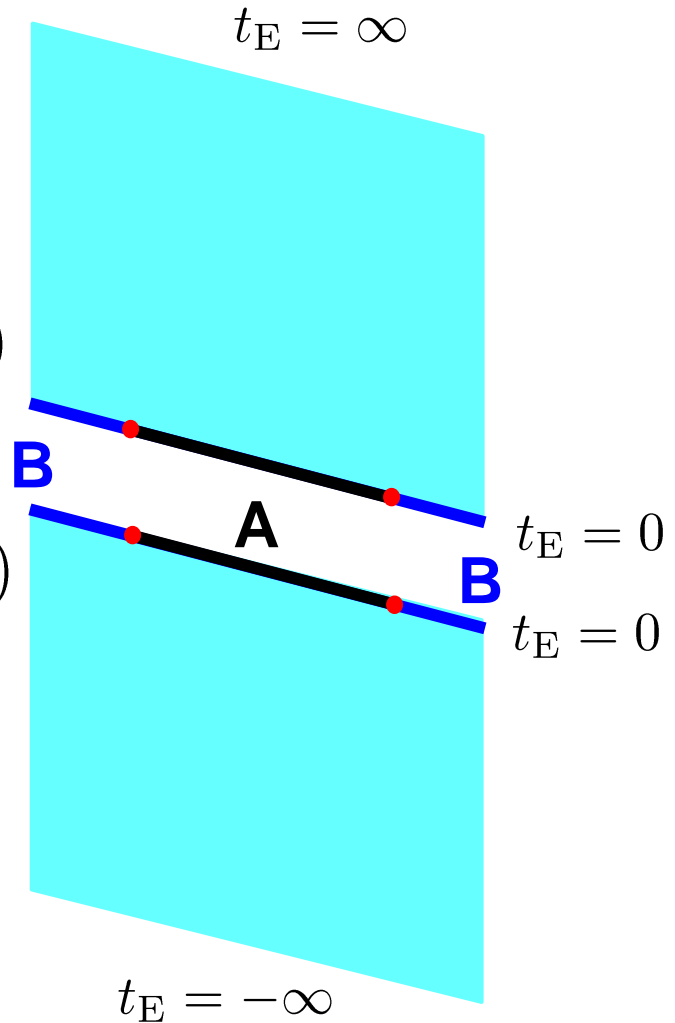
Replica trick:

0. analytically continue: $t_E = i t$

1. path integral representation
of ground state wave function

$$\Psi_0^\dagger(\phi_A, \phi_B)$$

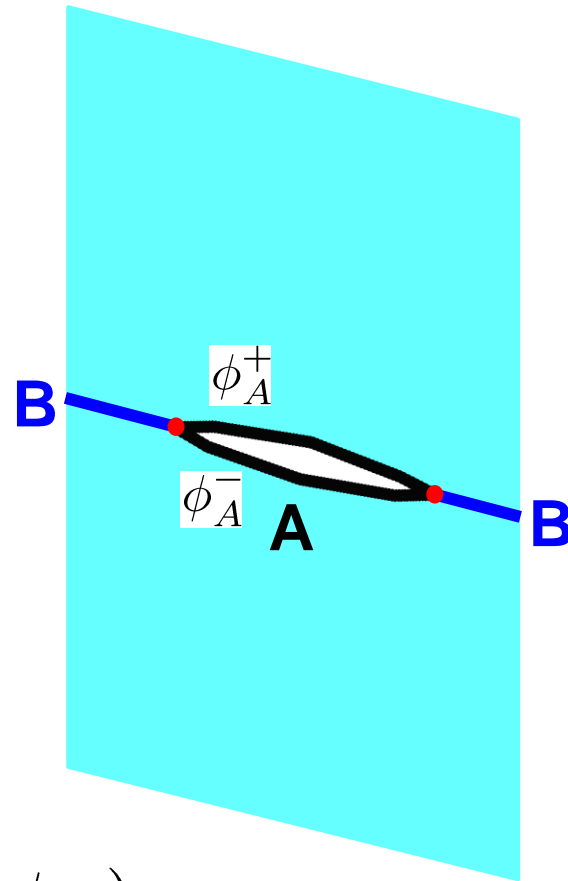
$$\Psi_0(\phi_A, \phi_B)$$



Replica trick:

0. analytically continue: $t_E = i t$
1. path integral representation of ground state wave function
2. trace over ϕ_B to construct density matrix $\rho_A(\phi_A^+, \phi_A^-)$

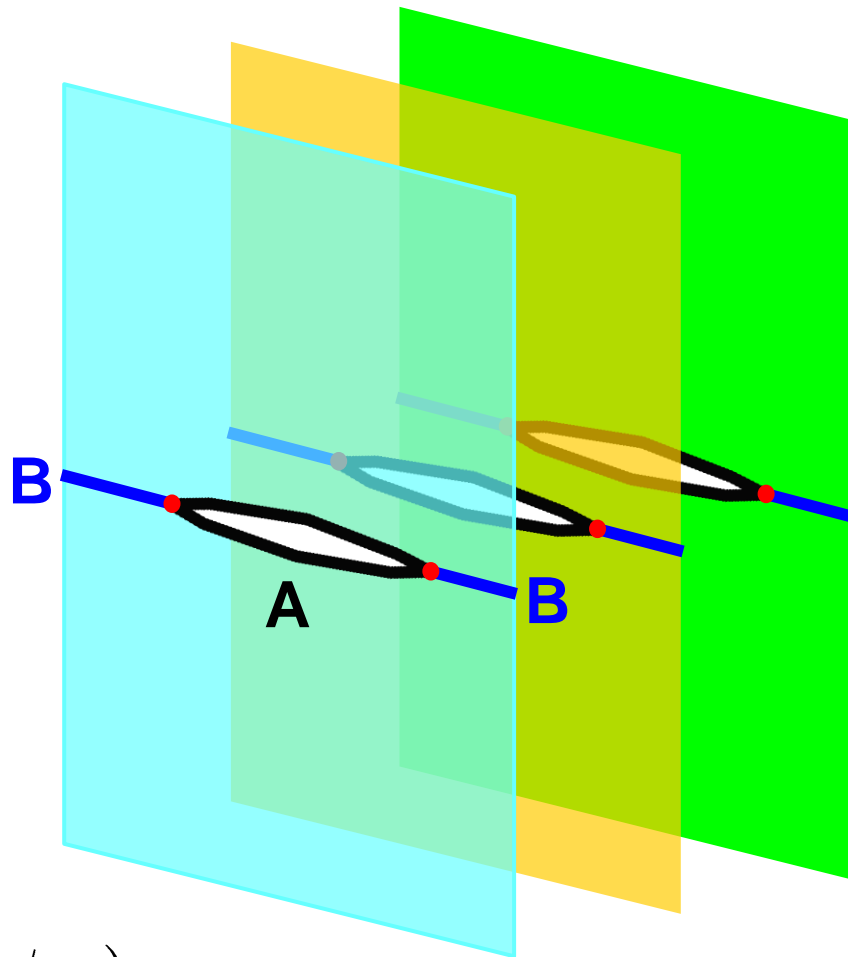
$$\begin{aligned} \rho_A(\phi_A^+, \phi_A^-) \\ = \text{Tr}_{\phi_B} \Psi^\dagger(\phi_A^+, \phi_B) \Psi(\phi_A^-, \phi_B) \end{aligned}$$



