

# Dynamical Systems and Uncertainty

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# Co-workers and Web resources

Istvan Szunyogh, Brian Hunt, Edward Ott,  
Eugenia Kalnay, Jim Yorke  
and many others!

**Thanks to:** NSF, NASA, ASU

**Papers, preprints, and codes:**

<http://www.weatherchaos.umd.edu>

<http://math.la.asu.edu/~eric>

# Principal papers

Preprints: [www.weatherchaos.umd.edu](http://www.weatherchaos.umd.edu)

- Initial papers:
- E. Ott *et al.*, *Tellus A* **56** (2004), 415–428
  - I. Szunyogh *et al.*, *Tellus A* **57** (2005), 528–545

Refined mathematical implementation: B. R. Hunt, E. K., I. Szunyogh, *Physica D* **230** (2007) 112–126

Results with real data: I. Szunyogh, E.K. *et al.*, *Tellus A* **60** (2008) 113–130

# Outline of the lecture

- Relevant theory of ordinary differential equations
- A little bit about chaos
- Forecast ensembles and local low dimensionality

# Picard's existence and uniqueness theorem

- Suppose  $\mathbf{f}(\mathbf{x}, t)$  is **Lipschitz continuous** in a neighborhood  $N$  of  $(\mathbf{x}_0, t_0)$ , i.e.,

$$\|\mathbf{f}(\mathbf{x}, t) - \mathbf{f}(\mathbf{y}, t)\| \leq L \|\mathbf{x} - \mathbf{y}\|$$

for some constant  $L$  whenever  $\mathbf{x}, \mathbf{y} \in N$

- Then the initial value problem

$$\mathbf{x}' = \mathbf{f}(\mathbf{x}, t) \quad \text{with} \quad \mathbf{x}(t_0) = \mathbf{x}_0$$

has a unique solution in an interval around  $t_0$  (the size of which depends on  $N$  and  $\mathbf{f}$ )

- Implication:

**Perfect initial data  $\implies$  Perfect predictability**

# Gronwall's inequality

- Given  $\mathbf{f}$ ,  $N$ , and  $L$  as before, and suppose that  $\hat{\mathbf{x}}(t_0) = \hat{\mathbf{x}}_0$  approximates  $\mathbf{x}(t_0) = \mathbf{x}_0$
- Then

$$\|\mathbf{x}(t) - \hat{\mathbf{x}}(t)\| \leq \|\mathbf{x}_0 - \hat{\mathbf{x}}_0\| e^{L(t-t_0)}$$

- This is the best estimate that we can expect in general
- Example:

$$x' = Lx \quad \text{with} \quad x(0) = x_0 \quad \text{and} \quad \hat{x}(0) = \hat{x}_0$$

Then

$$|x(t) - \hat{x}(t)| = |x_0 - \hat{x}_0| e^{Lt}$$

# A hint of the difficulties

- **Uncertainties in initial conditions may amplify exponentially in time!**
- The details are highly equation dependent
- Example:  $x' = -Lx$  has the same Lipschitz constant, but

$$|x(t) - \hat{x}(t)| = |x_0 - \hat{x}_0| e^{-Lt} \rightarrow 0 \quad \text{as} \quad t \rightarrow \infty$$

- Under what circumstances do uncertainties grow?

# Simple case: Linear systems with constant coefficients

- Suppose  $\mathbf{x}' = \mathbf{A}\mathbf{x}$  where  $\mathbf{A} \in \mathbf{R}^{n \times n}$  has  $n$  distinct real eigenvalues
- The initial condition  $\mathbf{x}(0) = \mathbf{x}_0$  yields the solution

$$\mathbf{x}(t) = \mathbf{x}_0 e^{\mathbf{A}t} = c_1 e^{\lambda_1 t} \mathbf{v}_1 + \cdots + c_n e^{\lambda_n t} \mathbf{v}_n$$

where  $[\mathbf{x}_0]_V = (c_1, \dots, c_n)^T$  in the basis of eigenvectors

- Analogous results for repeated and complex eigenvalues



# Net result: Linear systems with constant coefficients

- Errors in initial conditions in the system  $\mathbf{x}' = \mathbf{A}\mathbf{x}$  grow exponentially with time whenever  $\mathbf{A}$  has a positive eigenvalue (or an eigenvalue with positive real part)
- This is a global result

