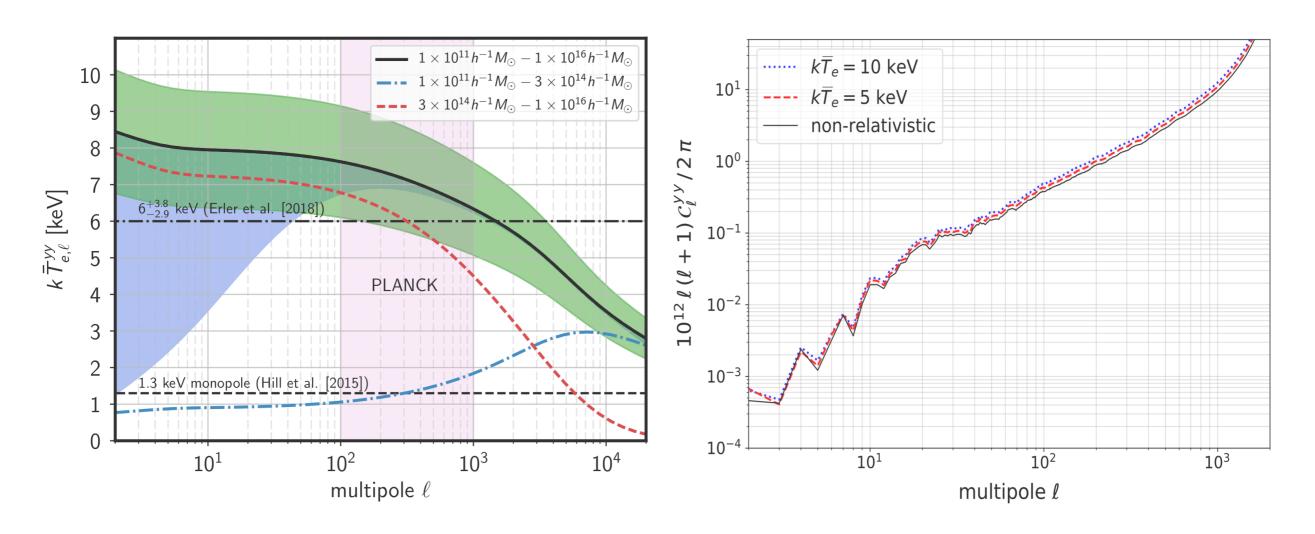
Can we neglect relativistic temperature corrections in the Planck SZ analysis?



Remazeilles, Bolliet, Rotti & JC, MNRAS, 483, 3459, 2019 Rotti, Bolliet, Remazeilles & JC, in preparation

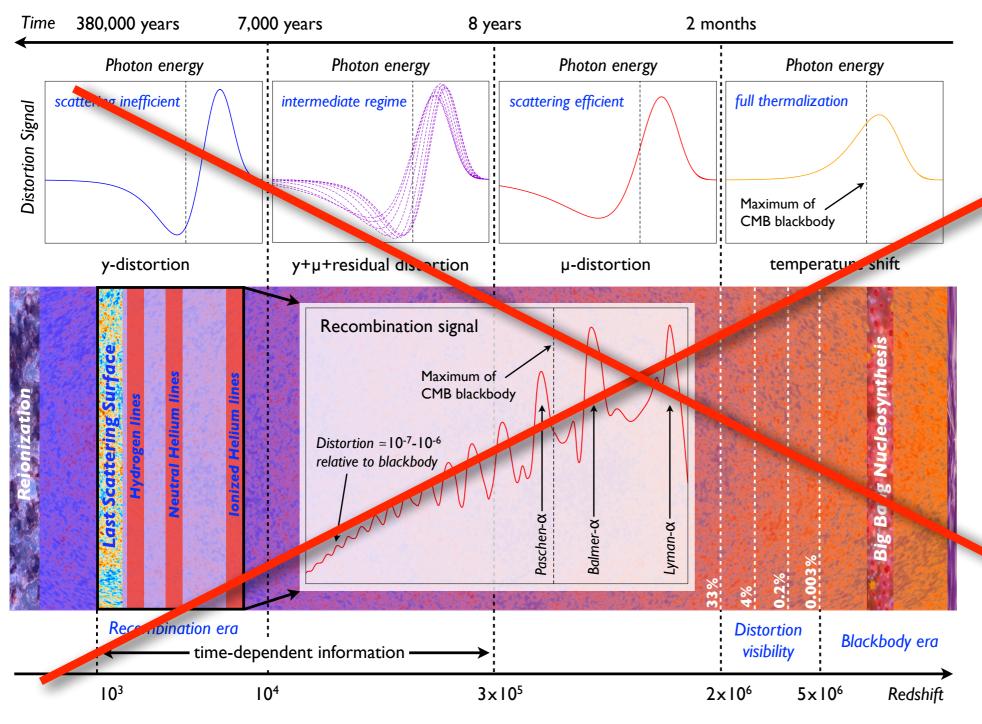


Jens Chluba

Cosmology - The Next Decade ICTS, Bangalore, January 22nd - 25th, 2019



Uniqueness of CMB Spectral Distortion Science



Guaranteed distortion signals in ACDM

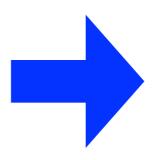
New tests of inflation and particle/dark matter physics

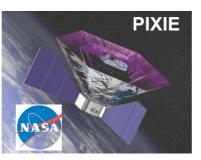
Signals from the reionization and recombination eras

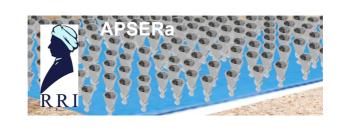
Huge discovery potential

Complementarity and synergy with CMB anisotropy studies

Chluba & Sunyaev, MNRAS, 419, 2012 Chluba et al., MNRAS, 425, 2012 Silk & Chluba, Science, 2014 Chluba, MNRAS, 2016