Replica trick:

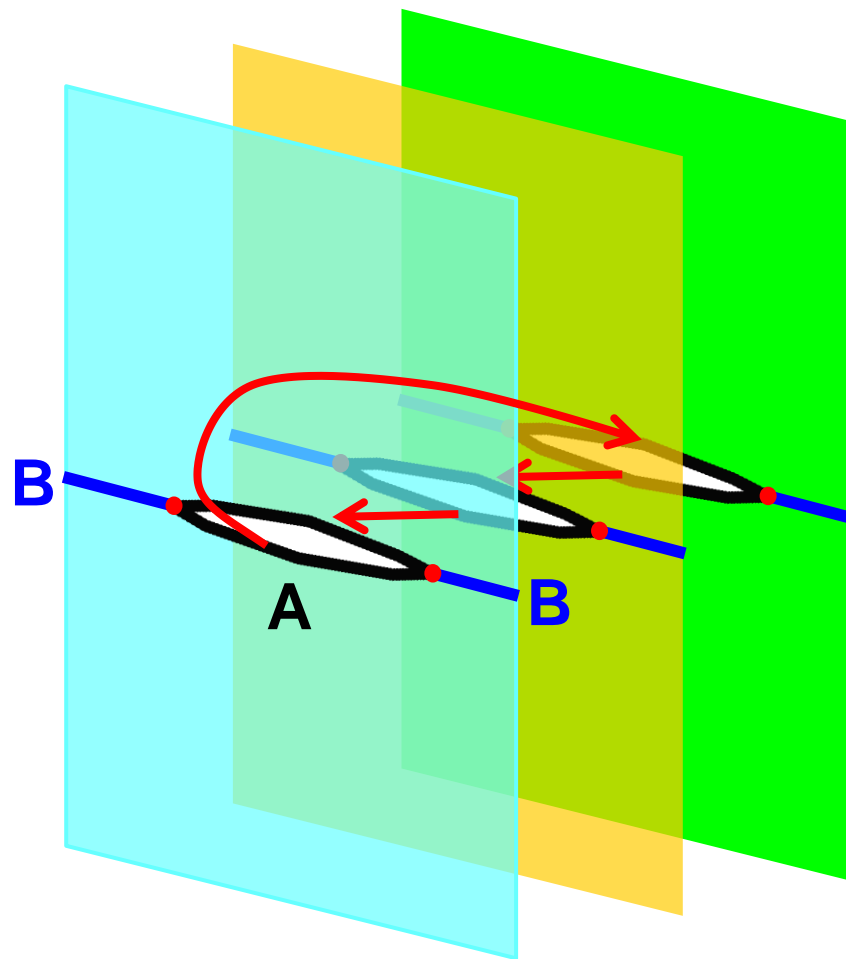
0. analytically continue: $t_E = i t$
1. path integral representation of ground state wave function
2. trace over ϕ_B to construct density matrix $\rho_A(\phi_A^+, \phi_A^-)$

$$\begin{aligned} \rho_A(\phi_A^+, \phi_A^-) &= \text{Tr}_{\phi_B} \Psi^\dagger(\phi_A^+, \phi_B) \Psi(\phi_A^-, \phi_B) \end{aligned}$$



Replica trick:

0. analytically continue: $t_E = i t$
1. path integral representation of ground state wave function
2. trace over ϕ_B to construct density matrix $\rho_A(\phi_A^+, \phi_A^-)$
3. evaluate $\text{Tr}(\rho_A^n)$



$$\text{Tr}(\rho_A^n) = \text{Tr}_{\phi_A^i} [\rho_A(\phi_A^1, \phi_A^{n-1}) \cdots \rho_A(\phi_A^3, \phi_A^2) \rho_A(\phi_A^2, \phi_A^1)]$$

Replica trick:

0. analytically continue: $t_E = i t$

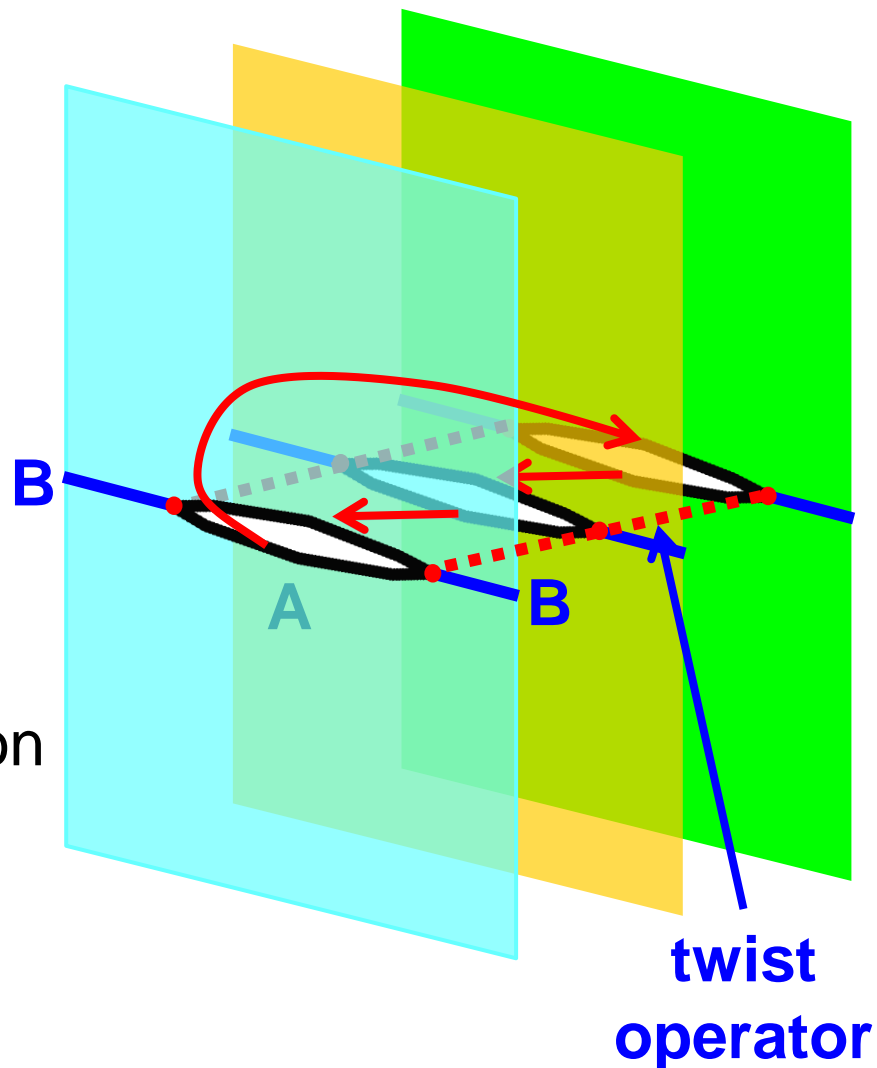
1. path integral representation of ground state wave function

2. trace over ϕ_B to construct density matrix $\rho_A(\phi_A^+, \phi_A^-)$

3. evaluate $\text{Tr}(\rho_A^n)$



evaluate euclidean partition function on n-fold cover of original space



Replica trick:

0. analytically continue: $t_E = i t$

1. path integral representation of ground state wave function

2. trace over ϕ_B to construct density matrix $\rho_A(\phi_A^+, \phi_A^-)$

3. evaluate $\text{Tr}(\rho_A^n)$

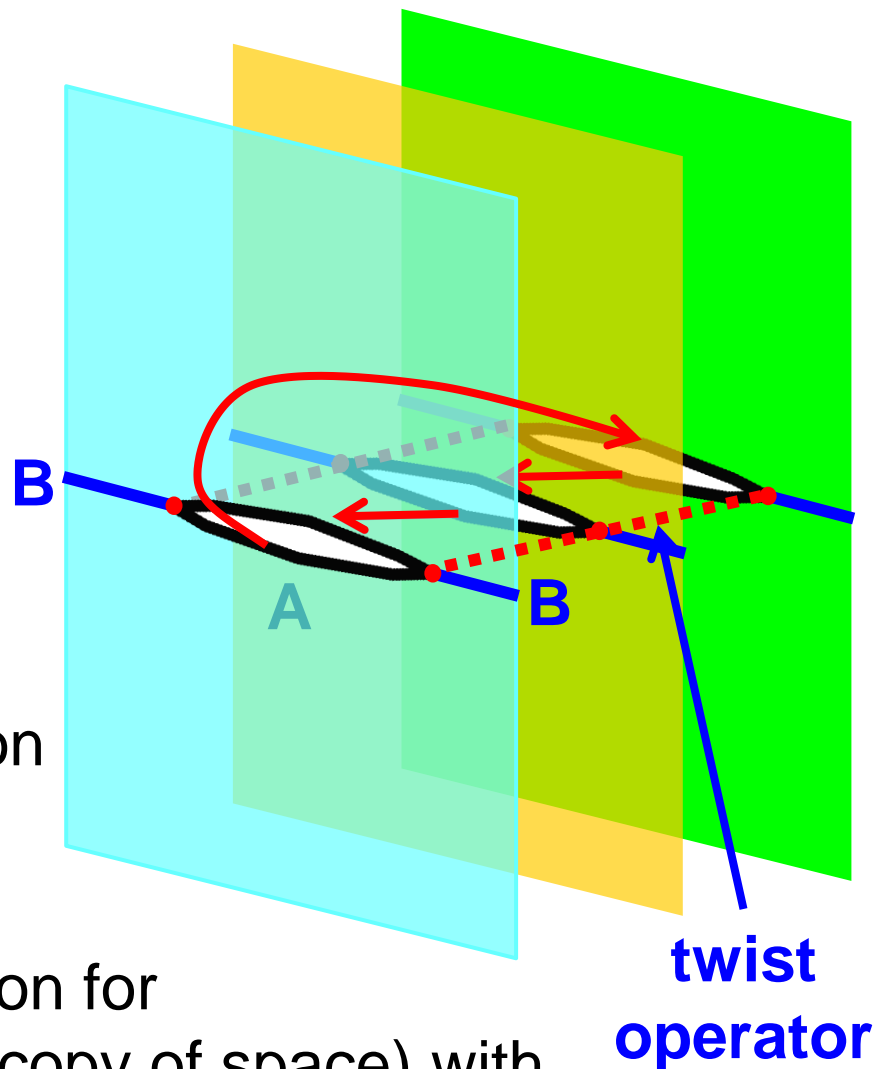


evaluate euclidean partition function on n-fold cover of original space

or

evaluate euclidean partition function for n copies of field theory (on single copy of space) with

twist operator inserted at boundary of region **A**



Why is deriving Holographic Entanglement Entropy hard?

- “standard” approach to calculate S_n relies on replica trick
- replica trick involves path integral of QFT on **singular** n-fold cover of background spacetime
- holographic slogan: “**its all geometry!**”
 - how do we deal with singularity in boundary???

Why is deriving Holographic Entanglement Entropy hard?

- “standard” approach to calculate S_n relies on replica trick
- replica trick involves path integral of QFT on **singular** n-fold cover of background spacetime
- holographic slogan: **“its all geometry!”**
 - how do we deal with singularity in boundary???
- “live with it!” → singularity extends into the bulk and it is effectively “extremized” as part of bulk gravity path integral

(Fursaev)

Why is deriving Holographic Entanglement Entropy hard?

- “standard” approach to calculate S_n relies on replica trick
- replica trick involves path integral of QFT on **singular** n-fold cover of background spacetime
- holographic slogan: **“its all geometry!”**
 - how do we deal with singularity in boundary???