# Harder case: Nonlinear systems

- **Local result:** Suppose  $\mathbf{x}_0$  is a fixed point for  $\mathbf{x}' = \mathbf{f}(\mathbf{x})$
- $\mathbf{x}_0$  is **hyperbolic** if the eigenvalues of the Jacobian matrix  $\mathbf{A} = \mathbf{Df}(\mathbf{x}_0)$  are all nonzero (or have nonzero real part)
- **Hartman-Grobman theorem:** There exists a change of coordinates that maps solutions of  $\mathbf{x}' = \mathbf{f}(\mathbf{x})$  onto solutions of the linear system  $\mathbf{x}' = \mathbf{Ax}$  in a neighborhood of  $\mathbf{x}_0$  whenever  $\mathbf{x}_0$  is hyperbolic

# Basic classification of hyperbolic fixed points

**Sink:** All eigenvalues negative (or negative real part)

**Saddle:** Some eigenvalues negative and some positive  
(or some negative and some positive real parts)

**Source:** All eigenvalues positive (or positive real part)

- Sinks are **stable**, i.e., insensitive to small initial errors.
- Saddles and sources are **unstable** (sensitive)

## Example: The damped nonlinear pendulum

Assume linear friction:

$$x'' + kx' + \sin x = 0 \quad \text{with} \quad k > 0$$

Define  $x_1 = \text{position}(= x)$  and  $x_2 = \text{velocity}(= x')$  The equivalent first-order system is

$$\begin{aligned}x_1' &= x_2 \\x_2' &= -kx_2 - \sin x_1\end{aligned}$$

There are two fixed points:

$$\begin{pmatrix} x_1 \\ x_2 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix} \quad \text{and} \quad \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} = \begin{pmatrix} \pi \\ 0 \end{pmatrix}$$

# Fixed point analysis

The linearized equation about each fixed point  $\mathbf{p}$  is

$$\mathbf{x}' = \mathbf{Df}(\mathbf{p}) \mathbf{x} = \begin{pmatrix} 0 & 1 \\ -\cos x_1 & -k \end{pmatrix} \mathbf{x}$$

At  $\mathbf{p} = (0, 0)$ :  $\mathbf{Df}(0, 0) = \begin{pmatrix} 0 & 1 \\ -1 & -k \end{pmatrix}$  with eigenvalues

$$\lambda_{\pm} = \frac{-k \pm \sqrt{k^2 - 4}}{2}$$

so  $\lambda_{\pm} < 0$  (or  $\operatorname{Re} \lambda_{\pm} < 0$ ). So  $(0, 0)$  is a **sink** (stable)

# Fixed point analysis, 2

- At  $\mathbf{p} = (\pi, 0)$ :  $\mathbf{Df}(\pi, 0) = \begin{pmatrix} 0 & 1 \\ +1 & -k \end{pmatrix}$  with eigenvalues

$$\lambda_{\pm} = \frac{-k \pm \sqrt{k^2 + 4}}{2}$$

- $\lambda_- < 0 < \lambda_+$ . Hence  $(\pi, 0)$  is a **saddle** (sensitive)
- Initial perturbations grow exponentially (at least initially)

# The geometry of uncertainty

- Suppose our knowledge of the initial condition  $\mathbf{x}_0$  is a “circle” of uncertainty (i.e., the underlying pdf is circularly symmetric and centered about  $\mathbf{x}_0$ )
- How does a dynamical system propagate the uncertainty?

# Linear example: $\mathbf{x} \mapsto \mathbf{A}\mathbf{x}$

- Basic formula: **matrix**  $\times$  **circle** = **ellipse**
- Key idea: The **singular value decomposition**

$$\mathbf{A}_{m \times n} = \mathbf{U}_{m \times n} \mathbf{S}_{n \times n} \mathbf{V}_{n \times n}^T$$

- $\mathbf{S} = \text{diag}(s_1, s_2, \dots, s_n)$  gives the **singular values**, which are the square roots of the eigenvalues of  $\mathbf{A}^T \mathbf{A}$
- By convention,  $s_1 \geq s_2 \geq \dots \geq s_n \geq 0$
- If  $C$  is the unit circle, then  $s_i$  is the length of the  $i$ th axis of  $\mathbf{A}C$