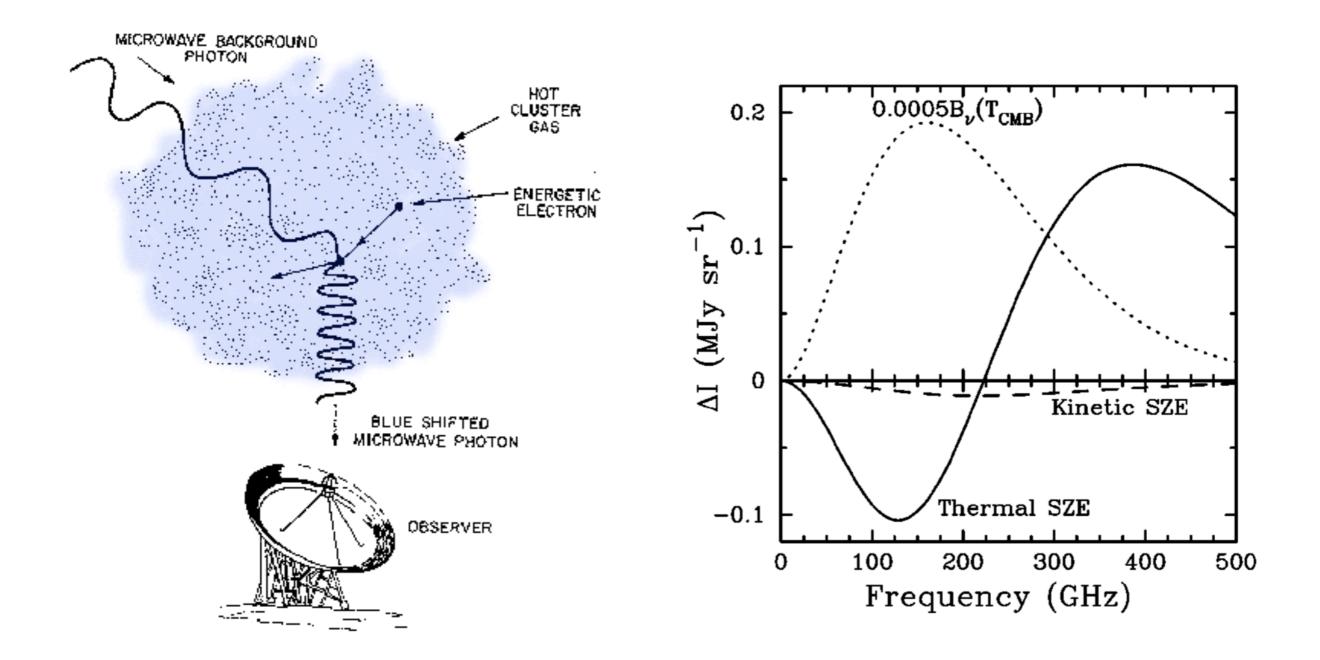






PRISTINE
COSMO
CMB-Bharat

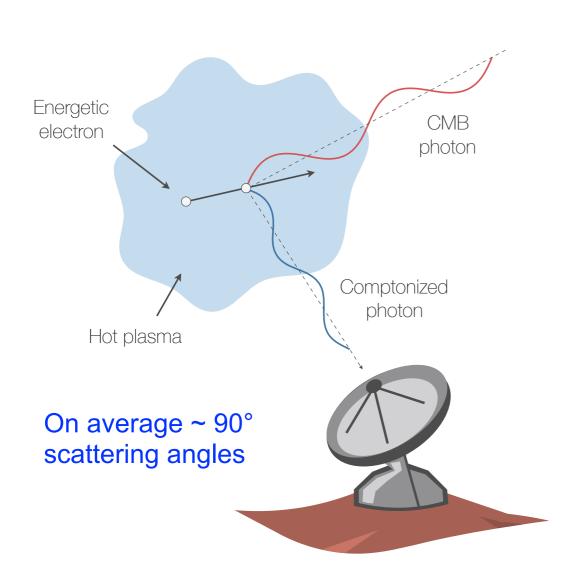
Classical SZ effects

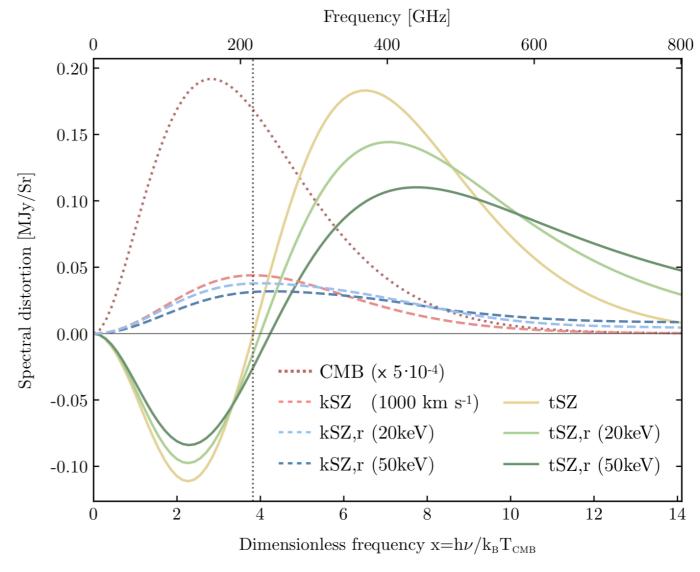


- SZ clusters are a great cosmological probe
- Many years of developments since its first prediction by Zeldovich & Sunyaev, 1969

Sunyaev & Zeldovich, 1980 Rephaeli, 1995 Birkinshaw, 1999 Carlstrom, Holder & Reese, 2002

New Comprehensive Review of SZ effects





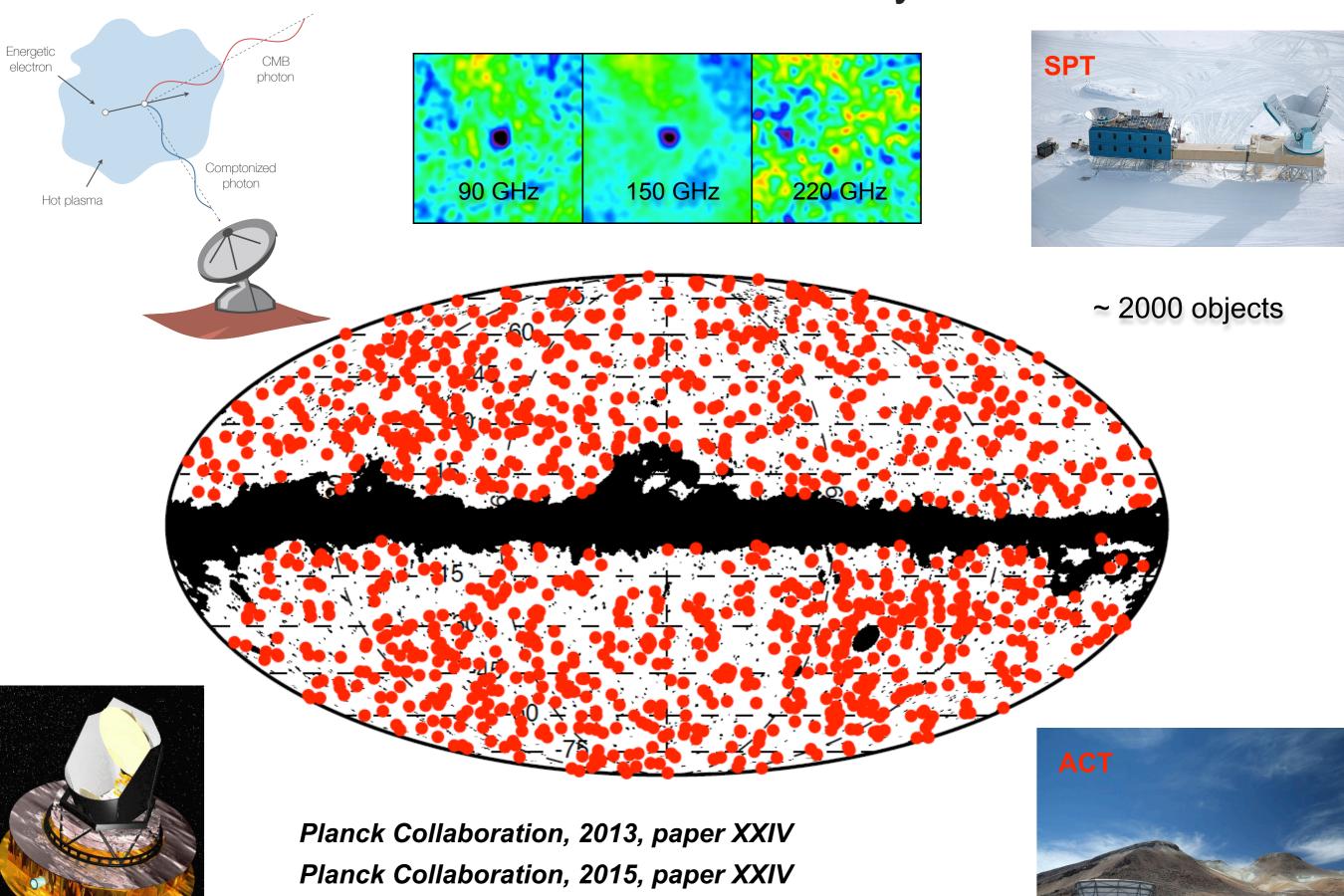
Astrophysics with the Spatially and Spectrally Resolved Sunyaev-Zeldovich Effects

A Millimetre/Submillimetre Probe of the Warm and Hot Universe

Tony Mroczkowski · Daisuke Nagai · Kaustuv
Basu · Jens Chluba · Jack Sayers · Rémi
Adam · Eugene Churazov · Abigail Crites ·
Luca Di Mascolo · Dominique Eckert · Juan
Macias-Perez · Frédéric Mayet · Laurence
Perotto · Etienne Pointecouteau · Charles
Romero · Florian Ruppin · Evan Scannapieco ·
John ZuHone