(Fursaev)
- “live with it!” → singularity extends into the bulk and it is effectively “extremized” as part of bulk gravity path integral
- problem: you get the wrong answer (at $n \neq 1$)

(Headrick)
- “smooth it out!” → find smooth bulk solution for new boundary metric; use conformal symmetry to “unwrap” singularity; (particularly “simple” for $d=2$: all boundary metrics locally conformally flat, all bulk sol’s locally AdS_3)

Why is deriving Holographic Entanglement Entropy hard?

- “standard” approach to calculate S_n relies on replica trick
 - replica trick involves path integral of QFT on **singular** n-fold cover of background spacetime
 - holographic slogan: **“its all geometry!”**
 - how do we deal with singularity in boundary???

(Fursaev)
 - “live with it!” → singularity extends into the bulk and it is effectively “extremized” as part of bulk gravity path integral
 - problem: you get the wrong answer (at $n \neq 1$)
 - “smooth it out!” → find smooth bulk solution for new boundary metric; use conformal symmetry to “unwrap” singularity; (particularly “simple” for $d=2$: all boundary metrics locally conformally flat, all bulk sol’s locally AdS_3)
 - “all of the above”
- (Lewkowycz & Maldacena)

Why is deriving Holographic Entanglement Entropy hard?

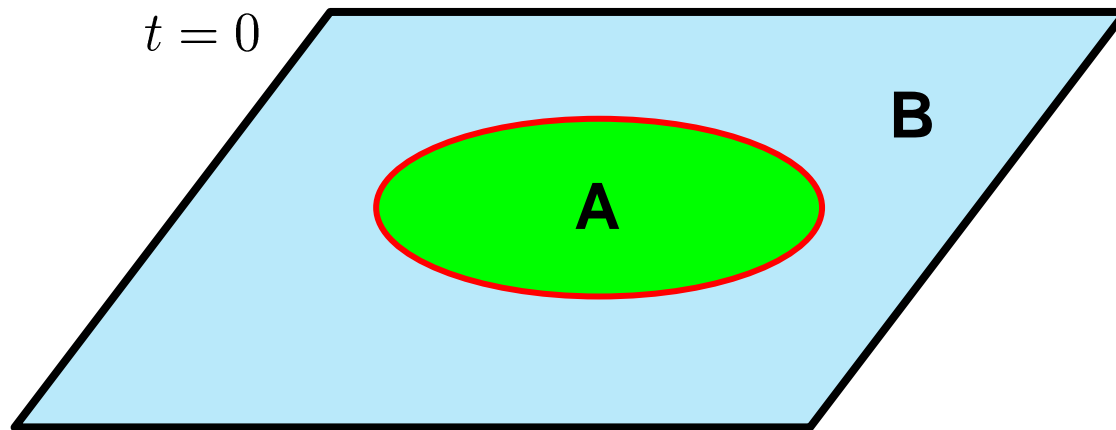
- “standard” approach to calculate S_n relies on replica trick
 - replica trick involves path integral of QFT on **singular** n-fold cover of background spacetime
 - holographic slogan: **“its all geometry!”**
 - how do we deal with singularity in boundary???

(Fursaev)
 - “live with it!” → singularity extends into the bulk and it is effectively “extremized” as part of bulk gravity path integral
 - problem: you get the wrong answer (at $n \neq 1$)
 - “smooth it out!” → find smooth bulk solution for new boundary metric; use conformal symmetry to “unwrap” singularity; (particularly “simple” for $d=2$: all boundary metrics locally conformally flat, all bulk sol’s locally AdS_3)
 - **need another calculation with simpler holographic translation***
- (Headrick)
- (Casini, Huerta & RM)
- (*realizing “smooth it out!” strategy in disguise)

Calculating Entanglement Entropy:

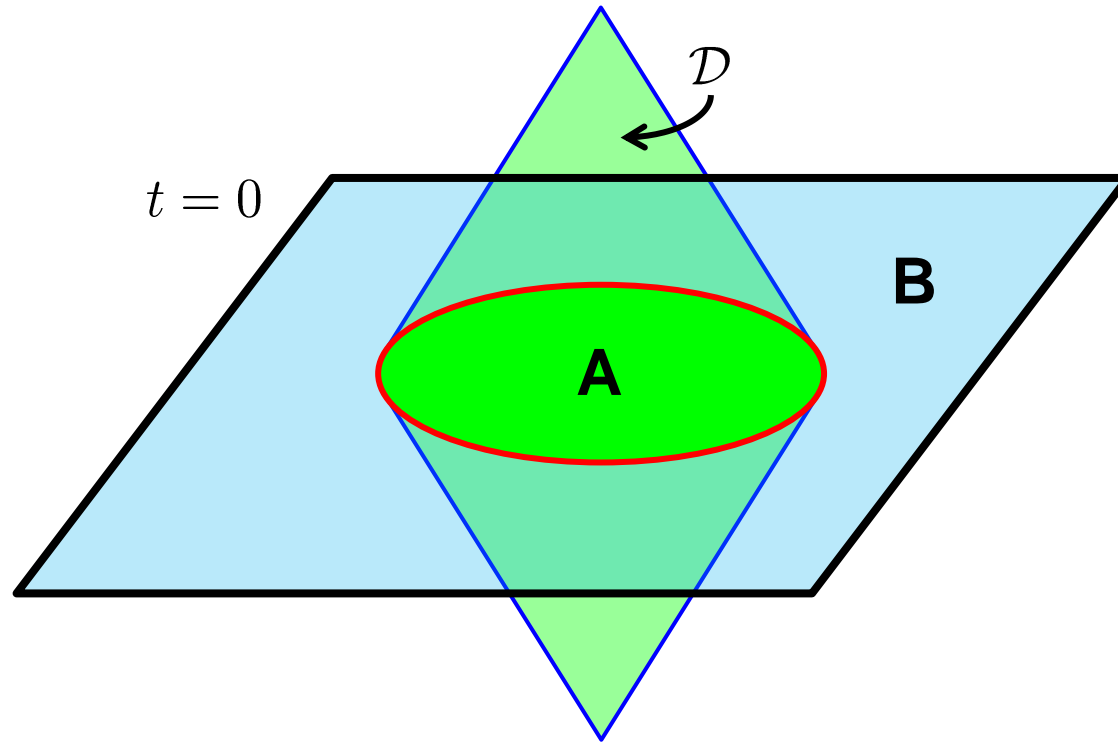
(Casini, Huerta & RCM)