## Example: the damped nonlinear pendulum with $k = 2$

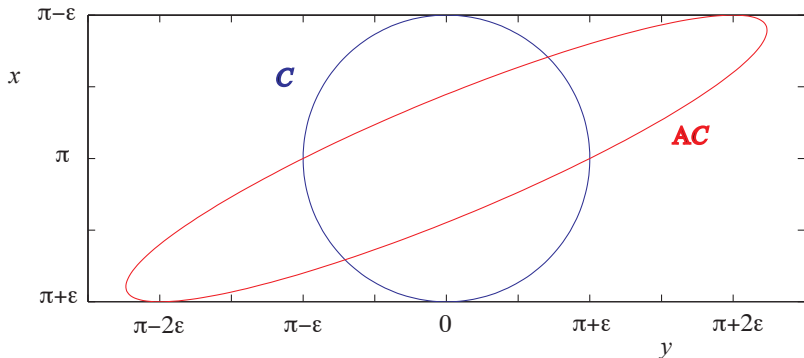
- Equation:  $x'' + 2x' + \sin x = 0$
- Consider the unstable fixed point at  $(\pi, 0)$
- The previous analysis shows that the eigenvalues are

$$\lambda_{\pm} = \frac{-k \pm \sqrt{k^2 + 4}}{2} = -1 \pm \sqrt{2}$$

- The Jacobian matrix is

$$\begin{pmatrix} 0 & 1 \\ -1 & -2 \end{pmatrix}$$

# Example: the damped nonlinear pendulum with $k = 2$



$$s_1 = 1 + \sqrt{2} \approx 2.4 \quad s_2 = \sqrt{2} - 1 \approx 0.4$$

# Forced, damped, nonlinear systems

- When damped nonlinear systems are forced strongly enough, they often become **chaotic**
- **In a chaotic process, every point is a sensitive point**
- Uncertainties in the initial condition of a chaotic process make it hard to predict—even if the process is deterministic

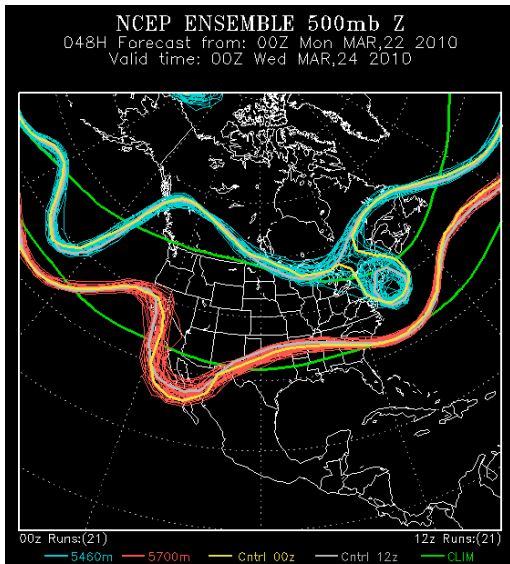
# Ensemble forecasting

- How does a nonlinear model propagate a “circle” of uncertainty?
- One procedure: Given  $\mathbf{x}' = \mathbf{f}(\mathbf{x})$ , integrate the variational equations  $\mathbf{U}' = \mathbf{Df}(\mathbf{x}) \mathbf{U}$
- **Not simple** to do if  $\mathbf{f}$  is big and complicated
- Simpler procedure: Integrate an **ensemble** of statistically equivalent initial conditions to approximate the uncertainty

## Ensemble forecasting, 2

- Ensemble forecasting is a Monte Carlo method: all state vectors are treated as random variables
- Take an initial condition that represents the “best guess” (maximum likelihood estimate) of the current atmospheric state
- Choose nearby initial conditions (relative to some norm) that sample the estimated error covariance of the best guess
- **Spaghetti plots** are one visualization method

# Spaghetti plot of a typical 72-hour forecast



# Key points

- The weather is a chaotic dynamical process
- Forecast uncertainty grows exponentially over short time scales...
- ...and varies considerably in time and space
- Lorenz's estimate: **The uncertainty in the global atmospheric state vector roughly doubles every 48 hours**
- Places an upper bound on the predictability of the weather:  $\sim 2$  weeks

# The role of data assimilation

- Without periodic corrections, the forecasts produced by a weather model would be no better than climatology
- Data assimilation combines model forecasts and atmospheric observations into an updated “best guess” of the global atmospheric state vector (and possibly its uncertainty)
- Some mathematical questions:
  - Find a useful representation of the forecast uncertainty
  - Use the available observational data efficiently
  - Estimate systematic errors (biases) in observations and forecast models
  - Where might extra observations be targeted to best advantage?