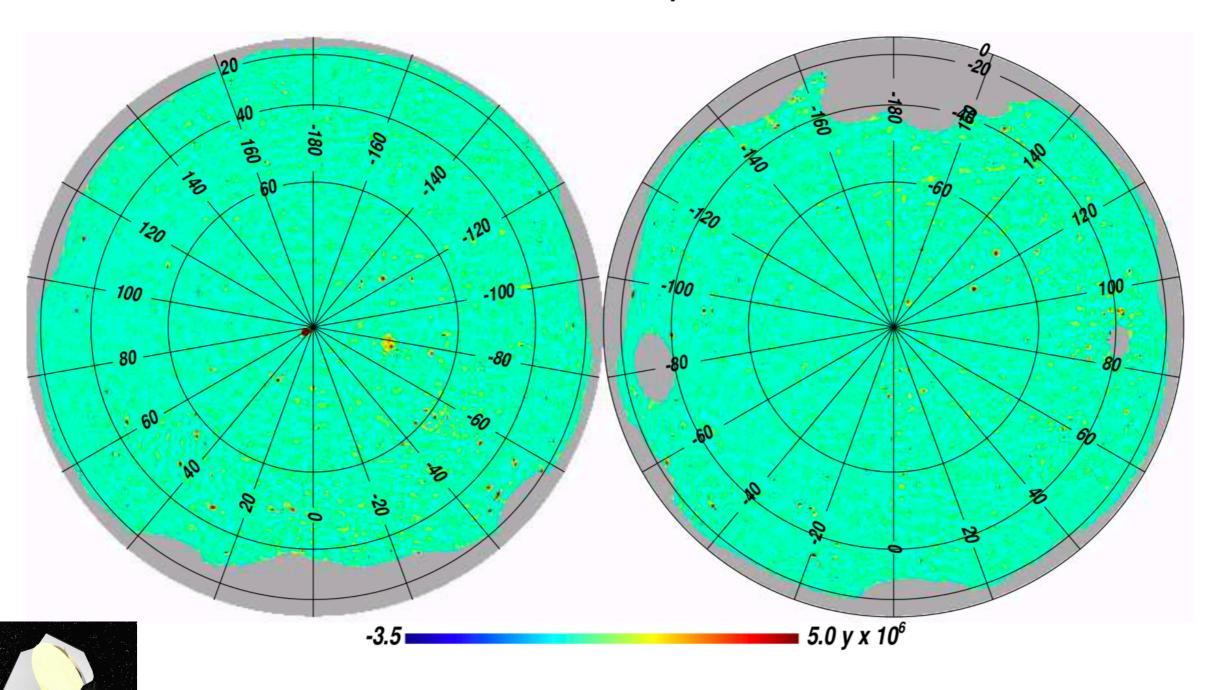
- Highlights high-resolution and high-sensitivity SZ
- Illuminates new directions
- Connection to simulations

Thermal SZ effect is now routinely observed!



All-sky Compton y-map from Planck

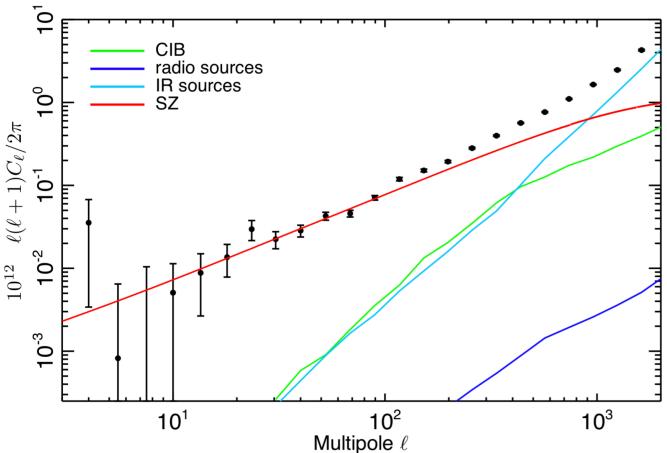
NILC tSZ map



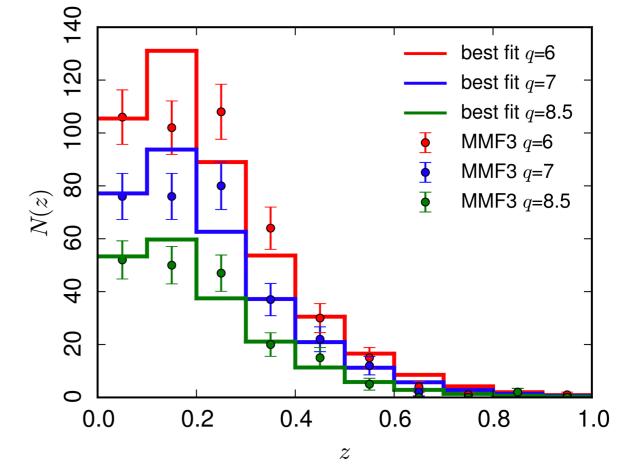
Planck Collaboration, 2013, paper XXI Planck Collaboration, 2015, paper XXII Map was produced by Mathieu Remazeilles

Planck SZ analysis



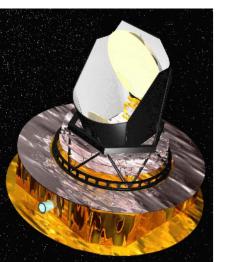


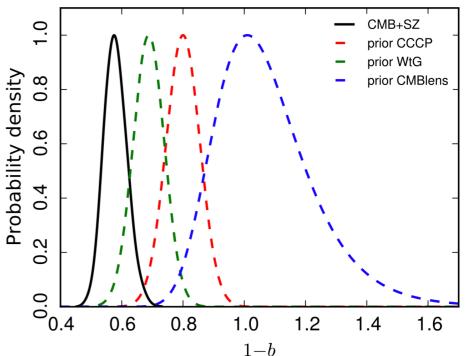
Number counts

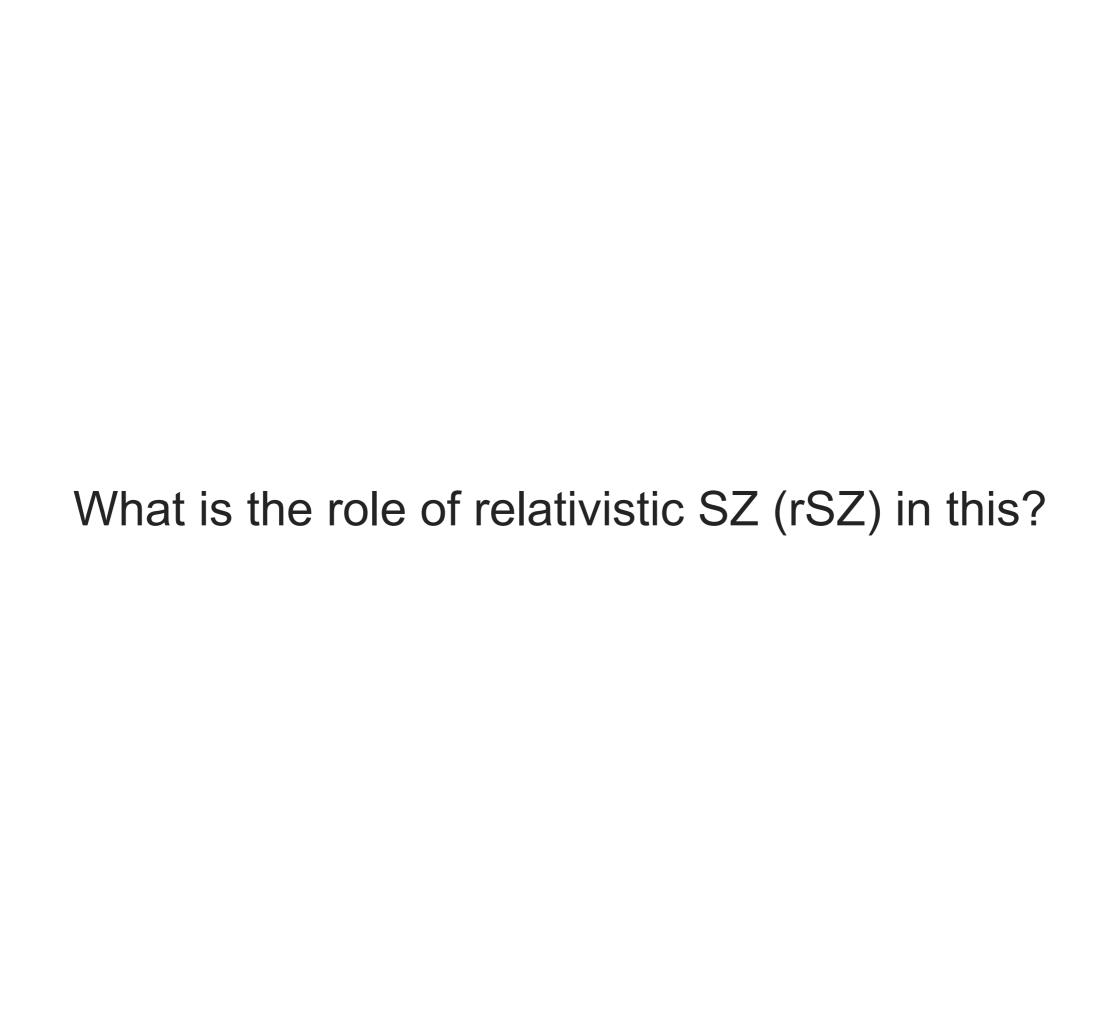


- SZ results on σ₈ in tension with CMB only result
- Hydrostatic mass bias
- Dependence on combination of data and modeling details