- take **CFT** in d-dim. flat space and choose $\Sigma = S^{d-2}$ with radius R
→ entanglement entropy: $S_{EE} = -Tr [\rho_A \log \rho_A]$



Calculating Entanglement Entropy:

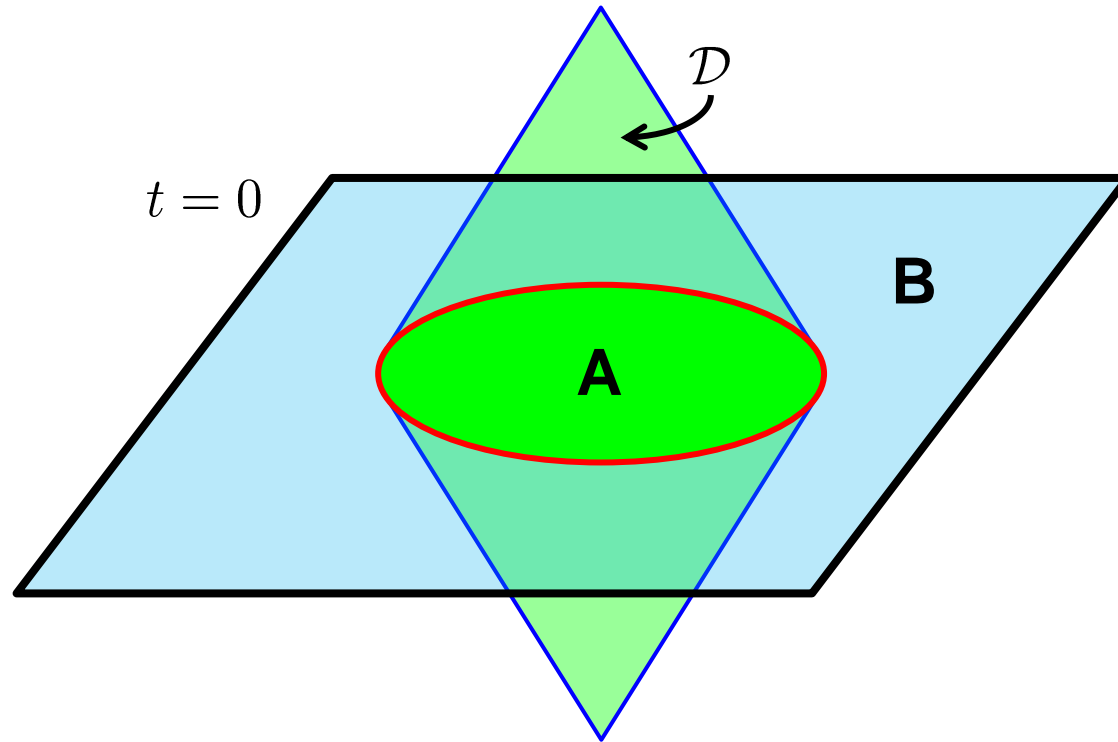
- take **CFT** in d -dim. flat space and choose $\Sigma = S^{d-2}$ with radius R
 → entanglement entropy: $S_{EE} = -Tr [\rho_A \log \rho_A]$



- density matrix ρ_A describes physics in entire causal domain \mathcal{D}

Calculating Entanglement Entropy:

- take **CFT** in d-dim. flat space and choose $\Sigma = S^{d-2}$ with radius R
 → entanglement entropy: $S_{EE} = -Tr [\rho_A \log \rho_A]$

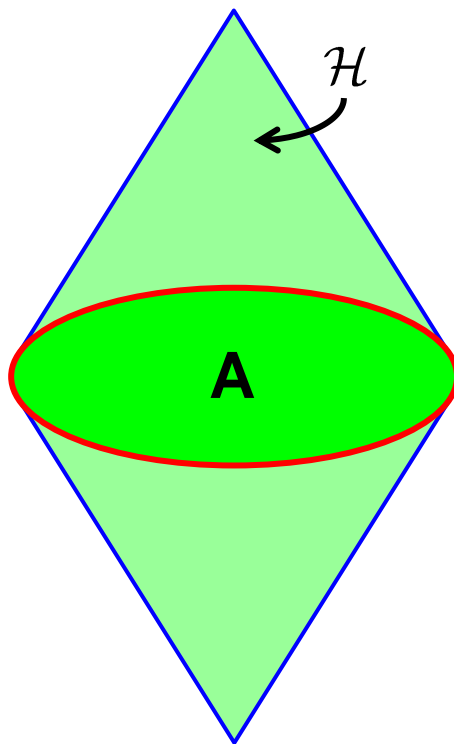


- density matrix ρ_A describes physics in entire causal domain \mathcal{D}
- conformal** mapping: $\mathcal{D} \rightarrow \mathcal{H} = R_t \times H^{d-1}$

General result for any CFT

(Casini, Huerta & RCM)

- take CFT in d-dim. flat space and choose $\Sigma = S^{d-2}$ with radius R
→ entanglement entropy: $S_{EE} = -Tr [\rho_A \log \rho_A]$



- conformal mapping: $\mathcal{D} \rightarrow \mathcal{H} = R_t \times H^{d-1}$

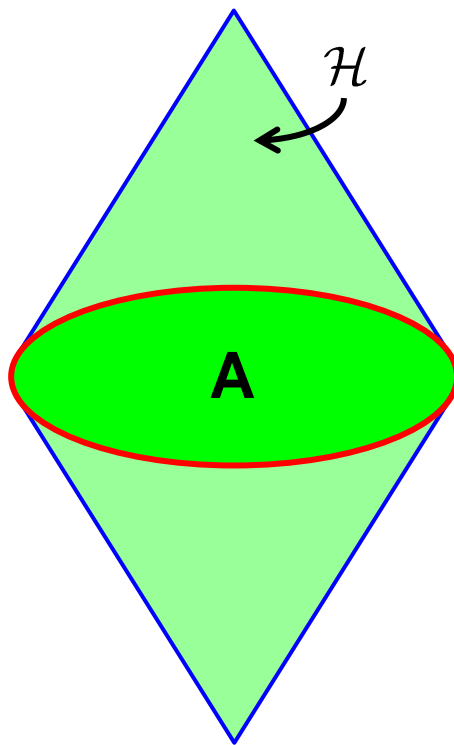
curvature scale: $1/R$

temperature: $T=1/2\pi R$!!