# Local low dimensionality

- Most geophysical models have millions of variables
- **But locally they often act like relatively low dimensional dynamical systems**
- A medium-resolution weather model has  $\sim 3000$  variables in a typical  $1000 \times 1000 \text{ km}^2$  synoptic region ( $\sim$ Texas, France)
- We can represent a forecast ensemble over such a region as a  $3000 \times k$  matrix  $\mathbf{X}_F$
- The singular values and vectors of  $\mathbf{X}_F$  can provide excellent low-rank approximations

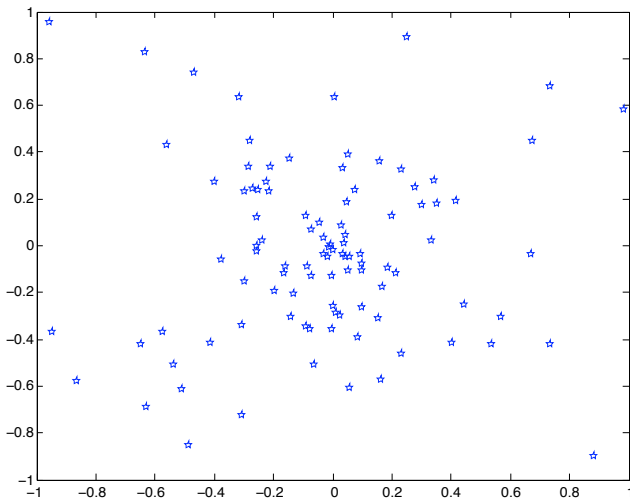
# The ensemble dimension

- The **ensemble dimension** (E-dimension) of an  $n \times k$  matrix is

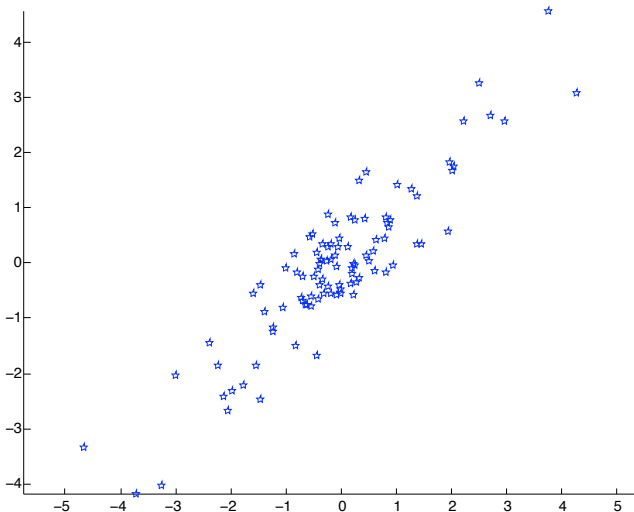
$$E \equiv \frac{(s_1 + s_2 + \cdots + s_k)^2}{s_1^2 + s_2^2 + \cdots + s_k^2}$$

- Measures the eccentricity of the “ellipse” of forecast uncertainty

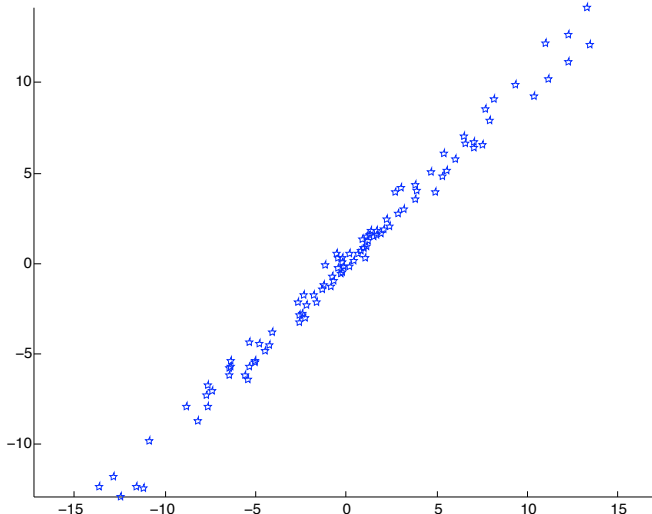
Example:  $s_1 = 3.78$ ,  $s_2 = 3.60$ ,  $E_{\text{dim}} = 1.99$



Example:  $s_1 = 19.24$ ,  $s_2 = 4.35$ ,  $E_{\text{dim}} = 1.43$



Example:  $s_1 = 83.65$ ,  $s_2 = 4.33$ ,  $E_{\text{dim}} = 1.10$



# Conclusions

- Linear dynamical systems can be analyzed globally
- Nonlinear dynamical systems allow local analyses, but global ones are hard
- Weather (or at least a typical numerical weather model) is chaotic
- Ensembles are a Monte Carlo method for quantifying forecast uncertainty
- Local low dimensionality of chaotic models can be exploited for data assimilation