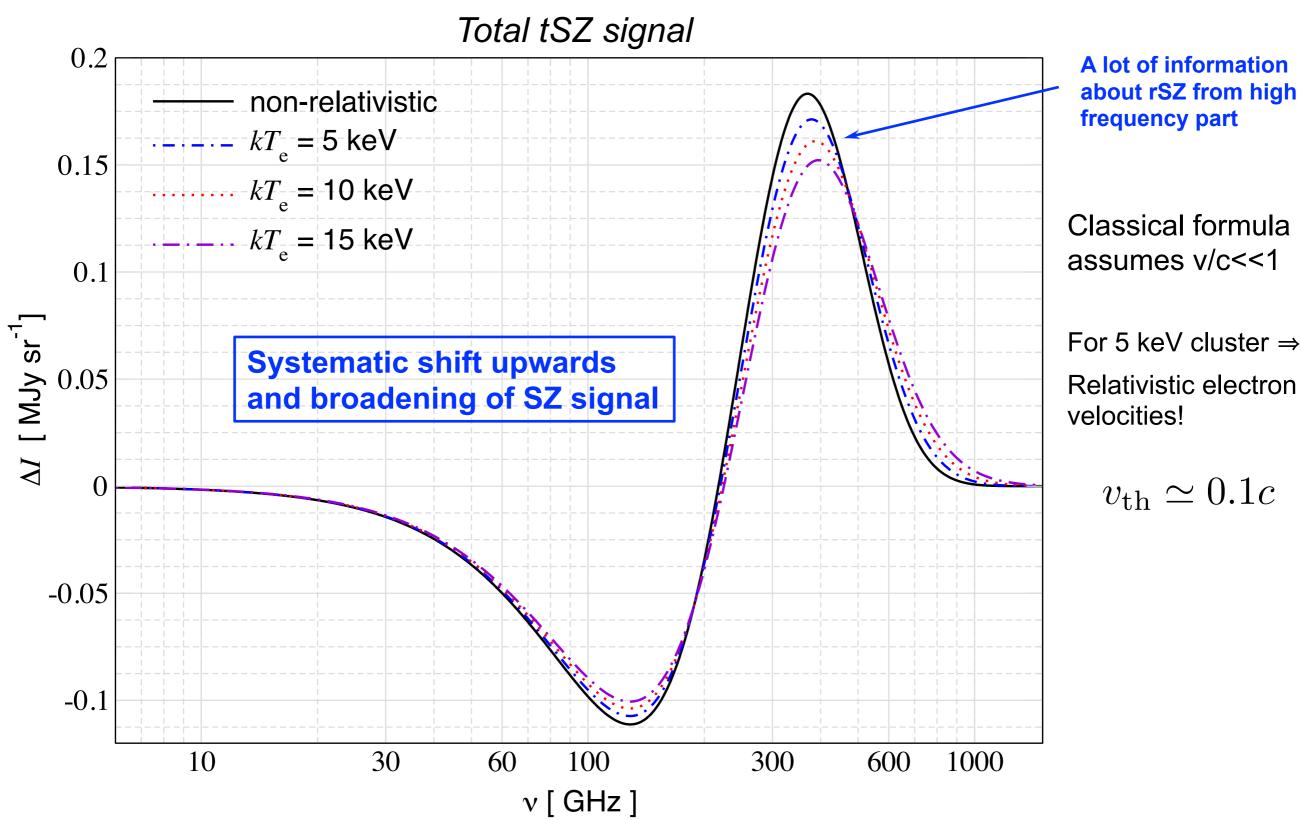
Planck Collaboration, 2015, paper XXIV Planck Collaboration, 2015, paper XXII







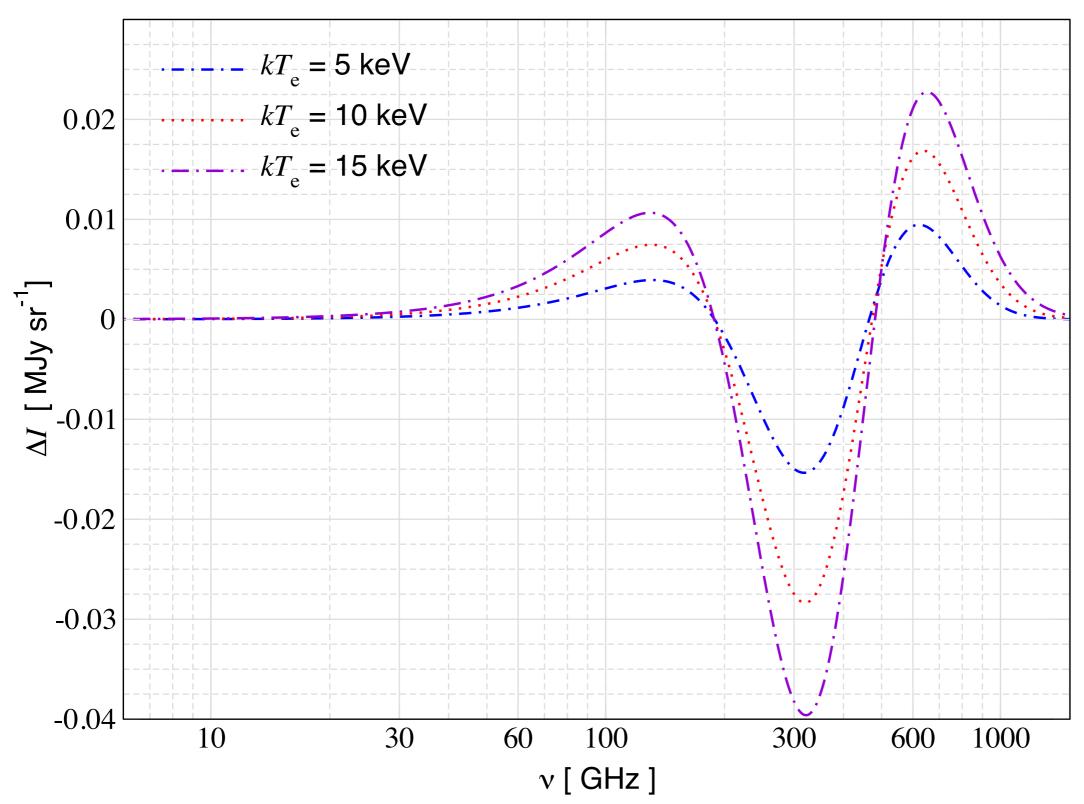
Effect of relativistic temperature corrections