General result for any CFT

(Casini, Huerta & RCM)

- take CFT in d-dim. flat space and choose $\Sigma = S^{d-2}$ with radius R
→ entanglement entropy: $S_{EE} = -Tr [\rho_A \log \rho_A]$



- conformal mapping: $\mathcal{D} \rightarrow \mathcal{H} = R_t \times H^{d-1}$

curvature scale: $1/R$

temperature: $T=1/2\pi R$!!

- for CFT: $\rho_{thermal} = U \rho_A U^{-1} \longrightarrow \boxed{S_{EE} = S_{thermal}}$

General result for any CFT

(Casini, Huerta & RCM)

- take CFT in d-dim. flat space and choose S^{d-2} with radius R
 - entanglement entropy: $S_{EE} = -Tr [\rho_A \log \rho_A]$
 - by conformal mapping relate to thermal entropy on $\mathcal{H} = R \times H^{d-1}$ with $\mathcal{R} \sim 1/R^2$ and $T=1/2\pi R$

$$S_{EE} = S_{thermal}$$

General result for any CFT

(Casini, Huerta & RCM)

- take CFT in d-dim. flat space and choose S^{d-2} with radius R
 - entanglement entropy: $S_{EE} = -Tr [\rho_A \log \rho_A]$
 - by conformal mapping relate to thermal entropy on $\mathcal{H} = R \times H^{d-1}$ with $\mathcal{R} \sim 1/R^2$ and $T=1/2\pi R$

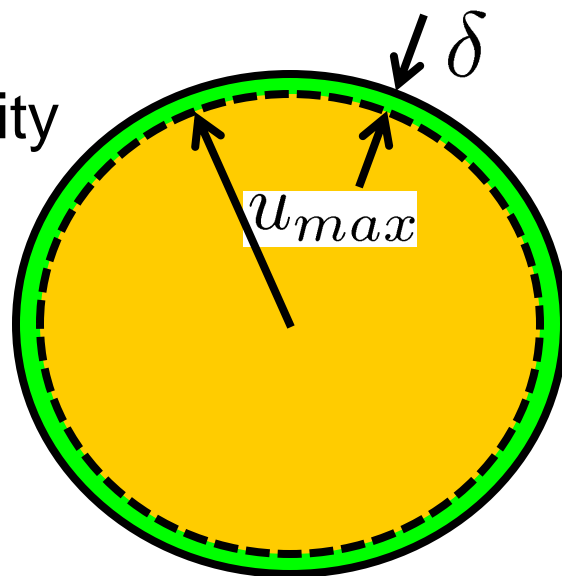
$$S_{EE} = S_{thermal}$$

- note both sides of equality are divergent

→ $S_{thermal}$ sums constant entropy density over infinite volume

- conformal map takes original UV cut-off to IR cut-off on H^{d-1}

$$u_{max} \simeq R/\delta$$



General result for any CFT

(Casini, Huerta & RCM)

- take CFT in d-dim. flat space and choose S^{d-2} with radius R
 - entanglement entropy: $S_{EE} = -Tr [\rho_A \log \rho_A]$
 - by conformal mapping relate to thermal entropy on $\mathcal{H} = R \times H^{d-1}$ with $\mathcal{R} \sim 1/R^2$ and $T=1/2\pi R$

$$S_{EE} = S_{thermal}$$

AdS/CFT correspondence:

- thermal bath in CFT = black hole in AdS

$$S_{EE} = S_{thermal} = S_{horizon}$$

- only need to find appropriate black hole

General result for any CFT

(Casini, Huerta & RCM)

- take CFT in d-dim. flat space and choose S^{d-2} with radius R
 - entanglement entropy: $S_{EE} = -Tr [\rho_A \log \rho_A]$
 - by conformal mapping relate to thermal entropy on $\mathcal{H} = R \times H^{d-1}$ with $\mathcal{R} \sim 1/R^2$ and $T=1/2\pi R$

$$S_{EE} = S_{thermal}$$

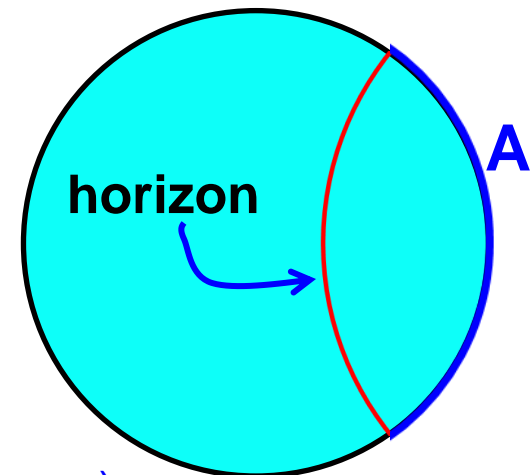
AdS/CFT correspondence:

- thermal bath in CFT = black hole in AdS

$$S_{EE} = S_{thermal} = S_{horizon}$$

- only need to find appropriate black hole
 - topological BH with hyperbolic horizon which intersects ∂A on AdS boundary

(Aminneborg et al; Emparan; Mann; . . .)



$$S_{EE} = S_{thermal} = S_{horizon}$$

- desired “black hole” is a hyperbolic foliation of empty AdS space

$$ds^2 = \frac{L^2}{z^2} (dz^2 - dt^2 + d\vec{x}^2)$$

$$S_{EE} = S_{thermal} = S_{horizon}$$

- desired “black hole” is a hyperbolic foliation of empty AdS space

$$ds^2 = \frac{L^2 d\rho^2}{(\rho^2 - L^2)} - \frac{\rho^2 - L^2}{R^2} d\tau^2 + \rho^2 d\Sigma_2^{d-1} \quad \longrightarrow \quad T = \frac{1}{2\pi R}$$

- bulk coordinate transformation implements desired conformal transformation on boundary

$$S_{EE} = S_{thermal} = S_{horizon}$$

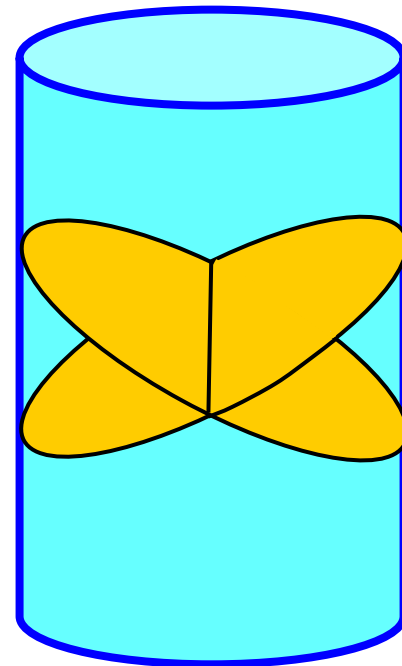
- desired “black hole” is a hyperbolic foliation of empty AdS space

$$ds^2 = \frac{L^2}{z^2} (dz^2 - dt^2 + d\vec{x}^2) d\tau^2 + \rho^2 d\Sigma_2^{d-1} \quad \longrightarrow \quad T = \frac{1}{2\pi R}$$

- bulk coordinate transformation implements desired conformal transformation on boundary

- Rindler coordinates of AdS space:

“Rinder-AdS coordinates”



$$S_{EE} = S_{thermal} = S_{horizon}$$

- desired “black hole” is a hyperbolic foliation of empty AdS space

$$ds^2 = \frac{L^2 d\rho^2}{(\rho^2 - L^2)} - \frac{\rho^2 - L^2}{R^2} d\tau^2 + \rho^2 d\Sigma_2^{d-1} \quad \longrightarrow \quad T = \frac{1}{2\pi R}$$

- apply Wald’s formula (for any gravity theory) for horizon entropy:

$$\begin{aligned} S &= -2\pi \int d^{d-1}x \sqrt{h} \frac{\partial \mathcal{L}}{\partial R^{\mu\nu}_{\rho\sigma}} \hat{\varepsilon}^{\mu\nu} \hat{\varepsilon}_{\rho\sigma} \\ &= \frac{2\pi}{\pi^{d/2}} \Gamma(d/2) a_d^* V(H^{d-1}) \end{aligned}$$

(RCM & Sinha)

$$S_{EE} = S_{thermal} = S_{horizon}$$

- desired “black hole” is a hyperbolic foliation of empty AdS space

$$ds^2 = \frac{L^2 d\rho^2}{(\rho^2 - L^2)} - \frac{\rho^2 - L^2}{R^2} d\tau^2 + \rho^2 d\Sigma_2^{d-1} \quad \longrightarrow \quad T = \frac{1}{2\pi R}$$

- apply Wald’s formula (for any gravity theory) for horizon entropy:

$$\begin{aligned} S &= -2\pi \int d^{d-1}x \sqrt{h} \frac{\partial \mathcal{L}}{\partial R^{\mu\nu}{}_{\rho\sigma}} \hat{\varepsilon}^{\mu\nu} \hat{\varepsilon}_{\rho\sigma} \\ &= \frac{2\pi}{\pi^{d/2}} \Gamma(d/2) a_d^* V(H^{d-1}) \end{aligned}$$

(RCM & Sinha)

where a_d^* contains all of the couplings from the gravity theory

$$\text{eg, } a_d^* = \frac{\pi^{d/2}}{\Gamma(d/2)} \frac{L^{d-1}}{\ell_P^{d-1}} \quad \text{for Einstein gravity}$$

$$S_{EE} = S_{thermal} = S_{horizon}$$

- desired “black hole” is a hyperbolic foliation of empty AdS space

$$ds^2 = \frac{L^2 d\rho^2}{(\rho^2 - L^2)} - \frac{\rho^2 - L^2}{R^2} d\tau^2 + \rho^2 d\Sigma_2^{d-1} \quad \longrightarrow \quad T = \frac{1}{2\pi R}$$

- apply Wald’s formula (for any gravity theory) for horizon entropy:

$$\begin{aligned} S &= -2\pi \int d^{d-1}x \sqrt{h} \frac{\partial \mathcal{L}}{\partial R^{\mu\nu}{}_{\rho\sigma}} \hat{\varepsilon}^{\mu\nu} \hat{\varepsilon}_{\rho\sigma} \\ &= \frac{2\pi}{\pi^{d/2}} \Gamma(d/2) a_d^* V(H^{d-1}) \end{aligned}$$

(RCM & Sinha)

where a_d^* = central charge for “A-type trace anomaly”

for even d

= entanglement entropy defines effective central charge

for odd d

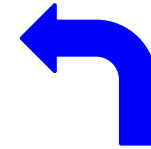
$$S_{EE} = S_{thermal} = S_{horizon}$$

- desired “black hole” is a hyperbolic foliation of empty AdS space

$$ds^2 = \frac{L^2 d\rho^2}{(\rho^2 - L^2)} - \frac{\rho^2 - L^2}{R^2} d\tau^2 + \rho^2 d\Sigma_2^{d-1} \quad \longrightarrow \quad T = \frac{1}{2\pi R}$$

- apply Wald’s formula (for any gravity theory) for horizon entropy:

$$S = \frac{2\pi}{\pi^{d/2}} \Gamma(d/2) a_d^* V(H^{d-1})$$



$$ds^2 = \frac{du^2}{1 + u^2} + u^2 d\Omega_2^{d-2}$$

$$S = a_d^* \frac{4\pi^{\frac{d-3}{2}}}{(d-2)\Gamma\left(\frac{d-1}{2}\right)} u_{max}^{d-2} + \dots$$

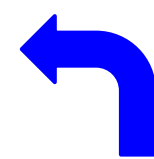
$$S_{EE} = S_{thermal} = S_{horizon}$$

- desired “black hole” is a hyperbolic foliation of empty AdS space

$$ds^2 = \frac{L^2 d\rho^2}{(\rho^2 - L^2)} - \frac{\rho^2 - L^2}{R^2} d\tau^2 + \rho^2 d\Sigma_2^{d-1} \quad \longrightarrow \quad T = \frac{1}{2\pi R}$$

- apply Wald’s formula (for any gravity theory) for horizon entropy:

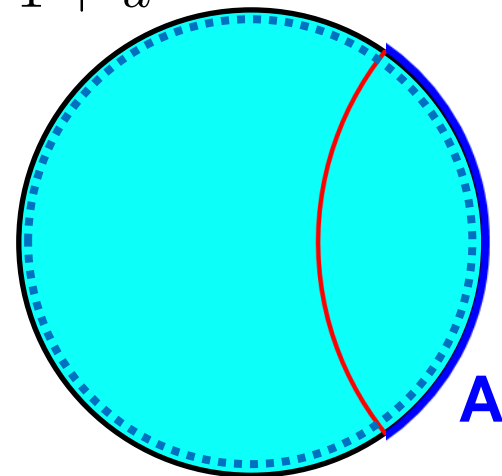
$$S = \frac{2\pi}{\pi^{d/2}} \Gamma(d/2) a_d^* V(H^{d-1})$$



intersection with standard regulator surface: $z_{min} = \delta$

$$ds^2 = \frac{du^2}{1 + u^2} + u^2 d\Omega_2^{d-2}$$

$$S = a_d^* \frac{4\pi^{\frac{d-3}{2}}}{(d-2)\Gamma(\frac{d-1}{2})} u_{max}^{d-2} + \dots$$



$$S_{EE} = S_{thermal} = S_{horizon}$$

- desired “black hole” is a hyperbolic foliation of empty AdS space

$$ds^2 = \frac{L^2 d\rho^2}{(\rho^2 - L^2)} - \frac{\rho^2 - L^2}{R^2} d\tau^2 + \rho^2 d\Sigma_2^{d-1} \quad \longrightarrow \quad T = \frac{1}{2\pi R}$$

- apply Wald’s formula (for any gravity theory) for horizon entropy:

$$S = \frac{2\pi}{\pi^{d/2}} \Gamma(d/2) a_d^* V(H^{d-1})$$

$$ds^2 = \frac{du^2}{1 + u^2} + u^2 d\Omega_2^{d-2}$$

intersection with standard regulator surface: $z_{min} = \delta$

$$S = a_d^* \frac{4\pi^{\frac{d-3}{2}}}{(d-2)\Gamma(\frac{d-1}{2})} \underbrace{\left(\frac{R}{\delta}\right)^{d-2}}_{\text{area law}} + \dots$$

“area law” for d-dimensional CFT

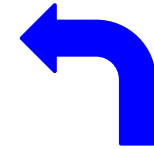
$$S_{EE} = S_{thermal} = S_{horizon}$$

- desired “black hole” is a hyperbolic foliation of empty AdS space

$$ds^2 = \frac{L^2 d\rho^2}{(\rho^2 - L^2)} - \frac{\rho^2 - L^2}{R^2} d\tau^2 + \rho^2 d\Sigma_2^{d-1} \quad \longrightarrow \quad T = \frac{1}{2\pi R}$$

- apply Wald’s formula (for any gravity theory) for horizon entropy:

$$S = \frac{2\pi}{\pi^{d/2}} \Gamma(d/2) a_d^* V(H^{d-1})$$



$$ds^2 = \frac{du^2}{1+u^2} + u^2 d\Omega_2^{d-2}$$

universal contributions:

$$S = \dots + (-)^{\frac{d}{2}-1} 4 a_d^* \log(2R/\delta) + \dots \quad \text{for even } d$$

$$\dots + (-)^{\frac{d-1}{2}} 2\pi a_d^* + \dots \quad \text{for odd } d$$

$$S_{EE} = S_{thermal} = S_{horizon}$$

universal contributions:

$$S = \dots + (-)^{\frac{d}{2}-1} 4 a_d^* \log(2R/\delta) + \dots \quad \text{for even } d$$

$$\dots + (-)^{\frac{d-1}{2}} 2\pi a_d^* + \dots \quad \text{for odd } d$$

- discussion extends to case with background: $R^{1,d-1} \rightarrow R \times S^{d-1}$
- for Einstein gravity, coincides with Ryu & Takayanagi result and horizon (bifurcation surface) coincides with R&T surface

→ no extremization procedure?!? but bifur. surface is extremal

- applies for classical bulk theories beyond Einstein gravity
- can imagine calculating “quantum” corrections (eg, Hawking rad)

“smooth it out”:

- consider **Euclidean** version of previous calculation
- conformal mapping for spherical entangling surface


→ Euclidean version gives one-to-one map: $R^d \leftrightarrow S^1 \times H^{d-1}$

→ with $\Delta\tau_E = n/T_0 = 2\pi R n$ ($n \in \mathbb{Z}$) get n-fold cover of R^d


$$S^1 \times H^{d-1} : \quad ds^2 = d\tau_E^2 + R^2 (du^2 + \sinh^2 u d\Omega_{d-2}^2)$$

$$\text{coord. transformation: } \exp(-u - i\tau_E/R) = \frac{R - r - it_E}{R + r + it_E}$$

$$\left[R^d \right]_n : \quad ds^2 = \Omega^2 \left[dt_E^2 + dr^2 + r^2 d\Omega_{d-2}^2 \right]$$

 n -fold cover

$$\Omega^2 = \frac{4R^4}{(R^2 - r^2 + t_E^2)^2 + 4r^2 t_E^2}$$



- (Euclidean) hyperbolic black holes provide smooth bulk metric

“smooth it out”:

- consider **Euclidean** version of previous calculation
- conformal mapping for spherical entangling surface

→ Euclidean version gives one-to-one map: $R^d \leftrightarrow S^1 \times H^{d-1}$

→ with $\Delta\tau_E = n/T_0 = 2\pi R n$ ($n \in \mathbb{Z}$) get n-fold cover of R^d

$$S^1 \times H^{d-1} : \quad ds^2 = d\tau_E^2 + R^2 (du^2 + \sinh^2 u d\Omega_{d-2}^2)$$

$$[R^d]_n : \quad ds^2 = \Omega^2 [dt_E^2 + dr^2 + r^2 d\Omega_{d-2}^2]$$

↑
n-fold cover

$$\Omega^2 = \frac{4R^4}{(R^2 - r^2 + t_E^2)^2 + 4r^2 t_E^2}$$

- (Euclidean) hyperbolic black holes provide smooth bulk metric for any n!!
- **special case**: symmetry emerges in certain conformal frame

Conclusions:

- AdS/CFT correspondence (gauge/gravity duality) has proven an excellent tool to study strongly coupled gauge theories
- holographic entanglement entropy is part of an interesting dialogue has opened between string theorists and physicists in a variety of fields (eg, condensed matter, nuclear physics, . . .)
- potential to learn lessons about issues in boundary theory
eg, readily calculate Renyi entropies for wide class of theories in higher dimensions
- potential to learn lessons about issues in bulk gravity theory
eg, holographic entanglement entropy may give new insight into quantum gravity or emergent spacetime

Lots to explore!