 $y = 10^{-4}$

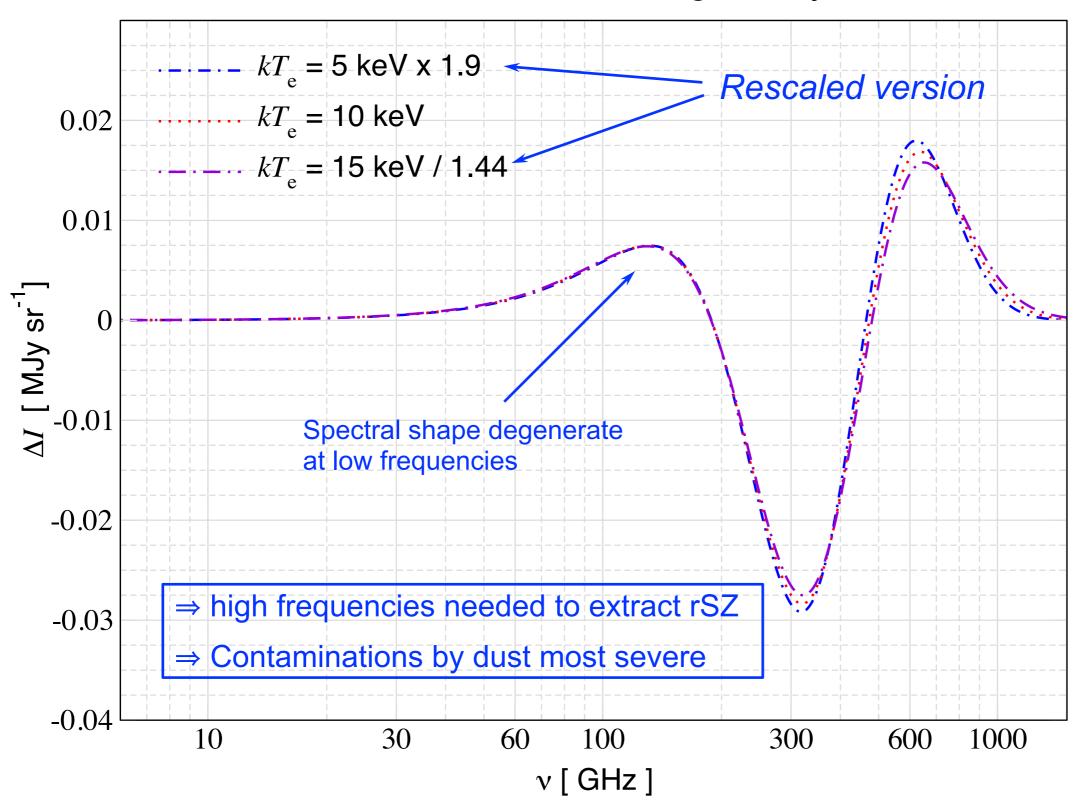
High frequencies are crucial for rSZ!

Relativistic correction signal only



High frequencies are crucial for rSZ!

Relativistic correction signal only



High-frequency spectrum has to be computed carefully...

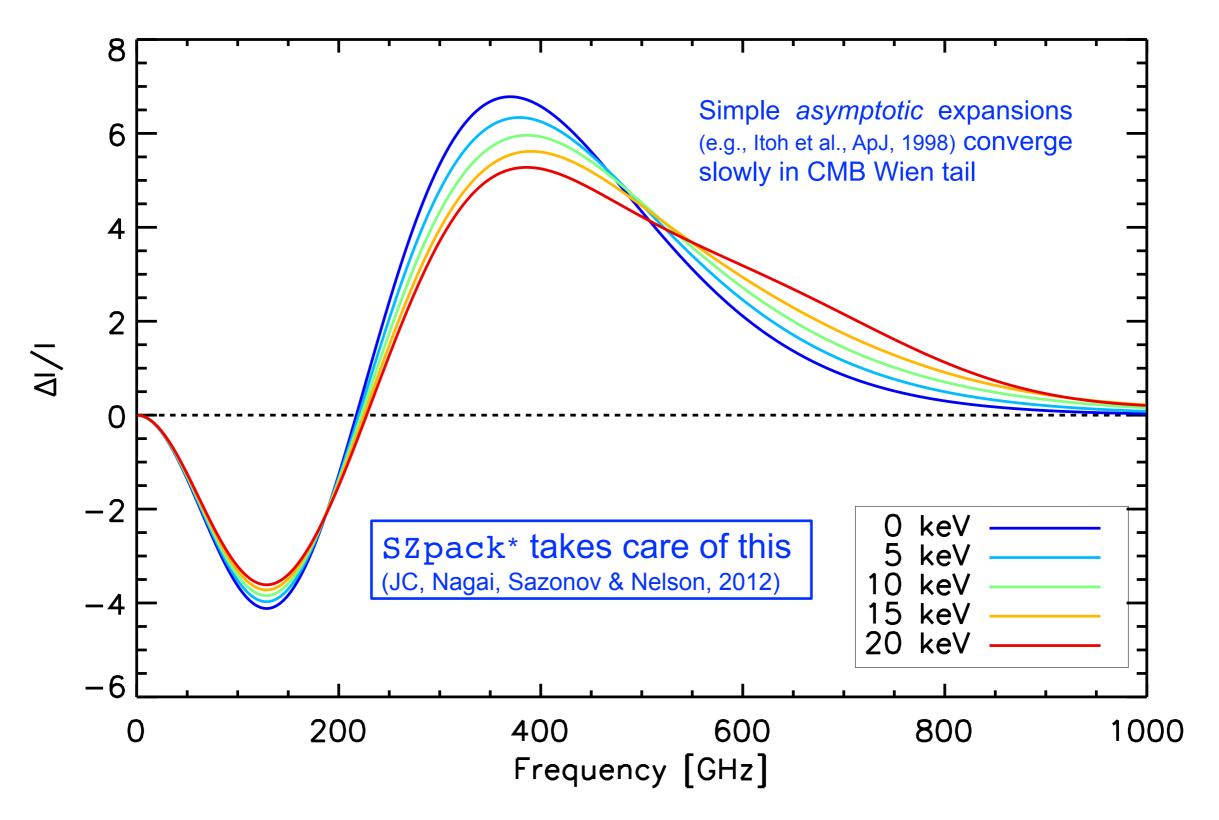
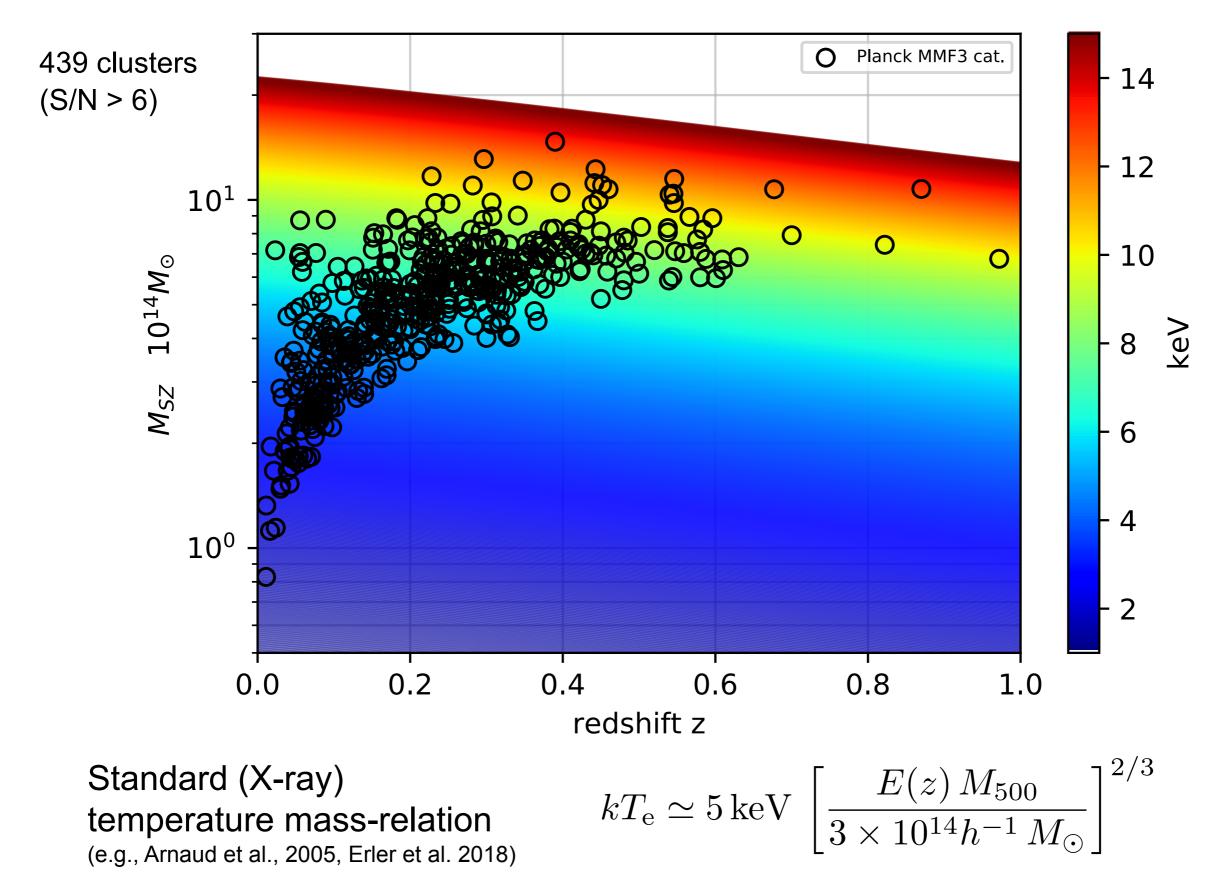
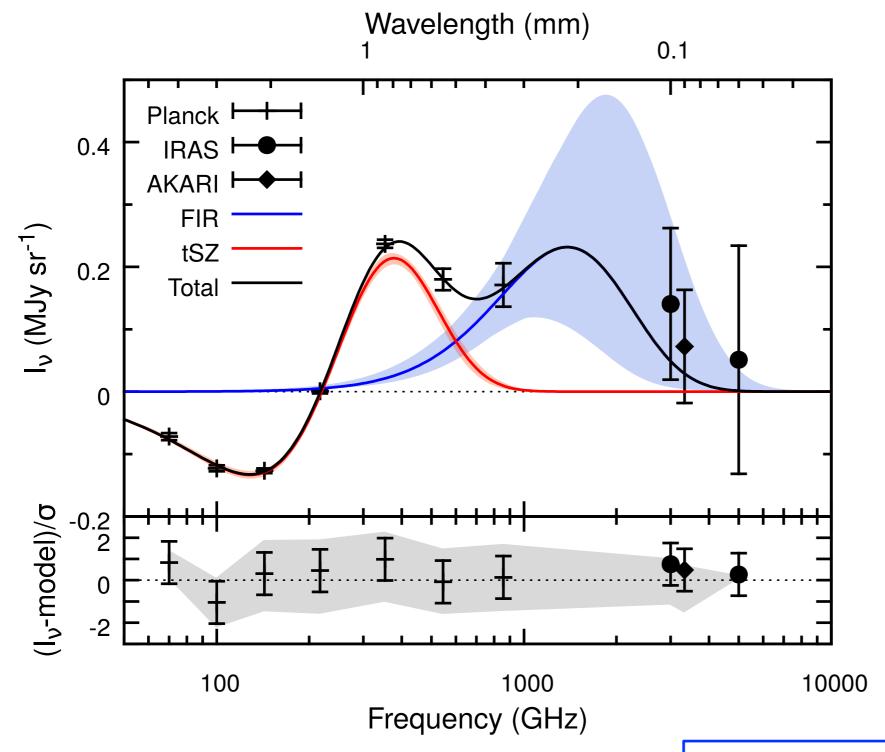


Figure from Hurier, 2016, ArXiv:1701.09020

Clusters seen by Planck are pretty hot!



Stacked Planck tSZ signal + foregrounds



- Matched filter approach
- Combination of data
- 772 clusters (PSZ2)

$$y_0 = 1.24^{+0.04}_{-0.04} \times 10^{-4}$$

$$k_B T_{SZ} = 4.36^{+2.13}_{-1.95} \text{ keV}$$

- rSZ at ~2.2 σ level
- In tension with Hurier, A&A, 2016 (claimed ~ 5 σ detection)
- 100 hottest clusters:

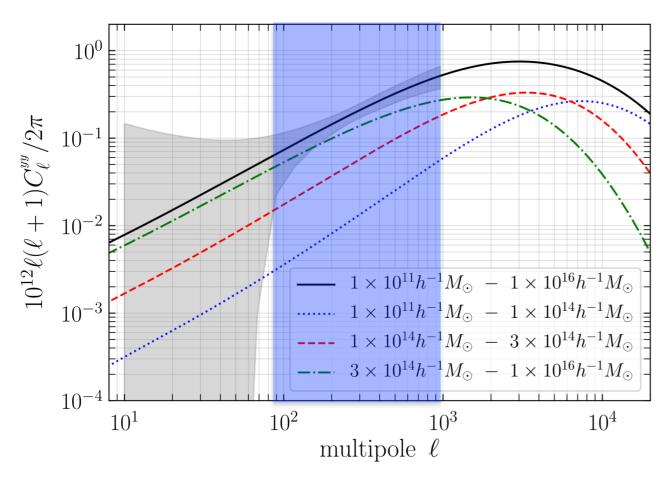
$$k_B T_{SZ} = 5.96^{+3.78}_{-2.93} \text{ keV}$$

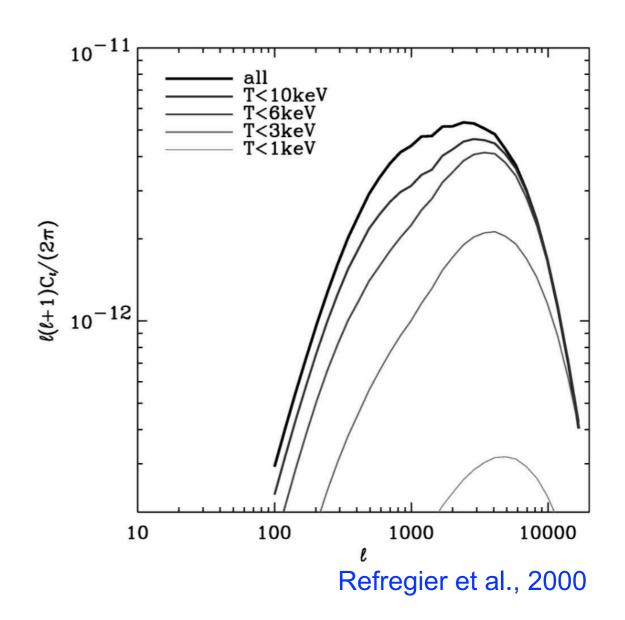
Erler, Basu, JC & Bertoldi, 2018, ArXiv:1709.01187

⇒ typical y-weighted temperature of ~ 4-7 keV quite reasonable

Theoretical SZ power spectrum computations

Computation with CLASS-SZ (Bolliet et al., 2017)





From Komatsu & Seljak, 2002:

$$C_{l} = \int_{0}^{z_{\text{max}}} dz \frac{dV}{dz} \int_{M_{\text{min}}}^{M_{\text{max}}} dM \underbrace{\frac{dn(M, z)}{dM}}_{\substack{\text{halo mass} \\ \text{function}}} \underbrace{|\tilde{y}_{l}(M, z)|^{2}}_{\substack{\text{cluster pressure} \\ \text{profile}}}$$

→ temperature of systems
 contributing to the multipole
 range relevant to *Planck*'s
 C_I analysis seems > 5 keV

How is the *y*-map obtained?

SZ Compton-y signal reconstruction: ILC

$$d(v, \overrightarrow{\theta}) = Y_0(v) \quad y(\overrightarrow{\theta}) + N(v, \overrightarrow{\theta})$$
Planck frequency signal foregrounds of interest + noise

• ILC = weighted linear combination of frequency maps:

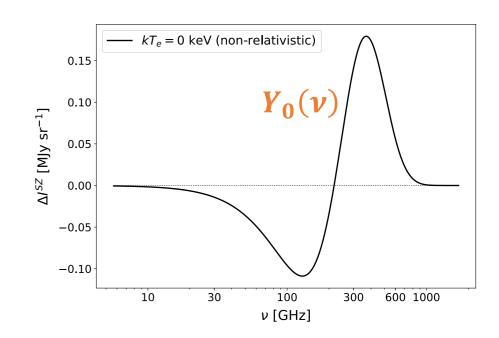
$$\widehat{y}(\overrightarrow{\theta}) = \sum_{\nu} w(\nu) d(\nu, \overrightarrow{\theta})$$

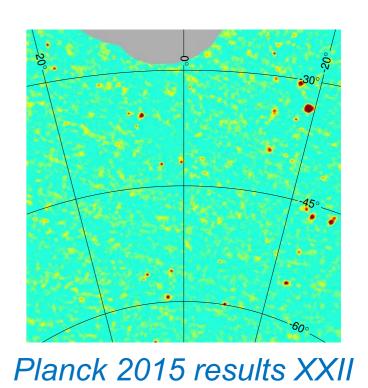
such that
$$\begin{cases} \langle \widehat{y}^2 \rangle = w^t \langle dd^t \rangle w \text{ minimum} \\ \sum_{\nu} w(\nu) Y_0(\nu) = 1 \end{cases} \tag{2}$$

• ILC weights : $w^t = \frac{Y_0^t C^{-1}}{Y_0^t C^{-1} Y_0}$ $(C \equiv \langle dd^t \rangle)$

$$\Rightarrow \widehat{y}(\overrightarrow{\theta}) = \underbrace{y(\overrightarrow{\theta})} + \underbrace{w^t N}_{recovered minimized signal residuals}$$

$$(2) \qquad by (1)$$





Miscalibrated Compton-y signal reconstruction

$$d(v, \overrightarrow{\theta}) = Y(v, T_e) \ y(\overrightarrow{\theta}) + N(v, \overrightarrow{\theta})$$
Planck frequency signal foregrounds of interest + noise

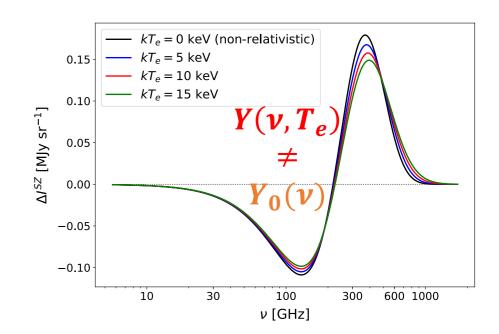
• ILC = weighted linear combination of frequency maps:

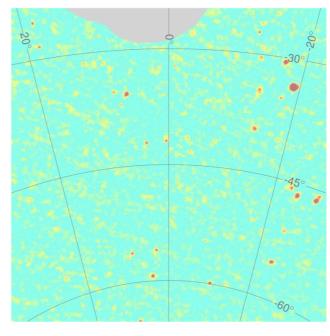
$$\widehat{y}(\overrightarrow{\theta}) = \sum_{\nu} w(\nu) d(\nu, \overrightarrow{\theta})$$

such that
$$\begin{cases} \langle \widehat{y}^2 \rangle = w^t \langle dd^t \rangle w \text{ minimum} \\ \sum_{\nu} w(\nu) Y_0(\nu) = 1 \end{cases} \tag{2}$$

• ILC weights :
$$w^t = \frac{Y_0^t C^{-1}}{Y_0^t C^{-1} Y_0}$$
 $(C \equiv \langle dd^t \rangle)$

$$\Rightarrow \widehat{y}(\overrightarrow{\theta}) = \frac{Y_0^t C^{-1} Y(T_e)}{Y_0^t C^{-1} Y_0} y(\overrightarrow{\theta}) + w^t N$$
Bias < 1



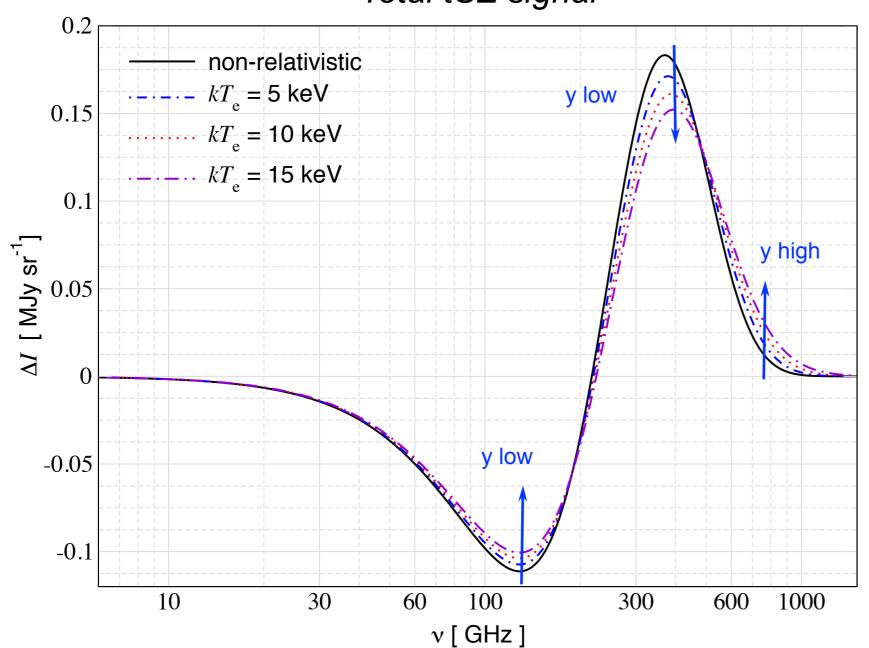


Planck 2015 results XXII

Underestimation of Compton-*y* **signal!**

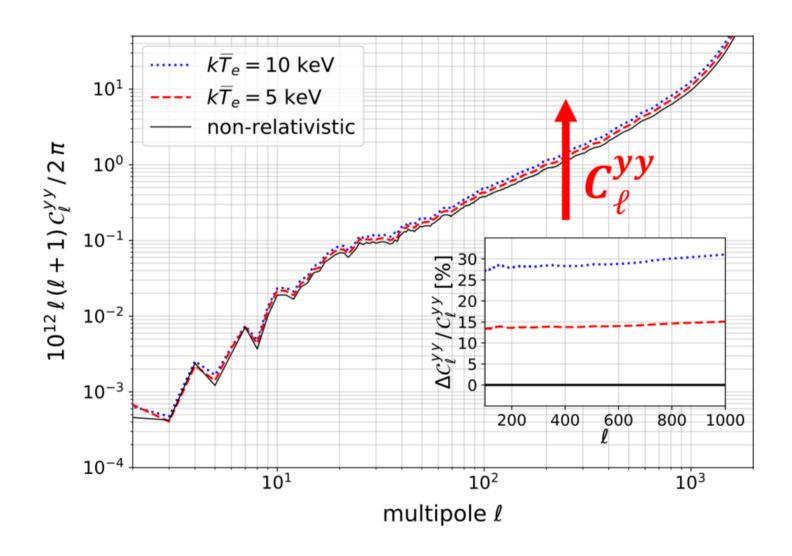
What is the net effect on the y-parameter?





- Obtained y-parameter is underestimated by ~7% for kT_e ~ 5 keV for Planck
- This is consistent with 353 GHz channel driving the effect for *Planck* data
- Also consistent with Erler et al. 2018 ILC analysis
- Total effect generally depends on frequency configuration and ability to subtract foregrounds (Rotti et al., in prep.)

Updating the Planck y-map power spectrum



Planck C_{ℓ}^{yy} increases with average cluster temperature \overline{T}_{e}

$$C_{\ell}^{yy} \propto \sigma_8^{8.1} \implies \frac{\Delta \sigma_8}{\sigma_8} \simeq 0.019 \left(\frac{k\bar{T}_e}{5 \text{ keV}}\right)$$

 $\simeq 1\sigma$ increase for $\bar{T}_e \simeq 5$ keV!

Remazeilles, Bolliet, Rotti, Chluba (2018)

Which effective electron temperature should be used?

• Single cluster / stacking → y-weighted temperature is relevant

$$kT_{\rm e}^y = \frac{\langle kT_{\rm e}y\rangle}{\langle y\rangle} = \frac{\int kT_{\rm e}^2N_{\rm e}\,\mathrm{d}l}{\int T_{\rm e}N_{\rm e}\,\mathrm{d}l}$$

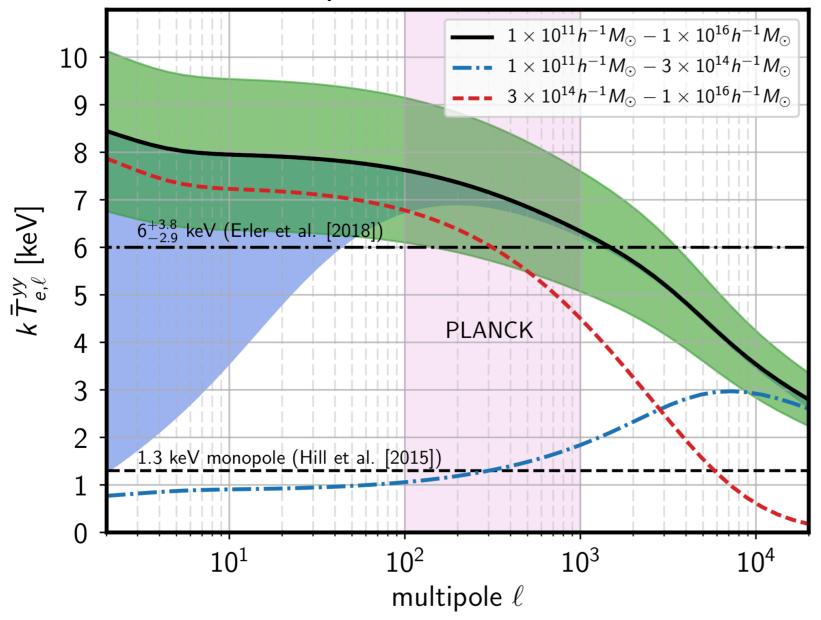
• SZ power spectrum analysis $\rightarrow y^2$ -weighted temperature is relevant

$$kT_{\mathrm{e},\ell}^{yy} = \frac{\langle kT_{\mathrm{e}}(M,z) | y_{\ell} |^{2} \rangle}{\langle |y_{\ell}|^{2} \rangle} \equiv \frac{C_{\ell}^{T_{\mathrm{e}},yy}}{C_{\ell}^{yy}}$$

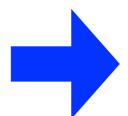
- Can be efficiently computed based on the halo model using CLASS-SZ
- Higher mass systems are up-weighted → higher effective temperature expected
- Scale-dependent quantity (but fixed temperature captures leading order effect)!

Theoretical estimate for the y^2 -weighted temperature

CLASS-SZ computation

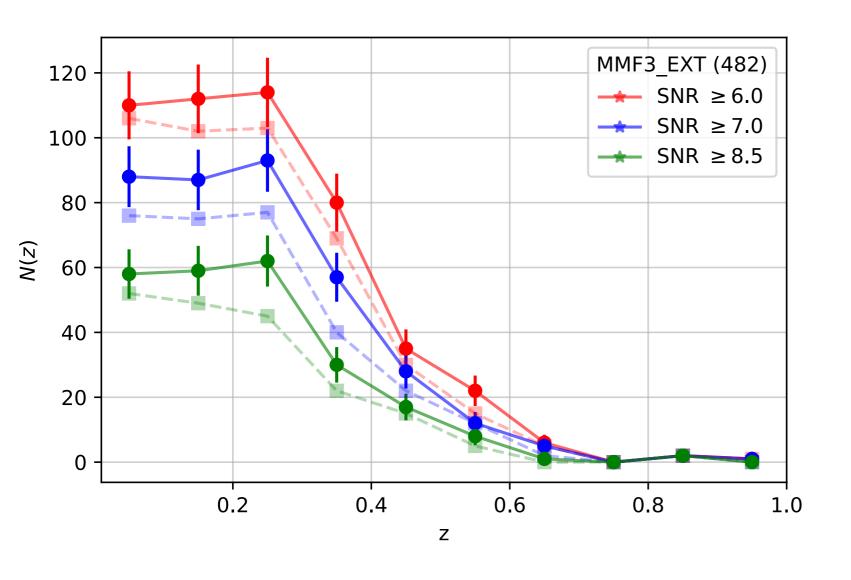


- Significant uncertainties...
- Contributions from diffuse y become important at large angular scales (e.g., Hansen et al., 2005)
- Assuming kT_e ~ 5 keV for Planck appears conservative
- Effective temperature is roughly constant for scales relevant to *Planck*

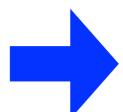


- rSZ plays a part in the σ_8 tension, alleviating it
- rSZ leads to systematic shift + increase of errors

Effect of rSZ on cluster number counts



- Rescaled S/N increases
- Effect of rSZ pushes some systems above S/N threshold
 - → # of clusters increases
- Caveats regarding noise rescaling
 - → filter scale also changed
- shape of N(z) seems better
- Still work in progress....



- rSZ again moves σ_8 into the right direction
- Preliminary effect found at ~1σ level

Rotti, Bolliet, Remazeilles & JC, in prep.

Future opportunities for rSZ studies

- Tens of thousands of clusters will be detected through the tSZ effect
- Unprecedented sensitivity, frequency coverage and angular resolution (e.g., SO, CCAT-prime, Millimetron, PICO...)
- Complements X-ray and lensing measurements



Individual systems

- y-maps, T_e-maps & velocity (?) maps
- Reconstruction of cluster profiles
- Non-thermal SZ (cosmic rays and turbulence)

Stacking analysis

- rSZ in mass bins
- Self-calibration of SZ temperature-mass relation
- Cosmology with new SZ observables (
 ← moments)

Statistical analysis

- rSZ power spectrum
- Cluster number counts in mass and redshift bins
- Higher order statistics (→ non-Gaussianity)

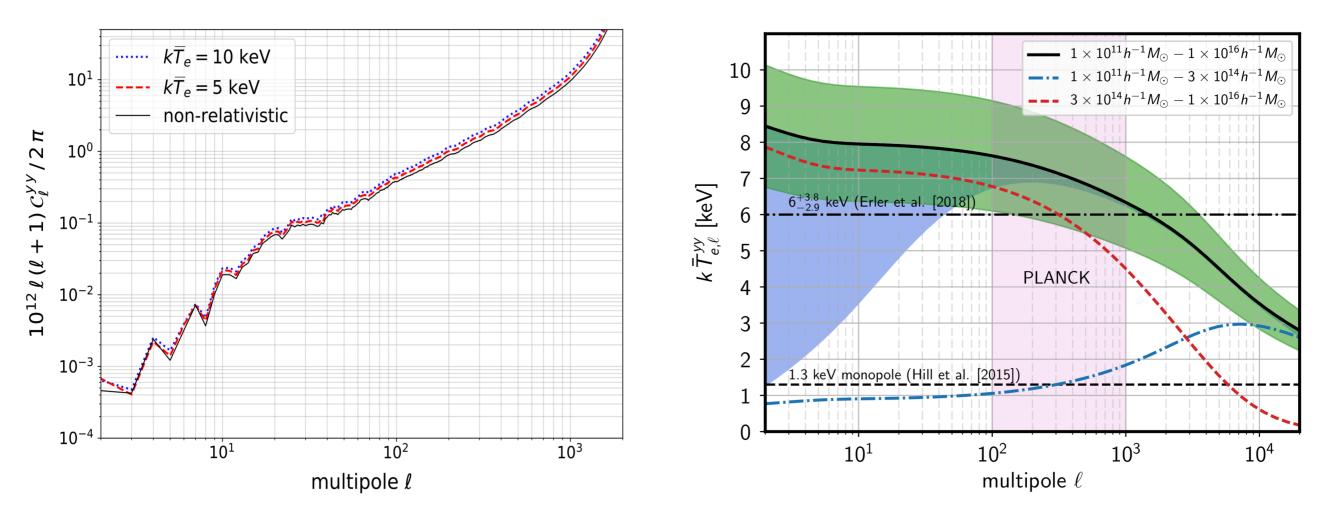


Highly relevant when using SZ clusters as a cosmological tool

 $(\leftrightarrow \text{neutrino masses}, \sigma_{8}, \text{dark energy})$

Conclusions

- rSZ still hard to see for individual clusters or in stacking analysis (e.g., Erler et al., 2017)
- rSZ causes an underestimation of the *y*-parameter for *Planck*
- rSZ plays a role in the σ_8 tension (shift upwards + increase of error)
- Careful analysis marginalizing over model uncertainties is required
- Net effect on SZ number counts still work in progress...



Remazeilles, Bolliet, Rotti & JC, MNRAS, 2019